DIS 2023

Michigan State University 27 – 31 March 2023



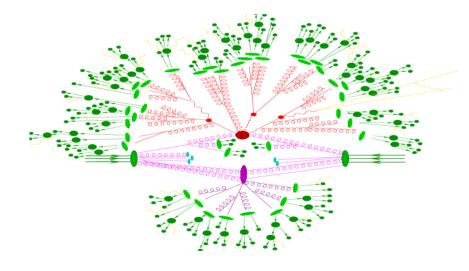
Underlying Event Measurements at ATLAS

Claire Gwenlan, Oxford

on behalf of the ATLAS collaboration

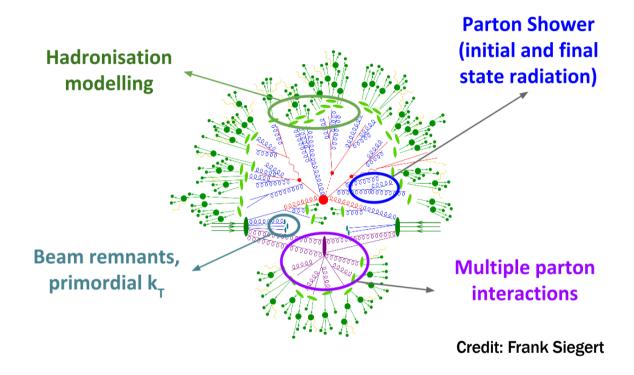






- Observables sensitive to colour reconnection in ttbar events (<u>arXiv:2209.07874</u>)
- Correlation of Y meson production with the underlying event (<u>ATLAS-CONF-2022-23</u>)

Underlying Event at the LHC

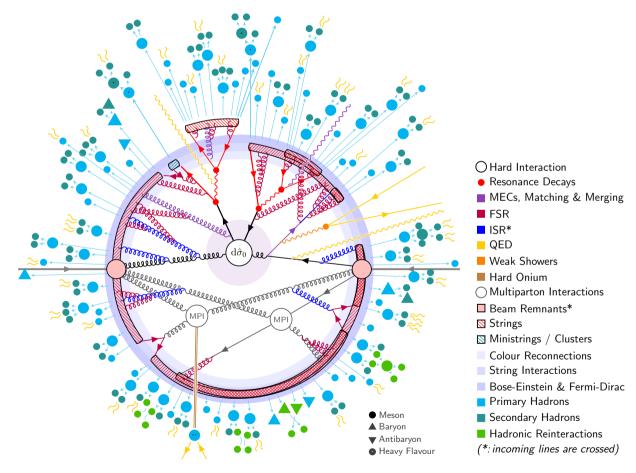


- underlying event: any hadronic activity not associated with hard scattering process
- contributions not always calculable in perturbative QCD → phenomenological models in MC,
 which must be tuned to data
- typically modelled with :
- multiple parton interactions
- initial and final state radiation
- colour reconnection (CR) with beam remnants



Measurements of observables sensitive to colour reconnection in tt events with the ATLAS detector

arXiv:2209.07874, accepted by EPJ

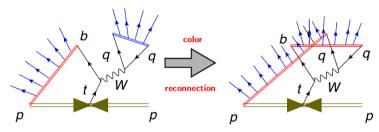


structure of a pp \rightarrow tt event, as modelled by PYTHIA.

Motivation

- CR experimentally motivated → first introduced in MC to describe rise of (pT) vs nch
- CR intimately connected to MPI
- new constraints on CR and MPI,
 will provide improved description everywhere
- specifically in context of top, modelling of CR is a dominant systematic uncertainty on top mass determination

Color reconnection affects the reconstruction of the top system

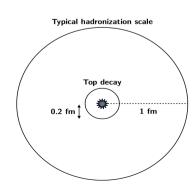


Ambiguity in the definition of the top mass: $m_{\rm top}^2 \neq (p_{\rm b} + p_{\rm j1} + p_{\rm j2})^2$ Credit: S Argyropoulos JHEP11 (2014) 043

 interesting unresolved question regarding involvement of top quark versus its decay products in CR

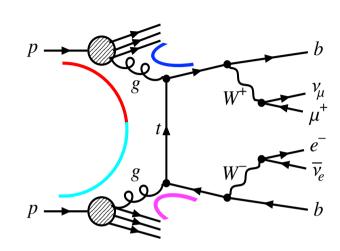
	$m_{\mathrm{top}} \; [\mathrm{GeV}]$
Result	172.21
Statistics	0.20
Method	0.05 ± 0.04
Matrix-element matching	0.40 ± 0.06
Parton shower and hadronisation	0.05 ± 0.05
Initial- and final-state QCD radiation	0.17 ± 0.02
Underlying event	0.02 ± 0.10
Colour reconnection	0.27 ± 0.07
Parton distribution function	0.03 ± 0.00
Single top modelling	0.01 ± 0.01
Background normalisation	0.03 ± 0.02
Jet energy scale	0.37 ± 0.02
b-jet energy scale	0.12 ± 0.02
Jet energy resolution	0.13 ± 0.02
Jet vertex tagging	0.01 ± 0.01
b-tagging	0.04 ± 0.01
Leptons	0.11 ± 0.02
Pile-up	0.06 ± 0.01
Recoil effect	0.39 ± 0.09
Total systematic uncertainty (without recoil)	0.67 ± 0.05
Total systematic uncertainty (with recoil)	0.77 ± 0.06
Total uncertainty (without recoil)	0.70 ± 0.05
Total uncertainty (with recoil)	0.80 ± 0.06

ATLAS-CONF-2022-058



ATLAS measurement of top quarks and CR

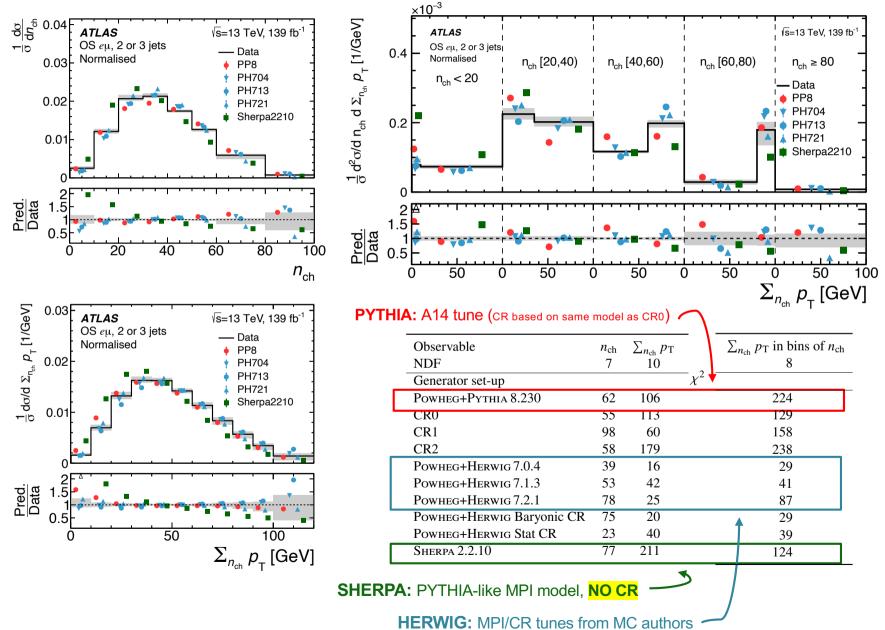
- select tt events in di-leptonic eμ channel
 (exactly 1 e, 1 μ and 2 or 3 jets, 2 of which b-tagged)
- measure inclusive[†] charged particle properties:
 † excluding leptons or jet tracks, pT > 0.5 GeV, |η| < 2.5
 - multiplicity n_{ch}
 - scalar sum of charged particle transverse momentum, $\Sigma_{n^{\mathrm{ch}}} p_{\mathrm{T}}$
 - $\Sigma_{n^{\text{ch}}} p_{\text{T}}$ in bins of n_{ch}



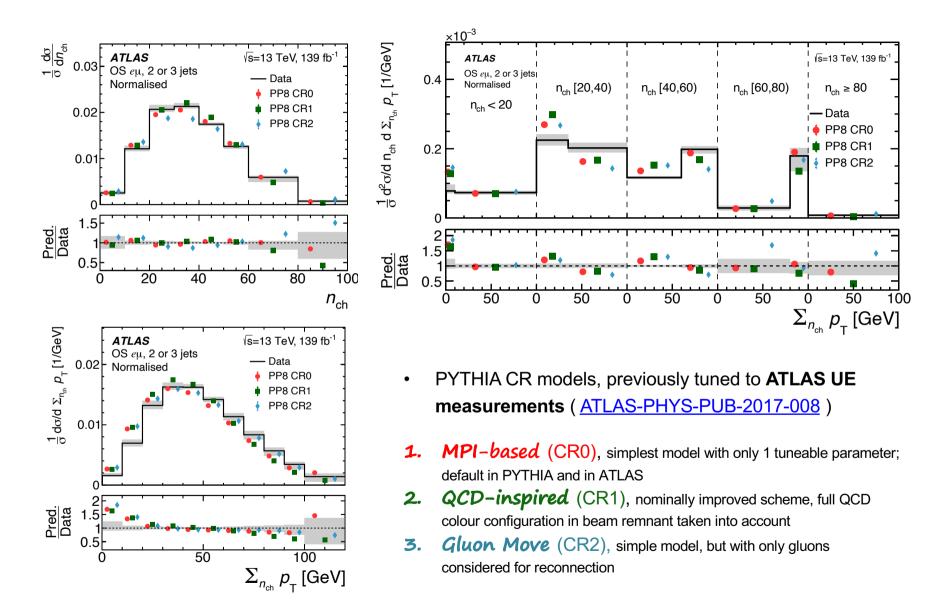
Credit: Z Citron for ATLAS, MPI@LHC

- pileup and fake-lepton contribution subtracted with MC templates
- compare MC with unfolded normalised differential cross sections:
- PYTHIA8: Lund string model for hadronisation several CR models
- HERWIG7: cluster hadronisation several CR models
- many tuneable parameters

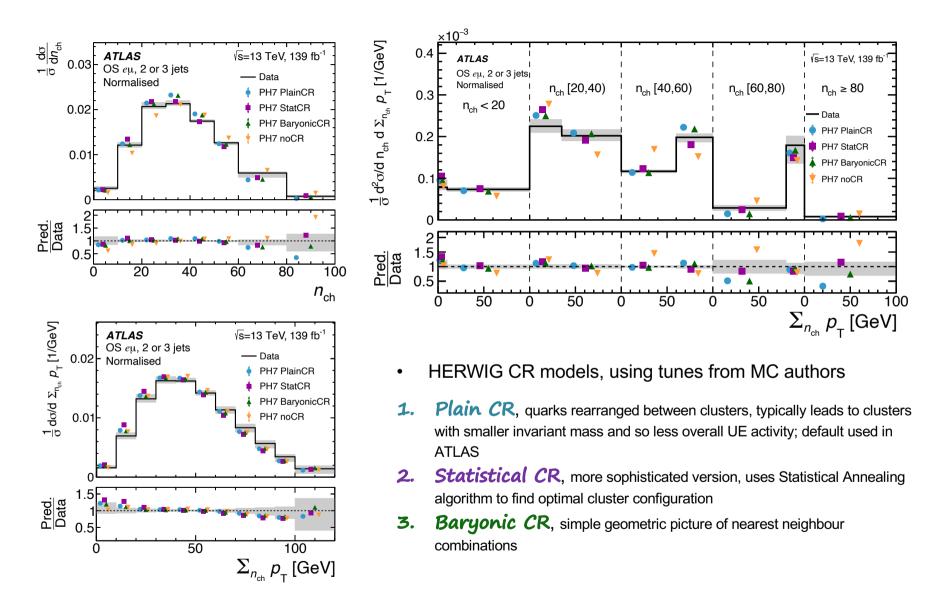
cf. different MC Generators



cf. PYTHIA8 CR Models



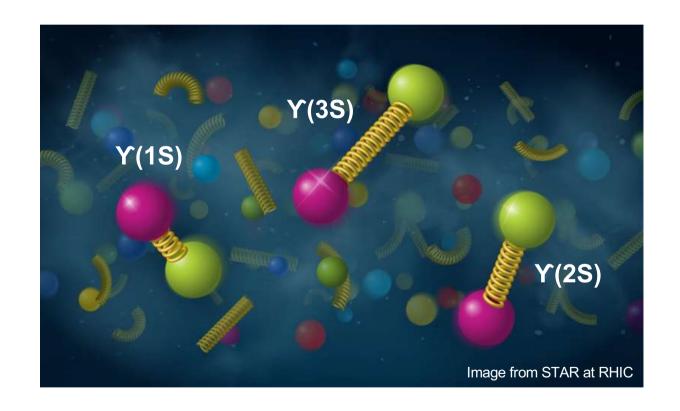
cf. HERWIG7 CR Models



overall best description is by the HERWIG CR models, though no panacea

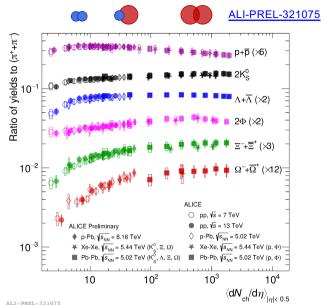
Correlation of Y meson production with the underlying event in pp collisions measured by the ATLAS experiment

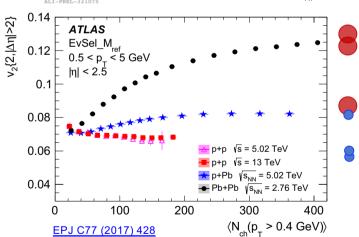
ATLAS-CONF-2022-23



Motivation

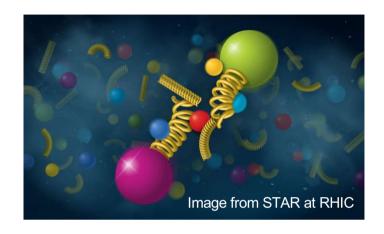
- Quark Gluon Plasma (QGP) in A+A well established,
 but smaller collision systems controversial
- → naively do not expect QGP formation in small systems
- → experimental observations of QGP-like behaviour in p+A and p+p calls this into question
- many studies of small systems demonstrating QGP-like phenomena in soft physics, EG.
 strangeness enhancement, multi-particle
 correlations in peripheral A+A, p+A, and also p+p
- → not many studies using hard probes
- ATLAS measurement (ATLAS-CONF-2022-23)
- study correlations between upsilon meson (Υ)
 production and inclusive charged particles to bridge soft-hard gap

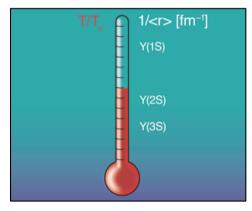




Why Y mesons?

- HI perspective: heavy quarkonia produced early in collision → experience full evolution
- measuring Y(nS) state suppression can be a thermometer for a QGP:
- 3 distinct Y (bb) states: Y(1S) Y(2S) Y(3S)
- sequentially "dissolve" or dissociate when potential between constituent quarks screened by colour charges of q, g in QGP
- the more loosely bound the state, the more easily it dissolves: Υ(3S) is more suppressed than Υ(2S) is more suppressed than Υ(1S)
- Y less affected by recombination effects than cc states, given small probability for b quarks
- → Y ideal probe to study QGP



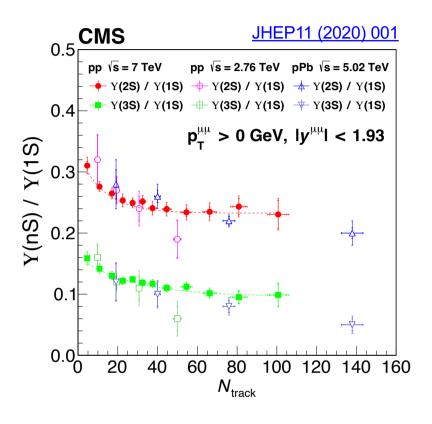


Adapted from EPJ C61 (2009) 705

CMS measurements of $\Upsilon(nS)$ and multiplicity

CMS

0.4



Y(nS) / $0.70 \le S_{\tau} < 0.85$ $0.85 \le S_{-} \le 1.00$ 0.1 $p_{-}^{\mu\mu} > 7 \text{ GeV}, |y^{\mu\mu}| < 1.2$ 20 40 60 80 100 120 N_{track} ... measured also as a function of

4.8 fb⁻¹ (7 TeV)

 $0.00 \le S_{T} < 0.55$

 $0.55 \le S_{T} < 0.70$

 $-0.70 \le S_{\tau} < 0.85$ $0.85 \le S_{\pm} \le 1.00$

 $0.00 \le S_{-} < 0.55$

 $0.55 \le S_{\tau} < 0.70$

140

 $\odot z$

Y(2S) / Y(1S)

- CMS results from 2014 (JHEP04 (2014) 103) and confirmed with more detail in 2020:
- Y(nS)/Y(1S) ratio decreases as a function of charged particle multiplicity

event topology

(sphericity)

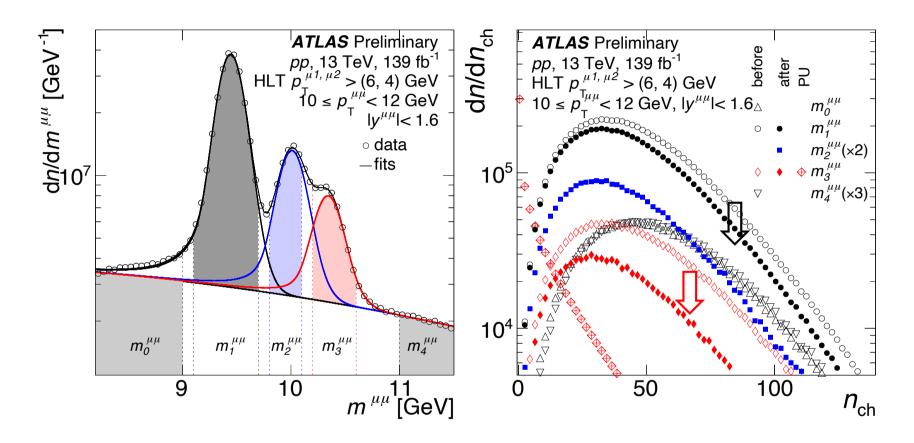
ATLAS approach to study Y(nS) correlation with UE

- ATLAS analysis is an "inversion" of CMS approach:
- instead of measuring "conventional" variables such as Y(nS) yields vs nch,
 ATLAS measures nch for different Y(nS)
- ATLAS full Run-2 √s=13 TeV pp dataset
- Y(nS) → μμ events with:
 - $8.2 \le m^{\mu\mu} < 11.8 \text{ GeV}$
 - $|y^{\mu\mu}| < 1.6$
- charged particle tracks
- not directly involved in formation of Y
 - 0.5 < pt < 10 GeV
 - $|\eta| < 2.5$



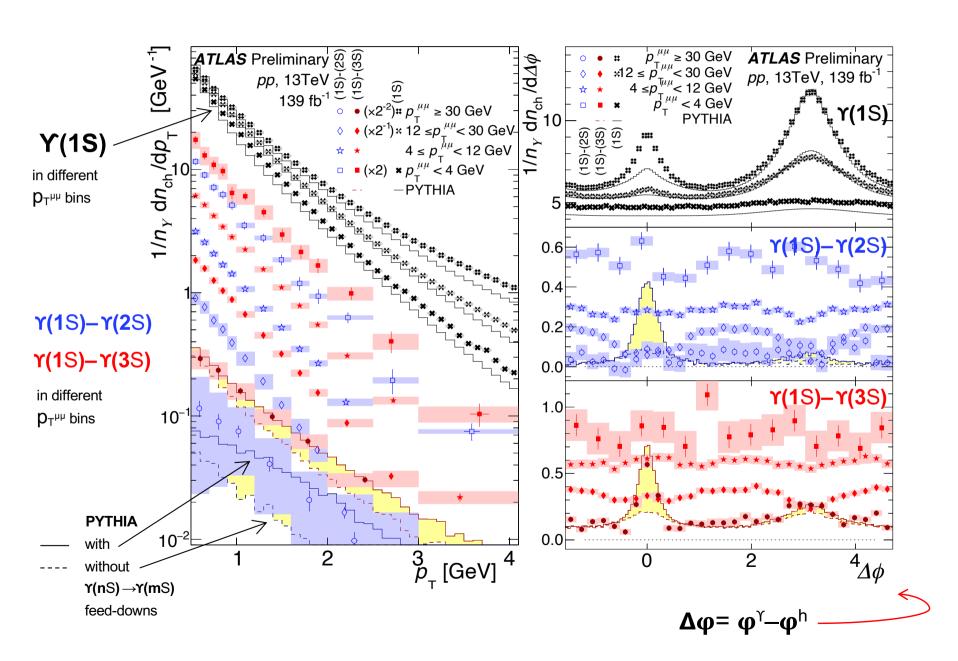
- search for modification of underlying event (soft) for different Y states (hard) in pp events – measure:
- multiplicity *nch*
- kinematic distributions: dnch/dpT and $dnch/d\Delta\phi$, where $\Delta\phi = \phi^{\gamma} \phi^{h}$
- analysis performed in p^{µµ} intervals and observables corrected for pileup

ATLAS measurement of $\Upsilon(nS)$ and underlying event



- well controlled background Y candidates in 5 mass regions → signal+BG fits and sideband subtraction
- pileup corrected on a statistical basis using event mixing technique, developed for EPJC 80 (2020) 64 (illustrated here with nch, but works also for dnch/dp τ and dnch/d $\Delta φ$)

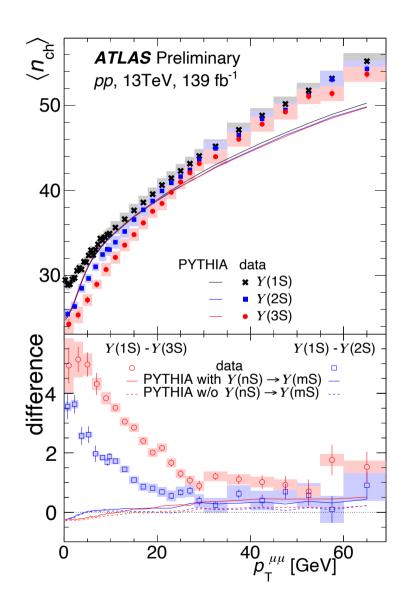
Kinematic distributions of $\Upsilon(1S) - \Upsilon(nS)$



Mean values of *n*_{ch} distribution

- significant differences observed for different Y(nS) states
- PYTHIA mismodels Y production no (nch) dependence of different states
- effect largest at $p_T^{\mu\mu}=0$:
- $\Upsilon(1S) \Upsilon(2S) \Delta \langle n_{ch} \rangle = 3.6 \pm 0.4$
- $\Upsilon(1S) \Upsilon(3S) \Delta \langle n_{ch} \rangle = 4.9 \pm 0.4$

 observations cannot be explained by mass differences, feed-down of states, or systematic uncertainties



summary

Observables sensitive to CR in ttbar events (arXiv:2209.07874):

ATLAS measurement of CR sensitive observables in top quark events gives detailed handle on CR and MPI

Correlation of Y meson production with the underlying event

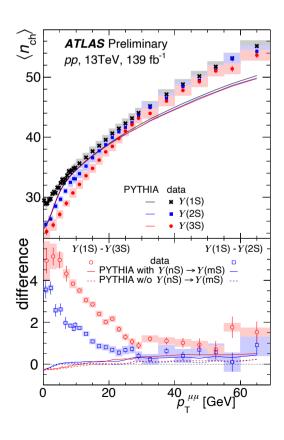
(<u>ATLAS-CONF-2022-23</u>):

significant differences in charged particle multiplicity observed

- **Y(1S)**: $\langle n_{ch} \rangle = 29.5 \pm 0.7$
- Y(2S): 12% less
- Y(3S): 17% less

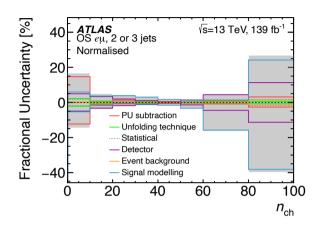
can be interpreted as suppression of excited states at higher multiplicity

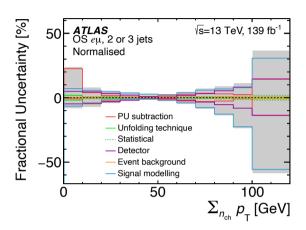
→ something interesting going on in **pp** that must be further explored

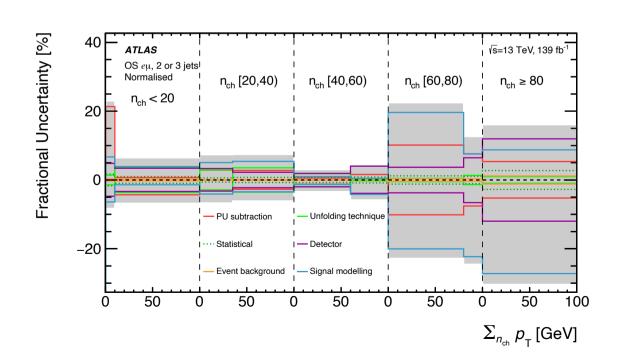


extras

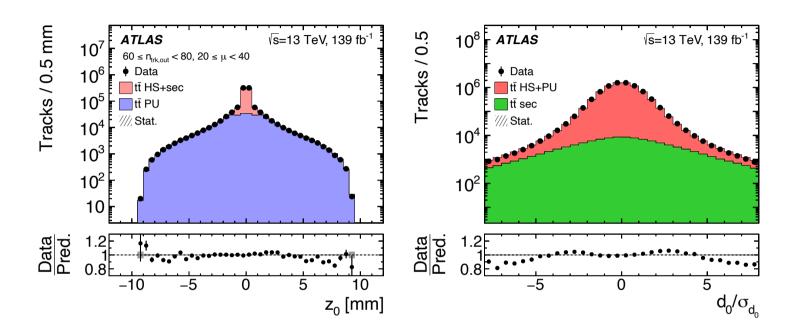
ATLAS measurement of top quarks and CR







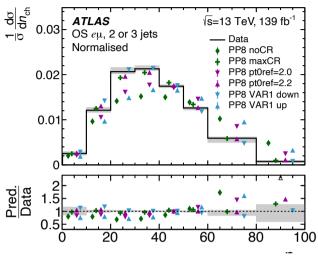
Track-based BG

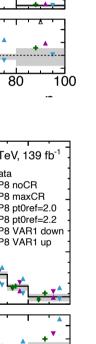


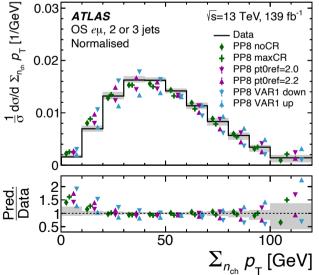
pileup scale factors c_{PU}

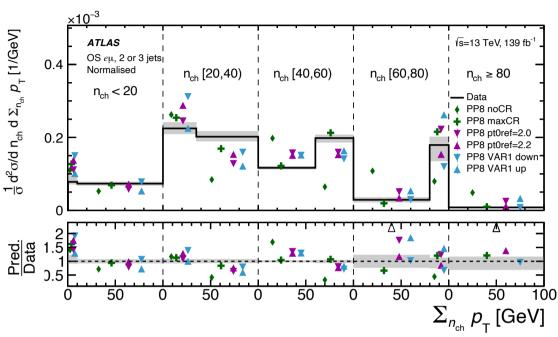
Region	$n_{\rm trk,out} < 20$	$20 \le n_{\rm trk,out} < 40$	$40 \le n_{\rm trk,out} < 60$	$60 \le n_{\rm trk,out} < 80$	$80 \le n_{\rm trk,out} \le 100$
$\mu < 20$	0.91	1.04	0.97	1.05	1.08
$20 \le \mu < 40$	0.91	1.08	1.08	1.07	1.11
$\mu \ge 40$	0.95	1.15	1.23	1.27	1.36

cf. PYTHIA8 - parameter variations



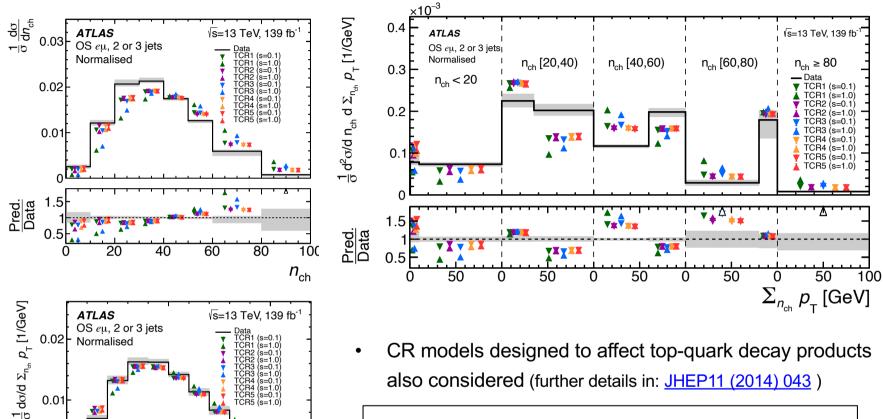






- CR range parameter R_{rec} is set to its maximal value of $R_{\text{rec}} = 10$, such that the reconnection probability reaches saturation (maxCR). The default is 1.71.
- CR is switched off, i.e. $R_{rec} = 0$ (noCR).
- MPI parameter p_{T0}^{ref} , by default set to 2.09, is lowered to 2.0 and raised to 2.2.
- UE activity is varied by using the Var1 eigentune of the A14 tune [21]. This eigentune includes variations of α_s and variations of R_{rec} .

cf. PYTHIA8 - top quark specific CR Models

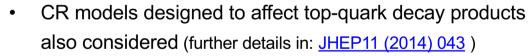


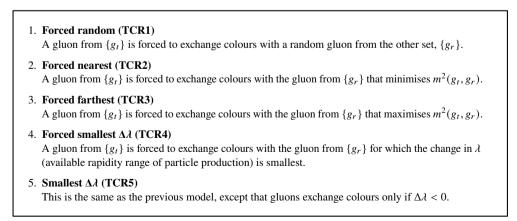
0.01

50

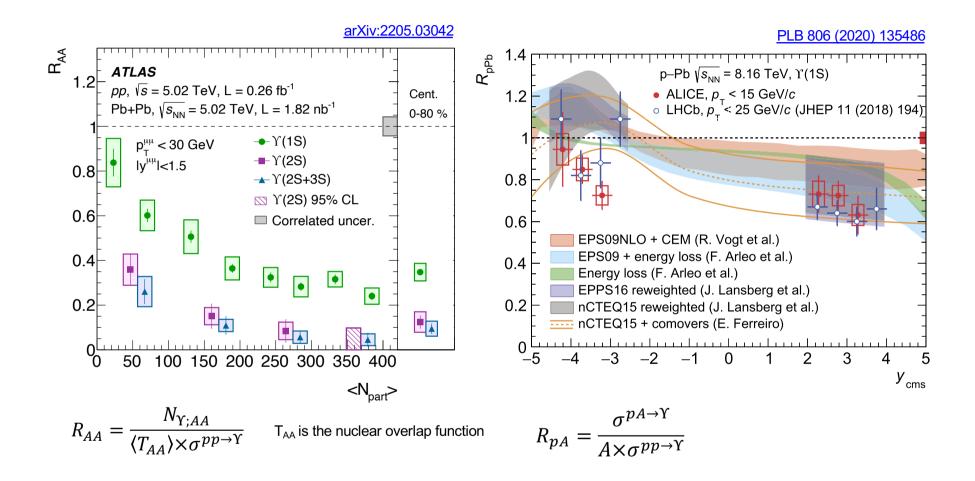
 $\Sigma_{n_{\rm ch}}\,\rho_{_{\rm T}}\,[{\rm GeV}]$

Pred. Data



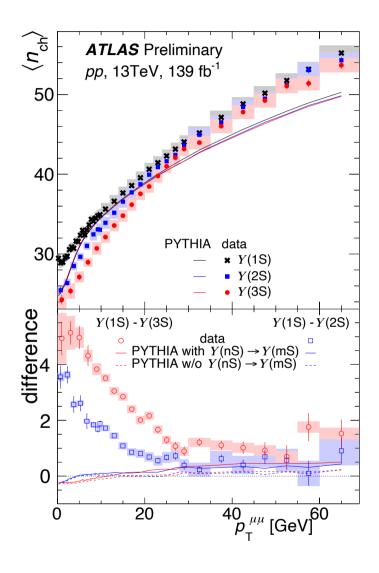


what is known about Y production and collectivity at the LHC?



- R_{AA}: nuclear modification factor compares **AA** to **pp**
- pA could give sense of influence of cold nuclear effects

naïve question



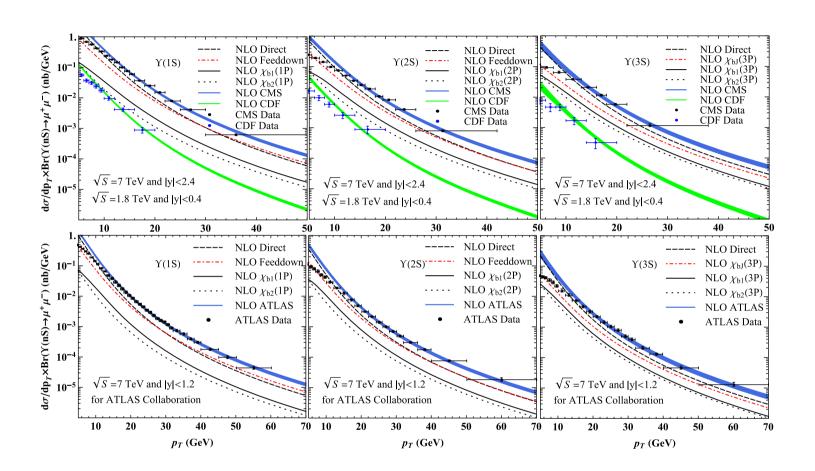
 could the observations correspond to an enhancement of nch for the Y(1S) ground state, rather than a suppression for the excited states?

- inclusive pp : $\langle n_{\rm ch} \rangle \approx 14$
- Drell-Yan 40 GeV $< m \le m_Z$ $\langle n_{\rm ch} \rangle = 24-28$
- Jets with leading particles $\langle n_{\rm ch} \rangle \approx 27$ $m < \frac{1}{2} m_{\rm Y}$

PLB 758 (2016) 67 EPJC 79 (2019) 666 JHEP 07 (2018) 032 JHEP 03 (2017) 157

Y(1S) appears consistent with these numbers, while Y(2S), Y(3S) lower, i.e. suppression of excited states

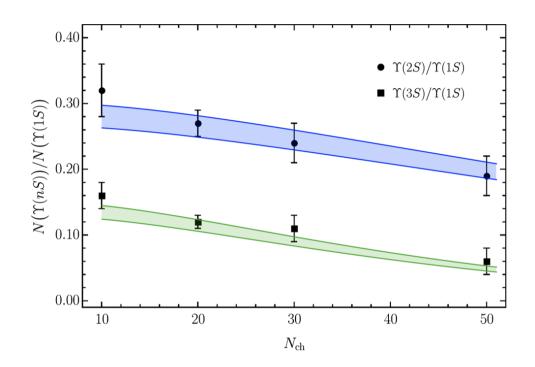
QCD cross section calculations



calculations show discrepancies with data, largest for higher Y(nS) and lower pT

Co-mover Interaction Model (CIM)

- CIM: quarkonia are broken by collisions with co-movers, i.e. final state particles with similar rapidities
- CIM typically used to explain
 p+A and A+A systems –
 matches CMS Y data



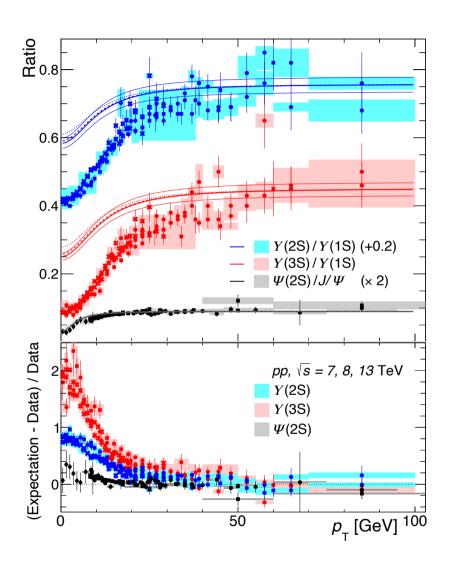
- CIM can be tested on new pp measurements to test if it can reproduce Y(nS)-Y(1S)
- cross sections
- multiplicity *n*ch
- kinematic distributions **P**T, Δφ, Δη

Quarkonia rations expected from mT-scaling

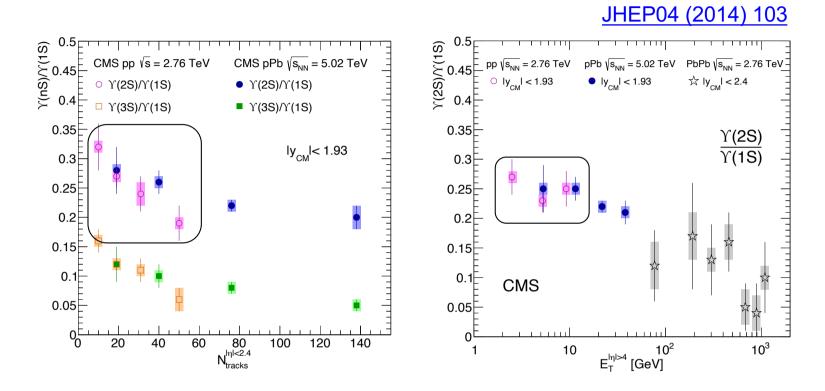
 mT-scaling allows one to define an expectation for the excited states relative to the ground states

$$\frac{\mathrm{d}^2 \sigma}{\mathrm{d}y \mathrm{d}m_{\mathrm{T}}} \propto \left(1 + \frac{m_{\mathrm{T}}}{nT}\right)^{-n}$$

- works well for light mesons at the LHC
- application to Y meson cross sections shows missing excited states at low pt
- Y(2S): missing factor of 1.6
- Y(3S): factor of 2.4

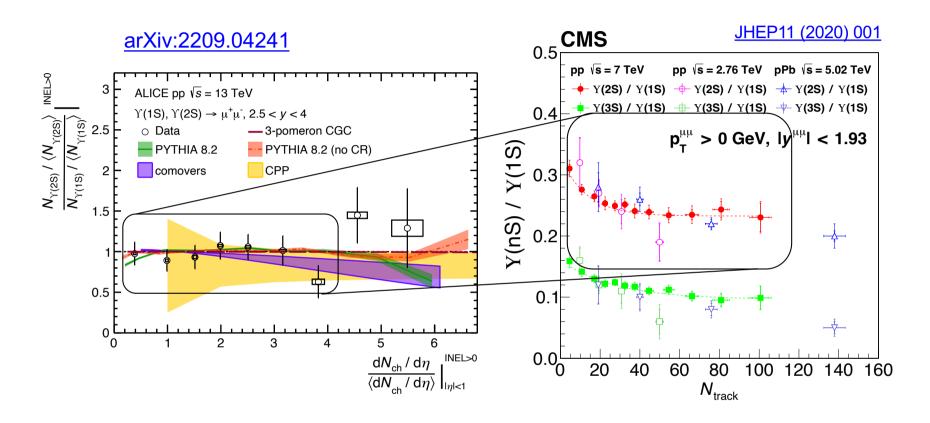


is there rapidity dependence?



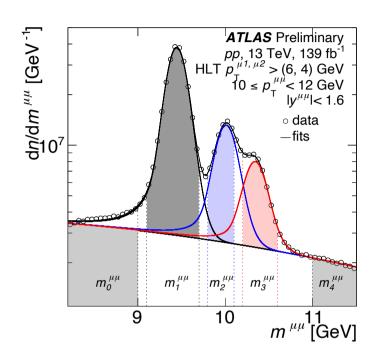
- CMS: dependence flattened when forward-midrapidity gap introduced
- (as noted in HP2018 summary: https://indico.cern.ch/event/634426/contributions/3003672/)
- could be due to loss of resolution ... ?

is there rapidity dependence?



- ALICE: forward Y(2S)/ Y(1S) vs tracks at midrapidity
- measurements do not clearly indicate rapidity dependence
- → Δη analysis should provide direct answer

Y(nS) and UE correlations — BG fit details



$$\begin{pmatrix} P(m_0^{\mu\mu}) \\ P(m_1^{\mu\mu}) \\ P(m_2^{\mu\mu}) \\ P(m_4^{\mu\mu}) \\ P(m_4^{\mu\mu}) \end{pmatrix} = \begin{pmatrix} 1 - f_{01} & f_{01} & 0 & 0 & 0 \\ k_1 (1 - s_1) & s_1 & 0 & 0 & (1 - k_1) (1 - s_1) \\ k_2 (1 - s_2 - f_{21} - f_{23}) & f_{21} & s_2 & f_{23} & (1 - k_2) (1 - s_2 - f_{21} - f_{23}) \\ k_3 (1 - s_3 - f_{32}) & 0 & f_{32} & s_3 & (1 - k_3) (1 - s_3 - f_{32}) \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_0 \\ P(\Upsilon(1S)) \\ P(\Upsilon(2S)) \\ P(\Upsilon(3S)) \\ P_4 \end{pmatrix}$$

$$k_n = \frac{\langle F_{\text{bkg}}(m) \rangle|_{m_4^{\mu\mu}} - \langle F_{\text{bkg}}(m) \rangle}{\langle F_{\text{bkg}}(m) \rangle|_{m_4^{\mu\mu}} - \langle F_{\text{bkg}}(m) \rangle}$$

 $P(m_n^{\mu\mu})$: measured distribution in mass interval $m_n^{\mu\mu}$

invert matrix to determine contributions coming from Y(nS) and from BG in the low (0) and high mass (4) BG intervals

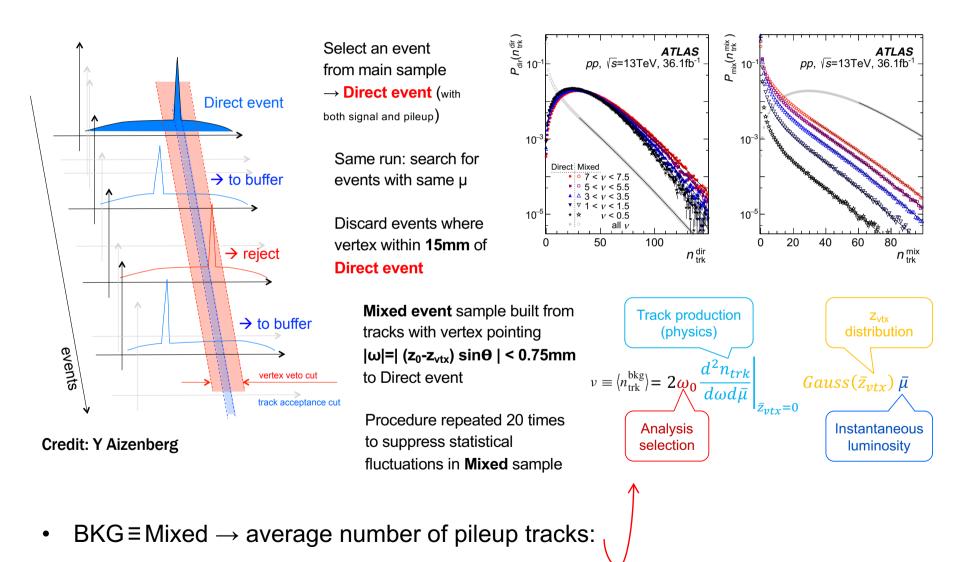
$$s_{n} = \frac{\int_{m_{n}^{\mu\mu}} N_{\Upsilon(nS)} F_{n}(m) dm}{\int_{m_{n}^{\mu\mu}} \operatorname{fit}(m) dm}$$

$$f_{nk} = \frac{\int_{m_{n}^{\mu\mu}} N_{\Upsilon(kS)} F_{k}(m) dm}{\int_{m_{n}^{\mu\mu}} \operatorname{fit}(m) dm}$$

$$k_{n} = \frac{\langle F_{\text{bkg}}(m) \rangle|_{m_{4}^{\mu\mu}} - \langle F_{\text{bkg}}(m) \rangle|_{m_{n}^{\mu\mu}}}{\langle F_{\text{bkg}}(m) \rangle|_{m_{4}^{\mu\mu}} - \langle F_{\text{bkg}}(m) \rangle|_{m_{0}^{\mu\mu}}}$$

correction for pileup

subtraction of pileup using event mixing technique EPJC 80 (2020) 64



31

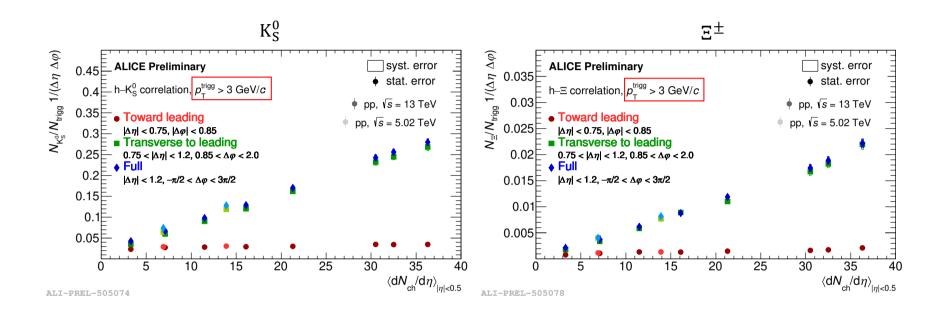
Y(nS) and UE correlations – systematics summary

	$p_{\rm T}^{\mu\mu} \le 4 {\rm GeV}$	$4 < p_{\mathrm{T}}^{\mu\mu} \le 12 \text{ GeV}$	$12 < p_{\rm T}^{\mu\mu} \le 30 \text{ GeV}$	$p_{\rm T}^{\mu\mu} > 30 \text{ GeV}$
$\Upsilon(1S)$	0.5 - 0.6	0.5 - 0.7	0.7 - 0.8	0.8 - 0.9
$\Upsilon(2S)$	0.6 - 0.6	0.5 - 0.7	0.7 - 0.8	0.8 - 1.0
$\Upsilon(3S)$	0.9 - 1.3	0.5 - 0.8	0.7 - 0.8	0.8 - 0.9
$\Upsilon(1S) - \Upsilon(2S)$	0.11 - 0.15	0.06 - 0.10	0.12 - 0.21	0.2 - 0.5
$\Upsilon(1S) - \Upsilon(3S)$	0.6 - 0.9	0.14 - 0.36	0.14 - 0.15	0.16 - 0.19

Table 1: Systematic uncertainties for measurements of $\langle n_{\rm ch} \rangle$ and their differences for different $\Upsilon(nS)$ states and for the difference between $\langle n_{\rm ch} \rangle$ measured for $\Upsilon(1S) - \Upsilon(nS)$. The values are the number of charged particles with $0.5 \le p_{\rm T} < 10$ GeV and $|\eta| < 2.5$.

shown for *n*_{ch} but propagated to all quantities

strange enhancement



- enhancement of strange hadrons is one of the signature pp collectivity results
- ALICE measurement (<u>ALI-PREL-505074</u>) shows strange hadron yields from out-of-jet and full samples increases with multiplicity and are consistent with each other
- toward jet yields do not depend on multiplicity
- → could suggest that it is an underlying event effect