

# Collider physics with no PDFs

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  - | More straightforward to compare to experimental data
- Physical basis previously discussed for example in Harland-Lang and Thorne 1811.08434, Hentschinski and Stratmann 1311.2825
- The novelty of our work:
  - | Momentum space
  - | Full three-flavor basis

# Straightforward example with only two observables

## Singlet and gluon approximation

$$(x, Q^2) = \sum_q^{n_f} [q(x, Q) + \bar{q}(x, Q^2)], \quad \text{where } n_f = 3$$

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## Singlet and gluon approximation

$$(x, Q^2) = \sum_q^{n_f} [q(x, Q) + \bar{q}(x, Q^2)], \quad \text{where } n_f = 3$$

Express  $F_2$  and  $F_L$  in terms of  $(x, Q^2)$  and gluon PDF:

$$F_2(x, Q^2) = x h e^2 i (x, Q^2), \quad \text{where } h e^2 i = \frac{1}{n_f} \sum_q^{n_f} e_q^2$$

$$\frac{F_L(x, Q^2)}{\frac{\alpha_s(Q^2)}{2\pi}} = x h e^2 i \left[ C_{F_L} + n_f C_{F_L g} g \right]$$

First non-zero order in  $\alpha_s \longrightarrow F_2, \frac{F_L}{\alpha_s} \neq 1$



## Straightforward example with only two observables

We need to invert  
this linear mapping

$$\longrightarrow \begin{bmatrix} F_2 \\ F_L / \frac{\alpha_s}{2\pi} \end{bmatrix} = \begin{bmatrix} xhe^2 i & 0 \\ xhe^2 i C_{FL} & xhe^2 i n_f C_{FLg} \end{bmatrix} \begin{bmatrix} \\ g \end{bmatrix}$$

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## Singlet and gluon PDF in physical basis

$$(x, Q^2) = \frac{1}{h e^2 i} \mathbf{1} \quad \tilde{F}_2$$

$$g(x, Q^2) = \int_x^1 \frac{dz}{z} \delta(1-z) \left\{ \frac{C_F}{4 T_R n_f h e^2 i} \begin{bmatrix} \frac{x}{z} \frac{d}{d \frac{x}{z}} & 2 \end{bmatrix} \frac{F_2 \left( \frac{x}{z}, Q^2 \right)}{\frac{x}{z}} \right. \\ \left. + \frac{2\pi}{\alpha_s(Q^2)} \frac{1}{8 T_R n_f h e^2 i} \begin{bmatrix} \frac{x^2}{z^2} \frac{d^2}{d \left( \frac{x}{z} \right)^2} & 2 \frac{x}{z} \frac{d}{d \frac{x}{z}} + 2 \end{bmatrix} \frac{F_L \left( \frac{x}{z}, Q^2 \right)}{\frac{x}{z}} \right\} \\ \frac{1}{n_f h e^2 i} \left\{ C_{g \tilde{F}_2^0} \tilde{F}_2^0 + C_{g \tilde{F}_2} \tilde{F}_2 + C_{g \tilde{F}_L^0} \tilde{F}_L^0 + C_{g \tilde{F}_L} \tilde{F}_L \right\}$$

Notation:

$$\tilde{F}_2(x, Q^2) = \frac{F_2(x, Q^2)}{x}, \quad \tilde{F}_L(x, Q^2) = \frac{2\pi}{\alpha_s(Q^2)} \frac{F_L(x, Q^2)}{x}, \\ \tilde{F}_{2,L}^0(x, Q^2) = x \frac{d}{dx} \tilde{F}_{2,L}(x, Q^2), \quad \tilde{F}_L^0(x, Q^2) = x^2 \frac{d^2}{dx^2} \tilde{F}_L(x, Q^2)$$

# Results for two observable basis

DGLAP equations of  $F_2$  and  $F_L/\alpha_s$  in physical basis:

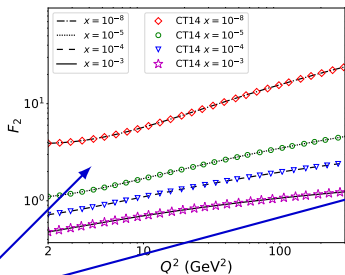
$$\frac{dF_2(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} x \left\{ 2P_{qq} \begin{bmatrix} C_{g\tilde{F}^{00}_L} & \tilde{F}^{00}_L + C_{g\tilde{F}^0_L} & \tilde{F}^0_L + C_{g\tilde{F}_L} & \tilde{F}_L \end{bmatrix} \right. \\ \left. + P_{qq} \begin{bmatrix} 1 & \tilde{F}_2 + 2P_{qq} & C_{g\tilde{F}^0_2} & \tilde{F}^0_2 + C_{g\tilde{F}_2} & \tilde{F}_2 \end{bmatrix} \right\}$$

$$\frac{d}{d \log(Q^2)} \left( \frac{F_L(x, Q^2)}{\frac{\alpha_s(Q^2)}{2\pi}} \right) = \\ \frac{\alpha_s(Q^2)}{2\pi} x \left\{ \begin{bmatrix} 2C_{F_L} & P_{qq} + C_{F_L g} & P_{gg} \end{bmatrix} \begin{bmatrix} C_{g\tilde{F}^{00}_L} & \tilde{F}^{00}_L + C_{g\tilde{F}^0_L} & \tilde{F}^0_L + C_{g\tilde{F}_L} & \tilde{F}_L \end{bmatrix} \right. \\ \left. + \begin{bmatrix} C_{F_L} & (P_{qq} \begin{bmatrix} 1 + 2P_{qq} & C_{g\tilde{F}_2} \end{bmatrix} + C_{F_L g} \begin{bmatrix} n_f P_{qq} & 1 + P_{gg} & C_{g\tilde{F}_2} \end{bmatrix}) \end{bmatrix} \tilde{F}_2 \right. \\ \left. + \begin{bmatrix} 2C_{F_L} & P_{qq} + C_{F_L g} & P_{gg} \end{bmatrix} C_{g\tilde{F}^0_2} \tilde{F}^0_2 \right\}$$

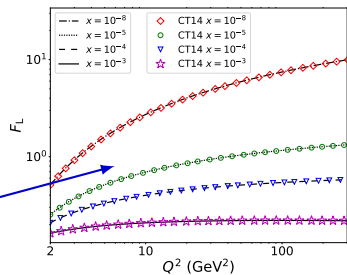
Scheme dependence of PDFs only at next order in  $\alpha_s$

! Should agree with evolution in terms of PDFs

# Comparison with conventional DGLAP evolution

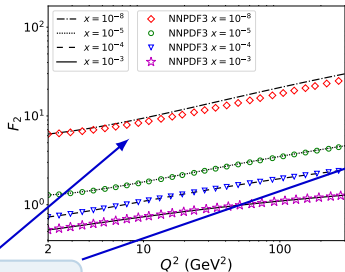


$F_2$  CT14 LO

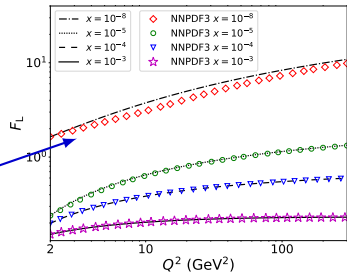


$F_L$  CT14 LO

Nice match



$F_2$  NNP3.0 LO



$F_L$  NNP3.0 LO

Small discrepancy  
after  $x = 10^{-6}$

# Six observable basis

- Full three-flavor basis:  $u, u, d, d, s = s$ , and  $g$   
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- Full three-flavor basis:  $u, u, d, d, s = s$ , and  $g$   
! Need six linearly independent DIS structure functions
- We choose structure functions: (again in first non-zero order in  $\alpha_s$ )

## Neutral current

$$F_2 = x \sum_q^{n_f} e_q^2 (q + \bar{q})$$

$$F_3 = 2 \sum_q^{n_f} (L_q^2 - R_q^2) (q - \bar{q}),$$

$$\text{where } L_q = T_q^3 - 2e_q \sin^2 \theta_W$$

$$\text{and } R_q = 2e_q \sin^2 \theta_W$$

$$\frac{F_L(x, Q^2)}{\frac{\alpha_s(Q^2)}{2\pi}} = x \left[ C_{F_L \bar{F}_2} \sum_q^{n_f} e_q^2 (q + \bar{q}) + \sum_q^{n_f} e_q^2 C_{F_L g} [g] \right]$$

## Charged current

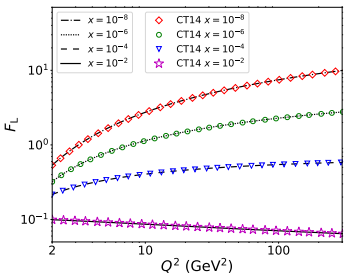
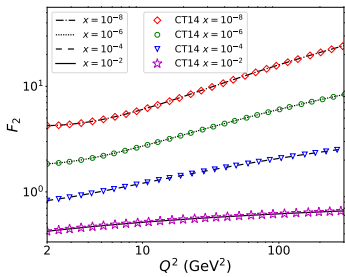
$$F_2^W = 2x(u + \bar{d} + \bar{s})$$

$$F_3^W = 2(u - \bar{d} - \bar{s})$$

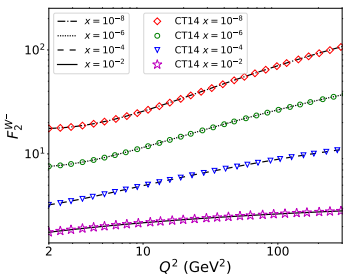
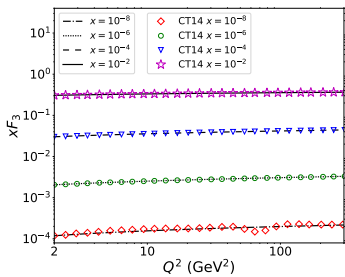
$$F_{2C}^W = 2x\bar{s}$$

Derive DGLAP equations for these  
Again, should agree with conventional DGLAP evolution

# Comparison with conventional DGLAP evolution



$F_2$  CT14 LO Match in full grid  $F_L$  CT14 LO



$xF_3$  CT14 LO

$F_2^W$  CT14 LO

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- We have implemented these equations numerically – agree with DGLAP-evolved PDFs
- Scheme dependence of PDFs starts to play part at NLO in  $\alpha_s$ 
  - ! By using a physical basis in second non-zero order in  $\alpha_s$ , we will be able to avoid scheme dependence
- What next:
  - | Establish physical basis at NLO
  - | Study how LHC cross sections, e.g. Drell-Yan, are expressed in physical basis
  - | Obtain physical basis including also heavy quarks

# Backup: Inverting the gluon PDF

Gluon PDF in mellin space

$$g(n) = \frac{1}{n_f C_{F_L g}(n)} \left[ \frac{1}{h e^2 i} \tilde{F}_L(n) \quad C_{F_L}(n) \quad (n) \right]$$

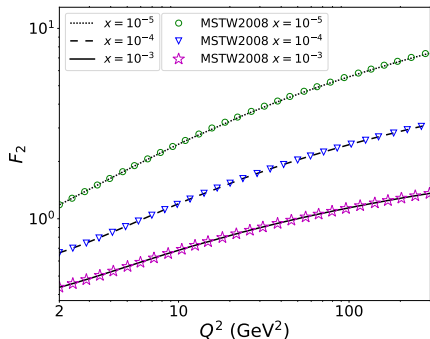
$$\frac{1}{C_{F_L g}(n)} = \frac{1}{8 T_R Z_0^n} \int_0^1 dz z^{n+2} \delta''(z - z_0),$$

where  $z_0 \in ]0, 1[$  is an arbitrary constant that cancels in final result.

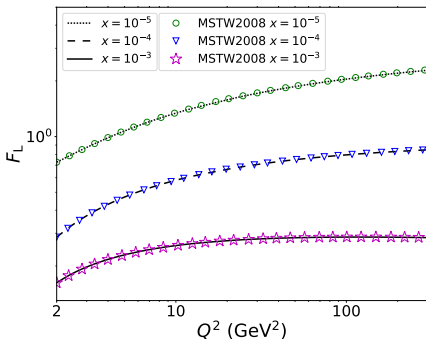
$$g(x, Q^2) = \int_x^1 \frac{dz}{z} \delta(1 - z) \left\{ \frac{C_F}{4 T_R n_f h e^2 i} \left[ \frac{x}{z} \frac{d}{d \frac{x}{z}} - 2 \right] \frac{F_2\left(\frac{x}{z}, Q^2\right)}{\frac{x}{z}} \right. \\ \left. + \frac{2\pi}{\alpha_s(Q^2)} \frac{1}{8 T_R n_f h e^2 i} \left[ \frac{x^2}{z^2} \frac{d^2}{d \left(\frac{x}{z}\right)^2} - 2 \frac{x}{z} \frac{d}{d \frac{x}{z}} + 2 \right] \frac{F_L\left(\frac{x}{z}, Q^2\right)}{\frac{x}{z}} \right\} \\ \frac{1}{n_f h e^2 i} \left\{ C_{g\tilde{F}^0_2} \tilde{F}^0_2 + C_{g\tilde{F}_2} \tilde{F}_2 + C_{g\tilde{F}^0_0_L} \tilde{F}^0_0_L + C_{g\tilde{F}^0_L} \tilde{F}^0_L + C_{g\tilde{F}_L} \tilde{F}_L \right\}$$

# Backup: Comparison with MSTW2008 PDF set

DGLAP evolution for  $F_2$  and  $F_L$  in two observable physical basis



$F_2$  MSTW LO



$F_L$  MSTW LO

## Backup: PDFs in six observable physical basis

$$\begin{bmatrix} F_2 \\ F_3 \\ F_2^W \\ F_3^W \\ F_{2c}^W \end{bmatrix} = \begin{bmatrix} xe_d^2 & xe_d^2 & xe_u^2 & xe_u^2 & 2xe_s^2 \\ 2A_d & 2A_d & 2A_u & 2A_u & 0 \\ 0 & 2x & 2x & 0 & 2x \\ 0 & 2 & 2 & 0 & 2 \\ 0 & 0 & 0 & 0 & 2x \end{bmatrix} \begin{bmatrix} d \\ \bar{d} \\ u \\ \bar{u} \\ \bar{s} \end{bmatrix},$$

where we have defined  $A_q$   $L_q^2$   $R_q^2$  in order to simplify equations.

$$xd(x, Q^2) = \frac{1}{A_u e_d^2 + A_d e_u^2} \left[ A_u F_2(x, Q^2) + \frac{e_u^2}{2} x F_3(x, Q^2) \frac{A_u(2e_u^2 + e_d^2)}{4} \frac{A_d e_u^2}{F_2^W} (x, Q^2) \right. \\ \left. \frac{A_u(2e_u^2 - e_d^2) + A_d e_u^2}{4} x F_3^W (x, Q^2) + \frac{A_u(e_d^2 - 2e_s^2)}{2} \frac{A_d e_u^2}{F_{2c}^W} (x, Q^2) \right]$$

$$x\bar{d}(x, Q^2) = \frac{1}{4} F_2^W (x, Q^2) - \frac{1}{4} x F_3^W (x, Q^2) - \frac{1}{2} F_{2c}^W$$

$$xu(x, Q^2) = \frac{1}{4} F_2^W (x, Q^2) + \frac{1}{4} x F_3^W (x, Q^2)$$

$$x\bar{u}(x, Q^2) = \frac{1}{A_u e_d^2 + A_d e_u^2} \left[ A_d F_2(x, Q^2) - \frac{e_d^2}{2} x F_3(x, Q^2) \frac{A_d(2e_d^2 + e_u^2)}{4} \frac{A_u e_d^2}{F_2^W} (x, Q^2) \right. \\ \left. + \frac{A_d(2e_d^2 - e_u^2) + A_u e_d^2}{4} x F_3^W (x, Q^2) + A_d(e_d^2 - e_s^2) F_{2c}^W (x, Q^2) \right]$$

$$x\bar{s}(x, Q^2) = xs(x, Q^2) = \frac{1}{2} F_{2c}^W (x, Q^2)$$

# Backup: Results in six observable basis

$$\frac{dF_2(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} \left[ P_{qq} \quad \tilde{F}_2 + 2 \sum_q^{n_f} e_q^2 P_{qg} \quad g \right], \quad \text{where } \tilde{F}_2(x, Q^2) = \frac{F_2(x, Q^2)}{x}$$

$$\frac{dx F_3(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} x P_{qq} \quad F_3$$

$$\frac{dF_2^{\text{W}}(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} x \left[ P_{qq} \quad \tilde{F}_2^{\text{W}} + 6 P_{qg} \quad g \right], \quad \text{where } \tilde{F}_2^{\text{W}}(x, Q^2) = \frac{F_2^{\text{W}}(x, Q^2)}{x}$$

$$\frac{dx F_3^{\text{W}}(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} x \left[ P_{qq} \quad F_3^{\text{W}} \quad 2 P_{qg} \quad g \right]$$

$$\frac{dF_{2c}^{\text{W}}(x, Q^2)}{d \log(Q^2)} = \frac{\alpha_s(Q^2)}{2\pi} x \left[ P_{qq} \quad \tilde{F}_{2c}^{\text{W}} + 2 P_{qg} \quad g \right], \quad \text{where } \tilde{F}_{2c}^{\text{W}}(x, Q^2) = \frac{F_{2c}^{\text{W}}(x, Q^2)}{x}$$

$$P_{qg} \quad g = \frac{C_F}{4 \sum_q^{n_f} e_q^2} \left[ \tilde{F}_2 + 2(z-1) \tilde{F}_2 \right] + \frac{1}{4 \sum_q^{n_f} e_q^2} \left[ \left( \frac{1}{x} \quad \frac{1}{2} \frac{d}{dx} \right) F_L(x, Q^2) + 1 \quad \tilde{F}_L \right]$$



## Backup: Results in six observable basis

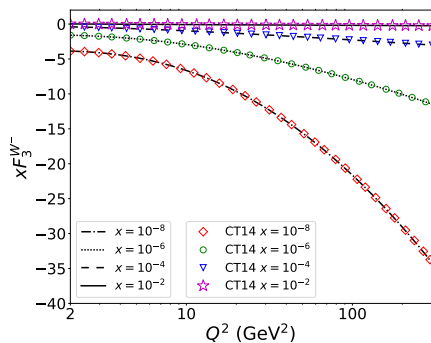
$$\frac{d}{d \log(Q^2)} \left( \frac{F_L(x, Q^2)}{\frac{\alpha_s(Q^2)}{2\pi}} \right) =$$

$$\frac{\alpha_s(Q^2)}{2\pi} x \left\{ \begin{aligned} & \left[ 2C_{F_L \tilde{F}_2} \quad P_{qg} + C_{F_L g} \quad P_{gg} \right] \left[ C_{g \tilde{F}^{00}_L} \quad \tilde{F}^{00}_L + C_{g \tilde{F}^0_L} \quad \tilde{F}^0_L + C_{g \tilde{F}_L} \quad \tilde{F}_L \right] \\ & + \left[ C_{F_L \tilde{F}_2} \quad \left( P_{qq} + 2P_{qg} \quad C_{g \tilde{F}_2} \right) + C_{F_L g} \quad P_{gg} \quad C_{g \tilde{F}_2} \right] \quad \tilde{F}_2 \\ & + \left[ 2C_{F_L \tilde{F}_2} \quad P_{qg} + C_{F_L g} \quad P_{gg} \right] \quad C_{g \tilde{F}^0_2} \quad \tilde{F}^0_2 + \sum_q^{n_f} e_q^2 C_{F_L g} \quad P_{qg} \quad \sum_q (q + \bar{q}) \end{aligned} \right\}$$

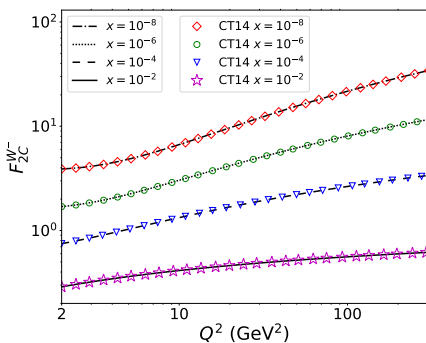
where  $\sum_q (q + \bar{q})$  can be expressed in physical basis (see slide 11/9)

# Backup: Comparison with conventional DGLAP evolution

Comparison for structure functions  $F_3^W$  and  $F_{2C}^W$  in six observable physical basis



$F_3^W$  CT14 LO



$F_{2C}^W$  CT14 LO