

A few days ago, I received an email from one of the organizers of this conference discussing a **particular scientific question**. With this presentation, I am answering this letter.

To do this, I slightly changed the initial title of my talk to "Whispering Gallery of a Neutral Particle with a Gravitational Shift" and yesterday, at the last desk of the conference, I completely rewrote my talk.



I am grateful to the organizers for the possibility to participate in the nice conference and always happy to come to this place (due to a different reason).

I would like to underline the excellent preparation of the meeting. Even a special wine with a quantum interference pattern (the quantum technology frontier !) has been specially produced for the participants of this meeting.

Lo Zefiro CHIANTI

03.07.23

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V.V. Nesvizhevsky

Lo Zefiro

CHIANTI

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The whispering gallery states (WGS) of a neutral particle is a useful phenomenon for searching for new physics beyond the Standard Model using precision quantum measurement technology (you see, the topic of my talk is precisely in line with the title of the conference ...)

Anna Nesvijevskaia:



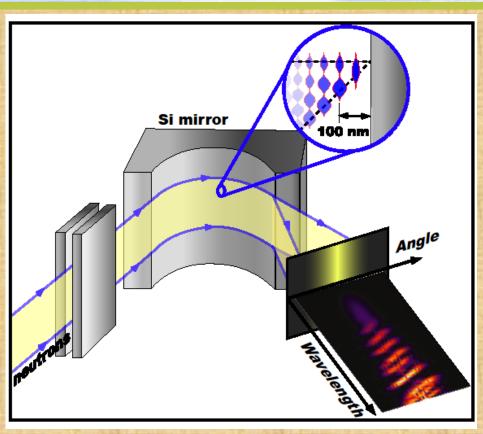
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What is the WGS of a neutral particle?

- a neutral particle (otherwise electromagnetic interactions dominate, and the WGS phenomenon does not exist, or at least precision measurements are impossible),
- the quantum reflection of this neutral particle from the surface (it is not very important what the nature of this quantum reflection is (a wave may always get reflected from a surface...),
- a curved mirror (to provide an effective centrifugal force),
- certain conditions on the parameters of the problem (observation time, particle velocity and mass, mirror radius, etc) essentially determined by the uncertainty relations UNSTITUT MAX VON LAUE - PAUL LANGEVIN V.V. Nesvizhevsky

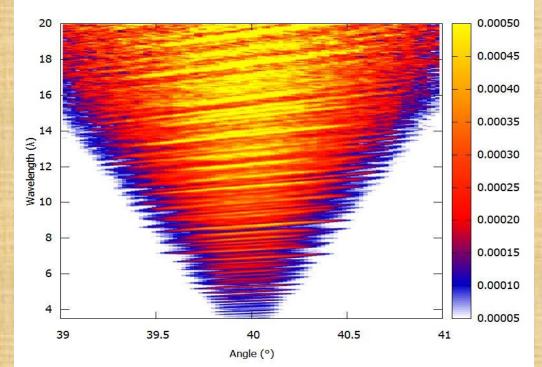




a **typical interference pattern** (shown to everyone at various conferences, but not published yet...

[V.V. N., et al, Neutron whispering gallery, Nature Physics 6:114, 2010].

The scheme from [H. Rauch, Neutrons in a whispering gallery, Nature Physics 6:79, 2010].



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New physics beyond the Standard Model, which can be studied using WGS, can be generalized as a search for extra fundamental short-range interactions (depending on the spin (axion-type) and not depending on the spin) between the particle and the mirror.

The WGS is a very competitive method over a certain range of characteristic distances of such interactions. Many modern physics problems could be probed:

- dark matter, dark energy, the origin of gravity, matterantimatter asymmetry, theories with extra light bosons, extra dimensions of space, etc,
- also various applications in quantum mechanics studies, surface studies, precision spectroscopy, etc
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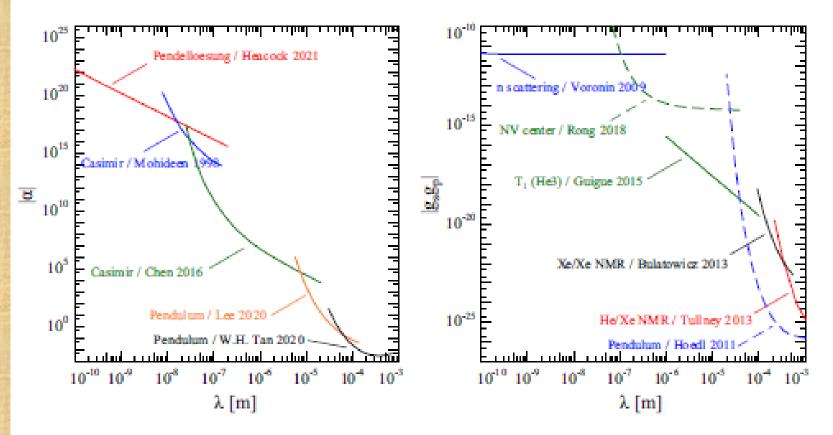


Figure 8: Left: Current most constraining limits for new spin-independent short-range forces. The first experiment (red line, [179]) uses free neutrons. At higher λ ranges, other techniques and probes are more sensitive [180, 181, 182, 183]. Right: Current most constraining limits for new spin-dependent interactions. Experiments [184, 185, 186, 187] use the neutron as the polarized particle, for the first of those it is a free neutron. In [188, 189], the polarized particle is an electron.

[H. Abele et al, *Particle physics at the European spallation source*, to be published in Physics Reports 2023] — the latest review.

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Many of the features of WGS and gravitational quantum states (GQS) do not depend on which neutral particle is used: neutron (all first experiments), hydrogen atom (GRASIAN collaboration), antihydrogen atom (GBAR collaboration), positronium, muonium (gravity of antimatter?), heavier neutral atoms, molecules, nanoparticles, nanodroptets etc.

Measurements with each of these particles have their own motivation and optimal sensitivity areas.

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Efficient quantum reflection results from almost ideal sharp **potential step** provided that the spatial size of the WGS is much larger than the spatial size of the diffuseness of the potential step, which is why WGS (and GQS) **depend little** on the properties of the mirror.

Moreover, the corresponding the small corrections (associated with the interaction of the particle with the mirror) can be accurately calculated (Serge Reynaud and collaborators) and experimentally studied.

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If the particle masses are (nearly) equal (neutron, hydrogen, antihydrogen), the WGS are practically identical.

If the mass is smaller (muonium, positronium), the states have a larger spatial size and the quantum reflection is more efficient.

If the mass is larger (heavier atoms, molecules, nanoparticles), the quantum states spatial sizes are smaller, the reflection is less efficient, the systematic effects are larger, however, a better sensitivity to extra fundamental interactions at smaller distances may appear.

Of course, the optimal **methods of observation/detection** depend on which particle is used and must be considered on a case-by-case basis.

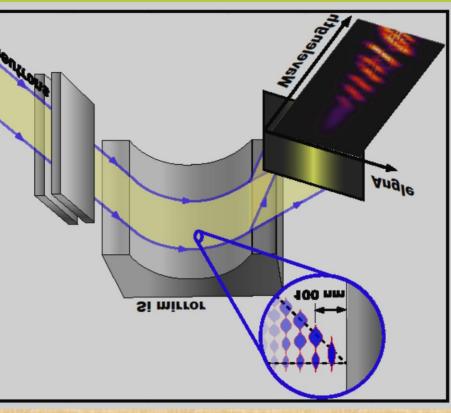
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With the characteristic parameters of all the WGS experiments discusse or performed, gravity is a small correction (hence "shift").

Moreover, this correction is usually suppressed by choosing the mirror orientation: the gravity axis is perpendicular to the axis of the centrifugal force.



However, the gravitational effect in WGS may be the goal of the experiment (as in the case of antimatter particles).

In this case, you need just to rotate the mirror by 90 degrees and examine this effect more closely. 03.07.23 V.V. Nesvizhevsky



[V.V. N. et al, *Quantum states of neutrons in the Earth's gravitational field*, Nature 415 (2002) 297]

First, let's analyze, for simplicity, GQS of ultracold neurons. This is an analogous phenomenon, differing only in the attraction of the particle particle to the surface:

gravity instead of centrifugal force

Similar triangular potential and theoretical formalism



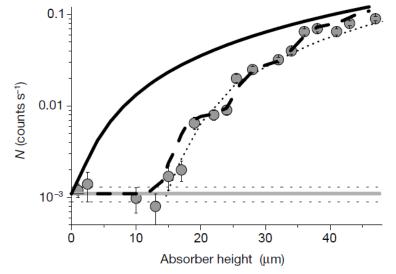


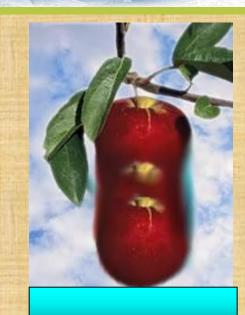
Figure 4 The neutron throughput versus the absorber height at low height values. The data points are summed up in intervals of 2 μ m. The dashed curve corresponds to a fit using the quantum-mechanical calculation, in which all level populations and the height resolution are fitted from the experimental data. The solid curve is again the full classical treatment. The dotted line is a truncated fit in which it is assumed that only the lowest quantum state—which leads to the first step—exists.

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- mass 1 at. un.,
- typical energies, $\varepsilon_0 \sim 0.6 \text{ peV}$,
- typical vertical sizes of GQS, $I_0 \sim 5.9 \mu m$,
- typical effective temperatures, ~10 nK,
- typical times of formation of GQS is $\tau_0 \sim 1.1$ ms



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- The surface reflects neutrons **elastically** due to the condition $l_0 >> 1$ Å (a typical interatomic distance). It means that the thermal motions of individual nuclei/atoms are effectively averaged out, and

- **specularly** provided the surface is well-polished and homogeneous (on the spatial scale of the order of I_0),

- A characteristic range of raising the potential should be $\langle \cdot |_0$ (the condition of sharpness of the reflecting potential). Otherwise, the raising part of potential modifies GQS, and this modification has to be known precisely.

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Two uncertainty ratios:

- Vertically, for the uncertainty of vertical position and vertical velocity (energy) l₀~5.9 μm v₀~2.4 cm/s (ε₀~0.6 peV).
- All these parameters of the problem are determined only by the value of gravitational acceleration at the local point of the measurement g~9.81 m/c². Parameters of the mirror are not involved, to a good accuracy.
- It is difficult to change g by a noticeable amount...

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You can tilt the mirror.

In this case, the component of motion along the mirror ("horizontal") is still classical and is **separated** from the vertical component ("quantum"), in a properly designed experiment.

The projection of the gravitational field will decrease. However, this is not very interesting, because the most interesting area in terms of sensitivity to shortrange interactions is "stronger" rather than "weaker" gravity.

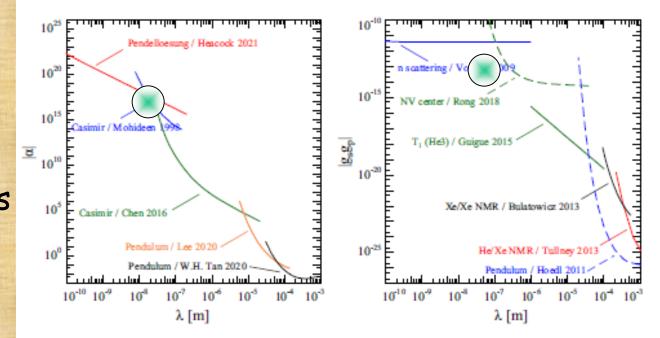


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You can move the mirror vertically with a constant acceleration.

However,

- the observation time will be limited, and
- the systematics will increase dramatically, because instead of a "perfectly uniform and known" gravitational field, there will be a "badly defined 3-dimensional non-constant and poorly measured" acceleration.



You can add a vertical magnetic field gradient.

Indeed, the gradient will affect the magnetic moment of the neutral particle (opposite signs of the force for the two spin states of the particle).

It is effectively possible

- to "increase" the gravitational interaction,
- "decrease" the gravitational interaction,
- compensate for the gravitational interaction ("no gravity"),
- and even simulate "anti-gravity".

The potential problem is the same as with the vertical acceleration of the mirror: it is **difficult to control** this effect with a reasonable accuracy.

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- Horizontally: the uncertainty relation for the energy of the quantum state ($\epsilon_0 \sim 0.6 \text{ peV}$) and the observation time ($\tau_0 \sim 1.1 \text{ ms}$).

- An ultracold neutron with a typical velocity of ~5 m/s travels ~0.55 cm in τ_0 ~1.1 ms. To provide condition τ >> τ_0 , we need the mirror length of at least 5.5 cm. A technically feasible length is, say, 55 cm.

- The corresponding accuracy of experiments with neutrons, in the flow-through mode, is $\sim 10^{-2}$. With certain efforts to take into account systematics, it may be improved to, say, $\sim 10^{-3}$. Then there are insurmountable **problems**: the systematics following from the uncertainty relation, and too low densities of neutrons in the phase space. Since the required increase in sensitivity is **still 2-3 orders** of magnitude, the task is difficult.

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GR **SI** One found way out of all these problems: instead of neutrons, use hydrogen atoms (GRASIAN).

- the density of hydrogen atoms in the phase space can be many orders of magnitude higher, which solves the problem of statistical sensitivity,
- it is easier to implement the trapping of atoms in the horizontal direction ([V.V. N. et al, A magneto-gravitational trap for studies of gravitational quantum states, Europ. Phys. J. C 80 (2020) 520], [J. Ahokas, et al, A large octupole magnetic trap for research with atomic hydrogen, Rev. Sci. Instr. 93 (2022) 023201]), thus increasing the observation time by orders of magnitude and solving the problem of systematic effects.

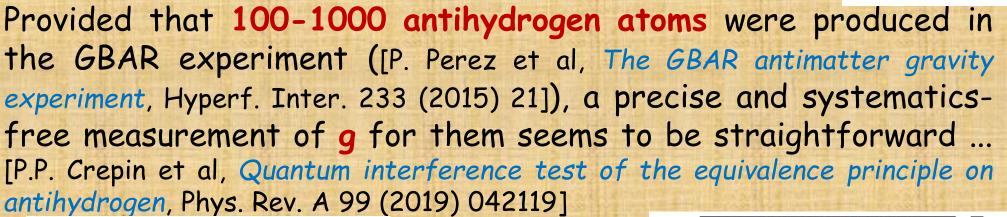
An option for a very competitive search for fundamental interactions at distances of ~1-10 µm.

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Even the density of antihydrogen atoms in phase space can be much larger than typical UCN densities.



0.04

0.03

0.02

0.01





What about extra fundamental interactions at shorter distances?

Whispering Gallery!

"Any" acceleration magnitude and better (than with mirror acceleration or magnetic field gradient) control over it.

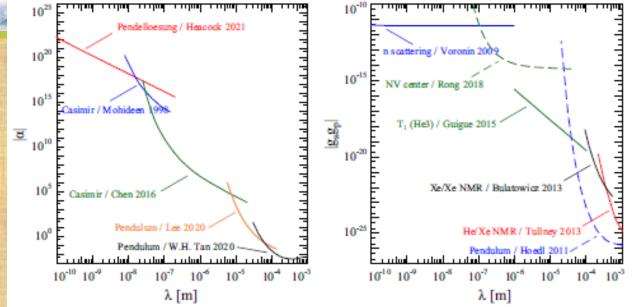


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The shape of a curved mirror is not known "exactly". But there is a good test experiment to control it: X-ray whispering gallery (no short-range forces and much higher statistical sensitivity). Very good experimental data are being analyzed now.

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Excellent result from WGS experiments with both polarized and unpolarized neutrons.

Competitive (best) constrains for characteristic distances of ~10 nm.

Significant potential for a further increase in sensitivity (no fundamental limitation of sensitivity improvement is found yet).

On a longer time-scale, it is possible to replace neutrons with **hydrogen atoms or heavier atoms** (GRASIAN) to further increase the sensitivity.

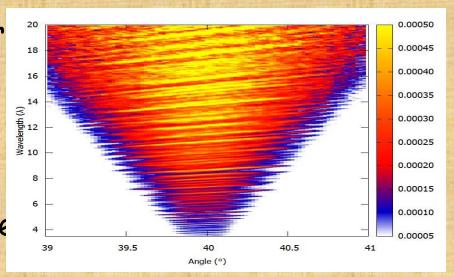
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Very well defined interference patter

What will happen to them if we add a gradient of a magnetic field or a gravitational field in a direction (close to) perpendicular to the surface



The interference pattern is deformed (shifted).

As always in interferometric experiments, a high sensitivity to additional interactions is expected.

Is it possible to measure this effect, for example, with neutrons?

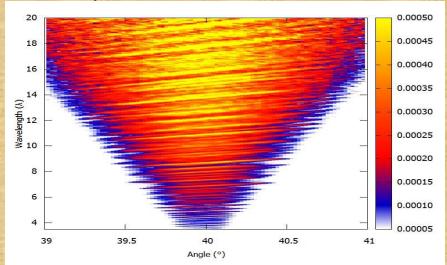
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In order of magnitude, the relative shift of the interference pattern lines is equal to the ratio of g to the centrifugal acceleration, or ~10⁻⁶ (I don't want to bore you with coefficients of the order of one).

A typical statistical accuracy for neutrons with the velocity of ~10³ m/s is 1/N x Δ S/S ~10⁻⁵, where N~10² (the number of interference lines) and Δ S/S is a typical statistical sensitivity for a line position.



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The **sensitivity** on top of the interference pattern (the slowest neutrons in the standard spectrum of cold neutrons) is already **sufficient**!



In a typical neutron experiment with a silicon mirror, the radius of the mirror was equal R = 2.5 cm. The radium is chosen so to provide a few sub-barrier WGS. The energy of WGS is $E_{WGS} =$ $\frac{\hbar}{\tau}$, where \hbar is the reduced Planck constant, and τ_{WGS} the characteristic time of formation of WGS equal $\tau_{WGS} = \sqrt[3]{\frac{2\hbar}{ma^2}}$, where m is the neutron mass and a is the centrifugal acceleration equal $a = \frac{V^2}{R}$, where V is the neutron longitudinal velocity. The characteristic size of WGS is $l_{WGS} = \sqrt[3]{\frac{\hbar^2}{2am^2}}$.

Can a similar measurement be performed with muonium?

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In order to evaluate such an experiment, we will use the same formulas but apply the following scaling factors: the ratio of particle masses is $\frac{m_{neutron}}{m_{muonium}} \sim 8.9$. First, we estimate the length of the cylindrical mirror surface L , which corresponds to the condition that most muonium atoms disappear L = $3 \cdot V_{muonium} \tau_{muonium} \sim 4 \, cm$, where $V_{muonium} = 6300 \pm 200 \, m/s$ is the muonium velocity and $\tau_{muonium} \sim 2.2 \,\mu s$ the muon's lifetime. Note about the same length as it is for neutrons, however, the need to use another detection scheme (annihilation on the mirror surface). A spatial resolution needed should be significantly better than $3 \cdot V_{muonium} \tau_{muonium}$, i.e. a fraction of millimeter.

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The reflectivity of muonium from a helium surface is shown here.

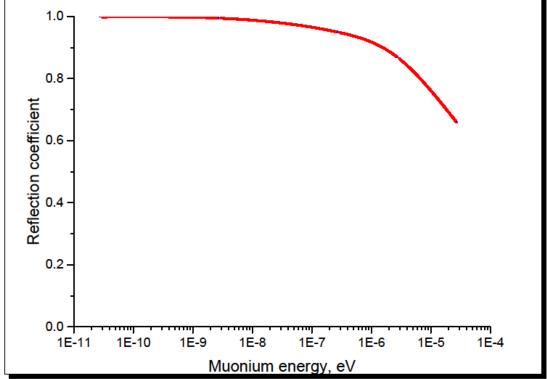
The effective potential height is $E_{QR,muonium} \sim 1 \mu eV$.

In order to have a few WGS, the centrifugal acceleration has to be $a_{muonium} \sim 1.5 \cdot 10^9 g$ and the mirror radius $R_{muonium} \sim 2.7 \ cm$ (close to neutron experiments). Statistics is also comparable to that in neutron experiments.

Thus, WGS of muonium can be observed. A gravitational shift probably not, due to the too high velocity $V_{muonium} = 6300 \pm 200 m/s$.

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Conclusion

The methods of WGS and GQS have passed a long way: from an issue of textbooks to the observation with UCNs and possible extension to other particles (atoms, antiatoms, etc), also to the applications (studies of gravitational interaction, short-range interactions, physics beyond Standard Model, extensions of quantum mechanics, surface studies, spectrometry, etc).

Interesting results have been obtained using these methods; Perspectives for further **progress are clear**.

A gravitational shift of WGS of neutrons can and should be measured; WGS of muonium can be measured; to measure a gravitational shift of WGS of muonium, the muonium velocity has to be further reduced by an order of magnitude.

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Simulation

I simulated (approximately) an interference pattern that can be measured using the current parameters of the muonium beam. The result is about the following.

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