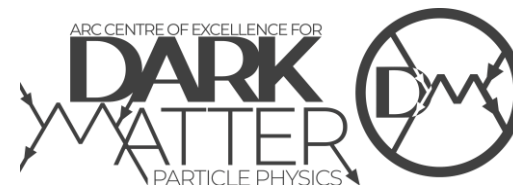


Dark Matter Searches with Astrophysics

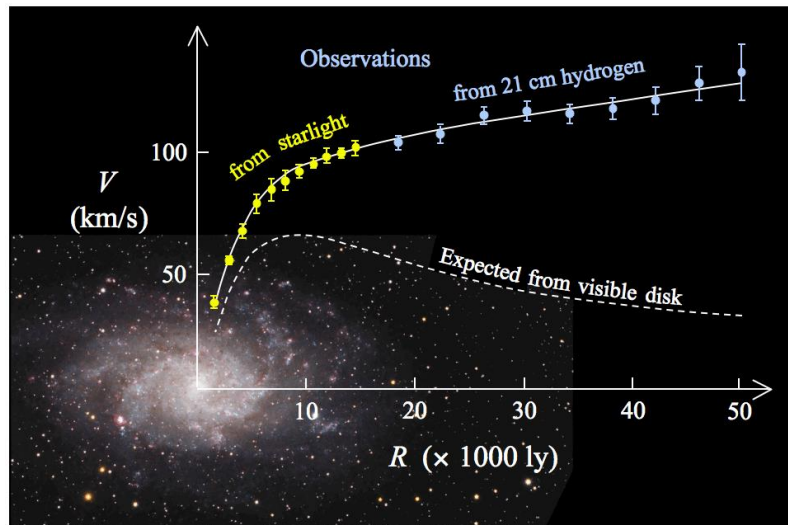
Nicole Bell

The University of Melbourne

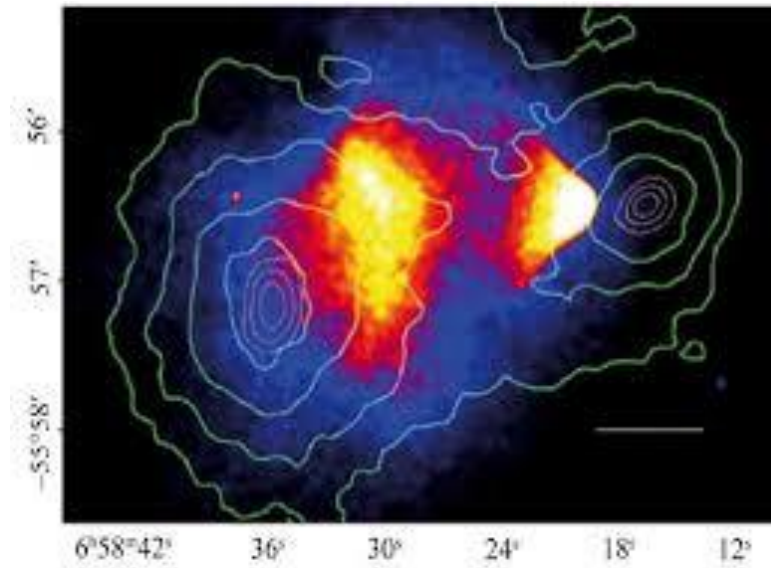


Evidence for dark matter

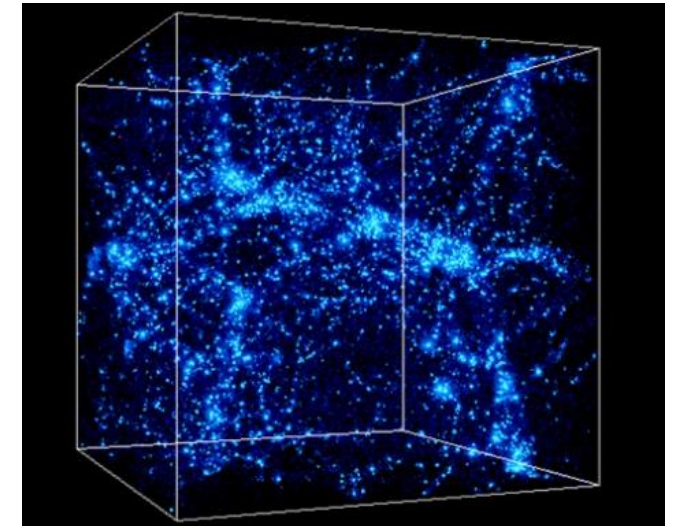
Astrophysical observations consistently point to the need for dark matter



Galaxy rotation curves



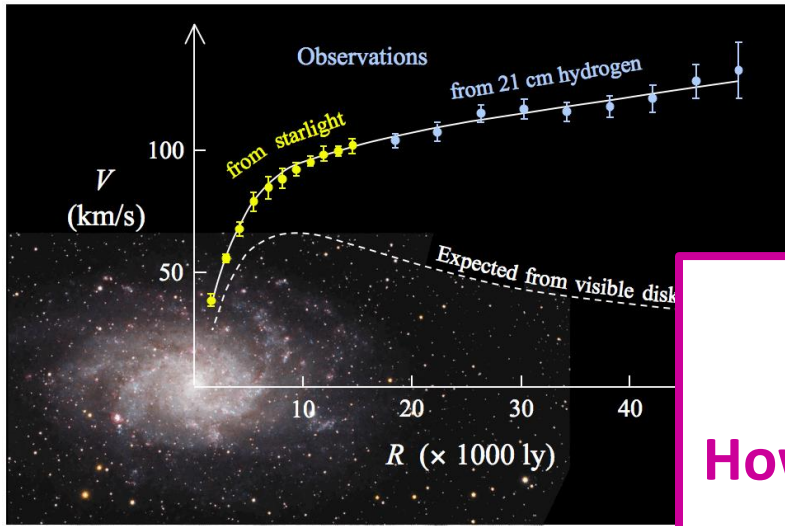
Clusters of galaxies



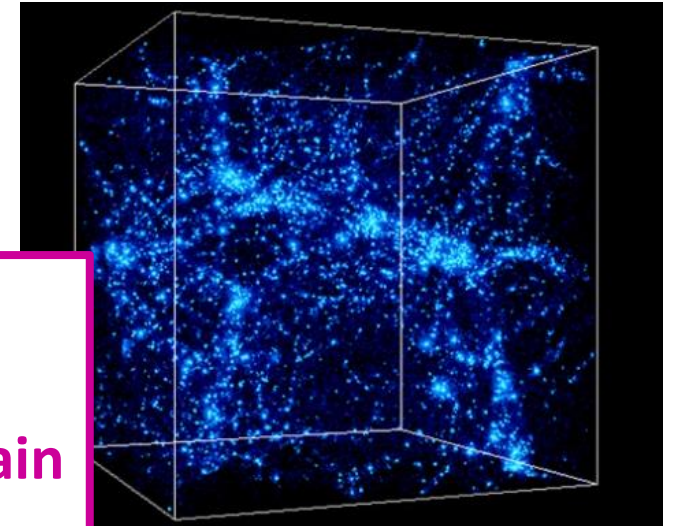
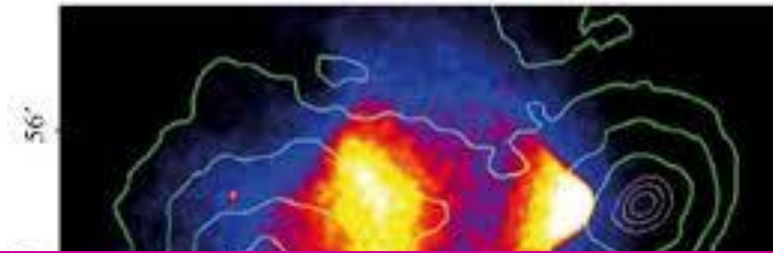
Large Scale Structure

Evidence for dark matter

Astrophysical observations consistently point to the need for dark matter



Galaxy rotation curves

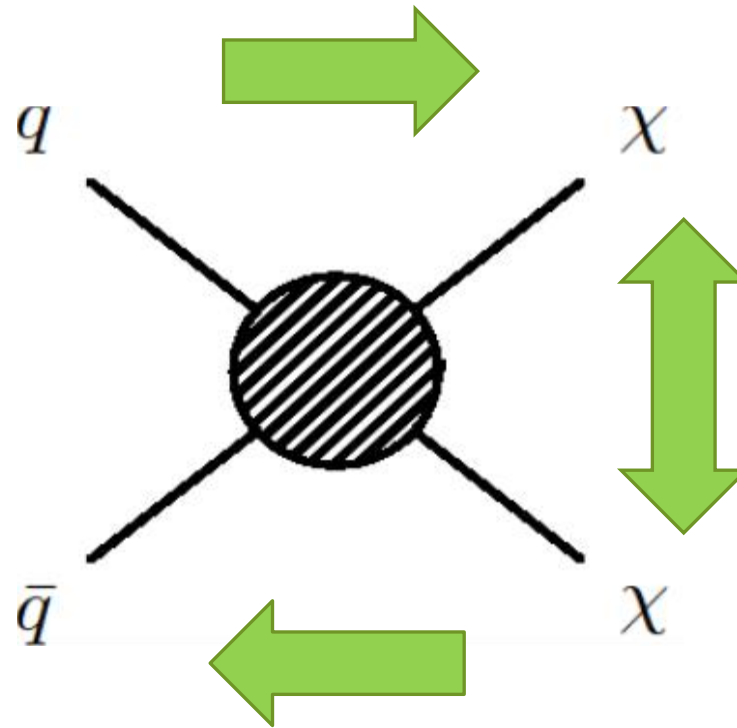


Large Scale Structure

How do we use astrophysics to constrain dark matter particle properties?

Looking for WIMP-type dark matter

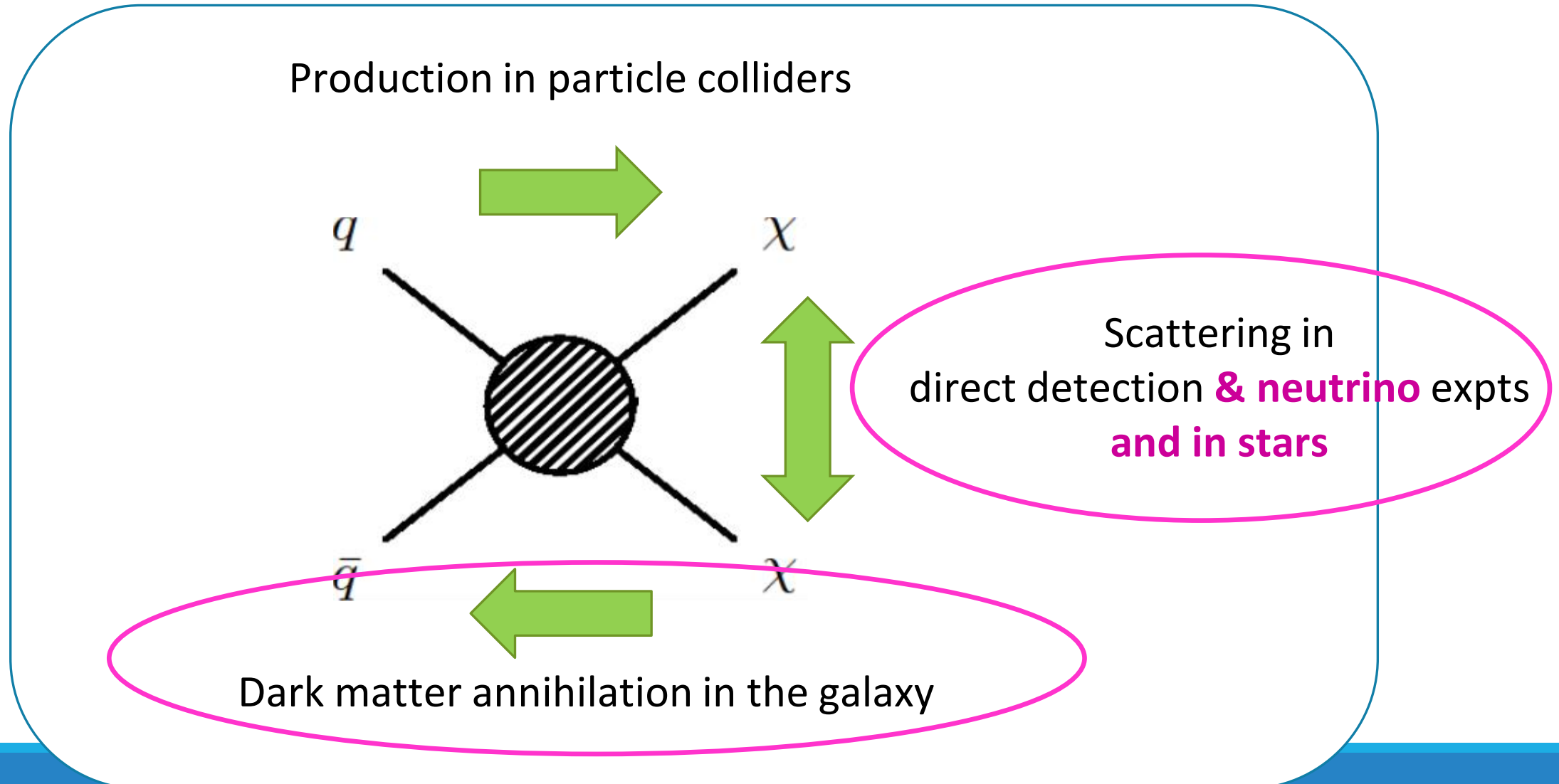
Production in particle colliders



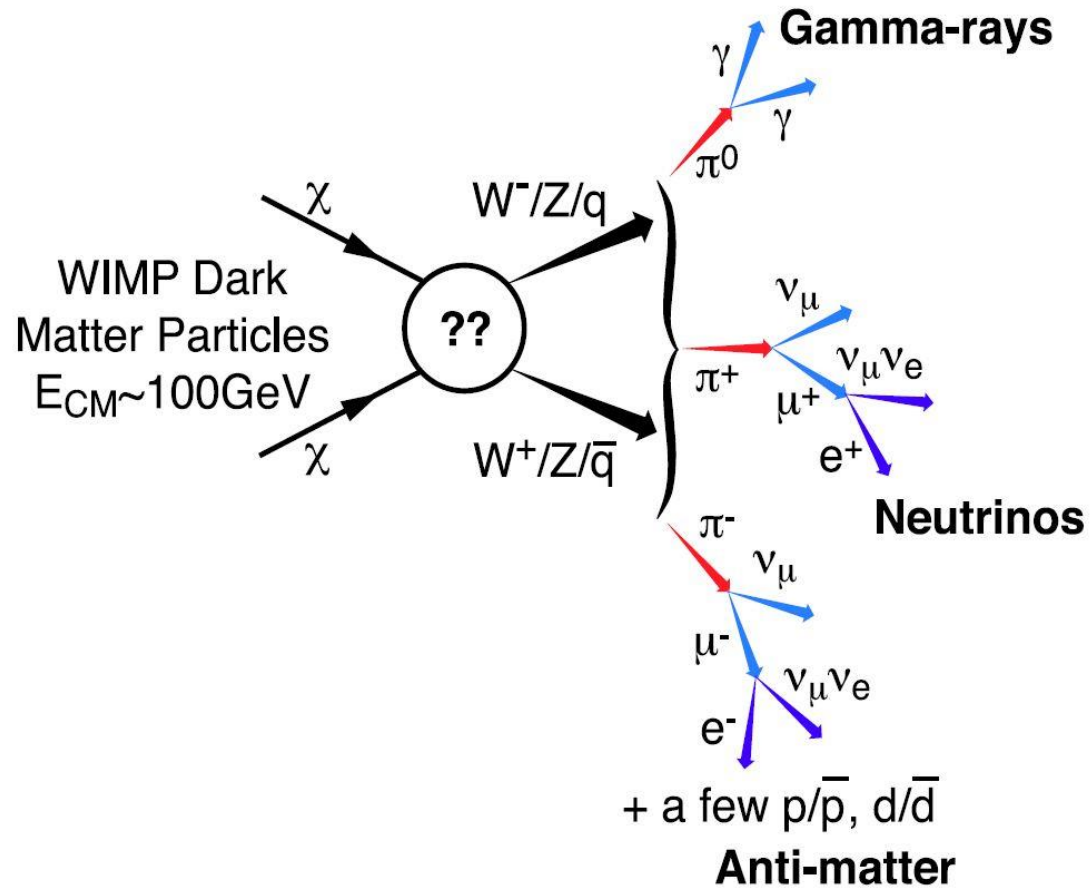
Scattering in
direct detection
experiments

Dark matter annihilation in the galaxy

Looking for WIMP-type dark matter



Indirect detection – Detecting dark matter annihilation in space



Indirect detection probes the dark matter annihilation cross-section

→ The most direct test of the thermal-relic dark matter paradigm

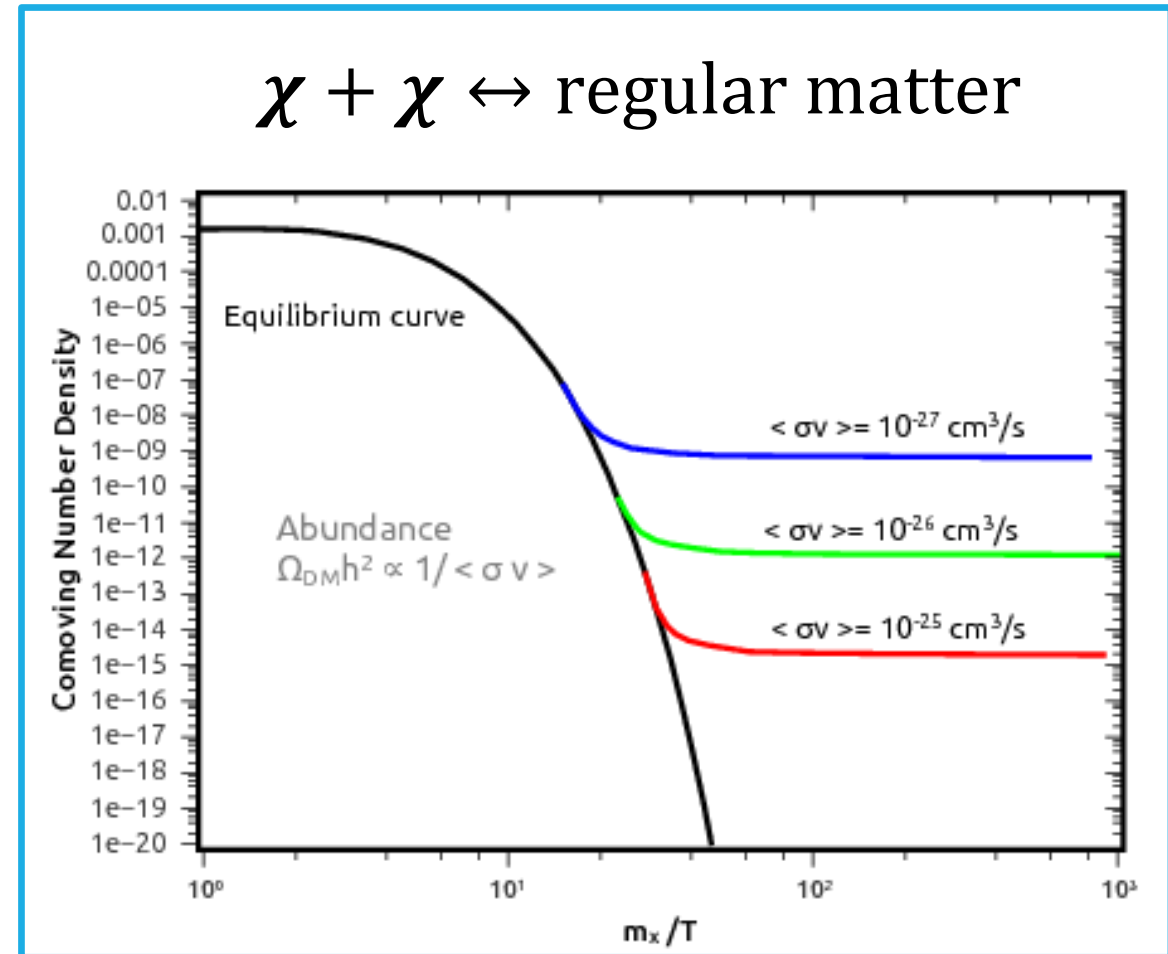
Thermal relic cross section (the WIMP miracle):

Relic DM density determined by the annihilation cross section:

$$\Omega_\chi \propto \frac{1}{\langle \sigma v \rangle_{ann}} \sim \frac{m_\chi^2}{g_\chi^4}$$

Required annihilation cross section:

$$\langle \sigma v \rangle_{ann} \sim 2 \times 10^{-26} \text{ cm}^3/\text{s}$$



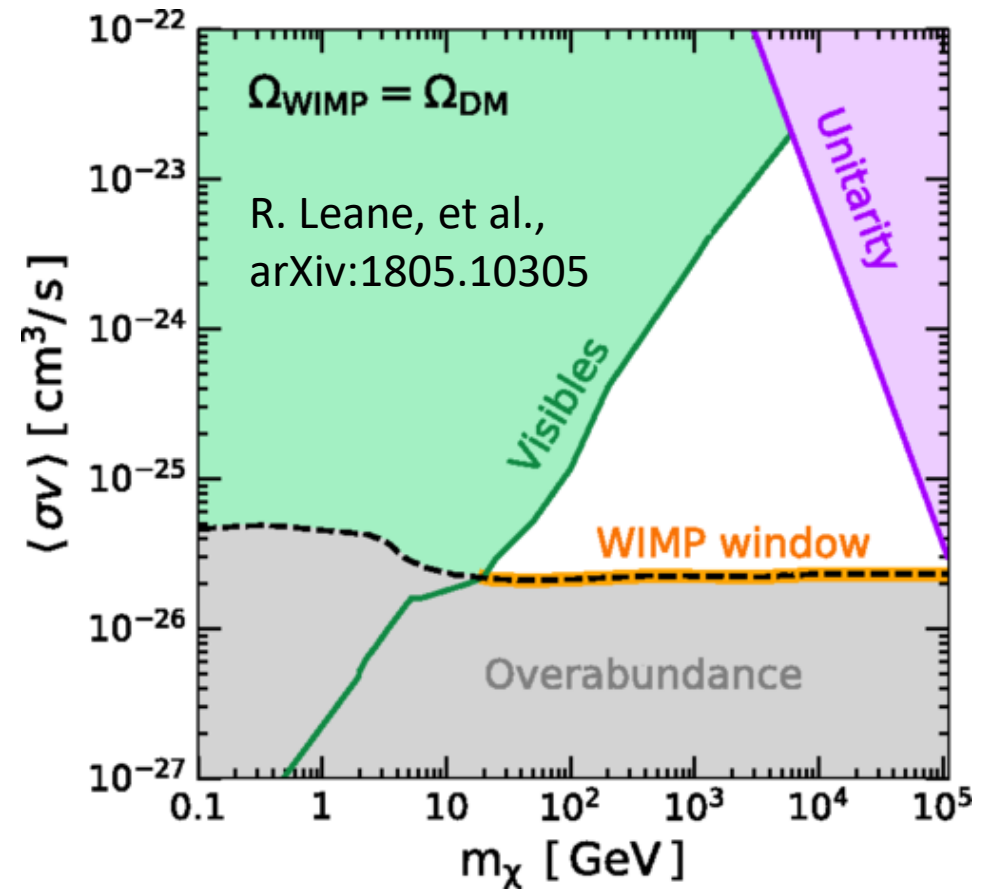
The WIMP window

Mass window for thermally produced WIMPs:

$m_\chi < 100 \text{ TeV}$ from Unitarity limit

$m_\chi > \text{MeV}$ to avoid upsetting BBN

→ We need to test thermal-relic annihilation cross sections across the full mass window



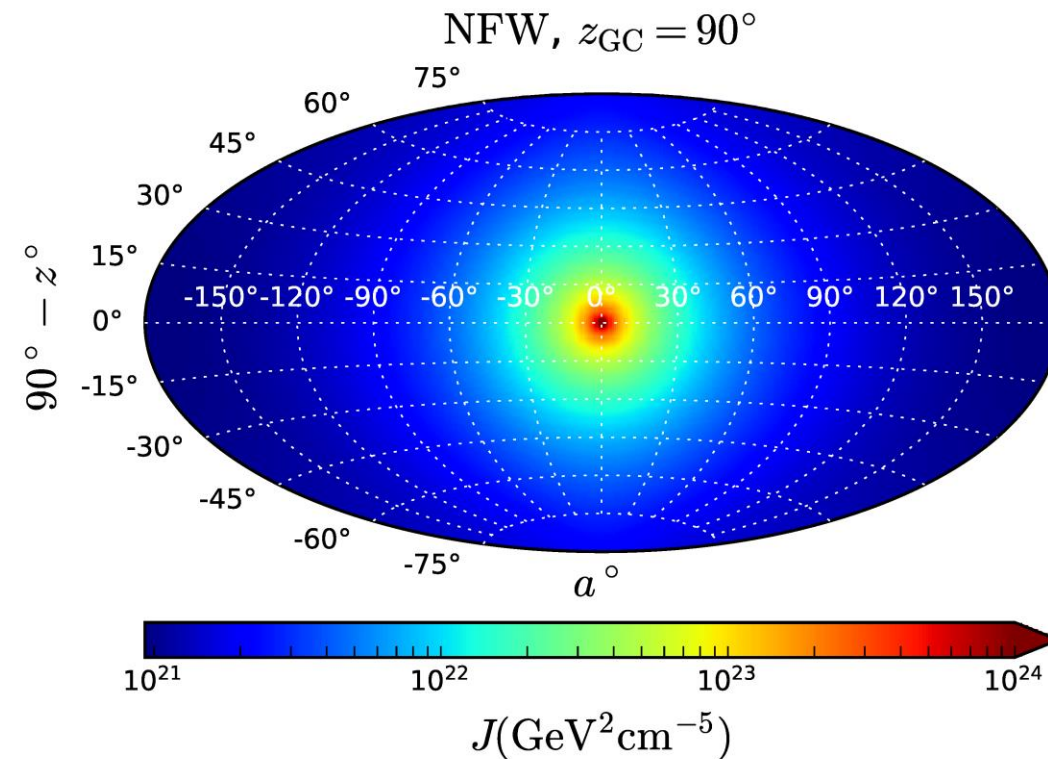
Dark matter annihilation signal

$$\frac{d\Phi_{\Delta\Omega}}{dE} = \langle\sigma v\rangle \frac{J_{\Delta\Omega}}{8\pi m_{DM}^2} \frac{dN}{dE}$$

Annihilation
cross section

Integral of
(density)² along
line of sight

Spectrum per
annihilation

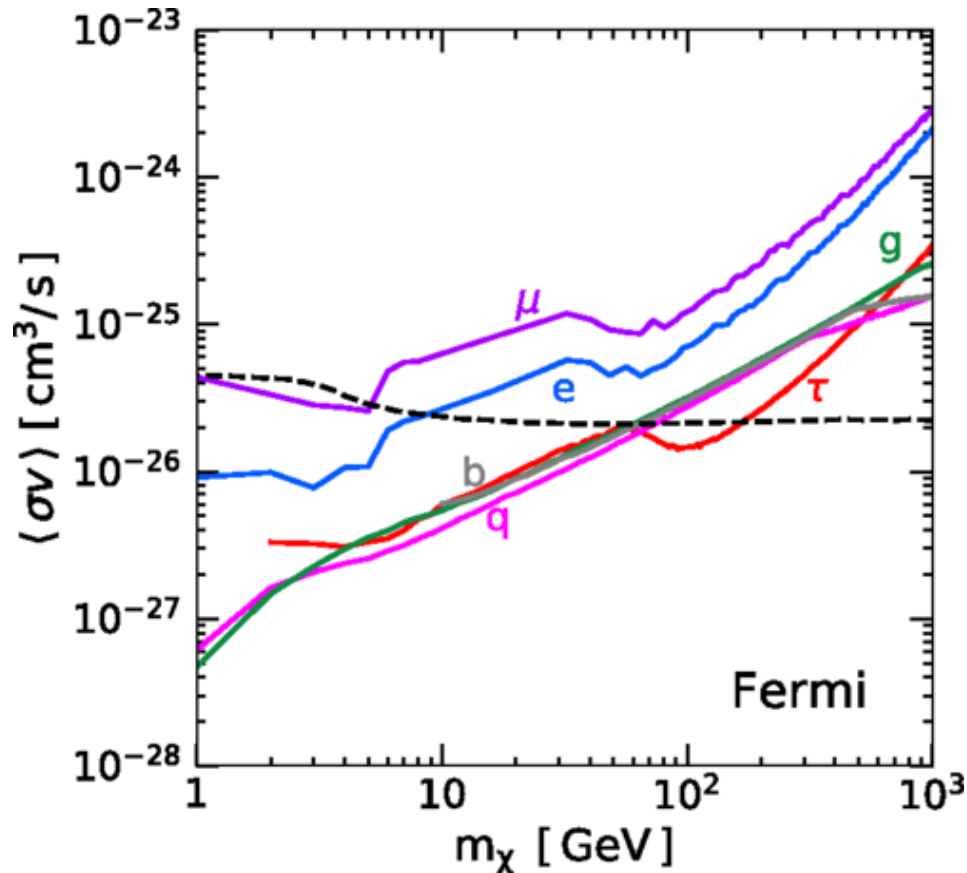


Bell, Dolan, Robles, arXiv: 2005.01950

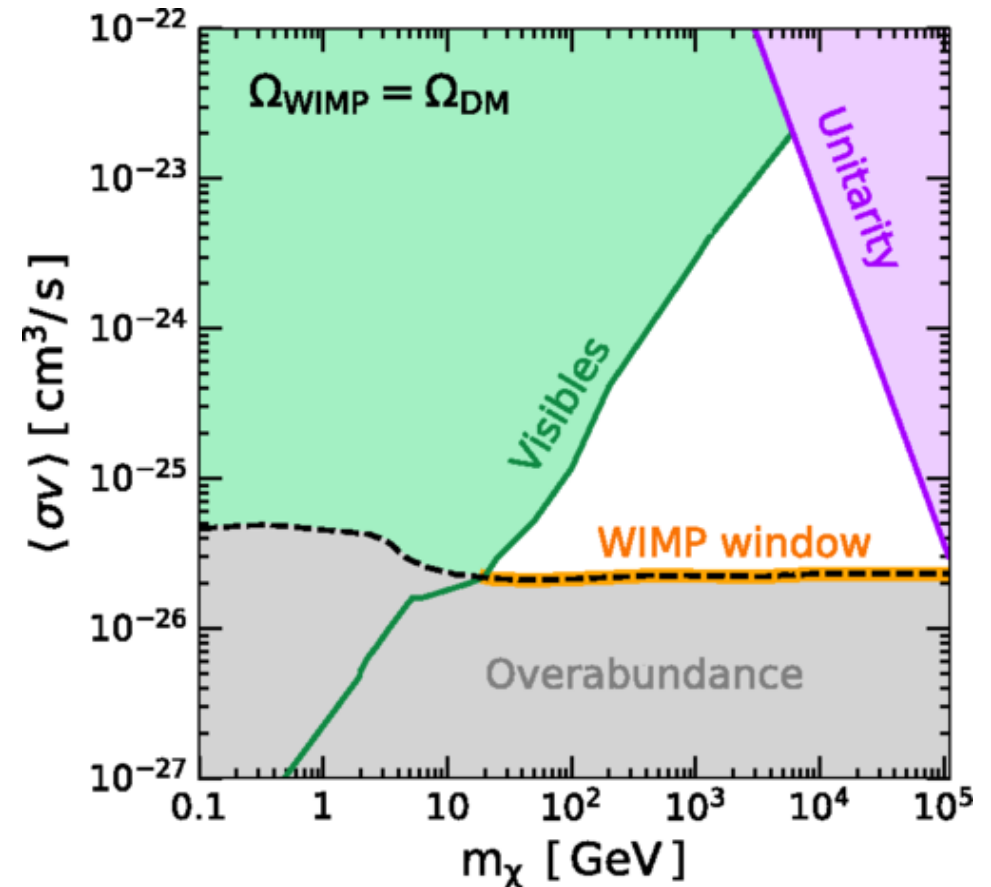
Indirect detection constraints

R. Leane, et al., arXiv:1805.10305

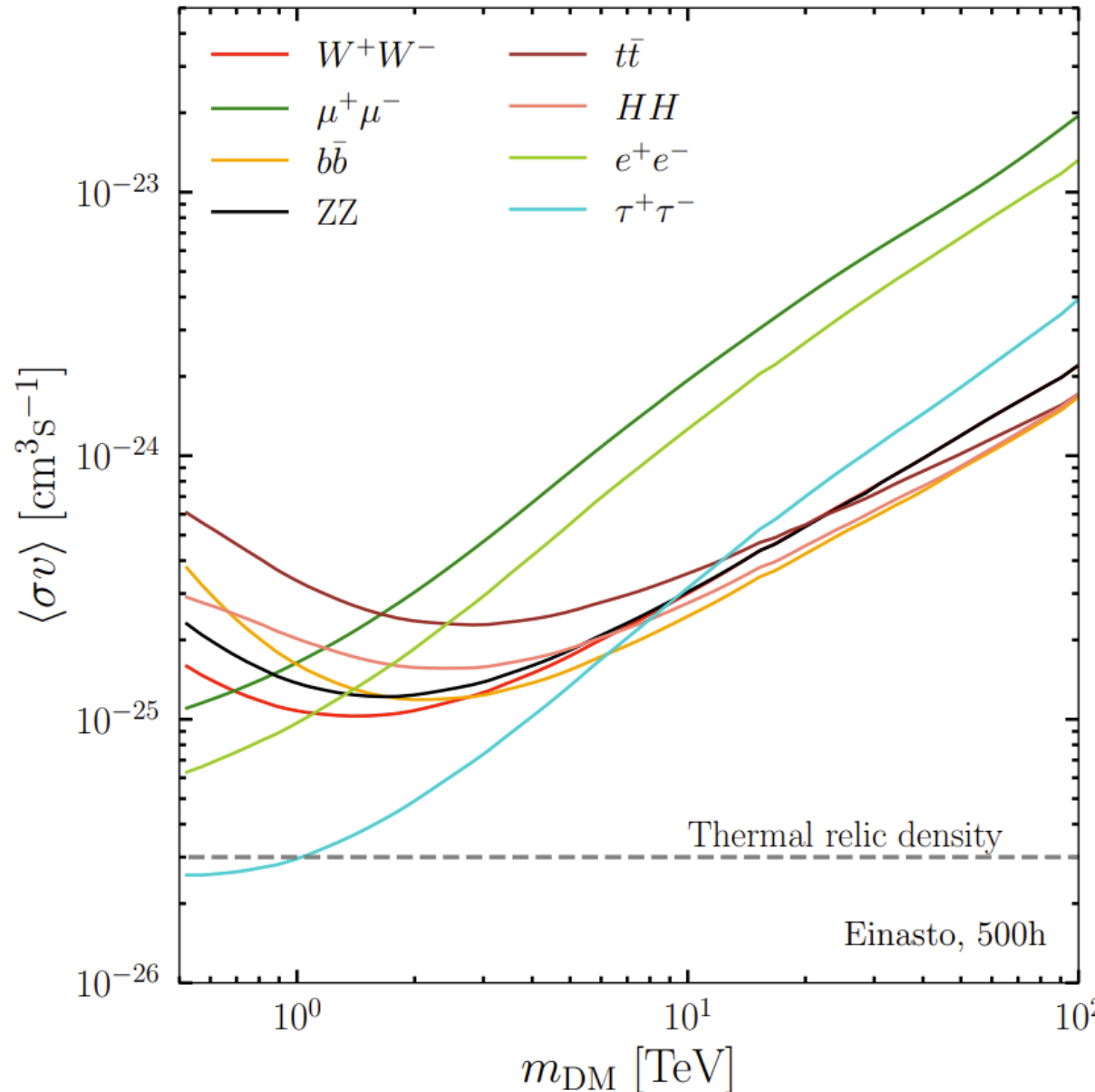
Fermi dSph limits



Annihilation to “visible” SM states



Closing the WIMP window: TeV gamma rays

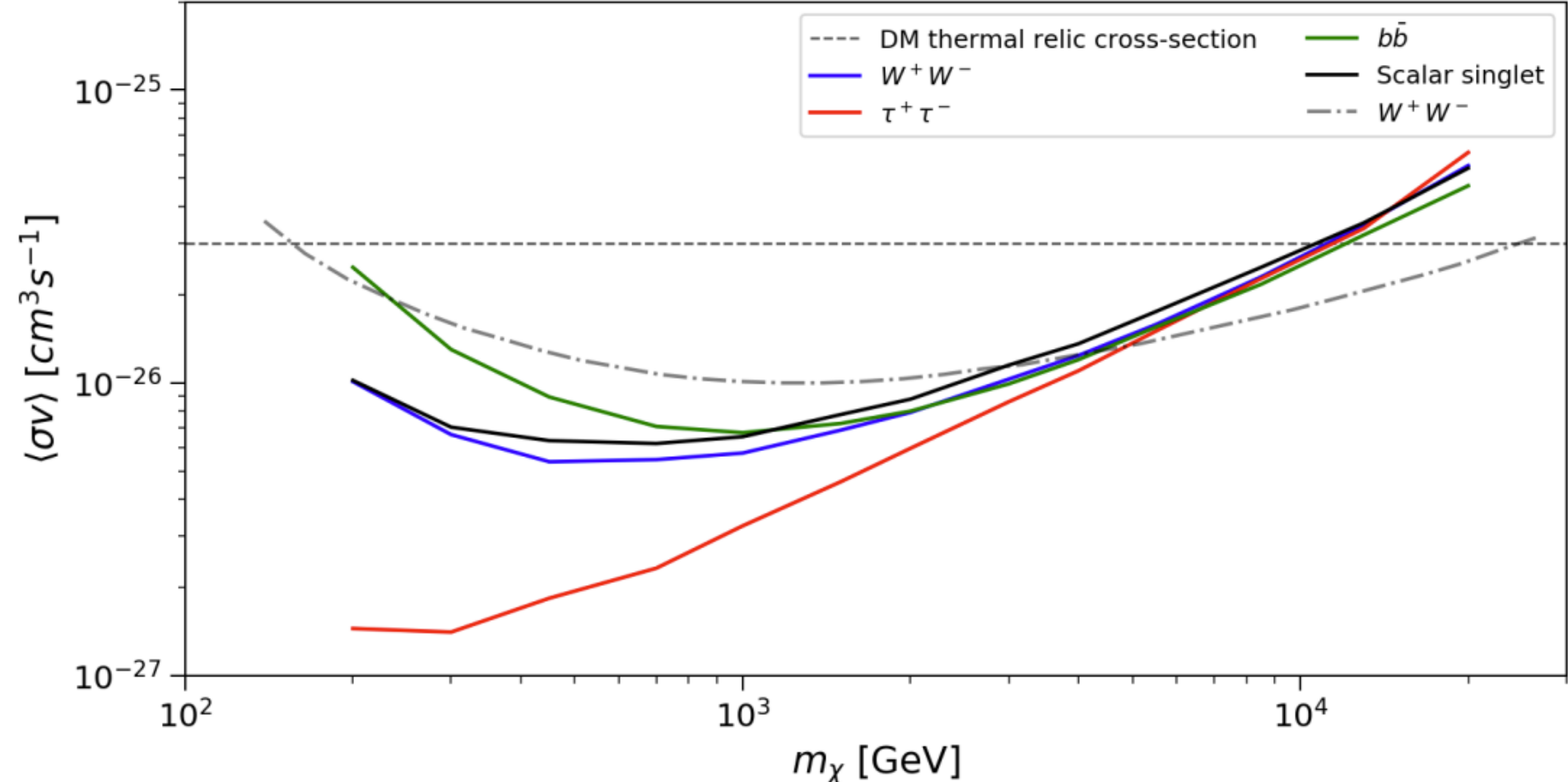


Projected sensitivity
for current generation
Cherenkov telescopes
(HESS-like)

Montanari, Moulin & Rodd,
[arXiv:2210.03140](https://arxiv.org/abs/2210.03140)

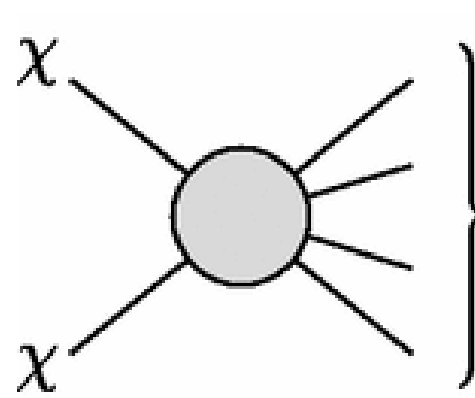
Closing the WIMP window: CTA projections

Mangipudi, Thrane & Balazs, arXiv:2112.10371



Closing the WIMP window: Neutrinos

- Indirect detection limits – typically neglect the possibility that dark matter may annihilate to “invisible” or hard-to-detect final states.



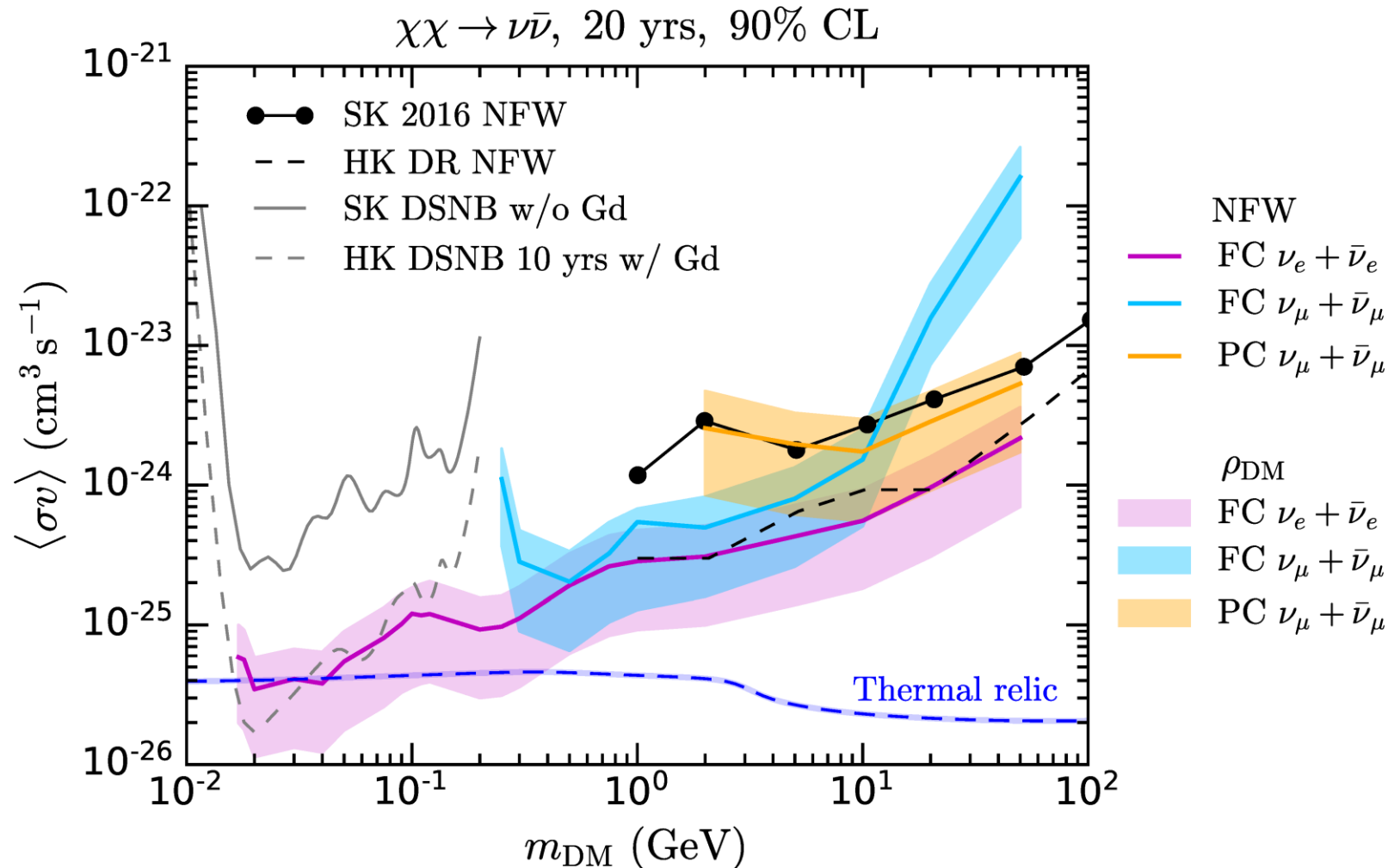
A Feynman diagram showing two incoming particles, labeled χ , interacting at a central grey circular vertex. From this vertex, several lines radiate outwards, representing various possible final states. A large curly bracket groups these lines and is labeled "All Standard Model final states".

$$\left. \begin{array}{c} \chi \\ \chi \end{array} \right\} \text{All Standard Model final states} = \begin{array}{l} \text{“Visible” states:} \\ \gamma\gamma, \bar{q}q, e^+e^-, \dots \\ + \\ \text{“Invisible” states:} \\ \bar{\nu}\nu \end{array}$$

Beacom, Bell,
Mack, PRL 2007

- **We must probe annihilation to neutrinos to fully test the WIMP hypothesis.**

Annihilation cross section limits: $\chi\chi \rightarrow \nu\bar{\nu}$

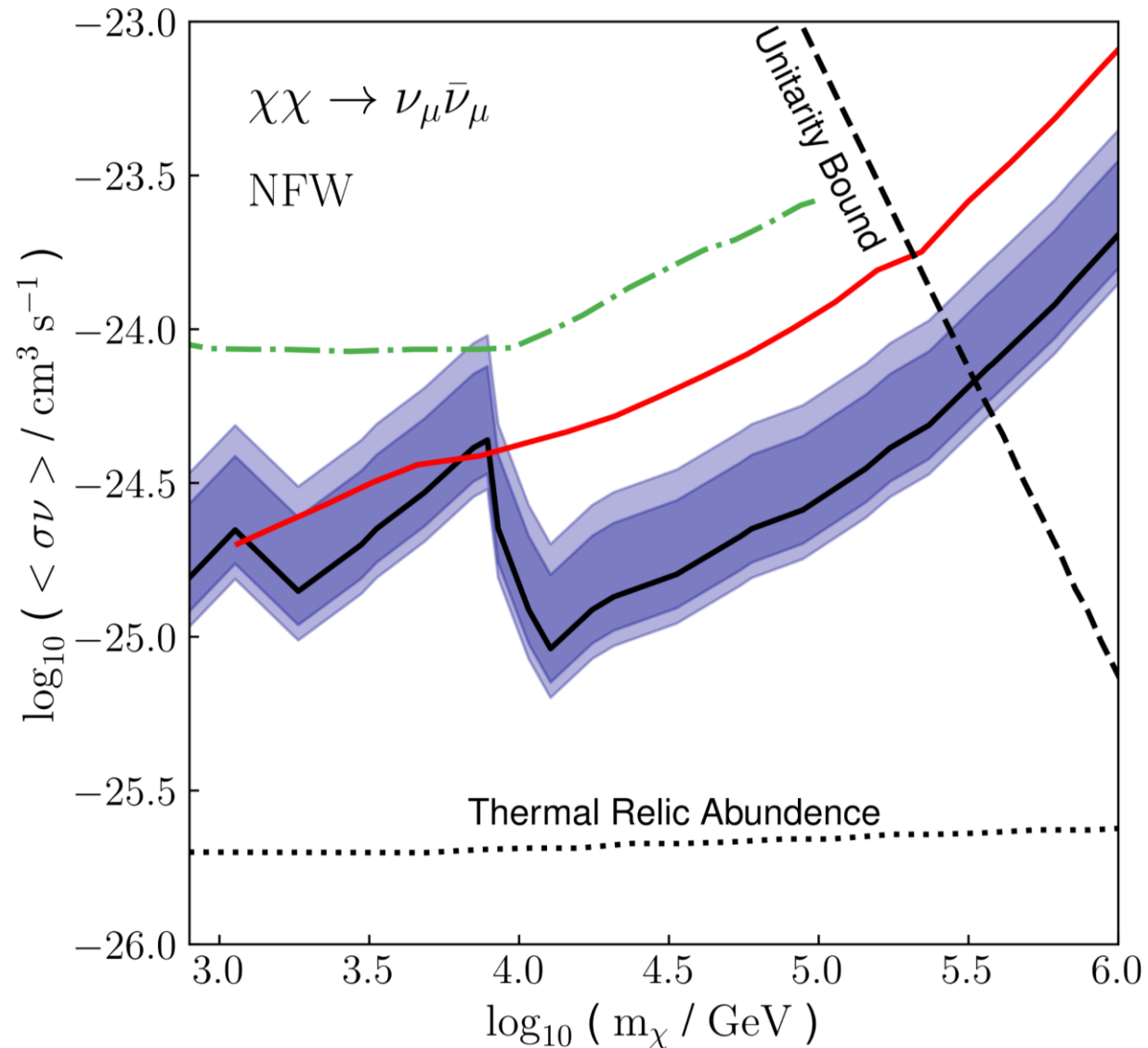


Thermal relic sensitivity
for $m_\chi \sim 30$ MeV

NFW – central lines
Isothermal – upper
Moore - lower

Bell, Dolan, Robles, arXiv: 2005.01950

Annihilation to $\nu\bar{\nu}$ – large mass

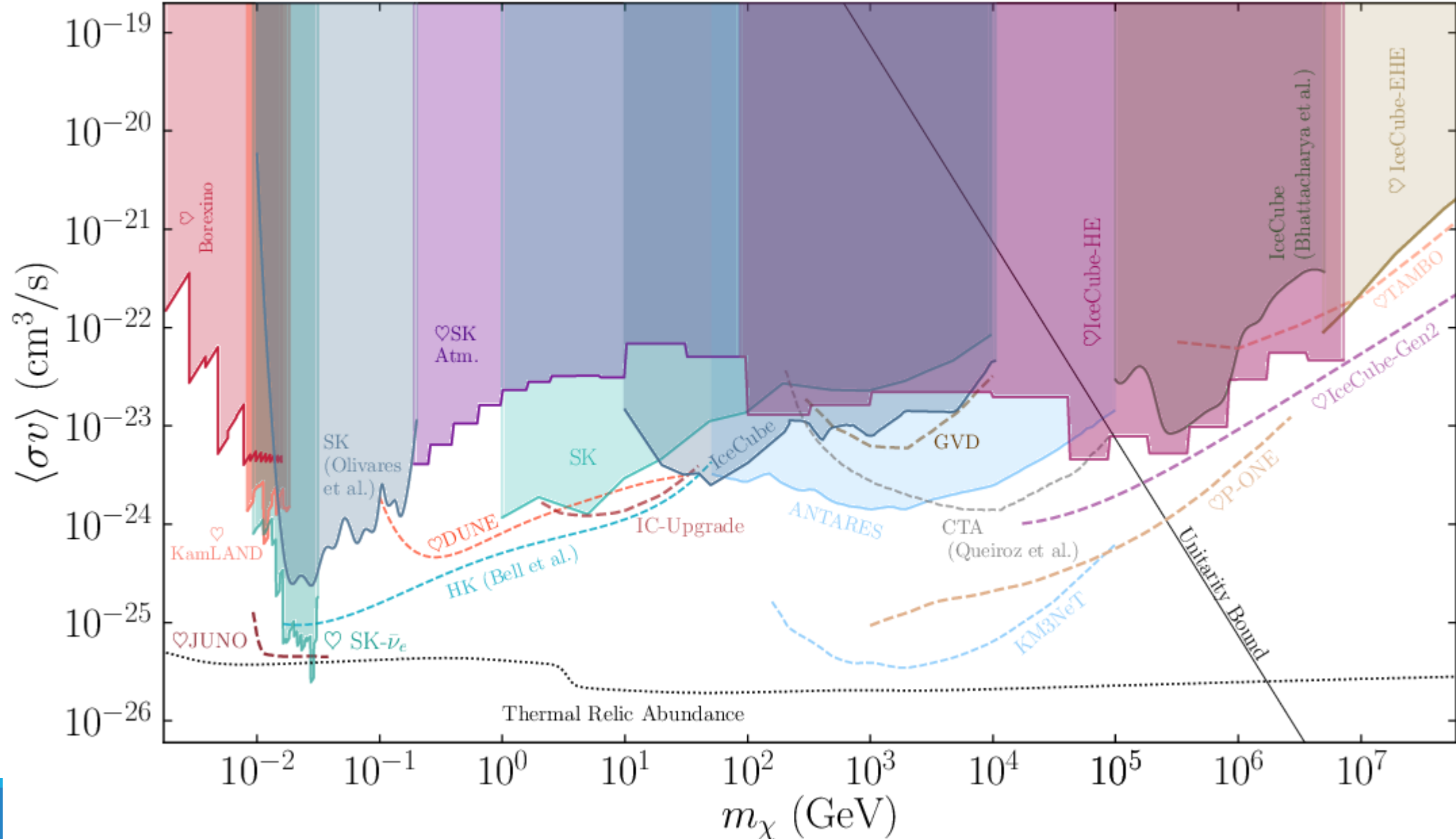


Desai, Li & Meighen-Berger arXiv:2302.10542

Projection for 10 years of P-OONE

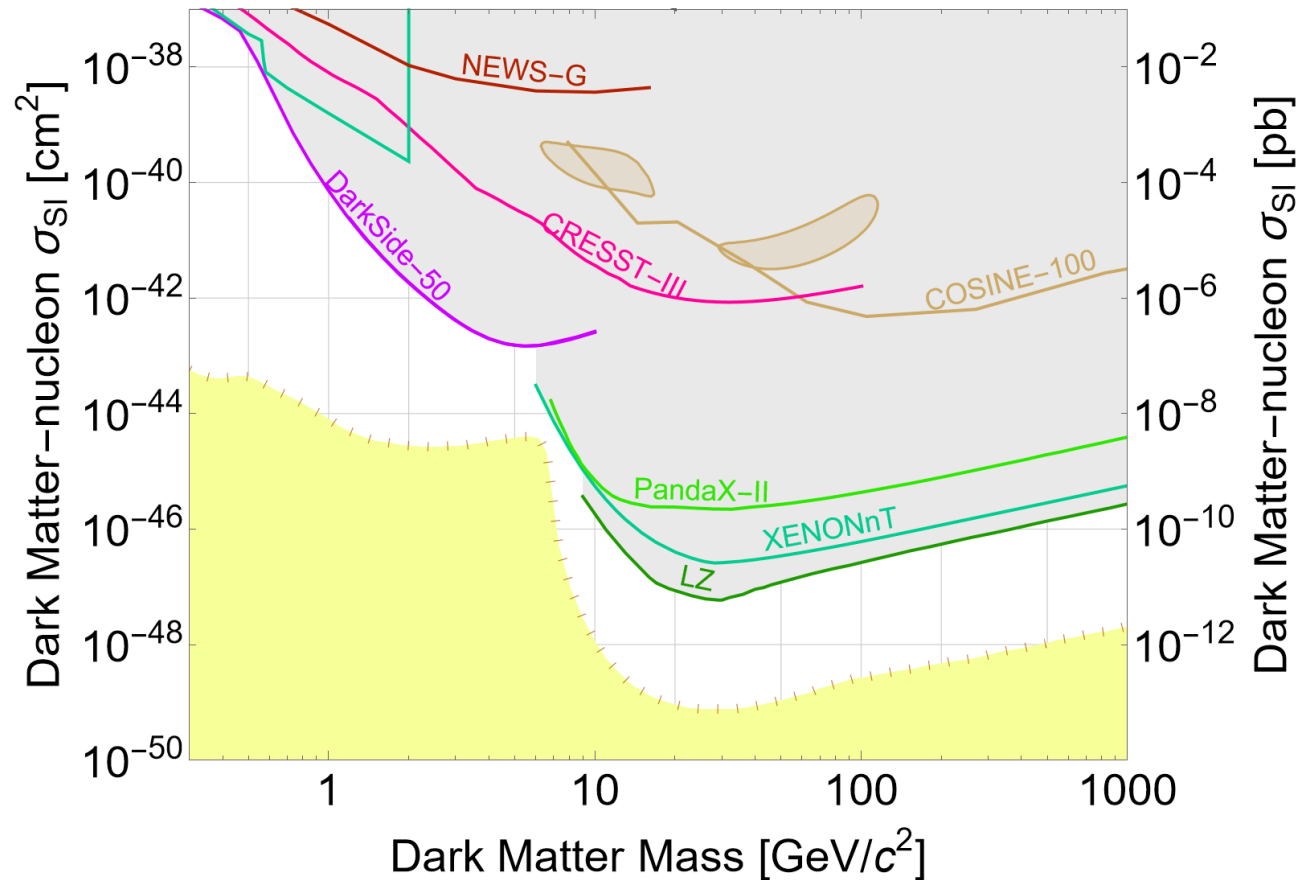
Annihilation to $\nu\bar{\nu}$

Arguelles et al, arXiv: 1912.09486



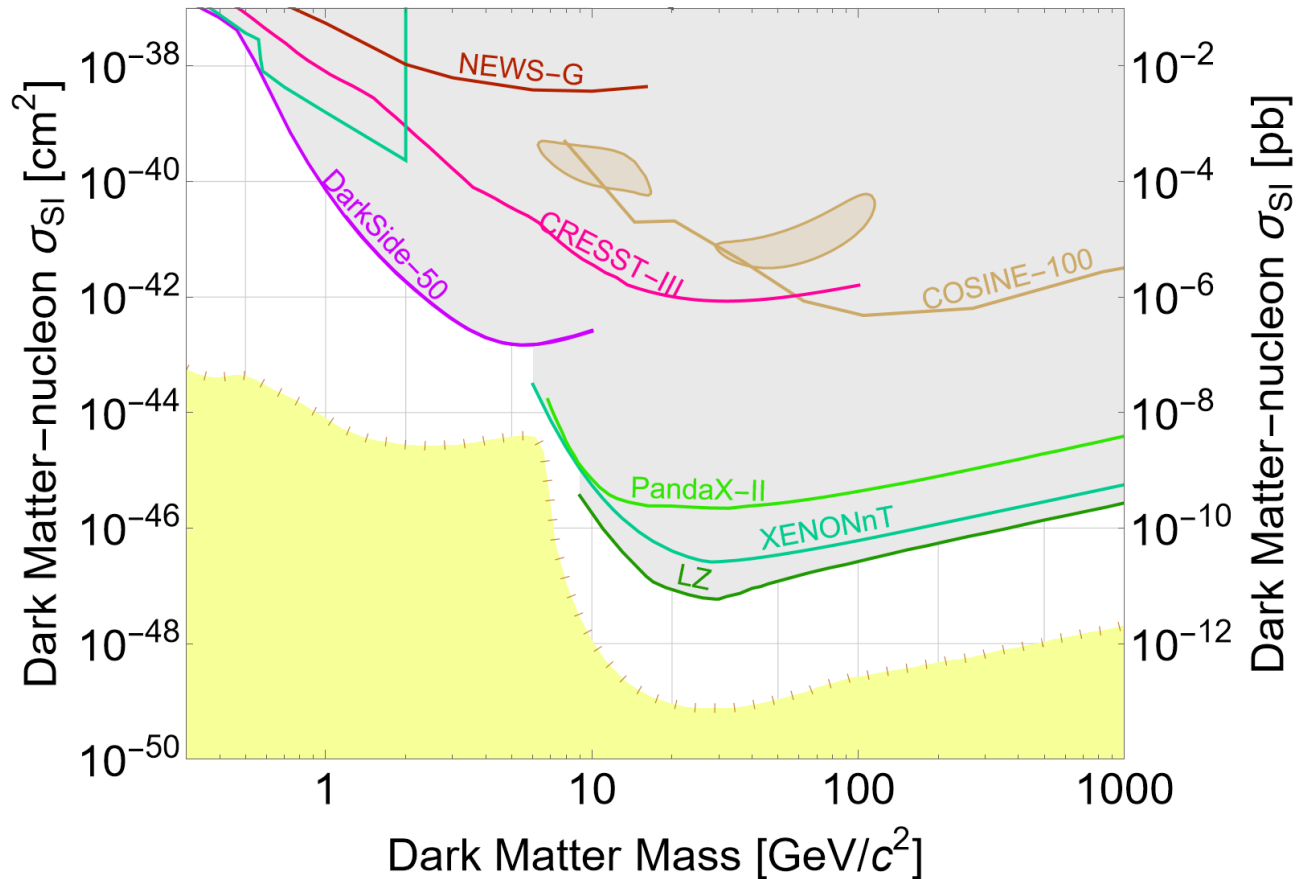
Dark matter scattering

Direct Detection bounds



Dark matter scattering

Direct Detection bounds



How can we improve or complement this with astrophysics?

Boosted Dark Matter

Halo dark matter

→ highly nonrelativistic $v \sim 10^{-3}c$

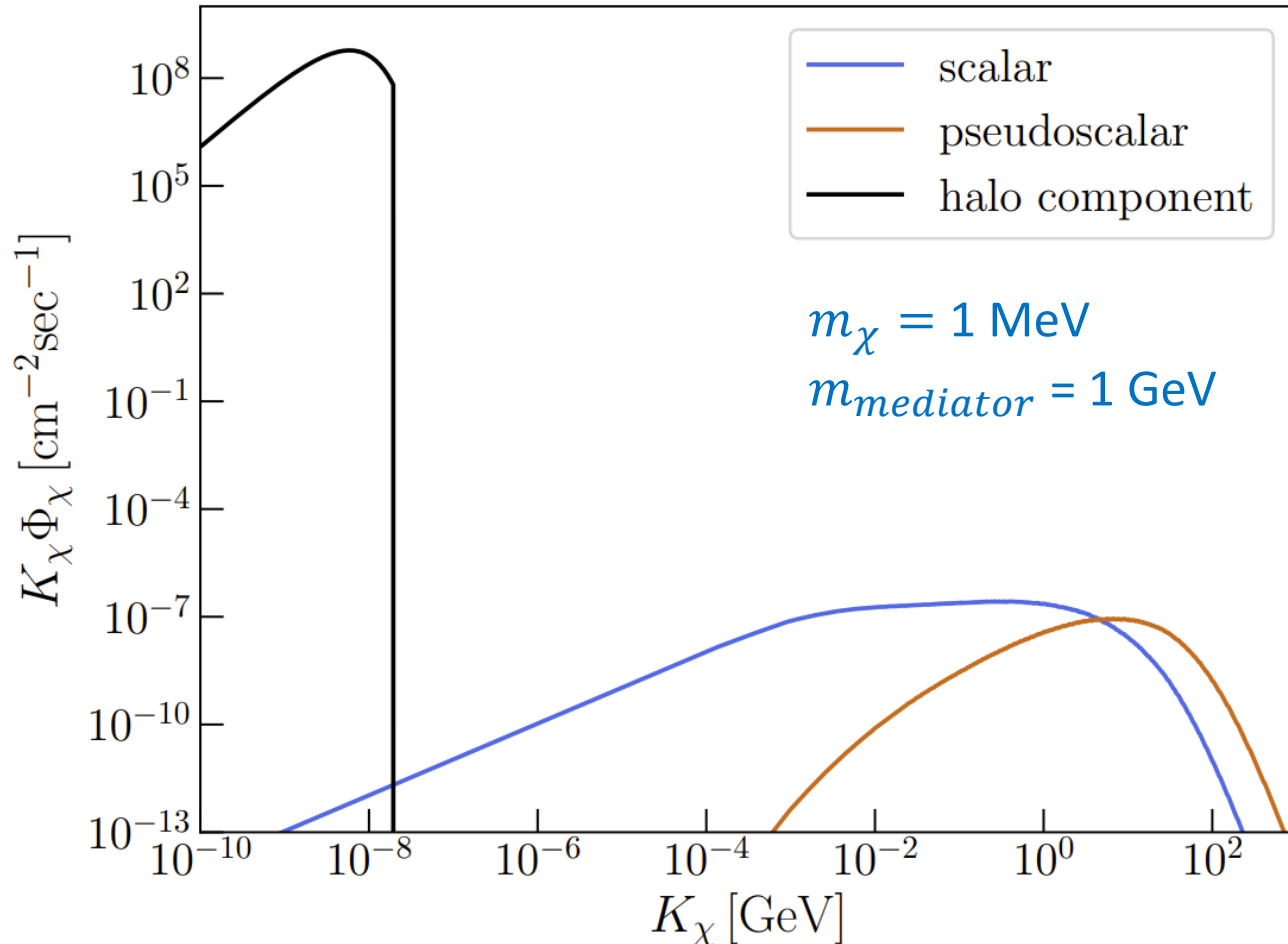
→ low energy recoils in direct detection experiments: $E_{R,max} = \frac{2\mu_T^2}{m_T} v_{max}^2$

Could there be a population of higher-energy dark matter?

- Boosted DM produced from decay/annihilation of heavier dark states
- **Cosmic-ray upscattered dark matter** (“inverse direct detection”)
- DM produced in cosmic ray interactions in the atmosphere (“CR beam dump”)
- Solar reflected dark matter
- Supernova dark matter (light dark matter produced in galactic supernova)

Cosmic ray up-scattered dark matter (CRDM)

Y. Ema et al, arXiv:2011.10939



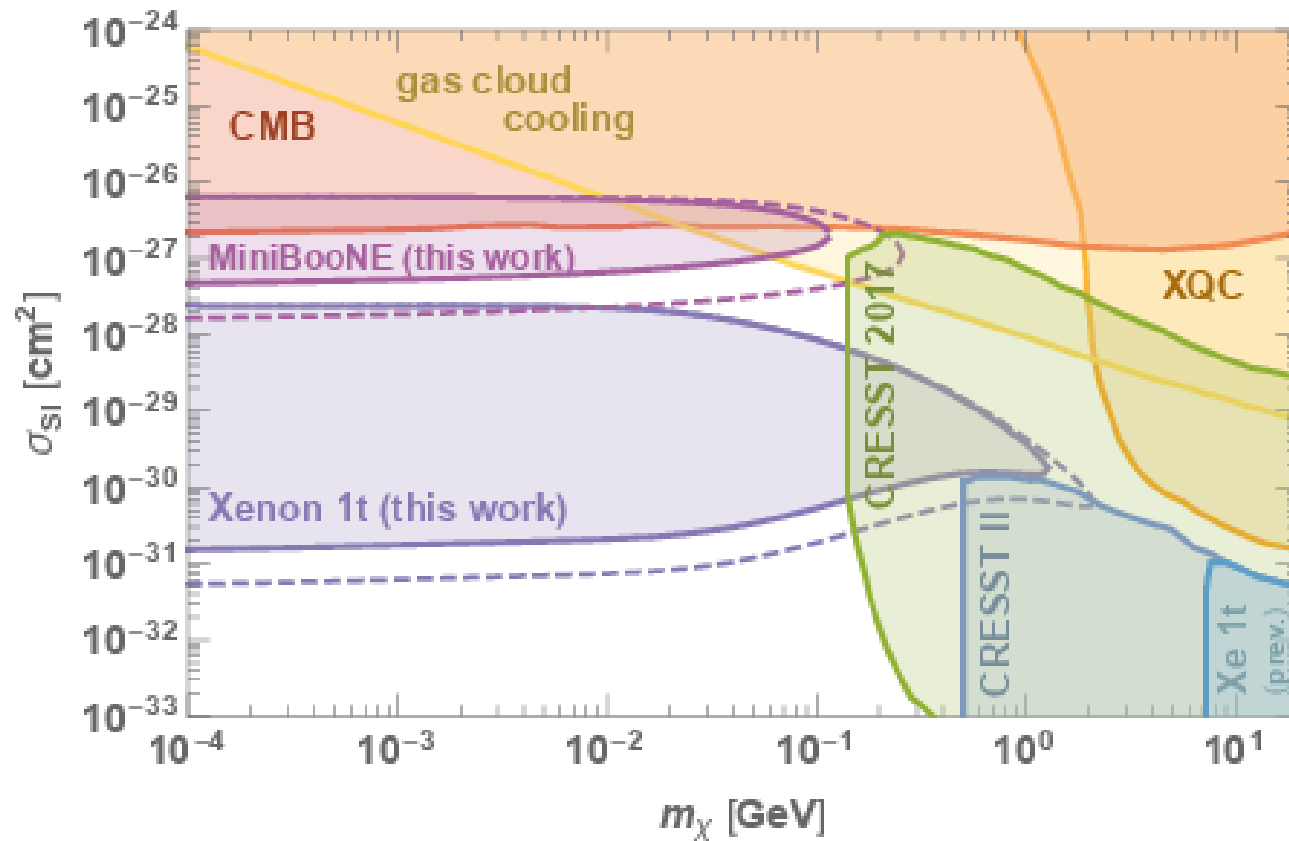
Assume the DM-nucleon scattering cross section is non-zero

→ cosmic rays will **unavoidably** scatter with DM, producing a (small) high energy DM flux.

→ Light boosted DM is visible in direct detection experiments

Cosmic ray up-scattered dark matter – sub-GeV masses

Bringmann & Pospelov, PRL 2019



Allows light dark to be constrained using existing experiments.

Note:

- these are BIG cross sections
- DM absorption in the earth imposes upper limit on the cross sections that can be probed

Cosmic ray up-scattered dark mater (CRDM)

Advantages of relativistic energies:

- Detectable signals for **light DM** in direct detection experiments
- Energetic enough to be seen in **neutrino experiments**
→ which have higher energy thresholds, but **significantly larger target mass**
- Removes velocity or momentum suppressions
→ e.g. Standard direct detection exp cannot see pseudoscalar-type interactions, as the cross section is suppressed as p^4

Disadvantages:

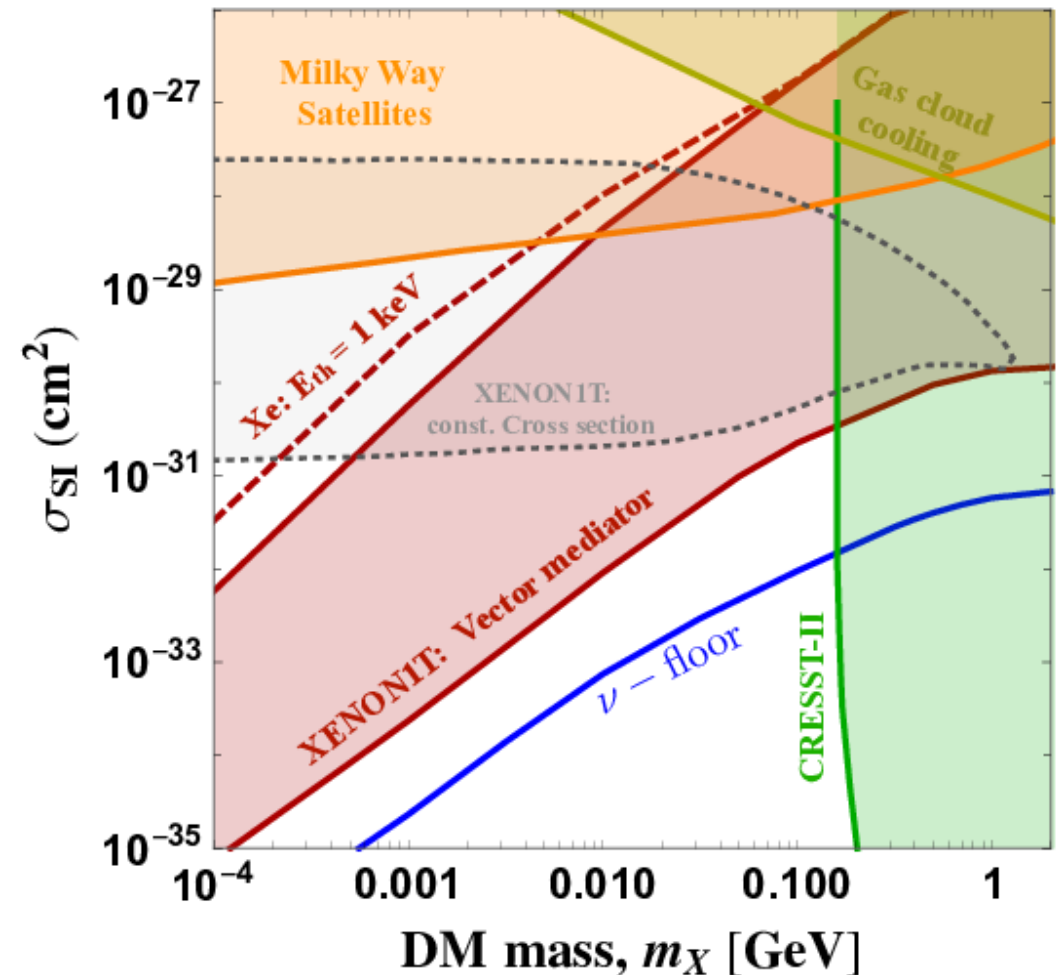
- Observable signals scale with **two** powers of the scattering cross section

Cosmic ray up-scattered dark matter

Dent, Dutta, Newstead, Shoemaker,
arXiv:1907.03782

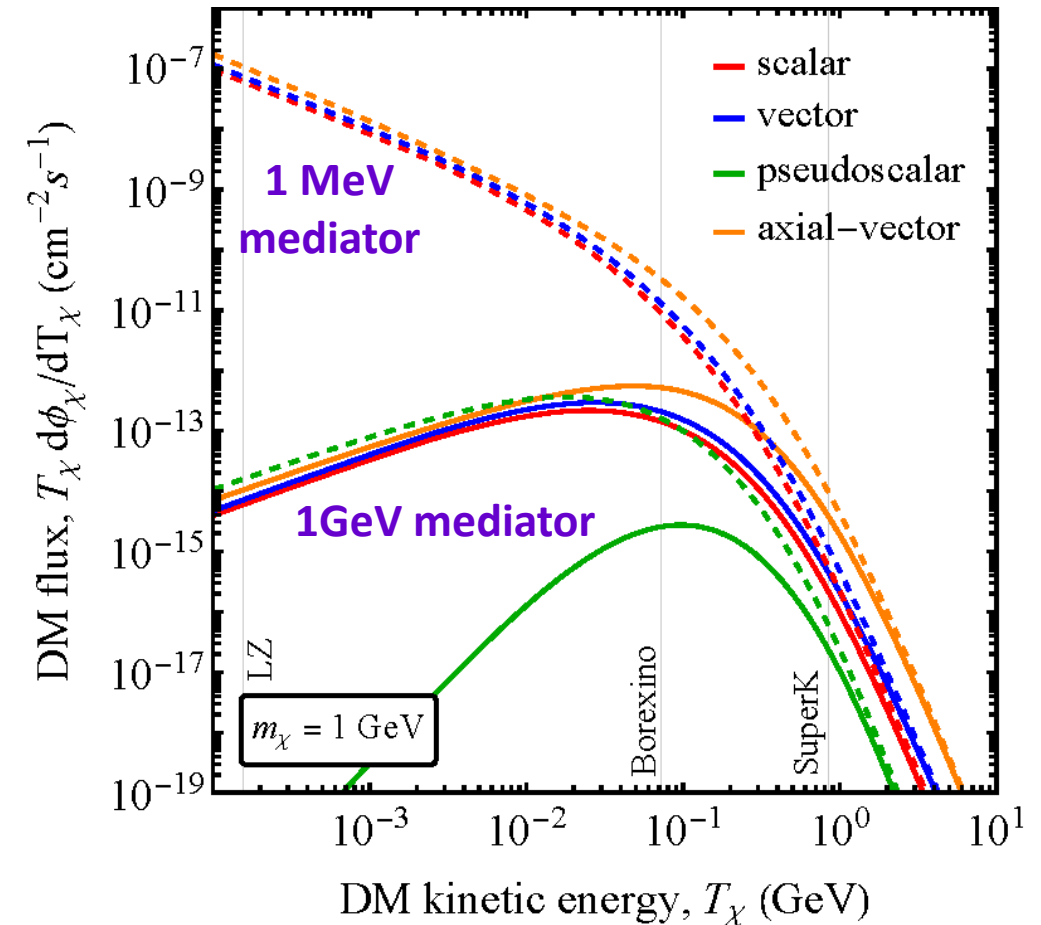
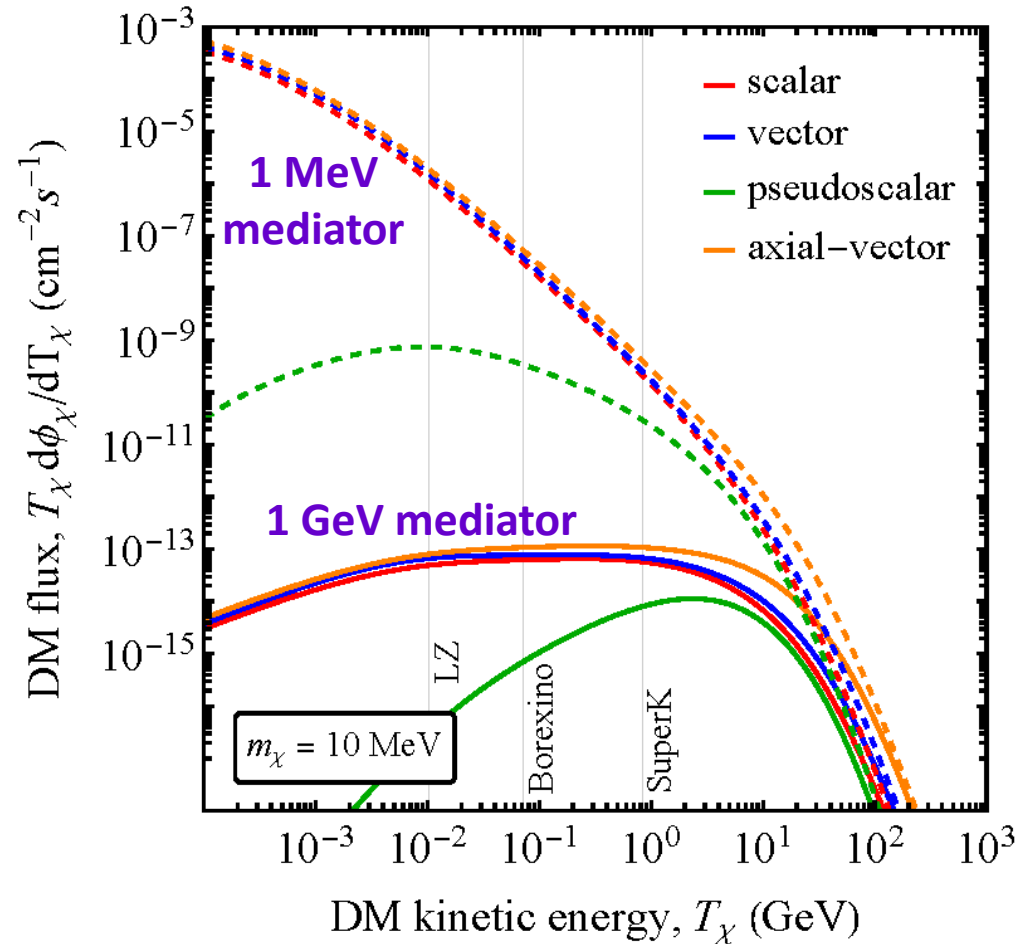
Very big cross-sections

- Light mediators?
 - Energy dep. of cross section matters
 - Other constraints
 - Composite DM?
- These limits are model dependent.



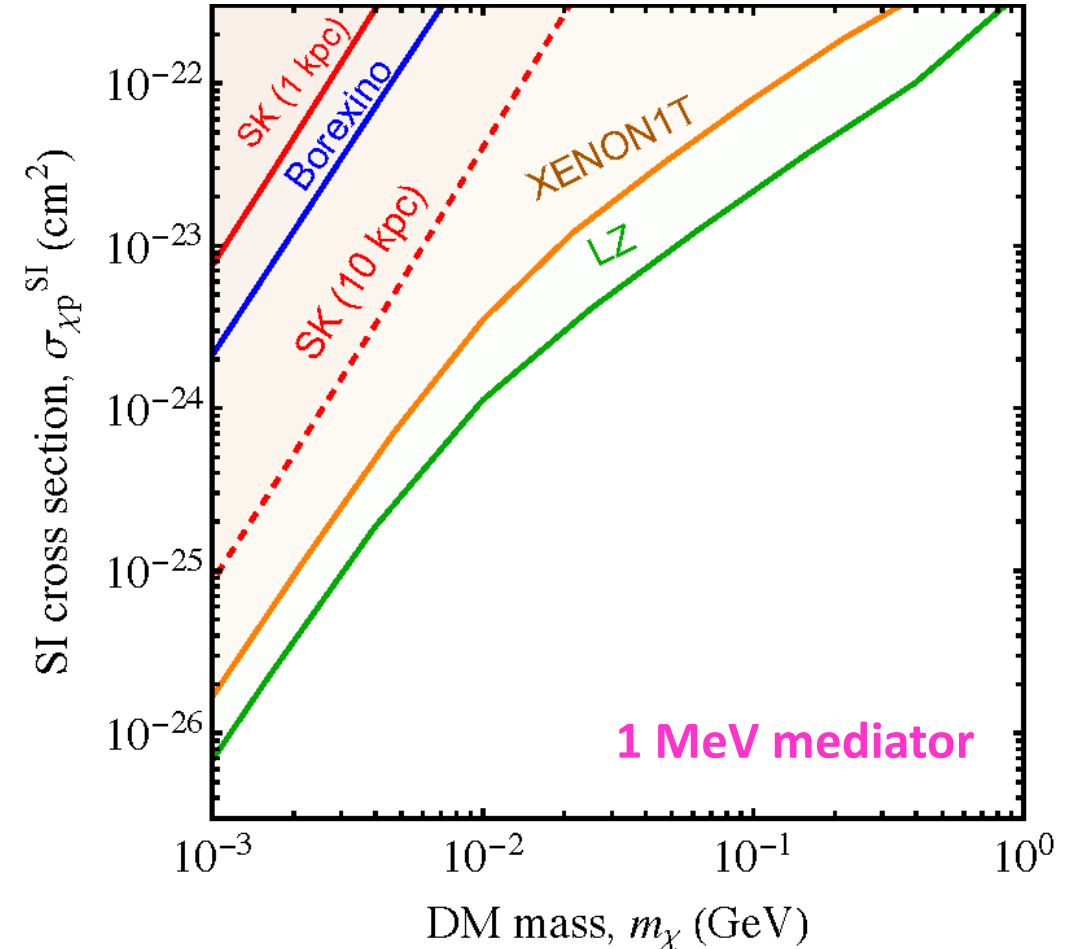
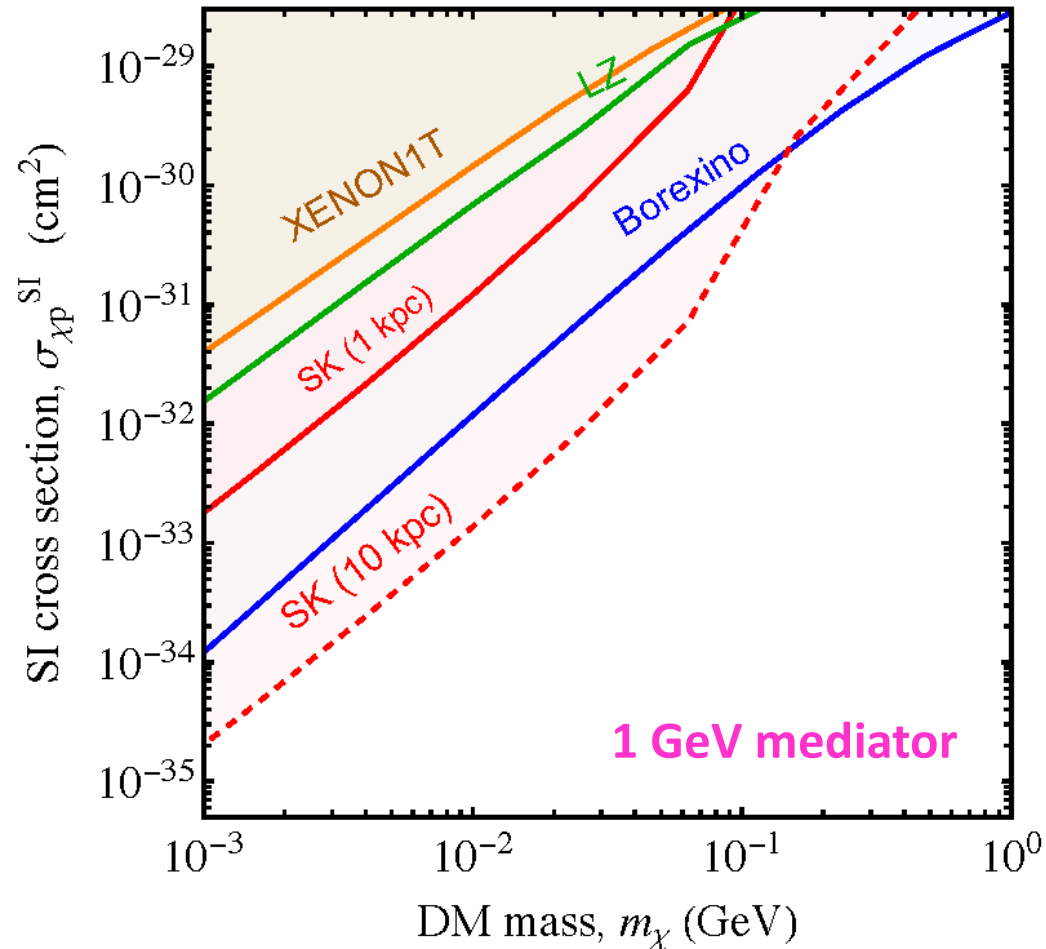
CR-upscattered DM: kinetic energy spectrum

Bell, Newstead and Shaukat Ali, in preparation



Boosted DM – neutrino vs direct detection exps.

Bell, Newstead and Shaukat Ali, in preparation

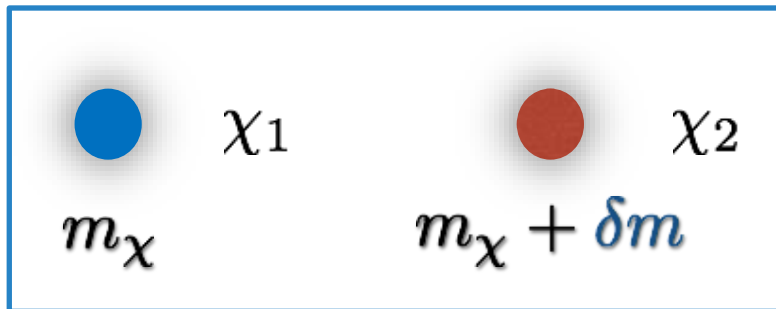


Boosted DM – Inelastic models

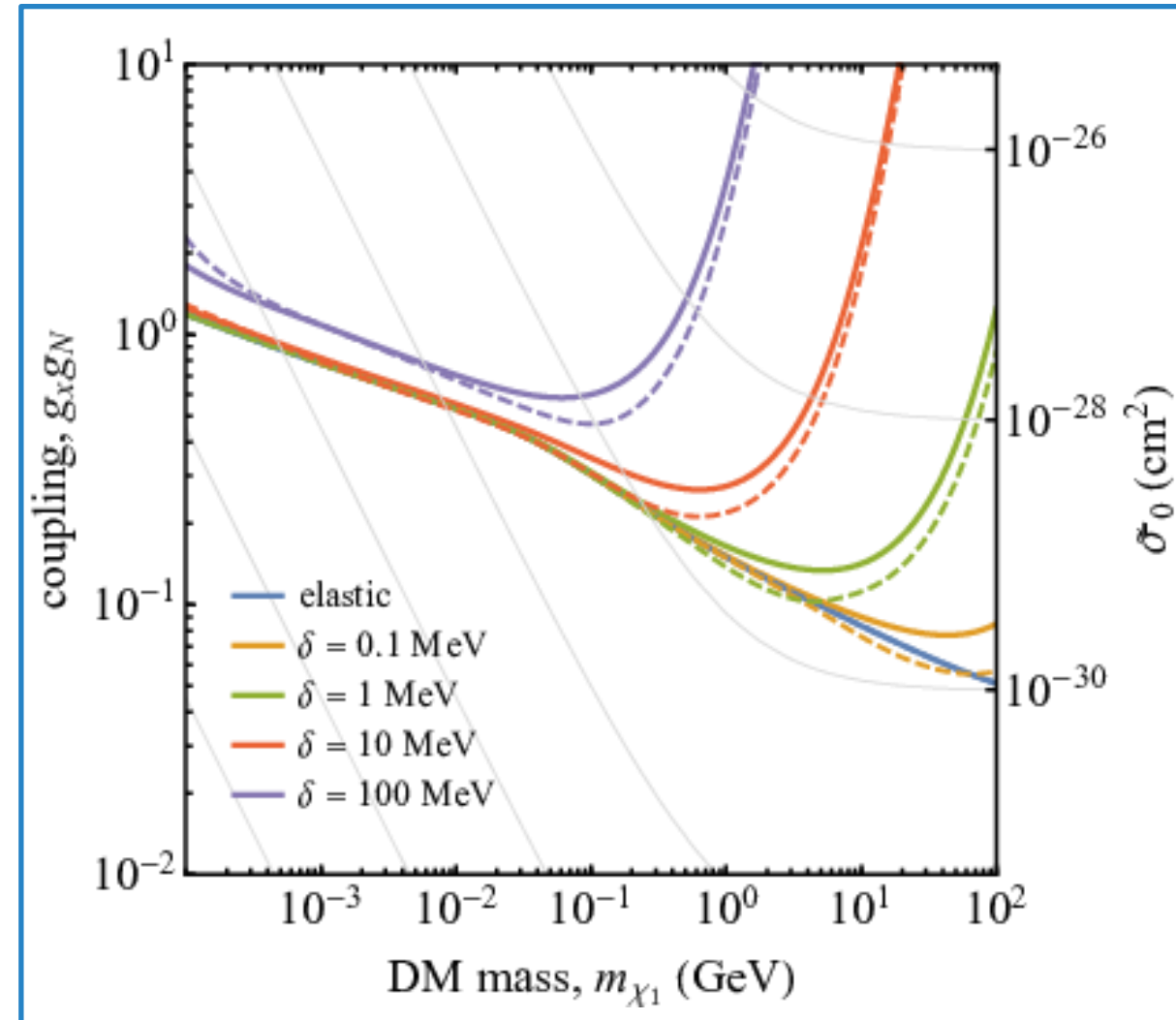
Bell, Dent, Dutta, Ghosh, Kumar, Newstead, Shoemaker arXiv:2108.00583

$$\chi_1 n \rightarrow \chi_2 n$$

Boosting to relativistic energies
→ enables inelastic scattering
with large mass gap



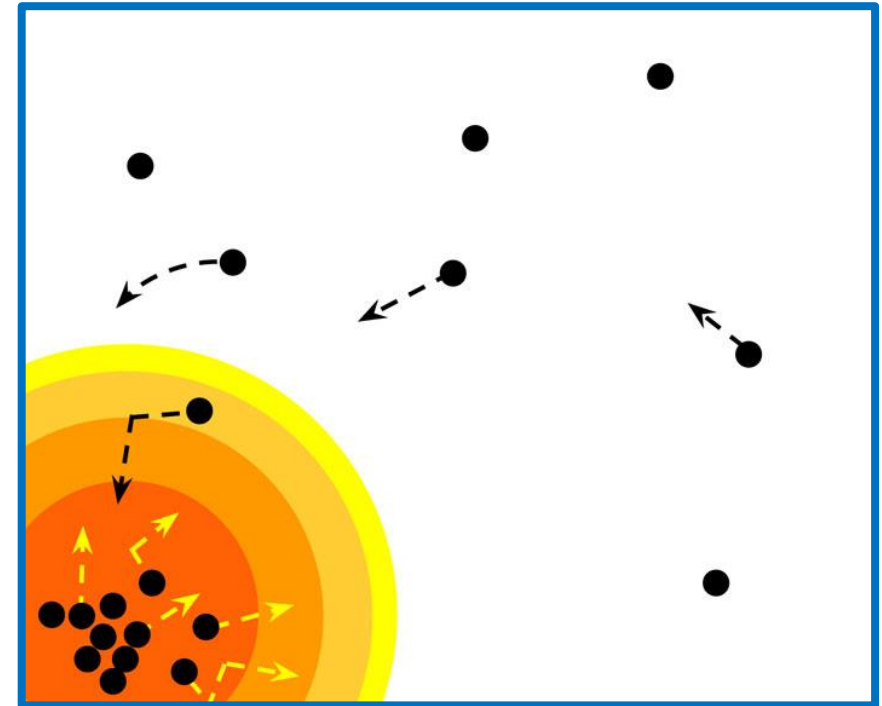
- Direct detection of non-rel DM, restricted to keV mass gaps $\delta m < \mathcal{O}(100)$ keV
- Boosted CR-DM: $\delta m \sim 100$ MeV



Dark Matter Capture in Stars

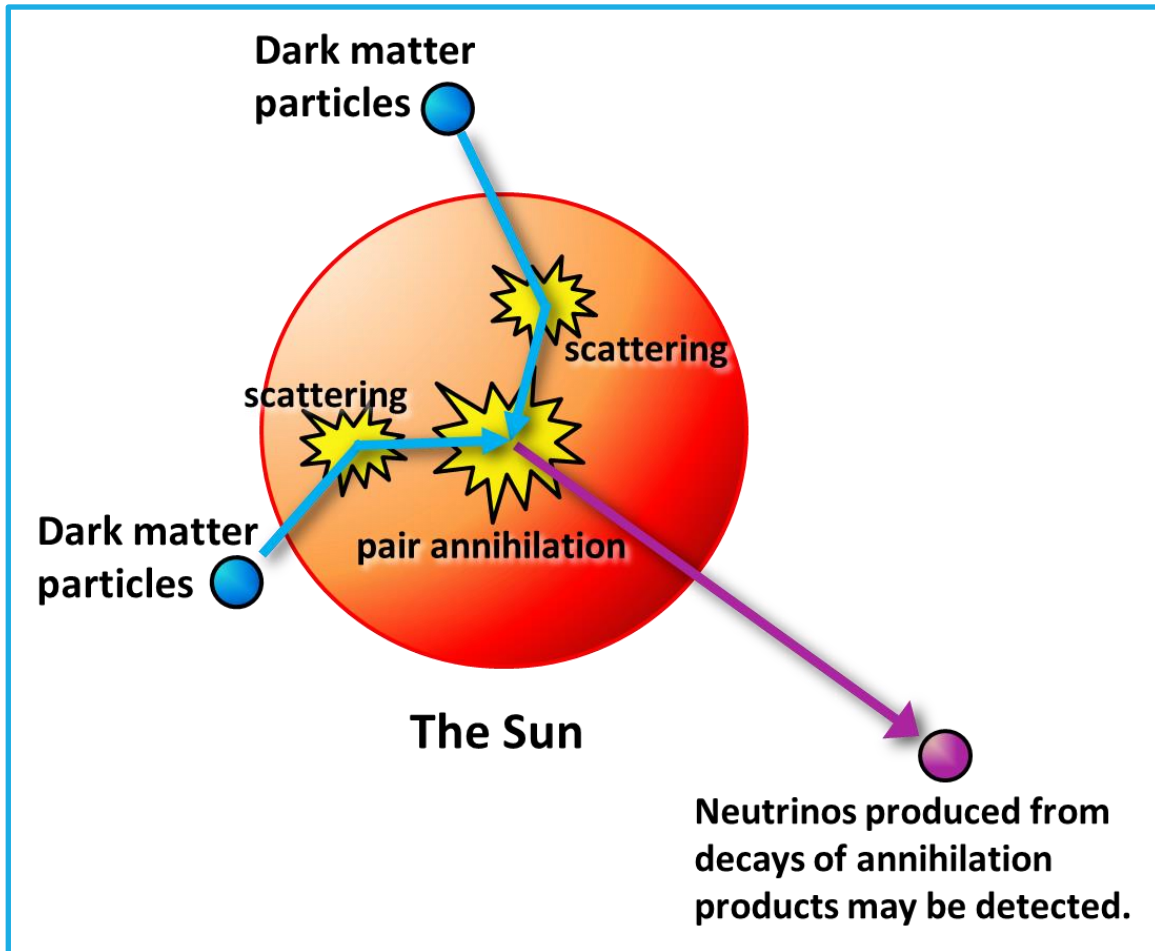
→ an alternative approach to Dark Matter Direct Detection experiments

- The Sun
- Neutron Stars
- White Dwarfs



Dark Matter Capture in Stars

→ *an alternative approach to Dark Matter Direct Detection experiments*



- Dark matter scatters, loses energy, becomes gravitationally bound to star
- Accumulates and annihilates in centre of the star → neutrinos escape

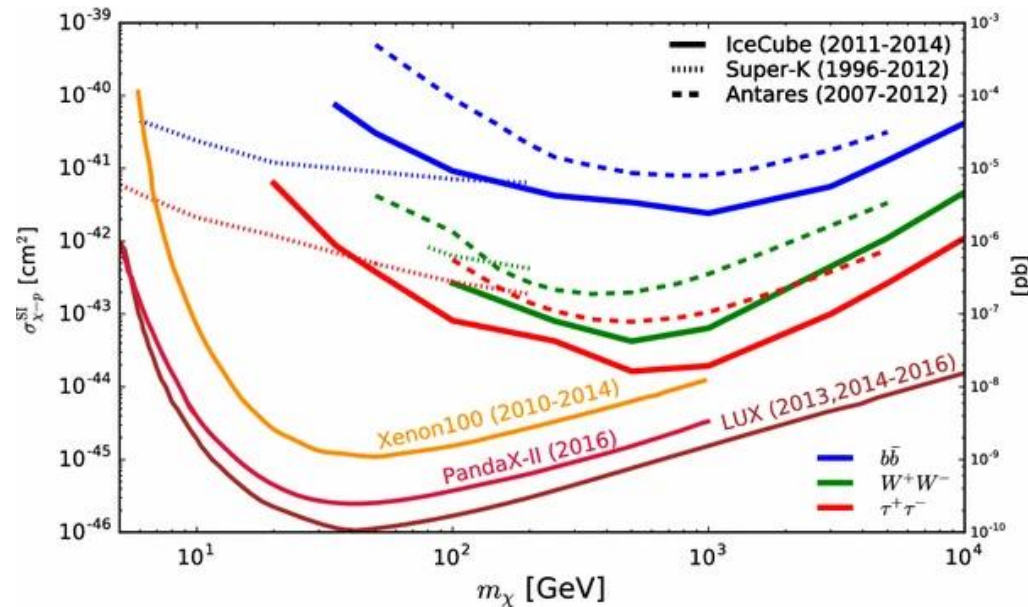
In equilibrium:

Annihilation rate = Capture rate

- controlled by DM-nucleon scattering cross section
- **probes the same quantity as dark matter direct detection experiments**

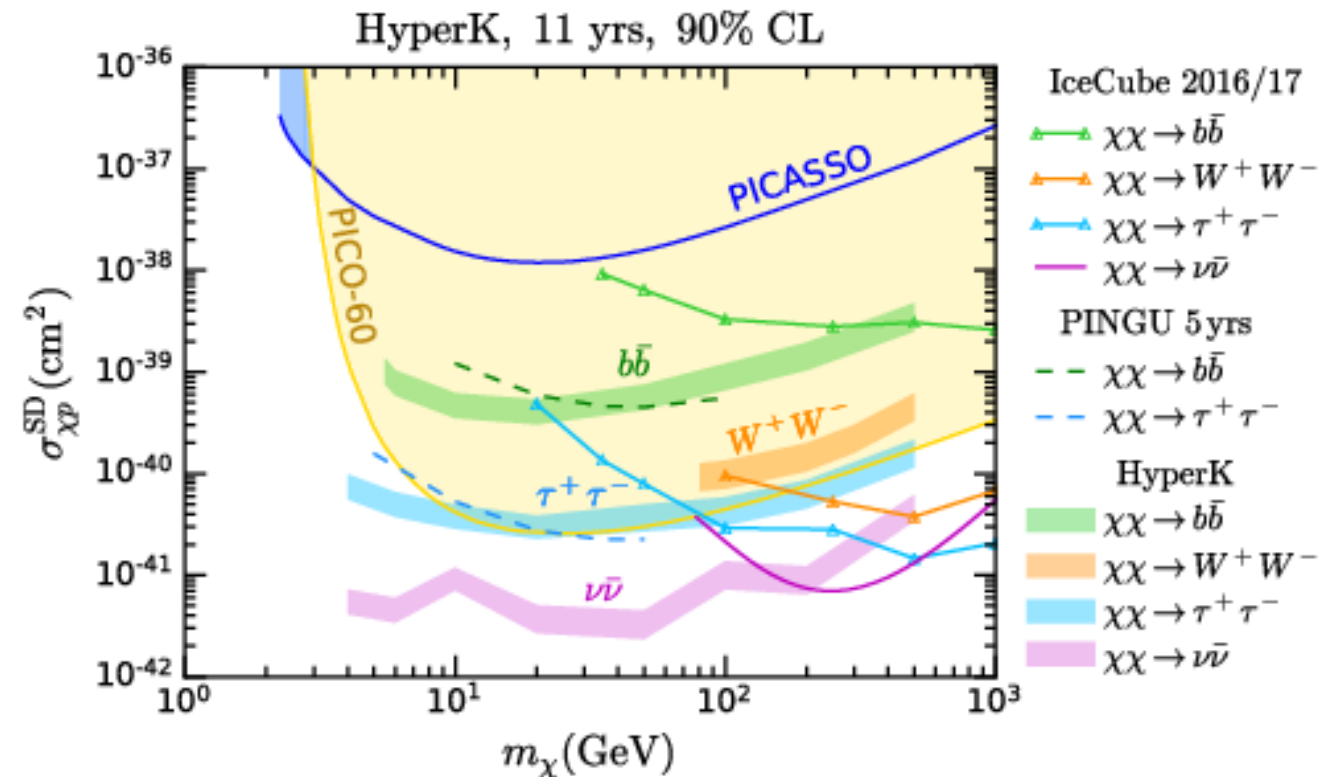
Dark matter annihilation in the Sun – Neutrinos

Spin-Independent (SI)



IceCube Collaboration, E. Phys. J. C 77 (2017)

Spin-Dependent (SD)

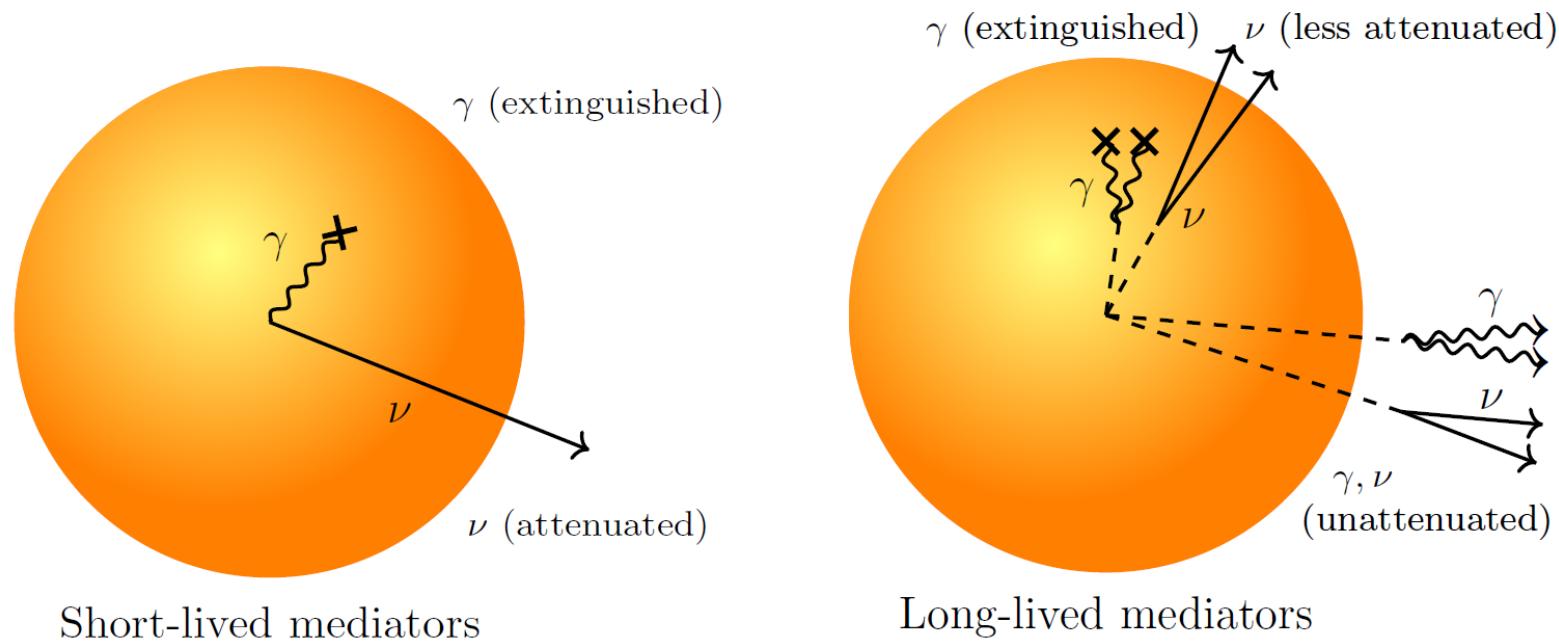


Bell, Dolan & Robles, arXiv:2107.04216

Gamma Rays from the Sun → long lived dark-sector particles

If captured DM annihilates to a light, long-lived mediator (e.g. a dark photon):

- Annihilation products can escape the Sun
- Decay between Sun and Earth → solar gamma rays or cosmic rays ([Batell arXiv:0910.1567](#))
- Decay beyond solar core → less attenuation of neutrino signal ([Bell & Petraki, JCAP 2011](#))

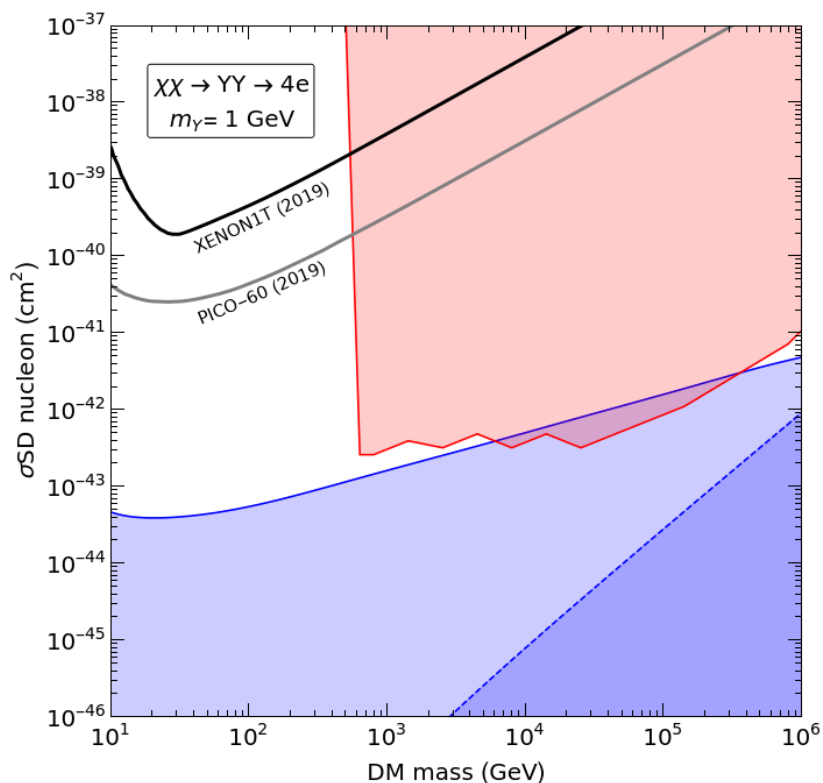


Leane, Ng & Beacom,
[arXiv:1703.04629](#)

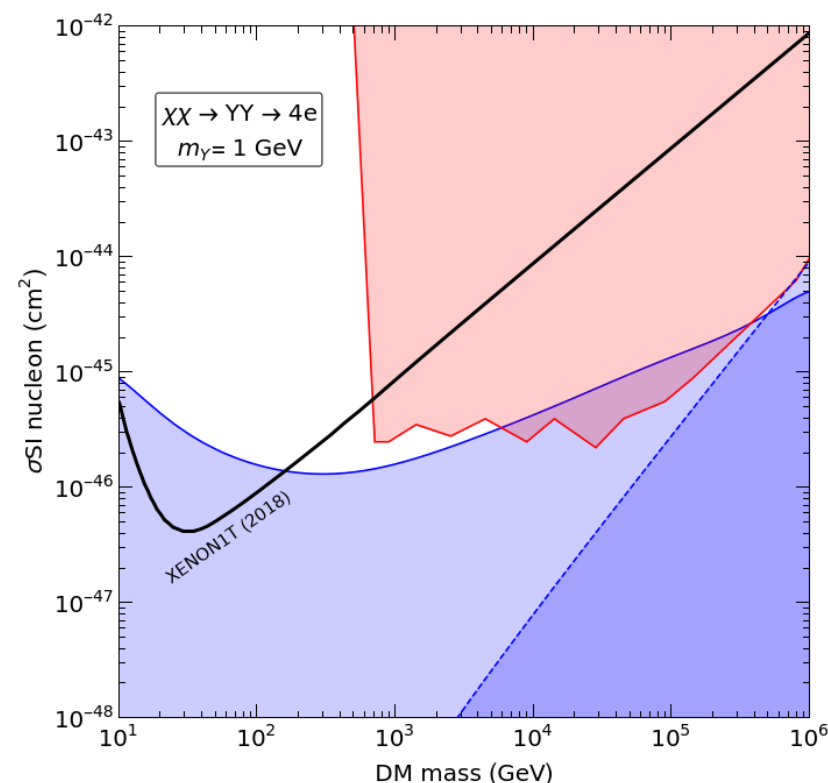
Gamma Rays from the Sun

HAWC gamma ray measurements provide strong constraints,
for both spin-dependent *and* spin-independent scattering

Spin-Dependent (SD)



Spin-Independent (SI)



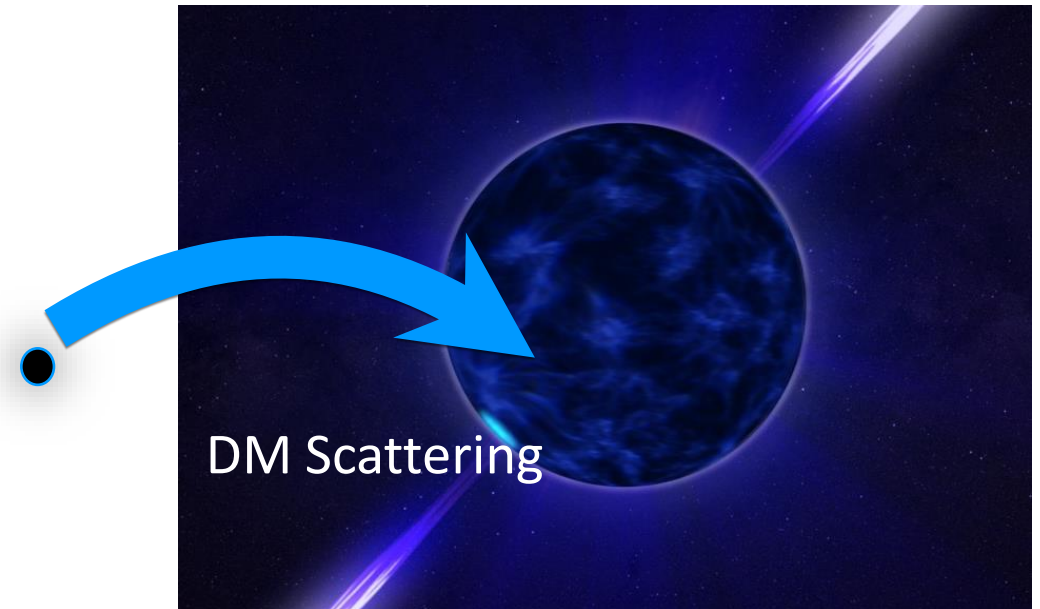
Bell, Dent
& Sanderson,
arXiv:2103.16794

Neutron Stars

Due to their extreme density, *neutron stars* capture dark matter *very* efficiently.

Capture probability saturates at order unity when the cross section satisfies the **geometric limit**

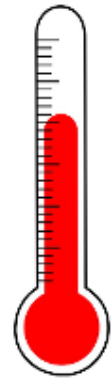
$$\sigma_{th} \sim \pi R^2 \frac{m_n}{M_*} \sim 10^{-45} \text{cm}^2$$



Neutron star heating

→ from dark matter scattering plus annihilation

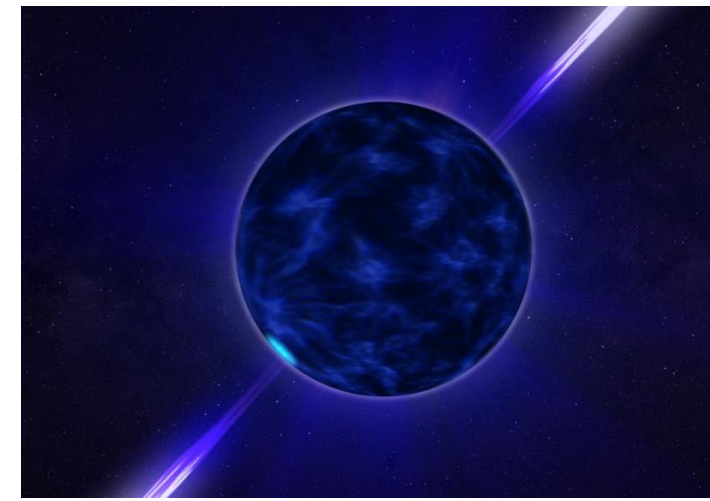
- **Capture** (plus subsequent energy loss)
→ DM *kinetic energy* heats neutron star ~ **1700K**
- **Annihilation** of thermalised dark matter
→ DM *rest mass energy* heats neutron star ~ **additional 700K**



Coollest known neutron star (PSR J2144-3933) has a temperature of $\sim 4.2 \times 10^4$ K

Old isolated neutron stars should cool to below 1000 K after ~ 10 Myr

DM capture in Neutron Stars



Completely different kinematic regime to direct detection experiments, because **DM is relativistic** upon infall to the NS:

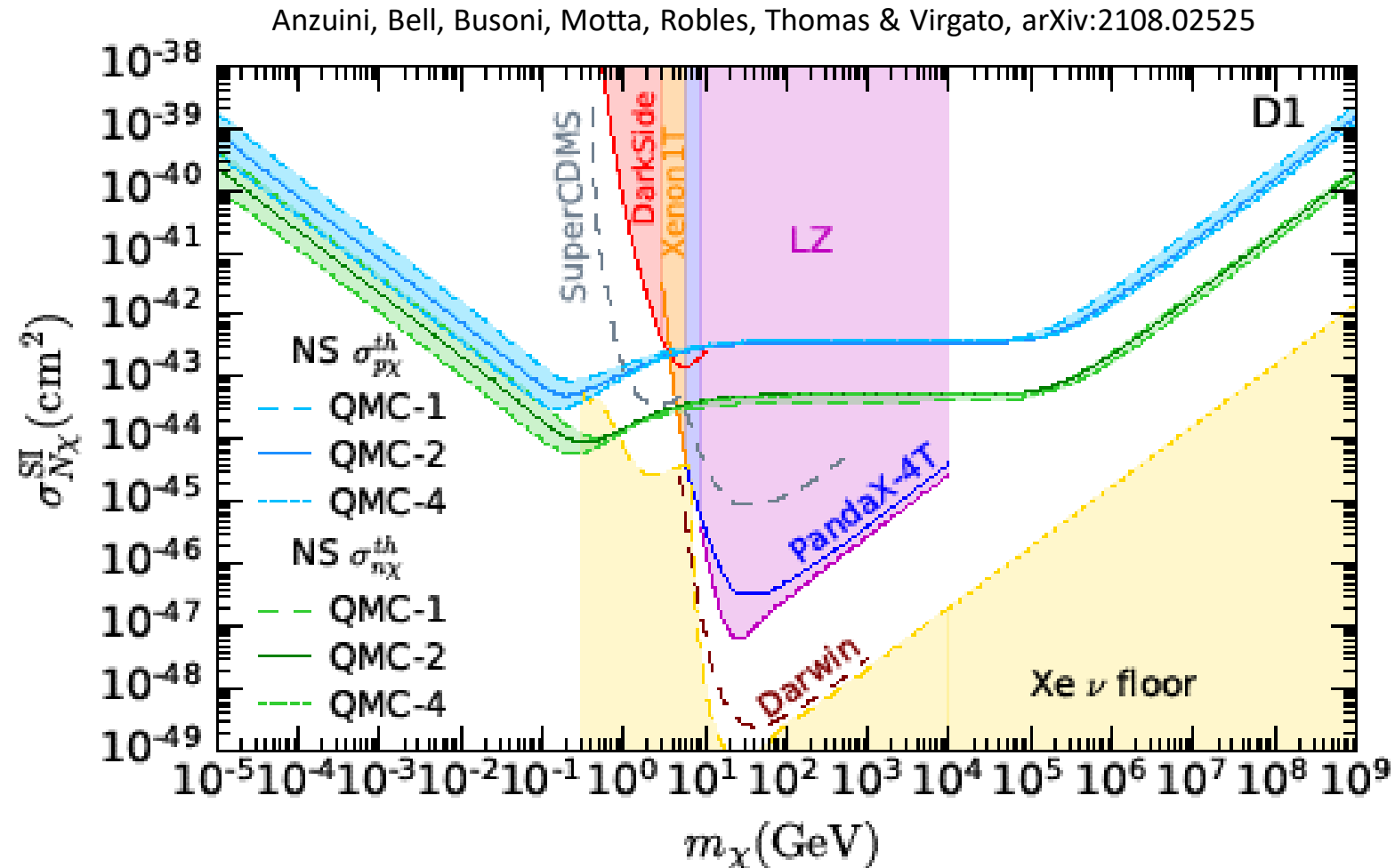
- **No velocity/momentum suppression**
→ *Sensitivity to interactions that direct detection experiments will never be able to see*
- **Must take momentum dependence of hadronic couplings into account**
$$c_n(q) = \frac{c_n^{(0)}}{(1 - q^2/Q_0^2)^2} \quad \text{with } Q_0 \sim 1 \text{ GeV}$$

→ which changes the capture rate by several orders of magnitude

Bell, Busoni, Motta, Robles, Thomas, Virgato, PRL 2021

NS Heating Sensitivity (projected limits)

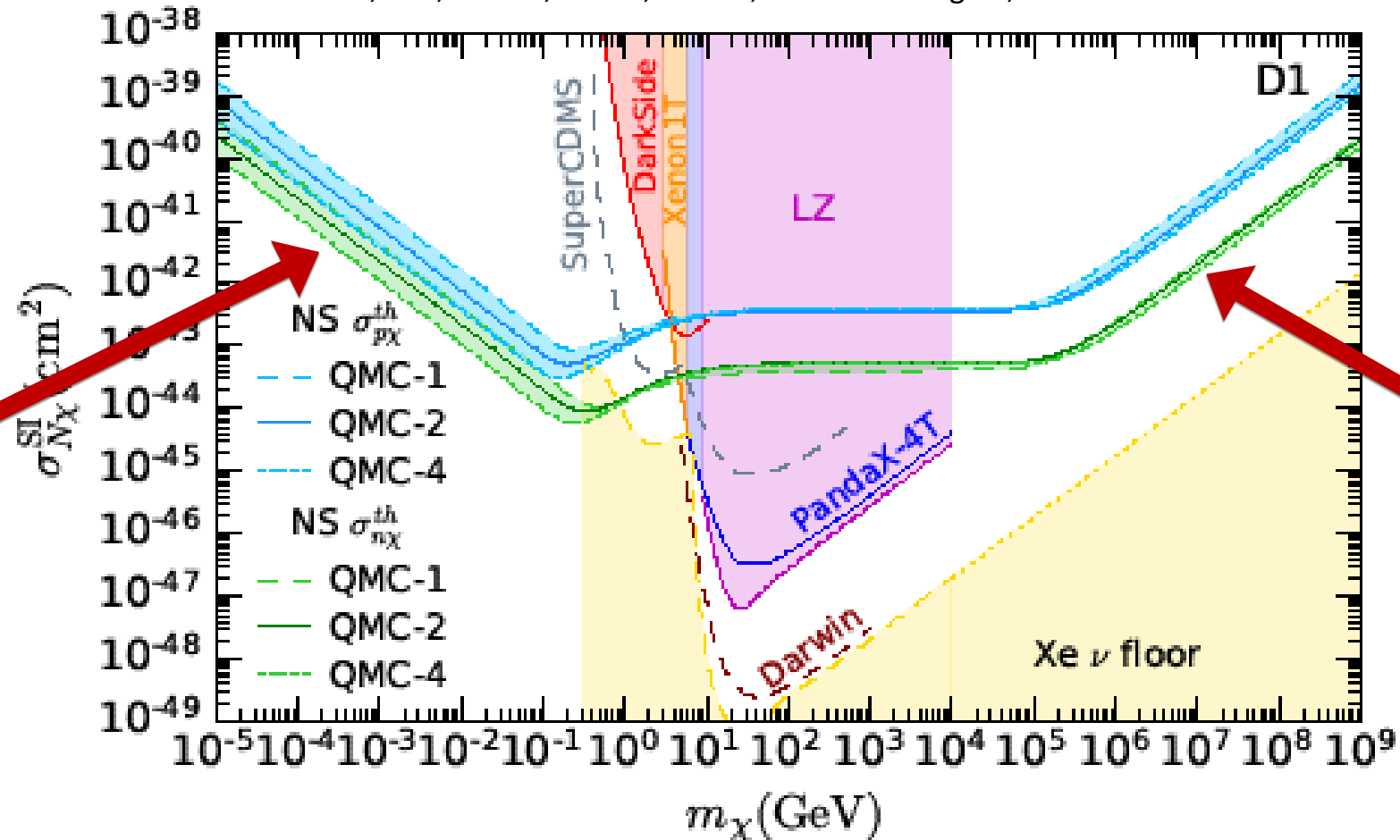
Ball-park sensitivity
= geometric
cross section
 $\sim 10^{-45} \text{ cm}^2$



NS Heating Sensitivity (projected limits)

Ball-park sensitivity
= geometric
cross section
 $\sim 10^{-45} \text{ cm}^2$

Anzuini, Bell, Busoni, Motta, Robles, Thomas & Virgato, arXiv:2108.02525

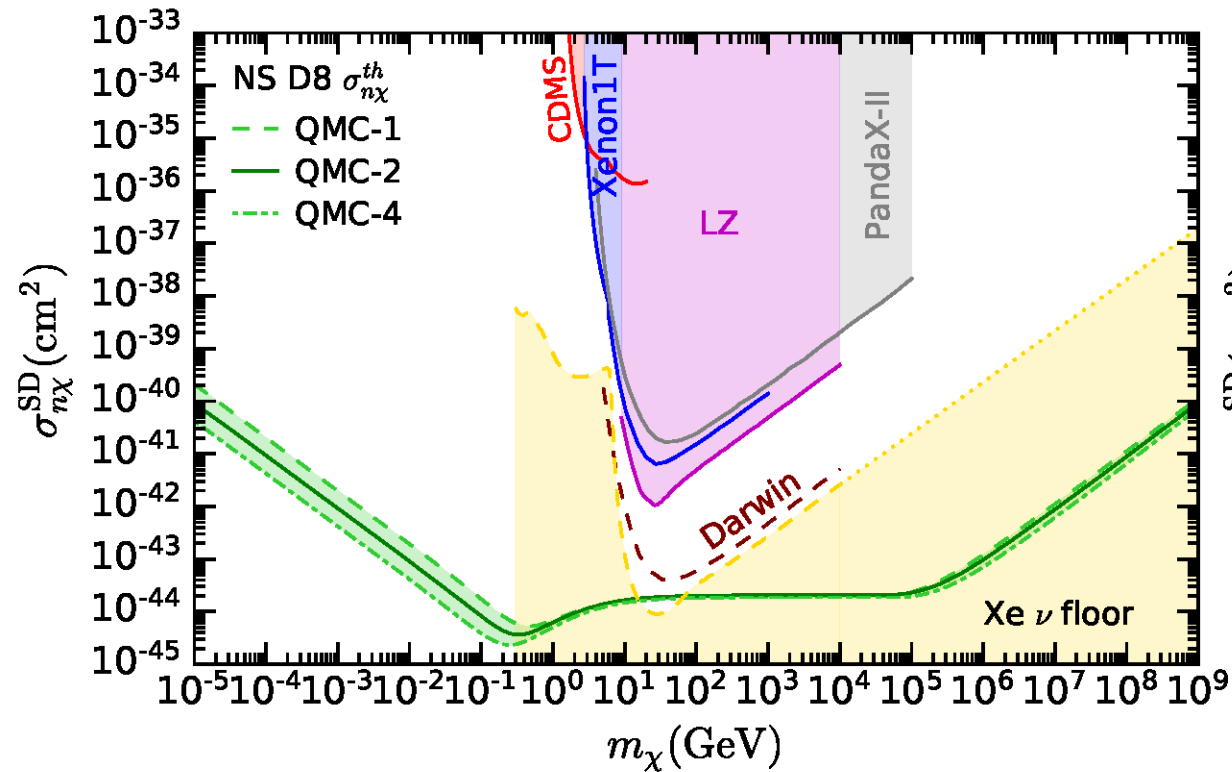


Pauli blocking
from degenerate
neutrons

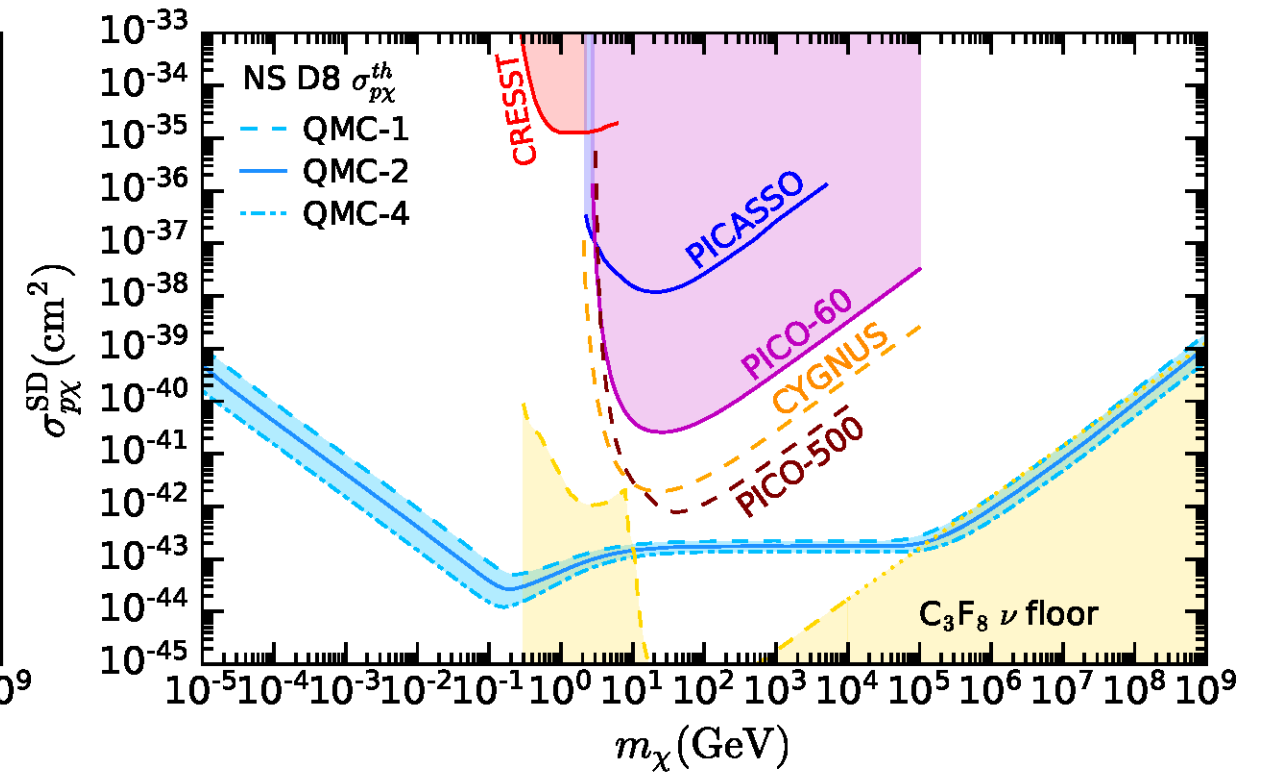
Momentum
transfer in
single collision
not sufficient
for capture

NS Heating Sensitivity: SD nucleon scattering

DM-neutron (SD scattering)

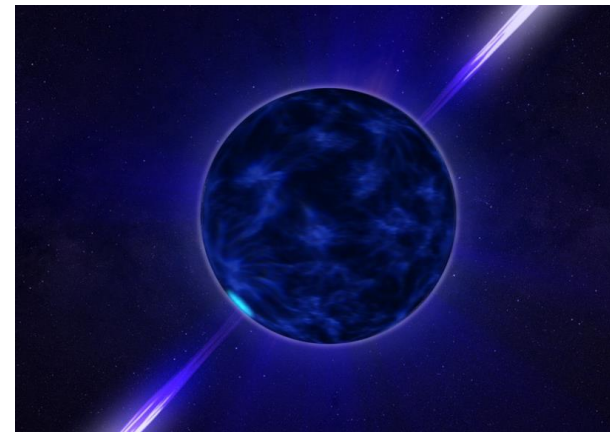


DM-proton (SD scattering)



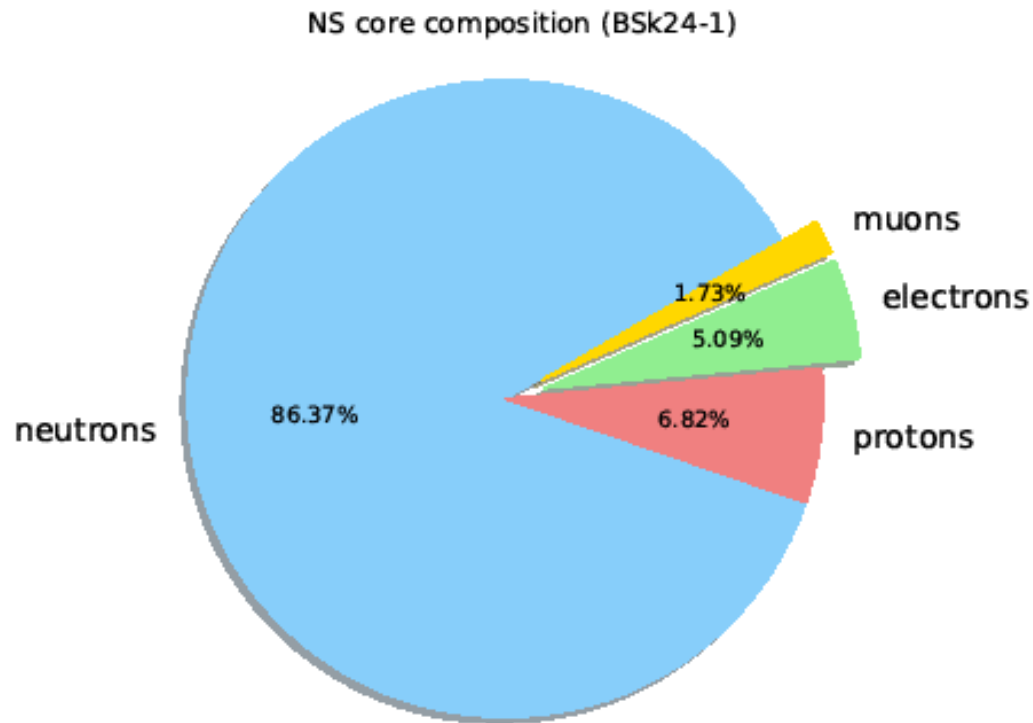
Anzuini, Bell, Busoni, Motta, Robles, Thomas and Virgato, arXiv:2108.02525

Leptons in Neutron Stars

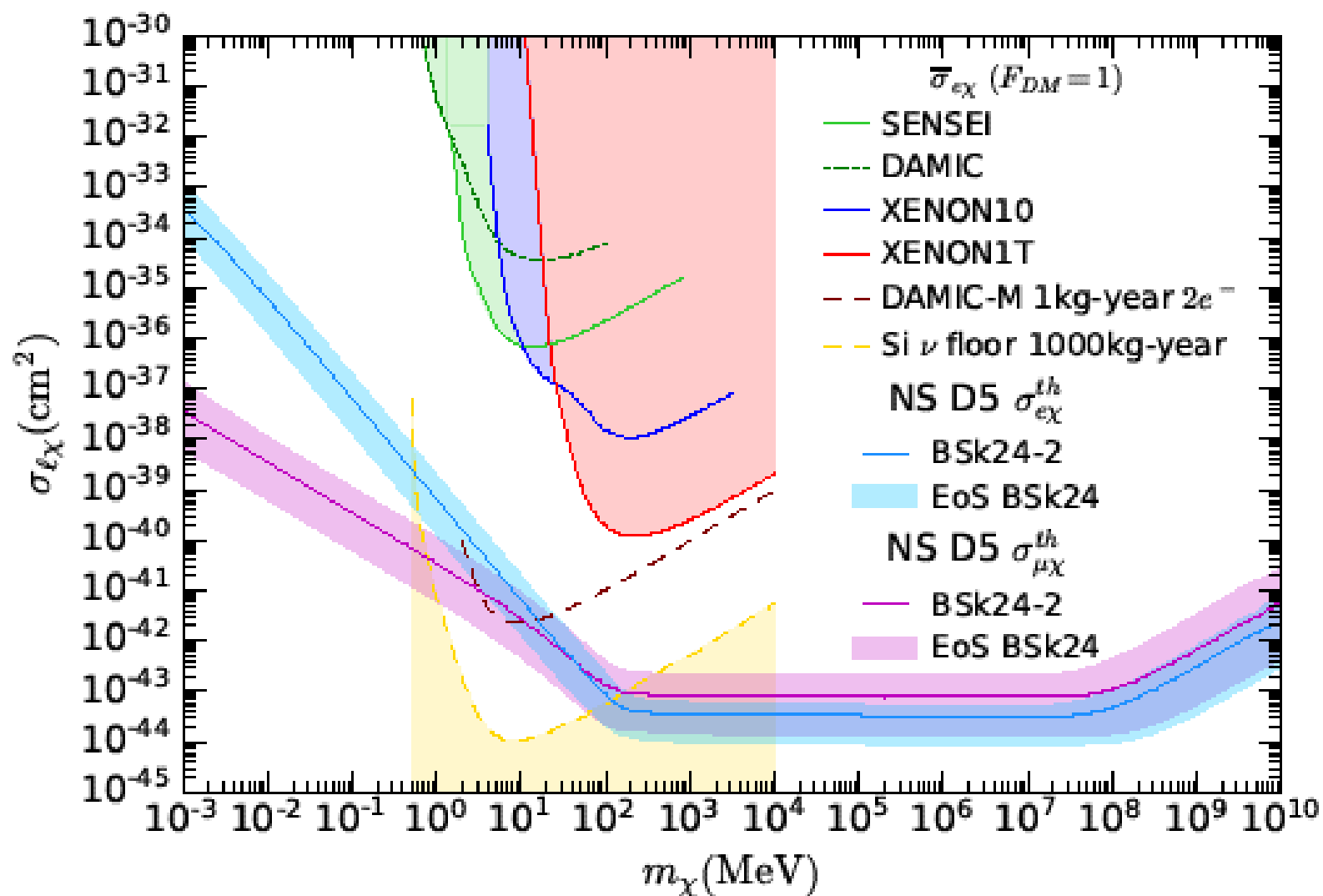


Beta equilibrium in the core determines the composition:

- Degenerate **neutrons**
- Smaller and approximately equal **electron** and **proton** abundances
- Small **muon** component



NS Heating Sensitivity: lepton scattering



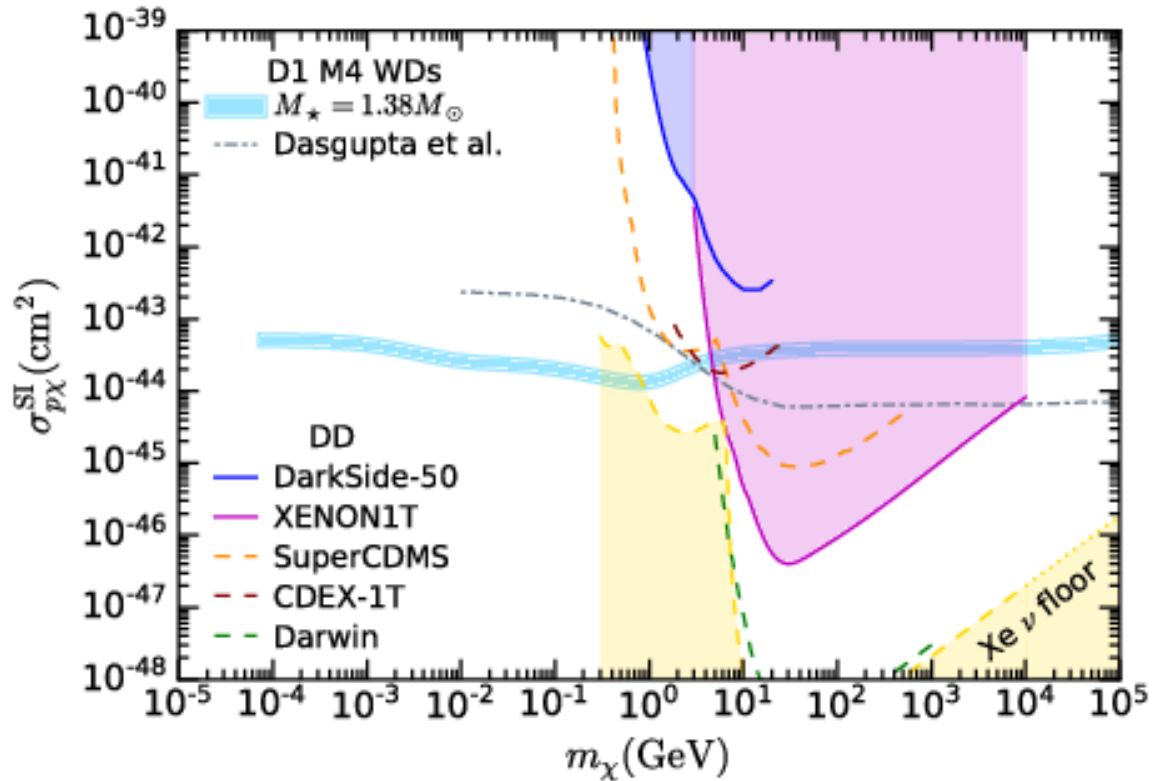
Bell, Busoni, Robles & Virgato arXiv:2010.13257

← Muon scattering

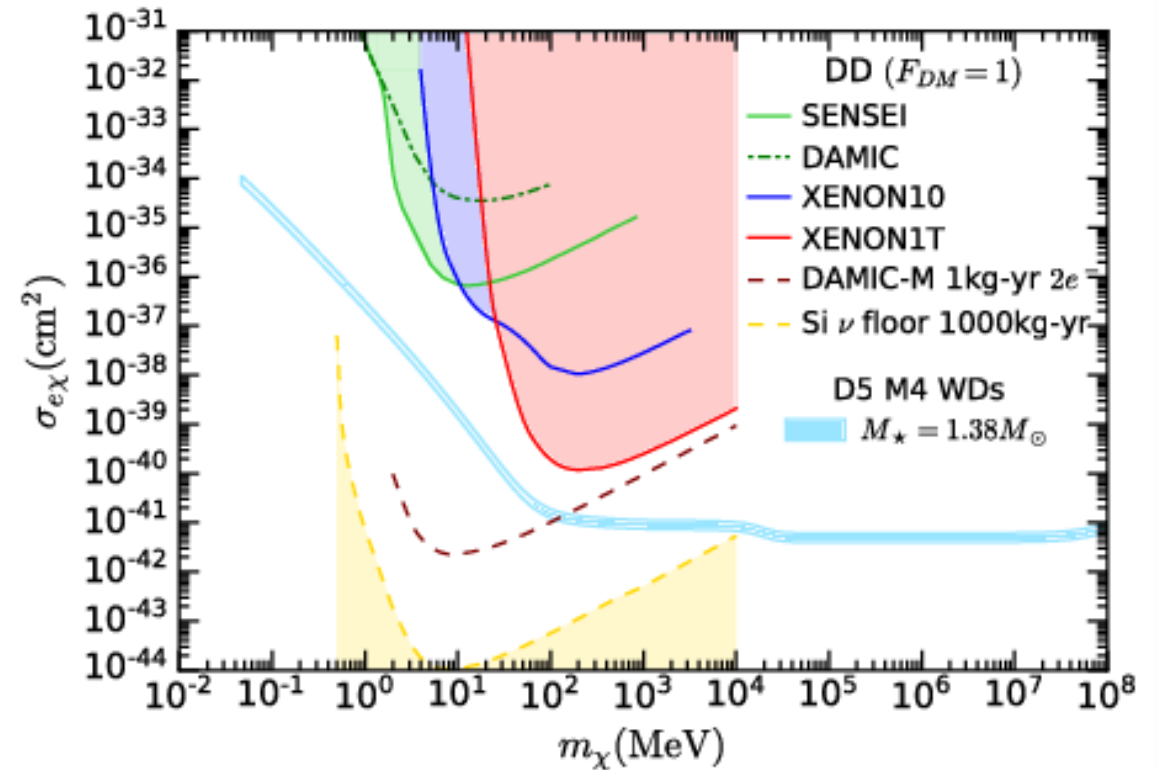
← Electron scattering

White dwarfs in M4 globular cluster

DM-nucleon scattering



DM-electron scattering



Bell, Busoni, Ramirez-Quezada, Robles & Virgato, arXiv:2104.14367

Summary

Testing the thermal-relic hypothesis with indirect detection

- Upcoming observations will make significant progress in closing the WIMP window
- Important to test DM annihilation to neutrinos

Boosted (high energy) dark matter

- Cosmic ray upscattering is inevitable
- Can be seen in neutrino experiments
- Limits are model dependent

Dark matter capture in stars

- Relativistic DM (boosted upon infall to neutron star)
- Can probe low mass dark matter, and look below the neutrino floor