# Dark Matter Induced Power in Quantum Devices

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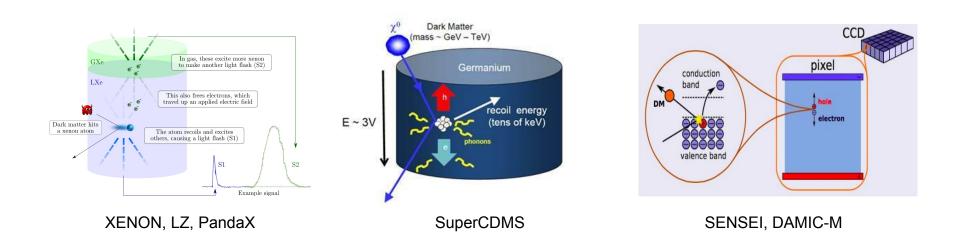
Seoul National University

arXiv:2210.09313 w/ R. Leane & N. Kurinsky

TAUP 2023 Vienna

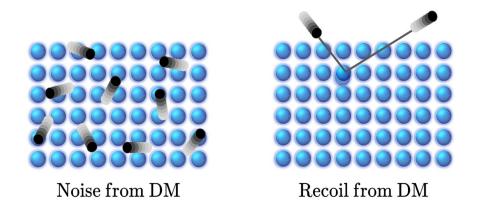


## **Conventional Direct Detection Experiments**



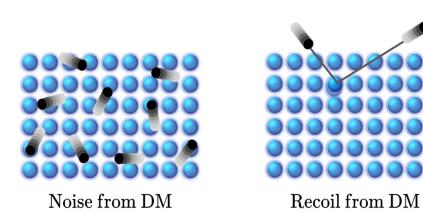
All looking for single scattering events from DM

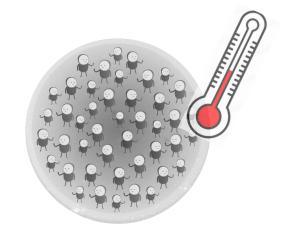
Instead of individual events, use power/heat deposited by DM



arXiv:2210.09313

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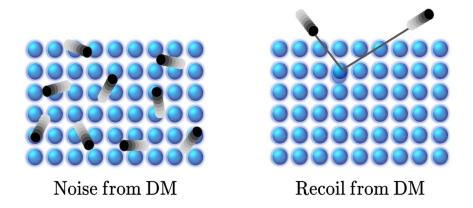




Using a very sensitive bolometers

arXiv:2210.09313

Instead of individual events, use power/heat deposited by DM

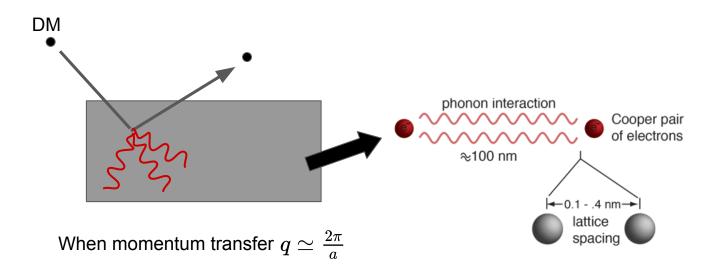


$$P_{
m DM} = \int d\omega \, \omega rac{d\Gamma}{d\omega}$$
  
Scattering rate

arXiv:2210.09313

## **Power Deposition Mechanism**

DM scattering creates phonons in this low mass regime: MeV - GeV

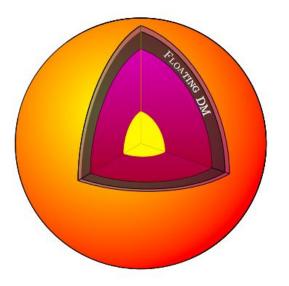


The excited phonons could directly be detected or

In a SC, the phonons will break Cooper pairs & release quasiparticles that can be detected

We already have devices that can measure very small power deposition in this way

#### **Captured Dark Matter on Earth**



arXiv:2012.03957, 2209.09834, 2303.01516

FIG. 1. Schematic of floating DM on the outer region of the celestial object as found in this work (dark shaded shell).

Over time, halo DM may get captured in the Earth and can get thermalized

## Conventional direct detection experiments do not have low enough threshold to probe thermalized DM

Halo DM

$$m_\chi=1~{
m GeV}$$

$$E_{
m kin}={\cal O}({
m keV})$$

$$v_{\scriptscriptstyle Y} = 230~{
m km}\,{
m s}^{-1}$$

Thermalized DM

$$m_\chi=1~{
m GeV}$$

$$E_{
m kin} = \mathcal{O}(10~{
m meV})$$

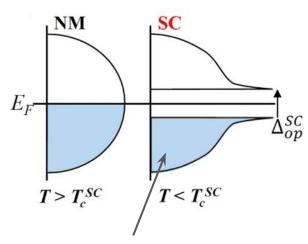
$$v_{\scriptscriptstyle Y} = 1~{
m km\,s^{-1}}$$

Low threshold experiments are ideal for them

## **Quantum Devices Based on Superconductor**

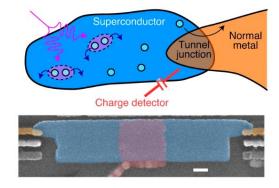
#### Low quasiparticle bkg device

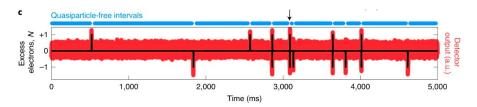
## Electron DoS



BCS ground state: sea of Cooper pairs



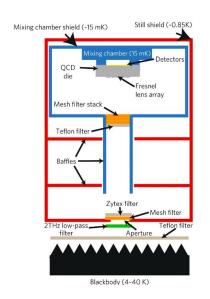


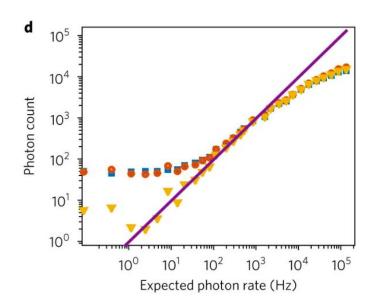


$$P = 6 \times 10^{-25} \text{ W}\mu\text{m}^{-3}$$

## **Quantum Devices Based on Superconductor**

Low noise bolometer, Cryogenic infrared sensor

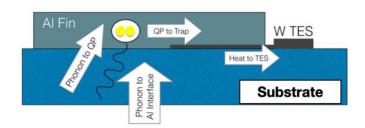




 $P = 1.7 \times 10^{-20} \text{ W}\mu\text{m}^{-3}$ 

## **Quantum Devices Based on Superconductor**

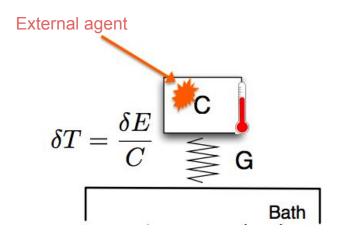
#### SuperCDMS Si detector covered with SC Al fins coupled to W TES



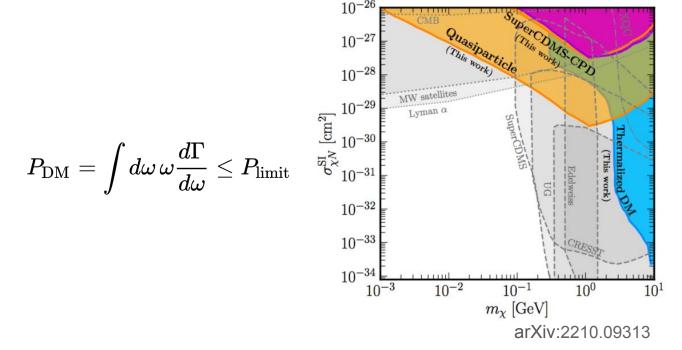
Athermal phonon sensor



$$P = 3 \times 10^{-21} \text{ W}\mu\text{m}^{-3}$$

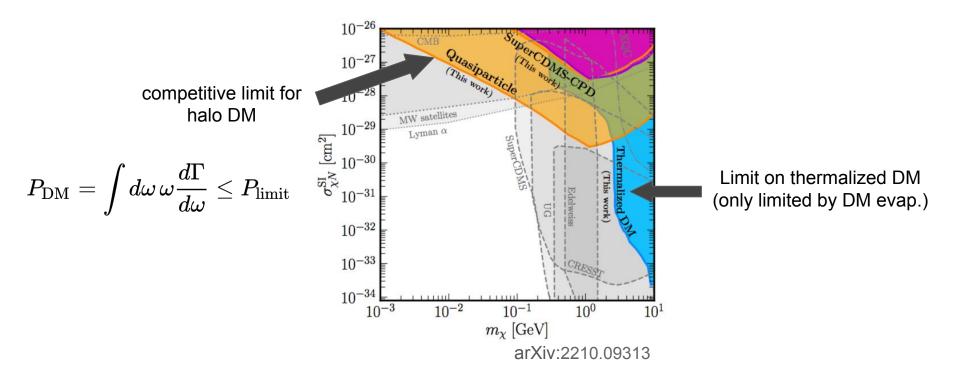


#### New Limits on DM-nuclear cross section



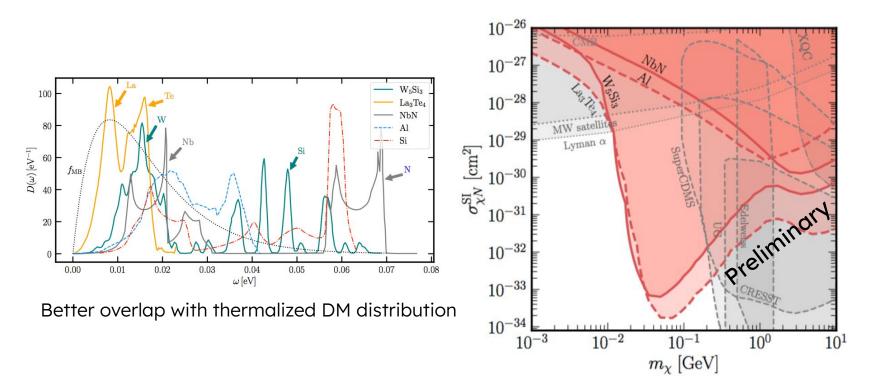
Unprecedented power sensitivity helps us put new limits on DM-nucleon cross section for both thermalized and halo population

#### New Limits on DM-nuclear cross section



Unprecedented power sensitivity helps us put new limits on DM-nucleon cross section for both thermalized and halo population

## Optimizing the absorber material

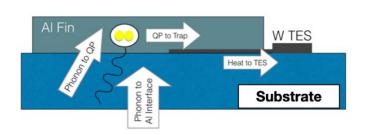


New materials with more phonon states at low energy

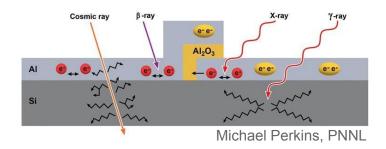
#### Challenges for detection:

- Power calibration of the calorimeter
- Neutron scattering
- ☐ Radioactivity & cosmic rays
- Unknown systematics

## Similarity w/ superconductor-based Qubit



Athermal phonon sensor used in SuperCDMS



SC Al-based Qubit chip

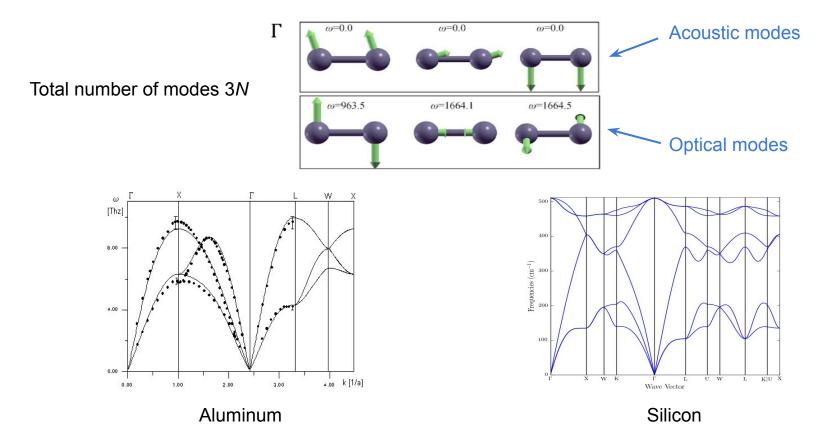
Technological similarity between Quantum sensors & Qubit chips

Cross-community collaboration will be critical

#### Final takeaways

- As DM searches dig deeper into the parameter space, it becomes more challenging
- New technique of measuring the power deposited by DM could be more useful in future
- Mesoscopic quantum devices are set to improve further in measuring small energy deposit
- Close technical connection with qubit development research
- Collaborative strategy will be beneficial

## Phonons are quanta of crystal lattice oscillation



#### **Dark Matter Interaction Rate**

Velocity-averaged interaction rate

$$\Gamma = rac{\pi \sigma_N n_\chi}{
ho_T \mu^2} \int d^3 v f_\chi(\mathbf{v}) \int rac{d^3 q}{(2\pi)^3} F_{
m med}^2(q) \underline{S(\mathbf{q},\omega_\mathbf{q})}$$

Structure Factor: detector response

$$rac{d\Gamma}{d\omega} = rac{\pi\sigma_N n_\chi}{
ho_T \mu^2} \int d^3v f_\chi(\mathbf{v}) \int rac{d^3q}{(2\pi)^3} F_{
m med}^2(q) S(\mathbf{q},\omega) \delta(\omega-\omega_\mathbf{q})$$

$$\omega_{f q}={f q}\cdot{f v}-rac{q^2}{2m_{
m cr}}=E_f-E_i$$
 Energy transfer in a single scattering

$$F_{
m med} = 1$$
  $F_{
m med} = (lpha m_e/q)^2$ 

Heavy mediator Light mediator

## Single & Multi-phonon Structure Factor

Coherent Approx. 
$$q < q_{
m BZ}$$

$$S_{
m LA} \simeq rac{2\pi}{V_c} rac{\sum A_d q^2}{2m_p \omega_{
m LA}} \delta(\omega - \omega_{
m LA})$$

Coherent Approx. 
$$q < q_{\rm BZ} \begin{cases} S_{\rm LA} \simeq \frac{2\pi}{V_c} \frac{\sum A_d q^2}{2m_p \omega_{\rm LA}} \delta(\omega - \omega_{\rm LA}) \\ S_{\rm LO} \simeq \frac{2\pi}{V_c} \frac{a^2 q^4}{32m_p \omega_{\rm LO}} \frac{A_1 A_2}{A_1 + A_2} \delta(\omega - \omega_{\rm LO}) \end{cases}$$

Knowledge from neutron

Incoherent Approx.

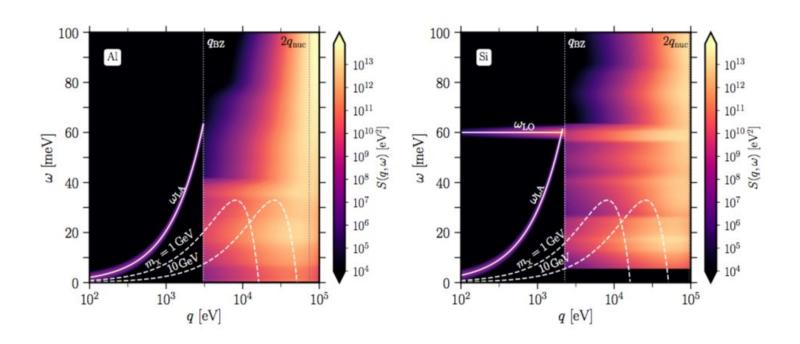
$$q>q_{
m BZ}$$

$$S_N \simeq rac{2\pi}{V_c} \sum_d f_d^2 e^{-2W_d(q)} \sum_n \left(rac{q^2}{2m_d}
ight)^2 rac{1}{n!} \Biggl(\prod_i \int d\omega_i rac{D_d(\omega_i)}{\omega_i} \Biggr) \, \delta(\omega - \sum_j \omega_j)$$



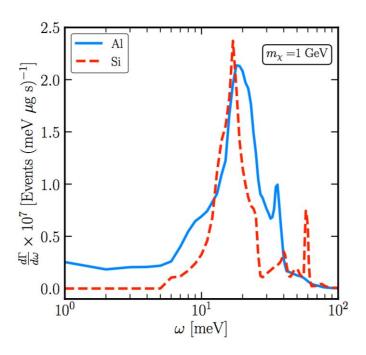
Approaches nuclear recoil in large  $\omega$  and q limit

#### Structure factors of AI & Si



Phonon structure factors  $S(q,\omega)$  in Al & Si are favorable for scattering with O(10 meV) energy DM

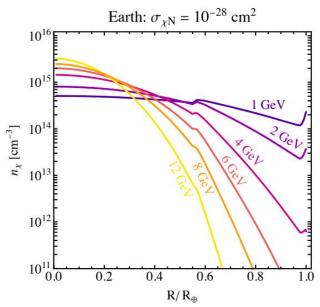
## Differential Scattering Rate in Al & Si



Phonon structure factors  $S(q,\omega)$  in Al & Si are favorable for scattering with O(10 meV) energy DM



## Floating Dark Matter on Earth



For DM mass 1-10 GeV and xsec > 10<sup>-35</sup> cm<sup>2</sup>, the thermalized population can get very dense near Earth's surface

However, these DM particles have very low energy,  $E_{\rm DM} \sim O(10 \text{ meV})$