Axion dark matter (theory and experiment)

Andreas Ringwald
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[Peccei,Quinn `77; Weinberg `78; Wilczek `78]

Solving two problems in one go

[Preskill, Wise, Wilczek 83; Abbott, Sikivie 83; Dine, Fischler 83,....]

Extensions of the Standard Model (SM) featuring the axion answer two fundamental questions in one go

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- Extensions of the Standard Model (SM) featuring the axion answer two fundamental questions in one go:
 - 1. Why do strong interactions conserve CP so accurately?

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Theta-term in QCD Lagrangian:

[Belavin et al. `75; 't Hooft 76; Callan et al. `76; Jackiw, Rebbi `76]

$$\mathcal{L}_{\mathrm{QCD}} \supset \overline{\theta} \, \frac{\alpha_s}{8\pi} G^b_{\mu\nu} \tilde{G}^{b,\mu\nu}$$

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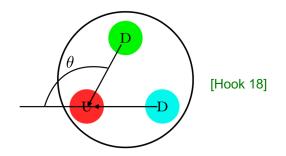
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- leads to electric dipole moment of neutron (nEDM):

[Crewther, Di Vecchia, Veneziano, Witten 79;...; Pospelov, Ritz 00]

$$d_n \sim \frac{m_u m_d}{m_u + m_d} \frac{1}{m_n^2} \,\bar{\theta} \,e \sim 10^{-16} \,\bar{\theta} \,e \,\mathrm{cm}$$



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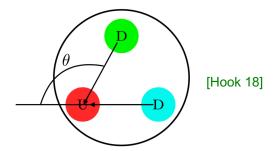
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• Experiment: [Abel et al. 20]

$$|d_n| < 1.8 \times 10^{-26} e \,\mathrm{cm}$$



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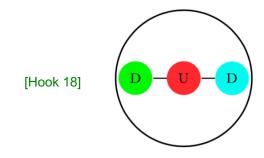
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 $\Rightarrow |\overline{\theta}| \lesssim 10^{-10}$



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- Extensions of the Standard Model (SM) featuring the axion answer two fundamental questions in one go:
 - 1. Why do strong interactions conserve CP so accurately?
 - 2. What is the nature of dark matter (DM)?



Minimal SM extension solving strong CP problem

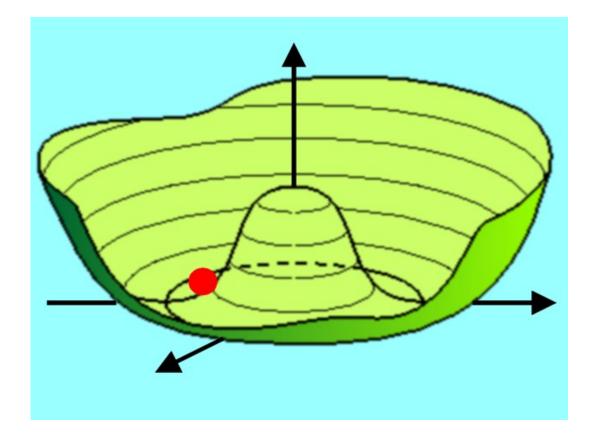
[Kim 79;Shifman,Vainshtein,Zakharov 80]

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• Add to SM a singlet complex scalar field σ , featuring a spontaneously broken global axial $U(1)_{PQ}$ symmetry, and an exotic quark $\mathcal Q$ charged under it:

$$\mathcal{L} \supset -\lambda_{\sigma} \left(|\sigma|^2 - \frac{v_{\mathrm{PQ}}^2}{2} \right)^2 + y \, \sigma \, \bar{\mathcal{Q}}_L \mathcal{Q}_R$$



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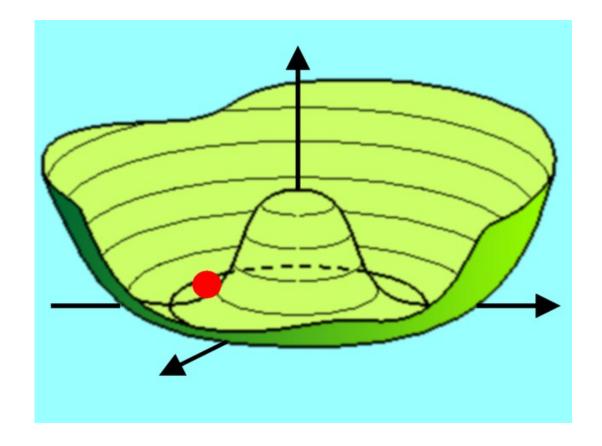
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- Features three particles beyond SM:
 - Excitation of angular field a: axion (NG-boson)

 - Excitation of modulus field ρ : $m_{\rho} = \sqrt{2\lambda_{\sigma}}v_{\mathrm{PQ}}$ Exotic quark: $m_{\mathcal{Q}} = \frac{y}{\sqrt{2}}v_{\mathrm{PQ}}$ $\left[\sigma(x) = \frac{1}{\sqrt{2}}\left(v_{\mathrm{PQ}} + \rho(x)\right)\mathrm{e}^{ia(x)/v_{\mathrm{PQ}}}\right]$

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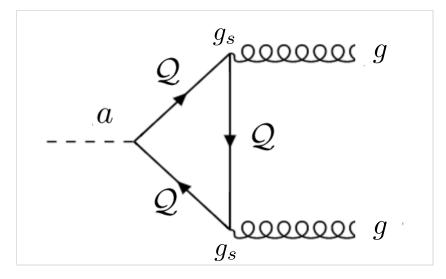
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[Chadha-Day, Ellis, Marsh, Science Advances 22]

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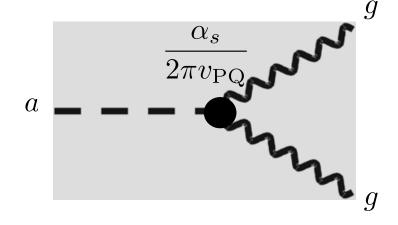
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dynamics of the axion is described by
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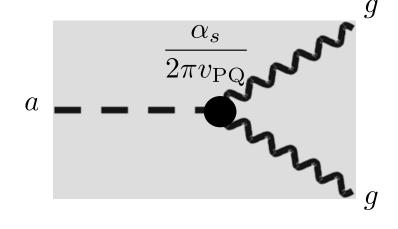
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 $\overline{ heta}$ -parameter can be eliminated: $a(x) o a(x) - \overline{ heta} \, v_{\mathrm{PQ}}$



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 m PQ}$
- Integrating out gluons (matching to chiral Lagrangian) results in non-trivial effective potential for axion field

Effective axion potential at energies below confinement scale [Di Vecchia, Veneziano `80; Leutwyler, Smilga 92]

$$V(\theta) = m_{\pi}^2 f_{\pi}^2 \left(1 - \frac{\sqrt{1 + z^2 + 2z \cos \theta}}{1 + z} \right)$$

$$z \equiv m_u / m_d \approx 1/2$$

predicts

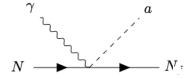
- $\langle \theta(x) \rangle = 0 \Rightarrow \mathsf{nEDM} \mathsf{vanishes}$
- axion mass:

$$m_a = \frac{\sqrt{V''(0)}}{v_{PQ}} = \frac{\sqrt{z}}{1+z} \frac{m_\pi f_\pi}{v_{PQ}} \approx 6 \ \mu \text{eV} \left(\frac{10^{12} \,\text{GeV}}{v_{PQ}}\right)$$

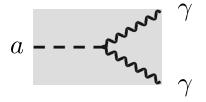
$$\mathcal{L}\supset -rac{1}{2}m_a^2a^2$$

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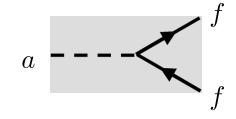
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Effective field theory below QCD scale

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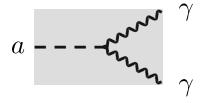
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$$f$$

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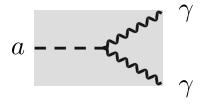


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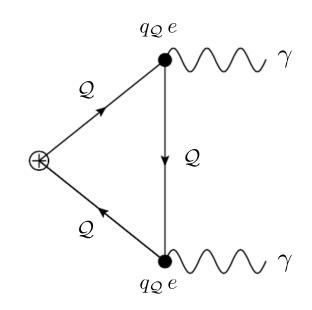
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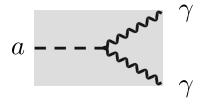
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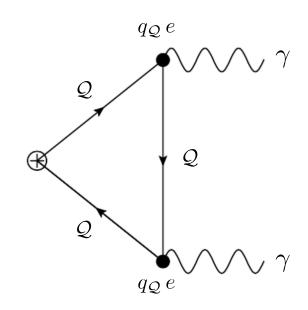
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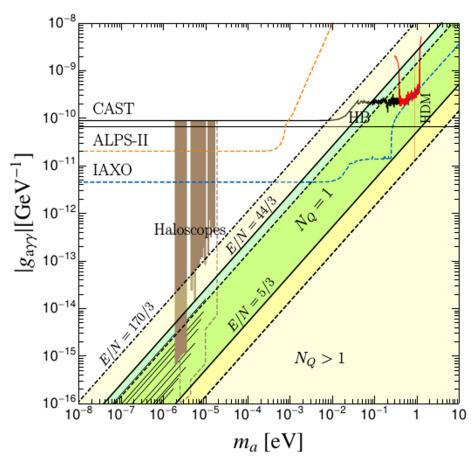
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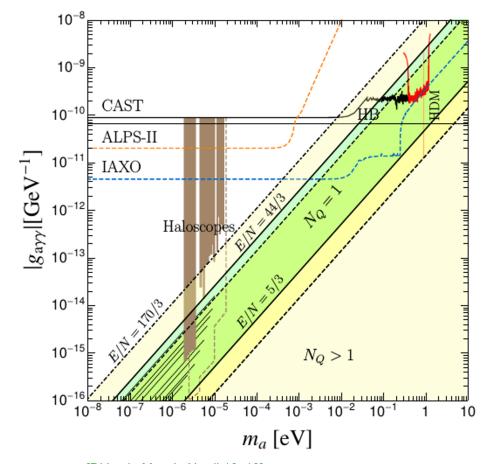


[Di Luzio, Mescia, Nardi 16, 18]

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- What if exotic quark carries a magnetic charge? [Sokolov,AR 21, 22, 23]
 - Its existence would explain not only strong CP conservation, but also charge quantisation [Dirac 1931, Schwinger 1966, Zwanziger 1968]



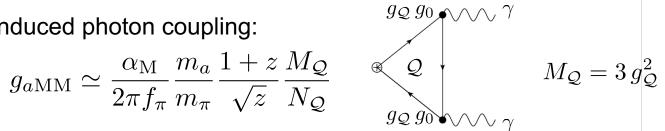
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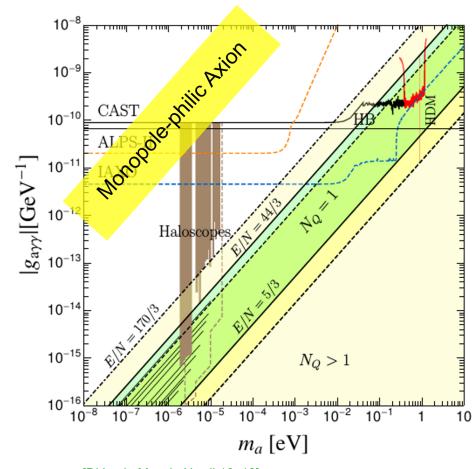
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 - Induced photon coupling:

$$g_{a{\rm MM}} \simeq \frac{\alpha_{\rm M}}{2\pi f_{\pi}} \frac{m_a}{m_{\pi}} \frac{1+z}{\sqrt{z}} \frac{M_{\mathcal{Q}}}{N_{\mathcal{Q}}}$$



strongly enhanced due to charge quantisation ($e g_0 = 6\pi n$):

$$\alpha_{\rm M} \equiv \frac{g_0^2}{4\pi} = \frac{9\pi}{e^2} = \frac{9}{4}\alpha^{-1} \quad \Rightarrow \quad \frac{g_{\rm aMM}}{g_{a\gamma\gamma}} = \left(\frac{3\,g_{\mathcal{Q}}}{2\,q_{\mathcal{Q}}}\right)^2 \alpha^{-2} \sim 10^5$$

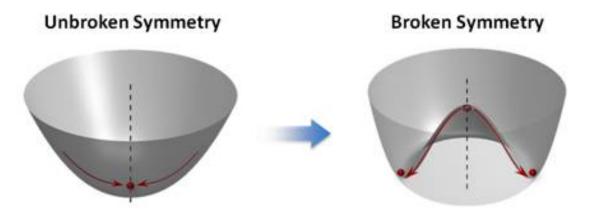


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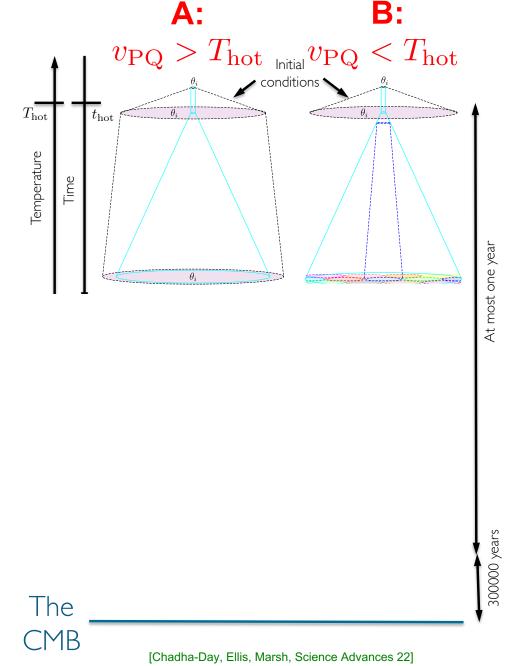
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[Peking University]

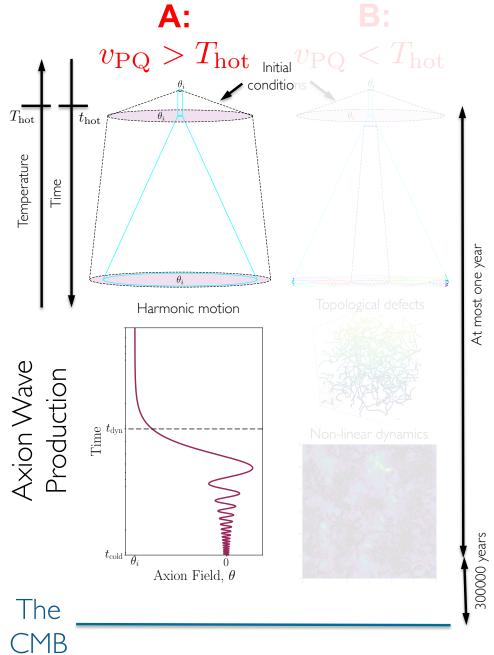
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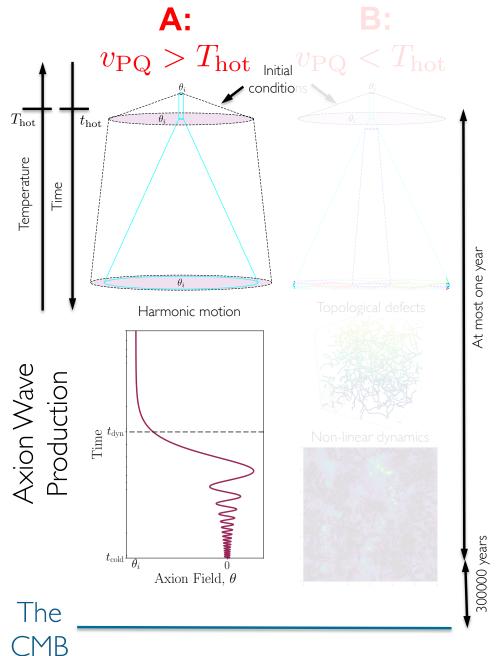
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 [Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,....]
 - DM axion mass not fixed by DM abundance:

$$\Omega_a h^2 pprox 0.12 \, \left(rac{6 \; \mu {
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ight)^{1.165} heta_i^2 \,, \qquad heta_i \equiv a(t_{
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 [Borsanyi et al. `16]

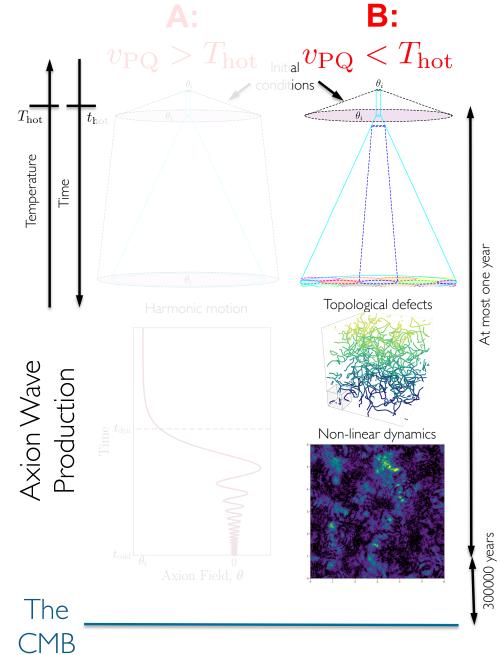


Standard production mechanisms

- Axion comes to existence as soon as Peccei-Quinn (PQ) symmetry is spontaneously broken
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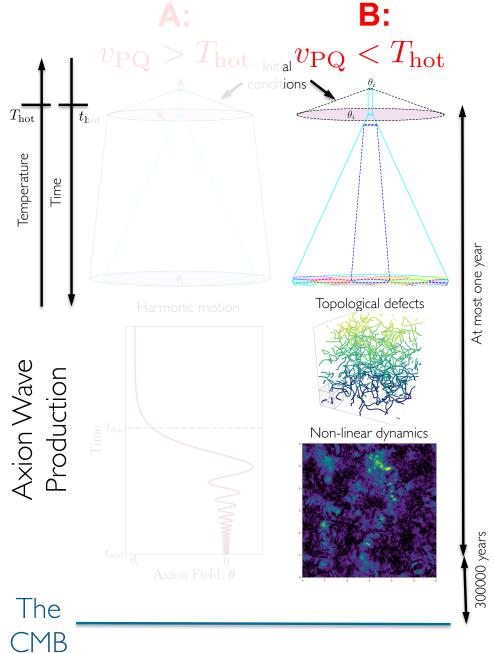
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- Scenario B ("post-inflationary PQ breaking"):
 - axion DM production by realignment mechanism and decay of topological defects (strings and domain walls)
 - required axion mass to explain 100% of DM abundance:

[Hiramatsu et al. 11,12,13; Kawasaki,Saikawa,Segikuchi 15; AR,Saikawa `16; Borsanyi et al. `16; Klaer,Moore `17; Gorghetto,Hardy,Villadoro `18; Buschmann et al. 19; Hindmarsh 19; Gorghetto,Hardy,Villadoro '20; Buschmann et al. 21;Beyer,Sarkar 23;...]

$$m_a \approx 26 \ \mu\text{eV} - 0.5 \ \text{meV}, \text{ for } N_{\text{DW}} = 2N_{\mathcal{Q}} = 1$$

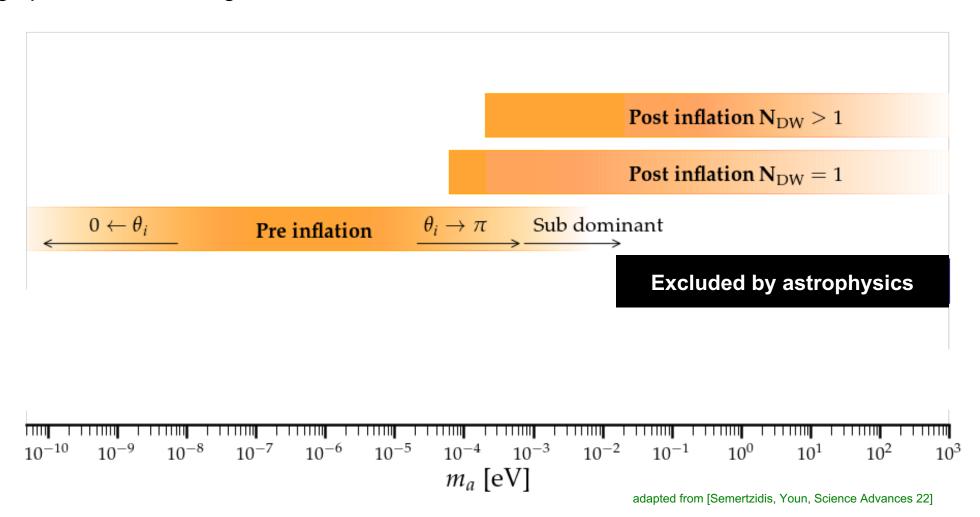
 $m_a \gtrsim \text{meV}, \text{ for } N_{\text{DW}} = 2N_{\mathcal{Q}} > 1$



Axion Dark Matter Experiments

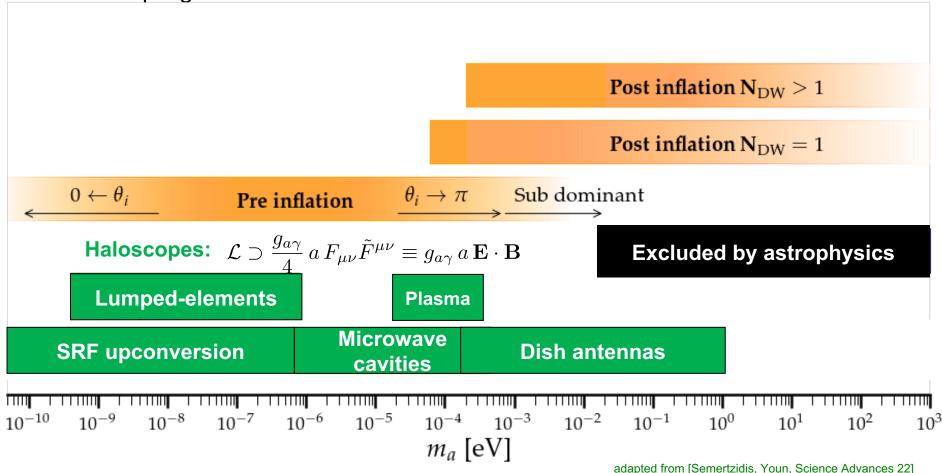
Variety of experimental techniques

Huge possible mass range:



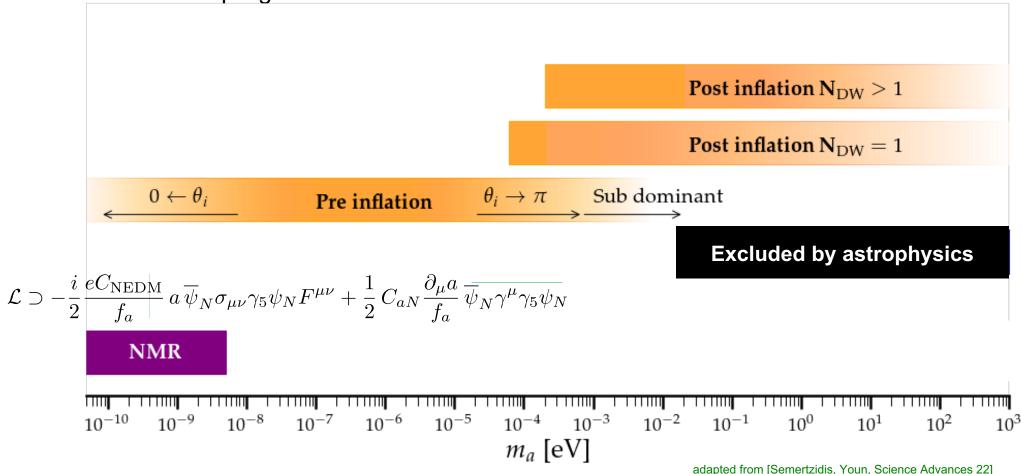
Variety of experimental techniques

 Huge possible mass range requires various experimental techniques to search for axion dark matter via different axion couplings:

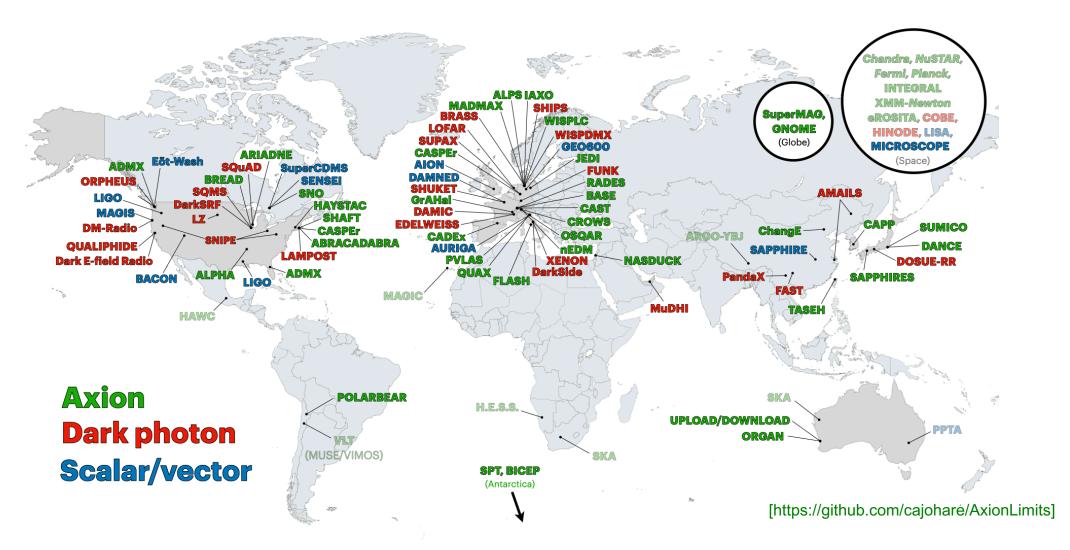


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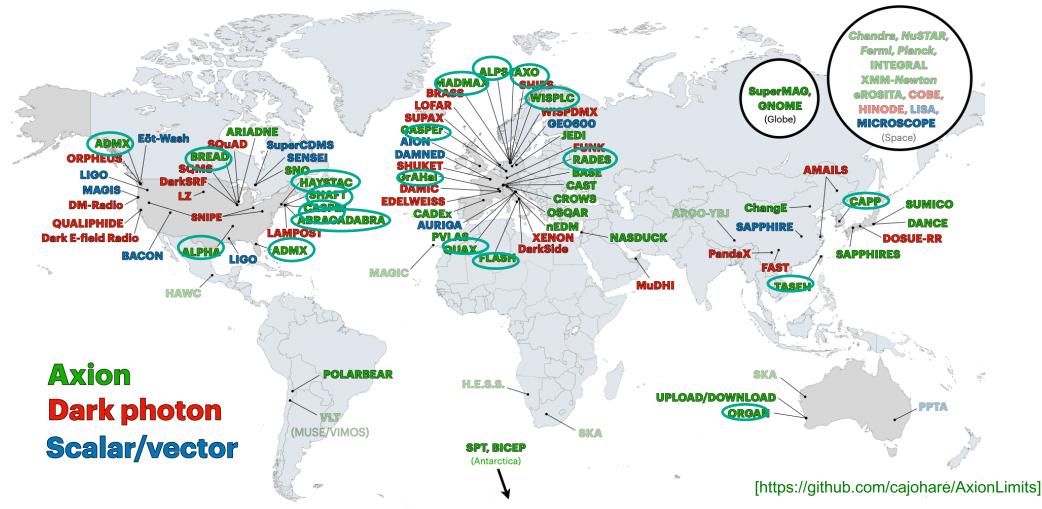


Enormous number of axion experiments worldwide



Enormous number of axion experiments worldwide

Can only talk about a subset of them ...



Wavy dark matter

- Axion DM experiments rely on the assumption that the dark matter halo of the Milky Way is comprised by axions
- The velocity dispersion of halo dark-matter axions is given by the galactic virial velocity, implying a macroscopic de Broglie wave length,

$$\lambda_{\rm dB} = 2\pi/(m_a v_a) \simeq {\rm km} \, (\mu {\rm eV}/m_a) (10^{-3} \, c/v_a)$$

Correspondingly, halo dark-matter axions behave as an approximately spatially homogeneous and monochromatic classical oscillating field,

$$a(t) \simeq \sqrt{2\rho_a} \cos(m_a t)/m_a$$

When presenting (projected) limits on halo dark-matter axion couplings it is assumed that

$$\rho_a = \rho_{\rm DM}^{\rm halo} \approx 0.45 \, {\rm GeV \, cm^{-3}}$$

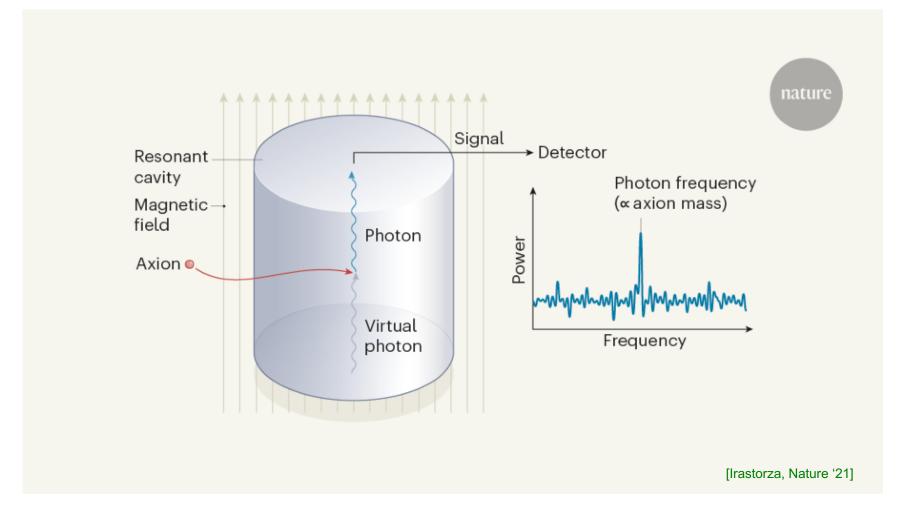
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Microwave cavities

• Concept: In microwave cavity placed in magnetic field, DM axion converts into photon

[Sikivie 83]

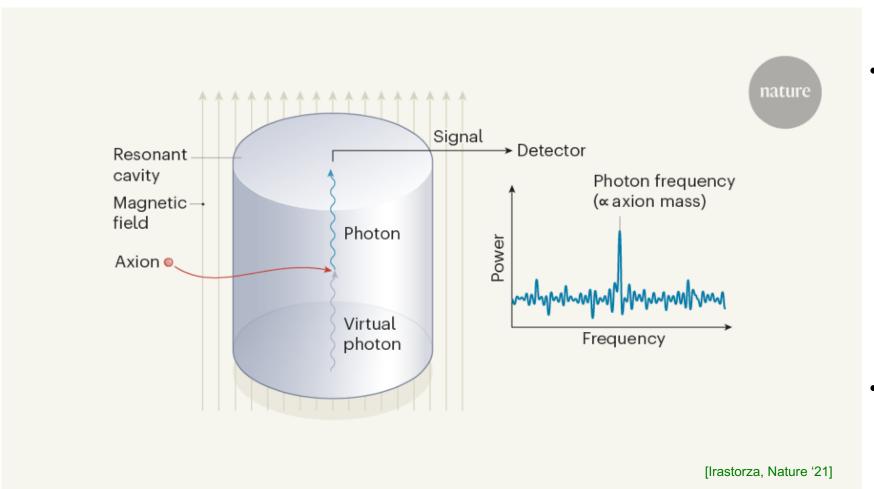


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[Sikivie 83]



 If axion mass matches resonance frequency of cavity,

$$m_a = 2\pi\nu_{\rm res} \sim 4\,\mu{\rm eV}\left(\frac{\nu_{\rm res}}{{\rm GHz}}\right)$$

power output

$$P_{\mathrm{out}} \sim g_{a\gamma}^2 \, \rho_{\mathrm{a}} \, B_0^2 \, V \, Q$$

enhanced by quality factor

$$Q \sim 10^5$$

Need to scan by tuning resonance frequency

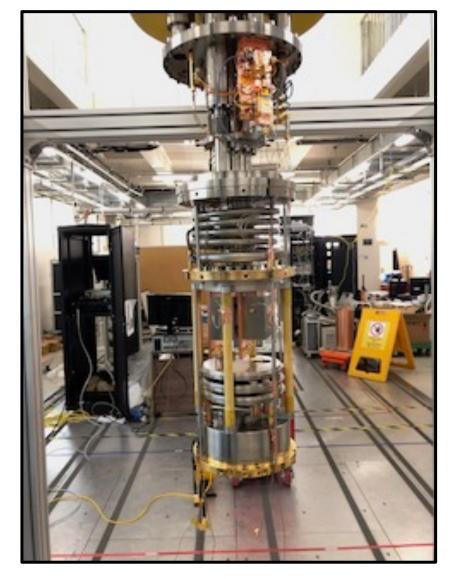
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Microwave cavities

Currently running:

- ADMX
- CAPP [Sung Woo Youn, DM par. session 2A]
- CAST-CAPP
- GrAHal
- HAYSTAC
- ORGAN
- QUAX
- RADES
- TASEH





ADMX

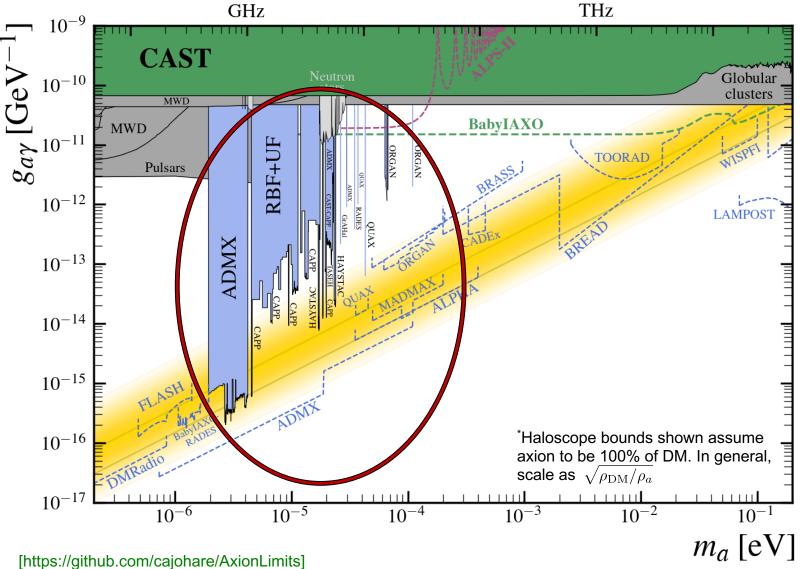
CAPP

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MWD

 10^{-10}

CAST

MWD

[https://github.com/cajohare/AxionLimits]

GHz

Microwave cavities

- Currently running:
 - **ADMX**
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 - **RADES**
 - **TASEH**
- More experiments proposed
- Within next decade, microwave cavities will dig deep into vanilla axion band in mass range

 $3a^{10-11}$ RBF+UF **TOORAD** Pulsars 10^{-12} **LAMPOS**7 **ADMX** 10^{-13} 10^{-14} 10^{-15} 10^{-16} *Haloscope bounds shown assume axion to be 100% of DM. In general, scale as $\sqrt{\rho_{\rm DM}/\rho_a}$ 10^{-17} 10^{-5} 10^{-4} m_a [eV]

THz

BabyIAXO

$$\mu eV \lesssim m_a \lesssim 100 \,\mu eV$$

Globular clusters

$$\mathcal{L} \supset \frac{g_{a\gamma}}{4} \, a \, F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma} \, a \, \mathbf{E} \cdot \mathbf{B}$$

Dish antennas

Dish antenna concept:

[Horns, Jaeckel, Lindner, Lobanov, Redondo, AR 13]

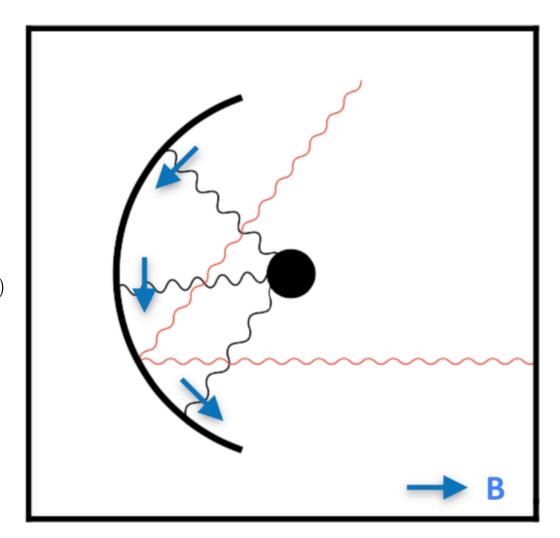
 Oscillating axion DM in a background magnetic field carries a small electric field component,

$$\vec{E}_a(t) = -g_{a\gamma}\vec{B}a(t)$$

• Metallic mirror placed in a magnet field pointing parallel to the mirror surface will emit a nearly monochromatic EM wave perpendicular to the mirror surface with a frequency $\nu=m_a/(2\pi)$ and a cycle-averaged power per unit area:

$$\mathcal{P}/\mathcal{A} \simeq 10^{-27} \frac{\text{W}}{\text{m}^2} \left(\frac{g_{a\gamma\gamma}}{10^{-14} \,\text{GeV}^{-1}} \right)^2 \left(\frac{10^{-4} \,\text{eV}}{m_a} \right)^2 \left(\frac{B}{10 \,\text{T}} \right)^2$$

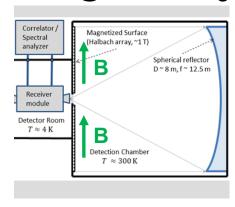
Broadband! No tuning necessary!



 $\mathcal{L} \supset \frac{g_{a\gamma}}{4} \, a \, F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma} \, a \, \mathbf{E} \cdot \mathbf{B}$

Dish antennas

BRASS @ U Hamburg



 Plane permanently magnetized conversion panel

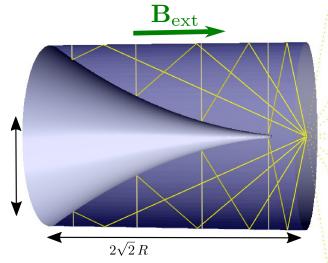
$$B = 0.8 \,\mathrm{T}$$
$$A = 4.7 \,\mathrm{m}^2$$

BRASS

Spherical reflector

[Bajjali et al., `23]



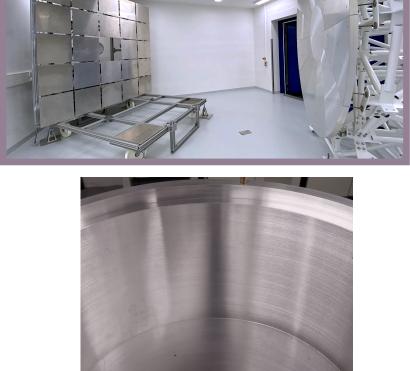


 Cylindric parabolic conversion panel allows use of solenoidal magnetic field

$$B \sim 10 \, \mathrm{T}$$

$$\mathcal{A} \sim 10\,\mathrm{m}^2$$

[Liu et al., 22]

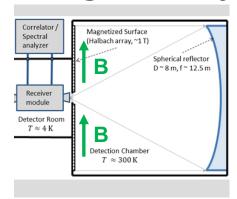




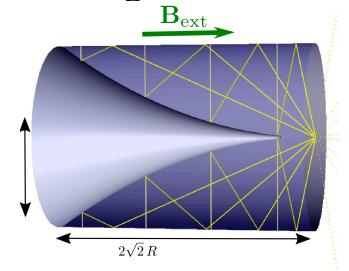
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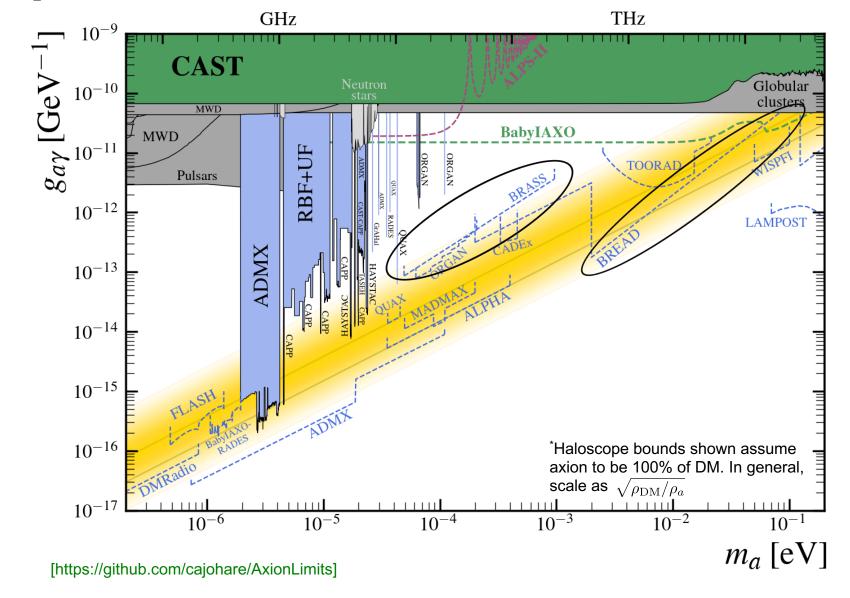
Dish antennas

BRASS @ U Hamburg



BREAD @ Fermilab





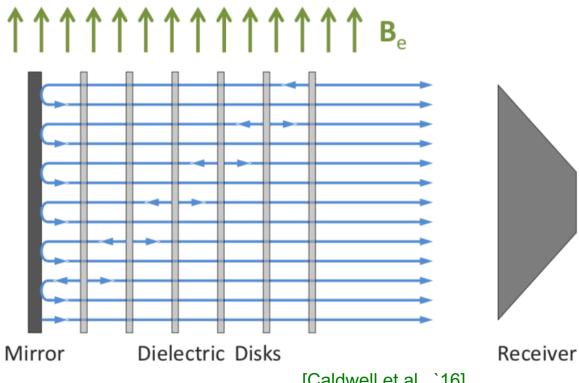
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Dish antennas

- Boosted dish antenna aka open dielectric resonator concept:
 - Add stack of dielectric disks with $\sim \lambda/2$ spacing in front of mirror (all immersed in magnetic field) [Jaeckel,Redondo 13]
 - Constructive interference of photon part of wave function

[Millar, Raffelt, Redondo, Steffen 16]

[Baryakhtar, Huang, Lasenby 18]



[Caldwell et al. `16]

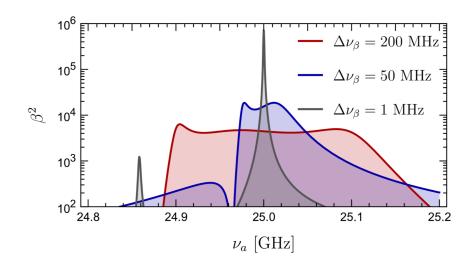
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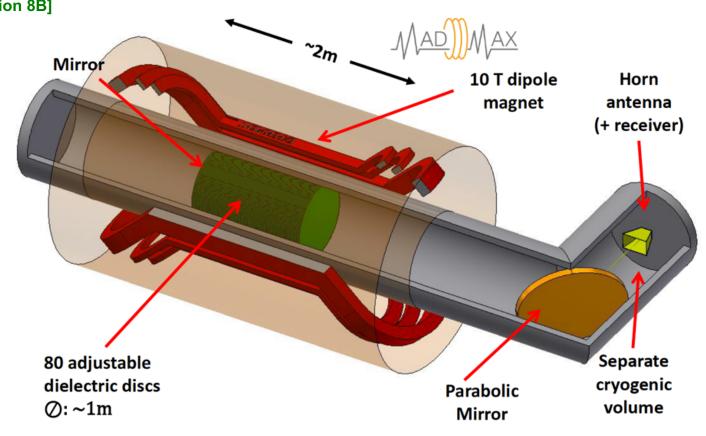
Dish antennas

MADMAX @ DESY [Juan Maldonado, poster session A]
[Juan Maldonado, DM parallel session 8B]

Conceptional design [Bruns et al. 19]

- 10 T magnet
- Large number of adjustable dielectric disks
 - Tunable frequency and bandwidth



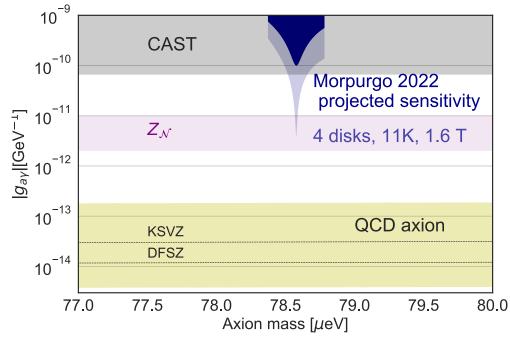


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- New prototype magnet could allow further exploration of so far unexplored parameter range





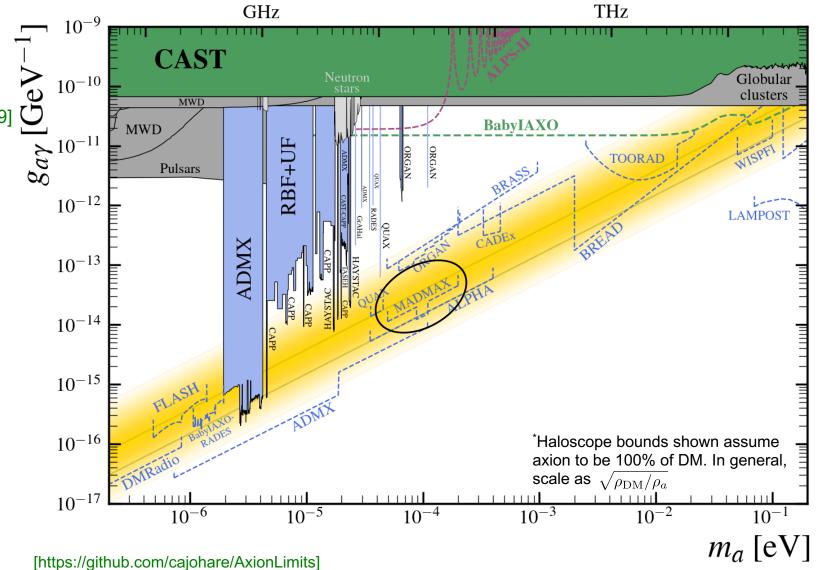
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- Full MADMAX expected to start data taking in 2030



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Plasma haloscope

Plasma haloscope concept:

[Lawson, Millar, Pancaldi, Vitagliano, Wilczek, 19]

 In a magnetized plasma, oscillating axion DM induces plasmon excitations,

$$\mathbf{E} = -g_{a\gamma}\mathbf{B}_{e}a\left(1 - \frac{\omega_{p}^{2}}{\omega_{a}^{2} - i\omega_{a}\Gamma}\right)^{-1}$$

 resonant enhancement, when plasma frequency matches axion mass,

$$\omega_p = \omega_a \approx m_a$$

limited by losses (□)

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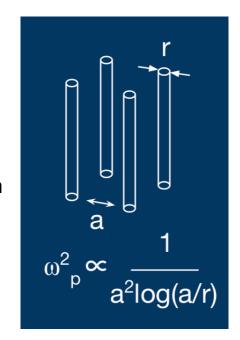
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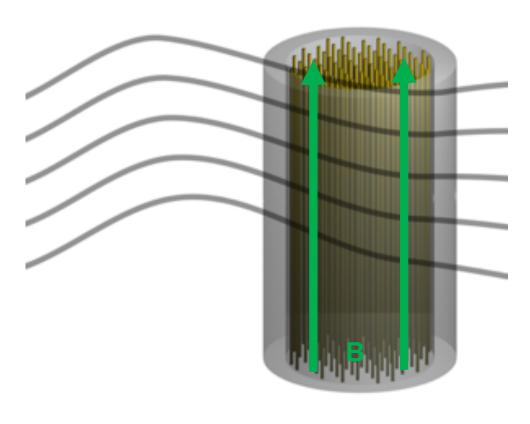
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- limited by losses (Γ)
- A plasma with tunable plasma frequency in the GHz range can be realised by a wire array with variable interwire spacing ("wire metamaterial")





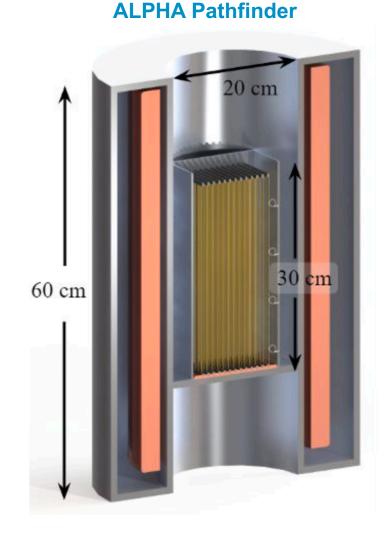
adapted from [Dunne, LBNL Instr. Sem. Feb 22]

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Plasma haloscope

ALPHA @ Oakridge

- Goal: building tunable, cryogenic plasma haloscope
- ALPHA Pathfinder
 - 20 cm bore 8T solenoid
 - several wire material prototypes constructed
 - tuning with low power piezoelectric translation
 - sensitive to mass range around 10–20 GHz
 - data run expected ~2026





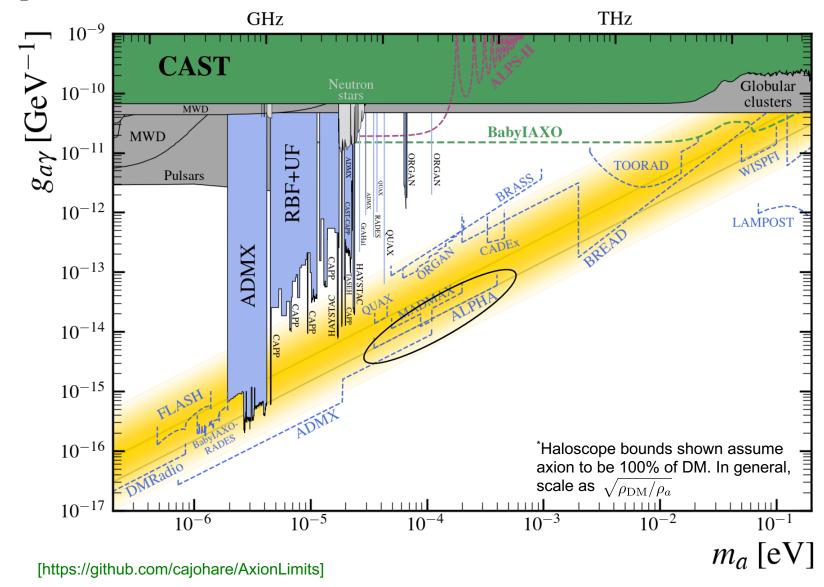
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- Full ALPHA designed to dig deep into axion band



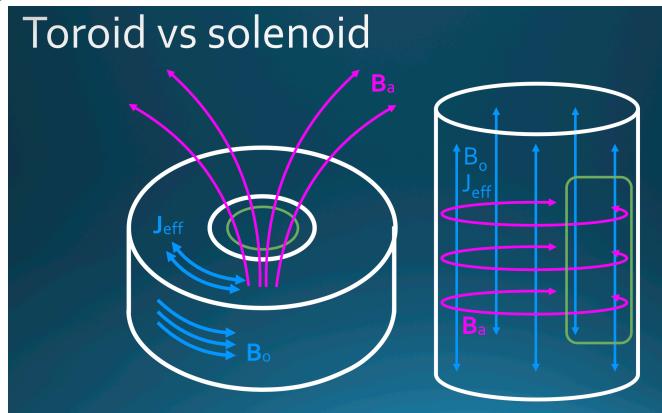
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Lumped-element detectors

[Sikivie, Sullivan, Tanner 14; Kahn, Safdi, Thaler `16]

Concept:

- Toroidal (solenoidal) magnet with fixed field B₀:
 - Axion DM generates oscillating effective current
 J_{eff} parallel to B₀
 - ... generating oscillating magnetic flux B_a through center (azimuthal magnetic flux)
 - ... which can be read out by pickup structure



$$\mathbf{J}_{eff} = g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B}$$

[Salemi '21]

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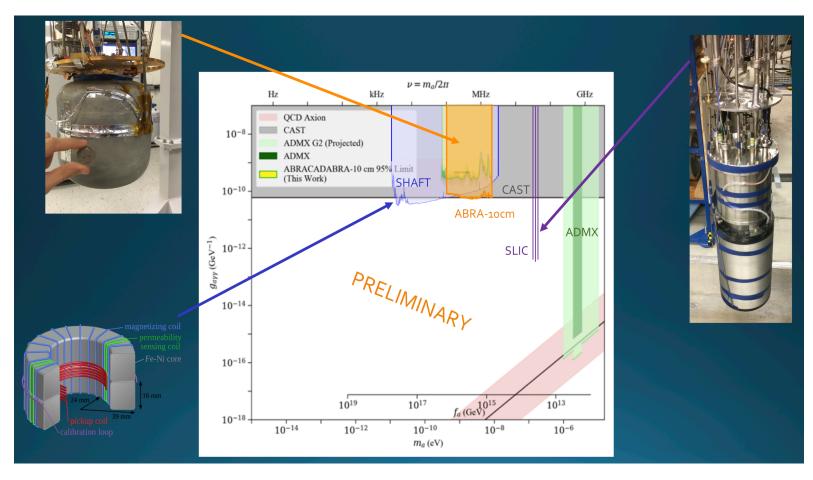
Lumped-element detectors

Pilot experiments:

ABRACADABRA [Ouellet et al. 19]

ADMX SLIC [Crisosto et al. 20]

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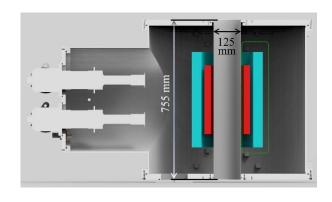


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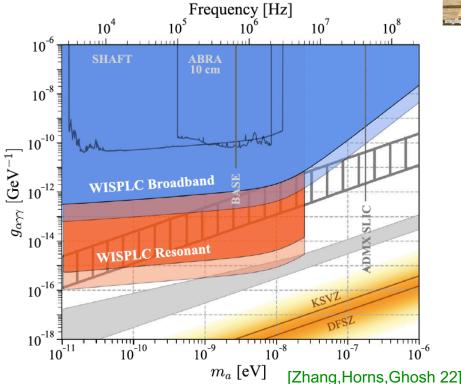
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- Pilot experiments:
 - ABRACADABRA [Ouellet et al. 19]
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- Next generation soon at start:
 - WISPLC [Zhang, Horns, Ghosh 22]



WISPLC





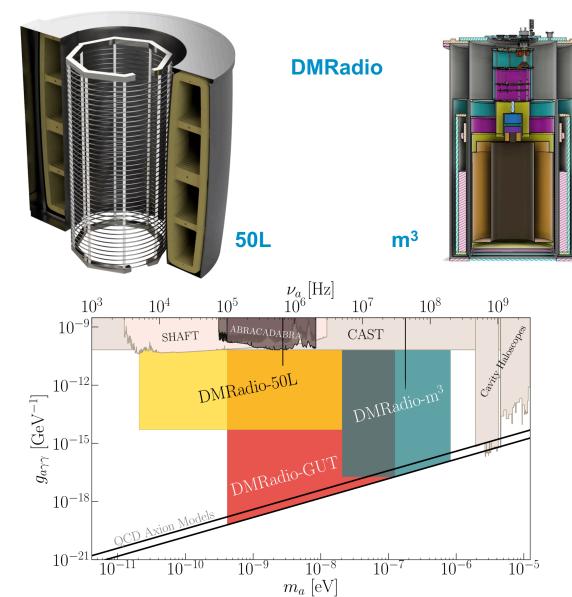
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[Chiara Salemi, DM par. session 2A] [Jessica Fry, DM par. session 2A] [Nicholas Rapidis, DM par. session 2A]

- DMRadio-50L
- DMRadio-m³
- DMRadio-GUT



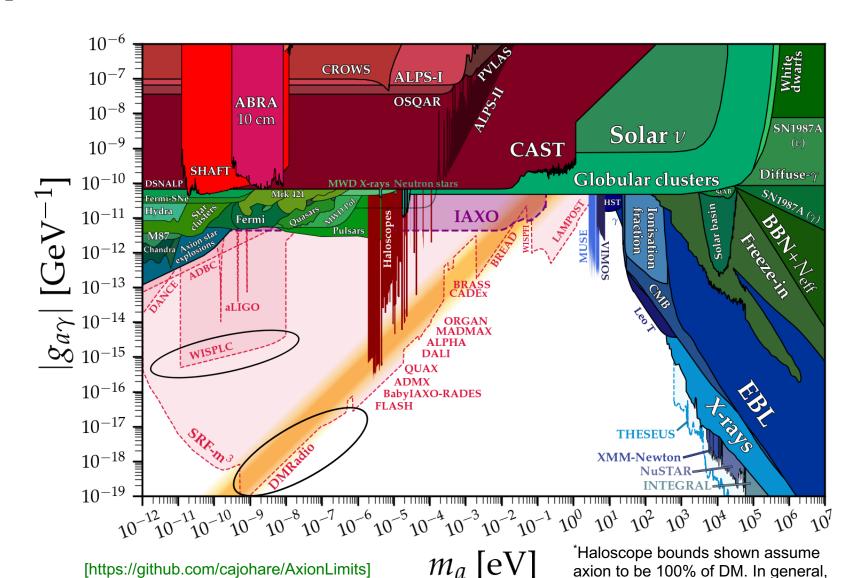
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- DMRadio-50L
- DMRadio-m³
- DMRadio-GUT
- Digging deep into unexplored territory at low mass

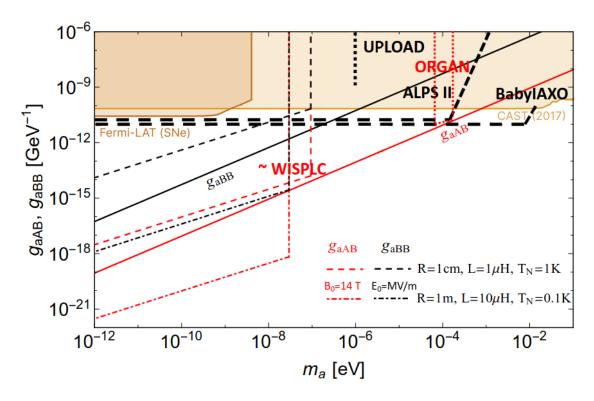


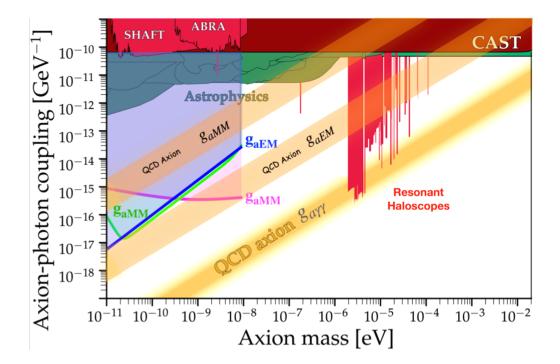
scale as $\sqrt{\rho_{\rm DM}/\rho_a}$

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Lumped-element detectors

Lumped-element detectors can also probe the new couplings of the monopole-philic axion:





*Haloscope bounds shown assume adapted from [Li,Zhang,Dai, 2211.06847] axion to be 100% of DM. In general, scale as $\sqrt{\rho_{\rm DM}/\rho_a}$

[Tobar, Sokolov, AR, Goryaev, 2306.13320]

 $\text{NMR Experiments} \qquad \mathcal{L} \supset -\frac{i}{2} \frac{e C_{\text{NEDM}}}{f_a} \, a \, \overline{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu} \equiv -\frac{i}{2} g_{aN\gamma} \, a \, \overline{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu}$

Searches for oscillating NEDMs

NMR Experiments $\mathcal{L} \supset -\frac{i}{2} \frac{e C_{\mathrm{NEDM}}}{f_{z}} a \overline{\psi}_{N} \sigma_{\mu\nu} \gamma_{5} \psi_{N} F^{\mu\nu} \equiv -\frac{i}{2} g_{aN\gamma} a \overline{\psi}_{N} \sigma_{\mu\nu} \gamma_{5} \psi_{N} F^{\mu\nu}$

Searches for oscillating NEDMs

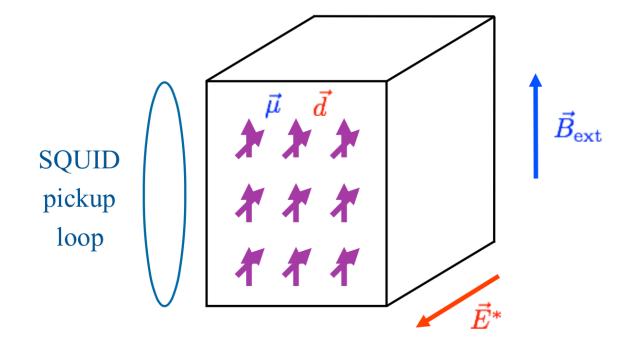
Concept:

Axion DM field induces oscillating NEDMs:

$$d_N(t) = g_{aN\gamma} \sqrt{2\rho_a} \cos(m_a t) / m_a$$

- Place a ferroelectric crystal (permanent electric polarisation fields \vec{E}^*) in external $\vec{B}_{\mathrm{ext}} \perp E^*$
- Nuclear spins are polarised along $\vec{B}_{\rm ext}$ and precess at Larmor frequency $\omega_L = 2\mu_N B_{\rm ext}$
- Interaction $\epsilon_S \vec{d}_N(t) \cdot \vec{E}^*$ of DM induced NEDM with the \vec{E}^* -field leads to resonant increase of transverse magnetisation of sample when $\omega_L=m_a$

[Graham, Rajendran 13; Budker et al. 14]

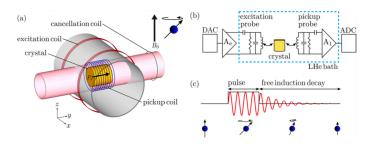


[Budker et al. 14]

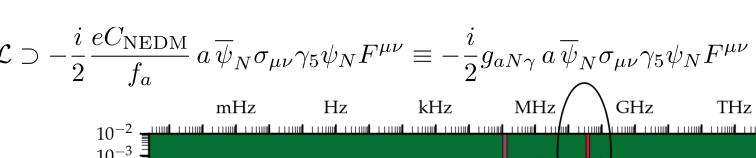
NMR Experiments

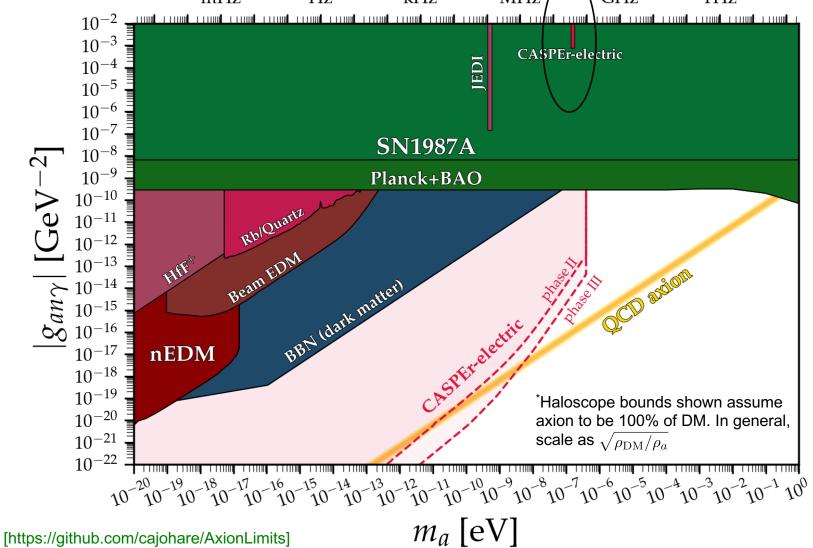
Searches for oscillating NEDMs

- CASPEr-Electric in Boston
 - Initial pathfinder experiment deep in the excluded region



[Aybas et al., 2101.01241]

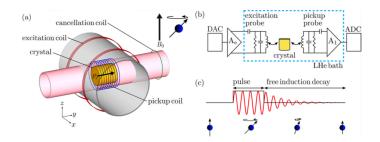




NMR Experiments

Searches for oscillating NEDMs

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 - Initial pathfinder experiment deep in the excluded region

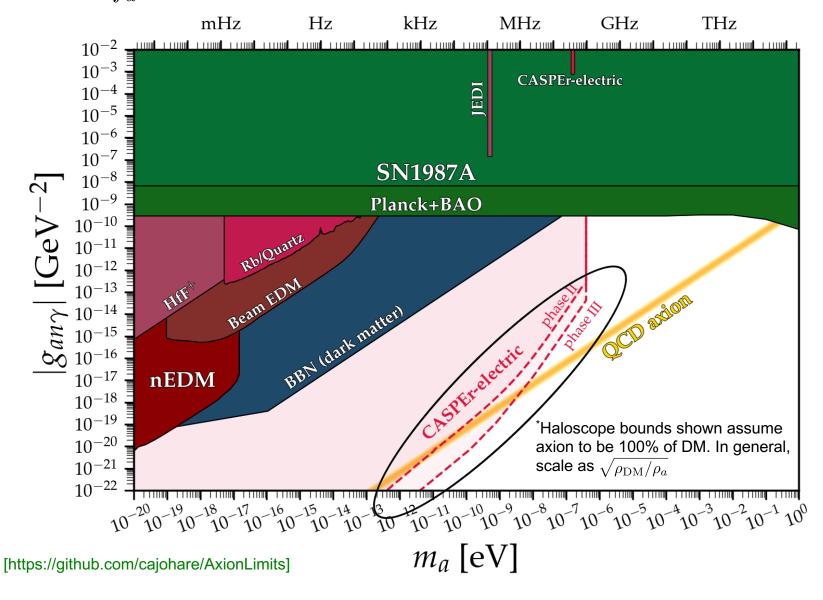


[Aybas et al., 2101.01241]

- Full experiment probes axion dark matter in mass range predicted by GUTs
- For nEDM bounds see poster session

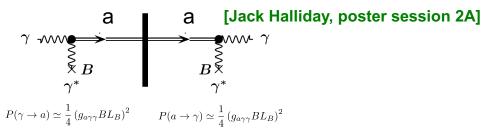
[Jacek Zejma, poster session A]

$$\mathcal{L} \supset -\frac{i}{2} \frac{eC_{\text{NEDM}}}{f_a} a \overline{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu} \equiv -\frac{i}{2} g_{aN\gamma} a \overline{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu}$$

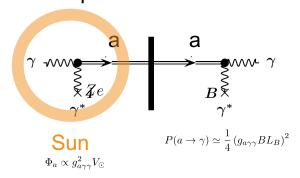


Further axion experiments

- So far only discussed experiments searching for halo DM axions
- Further experiments not relying on galactic halo being the source of axions:
 - Light-shining through walls: [Anselm 85; van Bibber 87]

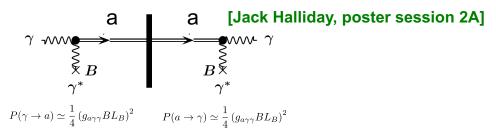


• Helioscopes: [Sikivie 83]

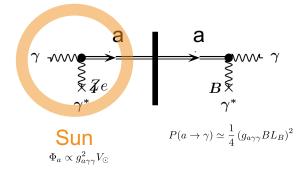


Further axion experiments

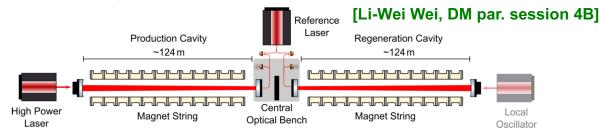
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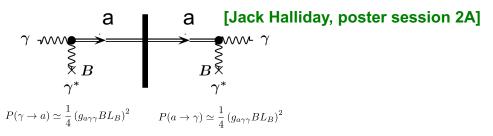
ALPS II @ DESY: data taking started in May 2023



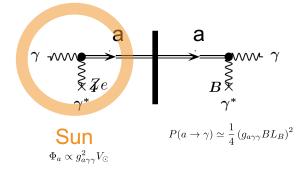


Further axion experiments

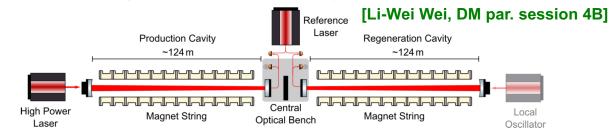
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ALPS II @ DESY: data taking started in May 2023



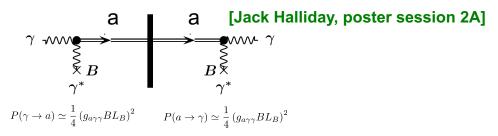


BabylAXO @ DESY: start data taking in 2028

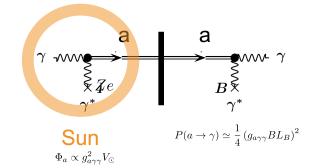


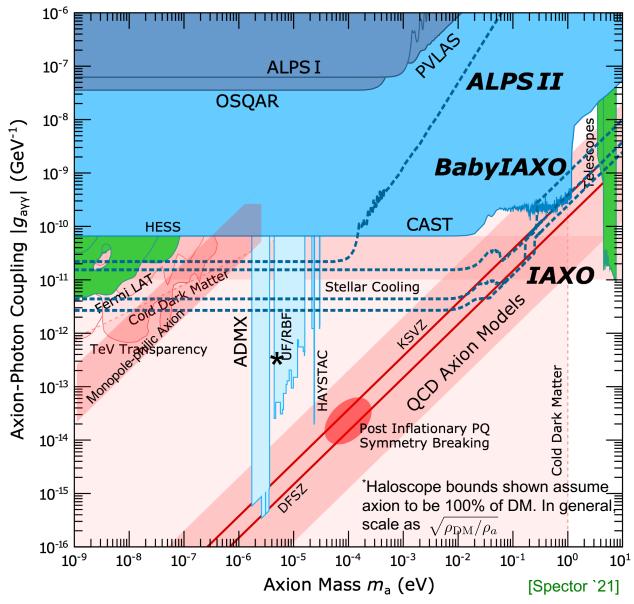
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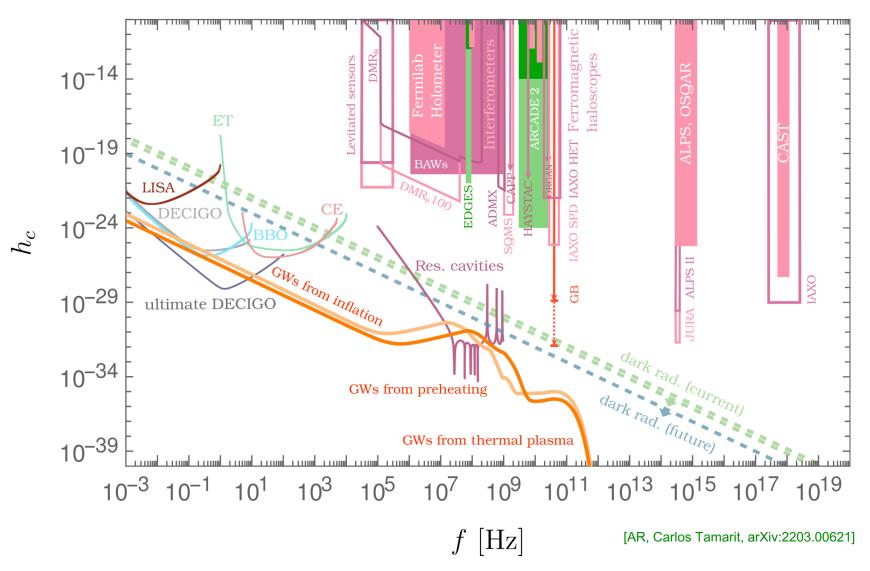
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Searches for High-Frequency Gravitational Waves

Axion haloscopes, LSW experiments, and helioscopes as HF-GW detectors



[Ejlli et al., 1908.00232]

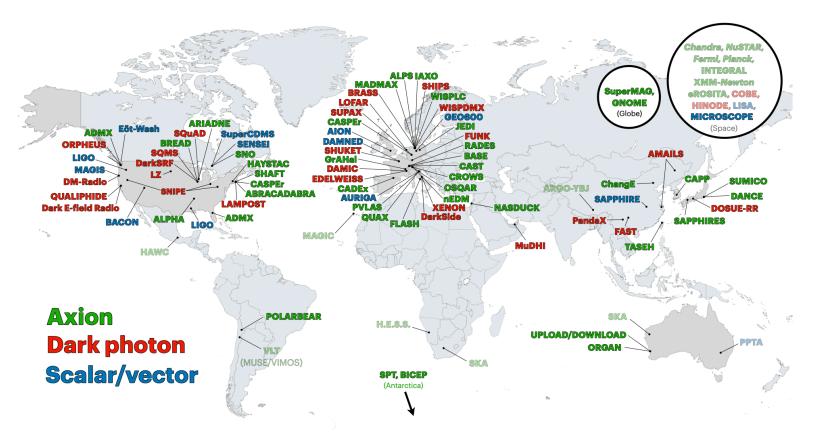
[AR et al., 2011.04731]

[Berlin et al., 2112.11465]

[Domcke et al., 2202.00695]

Summary

World-wide big experimental activity on axion searches exploiting different techniques and couplings:



• Many new experimental techniques developed ... often in very tight collaborations between phenomenologically oriented theorists and theoretically interested experimentalists

Stay tuned!