

Axion Clumps meeting Neutron Stars

1710.08910: Visinelli, SB, Redondo, Freese, Wilczek

2107.07399: Millar, SB, Lawson, Marsh

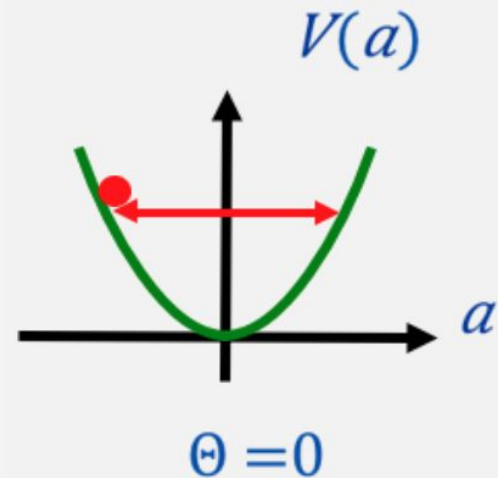
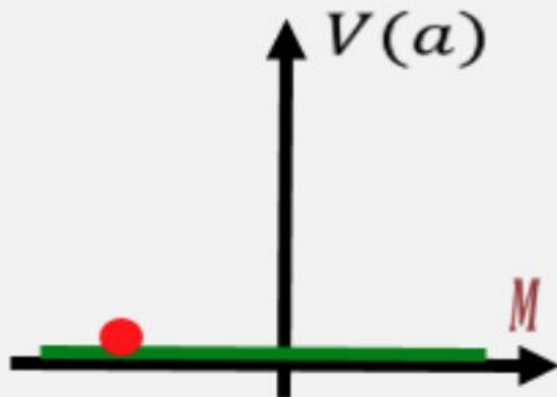
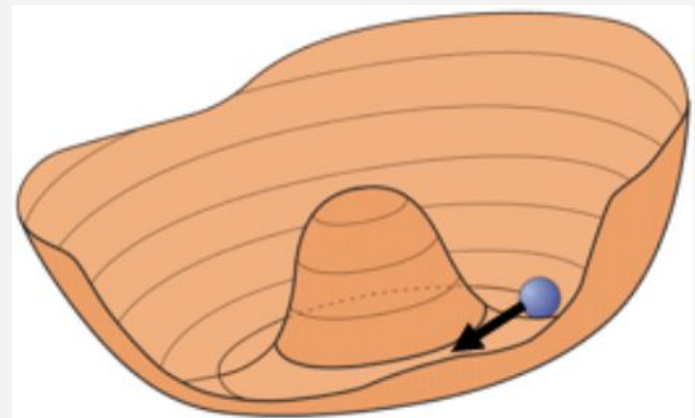
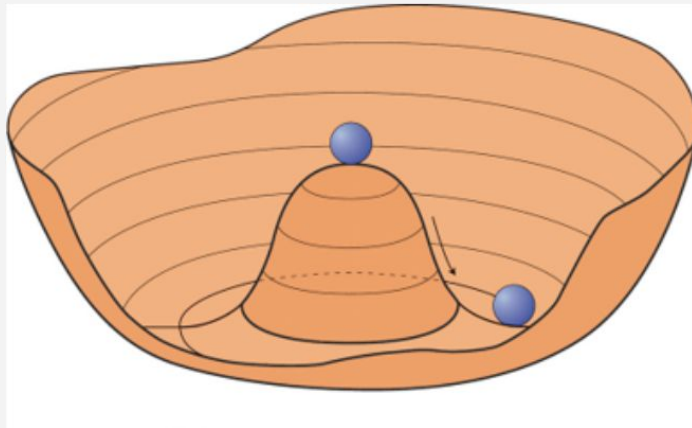
2212.08079: Witte, SB, Lawson, Marsh, Millar, Salinas

Sebastian Baum



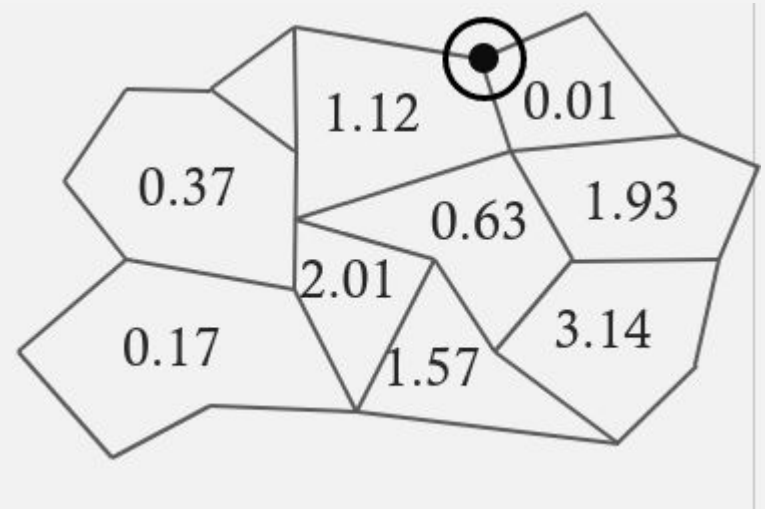
Axion Dark Matter Reminder: Re-alignment

- The axion field is "born" with random θ at the PQ phase transition
- At the QCD phase transition, the axion potential develops



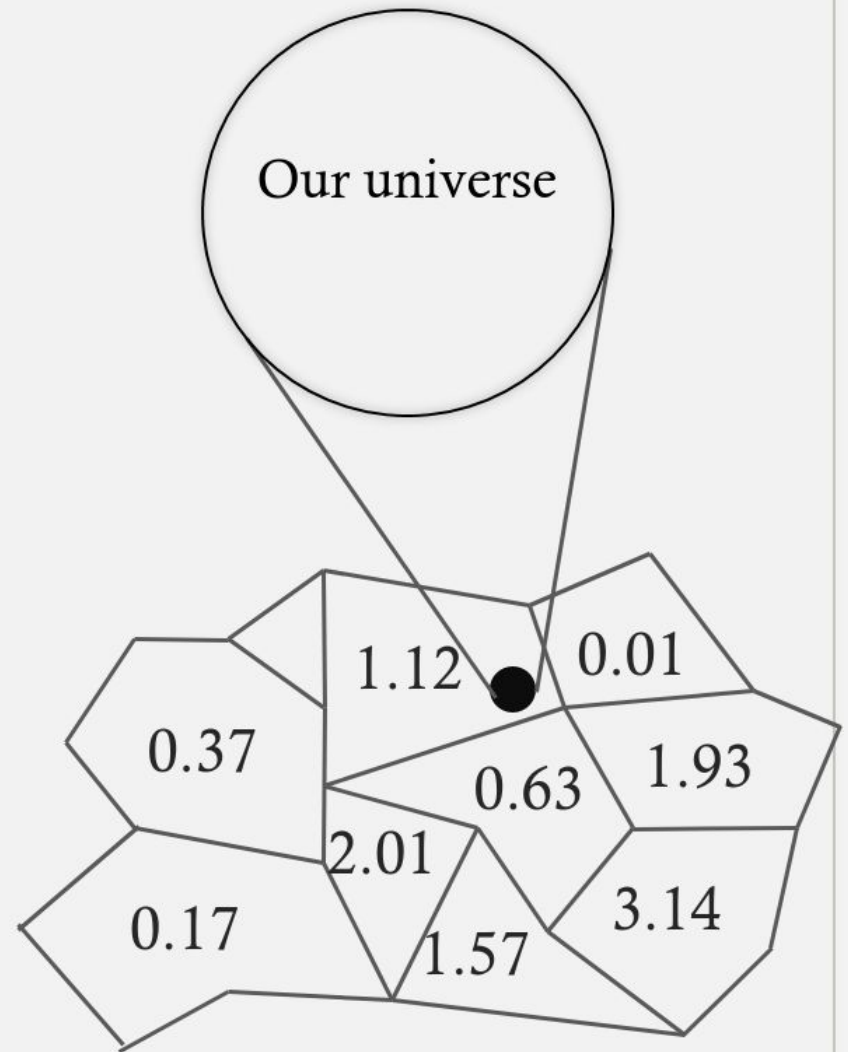
Axion DM: scenario 1

- PQ symmetry broken after the end of inflation
- Many patches with different misalignment angles in our Universe
- Topological defects (strings, domain walls, ...) exist in the early Universe

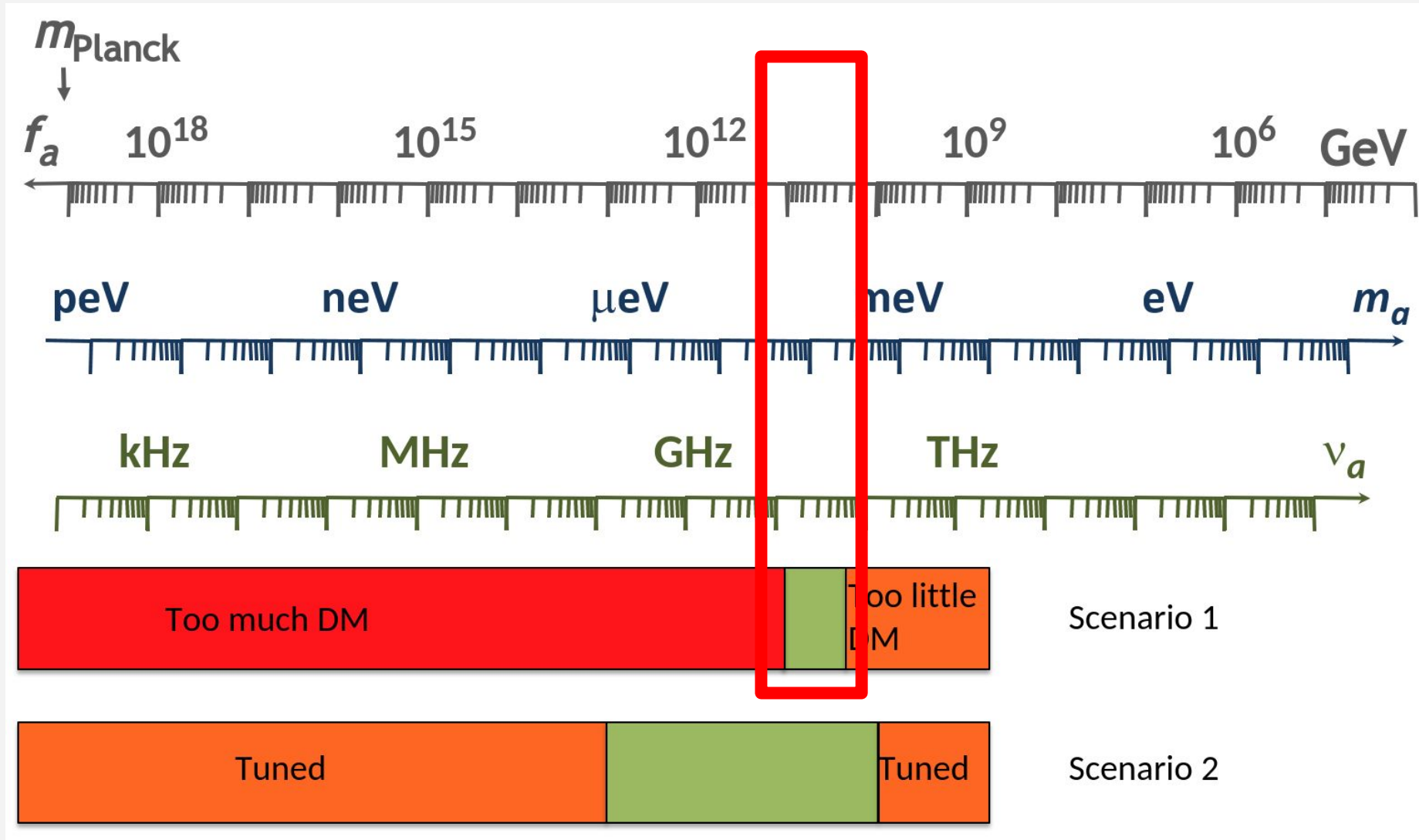


Axion DM: scenario 2

- PQ symmetry broken
(sufficiently early) before (the end of) inflation
- One misalignment angle in the whole observable Universe
- No topological defects



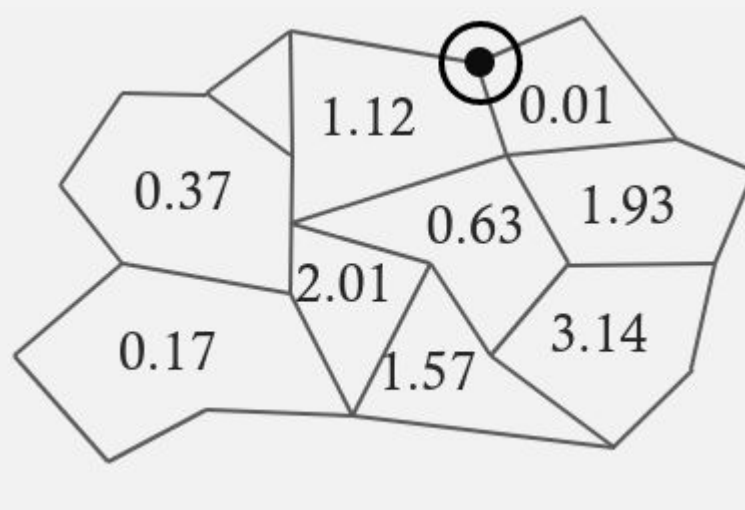
Axion Dark Matter



How much axion DM could be in compact objects?

In scenario 1 (PQ broken after inflation) the axion field in our Universe is “born” with $O(1)$ inhomogeneities

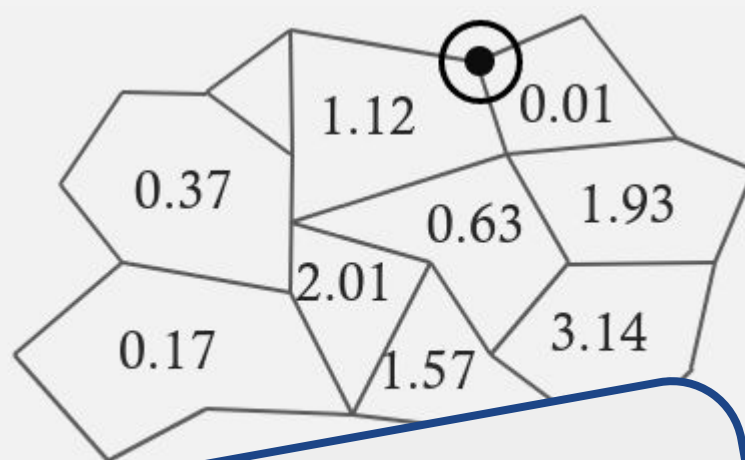
- Topological defects decay when the axion mass turns on
- “Axitons” decay quickly after the QCD phase transition finishes
- Axion overdensities become gravitationally bound and form “minicluster” around matter-radiation equality



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- Axion overdensities become gravitationally bound and form “minicluster” around matter-radiation equality



Most of axion DM could be in compact objects!

Levkov, Panin, Tkachev 1804.05857, Vaquero, Redondo, Stadler 1809.09241, Buschmann, Foster, Safdi 1906.00967, Eggemeier & Niemeyer 1906.01348, Eggemeier, Redondo, Dolag, Niemeyer, Vaquero 1911.09417, Eggemeier, O'Hare, Pierobon, Redondo, Wong 2212.00560, ...

Axion Miniclusters, Stars, Axitons, ...

Bosonic Dark Matter can collapse to high densities (occupation numbers) and attain (hydrostatic) equilibrium

Relevant contributions:

- Kinetic Pressure (repulsive)
- Gravity (attractive)
- Self-interactions (both)

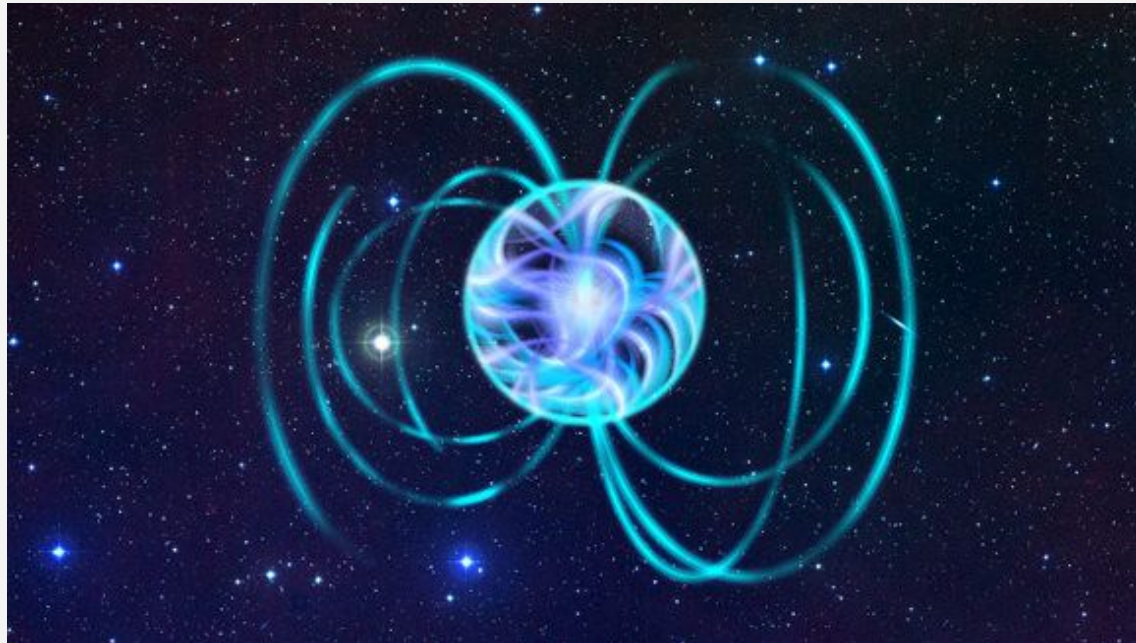
Summary I

In axion DM scenario 1 (PQ breaking after the end of inflation),
axion miniclusters and axion stars are ubiquitous

- Possible up to ~75% of axions bound in miniclusters
- At the center of miniclusters, axion stars form (quickly and efficiently!)
- Important open questions:
 - Evolution in hierarchical structure formation
 - Tidal disruption in Galaxies [Dokuchaev, Eroshenko, Tkachev \[1710.09586\]](#)
[Kavanagh, Edwards, Visinelli, Weniger \[2011.05377\]](#)
 - Dynamical evolution of axion stars [Fox, Weiner, Xiao \[2302.00685\]](#)
- Typical mass, size, density:
 - Axion minicluster: $M \sim 10^{-12} M_{\odot}$, $R \sim 10^9 \text{ km}$, $\rho \sim 10^3 \text{ GeV/cm}^3$
 - Axion star: $M \sim 10^{-13} M_{\odot}$, $R \sim 10^4 \text{ km}$, $\rho \sim 10^{17} \text{ GeV/cm}^3$

Enter the Neutron Star

- Radius: ~ 10 km
- Mass: $\sim 1.4 M_{\odot}$
- Magnetic field strength: $\sim 10^{14}$ G = 10^{10} T
- Plasmas mass of photons $\lesssim 50$ μ eV



(Resonant) axion photon conversion

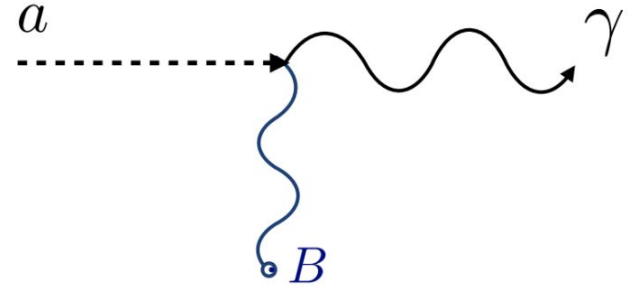
$$\mathcal{L} \sim g_{a\gamma\gamma} a E \cdot B$$



$$p_{a \rightarrow \gamma} \sim g_{a\gamma\gamma}^2 B^2 L^2$$

B: Magnetic Field

L: Length scale



(Resonant) axion photon conversion

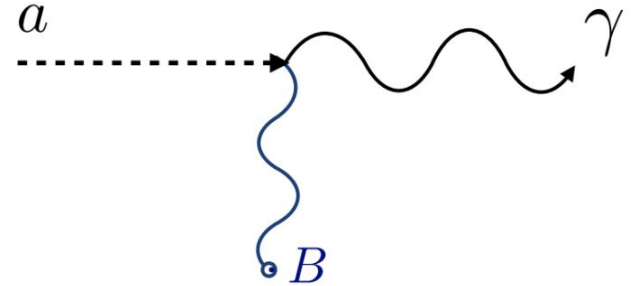
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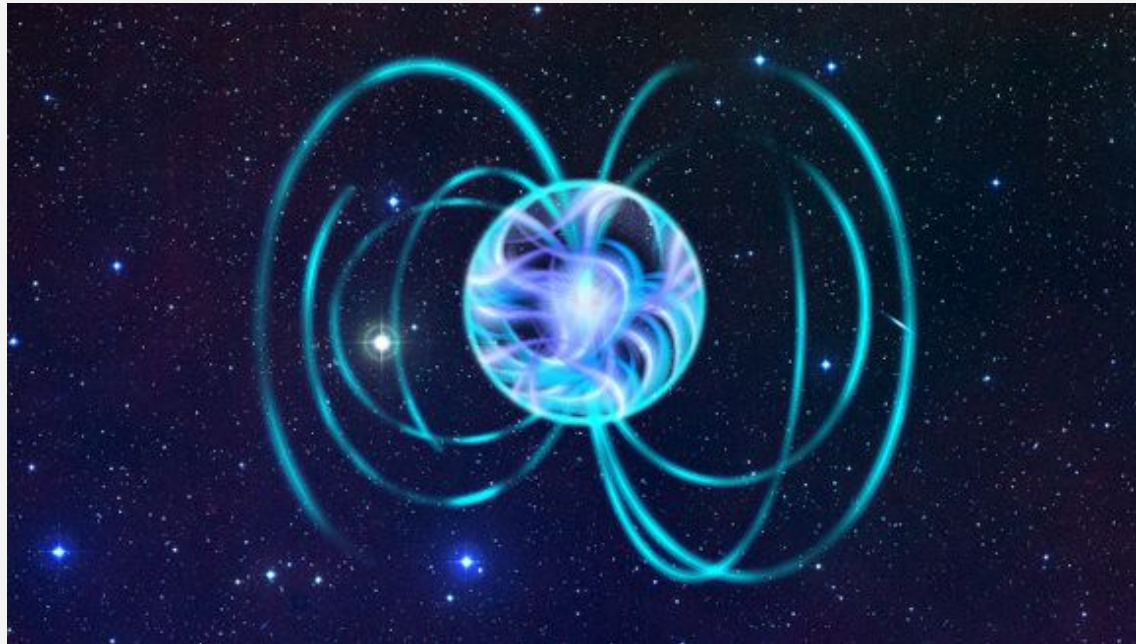


Two options for efficient axion-photon conversion of axion in astrophysical systems:

- Magnetic field with power at wavelength \sim axion mass
- Effective plasma mass of photons matching the axion mass

Enter the Neutron Star

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Axions meeting Neutron Stars

Pshirkov & Popov [0711.1264]

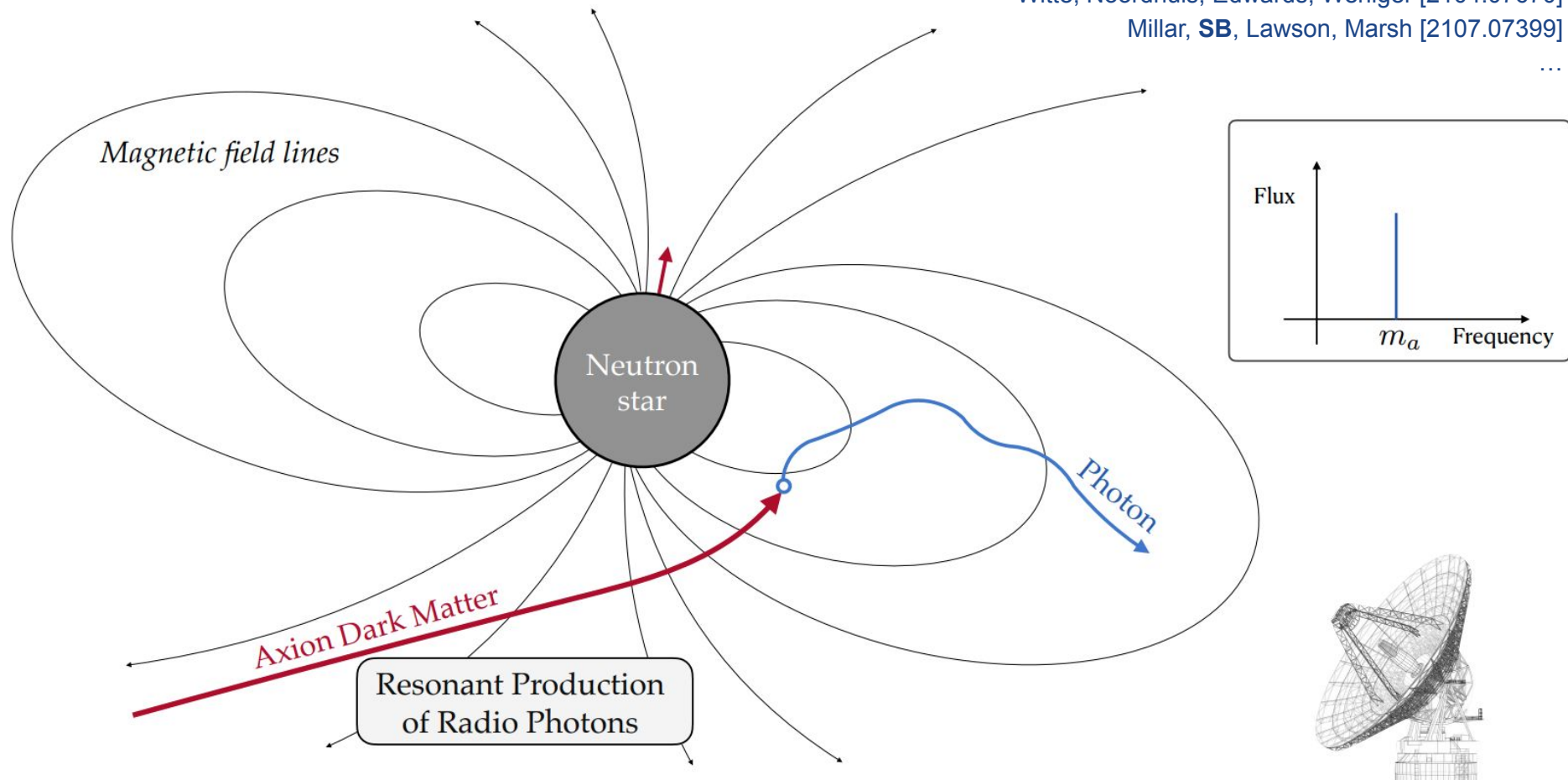
Hook, Kahn, Safdi, Sun [1804.03145]

...

Witte, Noordhuis, Edwards, Weniger [2104.07670]

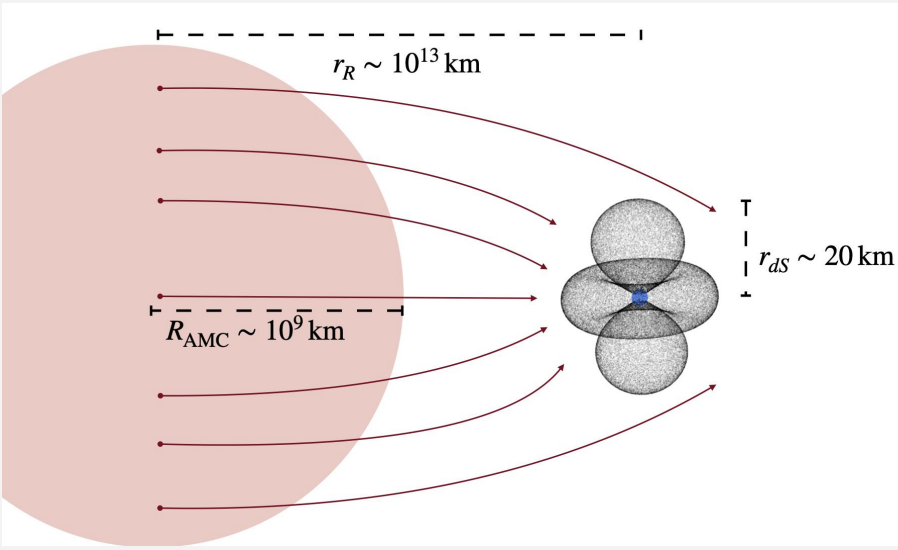
Millar, **SB**, Lawson, Marsh [2107.07399]

...



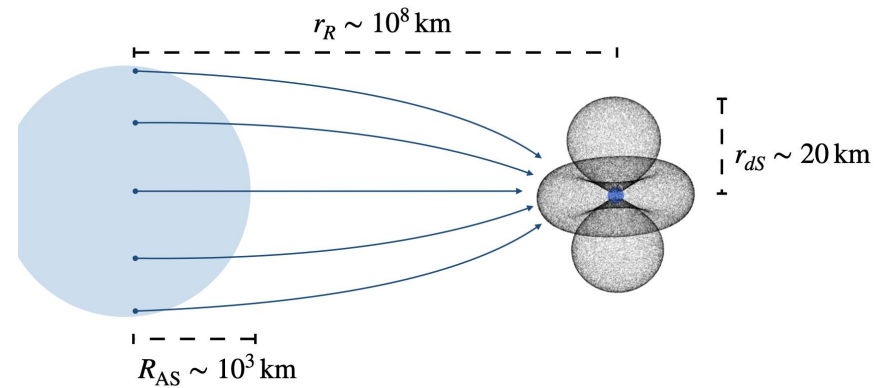
Axion Clumps & Neutron Stars

Axion Minicluster



Axion Star

Witte, **SB**, Lawson, Marsh, Millar, Salinas [2212.08079]



Axion Clumps & Neutron Stars

Three steps of the calculation:

- Gravitational infall of axions DM (see also [Bai, Du, Hamada \[2109.01222\]](#))
 - Treat axion clump as undisturbed until the Roche radius [$r_R = r_{AC} \times (2 M_{NS}/M_{AC})^{1/3}$]
 - Treat axions as free particles in the NS potential insides of the Roche radius
- Axion-photon conversion
 - Use 3D resonant conversion calculation from [Millar, **SB**, Lawson, Marsh \[2107.07399\]](#)
- Tracing the “photons” out of the plasma
 - Ray-tracing calculation à la [Witte, Noordhuis, Edwards, Weniger \[2104.07670\]](#)

Evaluate the whole chain as a hybrid importance/weighted sampling Monte Carlo integral with trajectories drawn homogeneously over the resonant conversion surface/drawn from the axion velocity distribution at the Roche radius

Axion Clumps & Neutron Stars

Result of the calculation in [Witte, SB, Lawson, Marsh, Millar, Salinas \[2212.08079\]](#):

For an axion clump – neutron star encounter, characterized by

- neutron star configuration (M_{NS} , magnetic field, plasma),
- axion clump (M_{AC} , R_{AC} , density profile, velocity distribution),
- encounter geometry (v_{AC} , b_{AC} , θ_v),

we compute, as functions of direction and time:

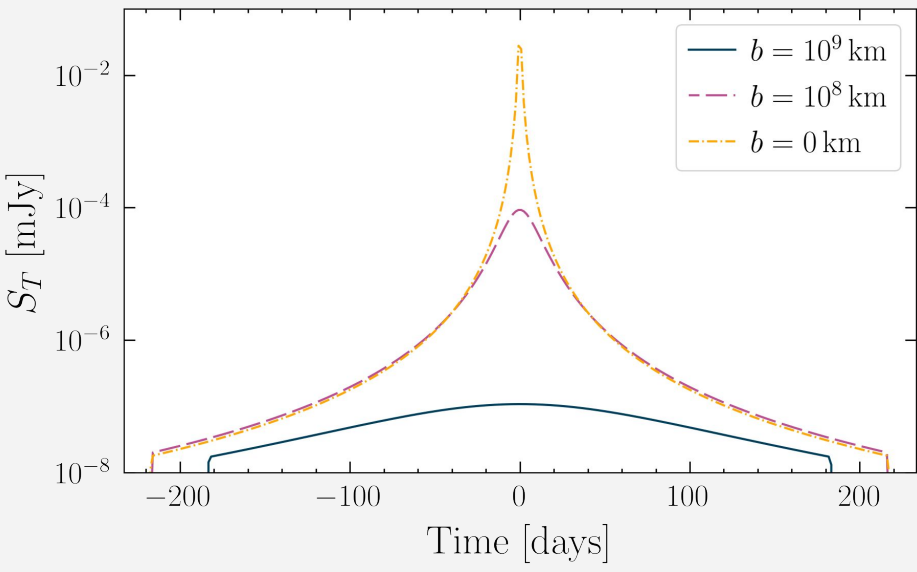
- Flux density at a telescope
- Characteristic linewidth

$$S_T(\theta, \phi) = \frac{dP}{d\Omega}(\theta, \phi) \frac{1}{\delta f_T \times d_T^2}$$

$$f_\sigma = \left[\frac{\sum_{j \in \text{pixel}} (E_j - \bar{E})^2 R_{a \rightarrow \gamma, j}}{\bar{E}^2 \sum_{j \in \text{pixel}} R_{a \rightarrow \gamma, j}} \right]^{1/2}$$

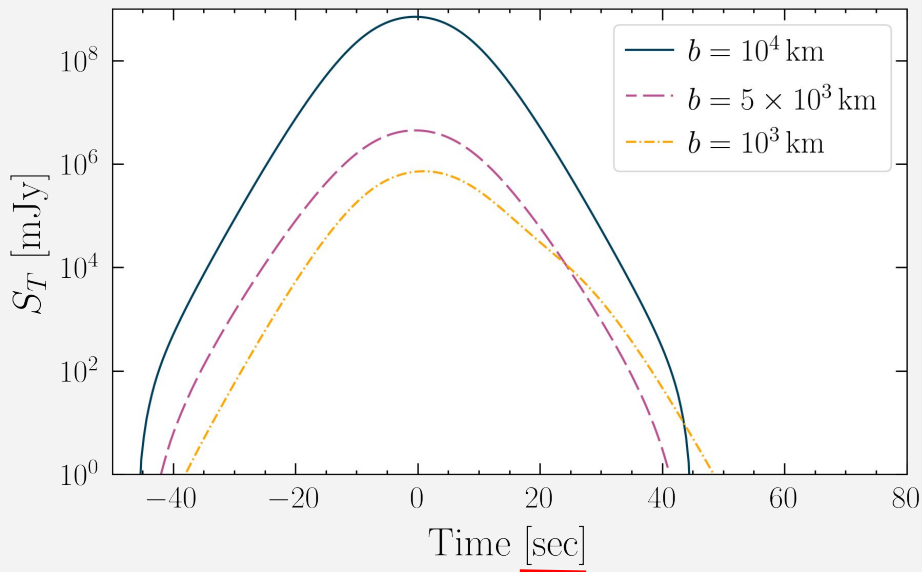
How long does the signal last?

Axion Minicluster



Axion Star

Witte, **SB**, Lawson, Marsh, Millar, Salinas [2212.08079]



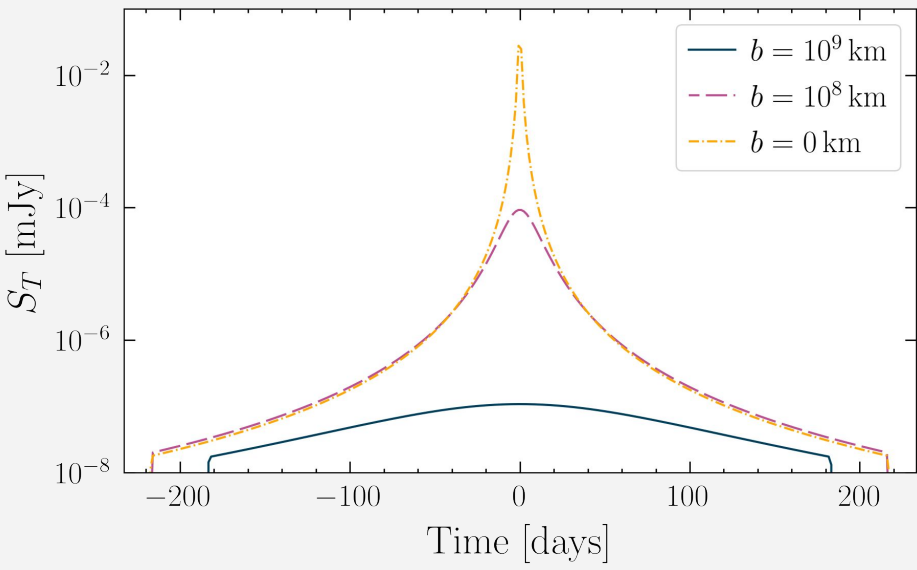
m_a	$26 \mu\text{eV}$
$g_{a\gamma\gamma}$	$10^{-14} \text{ GeV}^{-1}$
B_0	$1.6 \times 10^{14} \text{ G}$
P	3.76 s
θ_m	0.2 rad
θ_v	0.0 rad
δf_T	$10^{-5} \times m_a$
d_T	1 kpc

M_{AMC}	$10^{-12} M_\odot$
R_{AMC}	$1.86 \times 10^9 \text{ km}$
M_{AS}	$10^{-13} M_\odot$
R_{AS}	3905 km
ρ_{AMC}	NFW
b_{AMC}	$10^8 \hat{x} \text{ km}$
b_{AS}	$2 \times 10^3 \hat{y} \text{ km}$
$ v_{\text{AMC/AS}} $	100 km/s

How long does the signal last?

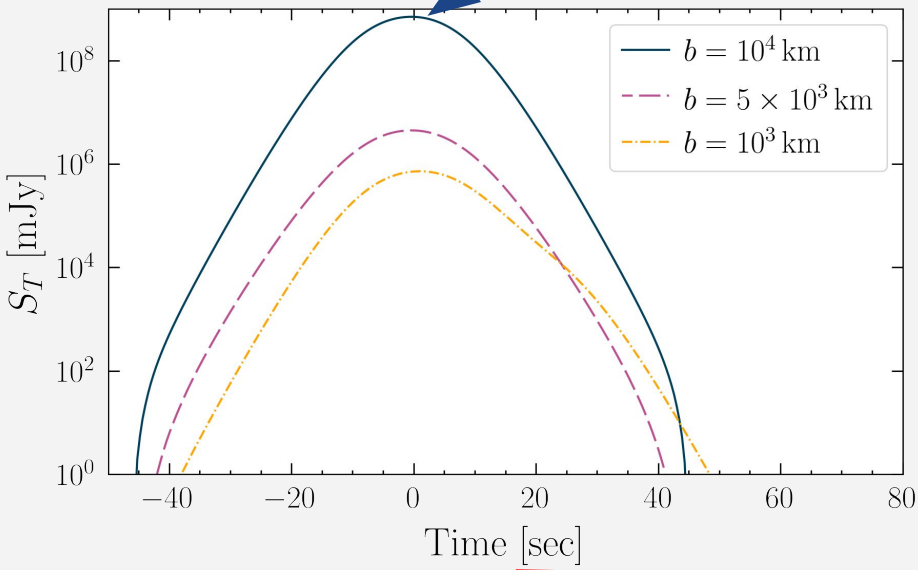
Still 10 mJy at 10 Mpc!

Axion Minicluster



Axion Star

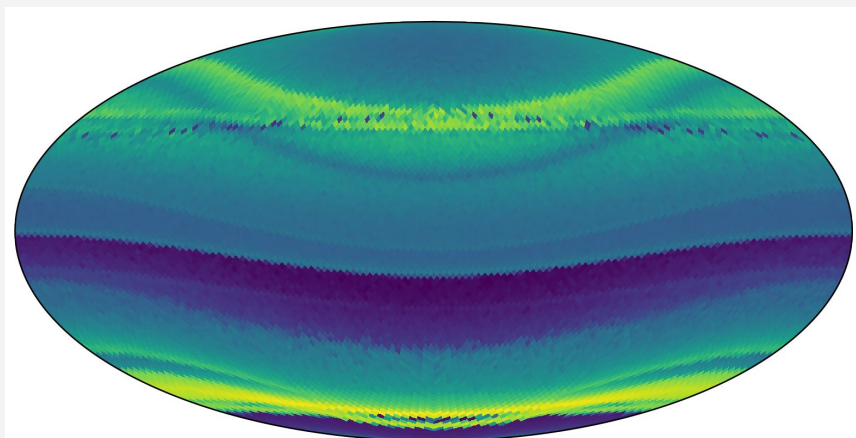
Witte, **SB**, Lawson, Marsh, Millar, Salinas [2212.08079]



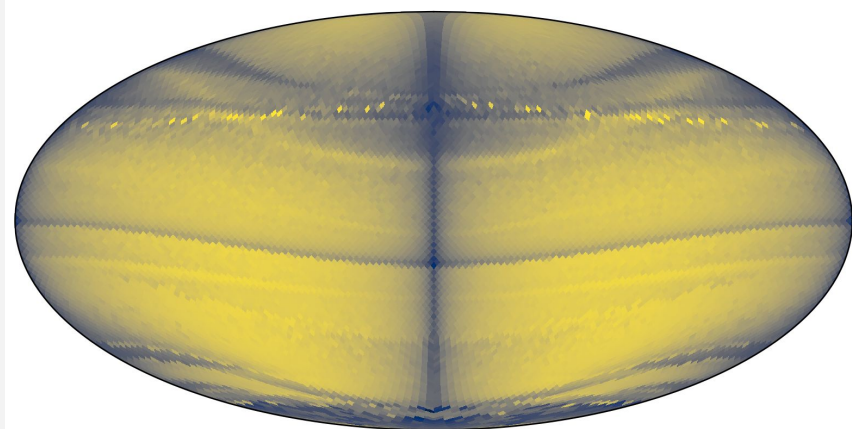
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Minicluster - Neutron Star Collision



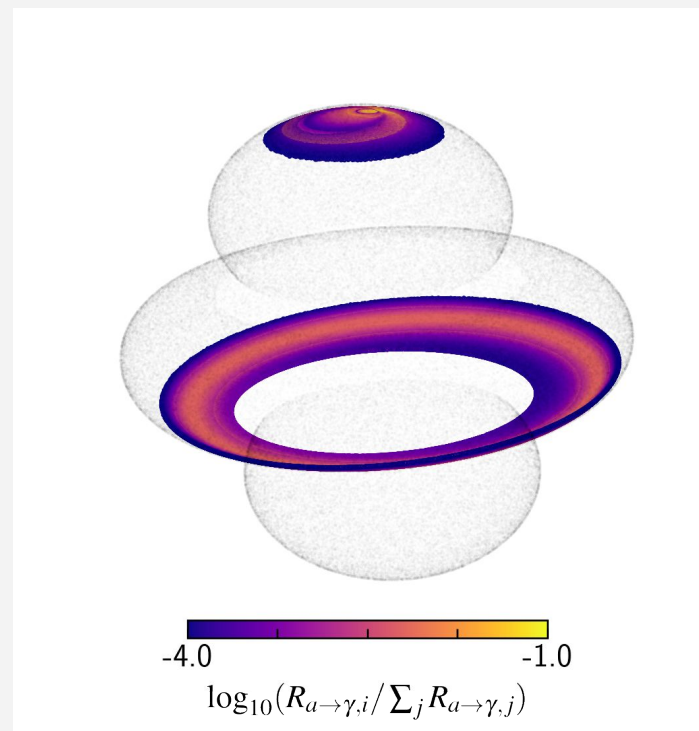
-7.18 $\log_{10}(S_T [\text{mJy}])$ -2.21



-8.5 $\log_{10}(f_\sigma)$ -5.5

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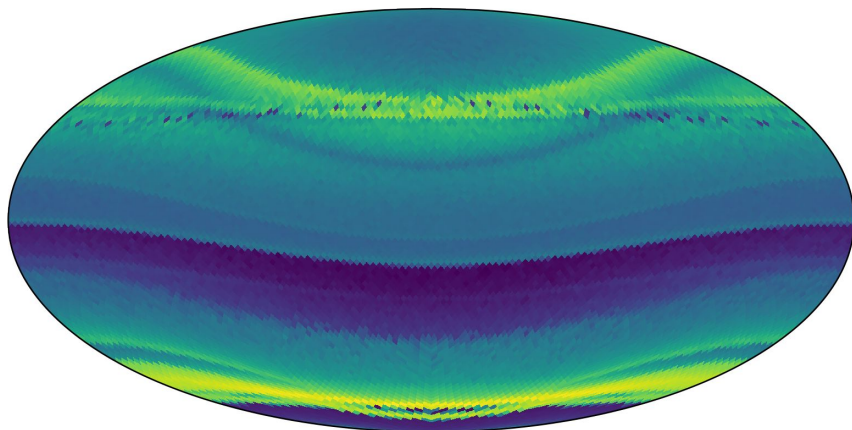
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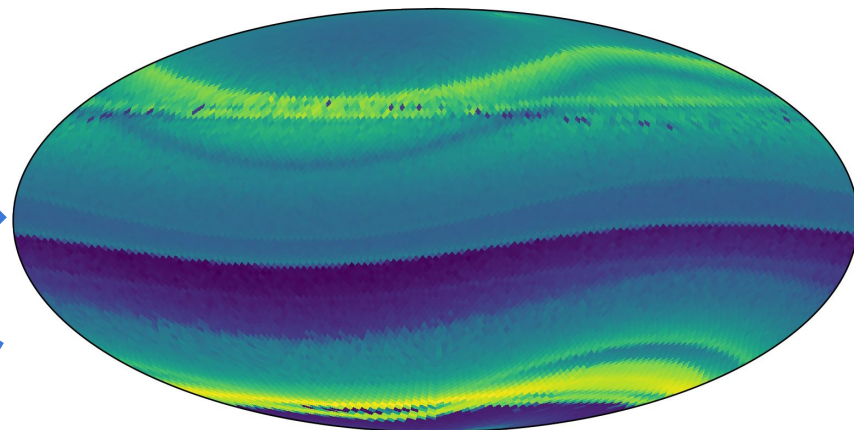
-4.0 $\log_{10}(R_{a \rightarrow \gamma, i} / \sum_j R_{a \rightarrow \gamma, j})$ -1.0

Minicluster - Neutron Star Collision: NS rotation

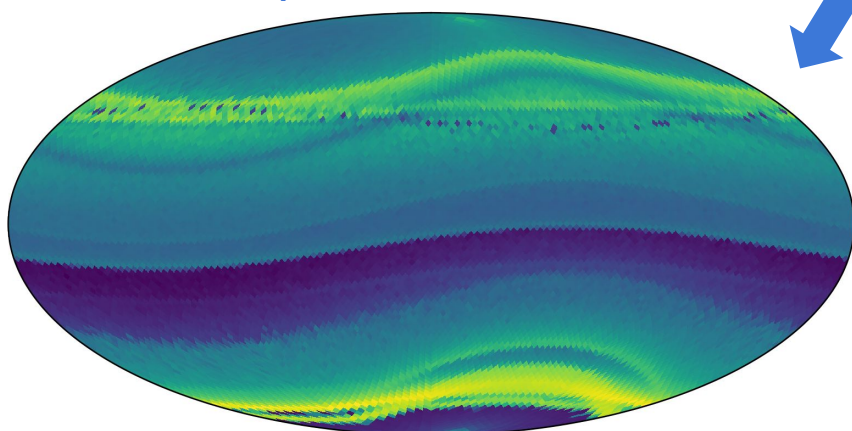
$$\psi = \Omega \times t = 0^\circ$$



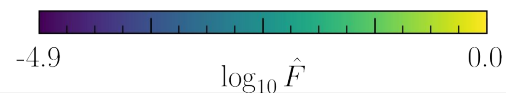
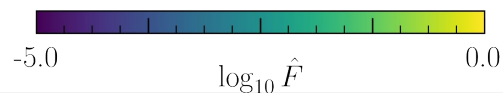
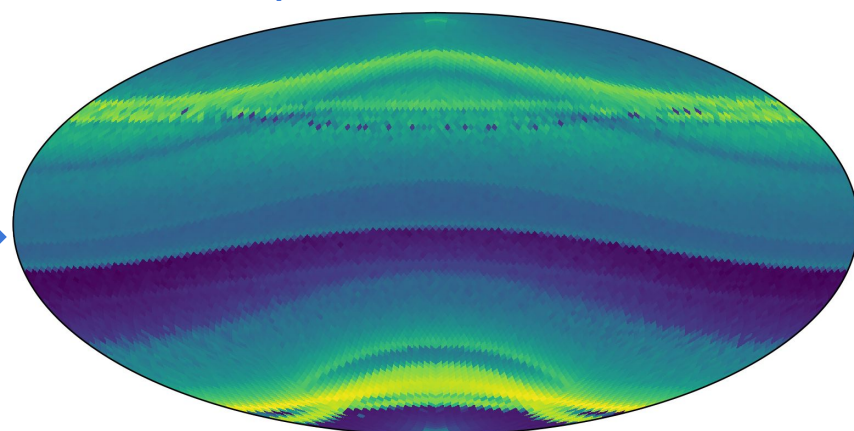
$$\psi = \Omega \times t = 36^\circ$$



$$\psi = \Omega \times t = 72^\circ$$

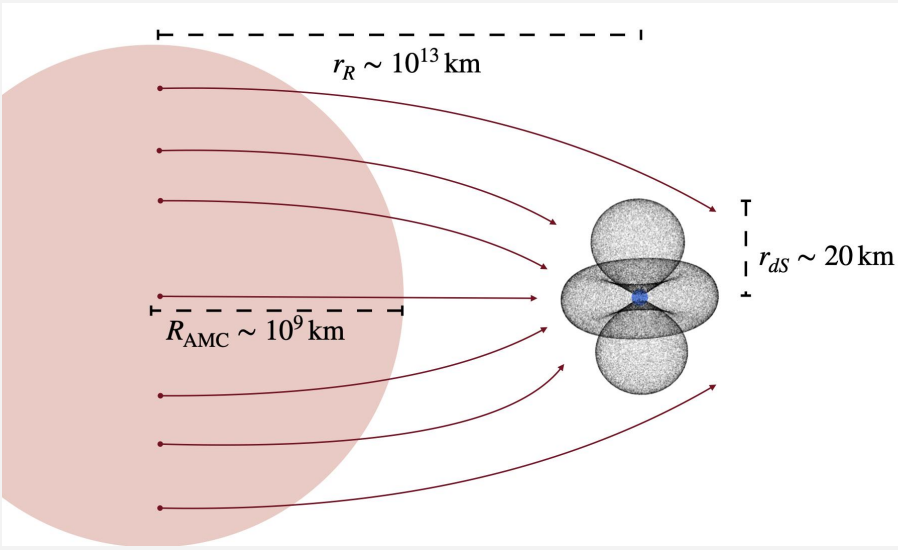


$$\psi = \Omega \times t = 108^\circ$$



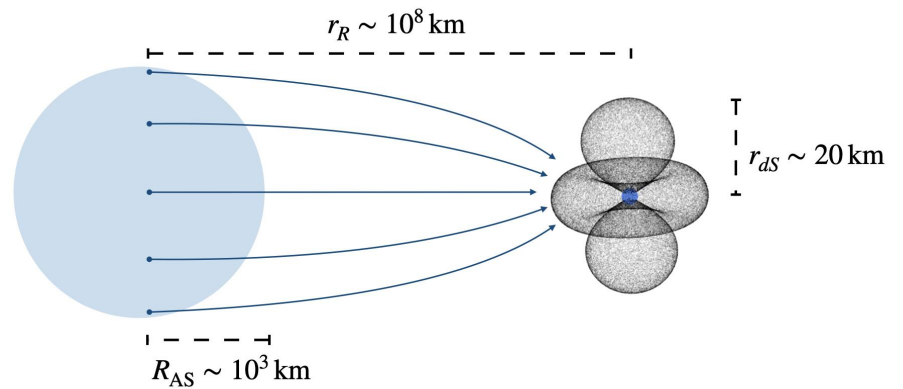
Axion Clumps & Neutron Stars

Axion Minicluster

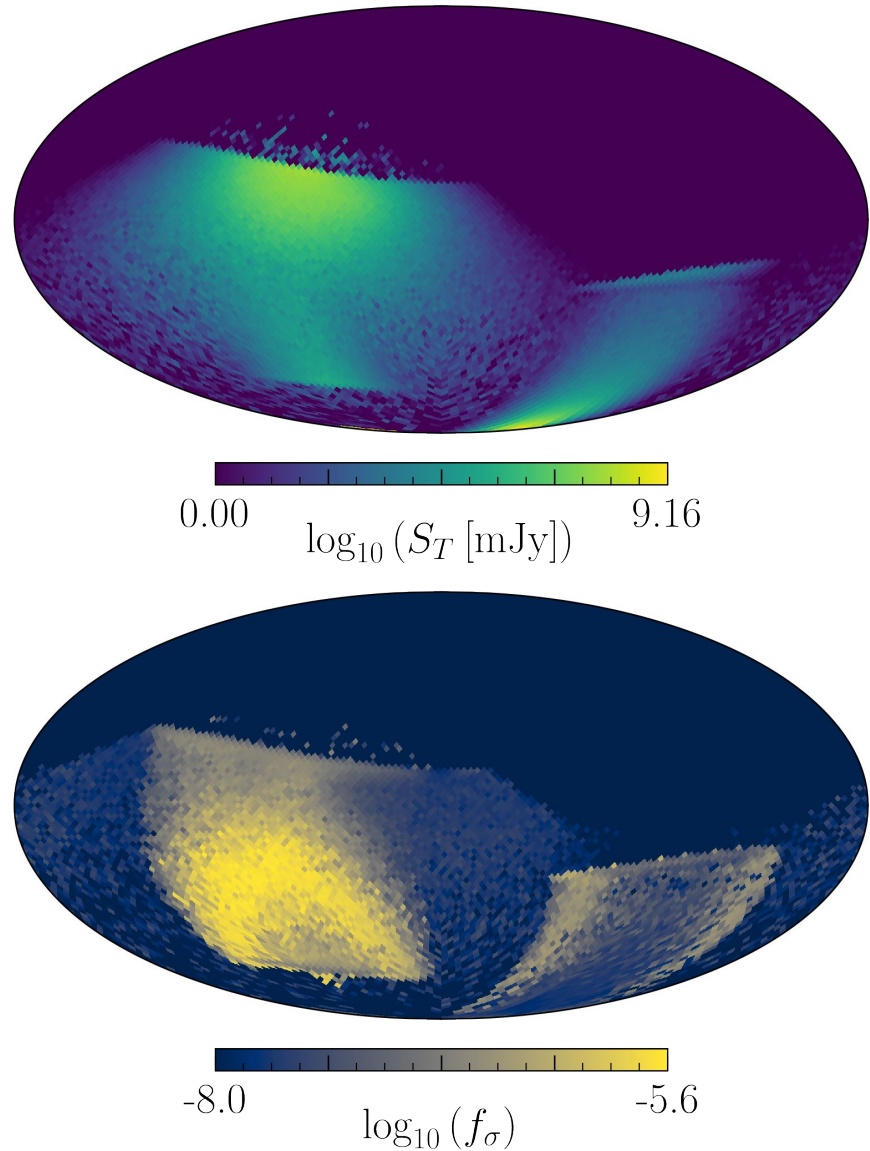


Axion Star

Witte, **SB**, Lawson, Marsh, Millar, Salinas [2212.08079]



Axion Star - Neutron Star Collision

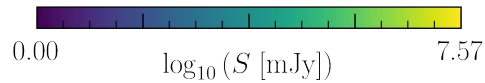
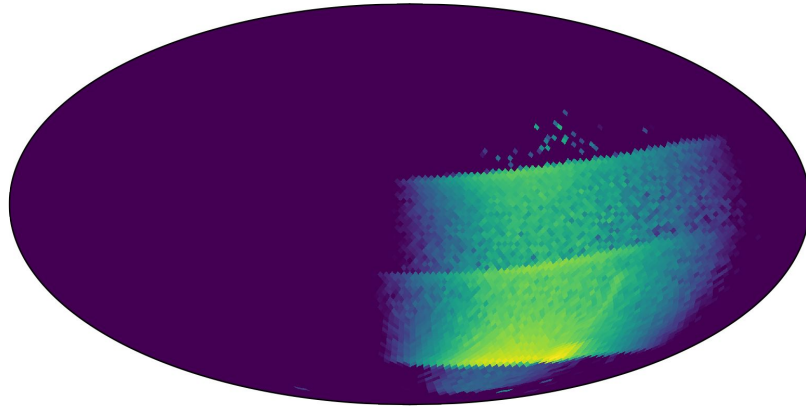


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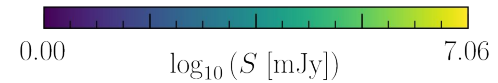
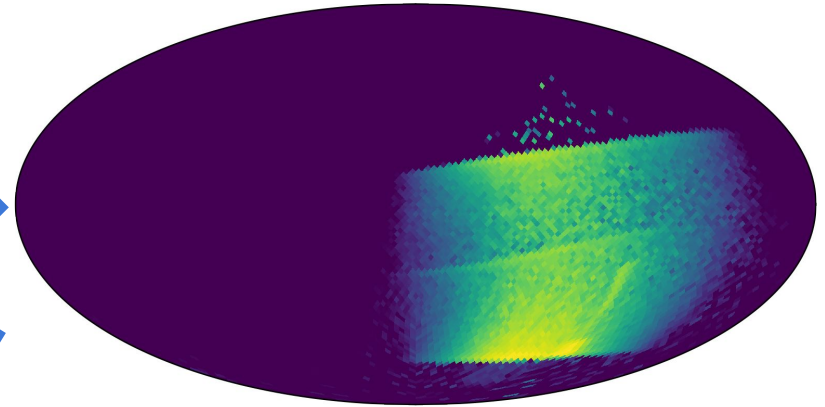
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Axion Star - Neutron Star Collision: NS rotation

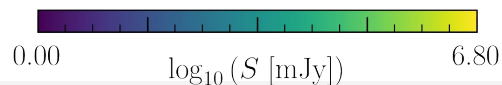
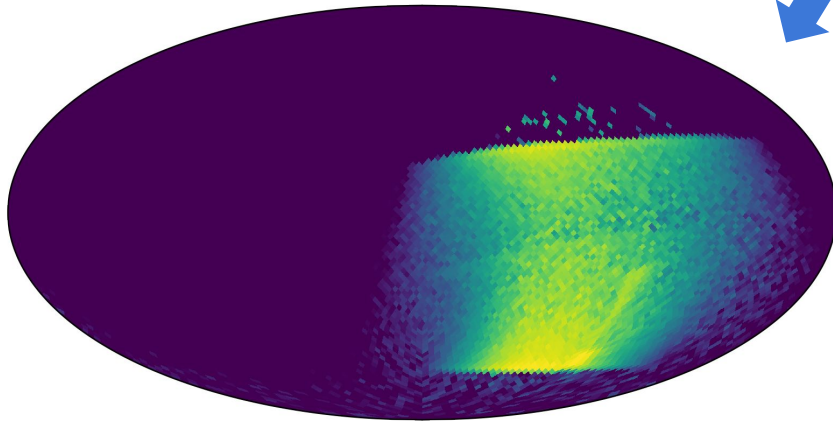
$$\psi = \Omega \times t = 0^\circ$$



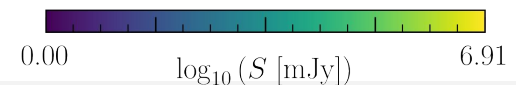
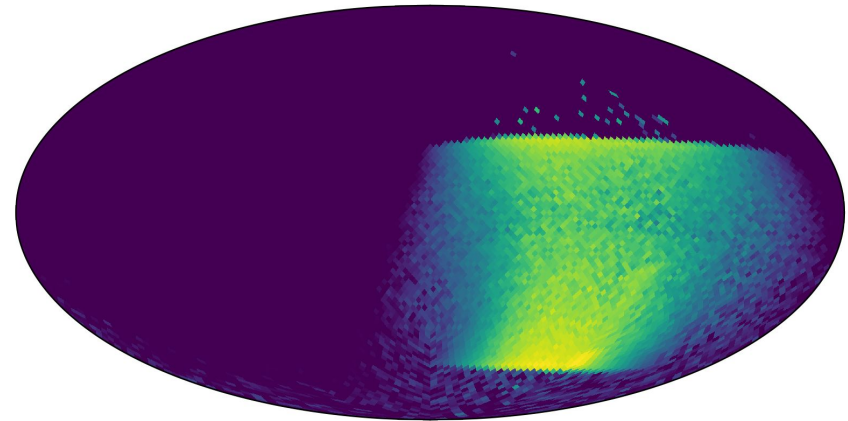
$$\psi = \Omega \times t = 22^\circ$$



$$\psi = \Omega \times t = 43^\circ$$

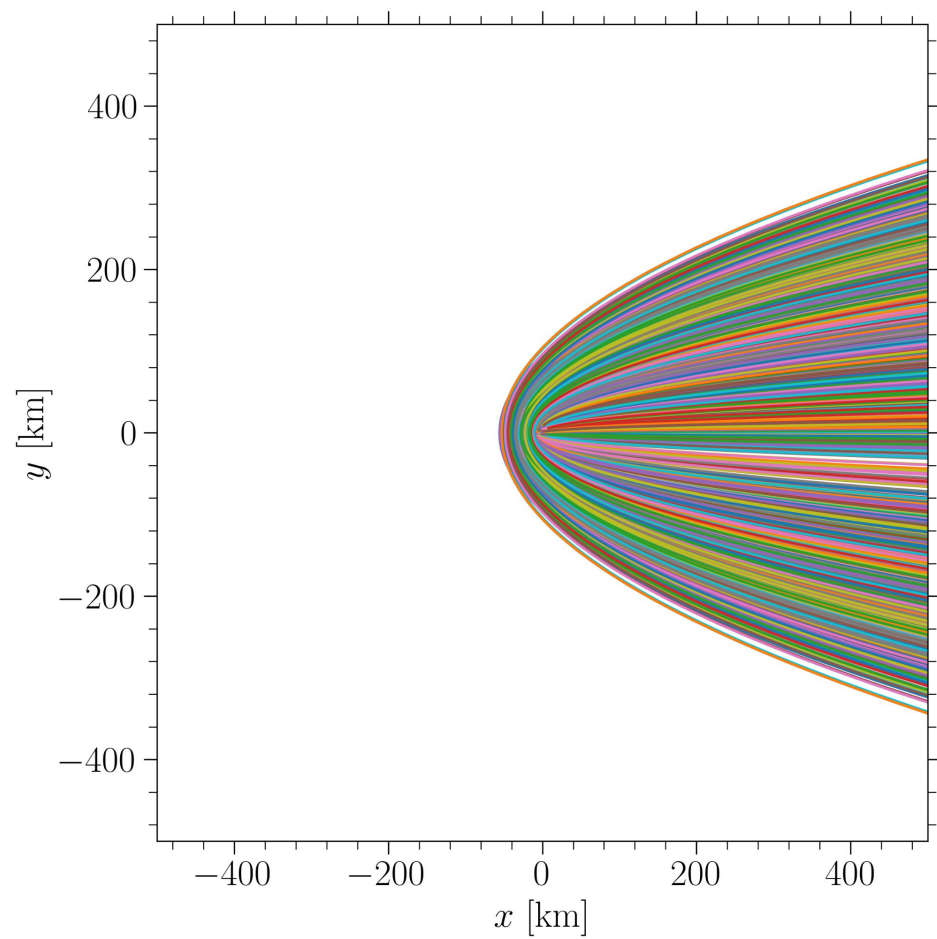
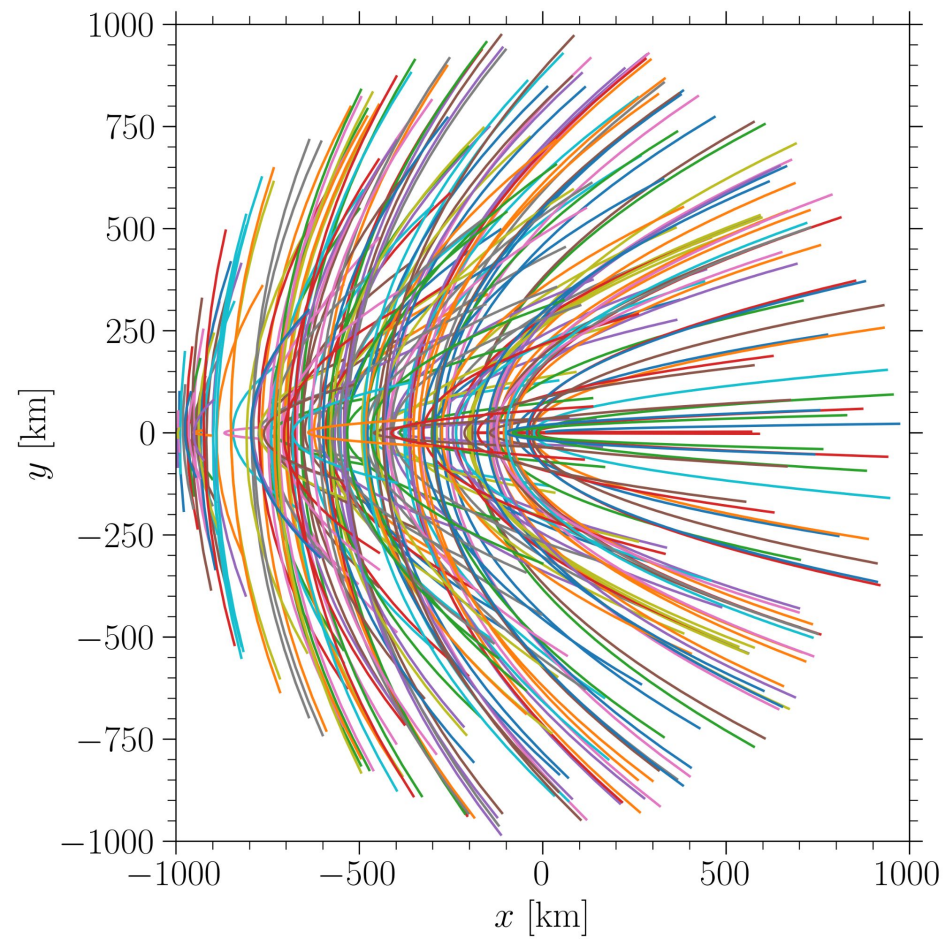


$$\psi = \Omega \times t = 65^\circ$$



Summary II

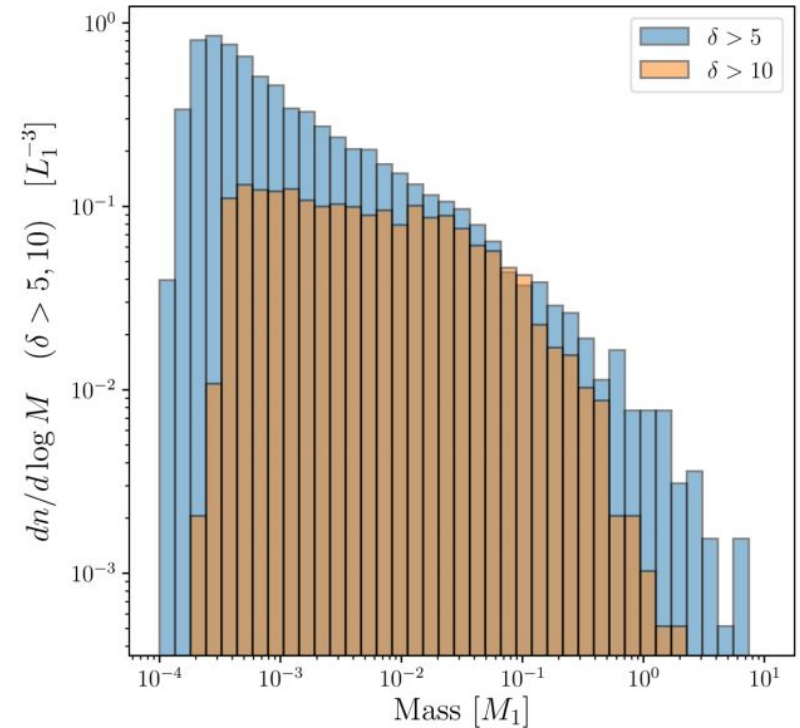
- Axion clump - Neutron Star collisions would generate spectacular radio signals
 - month-long signal for a $10^{-12} M_{\odot}$ AMC – NS collision
 - minutes for a $10^{-13} M_{\odot}$ AS – NS collision
 - AS – NS collisions could be bright enough to be observable anywhere within the Local Group
- Might be the only way to search for axion clumps
- Where and how to look? Population study needed. Inputs:
 - Neutron star population model
 - Axion minicluster mass function (+ profiles, distributions in galaxies, ...)
 - Axion star mass function (+ profiles, distributions in galaxies, ...)



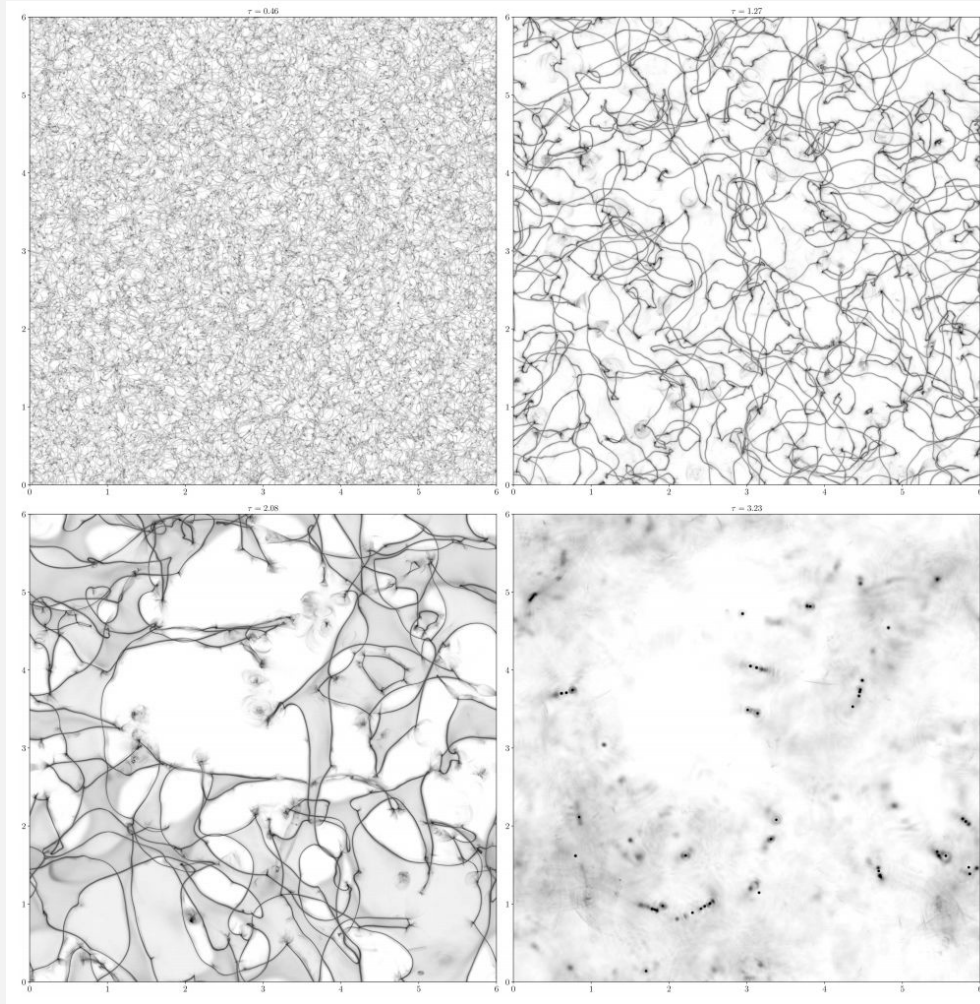
How much axion DM could be in compact objects?

Vaquero, Redondo, Stadler [1809.09241]

Minicluster “seed” mass function well after the end of the QCD phase transition

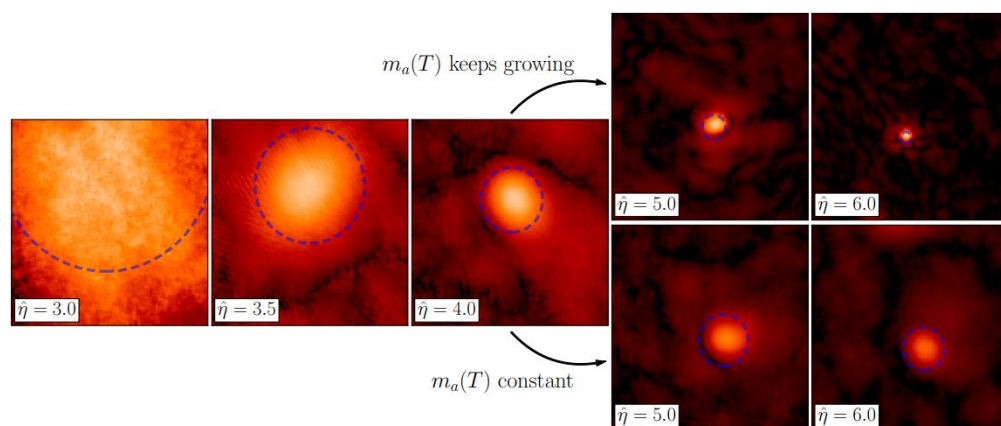
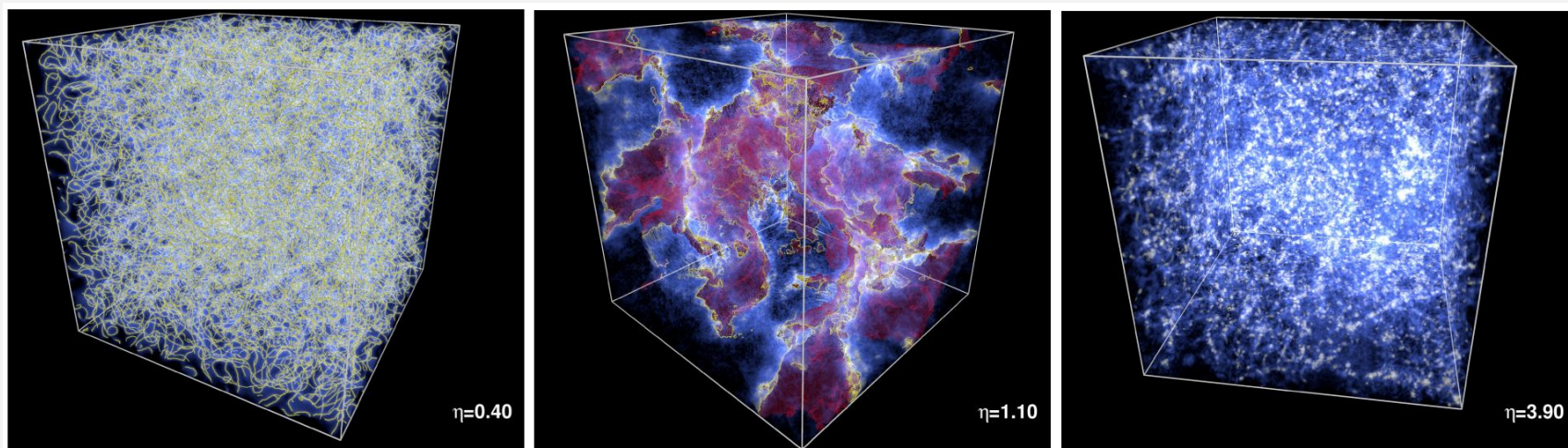


$$M_1 = 2.1 \times 10^{-12} M_\odot \frac{\Omega_{Ac} h^2}{0.12} \left(\frac{50 \mu\text{eV}}{m_a} \right)^{0.49}$$

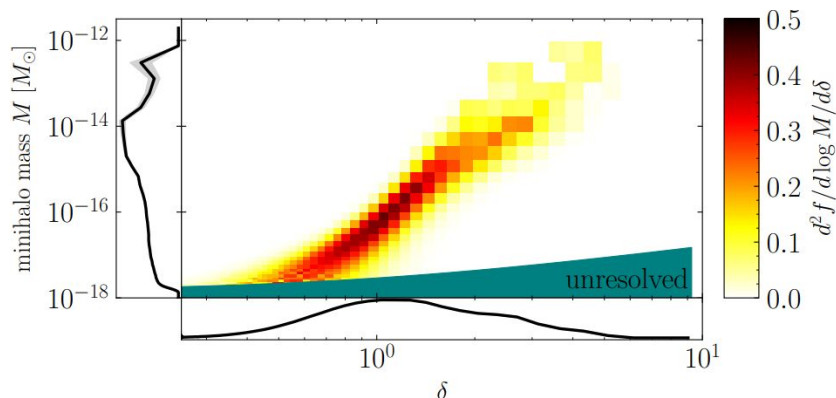


How much axion DM could be in compact objects?

Buschmann, Foster, Safdi [1906.00967]



Minicluster “seed” mass function at matter-radiation equality

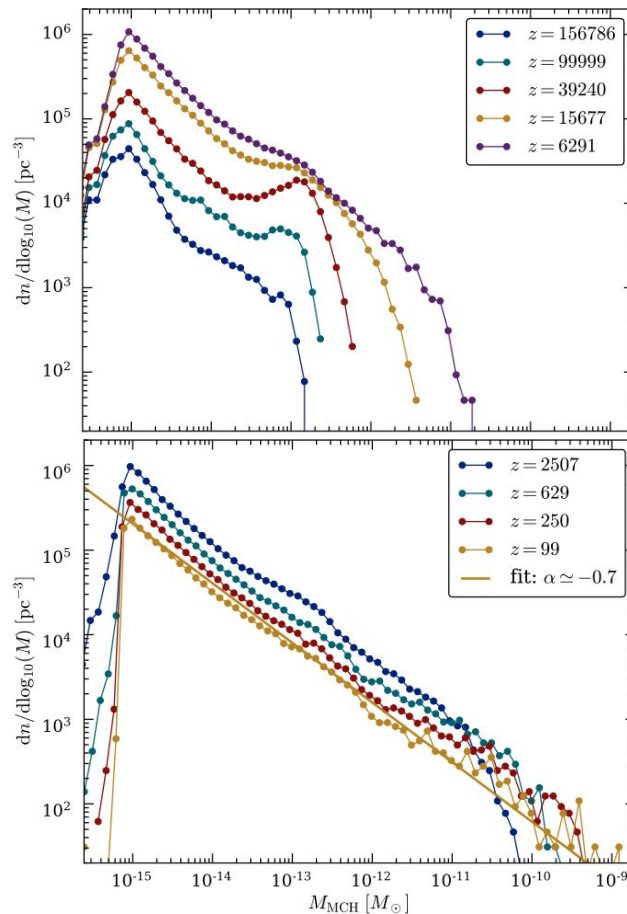


How much axion DM could be in compact objects?

Eggemeier, Redondo, Dolag, Niemeyer, Vaquero [1911.09417]

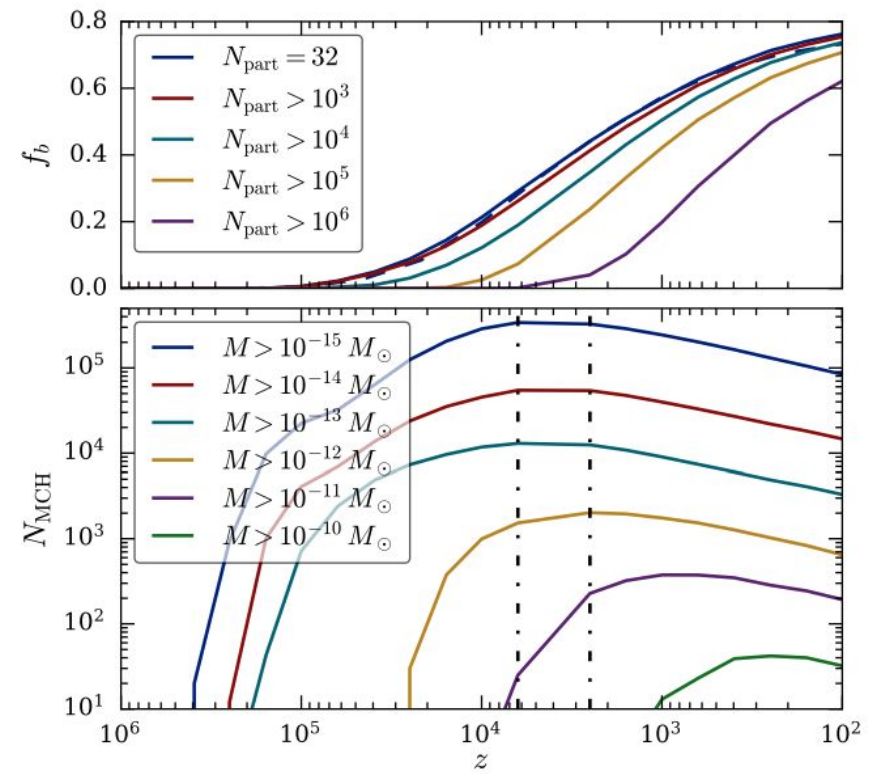
& before matter-radiation equality

Axion minicluster mass function



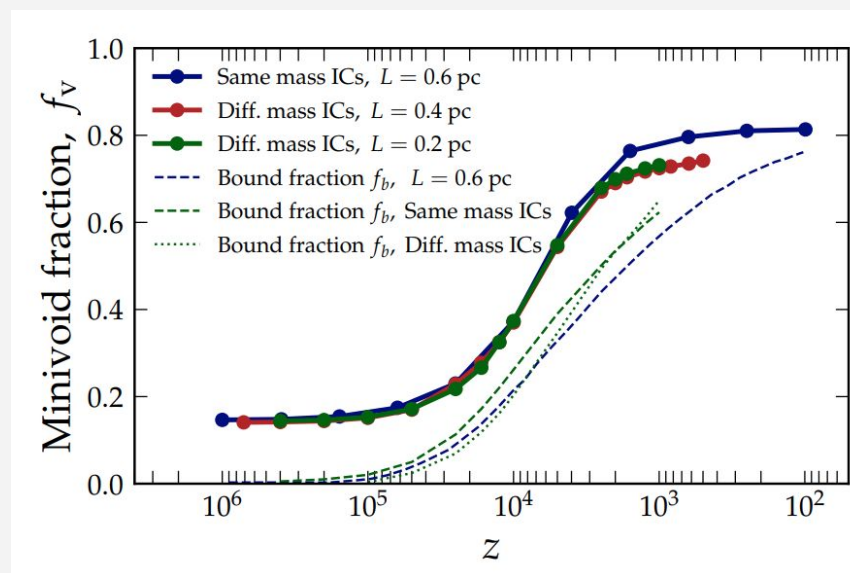
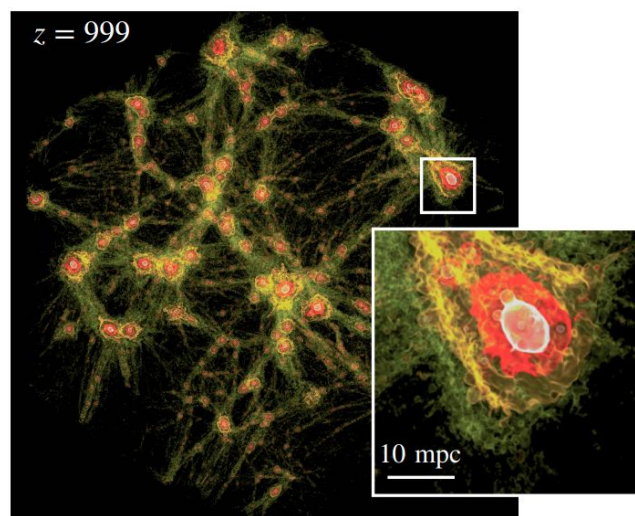
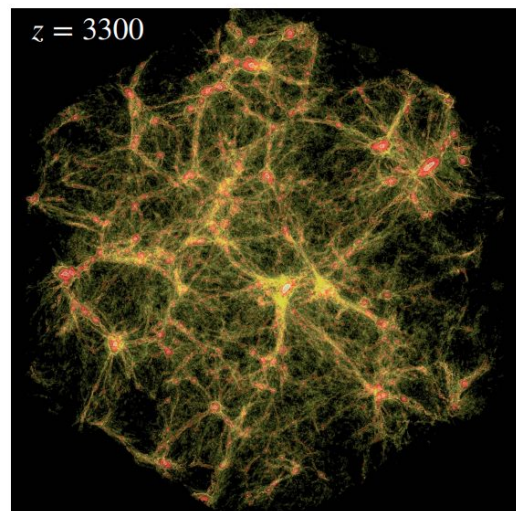
After

Bound fraction of axions



How much axion DM could be in compact objects?

Eggemeier, O'Hare, Pierobon, Redondo, Wong [2212.00560]



Axion Stars

- Dynamics of axion field ($a = f \theta$) controlled by action

$$S = \int d^4x \sqrt{-g} \left[\frac{f^2}{2} (\partial^\mu \theta) (\partial_\mu \theta) - V(\theta) \right]$$

- metric $g^{\mu\nu}$ determined by energy-momentum tensor of axion field
- Potential from chiral perturbation theory:

$$V(\theta) = \frac{\Lambda^4}{c_z} \left[1 - \sqrt{1 - 4c_z \sin^2(\theta/2)} \right] \quad \text{with} \quad c_z \approx \frac{m_u/m_d}{(1 + m_u/m_d)^2}$$

- First order self-interaction (quartic) is attractive!

Single harmonic approximation

- Approximation for $|\theta| \ll 1$

$$\theta = \Theta(r) \cos(\omega t)$$

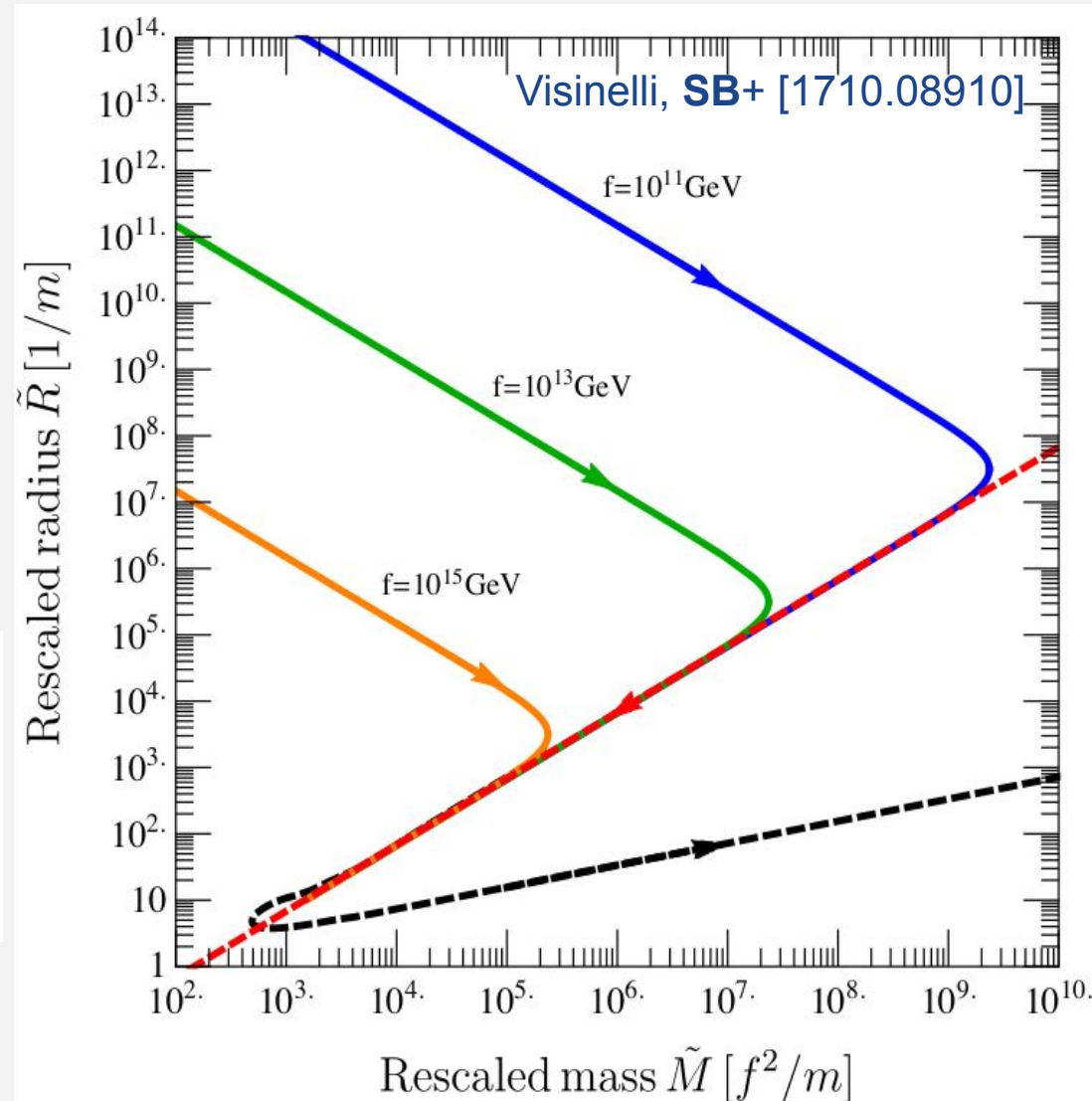
- average over time-period

$$2\pi\omega$$

Units for mass and radius:

$$\frac{f^2}{m} = 3 \times 10^{-20} M_{\odot} \left(\frac{10^{-5} \text{ eV}}{m} \right)^3,$$

$$\frac{1}{m} = 3 \times 10^{-11} R_{\odot} \left(\frac{10^{-5} \text{ eV}}{m} \right),$$



Axion star mass-radius diagram

Solutions can be understood from dimensional analysis:

For chiral potential

↙ $\lambda_\phi \approx -0.34$

$$U \propto \frac{f^2}{m} \left(-\frac{\beta \tilde{M}^2}{\tilde{R}} + \alpha_k \frac{\tilde{M}}{2\tilde{R}^2} + \frac{\alpha_4 \lambda_\phi}{4!} \frac{\tilde{M}^2}{\tilde{R}^2} \right)$$

Gravity

Quantum/Gradient
Pressure

quartic self-interaction

- Fitting to numerical results yields $\alpha_k = 9.9$ and $\alpha_4 = 1.7$
- Parameter $\beta = (f/m_{\text{pl}})^2$ measures strength of self-interaction vs gravity

Dilute Branch

Relevant contributions:

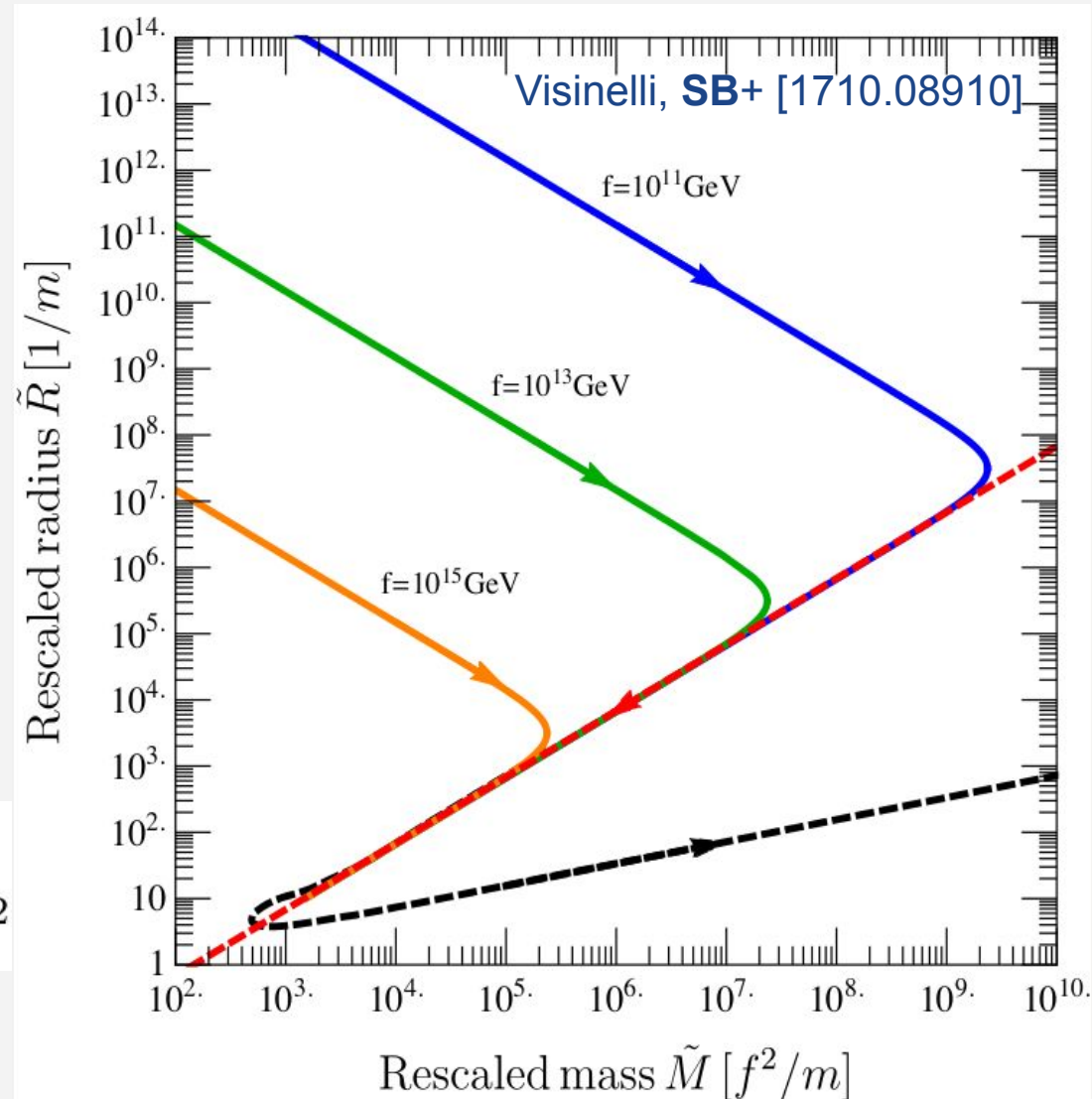
- quantum/gradient pressure
- gravity

From dimensional analysis:

$$\tilde{R}_+ \Big|_{\lambda_\phi \rightarrow 0} = \frac{\alpha_k}{\beta \tilde{M}}$$

Turning point at

$$|\theta_0| \sim 10^{-6} \times (10^{-5} \text{ eV}/m)$$
$$M \sim 10^{-11} M_\odot \times (10^{-5} \text{ eV}/m)^2$$



Axion Stars: the “Dilute Branch”

Relevant contributions:

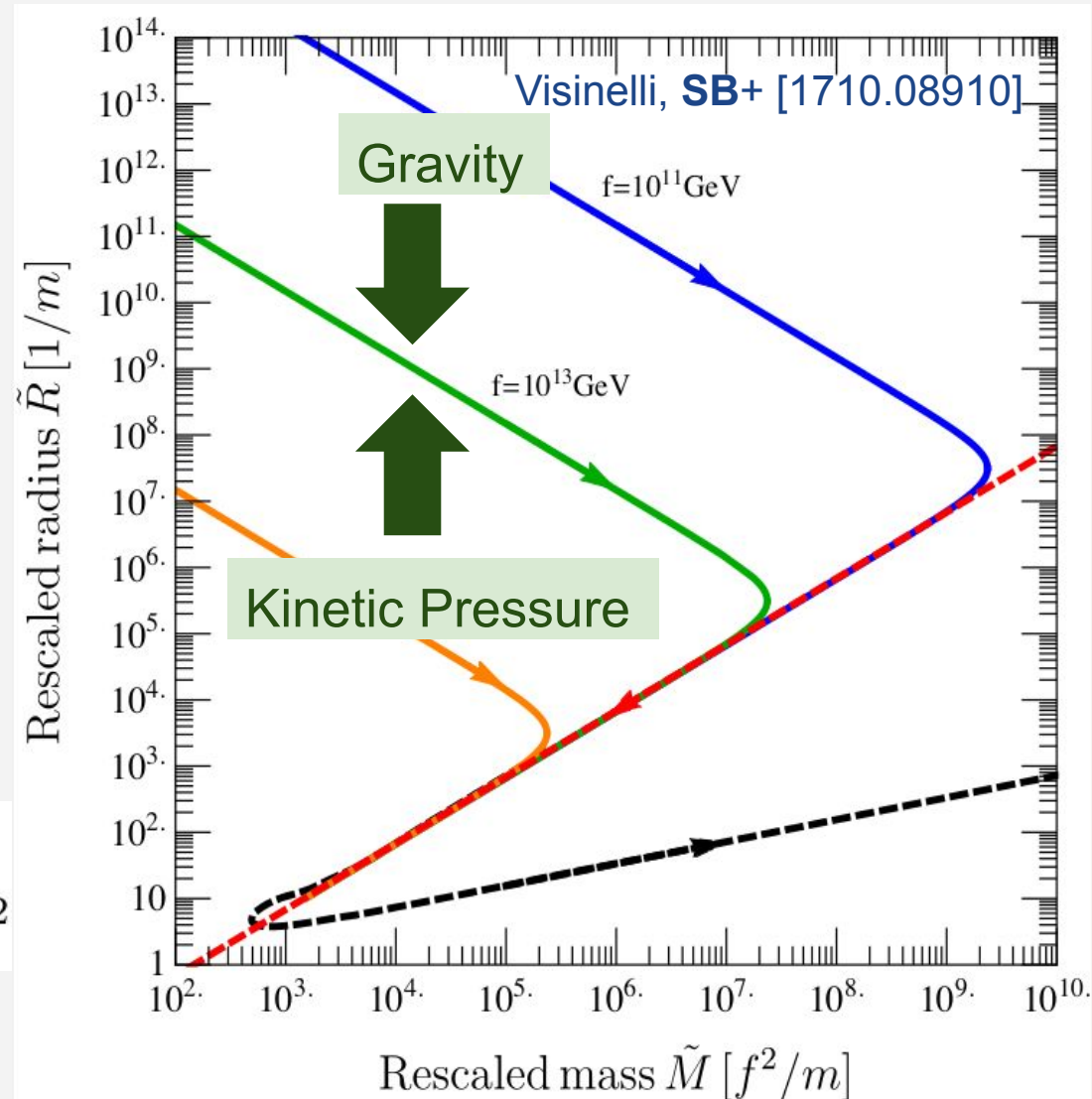
- Quantum/gradient pressure
- Gravity

From dimensional analysis:

$$\tilde{R}_+ \Big|_{\lambda_\phi \rightarrow 0} = \frac{\alpha_k}{\beta \tilde{M}}$$

Turning point at

$$|\theta_0| \sim 10^{-6} \times (10^{-5} \text{ eV}/m)$$
$$M \sim 10^{-11} M_\odot \times (10^{-5} \text{ eV}/m)^2$$



Critical Branch

Relevant contributions:

- Quantum/gradient pressure
- quartic self-interaction

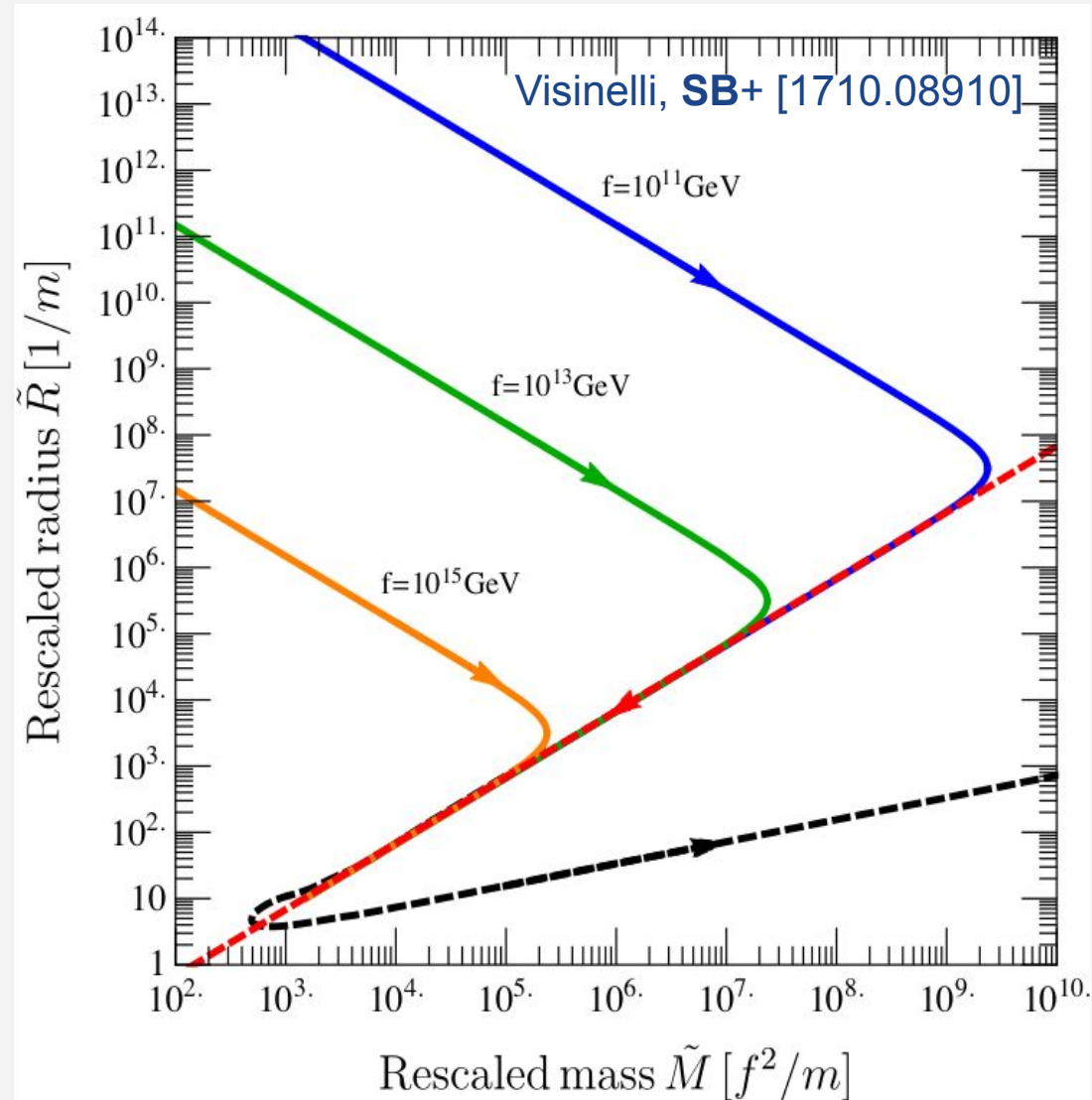
From dimensional analysis:

$$\tilde{R}_- \Big|_{G \rightarrow 0} = \frac{\alpha_4 |\lambda_\phi| \tilde{M}}{8\alpha_k}$$

Amplitude at center:

$$\Theta_0 \propto M^{-1}$$

$$\tilde{M}(\Theta_0 = 1) \approx \left(\frac{4\alpha_k}{\alpha_4 |\lambda_4|} \right)^{3/2}$$



Critical Branch

Relevant contributions:

- Quantum/gradient pressure
- quartic self-interaction

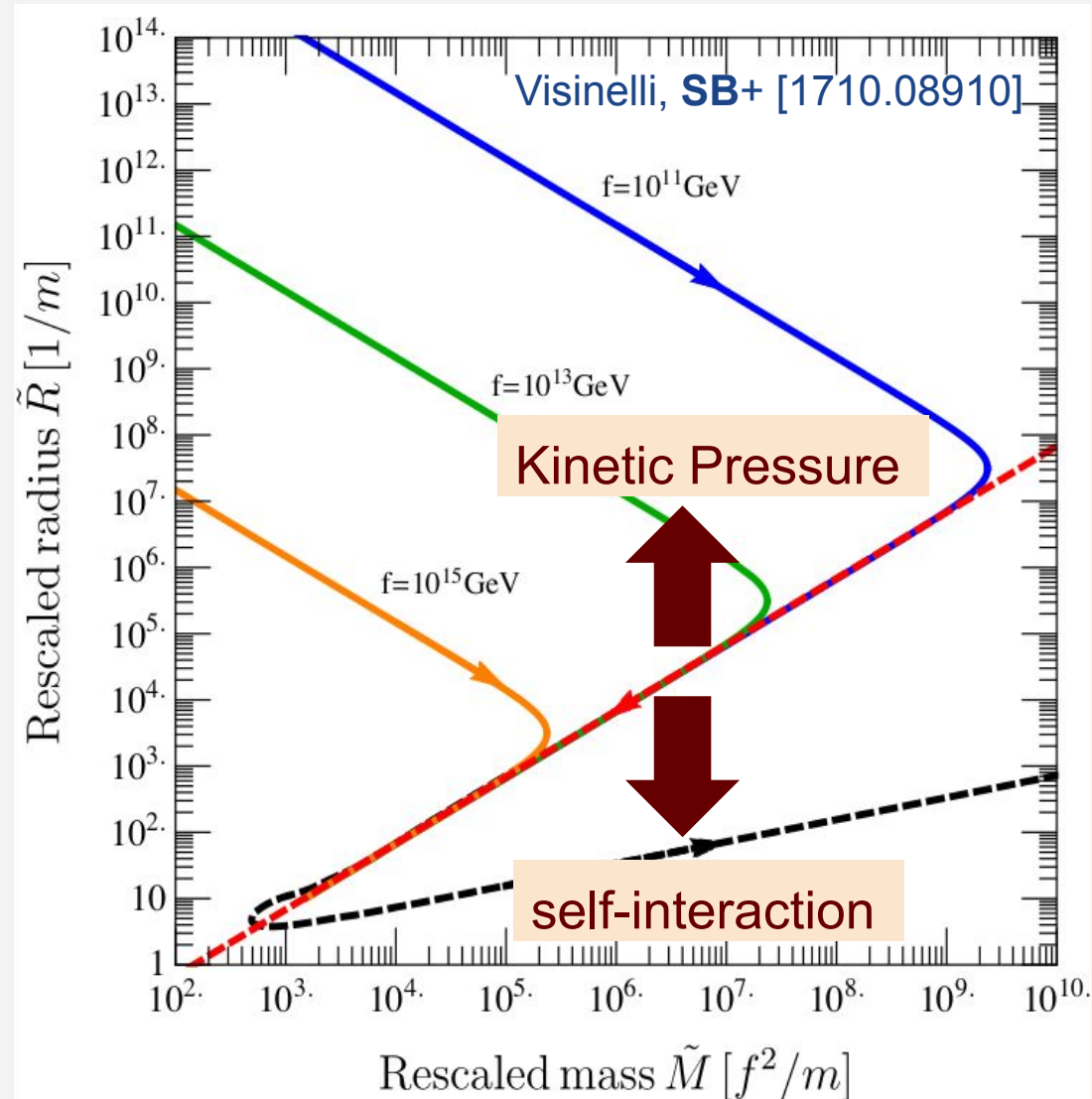
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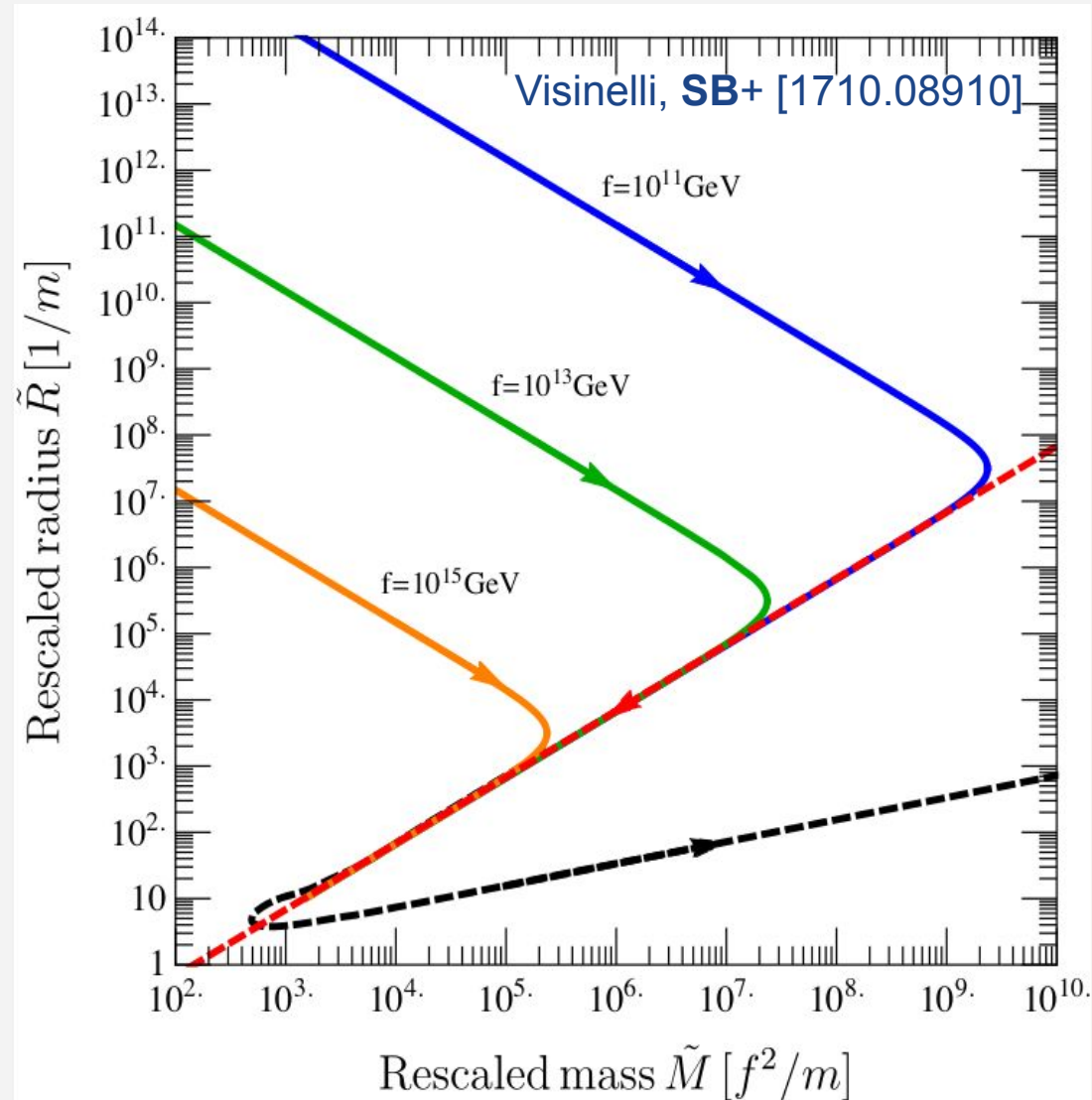
Dense Branch

Relevant contributions:

- Quantum/gradient pressure
- gradient energy
- all-order self-interaction

Saturated field: $\rho \sim \Lambda^4$

$$\tilde{R} \approx 0.6 \tilde{M}^{1/3}$$



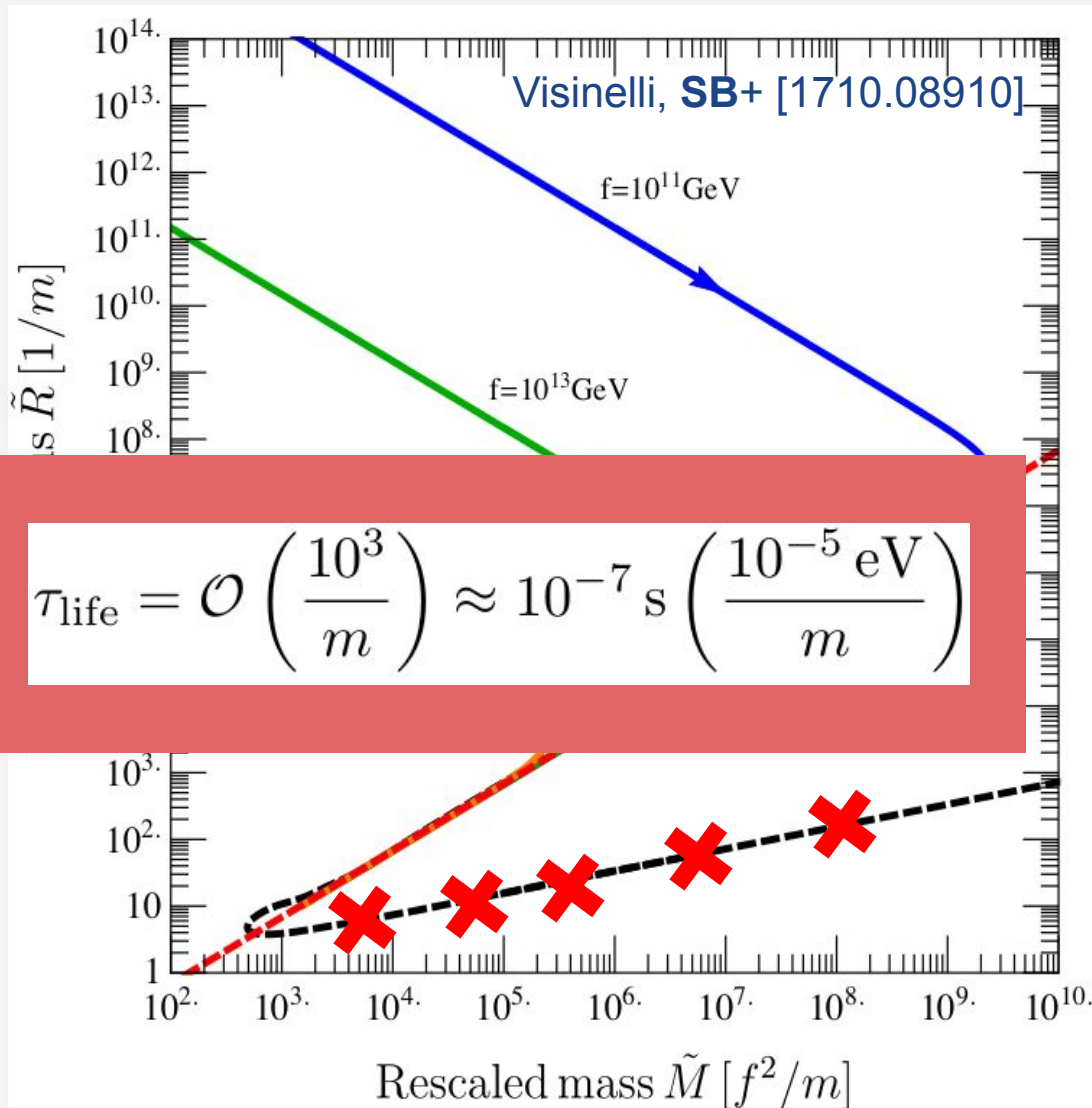
Dense Branch

Relevant contributions:

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- gradient energy
- all-order self-interaction

Saturated field: $\rho \sim \Lambda^4$

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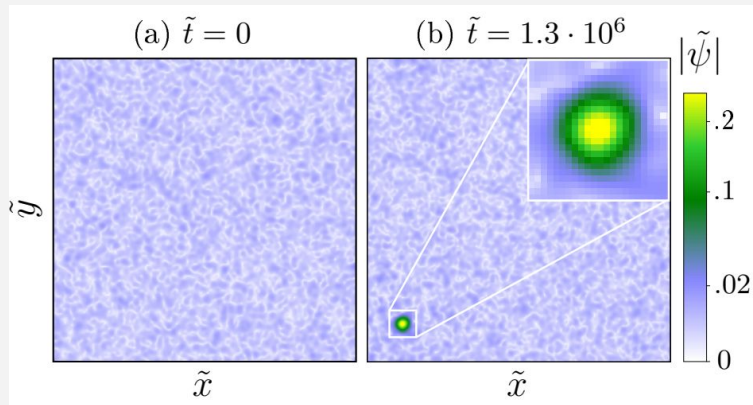
Axion Stars

- Only axion stars on the ‘dilute branch’ are stable, maximal mass $M \sim 10^{-11} M_{\odot} \times (10^{-5} \text{ eV}/m)^2$
- Dense axion stars either radiate relativistic axions until puffing out to the dilute branch, or might collapse to black holes

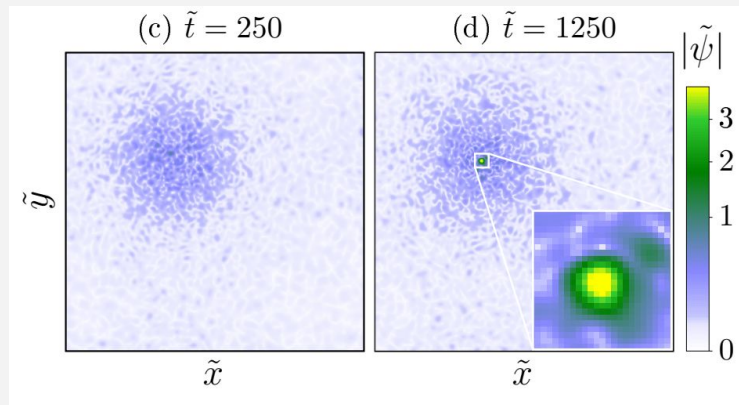
$$\tau_{\text{life}} = \mathcal{O} \left(\frac{10^3}{m} \right) \approx 10^{-7} \text{ s} \left(\frac{10^{-5} \text{ eV}}{m} \right)$$

Formation of axion stars in miniclusters

- In random field



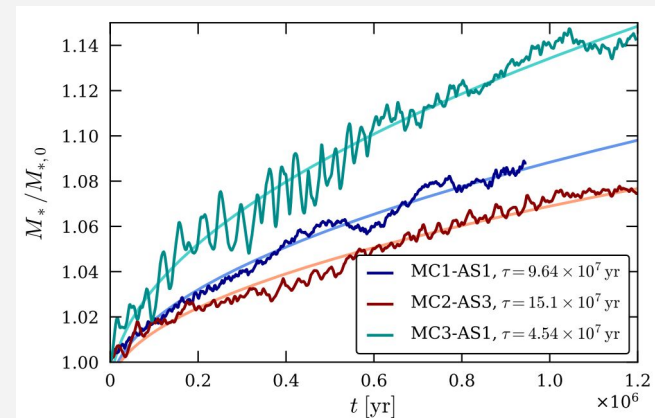
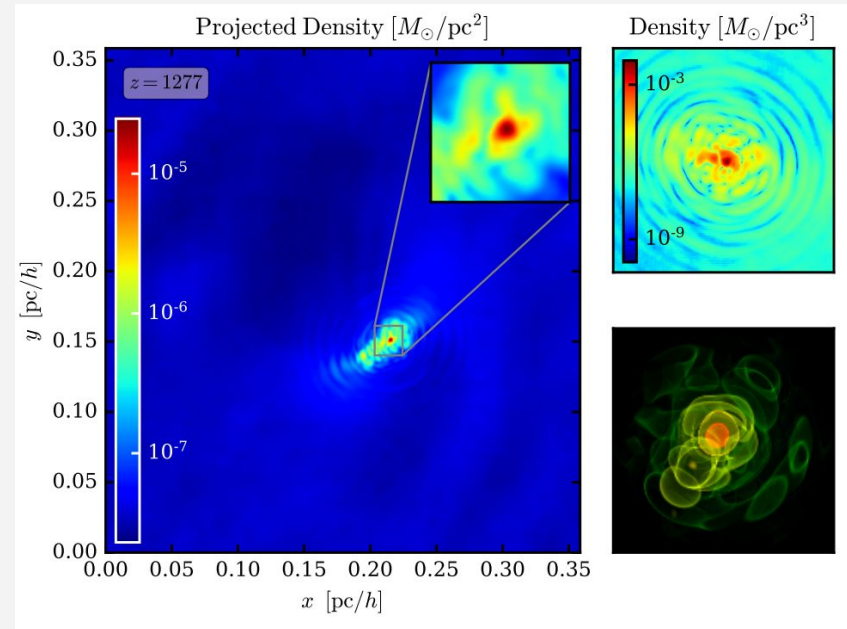
- In a minicluster



$$\tau_{gr} \sim \frac{10^9 \text{ yr}}{\Phi^3(1 + \Phi)} \left(\frac{M_c}{10^{-13} M_\odot} \right)^2 \left(\frac{m}{26 \mu\text{eV}} \right)^3$$

Levkov, Panin, Tkachev [1804.05857]

- In miniclusters



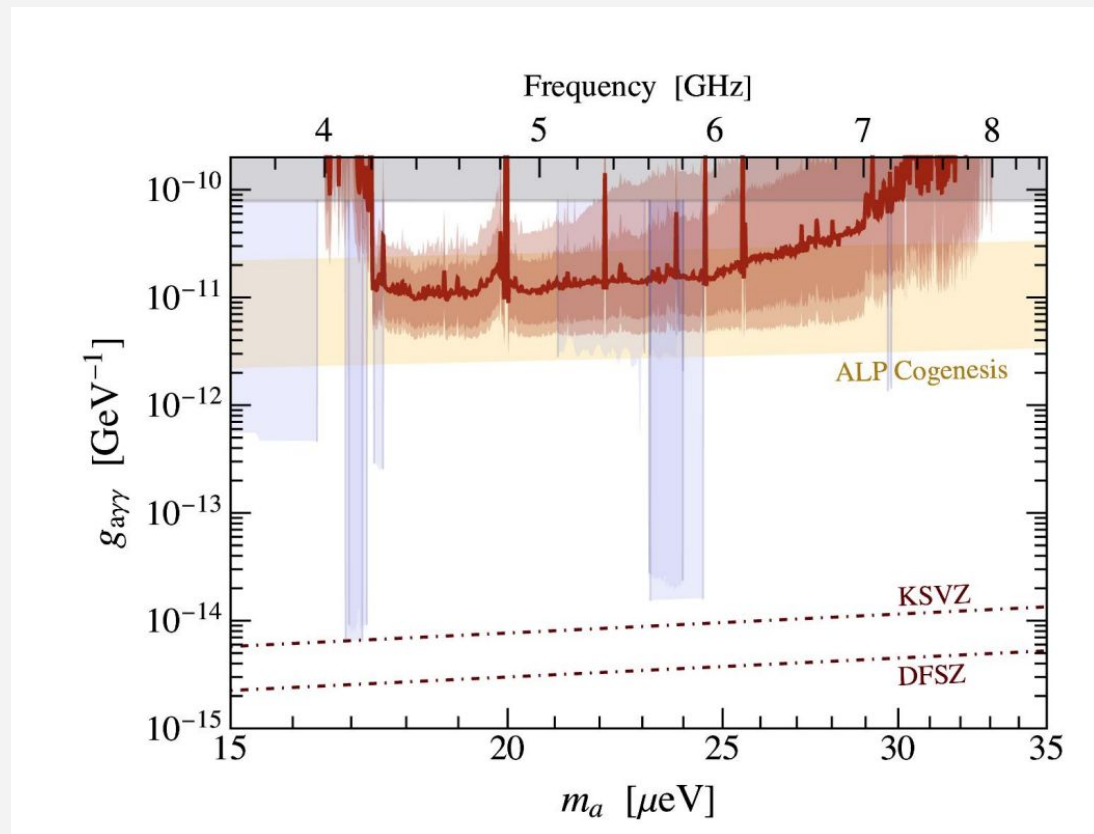
Eggemeier & Niemeyer [1906.01348]

Axions meeting Neutron Stars

Foster, Witte, Lawson, Linden, Gajjar, Weniger, Safdi [2202.08274]

Most up-to-date constraint:

- Observations of the galactic center (GBT Breakthrough Listen Data)



Minicluster - Neutron Star Collision: NS rotation

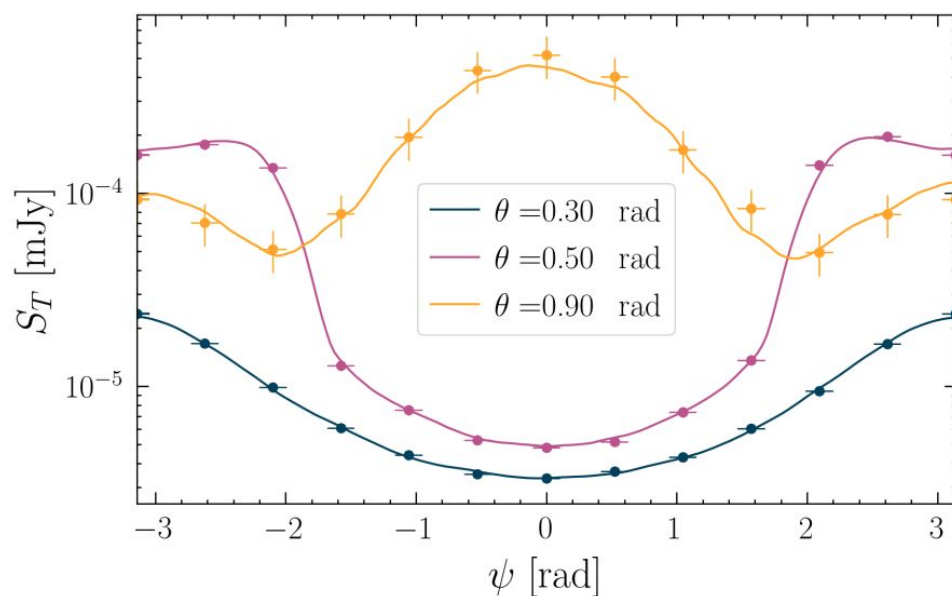
$$\psi = \Omega \times t = 0^\circ$$

Symbols: simulations at different ψ

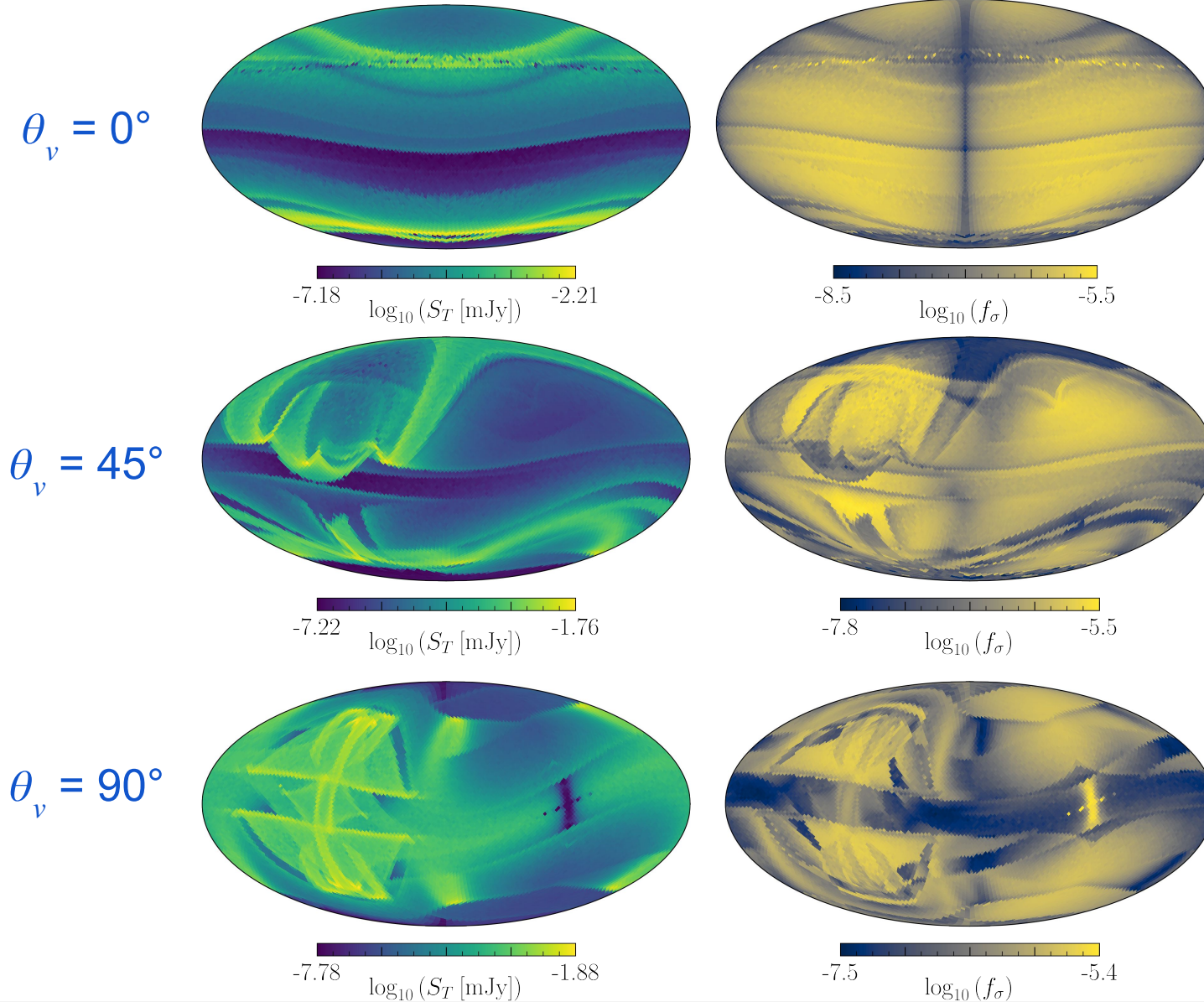
Line: azimuthally-averaged flux

$$\psi = \Omega \times t$$

$$t = 108^\circ$$



AMC - NS Collision: Varying the Encounter-Geometry



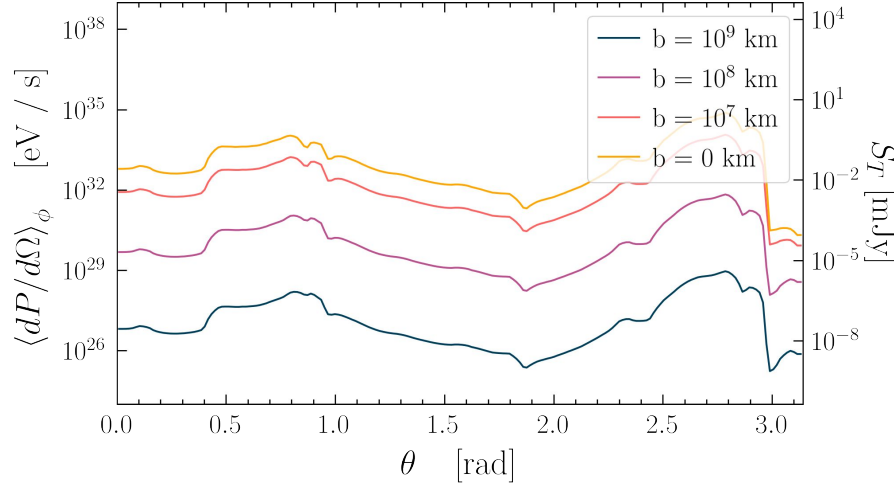
m_a	$26 \mu\text{eV}$
$g_{a\gamma\gamma}$	$10^{-14} \text{ GeV}^{-1}$
B_0	$1.6 \times 10^{14} \text{ G}$
P	3.76 s
θ_m	0.2 rad
θ_v	0.0 rad
δf_T	$10^{-5} \times m_a$
d_T	1 kpc

M_{AMC}	$10^{-12} M_\odot$
R_{AMC}	$1.86 \times 10^9 \text{ km}$
M_{AS}	$10^{-13} M_\odot$
R_{AS}	3905 km
ρ_{AMC}	NFW
b_{AMC}	$10^8 \hat{x} \text{ km}$
b_{AS}	$2 \times 10^3 \hat{y} \text{ km}$
$ v_{\text{AMC/AS}} $	100 km/s

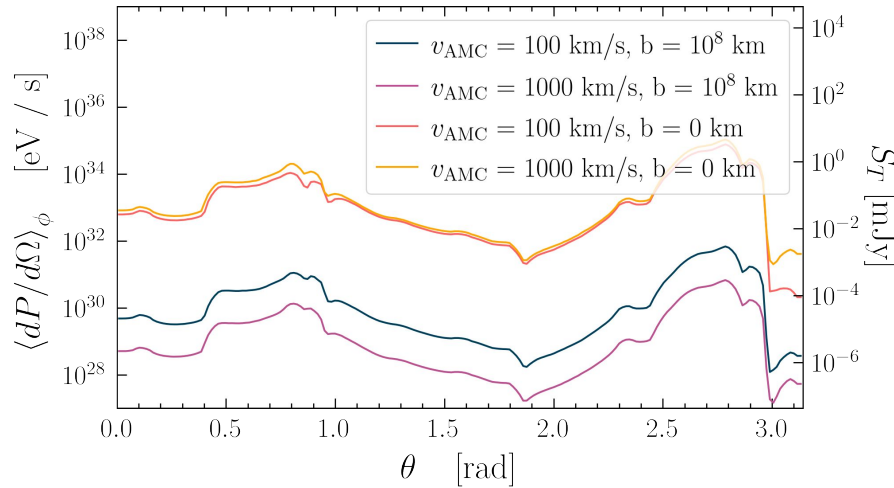
AMC - NS Collision: Varying the Encounter-Geometry

Azimuthally-averaged flux
 $\langle dP/d\Omega \rangle_\phi \equiv \frac{\omega}{2\pi} \int d\phi \frac{dP}{d\Omega}(\phi, \theta)$

Varying the impact parameter



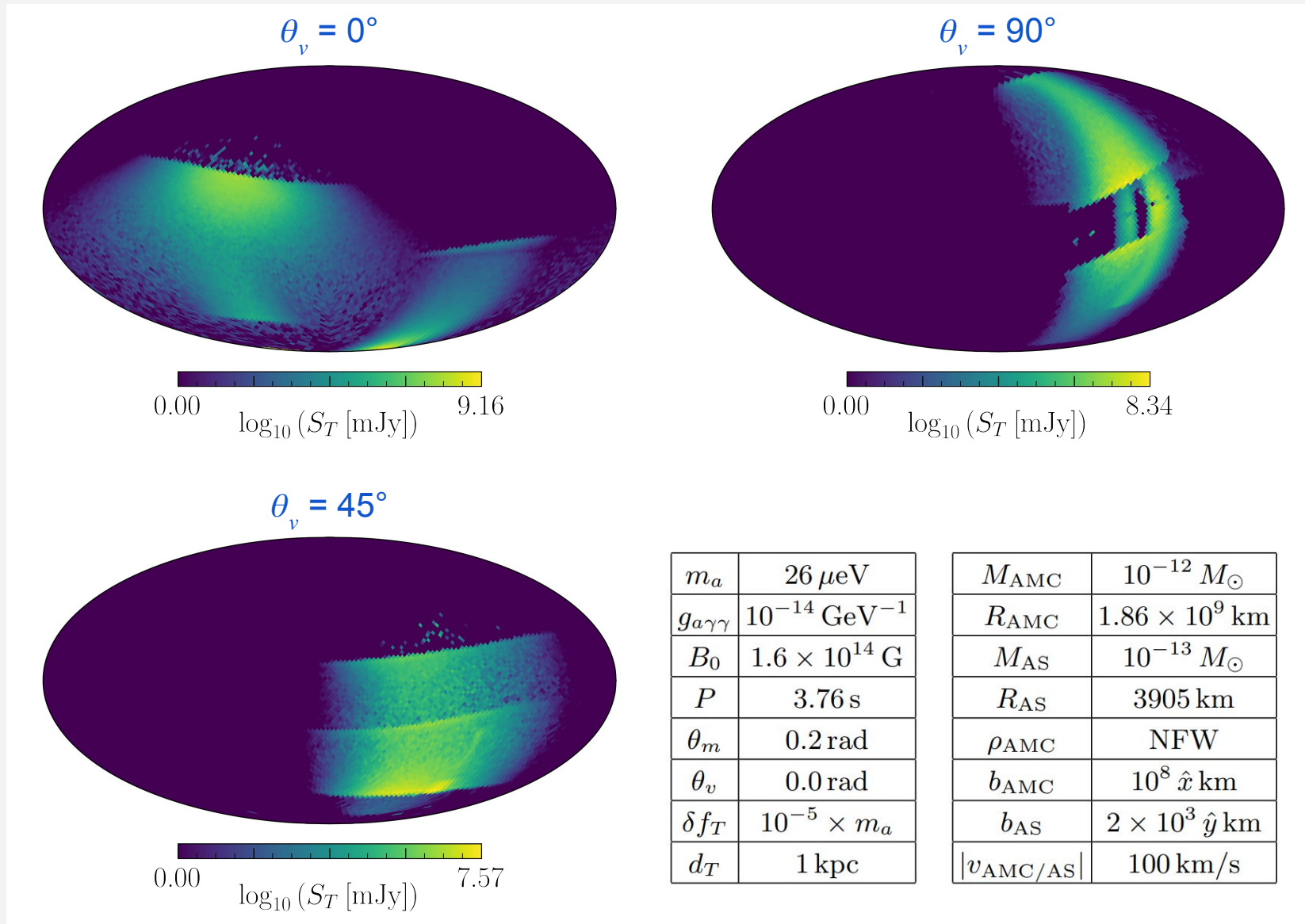
Varying the impact parameter & speed



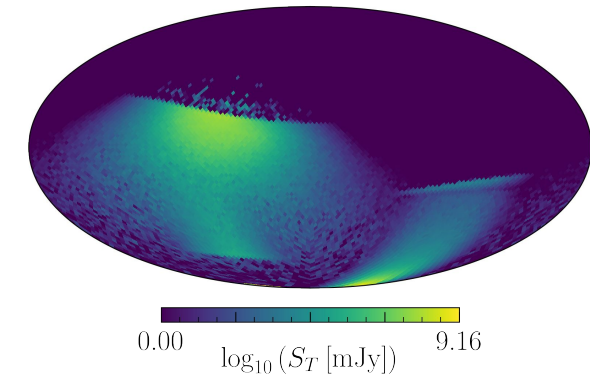
m_a	$26 \mu\text{eV}$
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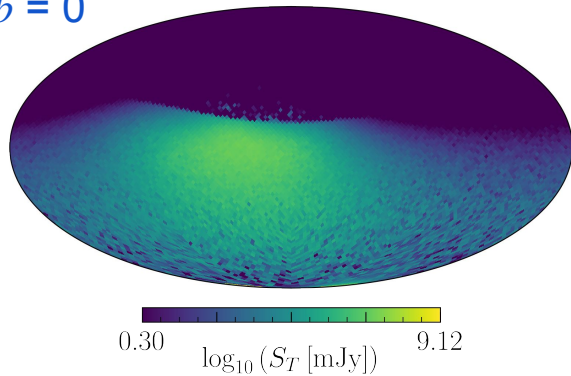
AS - NS Collision: Varying the Encounter-Geometry



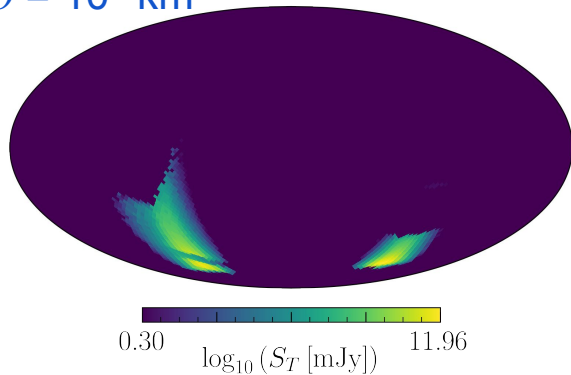
AS - NS Collision: Varying the Encounter-Geometry



$b = 0$



$b = 10^4 \text{ km}$

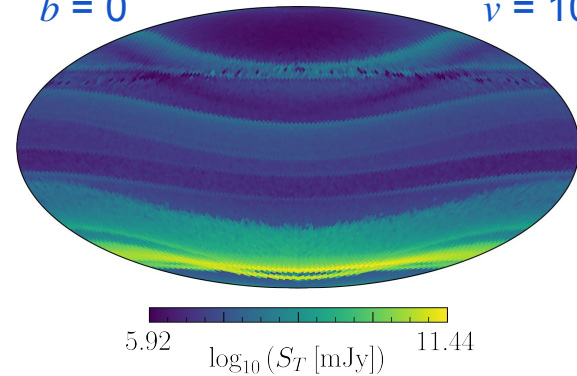


m_a	$26 \mu\text{eV}$
$g_{a\gamma\gamma}$	$10^{-14} \text{ GeV}^{-1}$
B_0	$1.6 \times 10^{14} \text{ G}$
P	3.76 s
θ_m	0.2 rad
θ_v	0.0 rad
δf_T	$10^{-5} \times m_a$
d_T	1 kpc

M_{AMC}	$10^{-12} M_\odot$
R_{AMC}	$1.86 \times 10^9 \text{ km}$
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R_{AS}	3905 km
ρ_{AMC}	NFW
b_{AMC}	$10^8 \hat{x} \text{ km}$
b_{AS}	$2 \times 10^3 \hat{y} \text{ km}$
$ v_{\text{AMC/AS}} $	100 km/s

$b = 0$

$v = 10^3 \text{ km/s}$



$b = 2 \times 10^3 \text{ km}$

$v = 10^3 \text{ km/s}$

