

Hadronisation of heavy quarks in pp collisions with ALICE at the LHC

Syaefudin Jaelani National Research and Innovation Agency of Indonesia - BRIN On behalf of the ALICE Collaboration

Charm and **beauty** quarks: $m_c \sim 1.3$ GeV/ c^2 , $m_b \sim 4.2$ GeV/ c^2 M Produced in hard-scattering processes M M the factorisation approach

dσ^c *dp*^c T $(x_1, x_2, \mu_R, \mu_F) \otimes D_{c \to D} (z = p_D/p_c, \mu_F)$

The production of heavy-flavour hadrons in hadronic collisions can be described by

$$
\frac{d\sigma^D}{dp_T^D} (p_T; \mu_F; \mu_R) = PDF (x_1, \mu_F) PDF (x_2, \mu_F) \otimes
$$

- **Charm** and **beauty** quarks: $m_c \sim 1.3$ GeV/ c^2 , $m_b \sim 4.2$ GeV/ c^2 M Produced in hard-scattering processes M
- M the factorisation approach

Test of perturbative QCD calculations M

The production of heavy-flavour hadrons in hadronic collisions can be described by

$$
\frac{d\sigma^D}{dp_T^D} (p_T; \mu_F; \mu_R) = \overline{PDF (x_1, \mu_F) PDF (x_2, \mu_F)} \otimes
$$
\nParton distribution functions
\n(PDFs)

- **Charm** and **beauty** quarks: $m_c \sim 1.3$ GeV/ c^2 , $m_b \sim 4.2$ GeV/ c^2 M Produced in hard-scattering processes M
- M the factorisation approach

The production of heavy-flavour hadrons in hadronic collisions can be described by

- M
- Measurement of fragmentation fractions (FF)

The yield ratios of hadrons are sensitive to the heavy-flavour hadronisation process

$$
\frac{d\sigma^D}{dp_T^D} (p_T; \mu_F; \mu_R) = \overline{PDF (x_1, \mu_F) PDF (x_2, \mu_F)} \otimes \frac{\overline{d\sigma^c}}{dp_T^c (x_1, x_2, \mu_R, \mu_F)} \otimes \overline{D_{c \to D} (z = p_D/p_c, \mu_F)}
$$
\nParton distribution functions
\n
$$
\overline{Hard scattering}
$$
\n(hadronisation)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(250)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(260)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(270)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(280)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(290)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(290)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(200)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(200)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(200)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(200)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F)}
$$
\n(200)
\n
$$
\overline{D} = \overline{PDF (x_1, \mu_F) P D F (x_2, \mu_F) P D F (x_2, \mu_F)
$$
\n(200)
\n

- **Charm** and **beauty** quarks: $m_c \sim 1.3$ GeV/ c^2 , $m_b \sim 4.2$ GeV/ c^2 M
- Produced in hard-scattering processes M
- M the factorisation approach

- Reference for p-Pb and Pb-Pb collisions M
- Test of pQCD calculations
- Study hadronisation \blacksquare

The production of heavy-flavour hadrons in hadronic collisions can be described by

$$
\frac{d\sigma^D}{dp_T^D} (p_T; \mu_F; \mu_R) = PDF(x_1, \mu_F) PDF(x_2, \mu_F) \otimes \frac{d\sigma^c}{dp_T^c} (x_1, x_2, \mu_R, \mu_F) \otimes D_{c \to D} (z = p_D/p_c, \mu_F)
$$
\nParton distribution functions\n\nHard scattering\n\nFragmentation function (hadronisation) (hadronisation)

pp collisions

ALICE Detector

Charm and beauty mesons

Non-prompt D-meson production measured down to low p_T (D ^+_S down to p_T = 2 GeV/*c*) -> access to beauty-meson production

- No significant p_T -dependence $\mathbf O$
- Good agreement with models that use FF tuned on leptonic collisions and with measurements at e^+e^- colliders $\mathbf O$
- Meson-to-meson yield ratios independent of $p_{\rm T}$ and collisions energies $\mathbf O$
- \bullet mechanisms

Charm and beauty fragmentation to mesons

Fragmentation fraction ratios for charm and beauty mesons are well described by PYTHIA8 predictions (with FF tuned on e^+e^-)

- \bullet
- No significant dependence on energy and collision systems Ő - From e^+e^- and ep to hadronic collisions

- $\Lambda_{\rm c}^{+}/{\rm D}^0$ ratios significantly higher than ${\rm e}^+{\rm e}^-$ results (LEP average: \bullet $0.113 \pm 0.013 \pm 0.006$ *EPJC 75 (2015) 19*
- *p*™ dependence observed, not present in $e^+e^−$ results
- PYTHIA8 Monash \bullet $\Lambda_{\rm c}^{+}/{\rm D}^0$ ratio at low $p_{\rm T}$ larger than what predicted by string fragmentation models tuned on e⁺e⁻ data
- Models which introduce a modified hadronisation with respect to \bullet in-vacuum fragmentation describe instead the data: PYTHIA8 (CR Mode 2), Catania, SHM+RQM, and QCM
- PYTHIA8 (CR Mode 2) \bullet Colour reconnection mechanisms beyond leading colour (BLC) approximation with new junction topologies that favour baryon formation *JHEP 1508 (2015) 003*
- Catania
	- thermalised system of u,d,s and gluons *PLB 821 (2021) 136622*
	- hadronisation via interplay of fragmentation and coalescence

- $\Lambda_{\rm c}^{+}/{\rm D}^0$ ratios significantly higher than ${\rm e}^+{\rm e}^-$ results (LEP average: \bullet $0.113 \pm 0.013 \pm 0.006$ *EPJC 75 (2015) 19*
	- *p*™ dependence observed, not present in $e^+e^−$ results
	- SH model + RQM
		- Quark hadronisation driven by statistical weights govern by hadron masses
		- Feed down from excited baryon states predicted by the Relativistic Quark Model (RQM) *PLB 795 (2019) 117-121*
	- QCM
		- Pure coalescence model
		- Charm is combined with co-moving light antiquark or two quarks *EPJC 78 (2018) 344*
- $\Lambda_{\rm c}^{+}/{\rm D}^{0}$ ratios qualitatively described by PYTHIA8 (CR Mode 2), \mathbf{O} Catania, SHM+RQM, and QCM

- Enhancement observed for heavier charm baryons $\mathbf O$
- $\Sigma_{\rm c}^{0,++}/\rm D^{0}$ largely enhanced with respect to $\rm e^+e^-$ measurements (~0.02 from Belle, PRD 97 (2005) 07) \bullet
- PYTHIA8 with CR-BLC, SHM+RQM, Catania, and QCM describe the $\Sigma_{\rm c}^{0,++}$ /D 0 ratio $\mathbf O$
- $\mathbf O$

SHM+RQM, Catania, and QCM describe the $\Lambda_c^+(\gets \Sigma_c^{0,+,++})/\Lambda_c^+$ ratio while PYTHIA8 with CR-BLC overestimates the data

 $\Xi_{\rm c}^{0,+}{\rm /D^0}$ higher than PYTHIA8 Monash, tuned to reproduced $\rm~e^+e^-$ results \bullet $\mathbf O$

11/7/2023 S. Jaelani - ICNFP 2023 12

Catania model (including hadronisation via coalescence) describes better the shape of the measured $\Xi_c^{0,+}/\rm{D}^0$

- $\Xi_{\rm c}^{0,+}{\rm /D^0}$ higher than PYTHIA8 Monash, tuned to reproduced $\rm~e^+e^-$ results \bullet
- Catania model (including hadronisation via coalescence) describes better the shape of the measured $\Xi_c^{0,+}/\rm{D}^0$ $\mathbf O$
- BR ($\Omega_c^0 \to \Omega^- \pi^+$) * Ω_c^0 / D^0 ratios show no p_T dependence \bullet
- Catania model closer to the measurement when decays from additional higher-mass resonances are considered $\mathbf O$

Charm fragmentation fractions and total charm cross section

ALI-PUB-500750

- Significant baryon enhancement with respect to e^+e^- and ϵ p collisions \bullet - Enhancement of a factor of \sim 3.3 for Λ_c^+
-
- bands

 $f(c \to H_c)$ different in pp and e^+e^- and ep collisions: f**ragmentation fractions are not universal across the collisions systems** $c\bar{c}$ production cross section measured at midrapidity lie on the upper edge of the FONLL and NNLO calculation uncertainty

Total beauty cross section

11/7/2023 S. Jaelani - ICNFP 2023 15

JHEP 05 (2021) 220

- dielectron, as well as with pQCD predictions (FONLL and NNLO)
- $\mathbf O$

The results from D-meson species are compatible within uncertainties among each other and with those obtained from

bb production cross section measured at midrapidity is found to be compatible with FONLL and NNLO calculations

Beauty baryon-to-meson ratios

- Good agreement between prompt and non-prompt $\Lambda_{\rm c}^{+}/{\rm D}^0$ ratios \bullet
- Similar baryon-to-meson enhancement compared to e^+e^- measurements
- Non-prompt $\Lambda_{\rm c}^+$ /D 0 ratios are well described by FONLL + PYTHIA8 model (when fragmentation fractions measured by LHCb are \bullet employed)
	- **->** access to beauty-baryon production mechanisms!

- $\rm D^+_s/D^0$ ratios do not show any $p_{\rm T}$ dependence and event multiplicity \bullet
- For $1 < p_{\rm T}$ < 12 GeV/*c,* clear hierarchy of the $\Lambda_{\rm c}^+/D^0$ ratios from high to low multiplicity events (5.3 σ significance) \bullet
- PYTHIA8 Monash does not reproduce the $p_{\rm T}$ trend of the $\Lambda_{\rm c}^+$ /D⁰ across different event multiplicities while PYTHIA8 CR-CLB describes the \bullet magnitude and the $p_{\rm T}$ trend of the $\Lambda_{\rm c}^{+}/{\rm D}^{0}$

Prompt D mesons and Λ_c^+ vs. event multiplicity c

D_s^+/D^0 and Λ_c^+/D^0 ratios measured in pp collisions at $\sqrt{s} = 13$ TeV in MB collisions and for different multiplicity classes

Prompt $\Lambda_{\rm c}^{+}/{\rm D}^0$ vs. event multiplicity

- pp, p-Pb, Pb-Pb shown together as a function of event \bullet multiplicity
- $p_{\rm T}$ -integrated, extrapolated down to $p_{\rm T}$ = 0, $\Lambda_{\rm c}^{+}/{\rm D}^0$ ratios do not dependent on **multiplicity**, **collision system** and **energy** within uncertainties
- Re-distribution of $p_{\rm T}$ that acts differently for baryons and mesons \bullet - $>$ No modification of overall p_T -integrated yield ratios
- Same mechanism in all collision systems? Modified \bullet hadronisation? Radial flow?
- Flat trend reproduced by models with hadronisation via \bullet **fragmentation** + **recombination** (Catania, TAMU)

Pb-Pb collisions

- **E** ALICE, pp at 13 TeV: Phys.Lett.B 829 (2022) 137065
- **E** ALICE, pp and p-Pb at 5.02 TeV: *Phys. Rev. Lett.* 127, 202301 Phys. Rev. C 104, 054905

E ALICE, Pb-Pb at 5.02 TeV: https://arxiv.org/abs/2112.08156

E STAR, Au-Au at 200 GeV: *Phys. Rev. Lett.* 124, 172301

A few more hints: D_{s1}^+ and D_{s2}^{*+} production

- First measurement of D_s^+ -resonance production in pp \bullet collisions at $\sqrt{s} = 13$ TeV
- No multiplicity dependence on D_{s1}^+/D_s^+ ratio $\mathbf O$
	- Reproduced by the statistical hadronisation model (SHM)
- Hint of decreasing trend of D_{s2}^{*+}/D_s^+ ratio with event $\mathbf O$ multiplicity?

- Interplay between hadron lifetime and hadronic rescattering?

- Hint of tension with SHM predictions
- Total uncertainties too large to conclude

Outlook: LHC Run 3 data taking

ntegrated luminosity [pb⁻ ALICE Performance 2022, pp \sqrt{s} = 13.6 TeV 20 2022-11-27 04:04:05 BARREL: $L = 17.6$ pb $MUON: L = 16.9$ pb $TRD: L = 2.0$ pb EMCAL: $L = 0.5$ pb **TRIGGERED** PHOS: $L = 0.3$ pb⁻¹ 14 12 10 27 Jun 27 Jul 26 Aug 25 Sep 25 Oct 24 Nov

- Larger interaction rate and upgrade of ALICE apparatus during LS 2 **->** Larger data samples in Run 3 than Run 2 (x10-100 depending on the collision system)
	- **->** Improved impact parameter resolution
	- - $>$ Lead to more precise measurements, and with an extended $p_{\rm T}$ reach, of the observables studied in Run 2
- Direct reconstruction of beauty mesons and baryons
- Better constraints to theoretical models of the strongly interacting medium and hadronisation

→Target samples of ALICE high-energy pp programme

- $L_{int} = 200$ pb⁻¹, B = 0.5 T
- $L_{int} = 3$ pb-1, B = 0.2 T
- **Provings 7 Target samples of ALICE high-energy PbPb programme**
	- $L_{int} = 13 \text{ nb}^{-1}$, $\sqrt{s_{NN}} = 5.3 \text{ TeV}$

- D-meson production well described using the fragmentation fraction from $\rm e^+e^ \bigcirc$ measurements
- Large enhancement of all charm-baryon production in pp collisions w.r.t. $\rm e^+e^-$ collisions \bigcirc In addition to simple fragmentation, other hadronisation mechanisms are needed to \bigcirc describe the measurements in pp collisions
-
- Dependence of the fragmentation fractions on collisions system is firmly established First measurement of D_s^+ -resonance production in pp collisions at $\sqrt{s} = 13 \text{ TeV}$
- \bigcirc

ALICE Collaboration ready to analyse Run 3 data to investigate the currently open questions

Thank you for your attention!

BACKUP SLIDES

[Phys. Lett. B 829 \(2022\) 137065](https://doi.org/10.1016/j.physletb.2022.137065)

CR beyond leading colourInitial state not insensitive to strong force (coloured partons, beam remnants)

-
- $MPI \rightarrow crucial$ to explain underlying event

CR beyond Leading Color approximation (CR-BLC)

"Simplified QCD" with 9 color indices to determine the string formation

String length minimization over all possible configurations, even those beyond the Leading Color topology \rightarrow Monash: only CR among LC

Enhanced leading color among MPIs and beam remnants

- Conditions for color reconnections:
	- Invariant mass of string *j*-th must overcome a threshold $m₀$ $C = m_{0i}/m_0 > 1$: enhanced reconnections
	- Causality: two strings must resolve each other between formation and hadronization, according to the time dilation due to the relative boost
	- \rightarrow Mode 0, 2, 3: different "severity" on this condition

Statistical approaches and coalescence

SHM+RQM PLB 795 (2019) 117-121

- Hadron formation driven by the mass at a hadronization temperature $T_H \to$ stat. weights $n_i \sim m_i^2 T_H K^2 (m_i/T_H)$ Strong feed-down from an augmented set of excited charm baryon states
- - PDG: $5 \Lambda_c$, $3 \Sigma_c$, $8 \Xi_c$, $2 \Omega_c$ \bigcirc
	- RQM: additional (not yet measured) 18 Λ_{μ} \bigcirc

Quark Coalescence Mechanism (QCM) Eur. Phys. J. C (2)

- Thermal weights to account for relative production vector mesons
- Hadron p_{T} spectrum from recombination of charm quarks from the hard scattering with equal-velocity light quarks in the nearby in phase-space

Catania coalescence model PLB 821, 136622

- Thermalised system of u, d, s and gluons
- Charm quark can hadronize either via fragmentation or coalescence
- Charm hadronization into ground and (PDG) excited states
	- The latter ones increase the abundance of the former ones \circ
	- Statistical "penalty" weight $[m_{H*}/m_{H}]^{3/2}$ \times exp(- $\Delta E/T$) \circ

 \mathbf{p} \cdots

