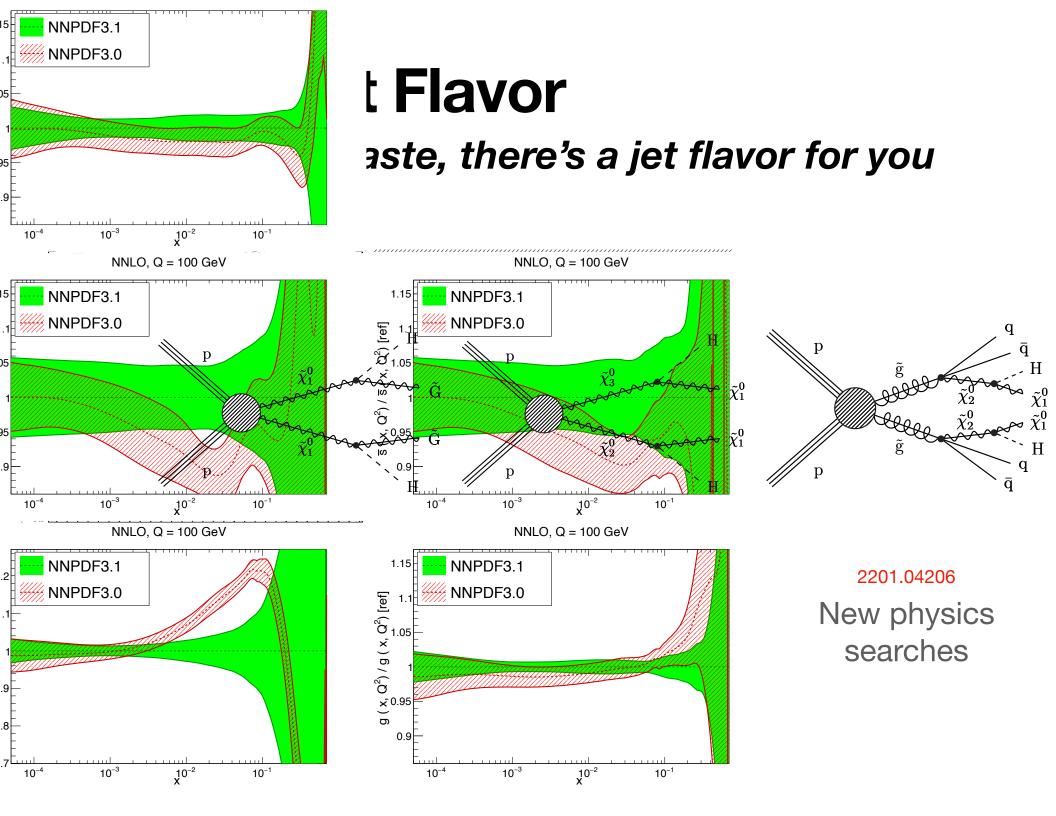
## A Fragmentation Approach to Jet Flavor Andrew Larkoski

with Simone Caletti, Simone Marzani, Daniel Reichelt 2205.01117

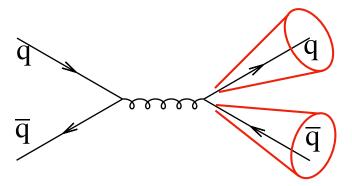
UCLA



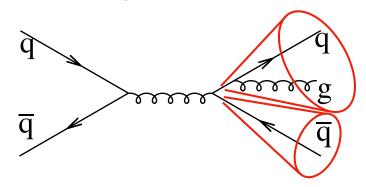
### **BSZ Flavor**

Come on, just sum the flavors of partons in the jet

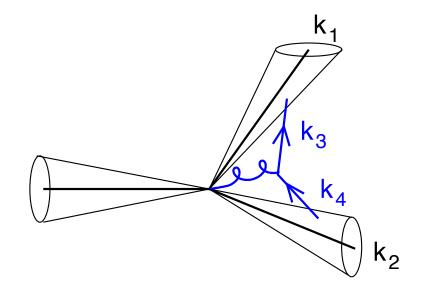
Unambiguous at LO



Unambiguous at NLO



Soft gluon splitting to quarks at NNLO spoils IR safety



## **BSZ Flavor**

#### **Oh.** What if you changed the clustering algorithm?

$$y_{ij}^{(F)} = \frac{2(1 - \cos \theta_{ij})}{Q^2} \times \begin{cases} \max(E_i^2, E_j^2), \\ \min(E_i^2, E_j^2), \end{cases}$$

softer of i, j is flavoured, softer of i, j is flavourless

Modification to Durham k<sub>T</sub> algorithm

Phys. Lett. B 269, 432 (1991)

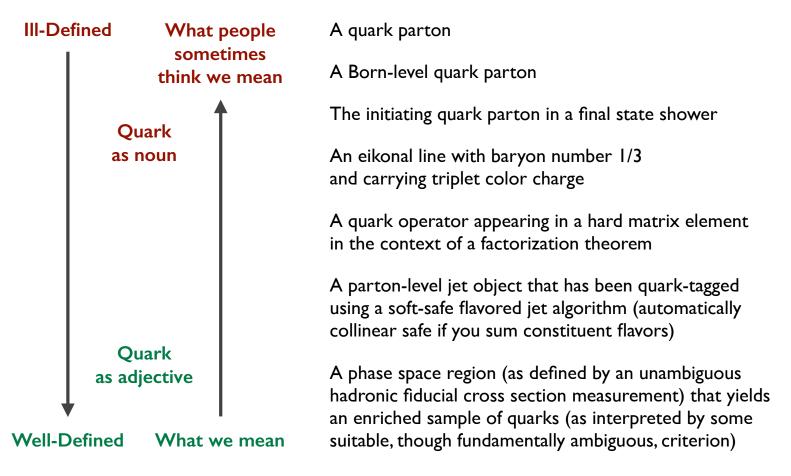
Only cluster soft quarks together if they have equal and opposite flavor

- Pros
  - It works and is IRC safe
  - Can be easily implemented in partonic fixed order calculation
- Cons
  - Modifies constituents of jets and no one uses k<sub>T</sub> to find jets anymore anti-k<sub>T</sub> flavor algorithm, 2205.11879
  - Complete non-starter in experiment; cannot begin to identify flavored jets

#### Les Houches Study

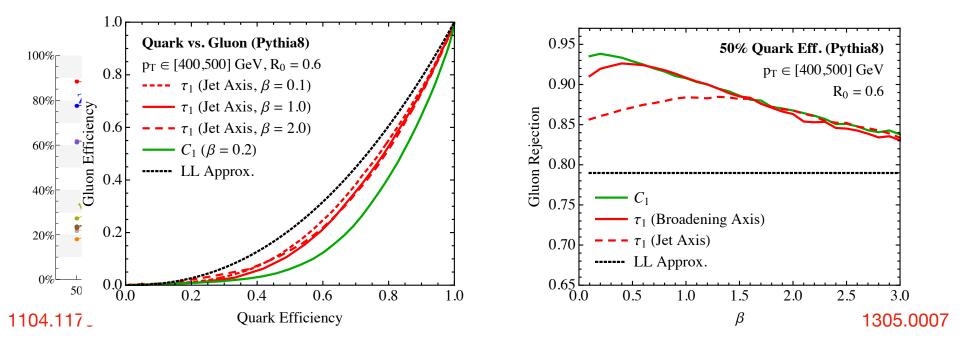
What is a Quark Jet?

From lunch/dinner discussions



#### Focus on in-principle observables

 Many, many results in the literature about "quark vs. gluon discrimination"



De facto but flawed definition:

jet flavor defined by requested process in event simulation

on where ALL Jets are Quark

<sup>6</sup> Fraction where ALL Jets are Gluon

#### Focus on in-principle observables

- Many, many results in the literature about "quark vs. gluon discrimination"
- By asymptotic freedom, jet flavor is unambiguous in deep UV

User requests short distance process in event simulator

#### Focus on in-principle observables

- Many, many results in the literature about "quark vs. gluon discrimination"
- By asymptotic freedom, jet flavor is unambiguous in deep UV
- Measurements are performed in the deep IR

Flow to IR governed by renormalization group/DGLAP

#### Focus on in-principle observables

- Many, many results in the literature about "quark vs. gluon discrimination"
- By asymptotic freedom, jet flavor is unambiguous in deep UV
- Measurements are performed in the deep IR
- Flow from UV to IR is not invertible

Jet boundary destroys one-to-one UV to IR map

#### Focus on in-principle observables

- Many, many results in the literature about "quark vs. gluon discrimination"
- By asymptotic freedom, jet flavor is unambiguous in deep UV
- Measurements are performed in the deep IR
- Flow from UV to IR is not invertible
- Give up on trying to get UV jet flavor; just focus on IR

• Only returns a QCD parton flavor (up, down, strange,...,gluon)

Simplifies classification

- Only returns a QCD parton flavor (up, down, strange,...,gluon)
- Can be applied to any set of partons

Does not require re-associating constituents of a jet

- Only returns a QCD parton flavor (up, down, strange,...,gluon)
- Can be applied to any set of partons
- IR safe, completely insensitive to soft particles

Ignores contribution to jet flavor from soft particles

- Only returns a QCD parton flavor (up, down, strange,...,gluon)
- Can be applied to any set of partons
- IR safe, completely insensitive to soft particles
- Inclusive over exactly collinear splittings

Absorb collinear divergences into fragmentation functions

- Only returns a QCD parton flavor (up, down, strange,...,gluon)
- Can be applied to any set of partons
- IR safe, completely insensitive to soft particles
- Inclusive over exactly collinear splittings
- Described by linear evolution equations

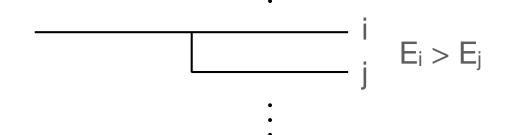
Use DGLAP as a guide and enables analytic solutions

The partonic flavor of a jet is defined to be the net flavor of the particle(s) whose momentum lies exactly along the WTA axis of the jet

Find closest pair of particles according to clustering metric d<sub>ij</sub>

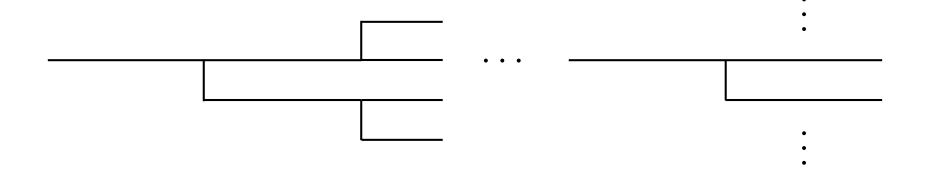
The partonic flavor of a jet is defined to be the net flavor of the particle(s) whose momentum lies exactly along the WTA axis of the jet

Sum energies and new momentum points along direction of harder particle



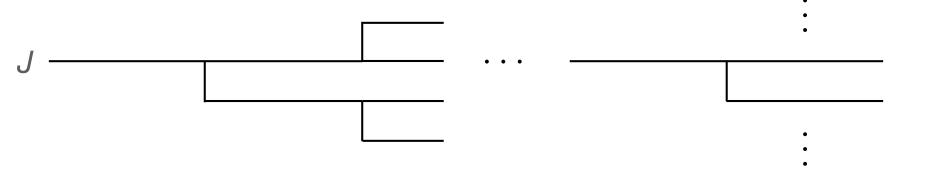
The partonic flavor of a jet is defined to be the net flavor of the particle(s) whose momentum lies exactly along the WTA axis of the jet

Continue pairwise combination until one particle remains



The partonic flavor of a jet is defined to be the net flavor of the particle(s) whose momentum lies exactly along the WTA axis of the jet

Final momentum lies along direction of particle in jet



## **Properties of Evolution Equations**

Ain't they beautiful?

$$Q^2 \frac{df_q(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ -C_F\left(2\log 2 - \frac{5}{8}\right) f_q(Q^2) + \frac{1}{3}T_R f_g(Q^2) \right]$$

$$Q^2 \frac{df_g(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ C_F\left(2\log 2 - \frac{5}{8}\right) - \left(C_F\left(2\log 2 - \frac{5}{8}\right) + \frac{2}{3}n_f T_R\right) f_g(Q^2) \right]$$

• Linear, inhomogeneous evolution equations

Vastly simpler than nonlinear evolution for hardest subjet, NGLs, track functions,...<sup>1411.5182</sup> hep-ph/0206076 1209.3019

## **Properties of Evolution Equations**

Ain't they beautiful?

$$Q^2 \frac{df_q(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ -C_F\left(2\log 2 - \frac{5}{8}\right) f_q(Q^2) + \frac{1}{3}T_R f_g(Q^2) \right]$$

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- Linear, inhomogeneous evolution equations
- Independent of the color of the gluon

Gluon emissions of gluons can't change WTA flavor

## **Properties of Evolution Equations**

Ain't they beautiful?

$$Q^2 \frac{df_q(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ -C_F\left(2\log 2 - \frac{5}{8}\right) f_q(Q^2) + \frac{1}{3}T_R f_g(Q^2) \right]$$

$$Q^2 \frac{df_g(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ C_F\left(2\log 2 - \frac{5}{8}\right) - \left(C_F\left(2\log 2 - \frac{5}{8}\right) + \frac{2}{3}n_f T_R\right) f_g(Q^2) \right]$$

- Linear, inhomogeneous evolution equations
- Independent of the color of the gluon
- Deep IR fixed points of jet flavor:

$$\lim_{Q_0^2 \to \infty} f_q(Q^2) = \frac{\frac{1}{3}T_R}{C_F \left(2\log 2 - \frac{5}{8}\right) + \frac{2}{3}n_f T_R} \approx 0.062149$$

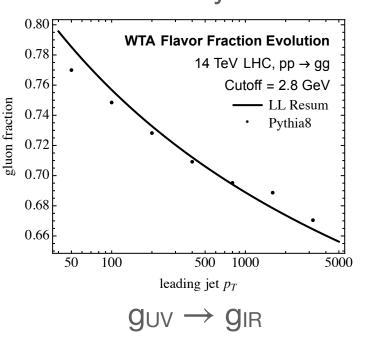
$$QCD \text{ with } n_f = 5$$

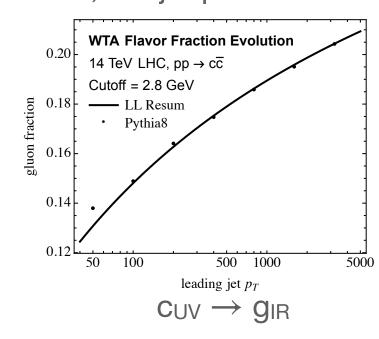
$$\lim_{Q_0^2 \to \infty} f_g(Q^2) = \frac{C_F \left(2\log 2 - \frac{5}{8}\right)}{C_F \left(2\log 2 - \frac{5}{8}\right) + \frac{2}{3}n_f T_R} \approx 0.37851$$

## **Comparison to Parton Shower**

#### Evolution of WTA Gluon Flavor from Jet p<sub>T</sub> to IR

 $pp \rightarrow cc \text{ or } pp \rightarrow gg \text{ events generated in Pythia8}$ No hadronization; parton shower terminated at fixed IR scale Study evolution as a function of the UV scale; the jet  $p_T$ 

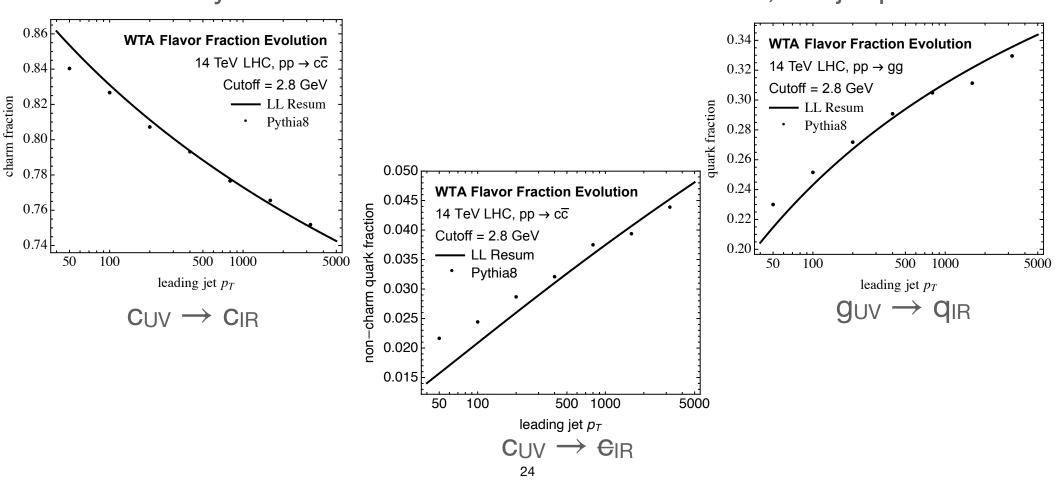




## **Comparison to Parton Shower**

#### **Evolution of WTA Quark Flavor from Jet p<sub>T</sub> to IR**

 $pp \rightarrow cc \text{ or } pp \rightarrow gg \text{ events generated in Pythia8}$ No hadronization; parton shower terminated at fixed IR scale Study evolution as a function of the UV scale; the jet  $p_T$ 



## Summary

- Collinear divergences aren't that scary!
- Let fragmentation functions be your friend
- Other things:
  - Can WTA flavor be embedded in a factorization theorem?
  - FO matches to pdfs all the time; can they match to WTA flavor?
  - WTA axis is extremely robust to contamination
  - Observables measured about WTA axis may connect partonic flavor to realistic hadronic jets

### Bonus

#### **Derivation of Evolution Equations** Just DGLAP + Hardest Energy Constraint

$$Q^2 \frac{df_q(x,Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \int_x^{\min[1,2x]} \frac{dz}{z} \left[ P_{qg\leftarrow q} \left(\frac{x}{z}\right) f_q(z,Q^2) + P_{q\bar{q}\leftarrow g} \left(\frac{x}{z}\right) f_g(z,Q^2) \right]$$

#### **Derivation of Evolution Equations** Just DGLAP + Hardest Energy Constraint

$$Q^{2} \frac{df_{q}(x,Q^{2})}{dQ^{2}} = \frac{\alpha_{s}}{2\pi} \int_{x}^{\min[1,2x]} \frac{dz}{z} \left[ P_{qg \leftarrow q} \left( \frac{x}{z} \right) f_{q}(z,Q^{2}) + P_{q\bar{q}\leftarrow g} \left( \frac{x}{z} \right) f_{g}(z,Q^{2}) \right]$$

$$z = \left[ \begin{array}{c} z - x \\ z \\ z \end{array} \right] \left[ \begin{array}{c} z - x \\ z \\ z \end{array} \right] \left[ \begin{array}{c} z - x \\ z \\ z \end{array} \right] \left[ \begin{array}{c} z - x \\ z \\ z \end{array} \right] \left[ \begin{array}{c} z \\ z \end{array} \right] \left[ \begin{array}{c} z \\ z \\ z \end{array} \right] \left[ \begin{array}{c} z$$

#### **Derivation of Evolution Equations** Integrate Over Energy Fractions

$$Q^{2} \frac{df_{q}(x,Q^{2})}{dQ^{2}} = \frac{\alpha_{s}}{2\pi} \int_{x}^{\min[1,2x]} \frac{dz}{z} \left[ P_{qg \leftarrow q} \left(\frac{x}{z}\right) f_{q}(z,Q^{2}) + P_{q\bar{q} \leftarrow g} \left(\frac{x}{z}\right) f_{g}(z,Q^{2}) \right]$$
$$\int_{0}^{1} dx f_{i}(x,Q^{2}) = f_{i}(Q^{2})$$

Fraction of jets with parton *i* along WTA axis

$$Q^2 \frac{df_q(Q^2)}{dQ^2} = \int_{1/2}^1 dy \left[ P_{qg \leftarrow q}(y) f_q(Q^2) + P_{q\bar{q} \leftarrow g}(y) f_g(Q^2) \right]$$
  
Just need reduced moments of splitting functions

#### **Derivation of Evolution Equations** *Final WTA Fraction Evolution Equations*

$$Q^{2} \frac{df_{q}(Q^{2})}{dQ^{2}} = \frac{\alpha_{s}}{2\pi} \left[ -C_{F} \left( 2\log 2 - \frac{5}{8} \right) f_{q}(Q^{2}) + \frac{1}{3} T_{R} f_{g}(Q^{2}) \right]$$

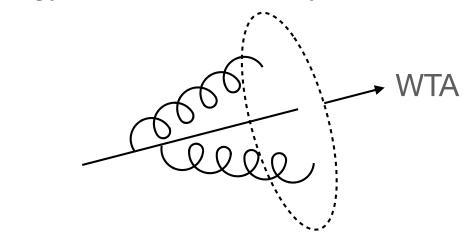
Quark flavor q evolution equation

$$Q^2 \frac{df_g(Q^2)}{dQ^2} = \frac{\alpha_s}{2\pi} \left[ C_F\left(2\log 2 - \frac{5}{8}\right) - \left(C_F\left(2\log 2 - \frac{5}{8}\right) + \frac{2}{3}n_f T_R\right) f_g(Q^2) \right]$$

Gluon flavor evolution equation

#### **Measuring WTA Axis**

Energy fraction carried by WTA axis



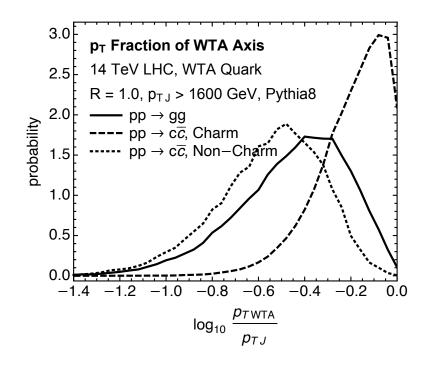
$$Q^{2} \frac{df_{q_{i}}(x,Q^{2})}{dQ^{2}} = \frac{\alpha_{s}}{2\pi} \int_{x}^{\min[1,2x]} \frac{dz}{z} \left[ P_{qg \leftarrow q} \left(\frac{x}{z}\right) f_{q_{i}}(z,Q^{2}) + P_{q\bar{q} \leftarrow g} \left(\frac{x}{z}\right) f_{g}(z,Q^{2}) \right]$$

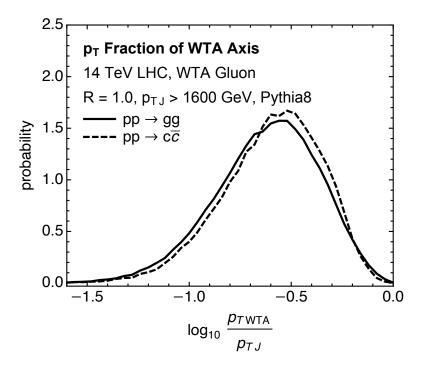
$$Q^{2} \frac{df_{g}(x,Q^{2})}{dQ^{2}} = \frac{\alpha_{s}}{2\pi} \int_{x}^{\min[1,2x]} \frac{dz}{z} \left[ P_{gq \leftarrow q} \left(\frac{x}{z}\right) \sum_{i=1}^{n_{f}} \left( f_{q_{i}}(z,Q^{2}) + f_{\bar{q}_{i}}(z,Q^{2}) \right) + P_{gg \leftarrow g} \left(\frac{x}{z}\right) f_{g}(z,Q^{2}) \right]$$

Solve 2n<sub>f</sub>+1 coupled integro-differential equations Just as complicated as DGLAP; solve in conjugate space and invert

#### Measuring WTA Axis

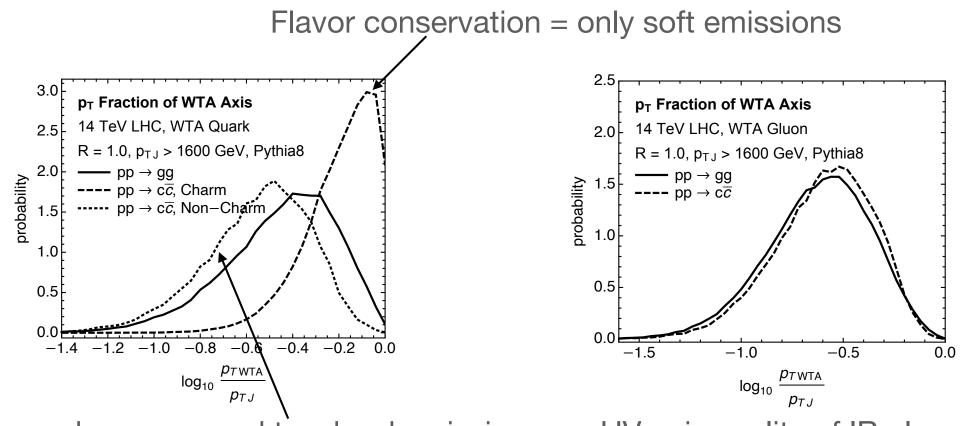
Simplest Solution: Parton Shower Monte Carlo





#### Measuring WTA Axis

Simplest Solution: Parton Shower Monte Carlo



Flavor change = need two hard emissions

UV universality of IR gluons

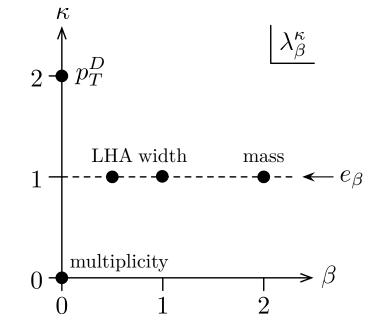
#### Measuring about WTA Axis

**Generalized Angularities** 

$$\lambda_{\beta}^{\kappa} = \sum_{i \in J} z_i^{\kappa} \theta_i^{\beta}$$

Flavor-Inclusive Distribution

$$p(\lambda_{\beta}) = \frac{d\sigma^{\text{fo}}}{d\lambda_{\beta}} \,\Delta_{\text{Sud}}(\lambda_{\beta})$$



Ansatz Flavor-Sensitive Distribution

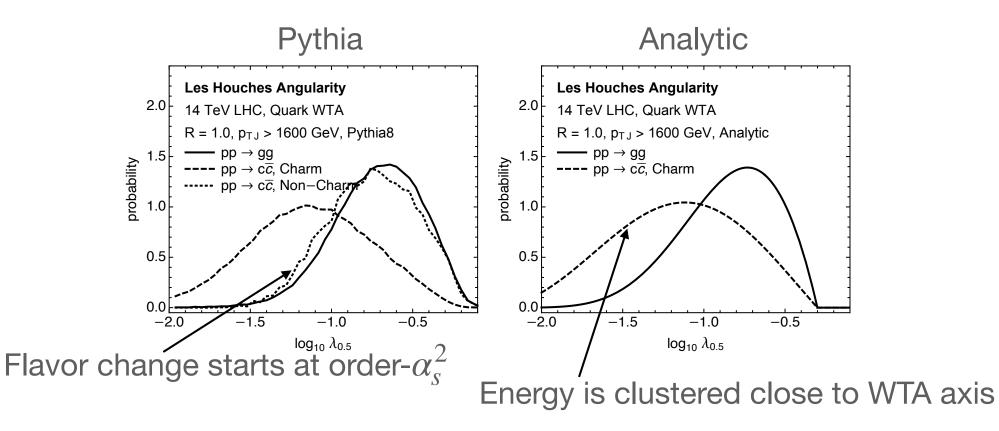
$$p_i(\lambda_\beta) = \frac{d\sigma_i^{\rm fo}}{d\lambda_\beta} \,\Delta_{\rm Sud}(\lambda_\beta)$$

Only resolves total color of collinear region WTA

1704.03878

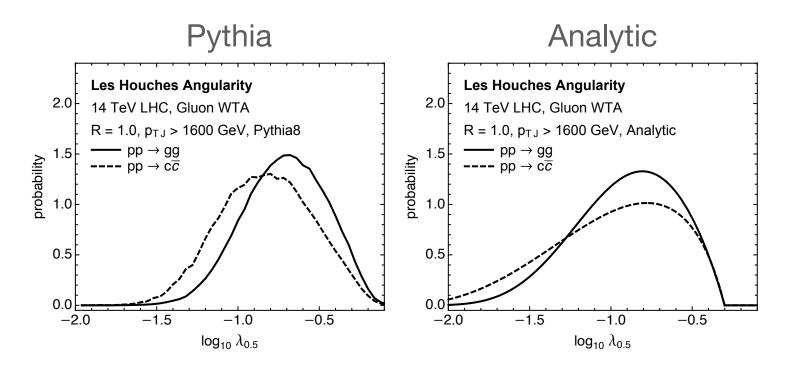
#### Measuring about WTA Axis

WTA Flavor Quark Jets



#### Measuring about WTA Axis

WTA Flavor Gluon Jets



Good qualitative agreement

Can this be systematic? Is there a factorization that enables calculation?