

# Small- $x$ Helicity Phenomenology

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# Proton Spin Puzzle

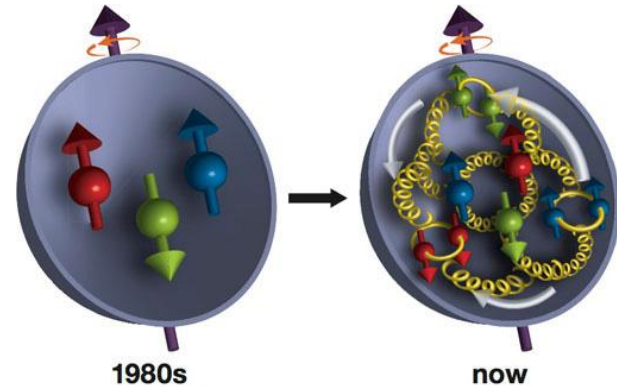
Jaffe-Manohar Spin Sum Rule:

$$\frac{1}{2} = S_q + L_q + S_g + L_g$$

$S_{q,g}$  = Helicity of quarks and gluons

$L_{q,g}$  = Orbital angular momentum

$S_q \sim 30\%$  of proton spin!



# Quark Helicity Parton Distribution Functions

$$S_q(Q^2) = \frac{1}{2} \int_0^1 dx \sum_q (\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2))$$

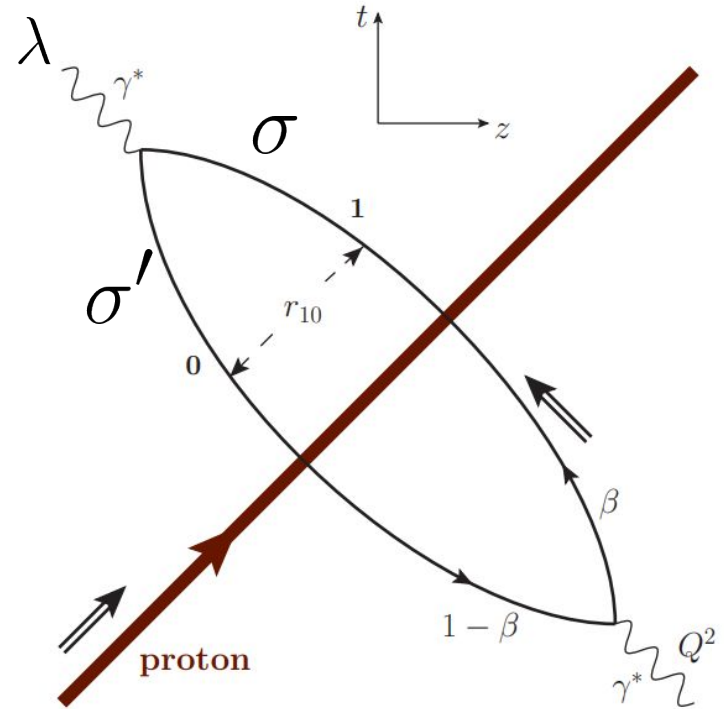
Helicity PDFs:

$$\Delta q = \text{[Diagram: Blue circle with red dot and right arrow]} - \text{[Diagram: Blue circle with red dot and left arrow]}$$

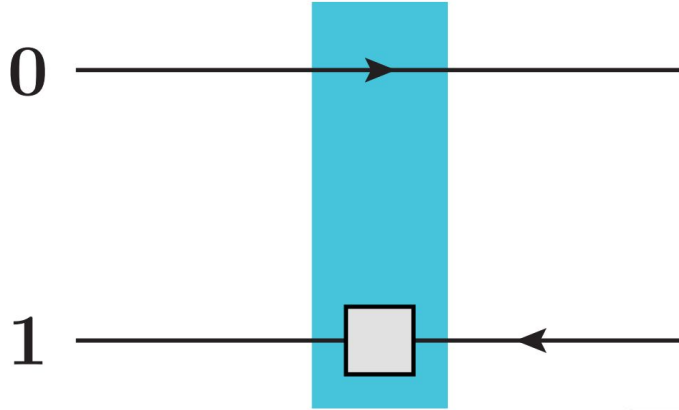
- $Q^2$  = resolution at which we probe the proton
- Bjorken  $x \sim \frac{1}{s}$ . We need theory to extrapolate to  $x=0$

# (Polarized) DIS in the (Polarized) Dipole Picture

$$g_1 \propto |\psi|^2 \otimes (Q + 2G_2)$$



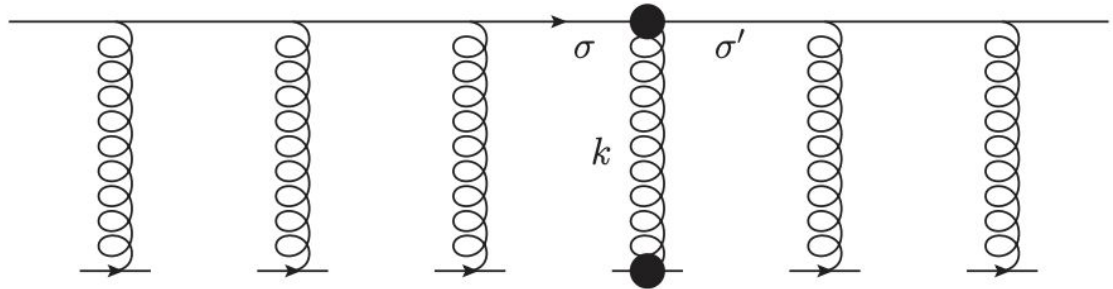
# (Polarized) DIS in the (Polarized) Dipole Picture



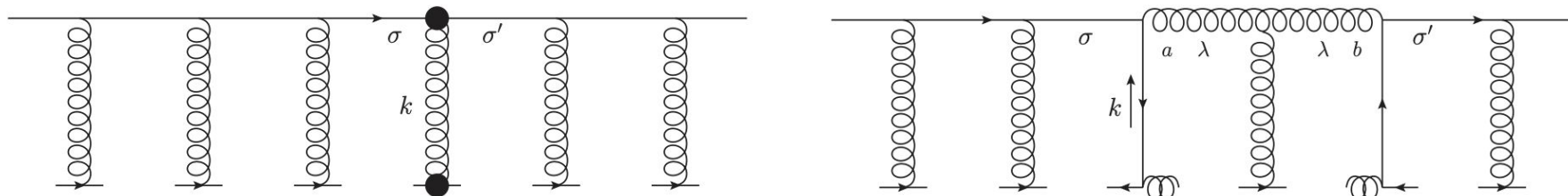
- Quark line undergoes one extra helicity exchange, which is **sub-eikonal**

- In pDIS, the electron and proton have their helicity specified
- Cross-section now dependent on **Polarized Dipole Amplitudes:**

$$Q_q, G_2, \tilde{G}$$



# Polarized Wilson Lines



$$\vec{\mu} \cdot \vec{B}$$

$$\mu B_z \sim F_{12}$$

• Chromo-magnetic field

$$\bar{\psi} \gamma^+ \gamma^5 \psi$$

• Axial Current

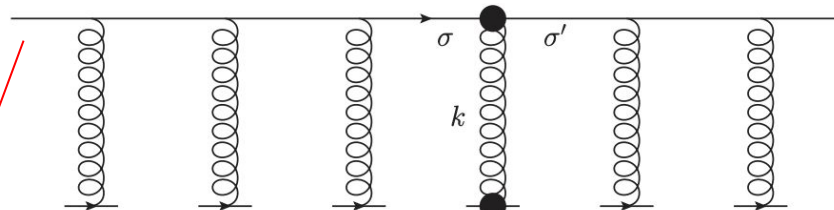
Polarized Dipole Amplitudes:

“ $Q_q$ ”

, “ $\tilde{G}$ ”

# Polarized Wilson Lines

- Quark propagator



$$\int \frac{dk^+}{2\pi} e^{ikx} \frac{\not{k}}{2k^+k^- - k_\perp^2}$$

- Sub-eikonal phase expansion

$A_\perp$

- Polarized gluon vertex

$$e^{-ix^- \frac{k_\perp^2}{2k^-}} \approx 1 - ix^- \frac{k_\perp^2}{2k^-} \Rightarrow \partial_\perp^2 \rightarrow D_\perp^2 \rightarrow "G_2"$$

# Calculating Helicity Distributions

$$\Delta q + \Delta \bar{q} = \frac{1}{N_c} \int_0^{\eta_{max}} d\eta \int_{s_{10}^{min}}^{\eta} ds_{10} \frac{1}{\alpha_s(s_{10})} (Q_q(s_{10}, \eta) + 2G_2(s_{10}, \eta))$$

- We incorporate running coupling that runs with size of the dipole
- $\eta \sim$  Longitudinal momentum
- $s_{10} \sim$  Transverse separation of Dipole



# Large Nc&Nf Helicity Evolution

In the large Nc&Nf, Nc/Nf fixed limit, the evolution equations for the polarized dipole amplitudes close:

$$Q_q(s_{10}, \eta) = Q_q^{(0)}(s_{10}, \eta) + \int_{s_{10}+y_0}^{\eta} d\eta' \int_{s_{10}}^{\eta'-y_0} ds_{21} \left[ Q_q(s_{21}, \eta') + 2\tilde{G}(s_{21}, \eta') + 2\tilde{\Gamma}_{s_{10}, s_{21}, \eta'} \right. \\ \left. - \bar{\Gamma}_f(s_{10}, s_{21}, \eta') + 2G_2(s_{21}, \eta') + 2\Gamma_2(s_{10}, s_{21}, \eta') \right] \\ + \frac{1}{2} \int_{y_0}^{\eta} d\eta' \int_{\max\{0, s_{10}+\eta'-\eta\}}^{\eta'-y_0} ds_{21} \left[ Q_q(s_{21}, \eta') + 2G_2(s_{21}, \eta') \right]$$

+ 9 more

- 5 Polarized dipole amplitudes mix under evolution:  $Q_{u,d,s}, \tilde{G}, G_2$
- With 5 auxiliary dipoles:  $\Gamma_{u,d,s}, \tilde{\Gamma}, \Gamma_2$  - which impose lifetime ordering
- Small-x cutoff,  $y_0 \propto \ln 1/x_0$

# Large Nc&Nf Helicity Evolution

- **5 Polarized dipole amplitudes** mix under evolution:  $Q_{u,d,s}, \tilde{G}, G_2$
- With 5 auxiliary dipoles:  $\Gamma_{u,d,s}, \tilde{\Gamma}, \Gamma_2$
- For a total of 10 equations that form a **closed system**
- Undetermined initial conditions:  $Q_{u,d,s}^{(0)}, \tilde{G}^{(0)}, G_2^{(0)}$

Recap:

$$\frac{1}{2} = S_q + L_q + S_g + L_g$$

$$S_q(Q^2) = \frac{1}{2} \int_0^1 dx \sum_q (\Delta q(x, Q^2) + \Delta \bar{q}(x, Q^2))$$

$$S_g(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

$$\Delta q + \Delta \bar{q} = \frac{1}{N_c} \int_0^{\eta_{max}} d\eta \int_{s_{10}^{min}}^{\eta} ds_{10} \frac{1}{\alpha_s(s_{10})} (Q_q(s_{10}, \eta) + 2G_2(s_{10}, \eta))$$

$$\Delta G(x, Q^2) = \frac{2N_c}{\alpha_s(Q^2)} G_2 \left( \sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{\Lambda^2}, \sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{x\Lambda^2} \right)$$

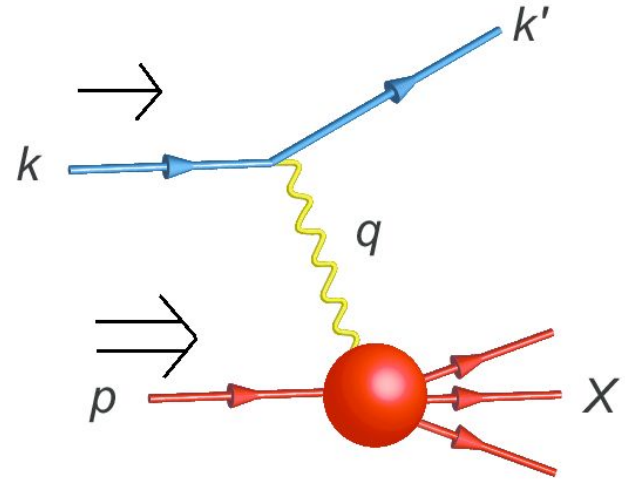
Large  $N_c$  &  $N_f$  Helicity Evolution

$$Q_q^{(0)}, \tilde{G}^{(0)}, G_2^{(0)}$$

# Observables - Double Spin Asymmetries in DIS

$$A_{||} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}} \propto A_1 \propto g_1^{p,n}$$

- $\uparrow$  ( $\downarrow$ ) is positive (negative) helicity electron
- $\uparrow\uparrow$  ( $\downarrow\downarrow$ ) is positive (negative) helicity proton
- $A_1$  is virtual photoproduction asymmetry



# Describing Observables - pDIS

What enters into observables are linear combinations of hPDFs

$$\Delta q^+ = \Delta q + \Delta \bar{q}$$

$$\Delta q^- = \Delta q - \Delta \bar{q}$$

- **Three** relevant hPDFs in DIS:  $\Delta u^+$ ,  $\Delta d^+$ ,  $\Delta s^+$ , involving **five** amplitudes
- Data exist for **two** observables that contain these hPDFs in linearly independent combinations:  $g_1^p$  and  $g_1^n$

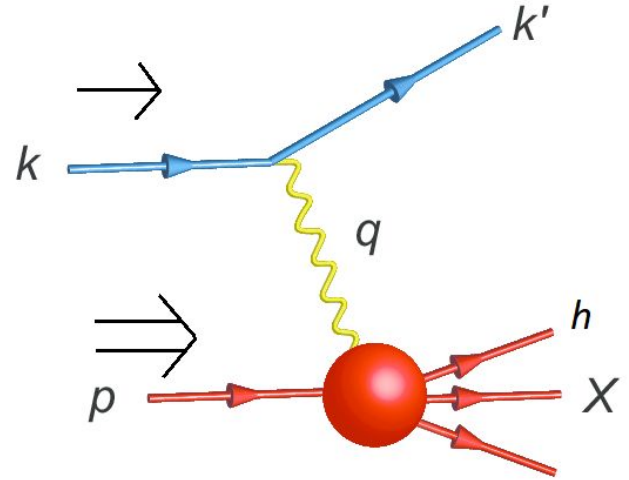
$$g_1^p(x, Q^2) = \frac{1}{2} \sum_q Z_q^2 \Delta q^+(x, Q^2)$$

- $Z_q$  is the quark charge fraction

# Observables - Double Spin Asymmetries in SIDIS

$$A_{||}(z) = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}} \propto g_1^h(z)$$

- $h$  is the tagged hadron
- $z$  is the momentum fraction of the virtual photon carried by the tagged hadron



# Describing Observables - pSIDIS

- 2 observables are not enough to describe 3 hPDFs.
- Expand our horizons to Semi-Inclusive DIS - all hPDFs are relevant here, both singlet,  $\Delta q^+$  and non-singlet,  $\Delta q^-$
- **Non-singlet distributions obey their own small-x evolution that has been solved**

$$\Delta q^- = \frac{N_c}{2\pi^3} \int d\eta \int ds_{10} Q_q^{NS}(s_{10}, \eta)$$

- $Q_q^{NS}$  is the non-singlet Polarized Dipole Amplitude - obeys its own evolution equation
- pSIDIS grants us access to the semi-inclusive, spin dependent structure functions  $g_1^h$

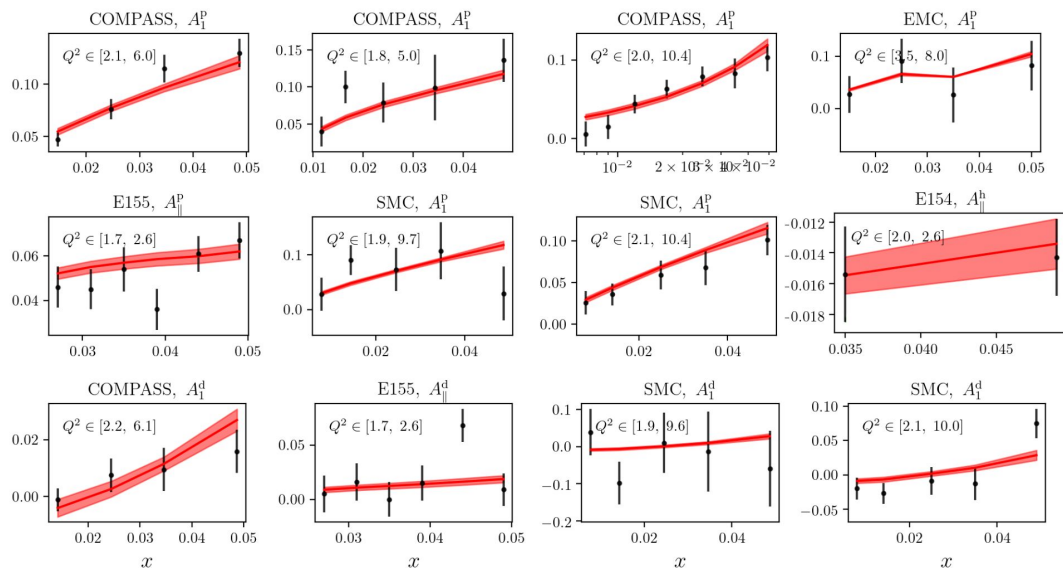
# $g_1^h$ Structure Functions

$$g_1^h(x, z, Q^2) = \frac{1}{2} \sum_q Z_q^2 \Delta q(x, z, Q^2) D_q^h(z, Q^2)$$

- $D_q^h$  are fragmentation functions - giving the probability quark  $q$  fragments into hadron  $h$
- $z$  Is the fraction of the virtual photons momentum carried by the hadron
- The flavour hPDF is obtained via  $\Delta q = \frac{1}{2}(\Delta q^+ + \Delta q^-)$
- In pSIDIS, we are able to scatter on 2 targets (proton, neutron), tag 2 outgoing hadrons (pion, kaon) that each have 2 charges -  $2 \times 2 \times 2 = 8$  new observables



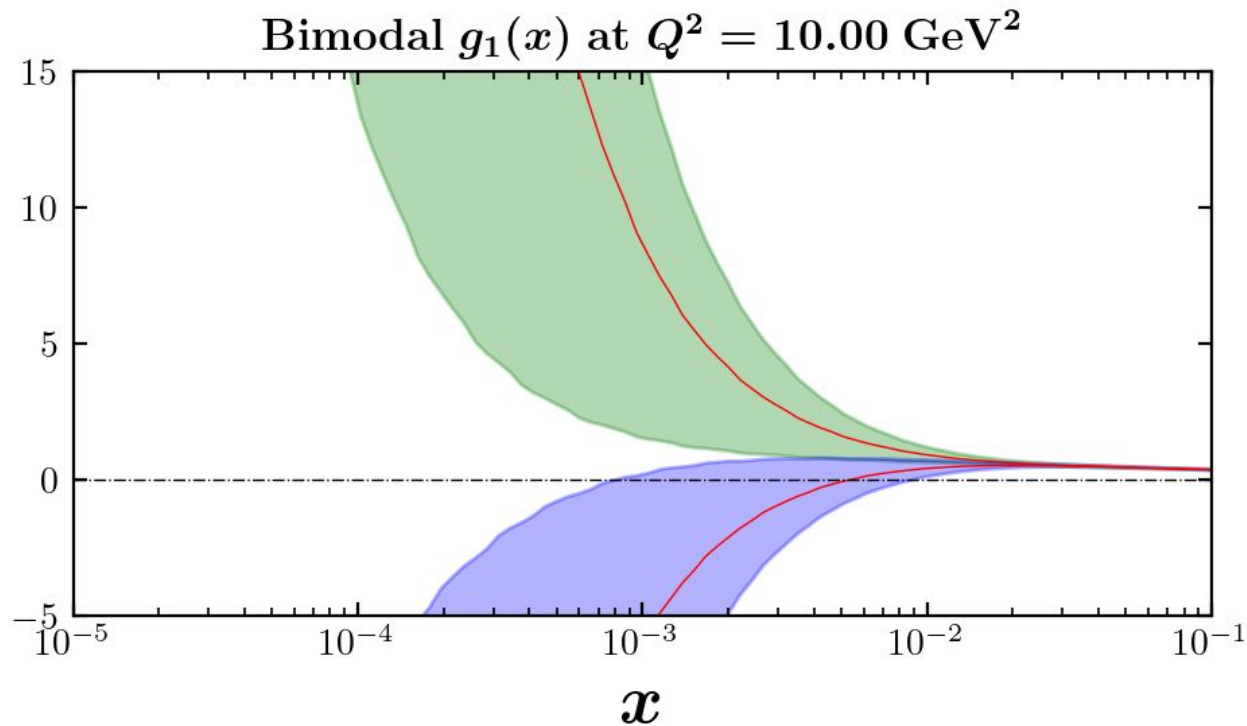
# Global fit of DIS - Data vs Theory



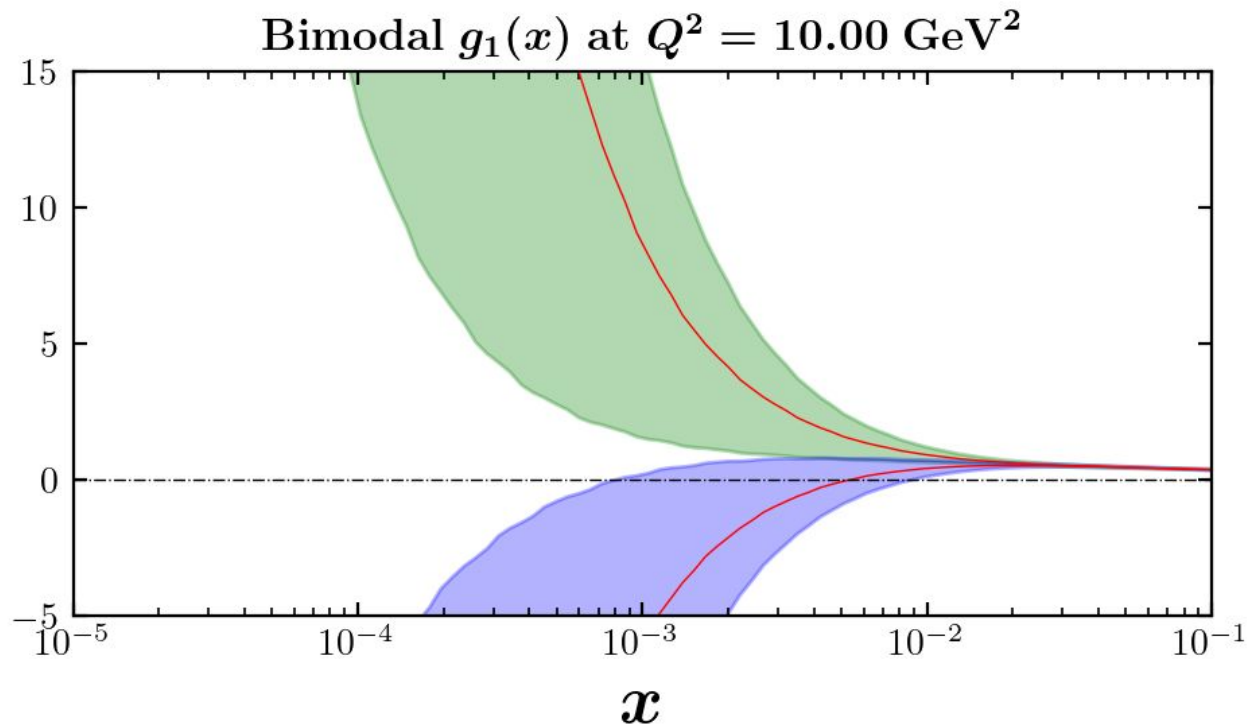
- Red curves - our theory
- Black dots - data
  - COMPASS
  - EMC
  - SMC
  - SLAC
  - HERMES
- Preliminary results

- Cut of  $0.005 < x < 0.1$
- Cut of  $1.69 \text{ GeV}^2 < Q^2 < 10.4 \text{ GeV}^2$
- Cut of  $0.2 < z < 1.0$
- Describing 234 data points
- With a  $\chi^2/npts = 1.04$

# (Preliminary) Extraction of $g_1^p$



# (Preliminary) Extraction of $g_1^p$



$\tilde{G}$  not  
sensitive to  
large- $x$  data,  
but responsible  
for small- $x$   
behaviour

# Constraining the rest of the Polarized Dipole Amplitudes

$$g_1^{p,n} \sim Q_u, Q_d, Q_s, G_2$$

$$g_1^h \sim Q_q, G_2, Q_q^{NS}$$

$$pp \rightarrow jets \sim G_2, \tilde{G}$$

- 2 observables, 4 polarized dipole amplitudes. Under constrained system
- 8 new observables, 3 new polarized dipole amplitudes. Exactly constrained - but  $\tilde{G}$  does not enter directly into observables
- Particle production might provide final constraints

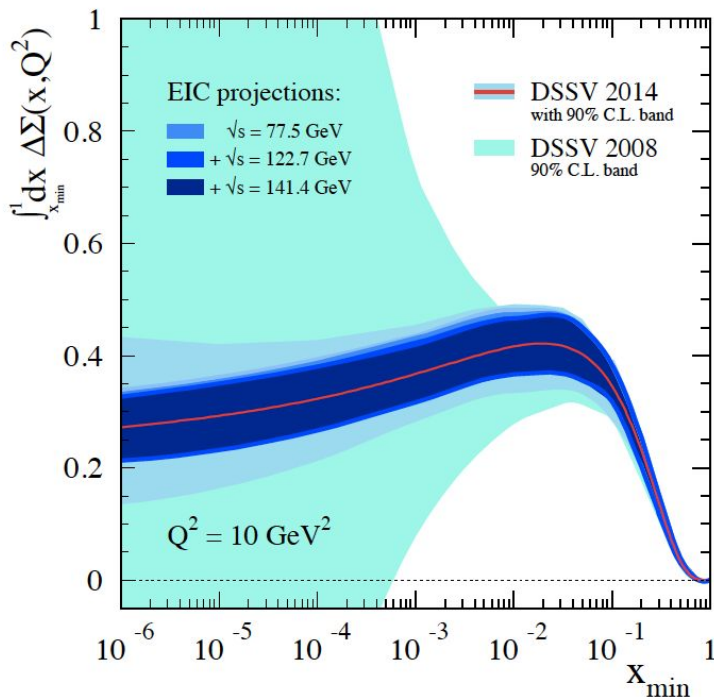
# Conclusions

- In order to resolve the spin puzzle, the small- $x$  behaviour of the hPDFs need to be understood
- This is accomplished using small- $x$  evolution
- Along with fitting to data
- Potentially a significant amount of spin is hiding in the small- $x$  region
- More work needs to be done to constrain small- $x$  behavior of the various polarized dipoles - especially  $G_2$  and  $\tilde{G}$
- Could be constrained by studying particle production in  $pp$  collisions as well as smaller- $x$  EIC data

# Backup Slides

# Quark hPDF - DGLAP extraction

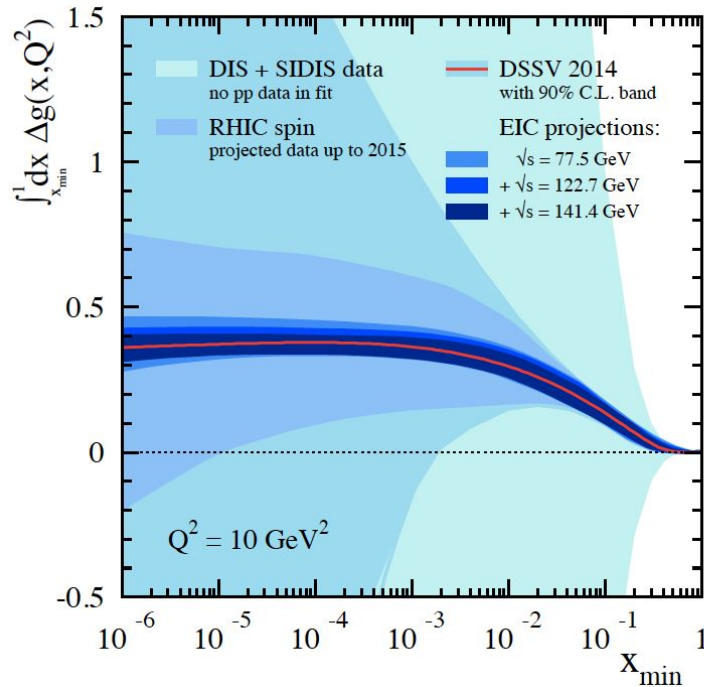
2 x  
(quark  
spin)



$$\Delta\Sigma = \sum_q (\Delta q + \Delta \bar{q})$$

- E. Aschenauer et al, [arXiv:1509.06489 \[hep-ph\]](https://arxiv.org/abs/1509.06489), (DSSV = de Florian, Sassot, Stratmann, Vogelsang, DGLAP-based helicity PDF extraction from data)
- Large uncertainty at small- $x$ !

# Gluon Helicity Parton Distributions Function



$$S_g(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

$\Delta G$  = Gluon Helicity PDF

- Uncertainty consistently blows up when extrapolating beyond data



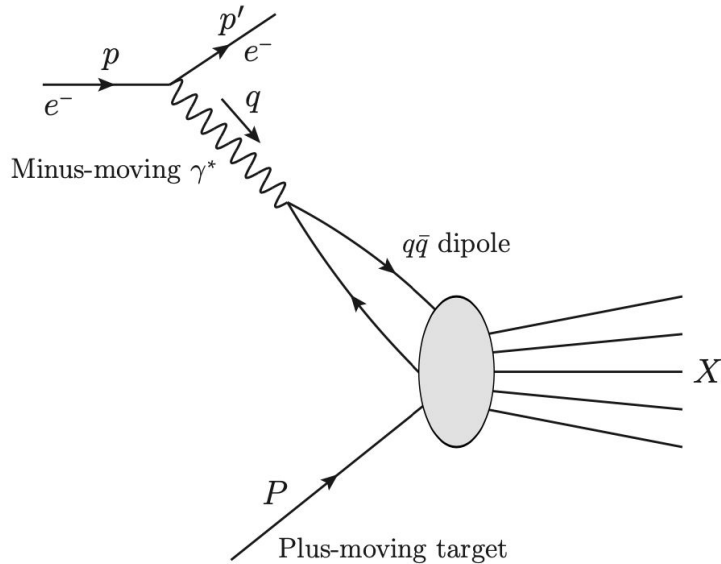
# The Plan

Any complete description of quark and gluon helicity needs to

- Describe existing data ( $5 \times 10^{-3} < x < 0.7$ )
- Predict future, e.g EIC, data ( $4 \times 10^{-3} < x < 5 \times 10^{-3}$ )
- Compare with said data
- Extrapolate down to  $x = 0$
- While maintaining good control over theoretical uncertainty

# Deep-Inelastic Scattering (DIS)

Probing the proton at small  $x$



- Electron of momentum  $p$  scatters off proton of momentum  $P$
- Transverse size given by virtuality of photon:

$$\frac{1}{x_{\perp}^2} \propto Q^2 = -q^2$$

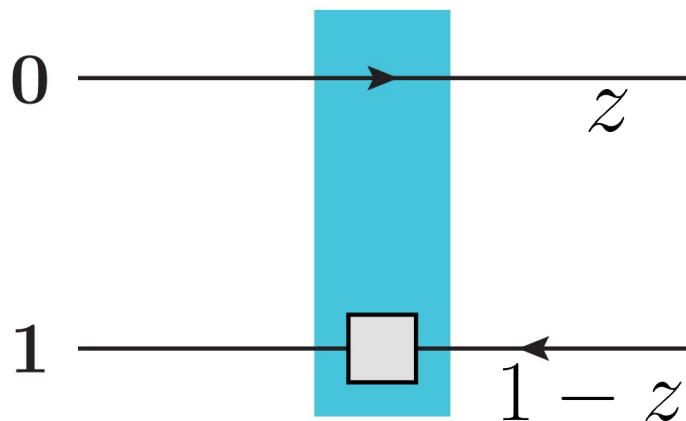
- Bjorken- $x$ :  $x = \frac{Q^2}{2P \cdot q} \approx \frac{Q^2}{s}$

## Calculating Helicity Distributions

$$\Delta G(x, Q^2) = \frac{2N_c}{\alpha_s(Q^2)} G_2 \left( \sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{\Lambda^2}, \sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{x\Lambda^2} \right)$$

- Jaffe-Manohar Gluon Helicity Distribution
- $\Lambda^2$  Infrared cutoff

# Polarized Dipole Amplitude - Degrees of Freedom



$$Q_q(s_{10}, \eta)$$

Polarized Dipole Amplitudes are functions of

- Transverse separation:

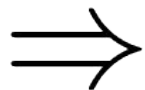
$$x_{10}^2 = (\underline{x}_1 - \underline{x}_0)^2$$

- Momentum Fraction times center of mass energy:  $zS$
- Rescaled variables:

$$\eta = \sqrt{\frac{N_c}{2\pi}} \ln \frac{zs}{\Lambda^2} \quad s_{10} = \sqrt{\frac{N_c}{2\pi}} \ln \frac{1}{x_{10}^2 \Lambda^2}$$

# Helicity Evolution

Using Light-Cone Operator Treatment, we need to resum all gluon exchanges that exchange helicity information



Resumming all terms containing:

$$\alpha_s \int_x^1 \frac{dz}{z} \int_{1/s}^{1/Q^2} \frac{d^2 x_{21}}{x_{21}^2}$$

Resum double log  
(DLA) terms:

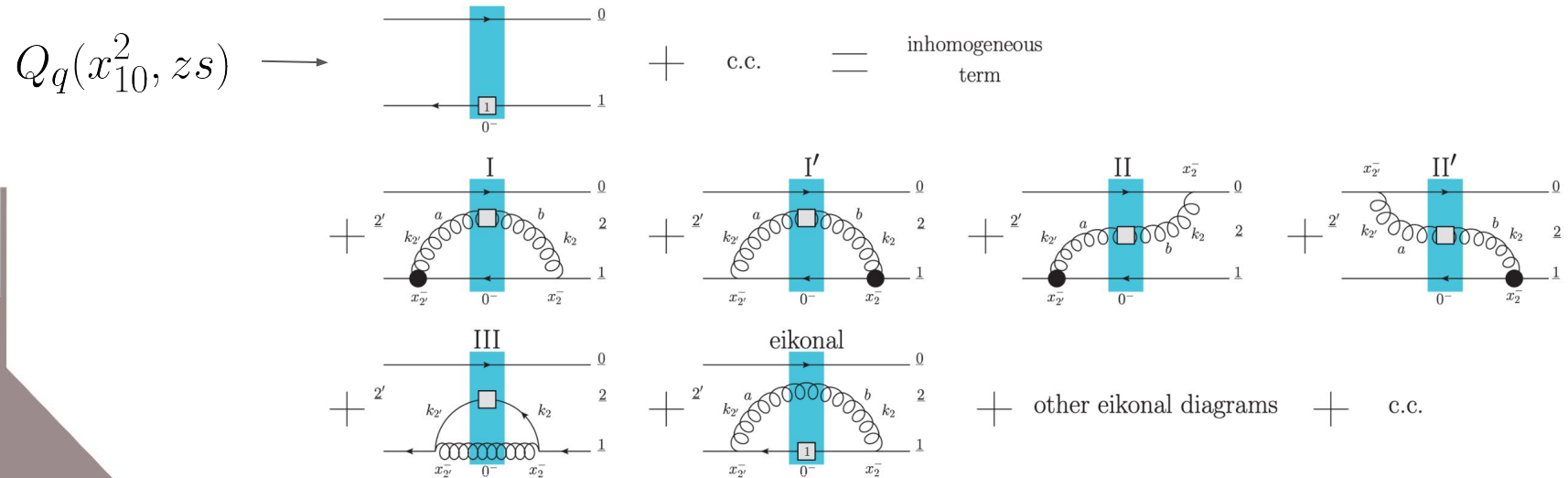
$$\alpha_s \ln^2(1/x)$$

Longitudinal part.  
Present in un-polarized  
evolution

Transverse part. UV  
exactly cancelled in  
un-polarized evolution

# Helicity Evolution

- Relate Polarized Dipole Amplitude to themselves at higher energies by resumming emission diagrams - resumming Double Log (DLA) contributions:  $\alpha_s \ln^2(1/x)$



# Sub-eikonal Expansion

- Expansion in energy or in  $x$

$$1/x,$$

Eikonal

$$F_1, F_2$$

$$x^0,$$

Sub-Eikonal

$$g_1^{p,n}, \Delta q, \Delta \bar{q}$$

$$x^1$$

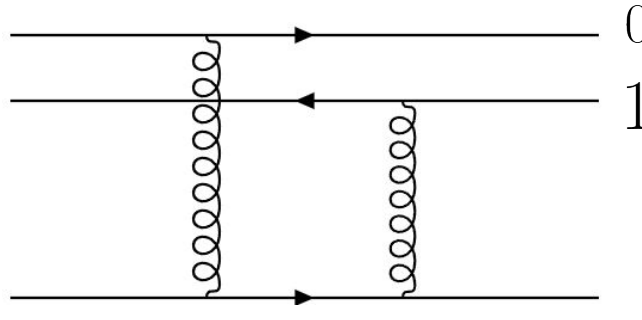
Sub-Sub-Eikonal

Transversity

- No eikonal terms contain any helicity information - Wilson lines are helicity independent
- Must calculate sub-eikonal terms to access helicity

# Inhomogeneous term

The inhomogeneous term is given by a Born-inspired ansatz:



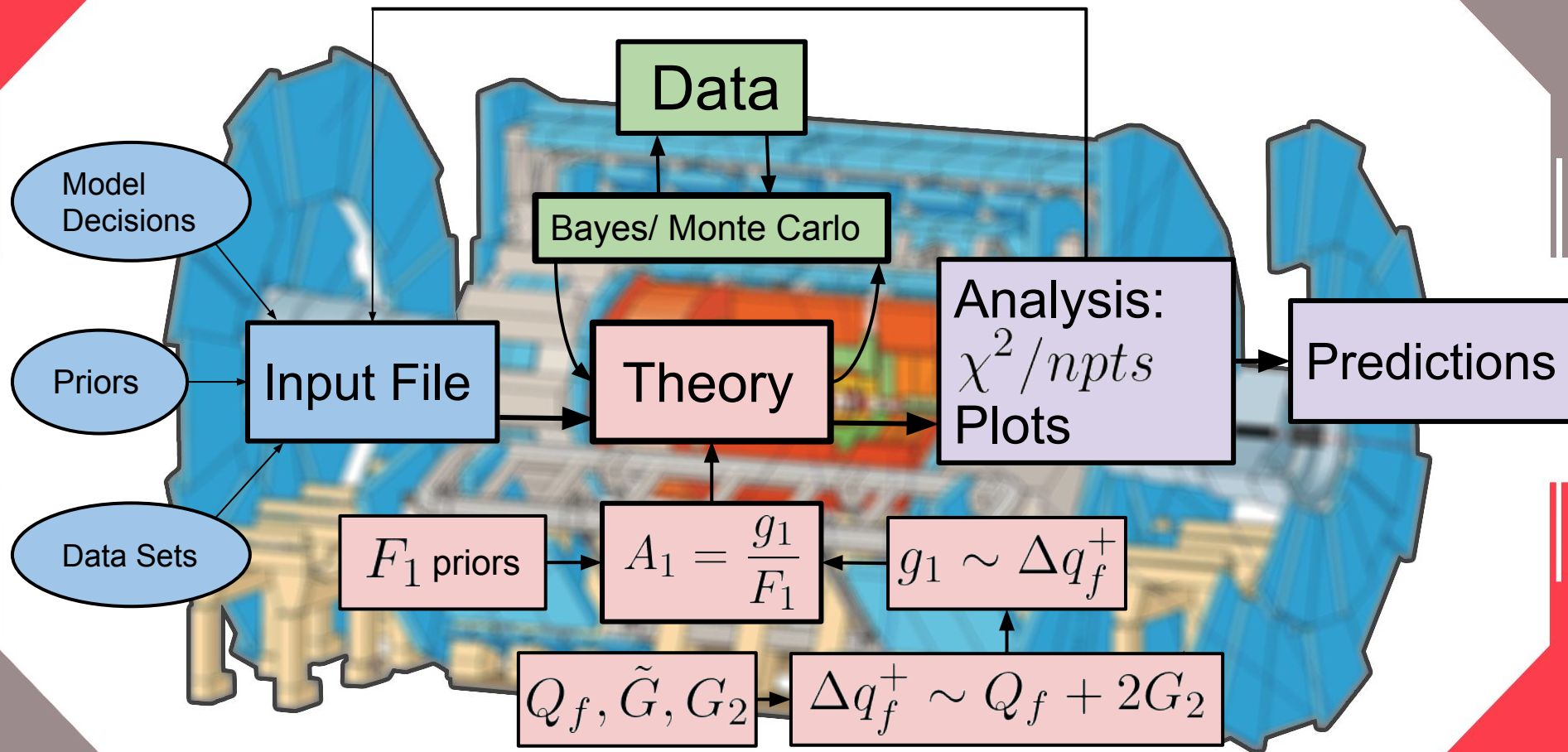
The diagram shows three horizontal lines representing particles. The top line is labeled '0' and has an arrow pointing right. The middle line is labeled '1' and has an arrow pointing left. The bottom line has an arrow pointing right. A vertical chain of eight circles connects the top and middle lines. A second vertical chain of eight circles connects the middle and bottom lines.

$$\propto \int_0^s \frac{dk_{\perp}^2}{k_{\perp}^2} (1 - e^{-\underline{k} \cdot \underline{x}_{10}}) = \pi \ln(s x_{10}^2)$$
$$\propto \eta - s_{10}$$

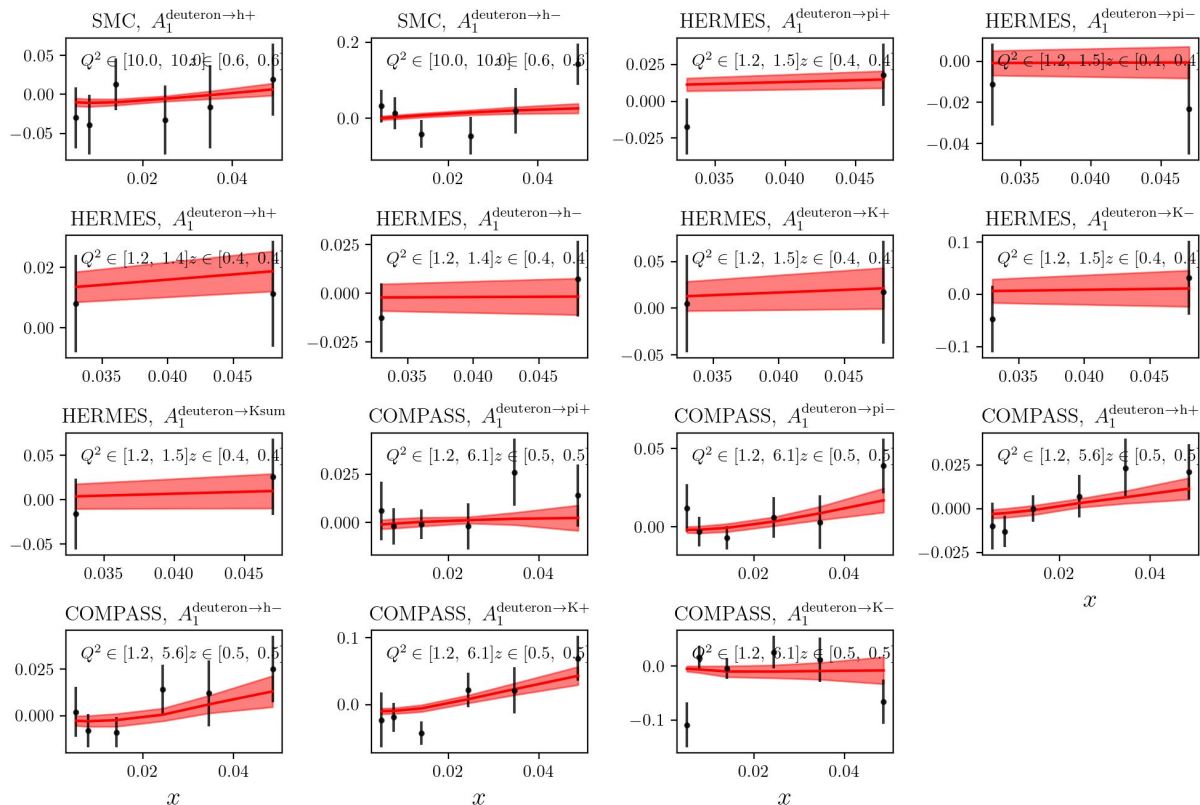
$$\Gamma_q^{(0)} = Q_q^{(0)} = a\eta + bs_{10} + c$$

- Same form of the other Dipole Amplitudes
- Parameters a,b,c need to be extracted from data

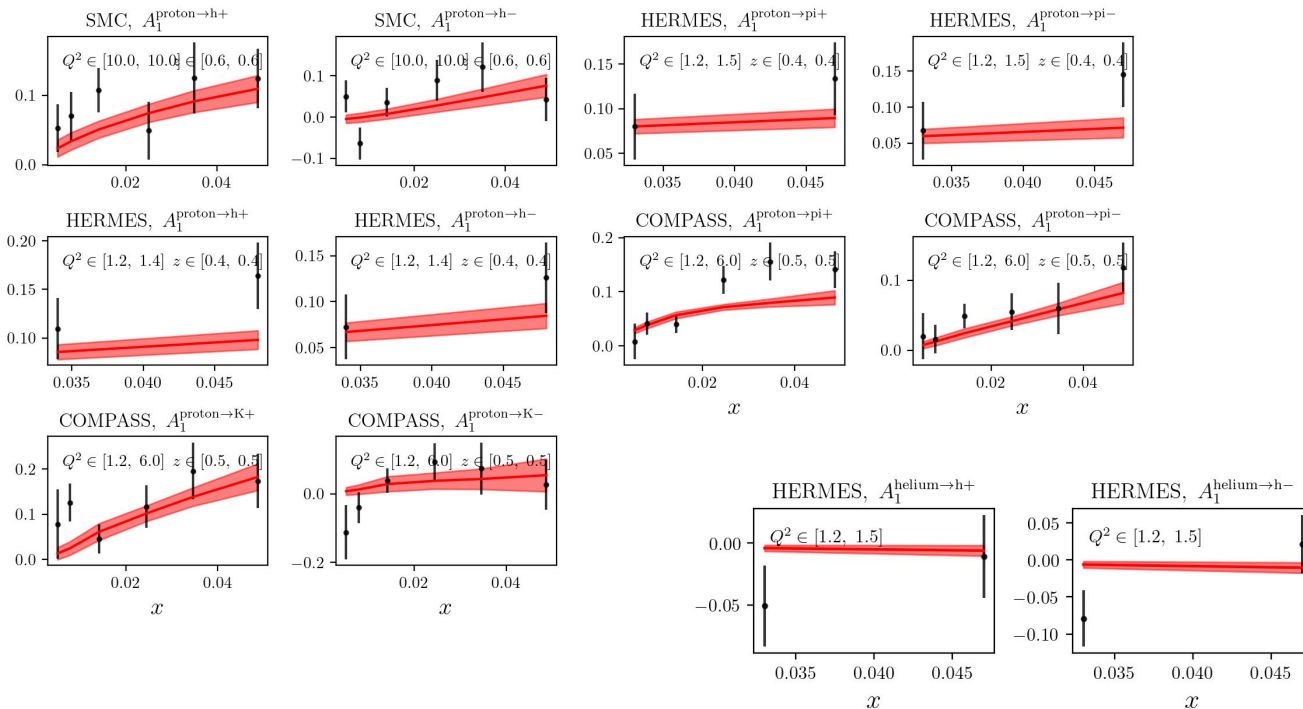




# Global fit of SIDIS - Data vs Theory

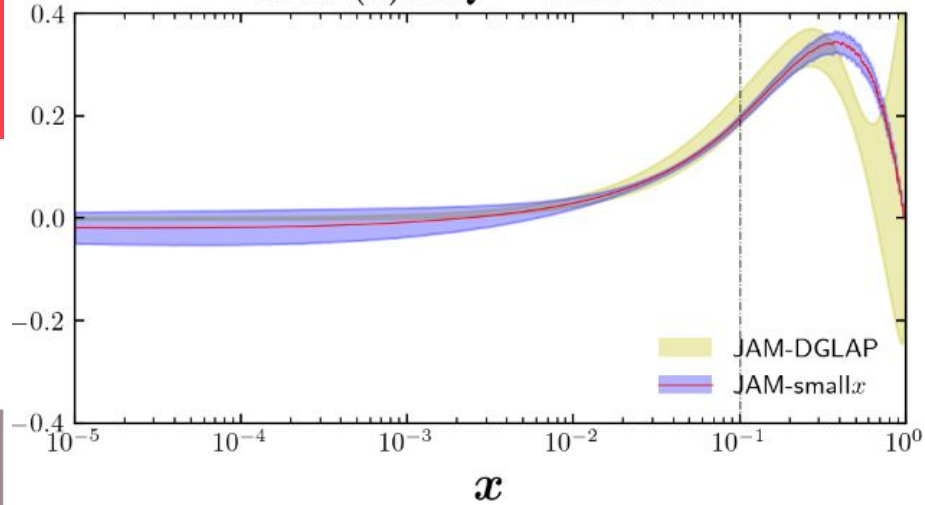


# Fitting SIDIS - Data vs Theory

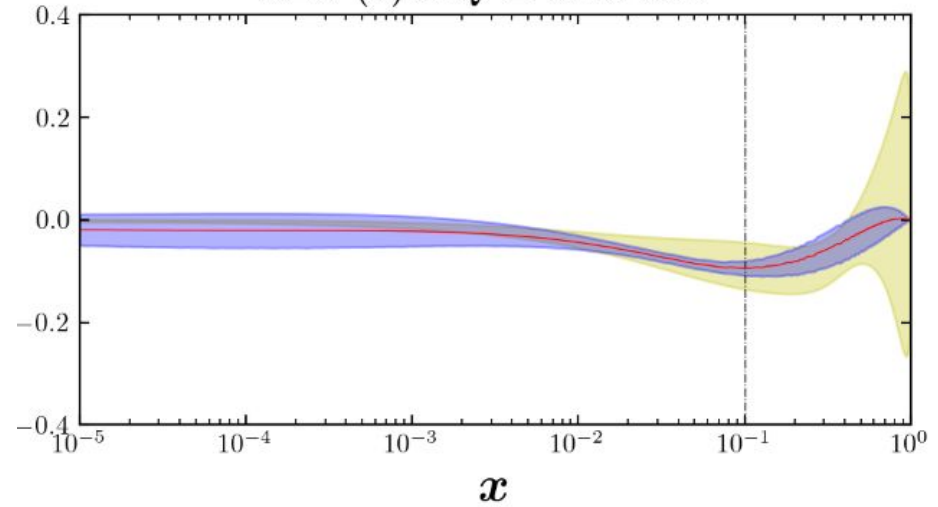


# hPDFs - (Preliminary)

$x\Delta u^+(x)$  at  $Q^2 = 10.00 \text{ GeV}^2$

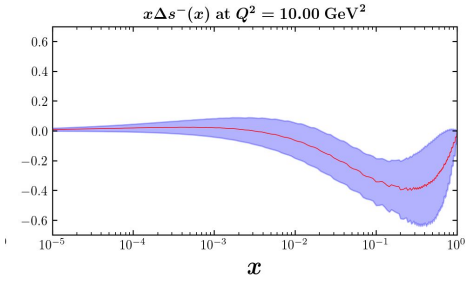
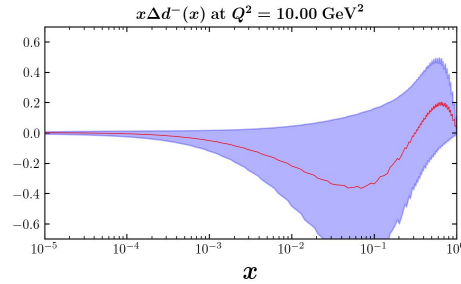
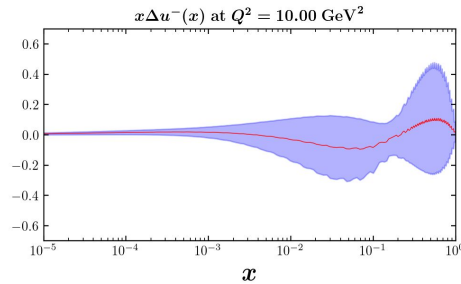
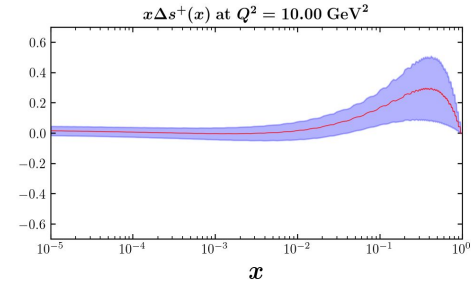
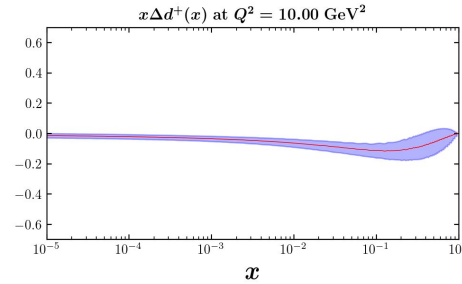
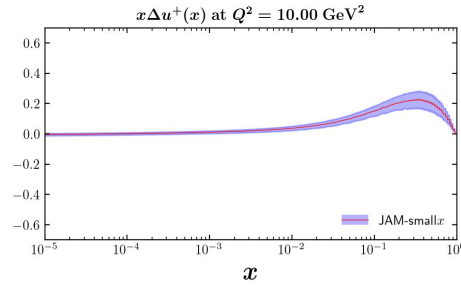


$x\Delta d^+(x)$  at  $Q^2 = 10.00 \text{ GeV}^2$



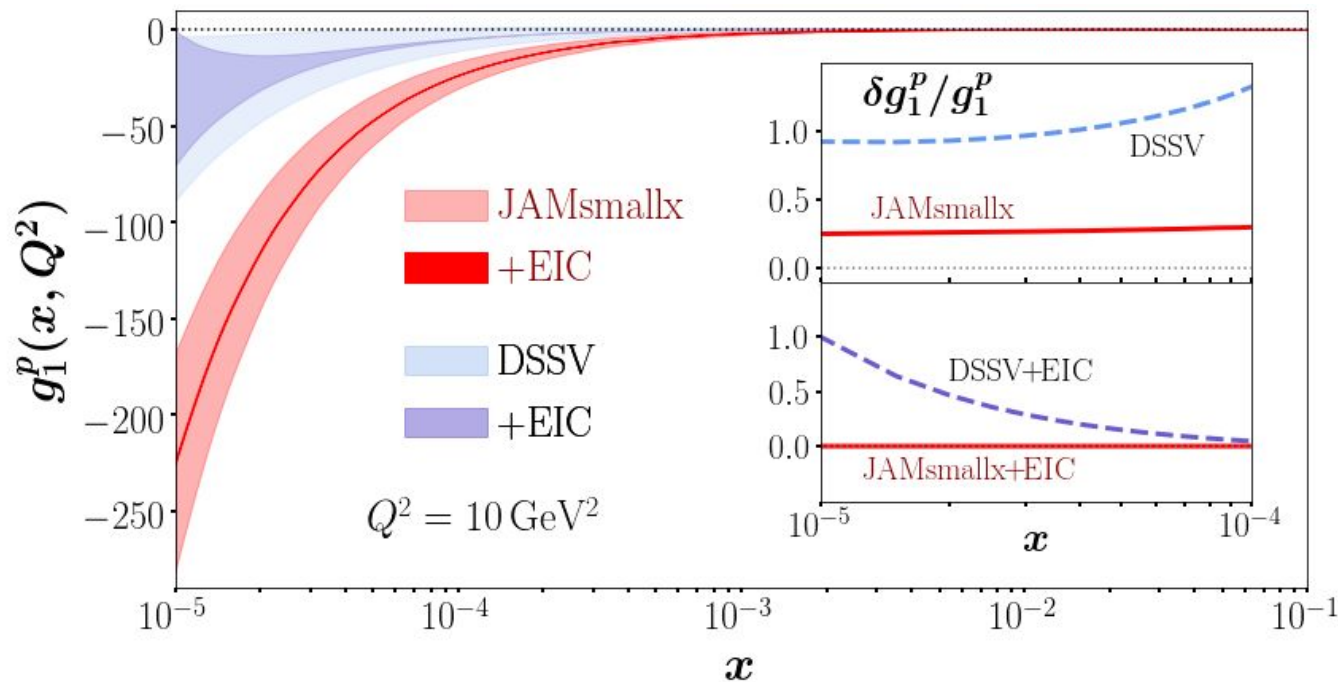
- DIS only: Strange distribution set to zero

# hPDFs - Preliminary



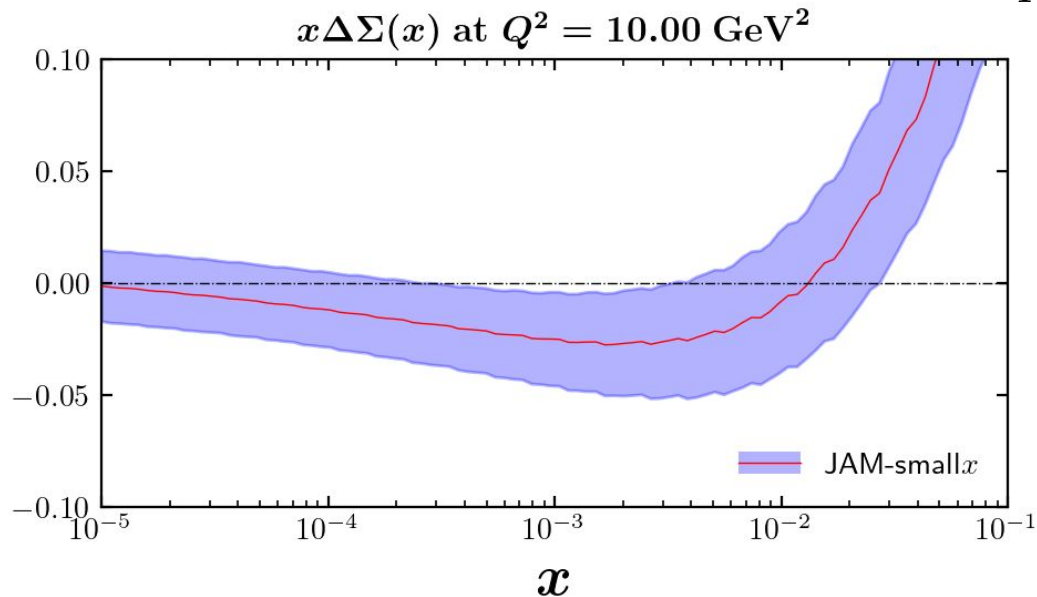
- Old version of evolution

# (Preliminary) Extraction of $g_1^p$

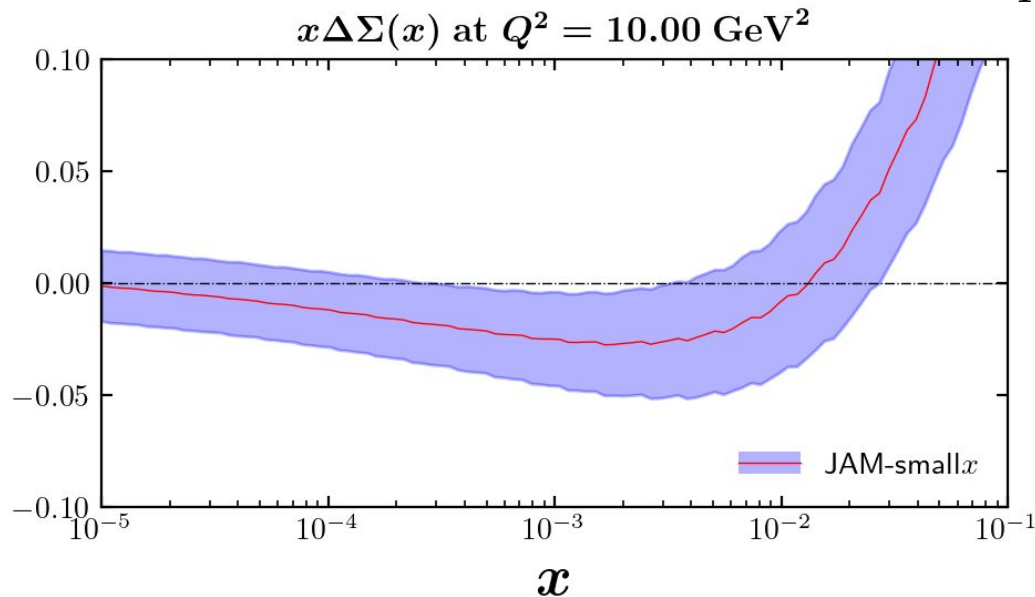


- DSSV uses DGLAP - rational function extrapolation of  $x$
- We use small- $x$  helicity evolution to predict the  $x$  behaviour
- Leads to control over uncertainty

# Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$



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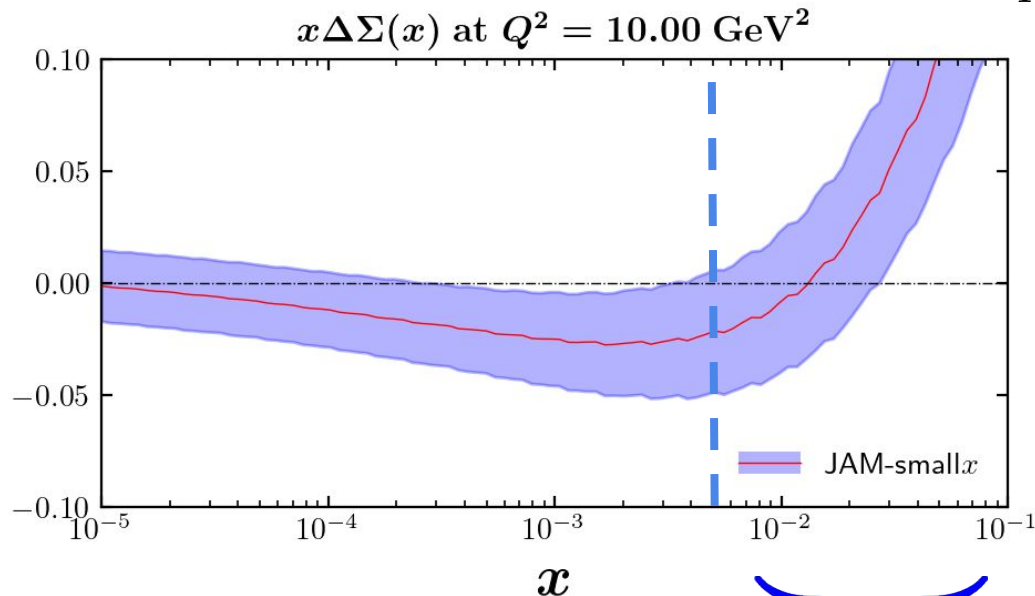


Large- $x$  region

$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$



# Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$

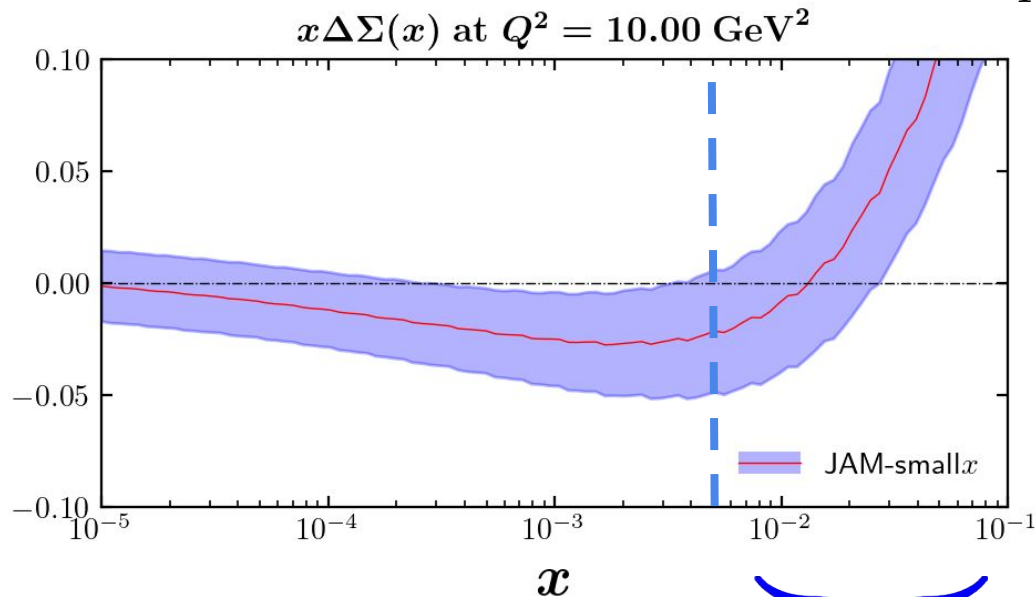


Large-x region

$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$

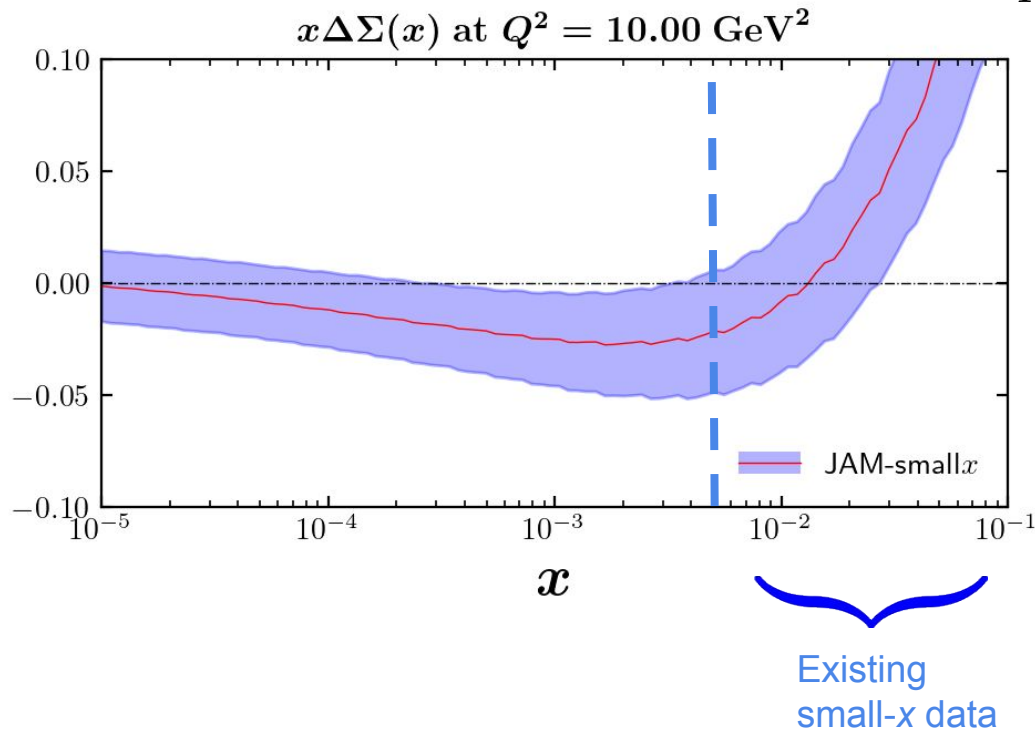
Existing  
small-x data

# Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$



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# Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$



$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$

Compare with:

$$\int_{10^{-5}}^{10^{-3}} dx \Delta\Sigma(x) = -0.1 \pm 0.1$$