

# *Study of cosmic rays in the ICARUS-T600 detectors*

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on behalf of the ICARUS Collaboration

Padova University, Physics and Astronomy Department and INFN – Sezione di Padova



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101003460, 101081478

## Outline

- SBN program at Fermilab
- ICARUS detectors system
- Preliminary study on cosmic rays light signal



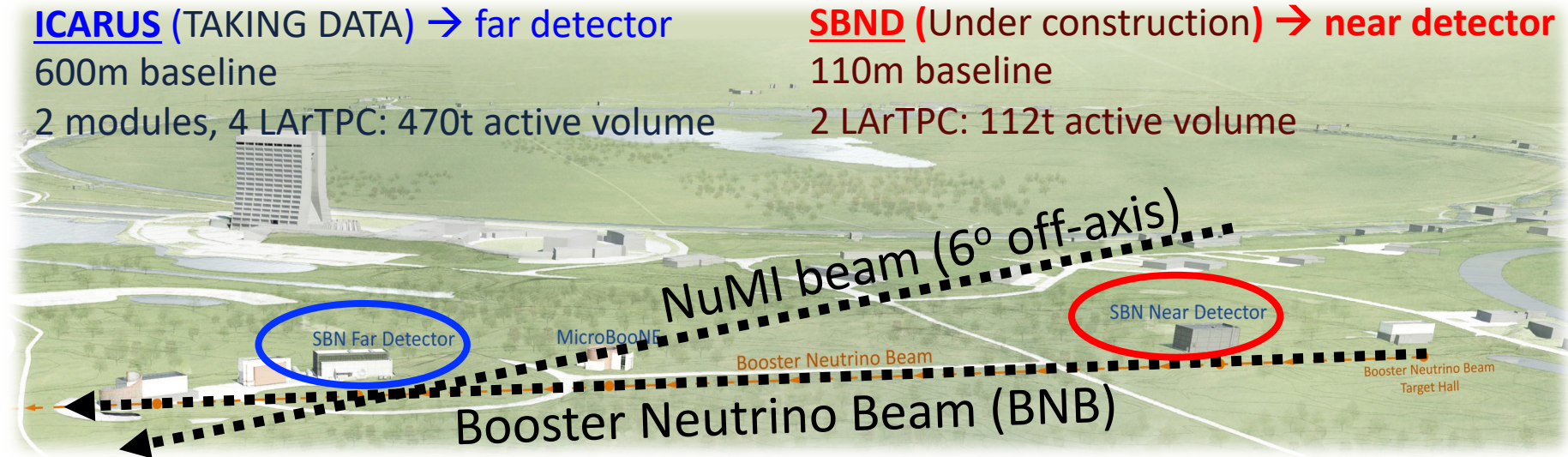
UNIVERSITÀ  
DEGLI STUDI  
DI PADOVA



Dipartimento di Fisica e  
Astronomia  
"Galileo Galilei"

# Short-Baseline Neutrino Program at Fermilab

See D. Mendez talk  
for more details



- **LAr-TPC** located **on-axis** of the **Booster Neutrino Beam (BNB)**
- **Searching for sterile- $\nu$  oscillations** both in **appearance** ( $\nu_e$ ) and **disappearance** ( $\nu_\mu$ ) channels to confirm/rule out previous anomalies from past experiments;
- **High-statistics  $\nu$ -Argon cross-section measurements and event identification/reconstruction studies:**
  - $\sim 10^6$  events/y in SBND  $< 1$  GeV from BNB;
  - $\sim 10^5$  events/y in ICARUS  $> 1$  GeV from **NuMi** off-axis beam.

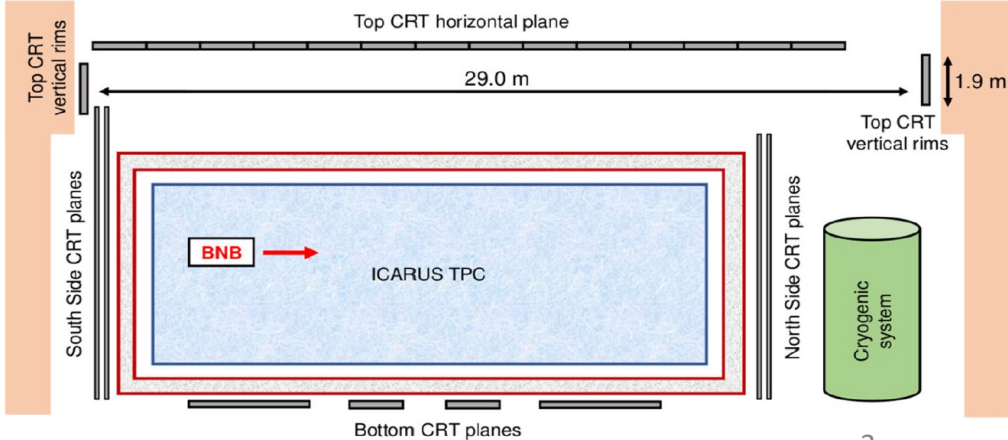
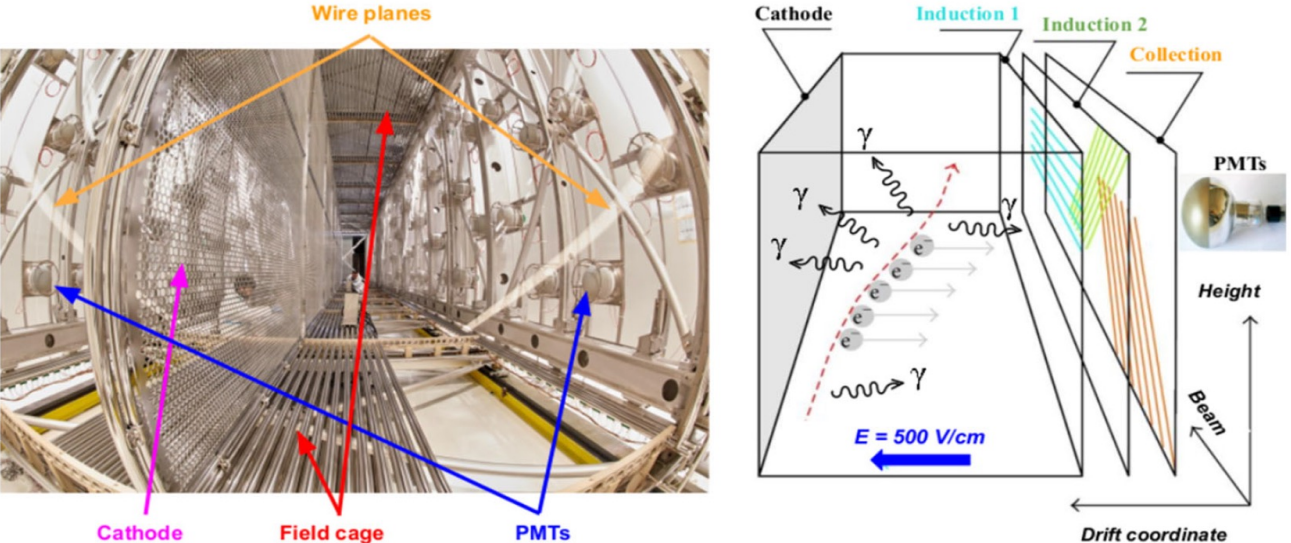
[Phys. Rep. 928:1-63 \(2021\)](#); [Ann.Rev.Nucl.Part.Sci. 69 363-387 \(2019\)](#)

# ICARUS detectors system

**Liquid Argon Time Projection Chambers (LArTPC)** are *high granularity continuously sensitive self-triggering* detectors with *3D imaging and calorimetric reconstruction* capabilities of events with complex topologies **ideal for  $\nu$ -physics** (proposed by C. Rubbia in 1977).

**ICARUS T600 → two identical cryostats**

- **two LArTPC per cryostat with a common cathode** ( $E_{drift} = 500 \text{ V/cm}$ );
- **3 “non-destructive” readout wire planes** with different orientation ( $0^\circ, \pm 60^\circ$ ) continuously read the ionization electrons ( $t_{drift} \sim 1\text{ms}$ ,  $v_{drift} \sim 1.6 \text{ mm/ms}$ );
- **90 PMTs per TPC** located **behind the wires** to collect scintillation light and provide the interaction time and the detector trigger.
- **Cosmic Ray Taggers** surround the cryostats, *tagging incoming cosmics with  $\sim 95\%$  efficiency*.



[Eur. Phys. J. C 83:467 \(2023\)](https://arxiv.org/abs/2308.00001)

# Scintillation light detection system

**360 PMTs** (8" Hamamatsu coated with TPB) installed **behind the TPC wire planes** (5% coverage, 15 ph.e./MeV) **allowing to:**

- Precisely **identify the interaction time of ionizing events** in the TPC (*time resolution  $\sim ns$* ).
- **Localize events** in the PMT plane (**spatial resolution  $< 50$  cm**).
- Roughly **determine the event topologies** for fast event selection.
- **Generate a trigger signal** for readout with a sensitivity to low energy events ( $\sim 100$  MeV).

The system was completed in 2019 and activated after the LAr filling in 2020.

[Eur. Phys. J. C 83:467 \(2023\)](#)



8" Hamamatsu  
R5912-MOD



The new ICARUS PMTs mounted behind the wires of one TPC.

# Trigger system and PMT data acquisition

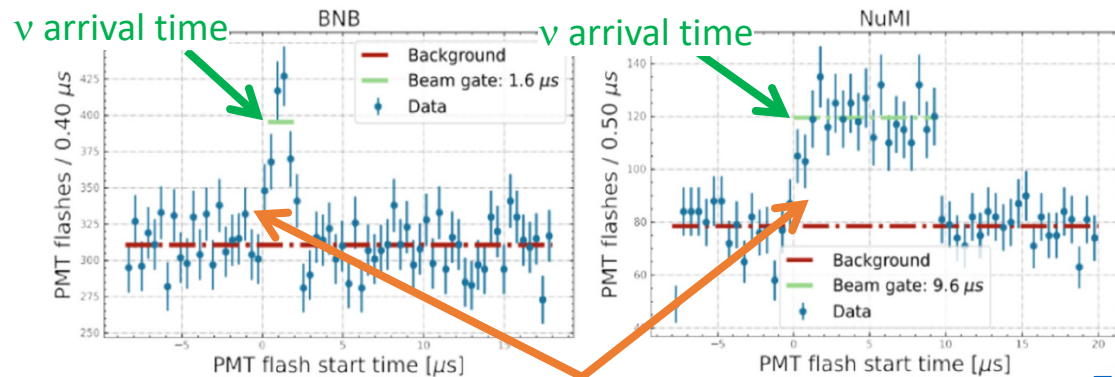
ICARUS main trigger signal is generated by the presence of light signals from PMTs in coincidence with BNB ( $1.6 \mu\text{s}$ ) and NuMI ( $9.6 \mu\text{s}$ ) beam spills defined using the Early Warning signals of proton beam extractions:

- Beam events are **collected requiring at least 5 fired PMT pairs** ( $M_j = 5$ ) inside one of 6 m longitudinal slices equipped with 30+30 opposite PMTs;
- PMT and CRT signals also recorded in 2 ms around the trigger to recognize **cosmics** crossing LAr-TPCs in 1 ms  $e^-$  drift time.



**Additional triggers** (*primitives*) to detect cosmic rays for calibration and background studies for the  $\nu$ -oscillation searches.

- *in beam spills w/o any request on the PMT signals*
- *outside of the beam spills*



Beam event excess visible in PMT light signals

[Eur. Phys. J. C 83:467 \(2023\)](#)

- **Trigger rate**  $\sim 0.7$  Hz (0.3, 0.15 and 0.25 Hz for BNB, NuMI and off beam respectively) setting the **PMT's threshold at 13 ph.e.**
- **Trigger efficiency** is now **under investigation** on data ( $>90\%$  above  $E_{\text{dep}} \sim 100$  MeV)

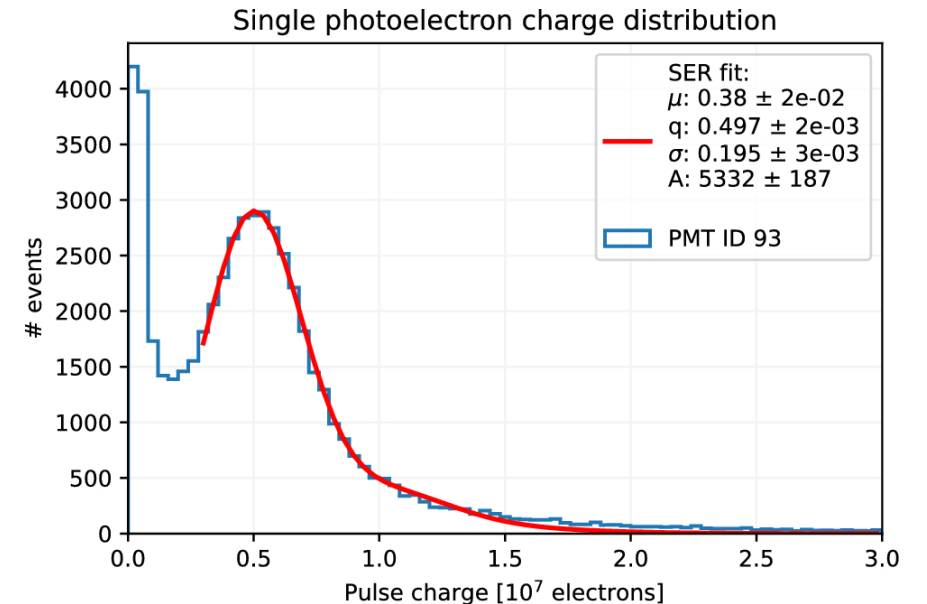
# Study of light signal from cosmic muon

## Monte Carlo simulation

**Scintillation photons** (both from neutrino interaction and from cosmics):

1. **generated** from energy deposition and particle type;
2. **propagated through LAr** → all information of photons (location, time, ...) reaching each PMT are stored exploiting a lookup tables (*photon library*) previously computed;
3. photon by photon the **single photon response** is added;
4. also **noise** is simulated and added to the simulated waveforms;

→ if this signal exceed  $\sim 0.6$  ph.e. **threshold** on a channel, the **waveform** is recorded inside **4  $\mu$ s window**.

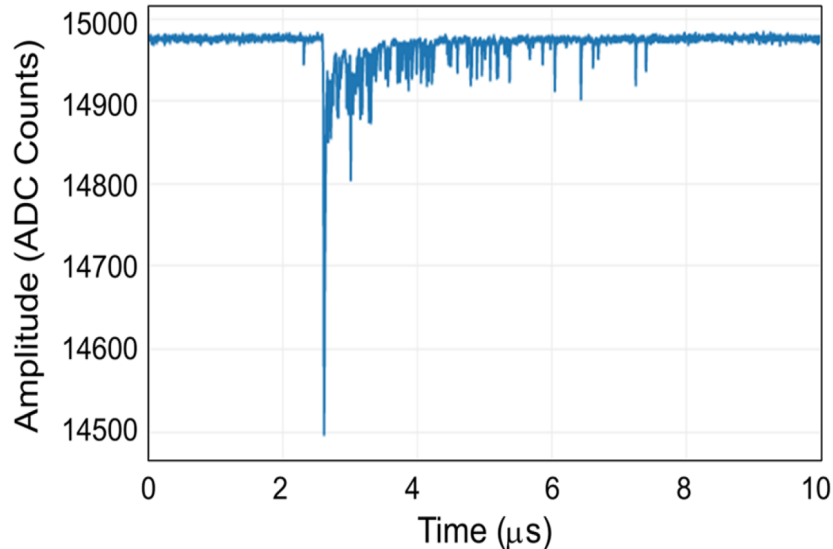


*Example of a typical single photo-electron (ph.e.) charge distribution*

[Eur. Phys. J. C 83:467 \(2023\)](#)

# Study of light signal from cosmic muon

## Number of photoelectron



*PMT signal as recorded by the light detection system electronics.*

The **number of photoelectron** (*n. ph.e.*) is determined independently for each PMT signal (both in DATA and MC):

- **AREA** of the signal:
  1. Subtraction of the **baseline**;
  2. Signals are **discriminated against a fixed level**;

$$\Rightarrow \text{AREA} = \sum_i V_{i,ADC} \text{ (in ADC counts)}$$

- **CHARGE** =  $F_E \cdot \text{AREA}$

→  $F_E$  depends on the electronics characteristics, and allows to convert the stored AREA (ADC unit) in charge:

$$F_E = \frac{k_{ADC \rightarrow V} \cdot \Delta t}{R} = \frac{0.122 \text{ mV/ADC} \cdot 2 \text{ ns}}{50 \Omega}$$

[Eur. Phys. J. C 83:467 \(2023\)](#)

$$n. \text{ ph. e.} = \frac{\text{CHARGE}}{e \text{ (= } 1.6 \cdot 10^{-19} \text{ C)}} \cdot \frac{1}{\text{GAIN}} = \frac{1}{e} \cdot \frac{1}{\text{GAIN}} \cdot F_E \cdot \text{AREA}$$

# Preliminary study of light signal: comparison **MC** vs. DATA

## Samples and selections

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### Analyzed samples:

- Reconstructed **DATA** → cosmic muons entering the TPC;
- Reconstructed **MC** → in time (with beam) single cosmic muons.



# Preliminary study of light signal: comparison **MC** vs. DATA

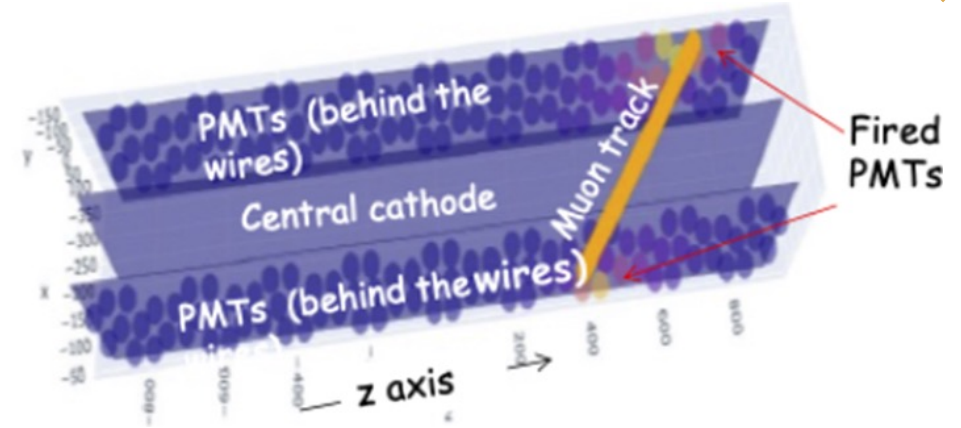
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### → Selections:

- **Tracks** (ionizing particle within LArTPC):
  - passing through the cathode and fully detected in the TPC in order to measure the time of the ionizing particle crossing the cathode ( $t_{\text{track}}$ )
  - longer than 50 cm.



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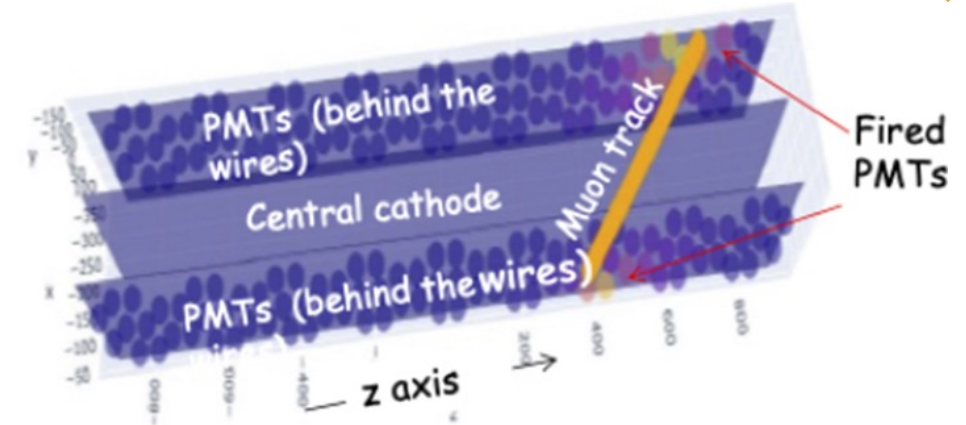
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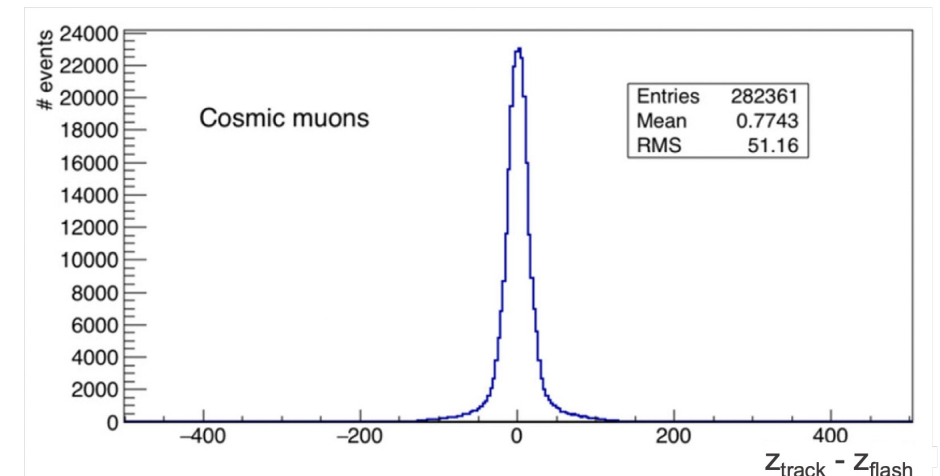
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  - passing through the cathode and fully detected in the TPC in order to measure the time of the ionizing particle crossing the cathode ( $t_{\text{track}}$ )
  - longer than 50 cm.
- **Flashes** (time-coincident (within  $\sim 1\mu\text{s}$ ) ph.e. signals across multiple PMTs aiming for the reconstruction of an interaction).  
The first fired PMT provides the flash time ( $t_{\text{flash}}$ ).
  - in coincidence with track:
    - time difference:  $\Delta t = t_{\text{track}} - t_{\text{flash}} (\lesssim 4 \mu\text{s})$
    - barycenter difference:  $\Delta z = z_{\text{track}} - z_{\text{flash}} (< 30 \text{ cm})$

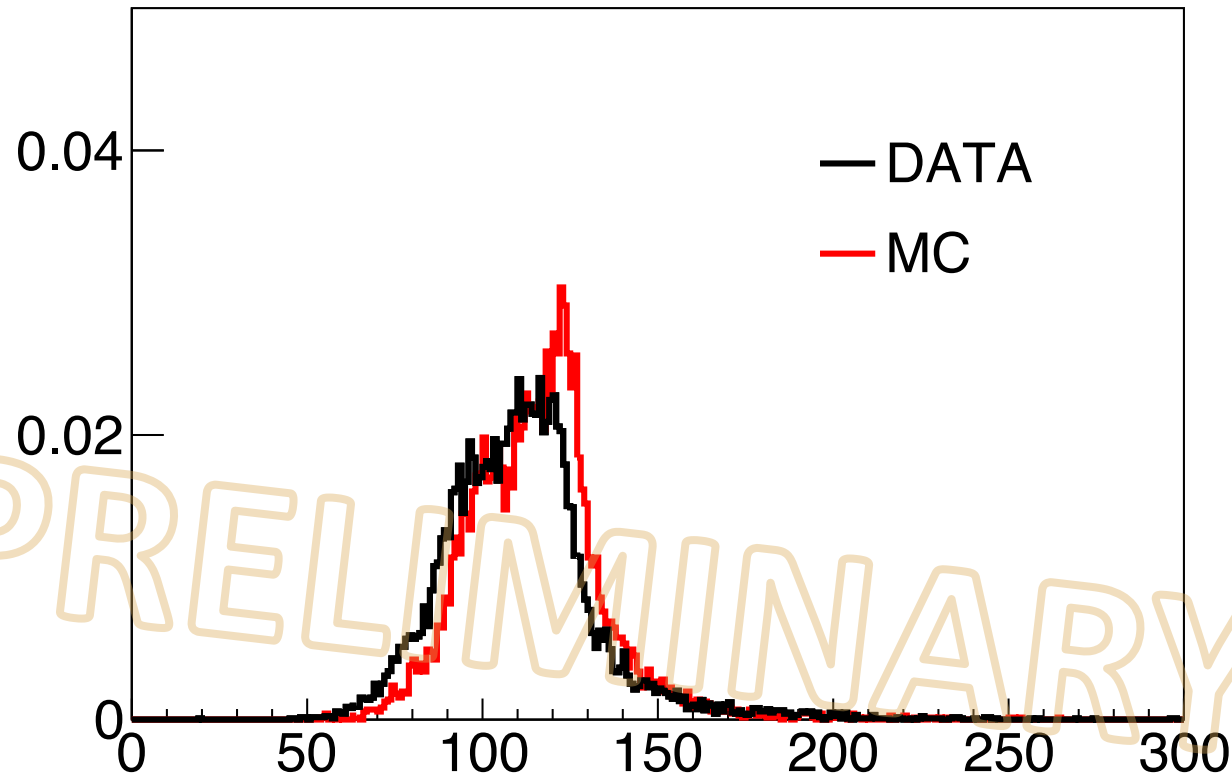


[Eur. Phys. J. C 83:467 \(2023\)](#)

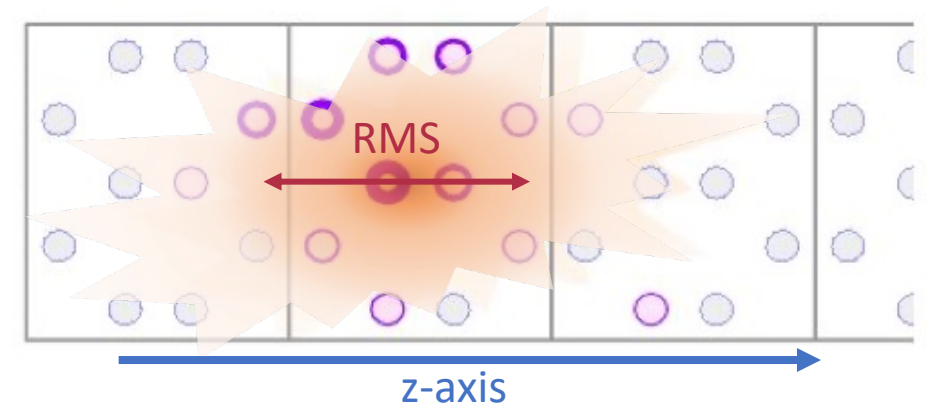


# Preliminary study of light signal: comparison **MC** vs. DATA

## Comparison of light signal extension along the beam direction



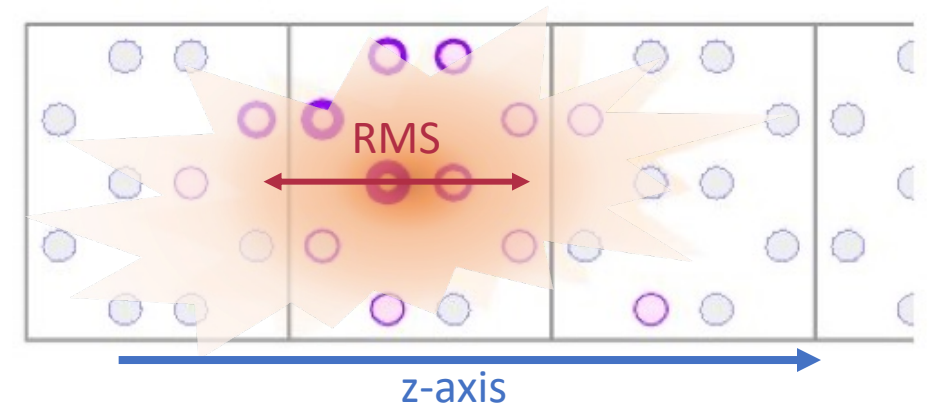
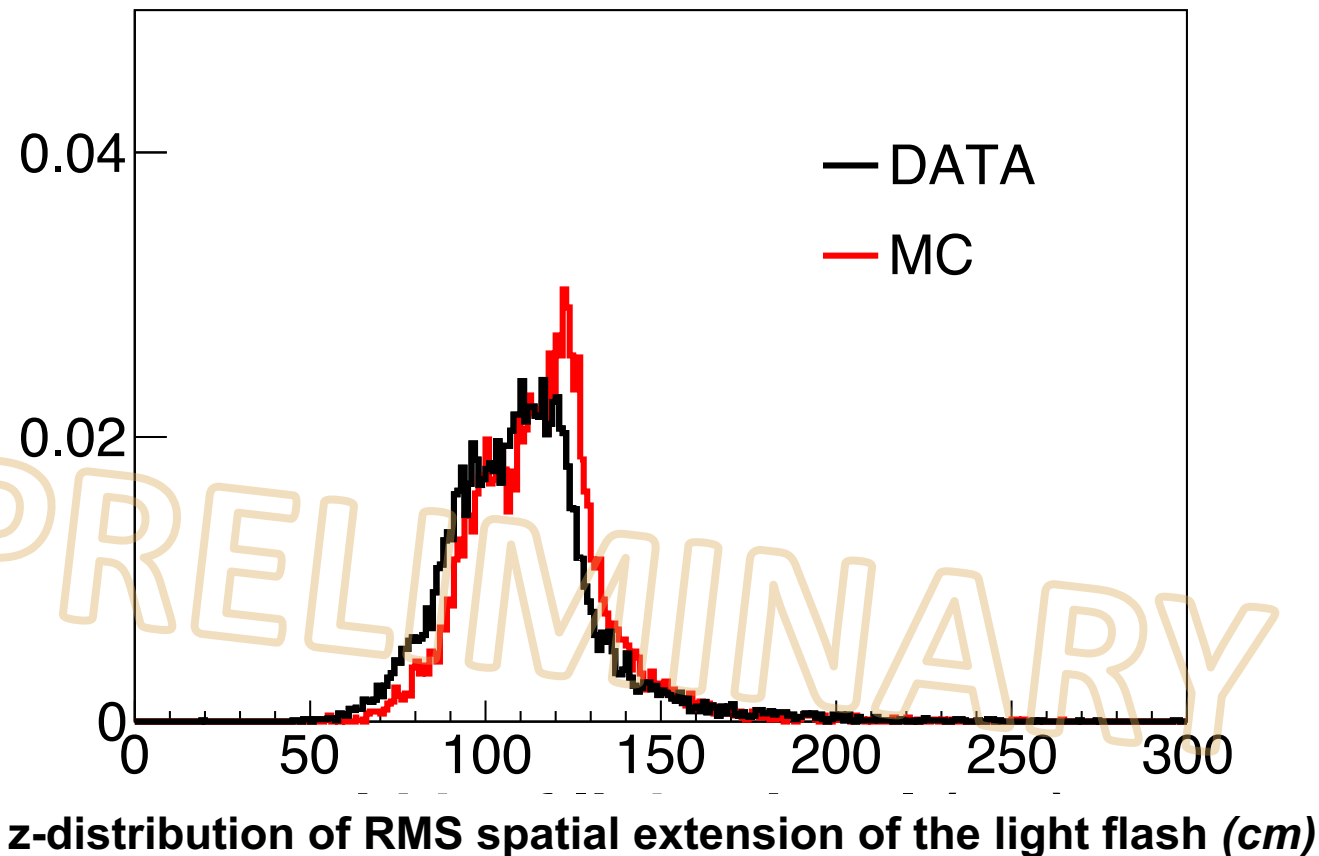
**z-distribution of RMS spatial extension of the light flash (cm)**



For each light flash, the position of its barycenter and its **RMS spatial extension along the beam axis z** are determined accounting for the number of **collected ph.e. on each fired PMT**.

# Preliminary study of light signal: comparison **MC** vs. DATA

## Comparison of light signal extension along the beam direction



For each light flash, the position of its barycenter and its **RMS spatial extension along the beam axis z** are determined accounting for the number of **collected ph.e. on each fired PMT**.

**Good agreement of MC prediction on the scintillation light extension with DATA.**

# Preliminary study of light signal: comparison **MC** vs. DATA

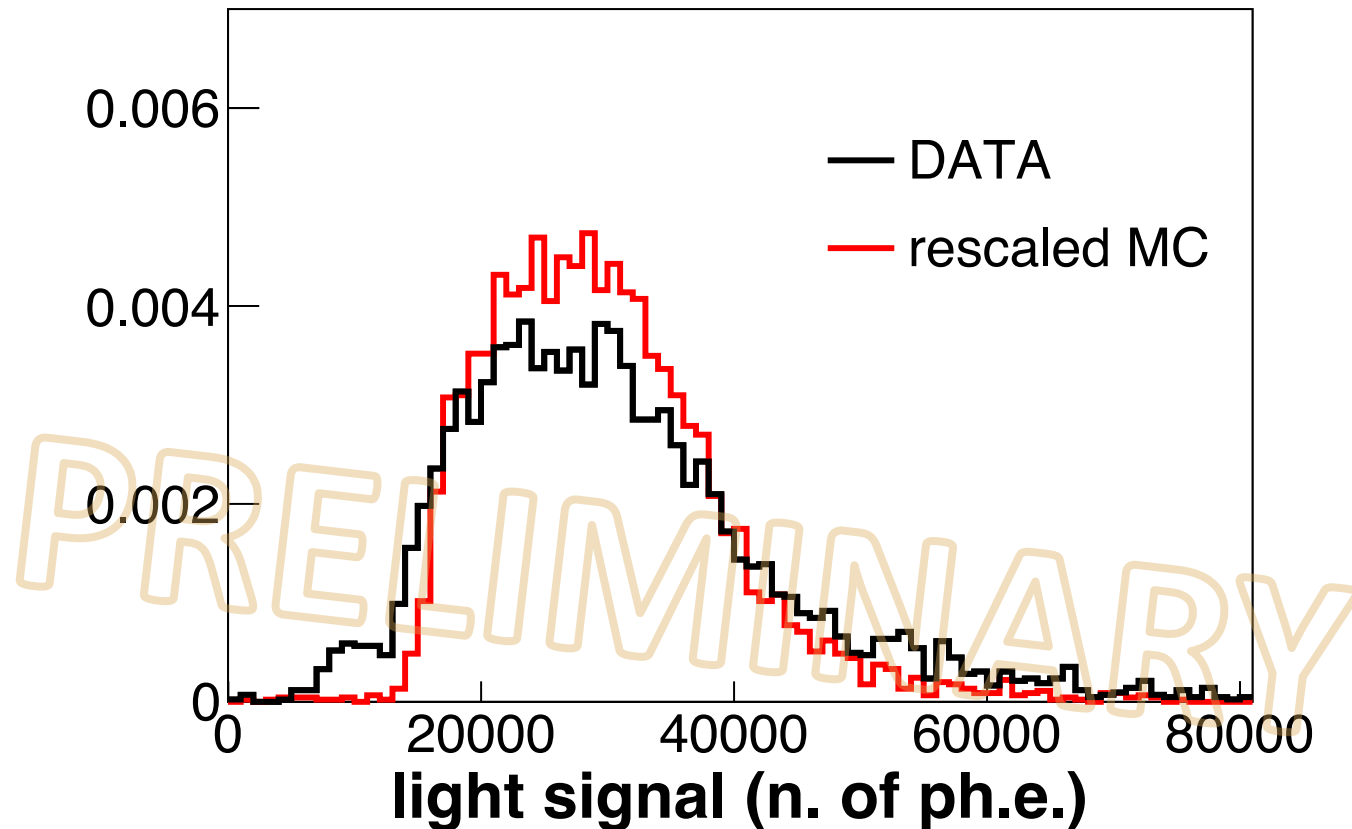
## Rescaling of the MC light signal (ph.e.)

At moment, several factors are **under study**:

- *quantum efficiency* at LAr temperature
- *effects of the refraction index* of LAr
- *reconstruction artifacts*
- ...



In the meanwhile, the **MC signal is rescaled** to the same average value as measured in **DATA** by fitting the ph.e. curve.



# Preliminary study of light signal: comparison **MC** vs. DATA

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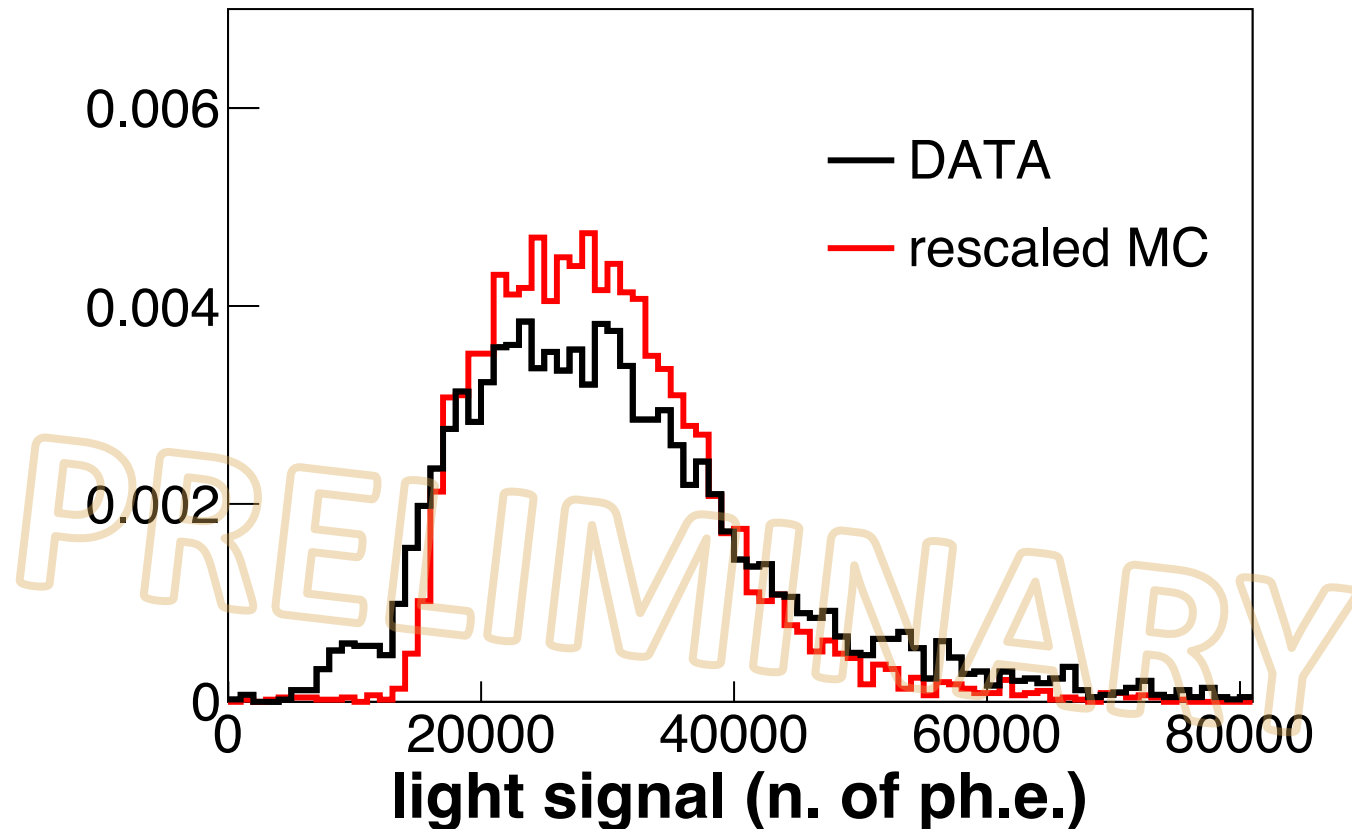
- *quantum efficiency* at LAr temperature
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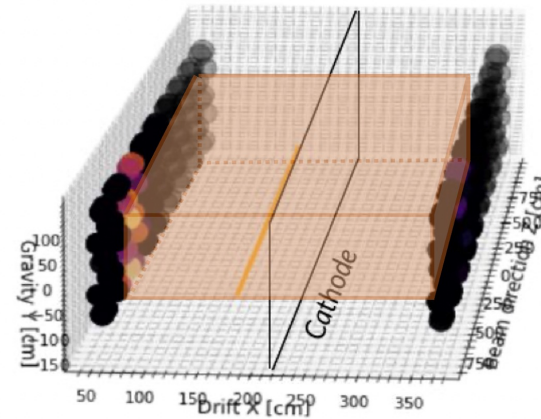
This allows to **open the possibility to study** with rescaled MC the **features of events associate with light** to be later applied to event's selection.



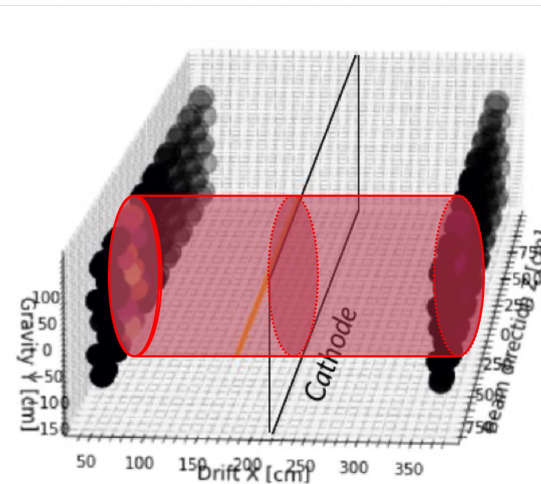
# What's next...

- **Deeper analysis** on selected detector regions (i.e. center, boundary, close to the fired PMTs, far from the to the fired PMTs, ...)
  - **First check:** focus on events far from detector boundary (i.e. trigger threshold, boundary acceptance, ... ) to be independent to reconstruction.
- **Deeper analysis** selecting specific tracks inclination (i.e. vertical, horizontal, ...) and length.
- **Check and measure of the relevant parameters of scintillation light** production and collection by PMTs in order to improve the MC simulation.

Example (sketch) of detector's region selection



Center of the detector:  
far from the boundary!



Close to the fired PMTs

# Summary

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The preliminary study of LIGHT SIGNAL related to cosmic muon in the LArTPC is presented.

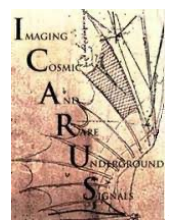
- The **MC** simulation of the **light signal extension along the beam direction** for cosmic muon is **well describing** the **collected DATA**.
- Relevant parameters in the optical simulation of the scintillation light, which affects the signal detected by PMTs are still under study.
- A **preliminary rescaling** of light simulated by **MC** to the collected DATA is applied **to study the features of events associate with light** in order **to develop event selection criteria and analysis**.





*Thank you for your attention!*





# The ICARUS collaboration

12 INFN groups, 12 US institutions, CERN,  
1 Mexican institution, 1 Indian Institution

*P. Abratenko<sup>19</sup>, A. Aduszkiewicz<sup>21</sup>, F. Akbar<sup>23</sup>, M. Artero Pons<sup>15</sup>, J. Asaadi<sup>24</sup>, M. Babicz<sup>2</sup>, W.F. Badgett<sup>5</sup>, L.F. Bagby<sup>5</sup>, B. Baibussinov<sup>15</sup>, B. Behera<sup>4</sup>, V. Bellini<sup>7</sup>, O. Beltramello<sup>2</sup>, R. Benocci<sup>13</sup>, J. Berger<sup>4</sup>, S. Berkman<sup>5</sup>, S. Bertolucci<sup>6</sup>, M. Betancourt<sup>5</sup>, K. Biery<sup>5</sup>, M. Bonesini<sup>13</sup>, T. Boone<sup>4</sup>, B. Bottino<sup>8</sup>, A. Braggiotti<sup>15</sup>, J. Bremer<sup>2</sup>, S. Brice<sup>5</sup>, V. Brio<sup>7</sup>, C. Brizzolari<sup>13</sup>, J. Brown<sup>5</sup>, H. Budd<sup>23</sup>, A. Campani<sup>8</sup>, A. Campos<sup>27</sup>, D. Carber<sup>4</sup>, M. Carneiro<sup>1</sup>, H. Carranza<sup>24</sup>, D. Casazza<sup>8</sup>, A. Castro<sup>3</sup>, M. Cicerchia<sup>15</sup>, S. Centro<sup>15</sup>, G. Cerati<sup>5</sup>, M. Chalifour<sup>2</sup>, A. Chatterjee<sup>26</sup>, D. Cherdack<sup>21</sup>, S. Cherubini<sup>11</sup>, N. Chitirasreemadam<sup>25</sup>, T. Coan<sup>18</sup>, A. Cocco<sup>14</sup>, M. R. Convery<sup>17</sup>, S. Copello<sup>16</sup>, A. De Roeck<sup>2</sup>, S. Di Domizio<sup>8</sup>, D. Di Ferdinando<sup>6</sup>, L. Di Noto<sup>8</sup>, M. Diwan<sup>1</sup>, S. Donati<sup>25</sup>, J. Dyer<sup>4</sup>, S. Dytman<sup>22</sup>, S. Dolan<sup>2</sup>, F. Dolek<sup>27</sup>, L. Domine<sup>17</sup>, R. Doubnik<sup>5</sup>, F. Drielsma<sup>17</sup>, C. Fabre<sup>2</sup>, A. Falcone<sup>13</sup>, C. Farnese<sup>15</sup>, A. Fava<sup>5</sup>, F. Ferraro<sup>8</sup>, F. Garcia<sup>17</sup>, C. Gatto<sup>14</sup>, M. Geynisman<sup>5</sup>, D. Gibin<sup>15</sup>, A. Gioiosa<sup>25</sup>, W. Gu<sup>1</sup>, M. Guerzoni<sup>6</sup>, A. Guglielmi<sup>15</sup>, S. Hahn<sup>5</sup>, A. Heggstuen<sup>4</sup>, B. Howard<sup>5</sup>, R. Howell<sup>23</sup>, J. Hrivnak<sup>2</sup>, C. James<sup>5</sup>, W. Jang<sup>24</sup>, L. Kashur<sup>4</sup>, W. Ketchum<sup>5</sup>, J.S. Kim<sup>23</sup>, D.H. Koh<sup>17</sup>, U. Kose<sup>2</sup>, J. Larkin<sup>1</sup>, G. Laurenti<sup>6</sup>, G. Lukhanin<sup>5</sup>, A. Maria<sup>26</sup>, C. Mariani<sup>27</sup>, C. Marshall<sup>23</sup>, S. Martinenko<sup>1</sup>, N. Mauri<sup>6</sup>, A. Mazzacane<sup>5</sup>, K.S. McFarland<sup>23</sup>, D.P. Mendez<sup>1</sup>, G. Meng<sup>15</sup>, A. Menegolli<sup>16</sup>, O.G. Miranda<sup>3</sup>, D. Mladenov<sup>2</sup>, A. Mogan<sup>4</sup>, N. Moggi<sup>6</sup>, N. Montagna<sup>6</sup>, A. Montanari<sup>6</sup>, C. Montanari<sup>5,b</sup>, M. Mooney<sup>4</sup>, G. Moreno Granados<sup>3</sup>, J. Mueller<sup>4</sup>, M. Murphy<sup>27</sup>, D. Naples<sup>22</sup>, M. Nessi<sup>2</sup>, T. Nichols<sup>5</sup>, S. Palestini<sup>2</sup>, M. Pallavicini<sup>8</sup>, V. Paolone<sup>22</sup>, R. Papaleo<sup>11</sup>, L. Pasqualini<sup>6</sup>, L. Patrizii<sup>6</sup>, G. Petrillo<sup>17</sup>, C. Petta<sup>7</sup>, V. Pia<sup>6</sup>, F. Pietropaolo<sup>2,a</sup>, F. Poppi<sup>6</sup>, M. Pozzato<sup>6</sup>, A. Prosser<sup>5</sup>, G. Putnam<sup>20</sup>, X. Qian<sup>1</sup>, A. Rappoldi<sup>16</sup>, R. Rechenmacher<sup>5</sup>, L. Rice<sup>22</sup>, E. Richards<sup>22</sup>, F. Resnati<sup>2</sup>, A.M. Ricci<sup>25</sup>, A. Rigamonti<sup>2</sup>, G.L. Raselli<sup>16</sup>, M. Rosemberg<sup>19</sup>, M. Rossella<sup>16</sup>, C. Rubbia<sup>9</sup>, G. Savage<sup>5</sup>, A. Scaramelli<sup>16</sup>, D. Schmitz<sup>20</sup>, A. Schukraft<sup>5</sup>, F. Sergiampietri<sup>2</sup>, G. Sirri<sup>6</sup>, J. Smedley<sup>23</sup>, A. Soha<sup>5</sup>, L. Stanco<sup>15</sup>, J. Stewart<sup>1</sup>, N.B. Suarez<sup>22</sup>, H. Tanaka<sup>17</sup>, M. Tenti<sup>6</sup>, K. Terao<sup>17</sup>, F. Terranova<sup>13</sup>, V. Togo<sup>6</sup>, D. Torretta<sup>5</sup>, M. Torti<sup>13</sup>, Y.T. Tsai<sup>17</sup>, S. Tufanli<sup>2</sup>, T. Usher<sup>17</sup>, F. Varanini<sup>15</sup>, S. Ventura<sup>15</sup>, M. Vicenzi<sup>1</sup>, C. Vignoli<sup>10</sup>, B. Viren<sup>1</sup>, D. Warner<sup>4</sup>, Z. Williams<sup>24</sup>, P. Wilson<sup>5</sup>, R.J. Wilson<sup>4</sup>, J. Wolfs<sup>23</sup>, T. Wongjirad<sup>19</sup>, A. Wood<sup>21</sup>, E. Worcester<sup>1</sup>, M. Worcester<sup>1</sup>, M. Wospakrik<sup>5</sup>, H. Yu<sup>1</sup>, J. Yu<sup>24</sup>, A. Zani<sup>12</sup>, C. Zhang<sup>1</sup>, J. Zennamo<sup>5</sup>, J. Zettlemoyer<sup>5</sup>, S. Zucchelli<sup>6</sup>, M. Zuckerbrot<sup>5</sup>*

**Spokeperson: C. Rubbia, GSSI**

1. Brookhaven National Lab., USA
2. CERN, Switzerland
3. CINVESTAV, Mexico,
4. Colorado State University, USA
5. Fermi National Accelerator Lab., USA
6. INFN Bologna and University, Italy
7. INFN Catania and University, Italy
8. INFN Genova and University, Italy
9. INFN GSSI, L'Aquila, Italy
10. INFN LNGS, Assergi, Italy
11. INFN LNS, Catania, Italy
12. INFN Milano, Milano, Italy
13. INFN Milano Bic. and University, Italy
14. INFN Napoli, Napoli, Italy
15. INFN Padova and University, Italy
16. INFN Pavia and University, Italy
17. SLAC National Accelerator Lab., USA
18. Southern Methodist University, USA
19. Tufts University, USA
20. University of Chicago, USA
21. University of Houston, USA
22. University of Pittsburgh, USA
23. University of Rochester, USA
24. University of Texas (Arlington), USA
25. INFN Pisa and University, Italy
26. Ramanujan Faculty Phys. Res. India
27. Virginia Tech Institute

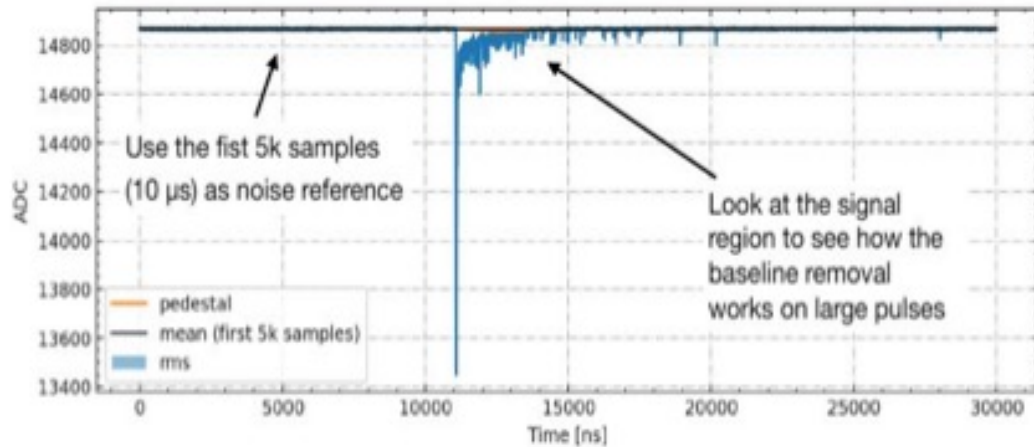
- a. On Leave of Absence from INFN Padova
- b. On Leave of Absence from INFN Pavia



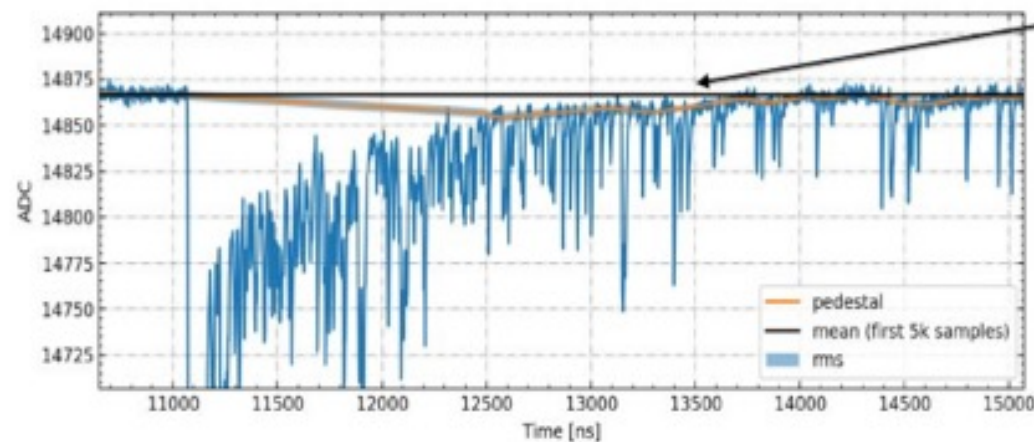
**SPARE SLIDES**

# Study of light signal (n.ph.e.)

## Number of photoelectron: AREA of the light signal



Zooming on the signal region it appears evident how this algorithm smoothen out some tail removing part of the charge contained in it



If we integrate the charge between 10 μs to 24 μs we find the following:

- Total charge integral > 5 ADC ( using mean subtraction ) : 94764
- Total charge integral > 5 ADC ( using pedestal removal ) : 84954

Ratio is 89 %

The number of ph.e. is determined independently for each PMT signal (both in DATA and MC):

- **AREA** of the signal:
  1. Subtraction of the **baseline**;
  2. Signals are **discriminated against a fixed level**;

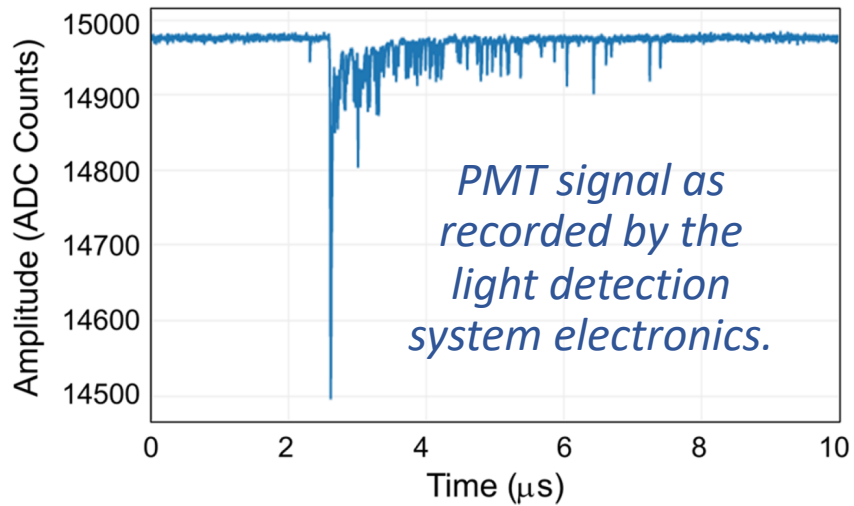
$$\Rightarrow \text{AREA} = \sum_i V_i \cdot k_{V \rightarrow \text{ADC}}$$

where:

- $i$  = index of sampling of the digitizer;
- $V$  = voltage read by the digitizer after the baseline subtraction.

# Study of light signal (n.ph.e.)

## Number of photoelectron



The number of ph.e. is determined independently for each PMT signal (both in DATA and MC):

- **AREA** of the signal:
  1. Subtraction of the **baseline**;
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where:

  - i = index of sampling of the digitizer;
  - V = voltage read by the digitizer after the baseline subtraction.

- **CHARGE** = AREA ·  $F_E$

- $F_E$  depends on the electronics characteristics:

$$F_E = \frac{k_{\text{ADC} \rightarrow \text{V}} \cdot \Delta t}{R} = \frac{0.122 \text{ mV/ADC} \cdot 2 \text{ ns}}{50 \Omega} = \frac{0.122 \cdot 2 \cdot 10^{-12}}{50} \left[ \frac{\text{V/ADC} \cdot \text{s}}{\Omega} \right]$$

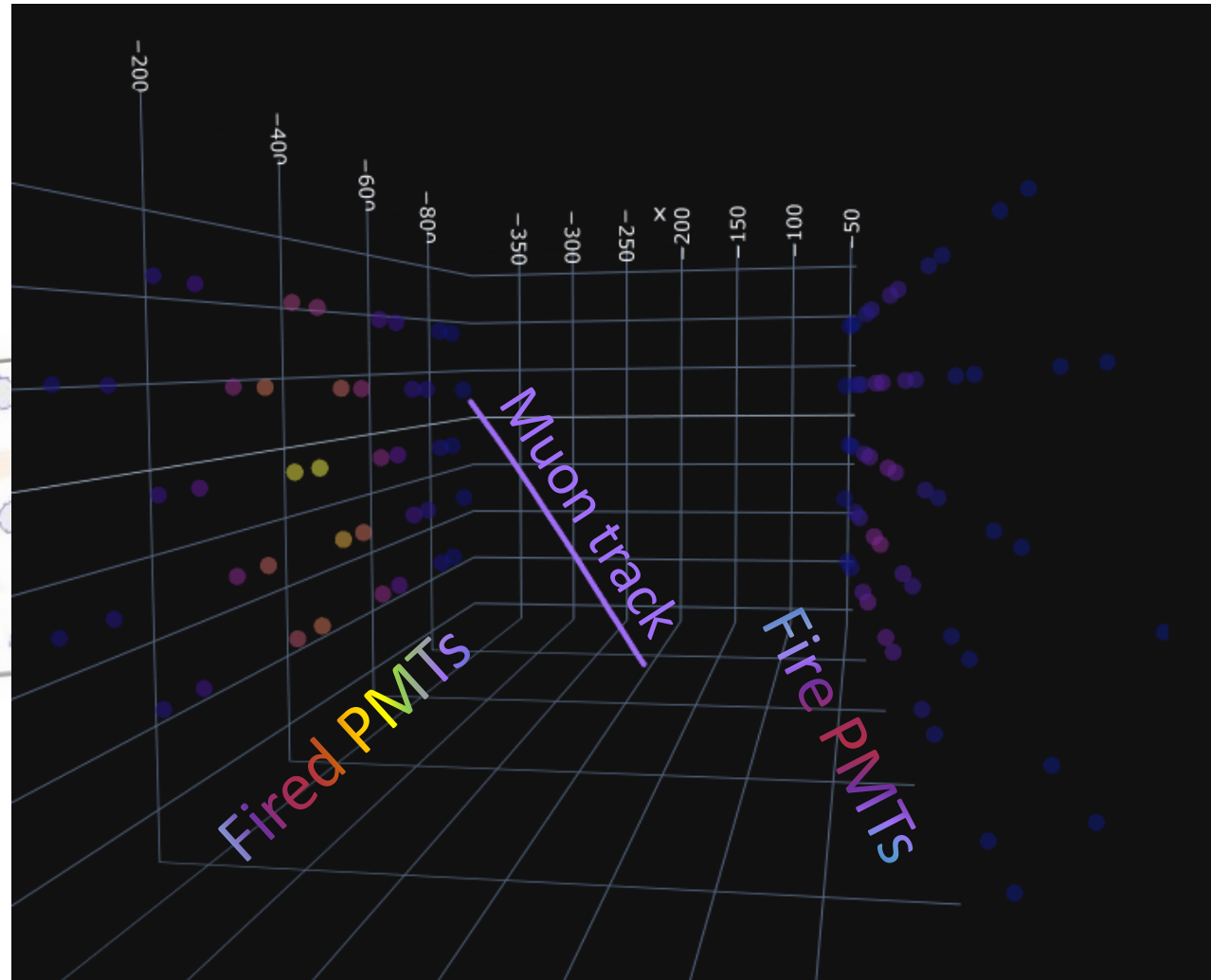
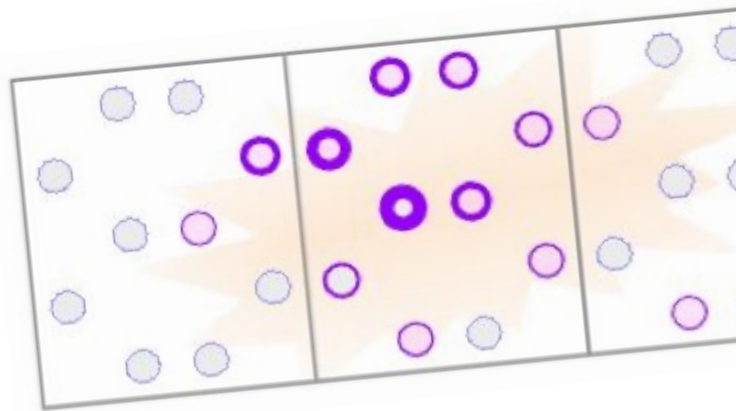
where: R = 50 Ohm → standard impedance of our circuit.

- Photoelectron number (**n.ph.e.**) = CHARGE /  $e (=1.6 \cdot 10^{-19} \text{ C})$ .

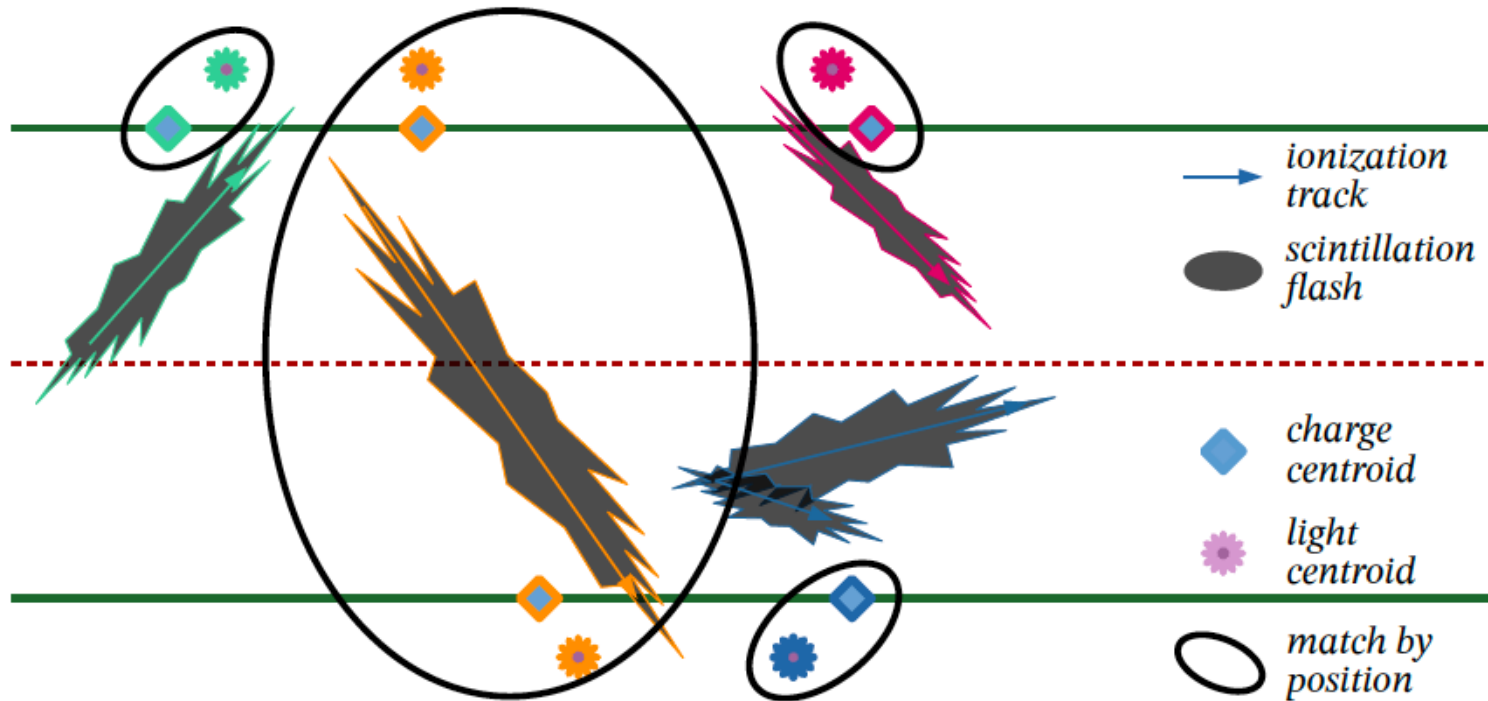
**n.ph.e.**

$$\begin{aligned} &= \frac{\text{CHARGE}}{e} \cdot \frac{1}{\text{GAIN}} \\ &= \frac{1}{e} \cdot \frac{1}{\text{GAIN}} \cdot FE \cdot \text{AREA} \\ &= \frac{0.122 \cdot 2 \cdot 10^{-12}}{50} \cdot \frac{1}{1.6 \cdot 10^{-19}} \cdot \frac{1}{\text{GAIN}} \cdot \text{AREA} \\ &= \left[ \frac{\text{V/ADC} \cdot \text{s}}{\Omega \cdot \text{C}} \cdot \text{ADC} \right] \end{aligned}$$

# Light signal study: Muon trajectory inside the TPC → light detected by PMTs



# Light signal study: Barycentre matching



1. **Reconstruction** of TPC tracks and PMT flashes;
2. **Compute the centroid** of each reconstructed **track** (within each TPC);
3. **Compute the centroids** of each reconstructed **flash** (within each TPC) at any time;
4. **Associate the two by proximity.**

# Study of light signal:

## Baricenter and width along beam axis

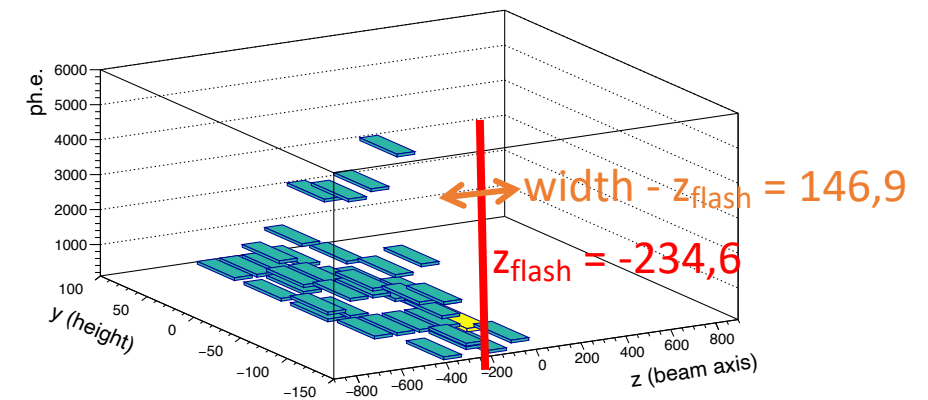
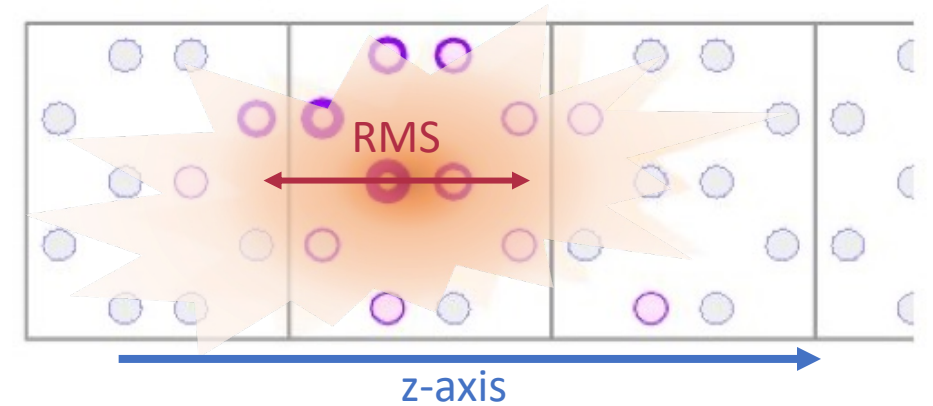
**Baricenter** along beam direction is the average of the fired PMT weighed with the quantity of reconstructed light in each fired PMT:

$$z_{flash} = \frac{\sum_{i = \text{fired PMT}} phe_i z_i}{\sum phe_i}$$

where  $z_i$  is the position of the fired PMT and  $w_i$  is the weight meaning the quantity of reconstructed light in each PMT.

Baricenter **width** along beam direction is the position **RMS of the fired PMTs, weighed with the quantity of reconstructed light in each fired PMT:**

$$width_{flash} = \sqrt{\frac{\sum_{i = \text{fired PMT}} phe_i (z_i - z_{flash})^2}{\sum phe_i}}$$



Example: run 9435, event 18513, flash\_id 0