

Frame-independent methods to access GPDs from lattice QCD

Martha Constantinou



Temple University

In collaboration with:

**S. Bhattacharya, K. Cichy, J. Dodson, X. Gao, A. Metz,
J. Miller, A. Scapellato, F. Steffens, S. Mukherjee, Y. Zhao**

DIS 2023

March 28, 2023



Motivation for GPDs studies

- ★ Crucial in understanding hadron tomography
- ★ Correlation between transverse position and longitudinal momentum of the quarks in the hadron and its mechanical properties

Motivation for GPDs studies

- ★ Crucial in understanding hadron tomography
- ★ Correlation between transverse position and longitudinal momentum of the quarks in the hadron and its mechanical properties

This talk:

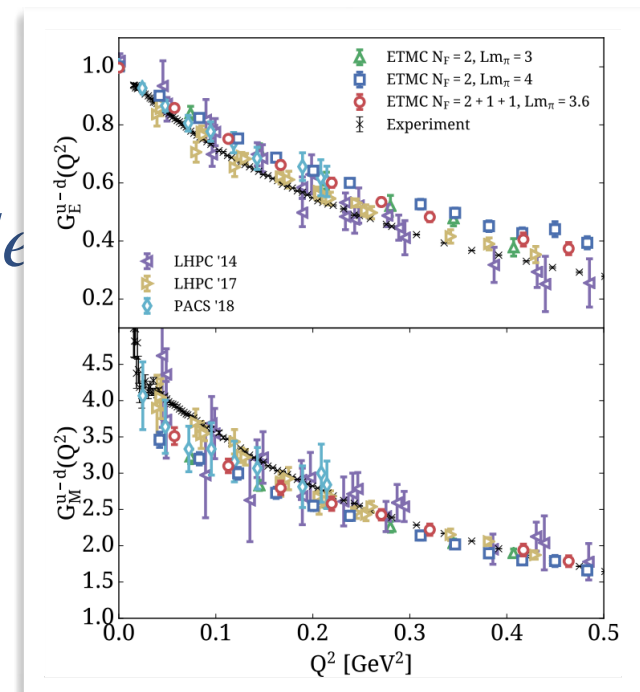
- ★ Prescription of how to access GPDs from *first principles* (lattice QCD):
 - with realistic computational resources
 - for a broad range of their variables
 - at fast convergence to light-cone GPDs

Motivation for GPDs studies

- ★ Crucial in understanding hadron tomography
- ★ Correlation between transverse position and longitudinal momentum of the quarks in the hadron and its mechanical properties

This talk:

- ★ Prescription of how to access GPDs from *first principles*
 - with realistic computational resources
 - for a broad range of their variables
 - at fast convergence to light-cone GPDs



Motivation for GPDs studies

- ★ Crucial in understanding hadron tomography
- ★ Correlation between transverse position and longitudinal momentum of the quarks in the hadron and its mechanical properties

This talk:

- ★ Prescription of how to access GPDs from *first principles* (lattice QCD):
 - with realistic computational resources
 - for a broad range of their variables
 - at fast convergence to light-cone GPDs

Motivation for GPDs studies

- ★ Crucial in understanding hadron tomography
- ★ Correlation between transverse position and longitudinal momentum of the quarks in the hadron and its mechanical properties

This talk:

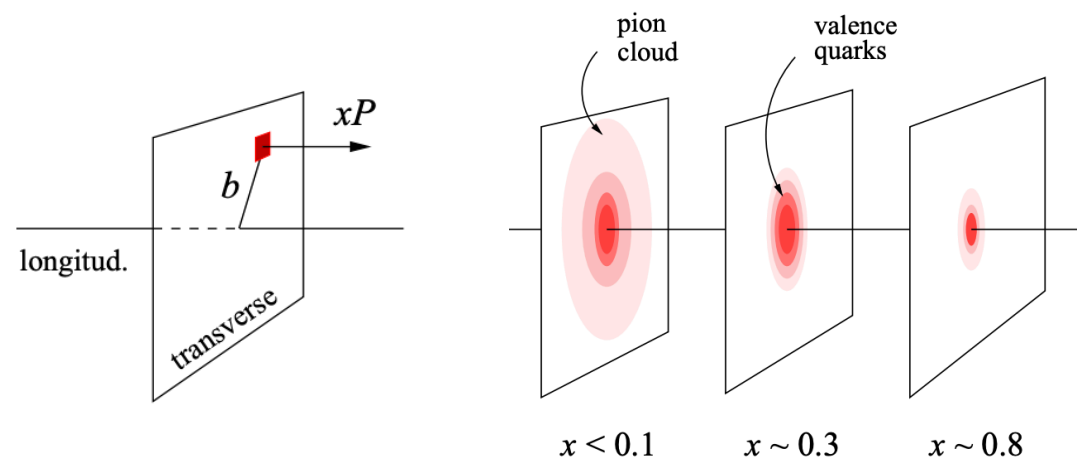
- ★ Prescription of how to access GPDs from *first principles* (lattice QCD):
 - with realistic computational resources
 - for a broad range of their variables
 - at fast convergence to light-cone GPDs

PHYSICAL REVIEW D **106**, 114512 (2022)

Generalized parton distributions from lattice QCD with asymmetric momentum transfer: Unpolarized quarks

Shohini Bhattacharya^{1,*} Krzysztof Cichy,² Martha Constantinou^{3,†} Jack Dodson,³ Xiang Gao,⁴ Andreas Metz,³
Swagato Mukherjee¹ Aurora Scapellato,³ Fernanda Steffens,⁵ and Yong Zhao⁴

Generalized Parton Distributions

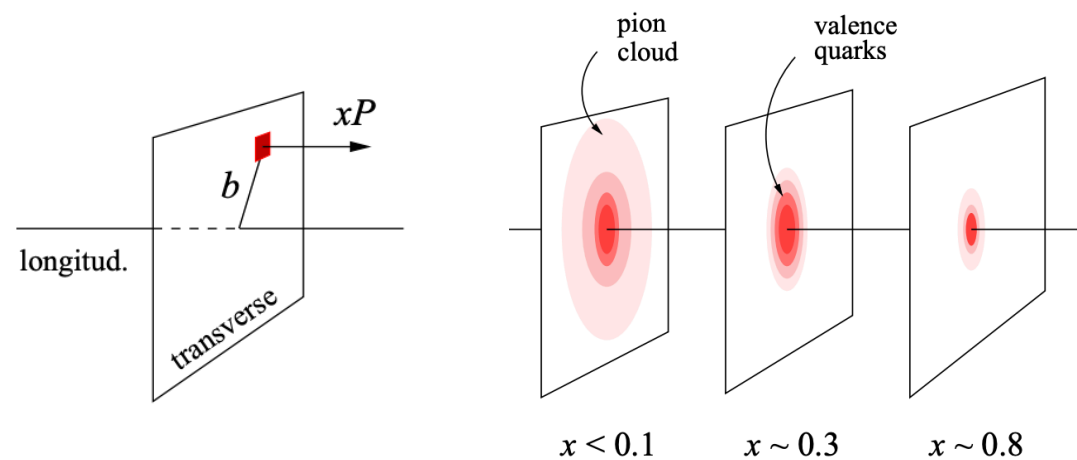


[H. Abramowicz et al., whitepaper for NSAC LRP, 2007]

$1_{\text{mom}} + 2_{\text{coord}}$ tomographic images of quark distribution in nucleon at fixed longitudinal momentum

3-D image from FT of the longitudinal mom. transfer

Generalized Parton Distributions



[H. Abramowicz et al., whitepaper for NSAC LRP, 2007]

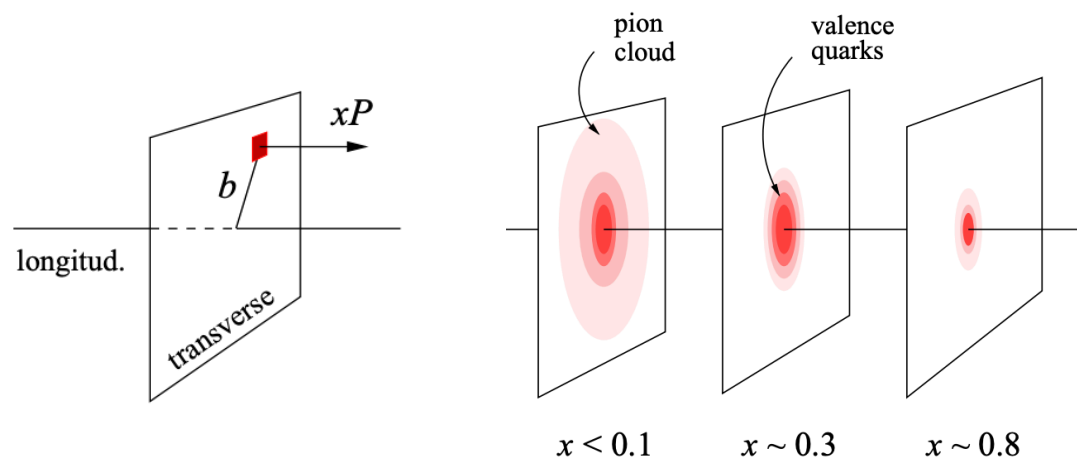
1_{mom} + 2_{coord} tomographic images of quark distribution in nucleon at fixed longitudinal momentum

3-D image from FT of the longitudinal mom. transfer

★ GPDs are not well-constrained experimentally:

- x-dependence extraction is not direct. DVCS amplitude: $\mathcal{H} = \int_{-1}^{+1} \frac{H(x, \xi, t)}{x - \xi + i\epsilon} dx$
(SDHEP [J. Qiu et al, arXiv:2205.07846] gives access to x)
- independent measurements to disentangle GPDs
- GPDs phenomenology more complicated than PDFs (multi-dimensionality)
- and more challenges ...

Generalized Parton Distributions



[H. Abramowicz et al., whitepaper for NSAC LRP, 2007]

1_{mom} + 2_{coord} tomographic images of quark distribution in nucleon at fixed longitudinal momentum

3-D image from FT of the longitudinal mom. transfer

★ GPDs are not well-constrained experimentally:

- x-dependence extraction is not direct. DVCS amplitude: $\mathcal{H} = \int_{-1}^{+1} \frac{H(x, \xi, t)}{x - \xi + i\epsilon} dx$
(SDHEP [J. Qiu et al, arXiv:2205.07846] gives access to x)
- independent measurements to disentangle GPDs
- GPDs phenomenology more complicated than PDFs (multi-dimensionality)
- and more challenges ...

Essential to complement the knowledge on GPD from lattice QCD

GPDs

**Through non-local matrix elements
of fast-moving hadrons**

GPDs on the lattice

- ★ GPDs: off-forward matrix elements of non-local light-cone operators
- ★ Off-forward correlators with nonlocal (equal-time) operators [\[Ji, PRL 110 \(2013\) 262002\]](#)

$$\tilde{q}_{\mu}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \gamma^{\mu} \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$
$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= Q_3/(2P_3) \end{aligned}$$

GPDs on the lattice

- ★ GPDs: off-forward matrix elements of non-local light-cone operators
- ★ Off-forward correlators with nonlocal (equal-time) operators [Ji, PRL 110 (2013) 262002]

$$\tilde{q}_{\mu}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \gamma^{\mu} \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$
$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= Q_3/(2P_3) \end{aligned}$$

- ★ Potential parametrization (γ^+ inspired)

$$F^{[\gamma^0]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^0 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{0\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$

$$F^{[\gamma^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^3 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{3\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$

GPDs on the lattice

- ★ GPDs: off-forward matrix elements of non-local light-cone operators
- ★ Off-forward correlators with nonlocal (equal-time) operators [Ji, PRL 110 (2013) 262002]

$$\tilde{q}_{\mu}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \gamma^{\mu} \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$

$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= Q_3/(2P_3) \end{aligned}$$

- ★ Potential parametrization (γ^+ inspired)

$$F^{[\gamma^0]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^0 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{0\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



reduction of power
corrections in fwd limit
[Radyushkin, PLB 767, 314, 2017]

$$F^{[\gamma^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^3 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{3\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



finite mixing with scalar
[Constantinou & Panagopoulos (2017)]

GPDs on the lattice

- ★ GPDs: off-forward matrix elements of non-local light-cone operators
- ★ Off-forward correlators with nonlocal (equal-time) operators [Ji, PRL 110 (2013) 262002]

$$\tilde{q}_{\mu}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \gamma^{\mu} \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$

$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= Q_3/(2P_3) \end{aligned}$$

- ★ Potential parametrization (γ^+ inspired)

$$F^{[r^0]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^0 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{0\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



reduction of power
corrections in fwd limit
[Radyushkin, PLB 767, 314, 2017]

$$F^{[r^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^3 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{3\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



finite mixing with scalar
[Constantinou & Panagopoulos (2017)]

- ★ Lorentz non-invariant parametrization
(typically symmetric frame to extract the “standard” GPDs)
- ★ Symmetric frame ($\vec{p}_f^s = \vec{P} + \vec{Q}/2, \vec{p}_i^s = \vec{P} - \vec{Q}/2$) requires separate calculations at each t

GPDs on the lattice

- ★ GPDs: off-forward matrix elements of non-local light-cone operators
- ★ Off-forward correlators with nonlocal (equal-time) operators [Ji, PRL 110 (2013) 262002]

$$\tilde{q}_{\mu}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \gamma^{\mu} \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$

$$\begin{aligned} \Delta &= P_f - P_i \\ t &= \Delta^2 = -Q^2 \\ \xi &= Q_3/(2P_3) \end{aligned}$$

- ★ Potential parametrization (γ^+ inspired)

$$F^{[\gamma^0]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^0 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{0\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



reduction of power
corrections in fwd limit
[Radyushkin, PLB 767, 314, 2017]

$$F^{[\gamma^3]}(x, \Delta; \lambda, \lambda'; P^3) = \frac{1}{2P^0} \bar{u}(p', \lambda') \left[\gamma^3 H_{Q(0)}(x, \xi, t; P^3) + \frac{i\sigma^{3\mu} \Delta_{\mu}}{2M} E_{Q(0)}(x, \xi, t; P^3) \right] u(p, \lambda)$$



finite mixing with scalar
[Constantinou & Panagopoulos (2017)]

- ★ Lorentz non-invariant parametrization
(typically symmetric frame to extract the “standard” GPDs)
- ★ Symmetric frame ($\vec{p}_f^s = \vec{P} + \vec{Q}/2, \vec{p}_i^s = \vec{P} - \vec{Q}/2$) requires separate calculations at each t

Light-cone GPDs using lattice correlators in non-symmetric frames

Theoretical setup

★ Parametrization of matrix elements in Lorentz invariant amplitudes

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Advantages

- Applicable to any kinematic frame and A_i have definite symmetries
- Lorentz invariant amplitudes A_i can be related to the standard H, E GPDs
- Quasi H, E may be redefined (Lorentz covariant):

Theoretical setup

★ Parametrization of matrix elements in Lorentz invariant amplitudes

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Advantages

- Applicable to any kinematic frame and A_i have definite symmetries
- Lorentz invariant amplitudes A_i can be related to the standard H, E GPDs
- Quasi H, E may be redefined (Lorentz covariant):

$$H(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = A_1 + \frac{\Delta_{s/a} \cdot z}{P_{avg,s/a} \cdot z} A_3 \quad E(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = -A_1 - \frac{\Delta_{s/a} \cdot z}{P_{avg,s/a} \cdot z} A_3 + 2A_5 + 2P_{avg,s/a} \cdot z A_6 + 2\Delta_{s/a} \cdot z A_8$$

Theoretical setup

★ Parametrization of matrix elements in Lorentz invariant amplitudes

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Advantages

- Applicable to any kinematic frame and A_i have definite symmetries
- Lorentz invariant amplitudes A_i can be related to the standard H, E GPDs
- Quasi H, E may be redefined (Lorentz covariant):

$$H(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = A_1 + \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3 \quad E(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = -A_1 - \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3 + 2A_5 + 2P_{avg,sla} \cdot z A_6 + 2\Delta_{sla} \cdot z A_8$$

Proof-of-concept calculation (zero quasi-skewness):

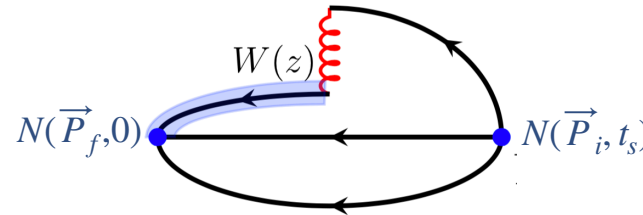
- symmetric frame:	$\vec{p}_f^s = \vec{P} + \frac{\vec{Q}}{2},$	$\vec{p}_i^s = \vec{P} - \frac{\vec{Q}}{2}$	$t^s = -\vec{Q}^2$
- asymmetric frame:	$\vec{p}_f^a = \vec{P},$	$\vec{p}_i^a = \vec{P} - \vec{Q}$	$t^a = -\vec{Q}^2 + (E_f - E_i)^2$

Parameters of calculation

★ Nf=2+1+1 twisted mass (TM) fermions & clover improvement

★ Calculation:

- isovector combination
- zero skewness
- $T_{\text{sink}}=1$ fm



Pion mass: 260 MeV
Lattice spacing: 0.093 fm
Volume: $32^3 \times 64$
Spatial extent: 3 fm

frame	P_3 [GeV]	\vec{Q} [$\frac{2\pi}{L}$]	$-t$ [GeV ²]	ξ	N_{ME}	N_{confs}	N_{src}	N_{tot}
symm	1.25	$(\pm 2, 0, 0), (0, \pm 2, 0)$	0.69	0	8	249	8	15936
non-symm	1.25	$(\pm 2, 0, 0), (0, \pm 2, 0)$	0.64	0	8	269	8	17216

★ Computational cost:

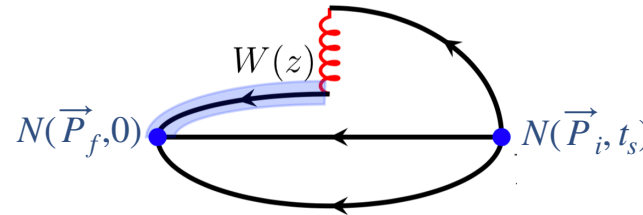
- symmetric frame 4 times more expensive than asymmetric frame for same set of \vec{Q} (requires separate calculations at each t)

Parameters of calculation

★ Nf=2+1+1 twisted mass (TM) fermions & clover improvement

★ Calculation:

- isovector combination
- zero skewness
- $T_{\text{sink}}=1$ fm



Pion mass: 260 MeV
Lattice spacing: 0.093 fm
Volume: $32^3 \times 64$
Spatial extent: 3 fm

frame	P_3 [GeV]	\mathbf{Q} [$\frac{2\pi}{L}$]	$-t$ [GeV ²]	ξ	N_{ME}	N_{confs}	N_{src}	N_{tot}
symm	1.25	$(\pm 2, 0, 0), (0, \pm 2, 0)$	0.69	0	8	249	8	15936
non-symm	1.25	$(\pm 2, 0, 0), (0, \pm 2, 0)$	0.64	0	8	269	8	17216

Small difference: $t^s = -\vec{Q}^2$ $t^a = -\vec{Q}^2 + (E_f - E_i)^2$

$$A(-0.64 \text{ GeV}^2) \sim A(-0.69 \text{ GeV}^2)$$

★ Computational cost:

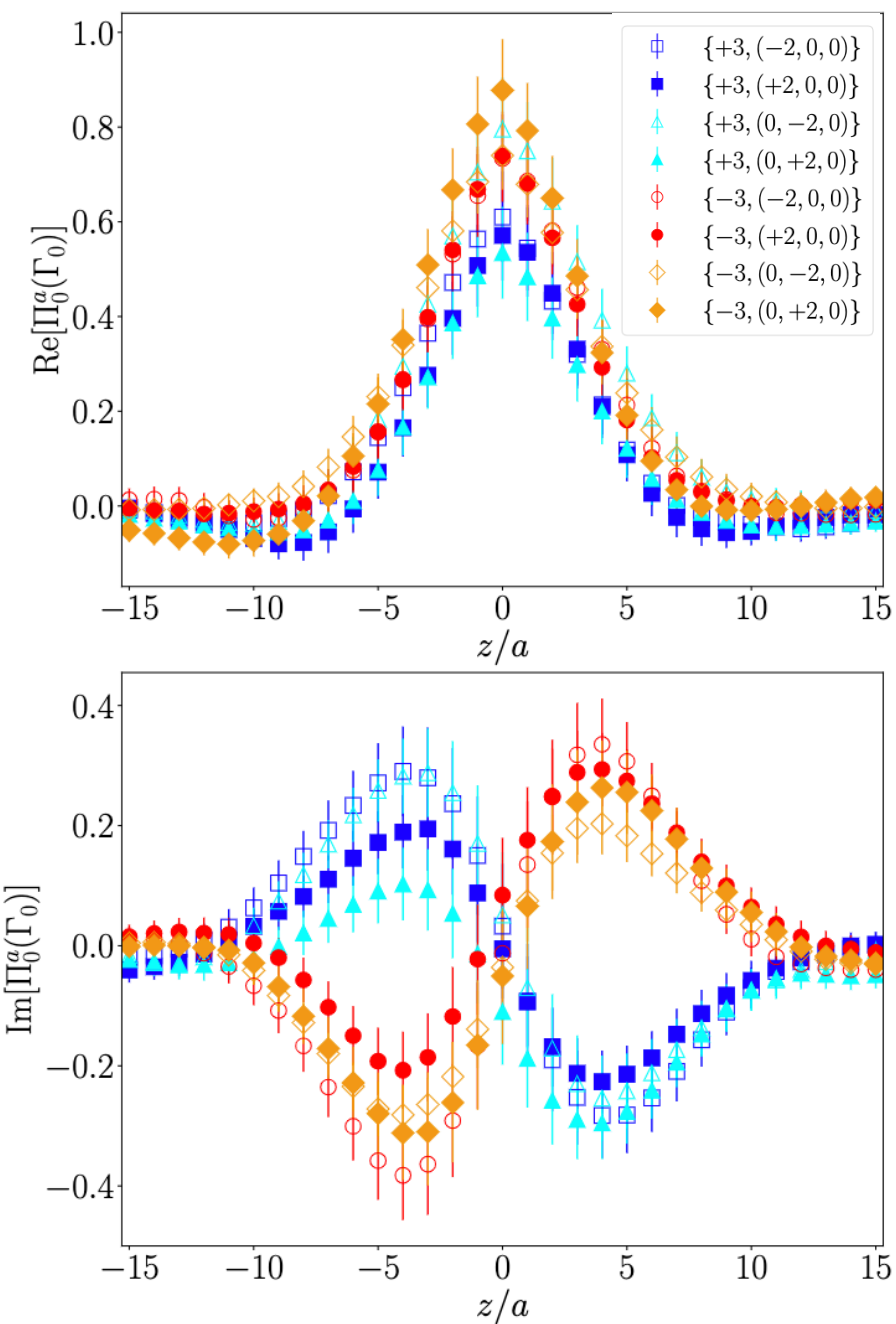
- symmetric frame 4 times more expensive than asymmetric frame for same set of \vec{Q} (requires separate calculations at each t)

Results: matrix elements

- ★ Eight independent matrix elements needed to disentangle the A_i asymmetric frame

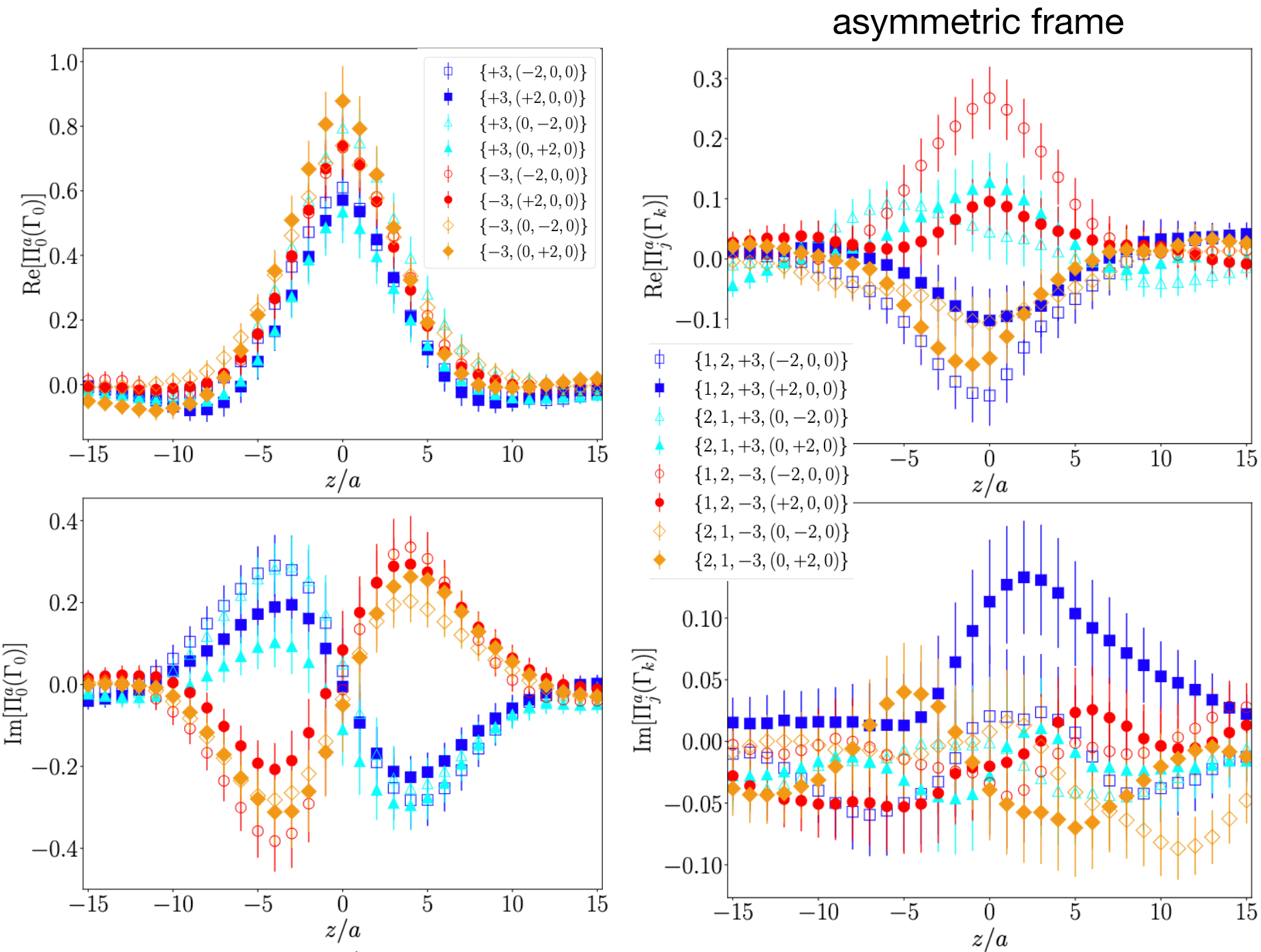
Results: matrix elements

- ★ Eight independent matrix elements needed to disentangle the A_i asymmetric frame



Results: matrix elements

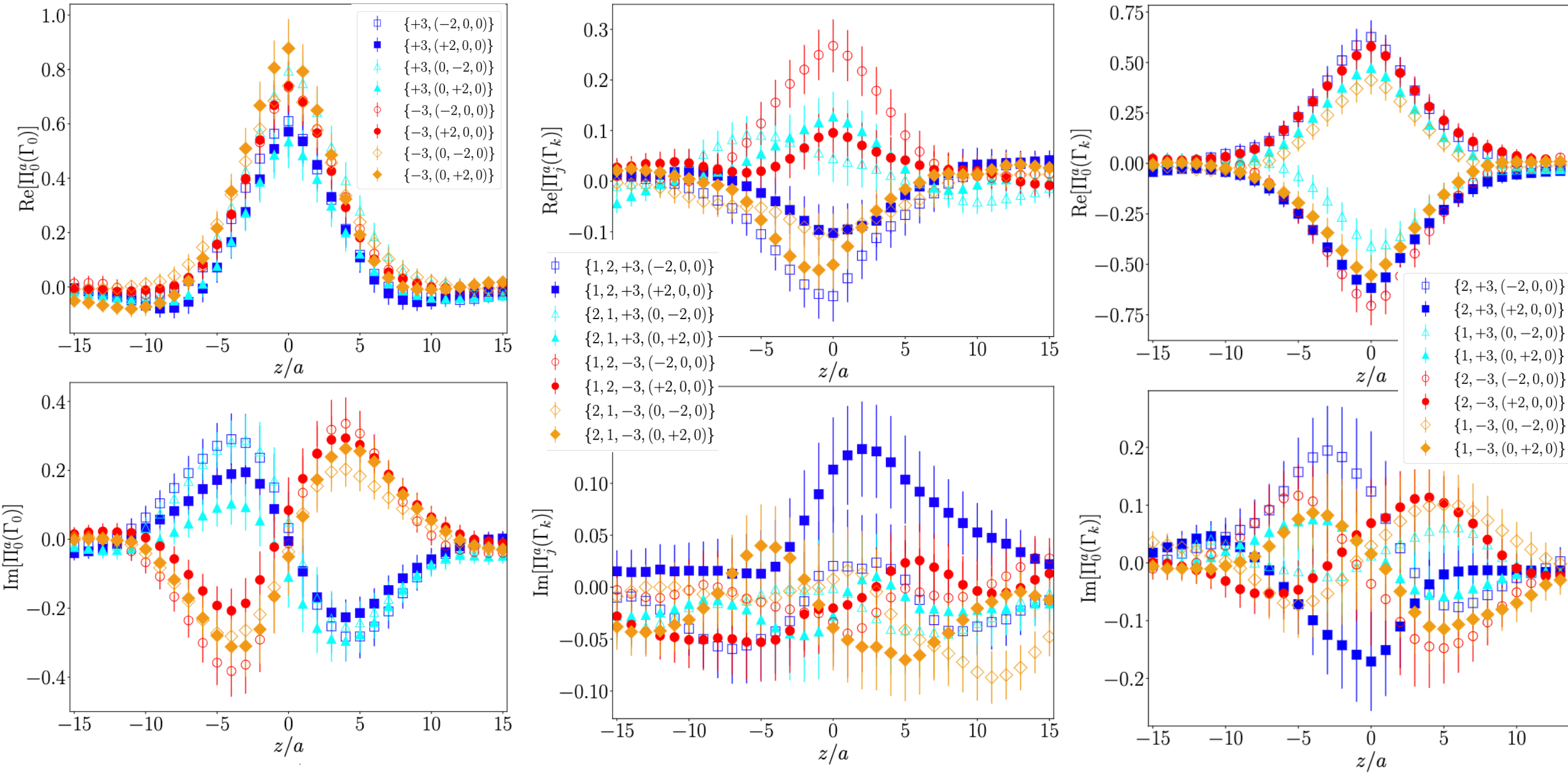
★ Eight independent matrix elements needed to disentangle the A_i



Results: matrix elements

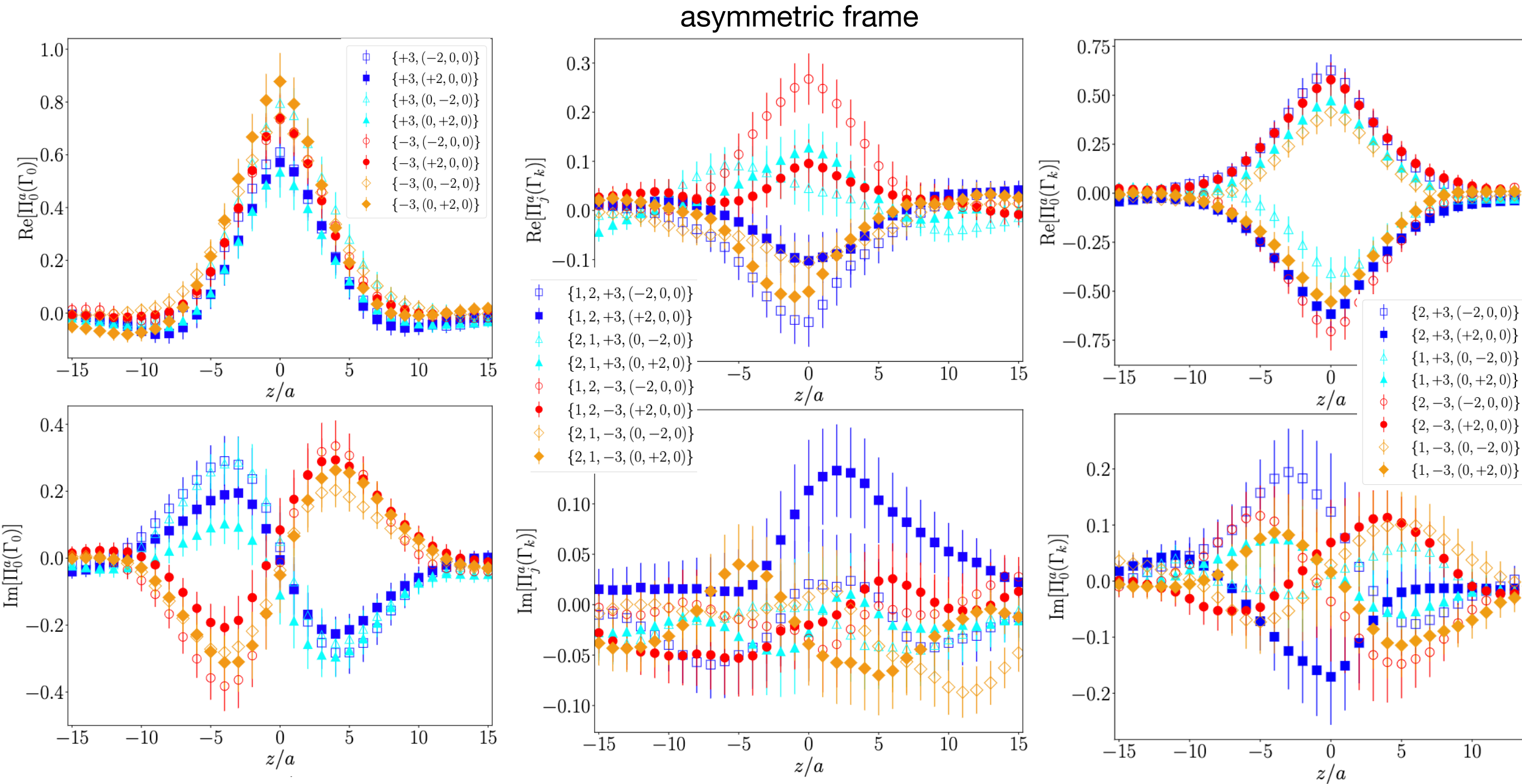
★ Eight independent matrix elements needed to disentangle the A_i

asymmetric frame



Results: matrix elements

★ Eight independent matrix elements needed to disentangle the A_i

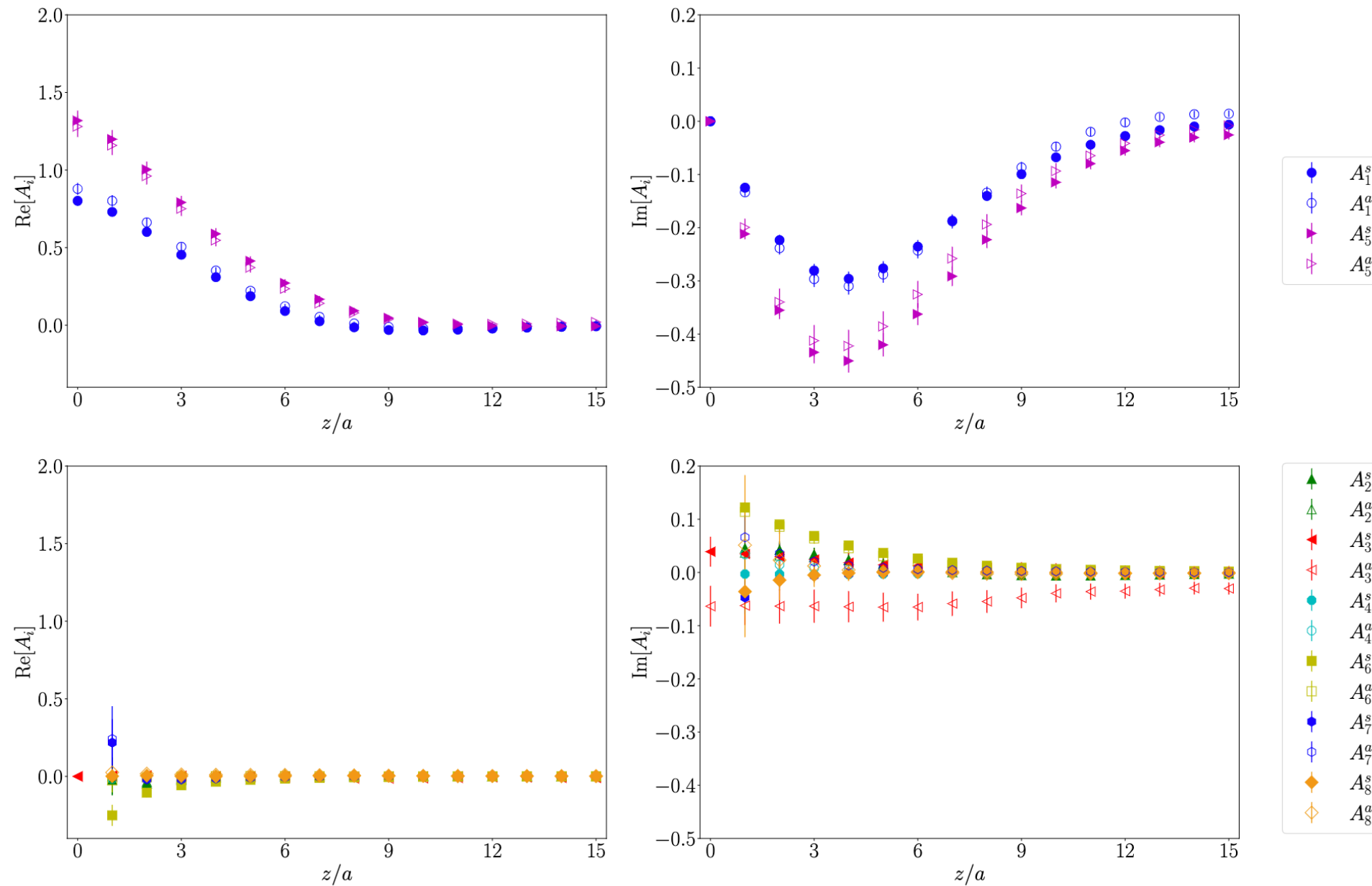


★ Asymmetric frame: ME do not have definite symmetries in $\pm P_3, \pm Q, \pm z$

★ Noisy ME lead to challenges in extracting A_i of sub-leading magnitude

How do the A_i compare between frames?

How do the A_i compare between frames?



- ★ A_1, A_5 dominant contributions
- ★ Full agreement in two frames for both Re and Im parts of A_1, A_5
- ★ Remaining A_i suppressed (at least for this kinematic setup and $\xi = 0$)

quasi-GPDs in terms of A_i

- ★ The mapping of A_i to the quasi-GPDs is not unique
- ★ Construction of a Lorentz invariant definition may be beneficial

$$(\xi = 0) \quad \Pi_H^{\text{impr}} = A_1$$

$$\Pi_E^{\text{impr}} = -A_1 + 2A_5 + 2zP_3A_6$$

- ★ All quasi-GPDs definitions converge to the same light-cone GPDs
(up to systematic effects)

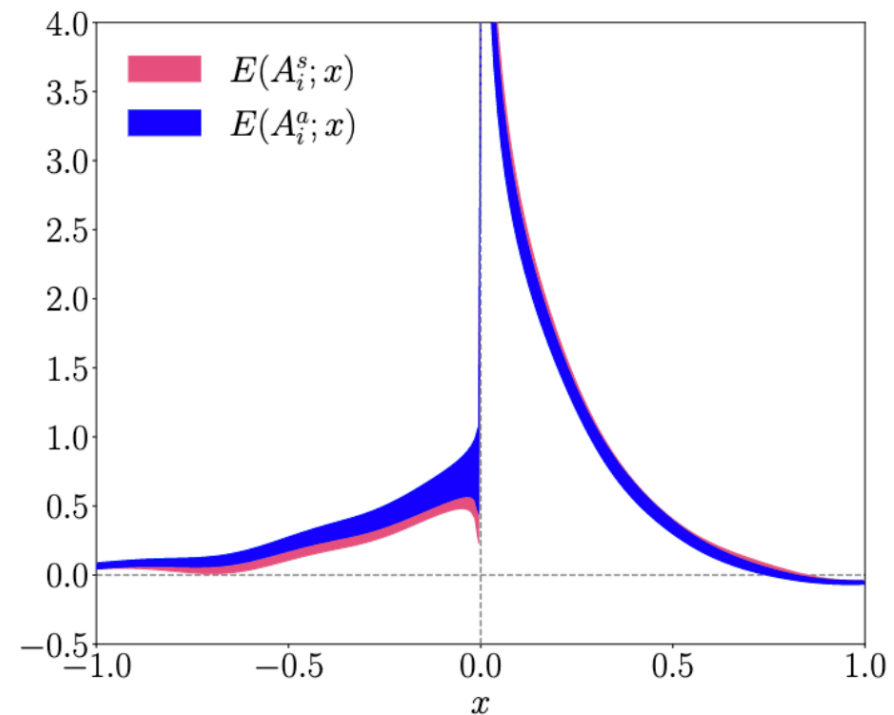
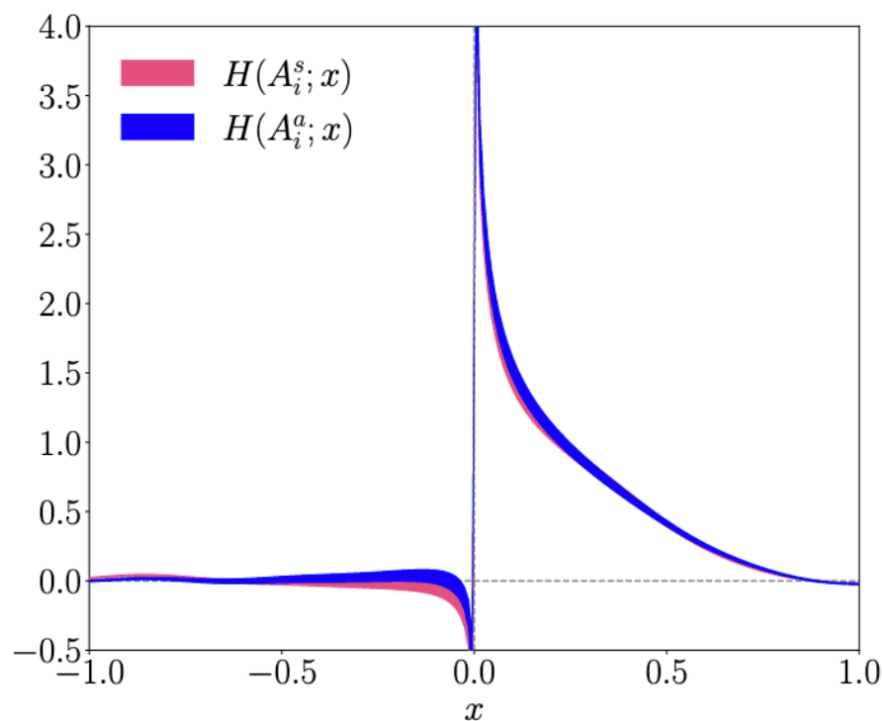
quasi-GPDs in terms of A_i

- ★ The mapping of A_i to the quasi-GPDs is not unique
- ★ Construction of a Lorentz invariant definition may be beneficial

$$(\xi = 0) \quad \Pi_H^{\text{impr}} = A_1$$

$$\Pi_E^{\text{impr}} = -A_1 + 2A_5 + 2zP_3A_6$$

- ★ All quasi-GPDs definitions converge to the same light-cone GPDs (up to systematic effects)



Agreement between frames for both quasi-GPDs (by definition)

Beyond exploration

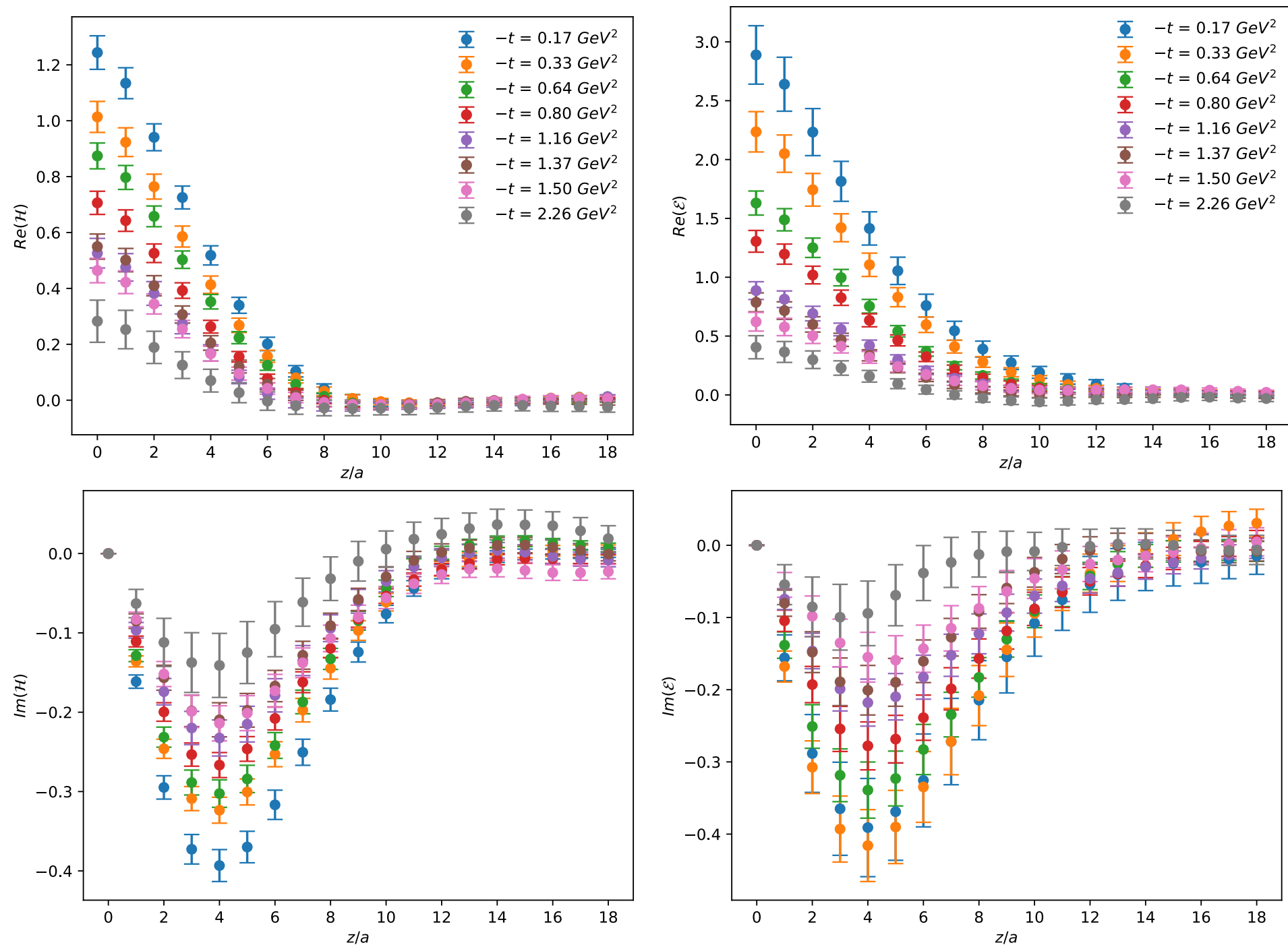
- ★ 11 values of $-t$ (3 in symm. frame and 8 in asymm. frame)
- ★ Separate calculation for each $-t$ value in symmetric frame
- ★ Two groups of $-t$ value in asymmetric frame: $\vec{Q} = (Q_x, 0, 0), (Q_x, Q_y, 0)$

frame	P_3 [GeV]	Δ [$\frac{2\pi}{L}$]	$-t$ [GeV ²]	ξ	N_{ME}	N_{confs}	N_{src}	N_{tot}
N/A	± 1.25	(0,0,0)	0	0	2	731	16	23392
symm	± 0.83	($\pm 2, 0, 0$), ($0, \pm 2, 0$)	0.69	0	8	67	8	4288
symm	± 1.25	($\pm 2, 0, 0$), ($0, \pm 2, 0$)	0.69	0	8	249	8	15936
symm	± 1.67	($\pm 2, 0, 0$), ($0, \pm 2, 0$)	0.69	0	8	294	32	75264
symm	± 1.25	($\pm 2, \pm 2, 0$)	1.39	0	16	224	8	28672
symm	± 1.25	($\pm 4, 0, 0$), ($0, \pm 4, 0$)	2.76	0	8	329	32	84224
asymm	± 1.25	($\pm 1, 0, 0$), ($0, \pm 1, 0$)	0.17	0	8	429	8	27456
asymm	± 1.25	($\pm 1, \pm 1, 0$)	0.33	0	16	194	8	12416
asymm	± 1.25	($\pm 2, 0, 0$), ($0, \pm 2, 0$)	0.64	0	8	429	8	27456
asymm	± 1.25	($\pm 1, \pm 2, 0$), ($\pm 2, \pm 1, 0$)	0.80	0	16	194	8	12416
asymm	± 1.25	($\pm 2, \pm 2, 0$)	1.16	0	16	194	8	24832
asymm	± 1.25	($\pm 3, 0, 0$), ($0, \pm 3, 0$)	1.37	0	8	429	8	27456
asymm	± 1.25	($\pm 1, \pm 3, 0$), ($\pm 3, \pm 1, 0$)	1.50	0	16	194	8	12416
asymm	± 1.25	($\pm 4, 0, 0$), ($0, \pm 4, 0$)	2.26	0	8	429	8	27456

- ★ Momentum transfer range is very optimistic
(some values have enhanced systematic uncertainties)

Unpolarized quasi-GPDs

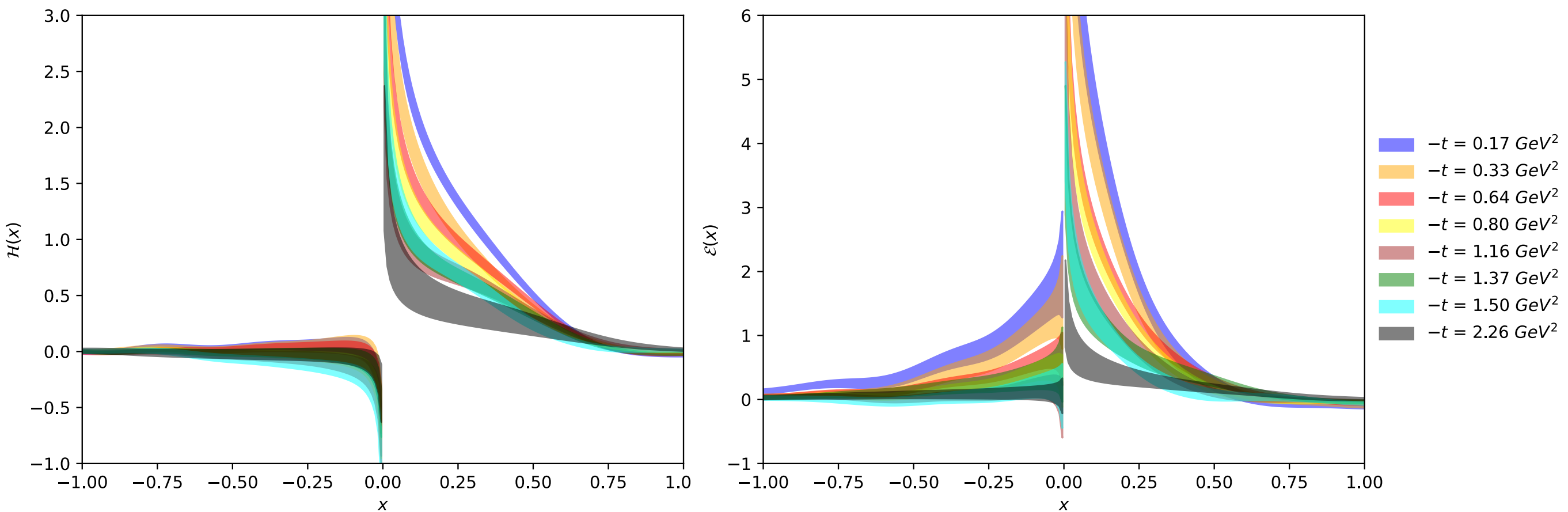
asymmetric frame



- ★ Impressive quality of signal quality
- ★ Behavior with increasing $-t$ as “expected” qualitatively

Unpolarized light-cone GPDs

- ★ quasi-GPDs transformed to momentum space
- ★ Matching formalism to 1 loop accuracy level
- ★ $\pm x$ correspond to quark and anti-quark region
- ★ Anti-quark region susceptible to systematic uncertainties.



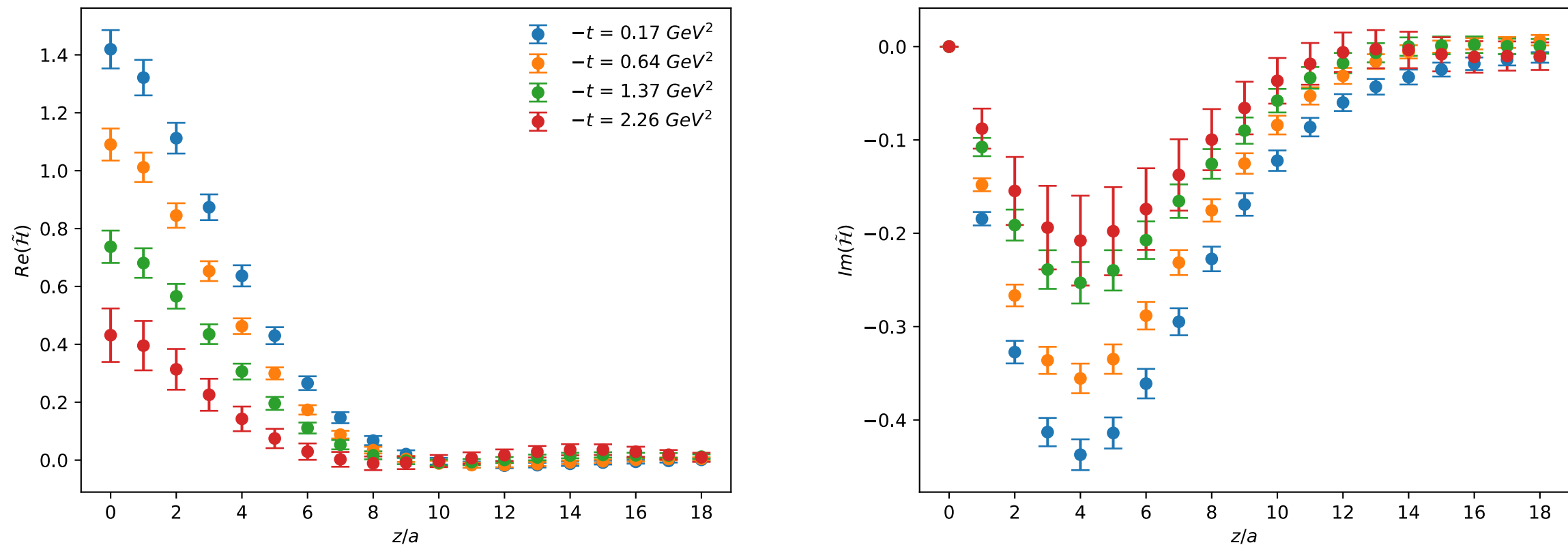
Helicity quasi-GPDs

- ★ Lorentz-invariant decomposition applicable to helicity case
- ★ At $\xi = 0$ only \widetilde{H} is accessible directly
(\widetilde{E} accessible from parametrization of the t dependence)

Helicity quasi-GPDs

- ★ Lorentz-invariant decomposition applicable to helicity case
- ★ At $\xi = 0$ only \widetilde{H} is accessible directly
(\widetilde{E} accessible from parametrization of the t dependence)

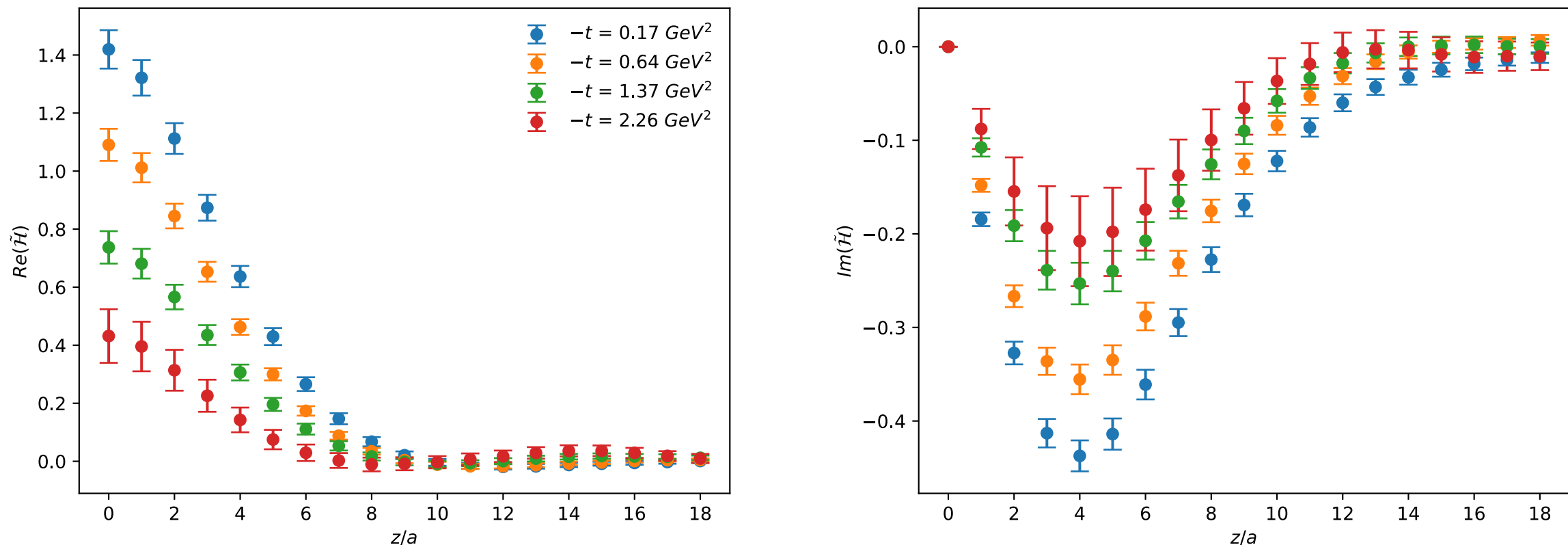
asymmetric frame



Helicity quasi-GPDs

- ★ Lorentz-invariant decomposition applicable to helicity case
- ★ At $\xi = 0$ only \widetilde{H} is accessible directly
(\widetilde{E} accessible from parametrization of the t dependence)

asymmetric frame



- ★ All values of t obtained at the cost of one
- ★ Preliminary analysis very encouraging!

How to lattice QCD data fit into the overall effort for hadron tomography

How to lattice QCD data fit into the overall effort for hadron tomography

- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence

How to lattice QCD data fit into the overall effort for hadron tomography

- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence



QUARK-GLUON TOMOGRAPHY COLLABORATION



U.S. DEPARTMENT OF
ENERGY

Office of
Science

Award Number:
DE-SC0023646

1. **Theoretical studies** of high-momentum transfer processes using perturbative QCD methods and study of GPDs properties
2. **Lattice QCD** calculations of GPDs and related structures
3. **Global analysis** of GPDs based on experimental data using modern data analysis techniques for inference and uncertainty quantification

Summary

- ★ Lattice QCD data on GPDs will play an important role in the pre-EIC era and can complement experimental efforts of JLab@12GeV
- ★ New proposal for Lorentz invariant decomposition has great advantages:
 - significant reduction of computational cost
 - access to a broad range of t and ξ
- ★ Future calculations have the potential to transform the field of GPDs
- ★ Mellin moments can be extracted utilizing quasi-GPDs data
- ★ Synergy with phenomenology is an exciting prospect!

Xiang Gao, Tue 11:50am

Summary

- ★ Lattice QCD data on GPDs will play an important role in the pre-EIC era and can complement experimental efforts of JLab@12GeV
- ★ New proposal for Lorentz invariant decomposition has great advantages:
 - significant reduction of computational cost
 - access to a broad range of t and ξ
- ★ Future calculations have the potential to transform the field of GPDs
- ★ Mellin moments can be extracted utilizing quasi-GPDs data
- ★ Synergy with phenomenology is an exciting prospect!

Xiang Gao, Tue 11:50am

Thank you



DOE Early Career Award (NP)
Grant No. DE-SC0020405