

Spin-1 Thermal Targets for Dark Matter Searches at Fixed Target Experiments

arXiv:2307.02207

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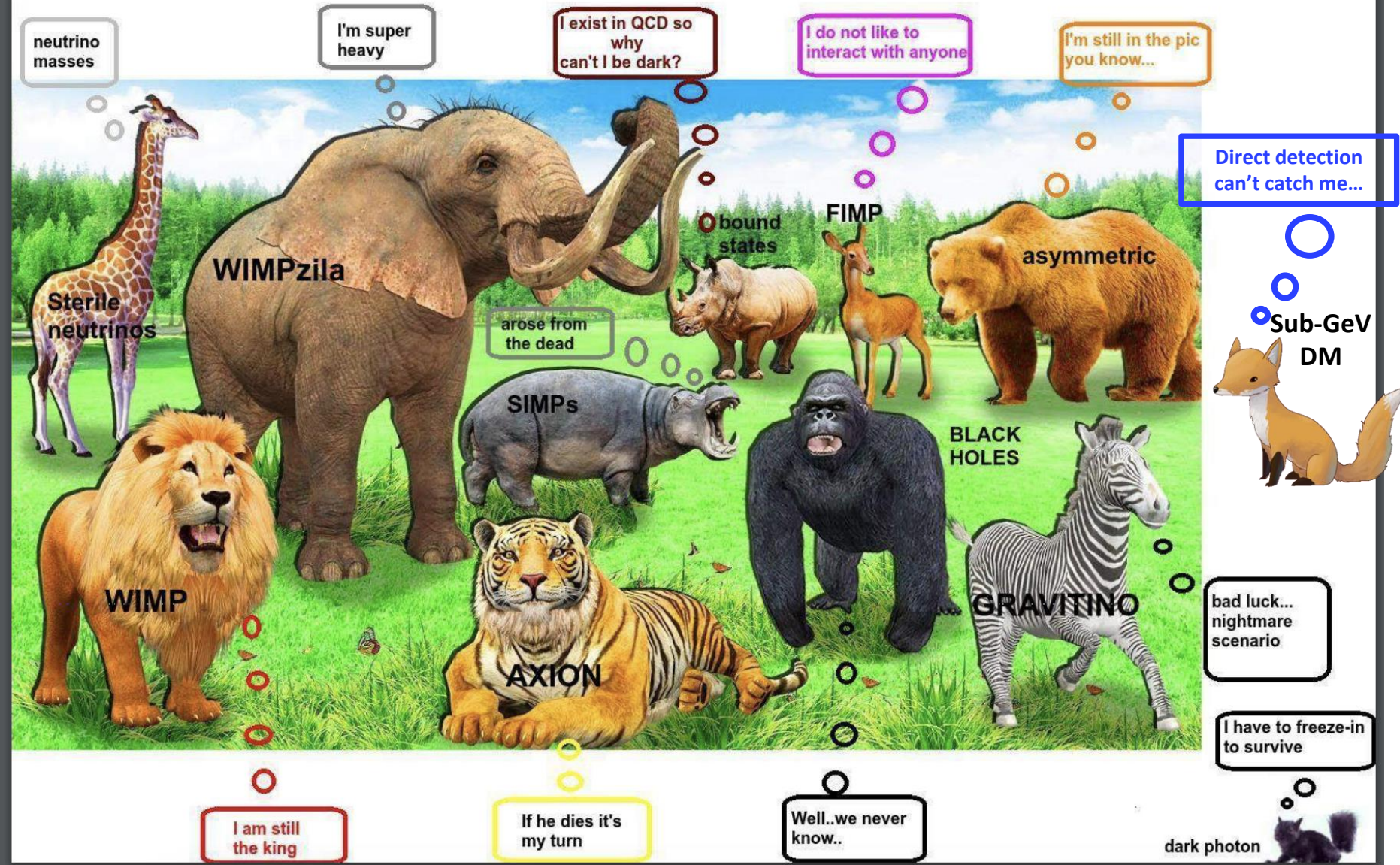
TAUP Vienna 2023



*Knut and Alice
Wallenberg
Foundation*

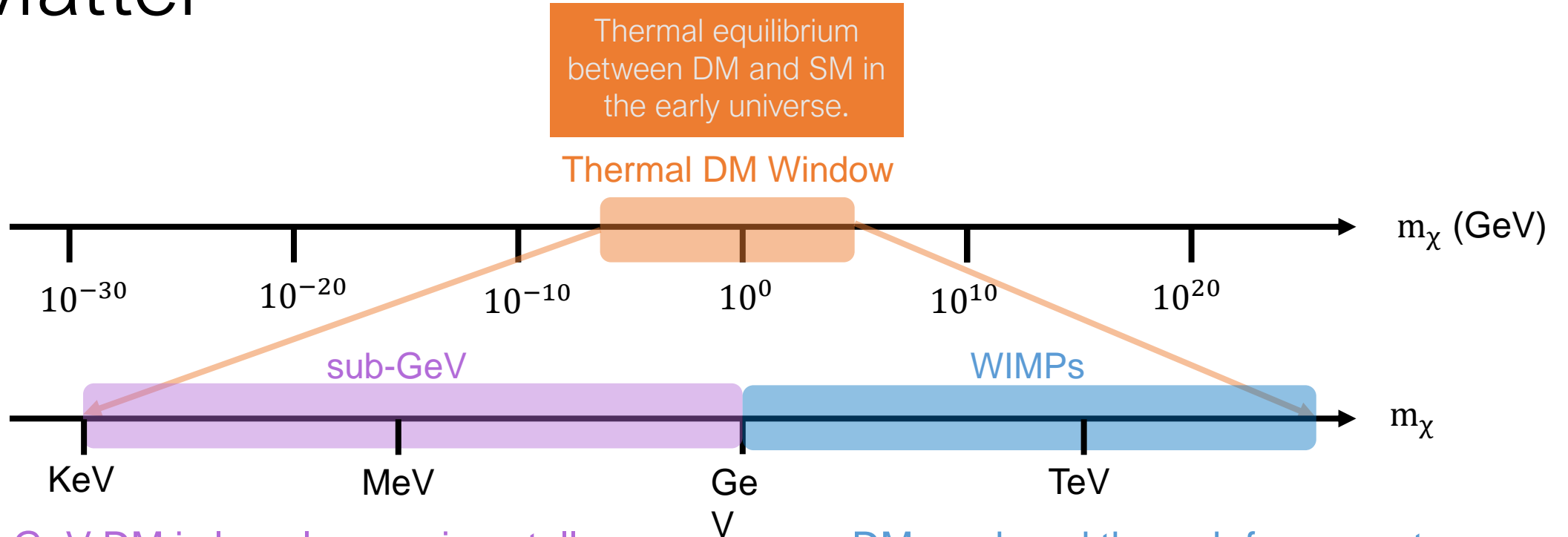


ZOO OF DARK MATTER CANDIDATES



- Was once in thermal equilibrium w SM
- mass: \sim MeV – GeV
- interacts w SM via a new hypothetical dark mediator particle

Sub-GeV Dark Matter

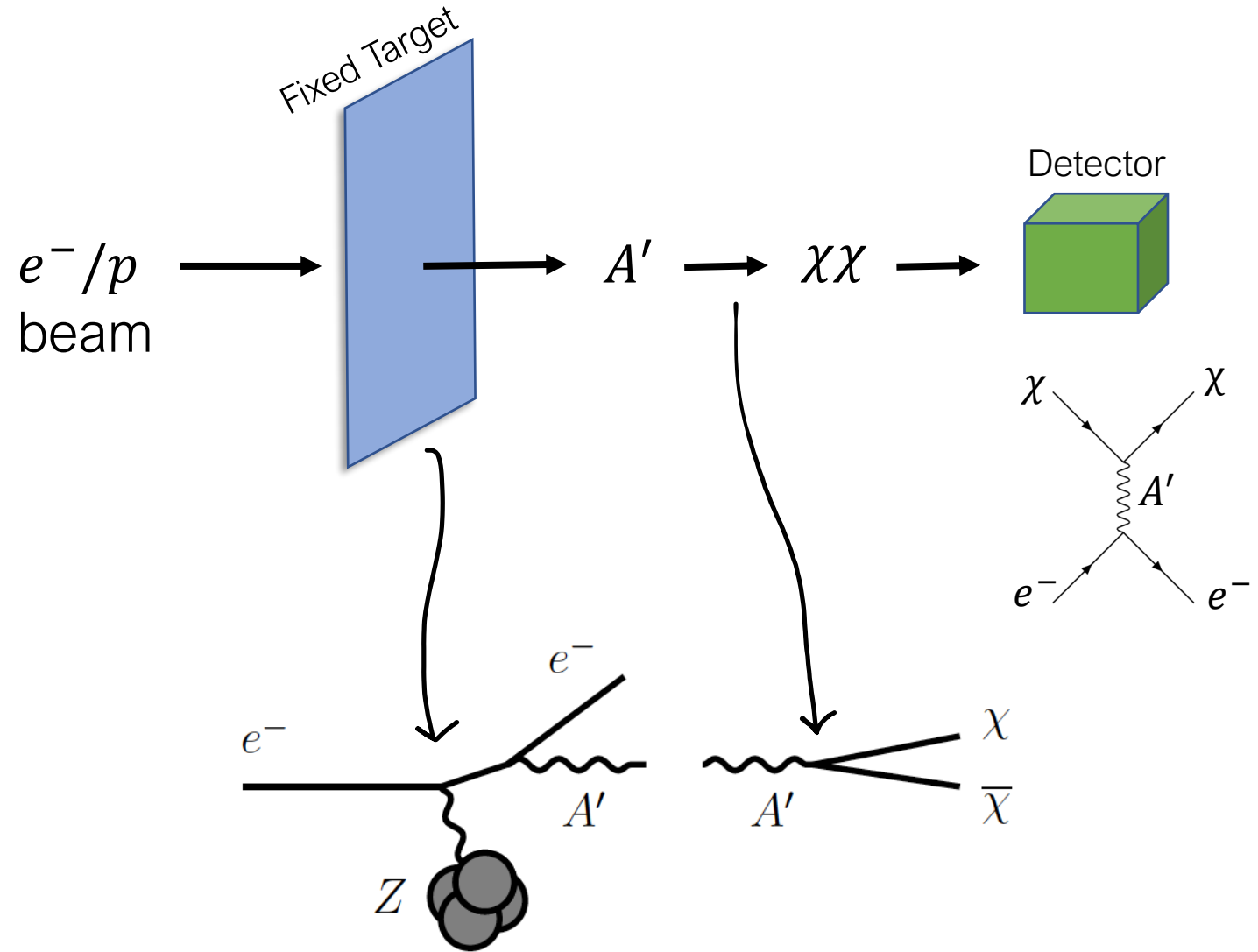


- Sub-GeV DM is largely experimentally **unexplored**..
 - Out of reach of nuclear recoil direct detection expts

- DM produced through freeze-out near weak scale
- GeV-TeV scale thermal DM already widely tested

Sub-GeV Dark Matter

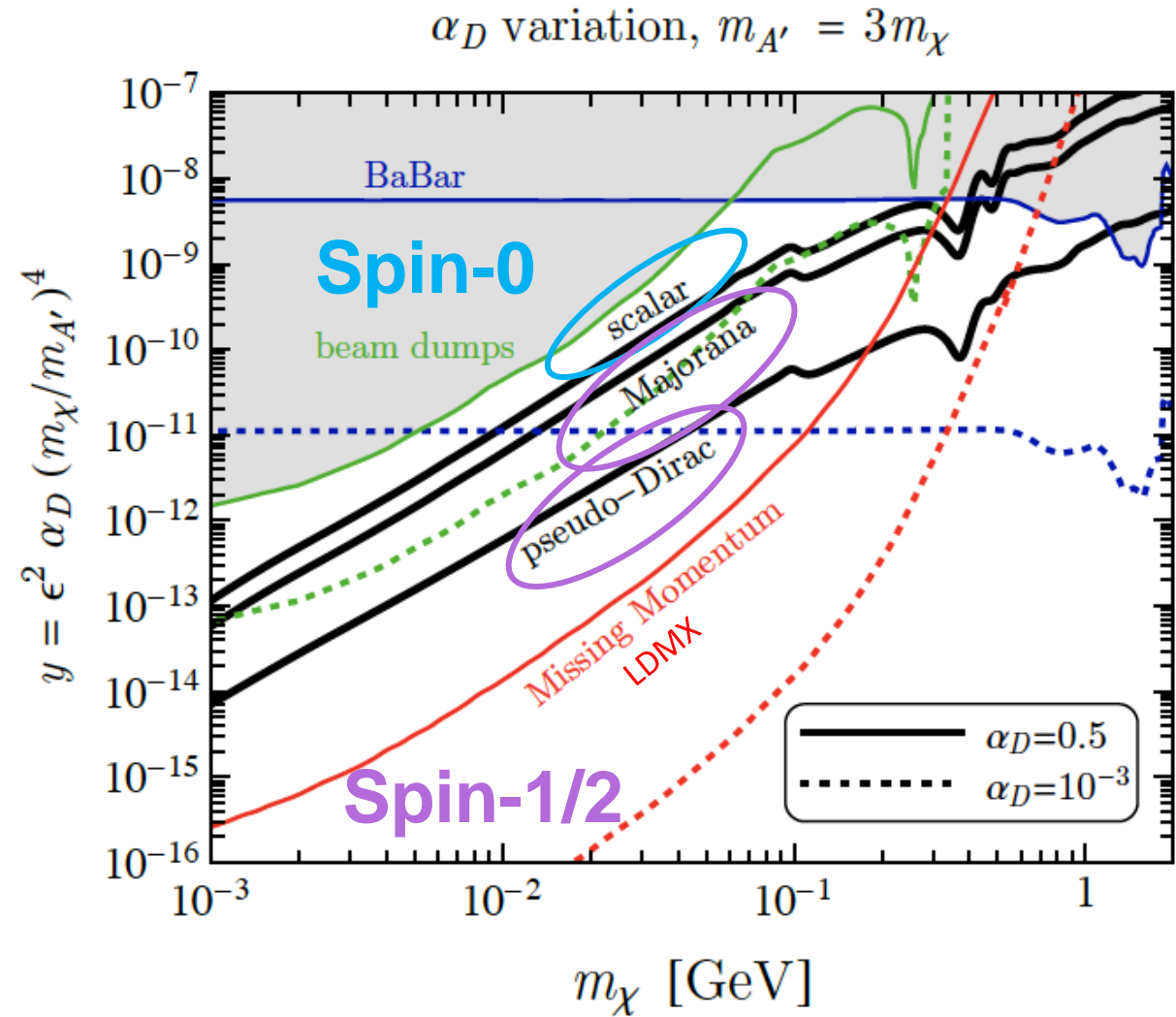
- Fixed target experiments can probe sub-GeV DM
- Future fixed target experiments such as **LDMX** will reach new sensitivities in the sub-GeV mass range.



See Lene Kristian's talk
next on LDMX!

Sub-GeV Dark Matter

- Fixed target experiments can probe sub-GeV DM
- Future fixed target experiments such as **LDMX** will reach new sensitivities in the sub-GeV mass range.
- How about **spin-1** DM?



arXiv:1807.0170

The Goal

- Broaden the existing studies on **sub-GeV DM** at **fixed target experiments**
- We consider,
 1. a set of simplified **spin-1 DM** candidates which have a **dark photon** mediator X, A'
 2. a renormalizable, UV complete SIMP **spin-1 DM** model with two **mediators** X, A', X_3

Limits/Projections

Electron Dark-Bremsstrahlung + (In)Visible Decay

- **LDMX [Light Dark Matter eXperiment]**

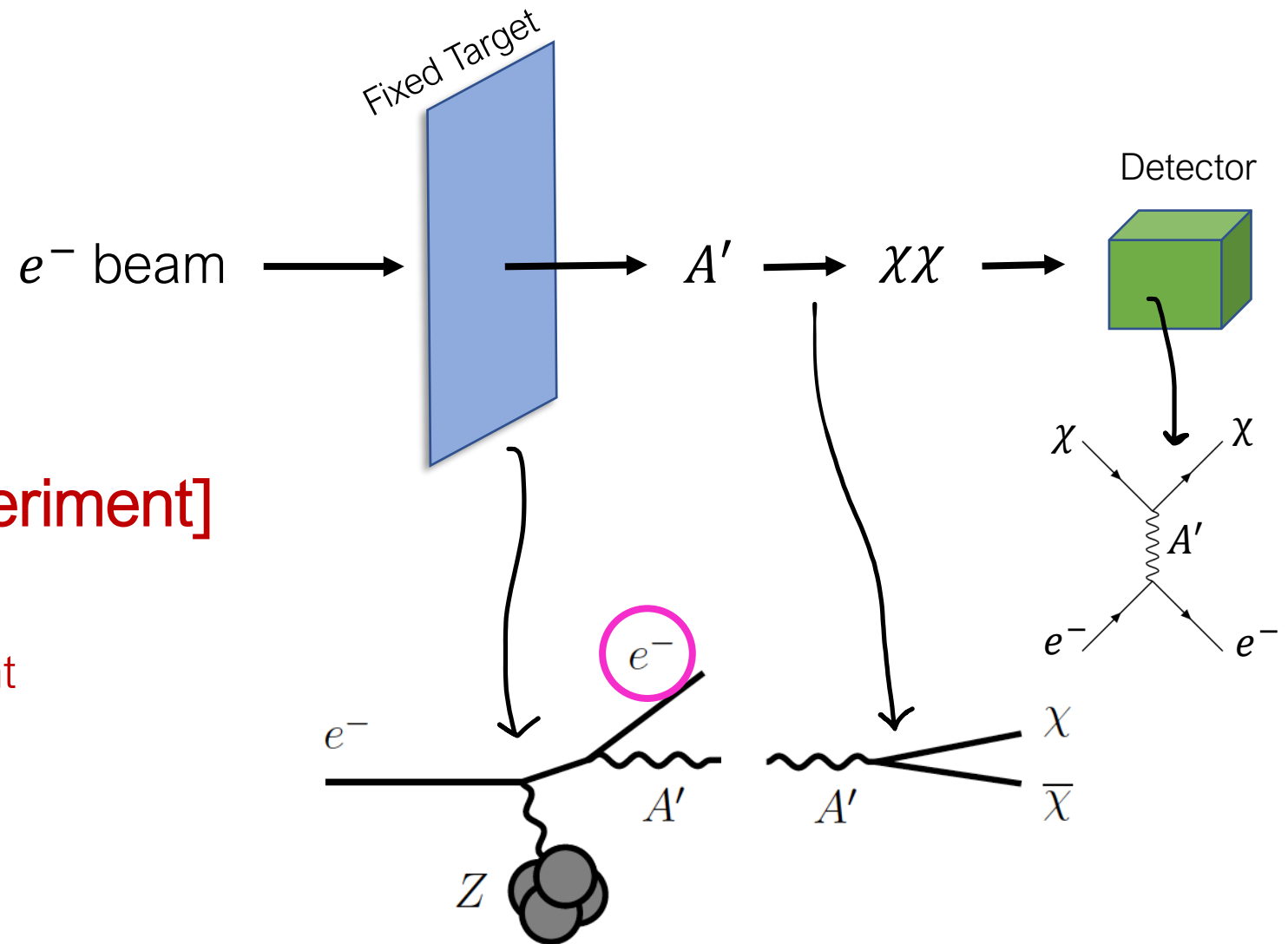
(arXiv:1808.05219)

- Future Experiment
- Missing Energy/Momentum experiment
- pT and E distributions of recoil e^-

- **NA64** (arXiv:1906.00176)

- Missing Energy
- E distributions of recoil e^-

- **SLAC E137** (arXiv:1406.2698)



Limits/Projections

Proton Beam Dumps with
Downstream DM Detector

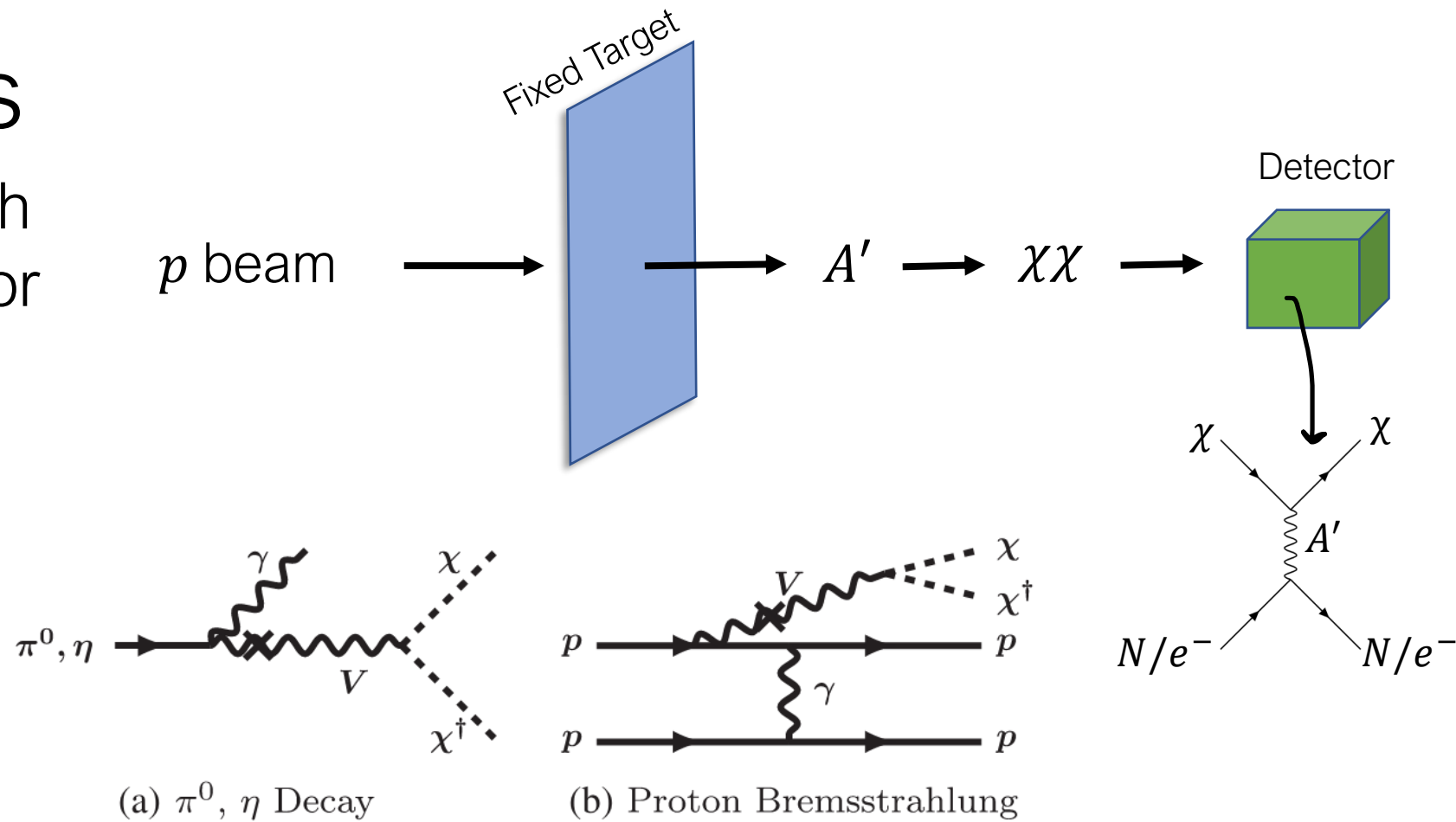
BdNMC software

LSND (arXiv:1107.4580)

- DM-electron scattering

MiniBooNE (arXiv:1702.02688)

- DM-electron/nucleon scattering



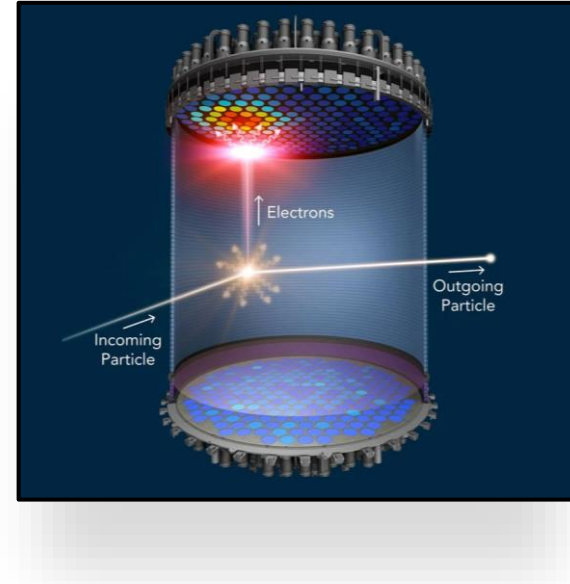
Limits/Projections

Monophoton Searches

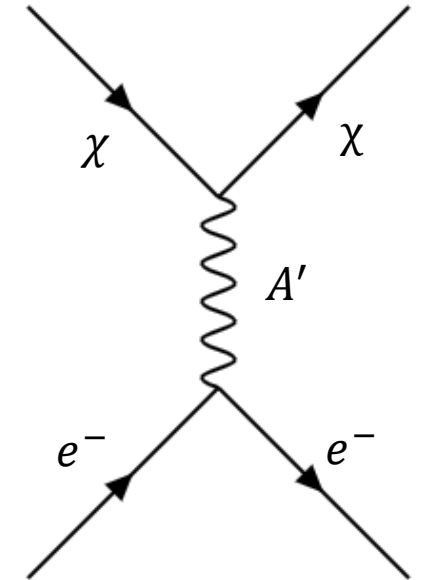
$$e^+e^- \rightarrow \gamma A', A' \rightarrow XX$$

- BaBar (arXiv:1702.03327)
- Belle II
 - Future experiment

Direct Detection (arXiv:2210.07305)



- Edelweiss
- Sensei
- Xenon1T
- Xenon10



Limits/Projections

Energy Injection

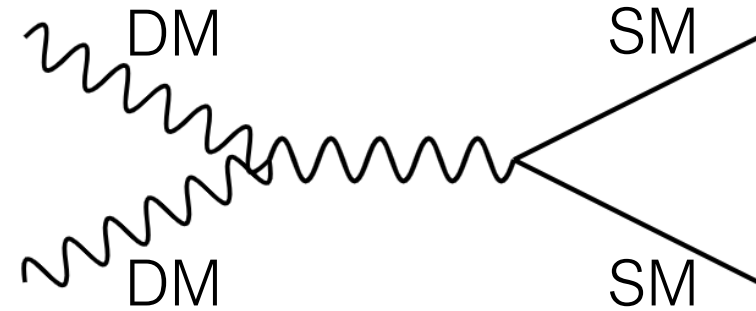
CMB

- anisotropies measurements by Planck constrain the annihilation parameter, P_{ann}

$$P_{ann} \equiv f(z) \frac{\langle \sigma v \rangle_{\chi\chi \rightarrow f\bar{f}}}{m_\chi}$$

$$P_{ann} \lesssim 3.2 \times 10^{-28} \text{cm}^3 \text{s}^{-1} \text{GeV}^{-1}$$

(Planck 2018)



$$\langle \sigma v \rangle = \underbrace{a}_{\langle \sigma v \rangle_s} + \underbrace{bv^2}_{\langle \sigma v \rangle_p} + \dots$$

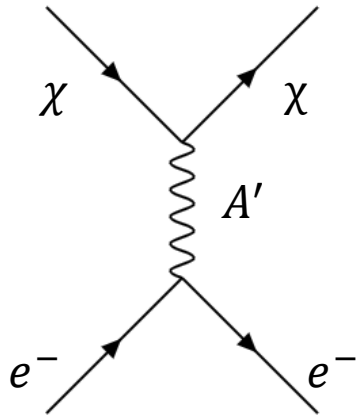
IGM temperature

- Lyman- α forest measurements, we require that the energy injected into the IGM does not overheat it at late times

Limits/Projections

Unitarity Violation

If the matrix elements in the theory are too large at tree level, additional fields or higher order diagrams are needed to restore unitarity of the S matrix.



arXiv:1510.02110: simplified *DM*

arXiv:2303.08351: spin-1 *DM* self scattering

$$|\mathrm{Im}(M_{ii}^J)|, 2|\mathrm{Re}(M_{ii}^J)| \leq 1$$

- Need to be careful of this for the **simplified spin-1 DM** models...
- But, the renormalizable and UV complete **SIMP spin-1 DM** model by construction does not violate unitarity!

Simplified Spin-1 Dark Matter Models

with a Dark Photon Mediator

$$-\mathcal{L} \supset (i b_5 X_\nu^\dagger \partial_\mu X^\nu A'^\mu + b_6 X_\mu^\dagger \partial^\mu X_\nu A'^\nu + b_7 \epsilon_{\mu\nu\rho\sigma} (X^{\dagger\mu} \partial^\nu X^\rho) A'^\sigma + h.c.) + h_3 A'_\mu \bar{f} \gamma^\mu f$$

b_5 : real

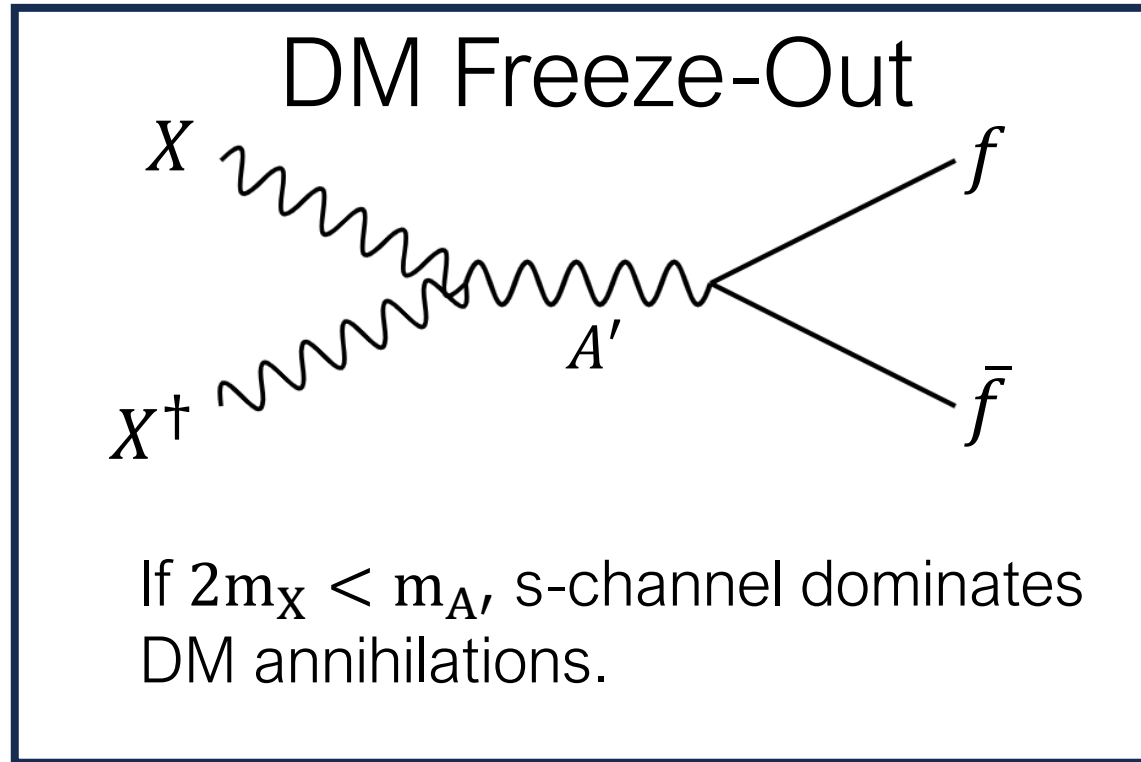
b_6 : complex

b_7 : complex

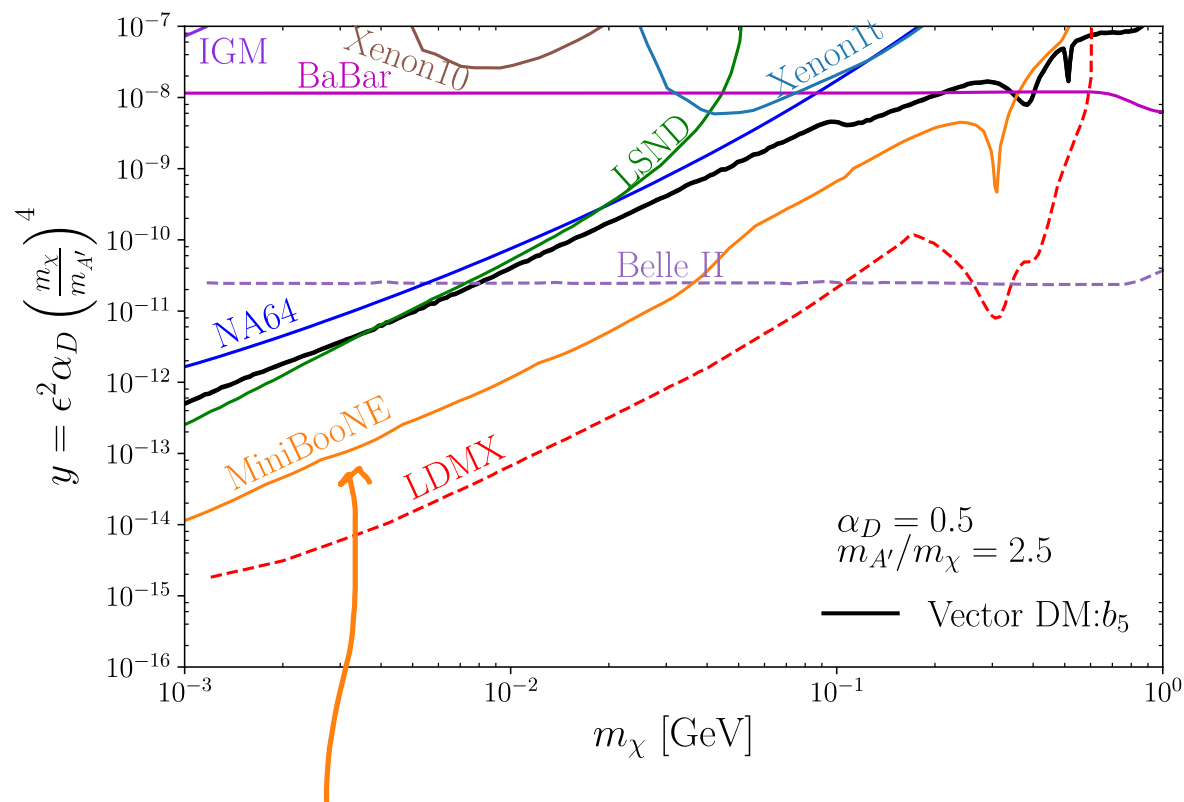
h_3 : real

$$h_3 \equiv \epsilon e$$

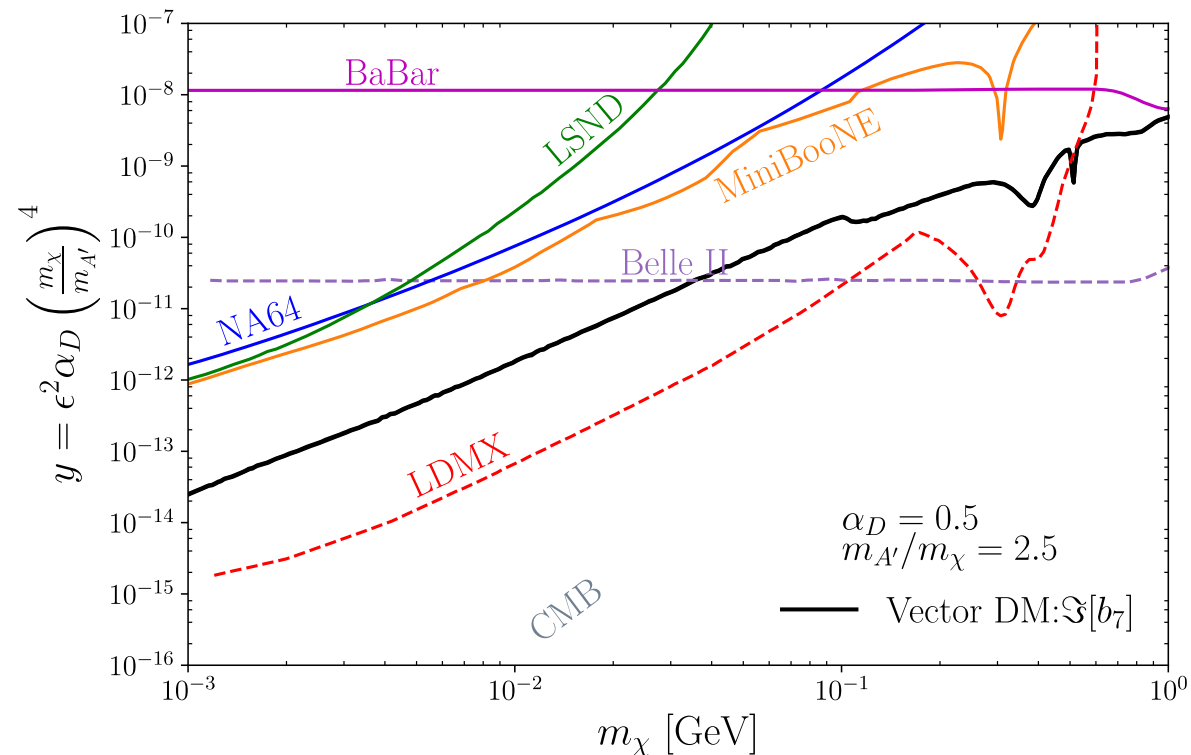
f : SM leptons and quarks
(excluding neutrinos)



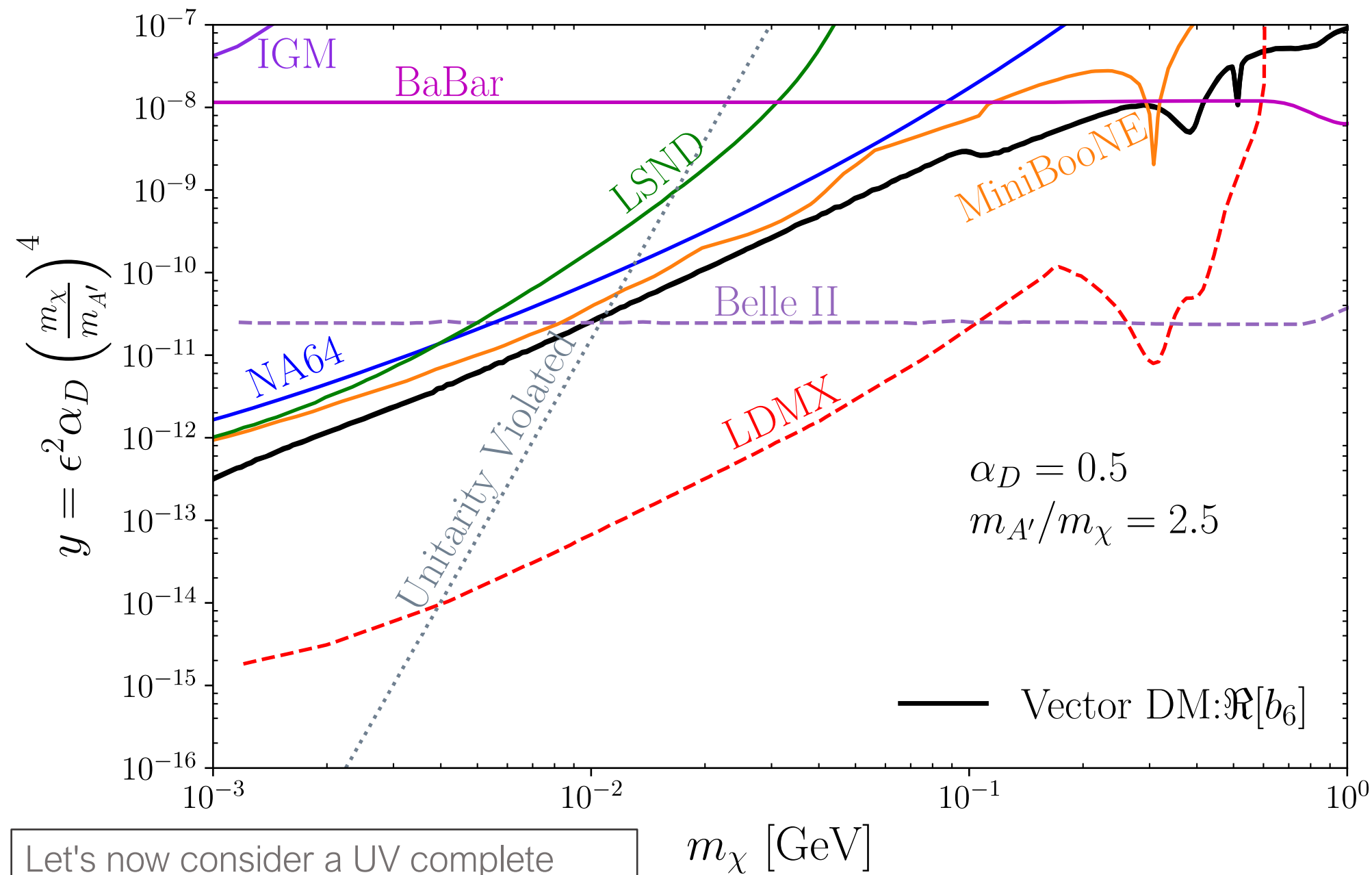
Strong Limits on Spin-1 Relic Targets



Large MiniBooNE limits
due to $\sigma_{\chi e^- \rightarrow \chi e^-}$ being
large when m_χ is small



Ruled out by current experiments and CMB



Spin-1 DM is
the first to be
probed by
LDMX!

Let's now consider a UV complete
model where unitarity is not violated!

$b_5 \rightarrow g_X$
 $\Im[b_6] \rightarrow -g_X/2$

SIMP Spin-1 Dark Matter

with Z' and X_3 as mediators

$SU_X(2) \times U_{Z'}(1)$

$$\begin{aligned}
 -\mathcal{L} \supset & -ig_X \cos \theta'_X \left[(\partial^\mu X^\nu - \partial^\nu X^\mu) X_\mu^\dagger \tilde{X}_{3,\nu} - (\partial^\mu X^{\nu\dagger} - \partial^\nu X^{\mu\dagger}) X_\mu \tilde{X}_{3,\nu} + X_\mu X_\nu^\dagger (\partial^\mu \tilde{X}_3^\nu - \partial^\nu \tilde{X}_3^\mu) \right] \\
 & -ig_X \sin \theta'_X \left[(\partial^\mu X^\nu - \partial^\nu X^\mu) X_\mu^\dagger \tilde{Z}'_\nu - (\partial^\mu X^{\nu\dagger} - \partial^\nu X^{\mu\dagger}) X_\mu \tilde{Z}'_\nu + X_\mu X_\nu^\dagger (\partial^\mu \tilde{Z}'^\nu - \partial^\nu \tilde{Z}'^\mu) \right] \\
 & -e\varepsilon \cos(\theta'_X) \tilde{Z}'_\mu \bar{f} \gamma^\mu f + e\varepsilon \sin(\theta'_X) \tilde{X}_{3\mu} \bar{f} \gamma^\mu f
 \end{aligned} \tag{1}$$

$\sin(\theta'_X) \ll \cos(\theta'_X)$

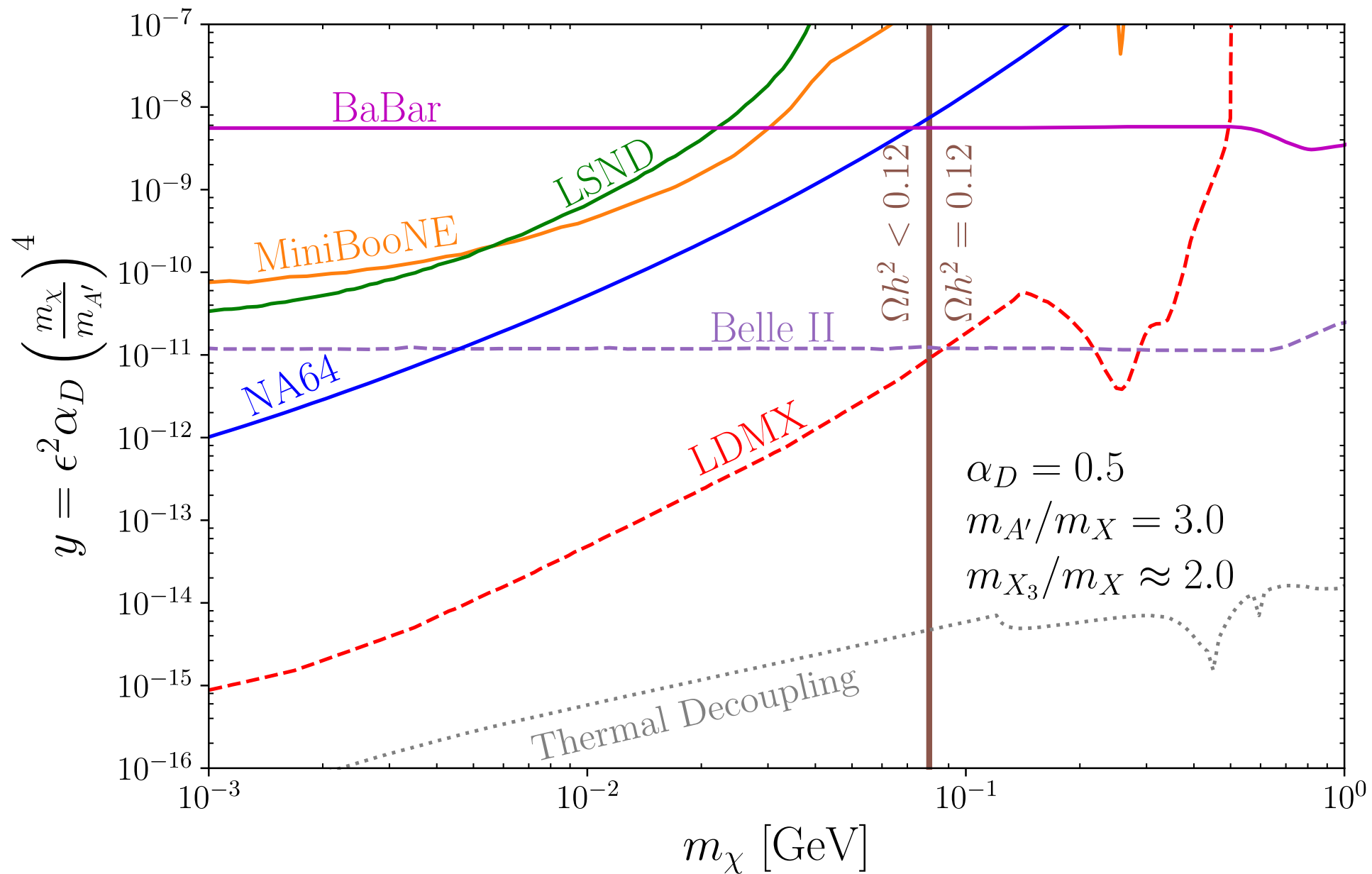
\rightarrow mainly Z' 's are
 produced at fixed
 target experiments

processes
 setting the
 relic density \rightarrow

$$\begin{aligned}
 & X_+ X_- \rightarrow f \bar{f} \\
 & X_+ X_+ X_- \rightarrow X_+ \tilde{X}_3 \\
 & X_+ X_- \rightarrow \tilde{X}_3 \tilde{X}_3
 \end{aligned}$$

Dominant since $g_X \gg e\varepsilon$

$$m_X^2 < m_{\tilde{X}_3}^2 < m_{\tilde{Z}'}^2$$



Weaker beam dump limits since DM – e^- scattering cross section is suppressed

Relic density is independent of ϵ , entire region consistent w Planck!

Summary

- Extending the current landscape of **sub-GeV DM** models considered in the context of **fixed target experiments**
- Spin-1 sub-GeV DM
 - where $m_{A'} > 2m_X$
- First model to be probed at upcoming LDMX!!

Future

- Consider $m_{A'} < 2m_X$
 - visible decays
- Additional UV complete spin-1 DM scenarios

Danke!

Backup Slides

SIMP spin-1 DM

$$SU_X(2) \times U_{Z'}(1)$$

$$\begin{aligned}
 -\mathcal{L} \supset & -ig_X \cos \theta'_X \left[(\partial^\mu X^\nu - \partial^\nu X^\mu) X_\mu^\dagger \tilde{X}_{3,\nu} - (\partial^\mu X^{\nu\dagger} - \partial^\nu X^{\mu\dagger}) X_\mu \tilde{X}_{3,\nu} + X_\mu X_\nu^\dagger (\partial^\mu \tilde{X}_3^\nu - \partial^\nu \tilde{X}_3^\mu) \right] \\
 & -ig_X \sin \theta'_X \left[(\partial^\mu X^\nu - \partial^\nu X^\mu) X_\mu^\dagger \tilde{Z}'_\nu - (\partial^\mu X^{\nu\dagger} - \partial^\nu X^{\mu\dagger}) X_\mu \tilde{Z}'_\nu + X_\mu X_\nu^\dagger (\partial^\mu \tilde{Z}'^\nu - \partial^\nu \tilde{Z}'^\mu) \right] \\
 & -e\varepsilon \cos(\theta'_X) \tilde{Z}'_\mu \bar{f} \gamma^\mu f + e\varepsilon \sin(\theta'_X) \tilde{X}_{3\mu} \bar{f} \gamma^\mu f
 \end{aligned} \tag{1}$$

- Dark spontaneous symmetry breaking by the VEVs of dark Higgs fields
- Dark Higgs Sector
 - Singlet scalar S
 - H_X
 - Kinetic mixing between Z' and hypercharge gauge bosons

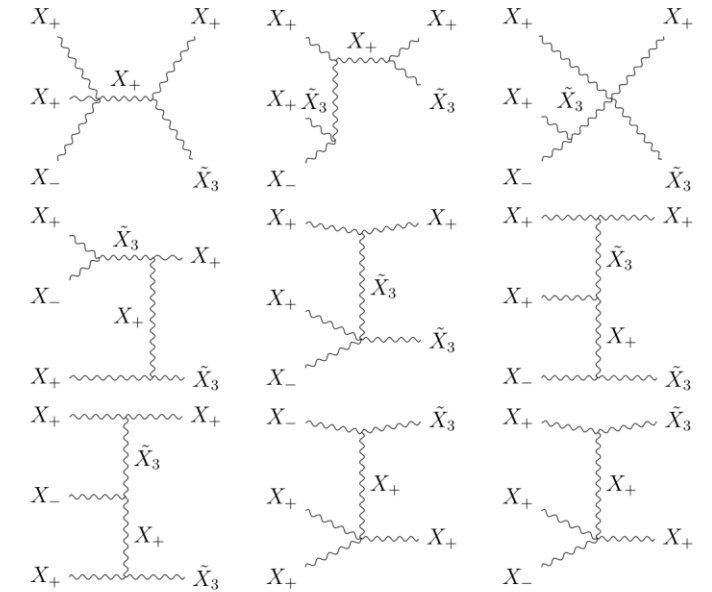
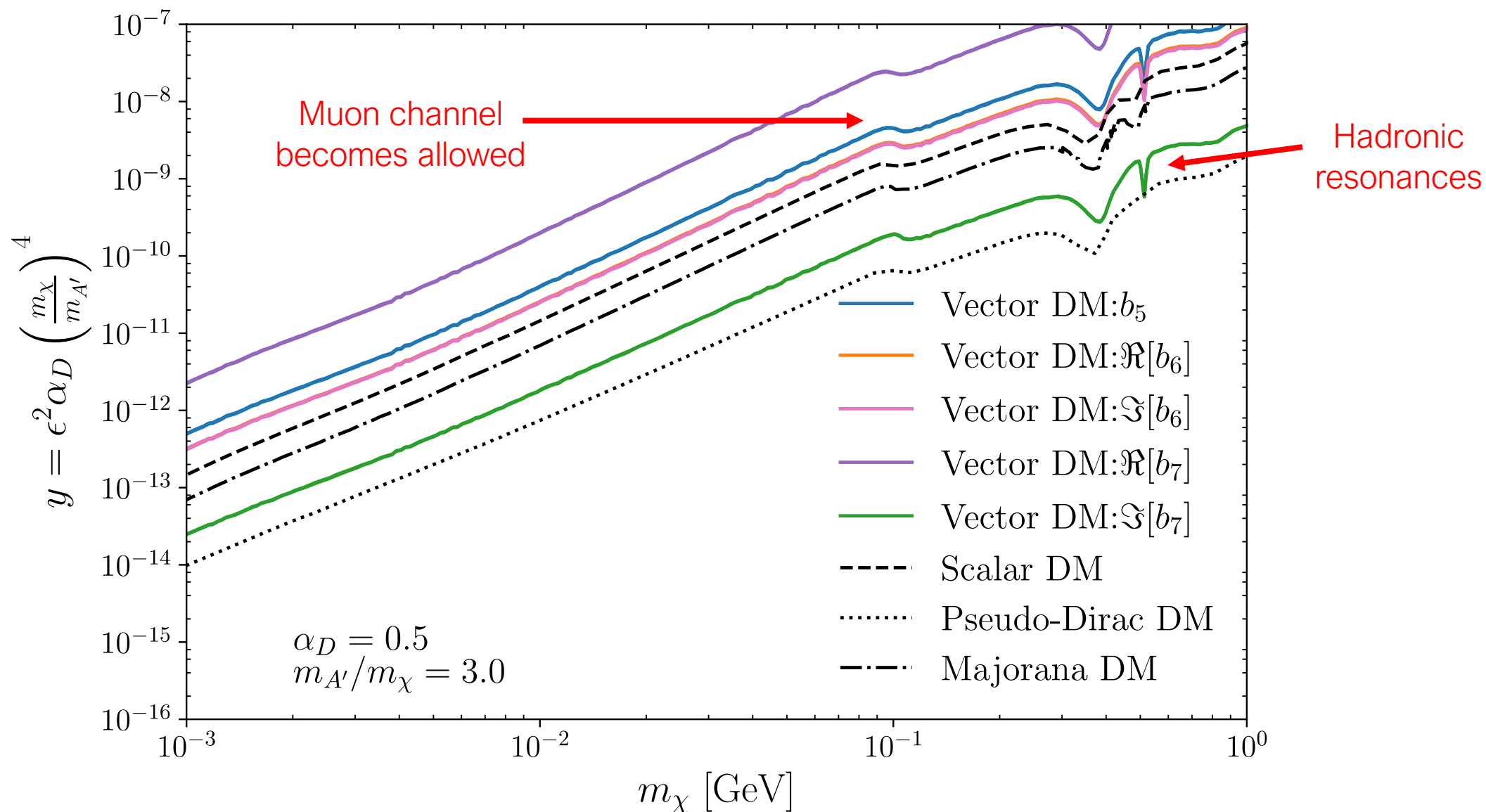


Figure 1: Feynman diagrams for $X_+ X_+ X_- \rightarrow X_+ \tilde{X}_3$.

Relic Targets of DM Models



Calculating Dark Matter Abundance

The Boltzmann Equation

$$\dot{n} + 3Hn = R$$

Universe's Expansion  Particle Physics 

- n : number density
- H : Hubble Rate (Universe's Expansion)
- R : Interaction Rate Density (# interactions per time and volume)
 - Includes all annihilations and productions
- More convenient to define Y and x
 - $Y \equiv \frac{n}{s}$, $x \equiv \frac{m}{T}$
 - s : entropy density

Ways of Producing Dark Matter

Γ : Interaction Rate (# interactions per time)

H : Hubble Rate (universe's expansion rate)

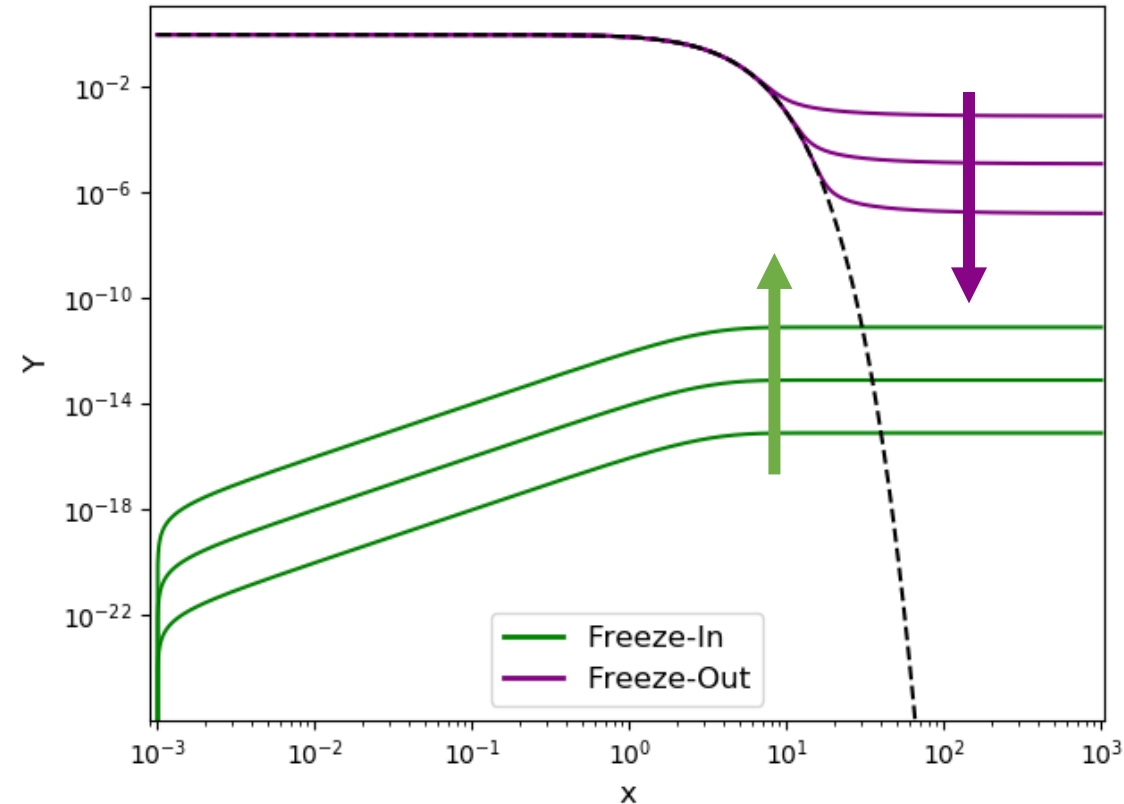
■ Freeze-In

- $\Gamma < H$ (decoupled)
- small interaction rates
- **never** thermalizes with bath

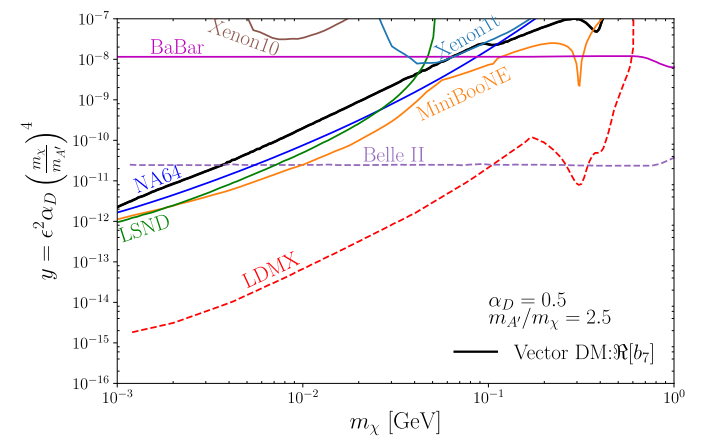
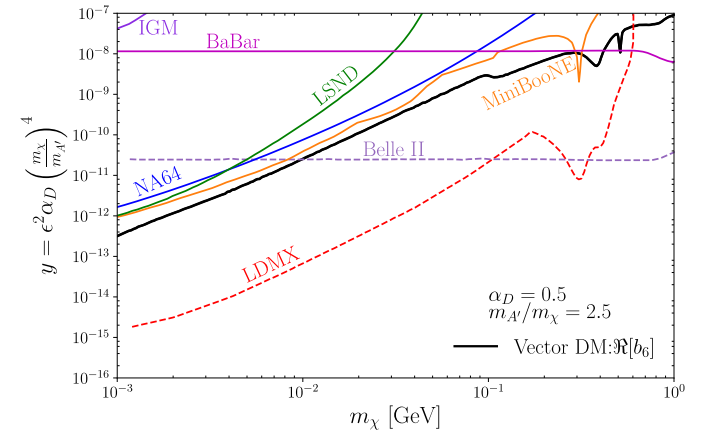
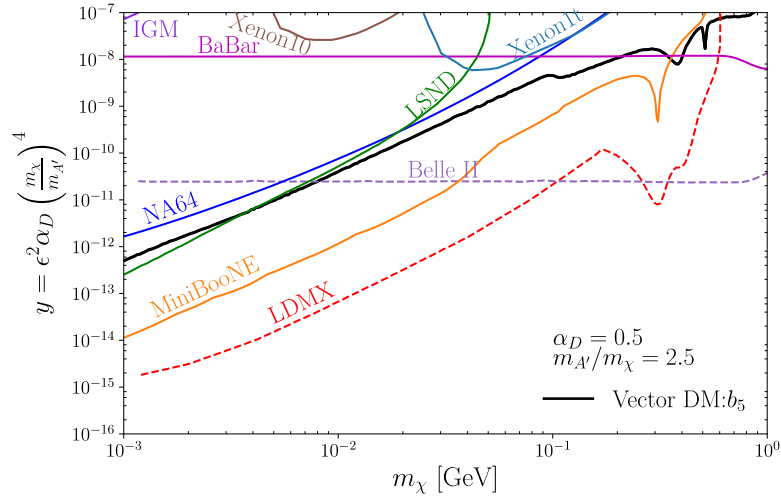
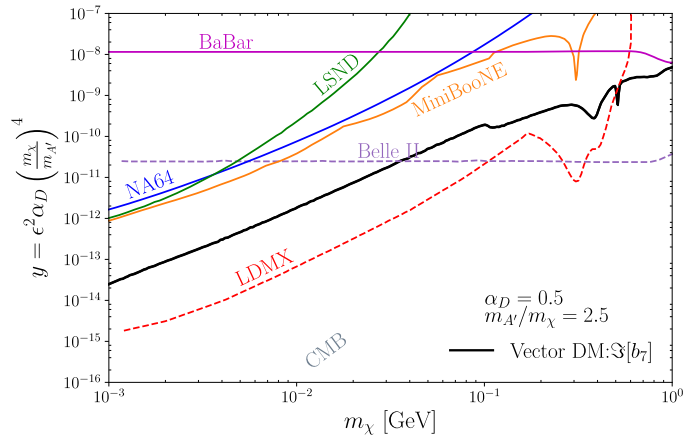
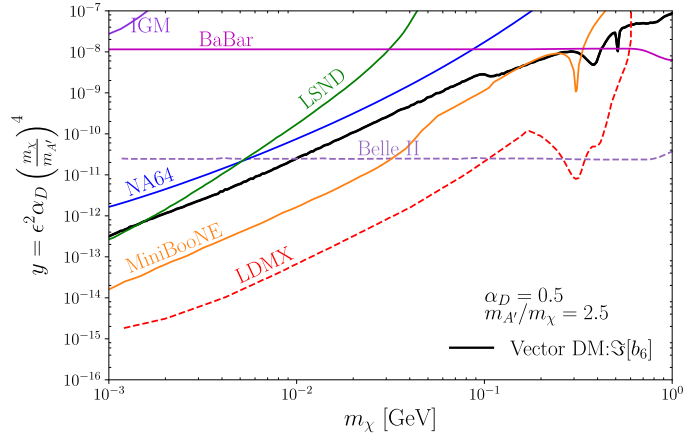
■ Freeze-Out

- $\Gamma > H$ (coupled)
- large interaction rates
- thermalizes with bath

This work!



→ Increasing Coupling



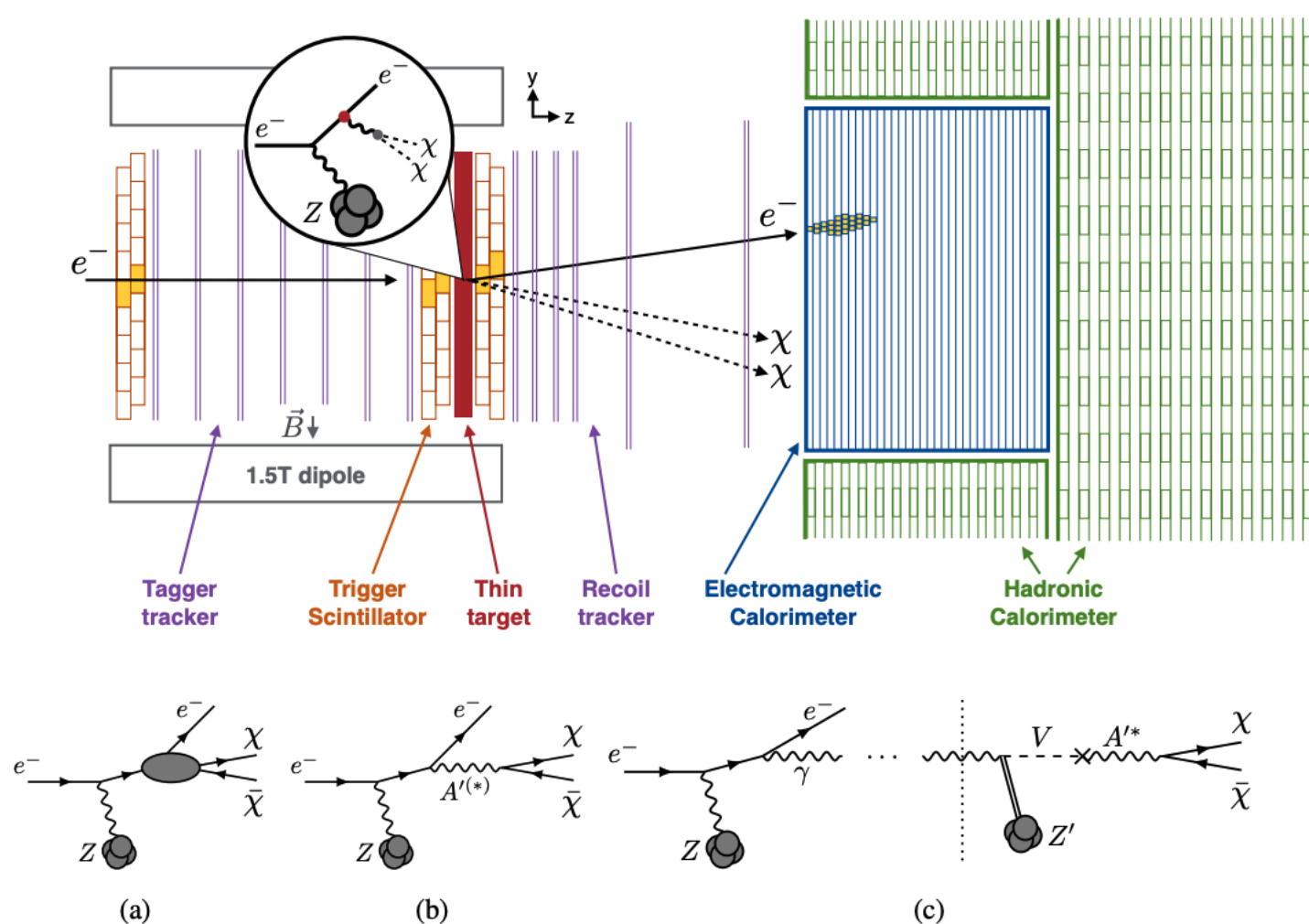
Confidence Intervals

- x% CL: If experiment is repeated many times, the intervals include the true parameter x% of the time
- Counting experiment, take Poisson distribution: $f(n; \nu) = \frac{\nu^n}{n!} e^{-\nu}$
- Uncertainty on number of background events
 - Neutrino flux, NCE cross section model ($\sigma_{\nu N \rightarrow \nu N}$), detector response
 - Nuisance parameters introduced

Experiments

Light Dark Matter eXperiment (LDMX)

- Future fixed target missing momentum exp
 - 2025: LESA delivers beam to LDMX allowing 4×10^{14} EOT
 - 2027: 10^{16} EOT
- e^- incident on a thin tungsten target
- Charged particle tracker and calorimeters to measure DM signature
 - Recoil electron pT accompanied by absence of other particle activity

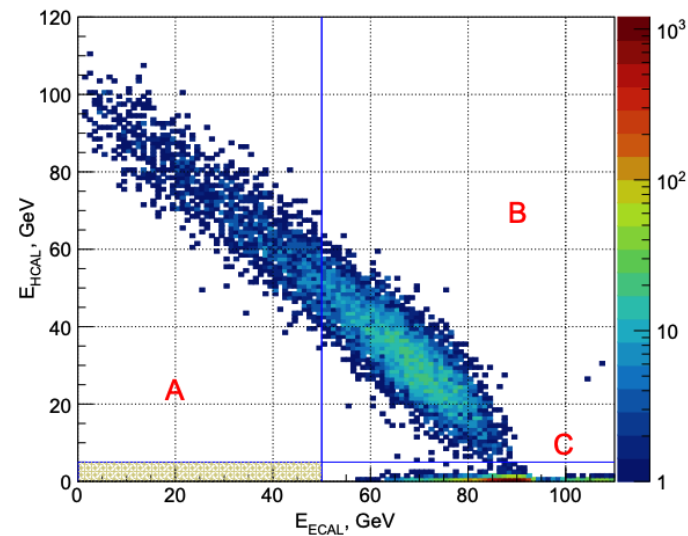
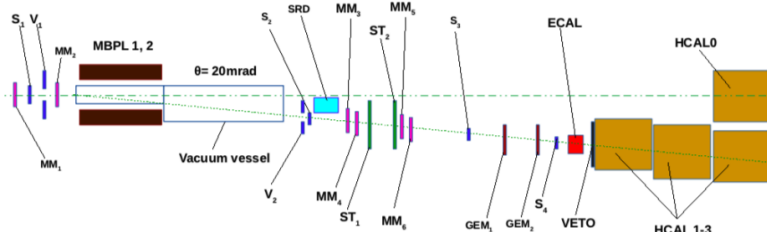


Electron Beam Dumps

NA64

arXiv:1710.00971

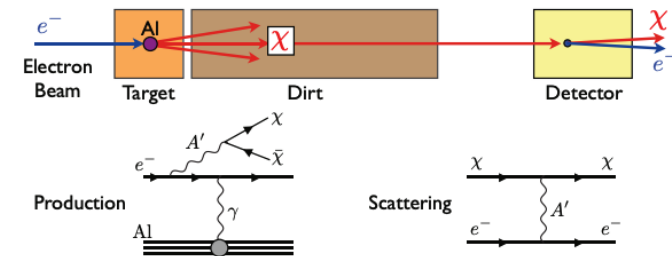
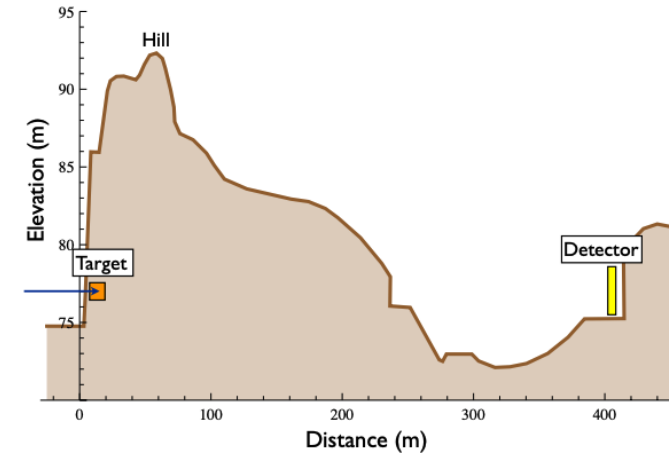
- 100 GeV electron beam incident on a lead target
- Event: single electron produced and missing energy



E137

arXiv:1406.2698

- DM produced from electron-target collisions
- 20 GeV beam incident on a set of aluminum plates interlaced with cooling water.
- Downstream detector



Proton Beam Dumps

arXiv:1107.4580

DM scatterings mimic neutrino scatterings!
(Neutral current-like scatterings)

LSND

arXiv:hep-ex/0101039

- pions produced by impacting an 800 MeV proton beam onto a water or metal target
- $\pi^0 \rightarrow A' \gamma, A' \rightarrow XX$

Mini-Boone

arXiv:1807.06137

- Designed to study short-baseline neutrino oscillations
- 8 GeV proton beam incident on a steel target
- Peak ~ 800 MeV (ρ mass)

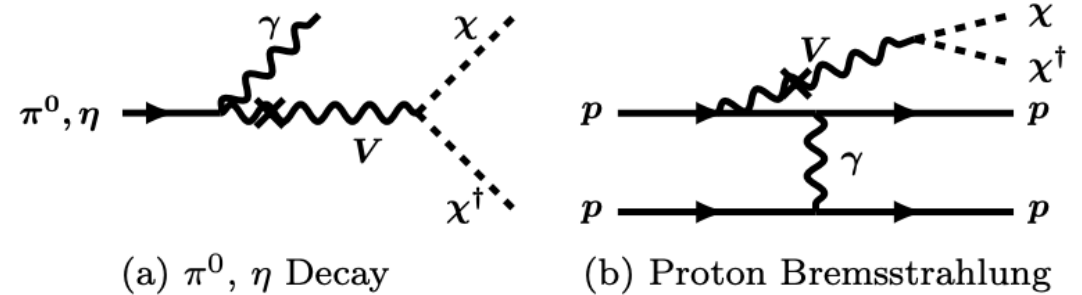
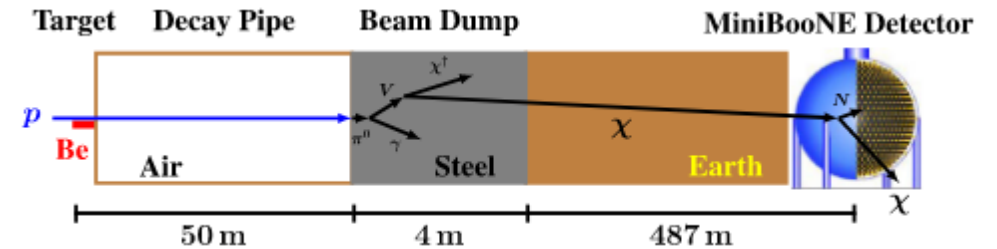


FIG. 2. DM production channels relevant for this search with an 8 GeV proton beam incident on a steel target.

Monophoton Searches $e^+e^- \rightarrow \gamma A', A' \rightarrow XX$

Search for single photon events in e^+e^- collision data



arXiv:1702.03327

- BABAR detector at PEP-II B-factory
- Large missing energy/momentum
- Exclusions for $m_{A'} \leq 8$ GeV

Belle-II

arXiv:1808.10567

- Experiment operated at SuperKEKB
- First data taken in 2019, more to come..
- 7 GeV electrons with 4 GeV positrons

Hadronic Resonances

- If DM freezes-out after the QCD phase transition (~ 150 MeV), DM annihilates to hadronic final states rather than to quarks.
 - Must consider for $m_X \lesssim 3 \text{ GeV}$

$$\sigma v_{XX \rightarrow A' \rightarrow \text{hadrons}} \approx R(s) \sigma v_{XX \rightarrow A' \rightarrow \mu^- \mu^+}$$

$$R(s) \equiv \sigma_{e^+ e^- \rightarrow \text{hadrons}} / \sigma_{e^+ e^- \rightarrow \mu^+ \mu^-}$$