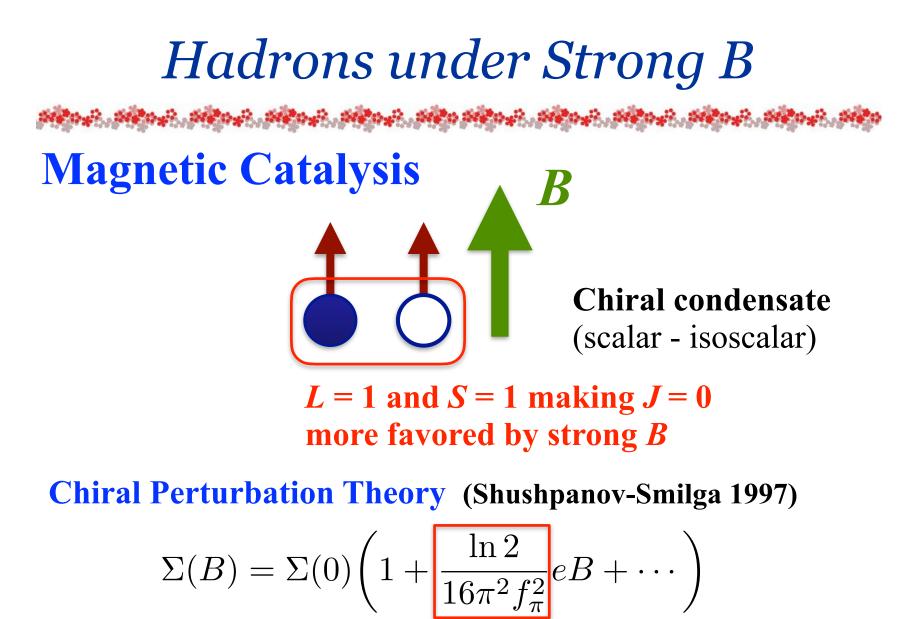
Magnetic Enhancement of Baryon Confinement Modeled via a Deformed Skyrmion

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Kenji Fukushima

The University of Tokyo

— 19th International Conference on QCD in Extreme Conditions (XQCD 2023) —

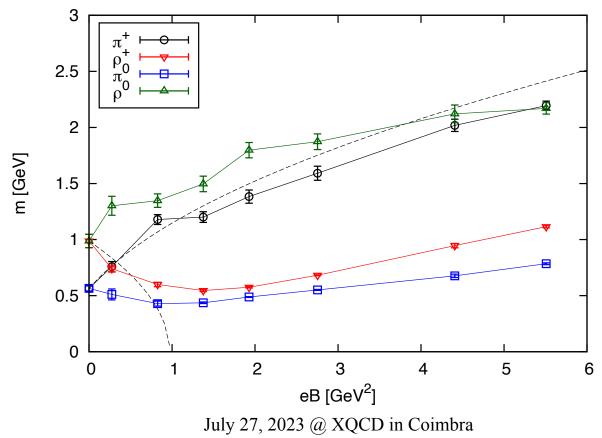


Positive coefficient

Hadrons under Strong B

Maria Ma Lattice-QCD

Hidaka-Yamamoto (2012)

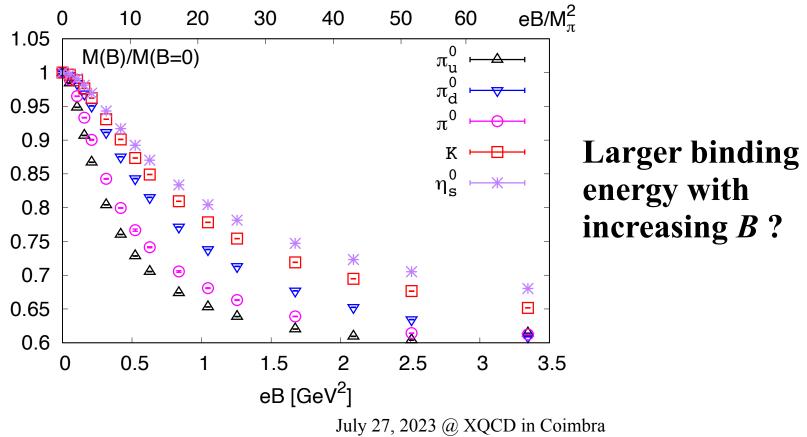


Hadrons under Strong B

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Lattice-QCD

Ding-Li-Tomiya-Wang-Zhang (2020)



Skyrmions under Strong B alla de alla de alla de alla alla de alla de alla de alla de alla de alla de Without *B* the baryon number is given by $\pi_3(\mathrm{SU}(2)) = \mathbb{Z}$ **Distinct homotopy connected?** (S^1 $\pi_1(U(1)) = \mathbb{Z}$ $\pi_3(\mathrm{SU}(2)) = \mathbb{Z}$ $B \rightarrow \infty$ B=0Phase transition, Crossover, ???

Skyrmions under Strong B $\Sigma = i \tau \cdot \Pi + \Pi_4 \qquad \Pi_1^2 + \Pi_2^2 + \Pi_3^2 + \Pi_4^2 = 1$

 z_0

-1

-1

 x^{0}

 $\pi_3(SU(2))$ winding is never unwound but it persist near the origin (where the *B* effect is small)

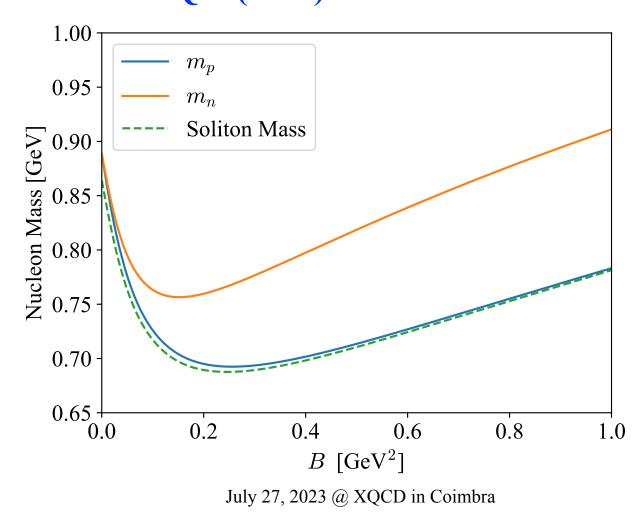
- Dominated by charged pions

Dominated by neutral pions

Winding between constant $\Pi_1^2 + \Pi_2^2$ and $\Pi_3^2 + \Pi_4^2$ surfaces

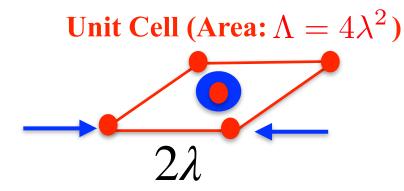
Skyrmions under Strong B

Chen-Fukushima-Qiu (2021)



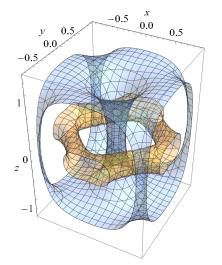
Skyrmions under Strong B Chen-Fukushima-Qiu (2021)

"Matter" candidates: Normal Crystal & Domain Wall



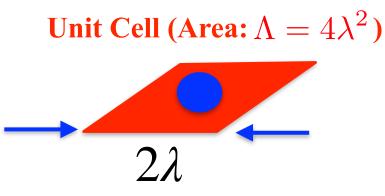
$$\begin{aligned} \Pi_4(0,0,0) &= - \ 1 \\ \Pi_4(\lambda,\lambda,0) &= + \ 1 \end{aligned}$$

The baryon number from $\pi_3(SU(2))$ is localized at the center and the edges.



Skyrmions under Strong B Chen-Fukushima-Qiu (2021)

"Matter" candidates: Normal Crystal & Domain Wall



 $\Pi_4(0,0,0) = -1$ $\Pi_4(x, y, 0) = -1$

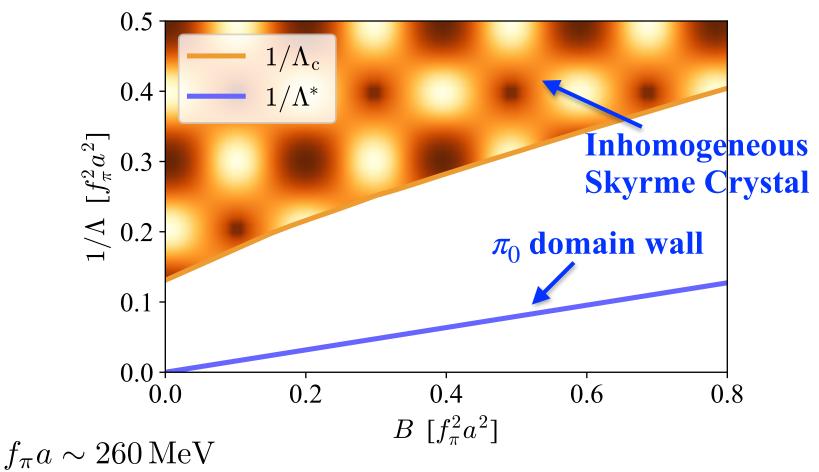
The baryon number from $\pi_1(U(1))$ is homogeneously distributed on the π_0 domain walls.

The baryon density and the magnetic flux are quantized

 $B\Lambda^*=2\pi$ Dirac quantization is strict!

Skyrmions under Strong B

Chen-Fukushima-Qiu (2021)

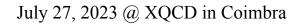


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The Electron-Ion Collider

A machine that will unlock the secrets of the strongest force in Nature



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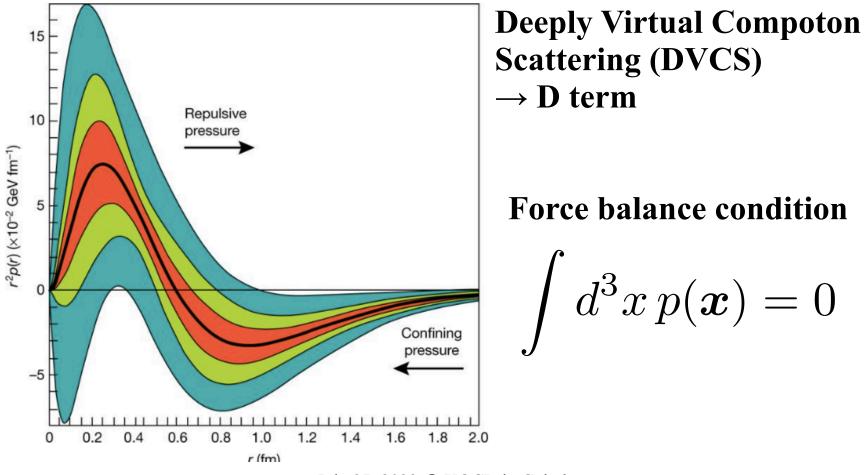
Quark and gluon confinement

Experiments at the EIC will offer novel insight into why quarks or gluons can never be observed in isolation, but must transform into and remain confined within protons and nuclei. The EIC—with its unique combinations of high beam energies and intensities—will cast fresh light into quark and gluon confinement, a key puzzle in the Standard Model of physics.

More

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Burkert-Elouadrhiri-Girod (2018)



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Chen-Fukushima-Qiu (2023)

Conservation law:

$$\partial^{\mu}T_{\mu\nu} = j^{\mu}_{Q}F_{\mu\nu}$$

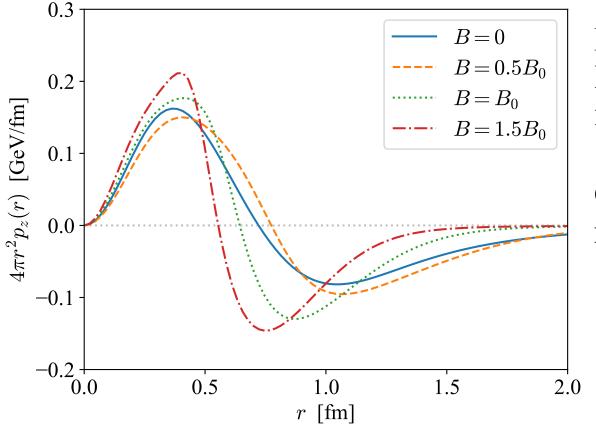
Spatial integration with x_{μ} **:**

$$\int d^3x \, T_{\mu\nu} = -\int d^3x \, x_\mu j_Q^\lambda \, F_{\lambda\nu}$$

Pressure sum rule:

$$P_z = \int d^3x \, p_z = \int d^3x \, T_{zz} = 0$$

Chen-Fukushima-Qiu (2023)

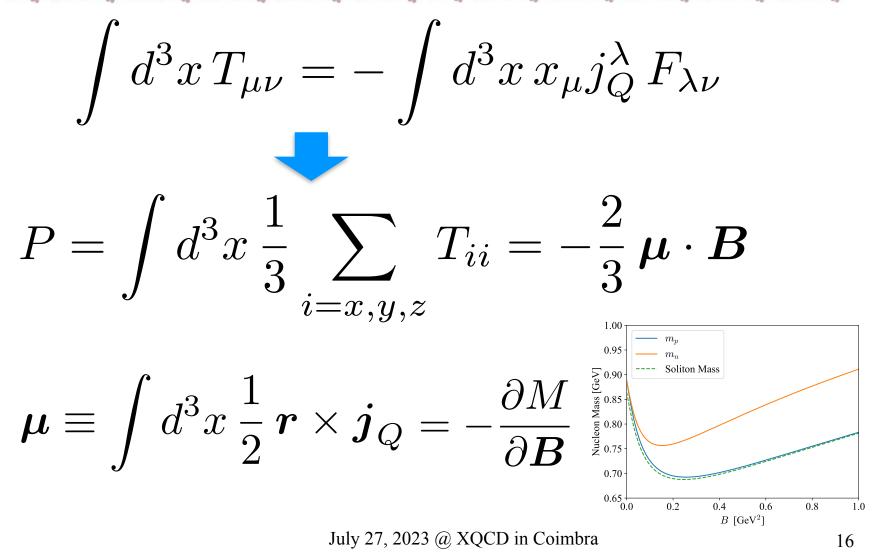


Pressure sum rule holds along the magnetic direction.

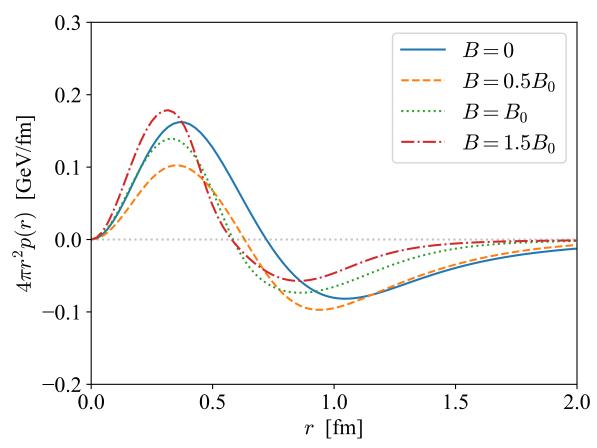
Oblate deformation is favored by *B*.

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Chen-Fukushima-Qiu (2023)



Less confining pressure is needed.

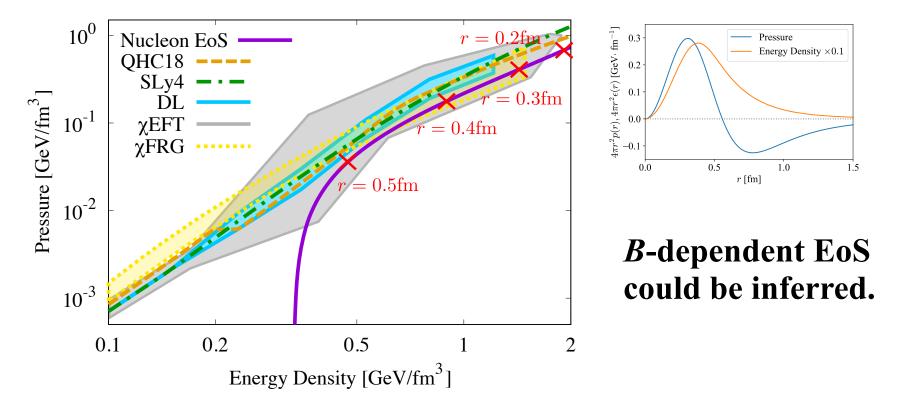
Confining force is provided by the magnetic pressure.

Depends on the sign of the magnetic moment?

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Possible Application

Fukushima-Kojo-Weise (2020)



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Conclusions

Charged and neutron meson masses have been calculated in lattice-QCD.

Pressure balance and confining pressure can be calculated / measured.

Whether *B* favors confinement or deconfinement can be judged based on the pressure sum rule.