

Prospects of inclusive reactions using ~~quantum computers~~ *real time calculations*



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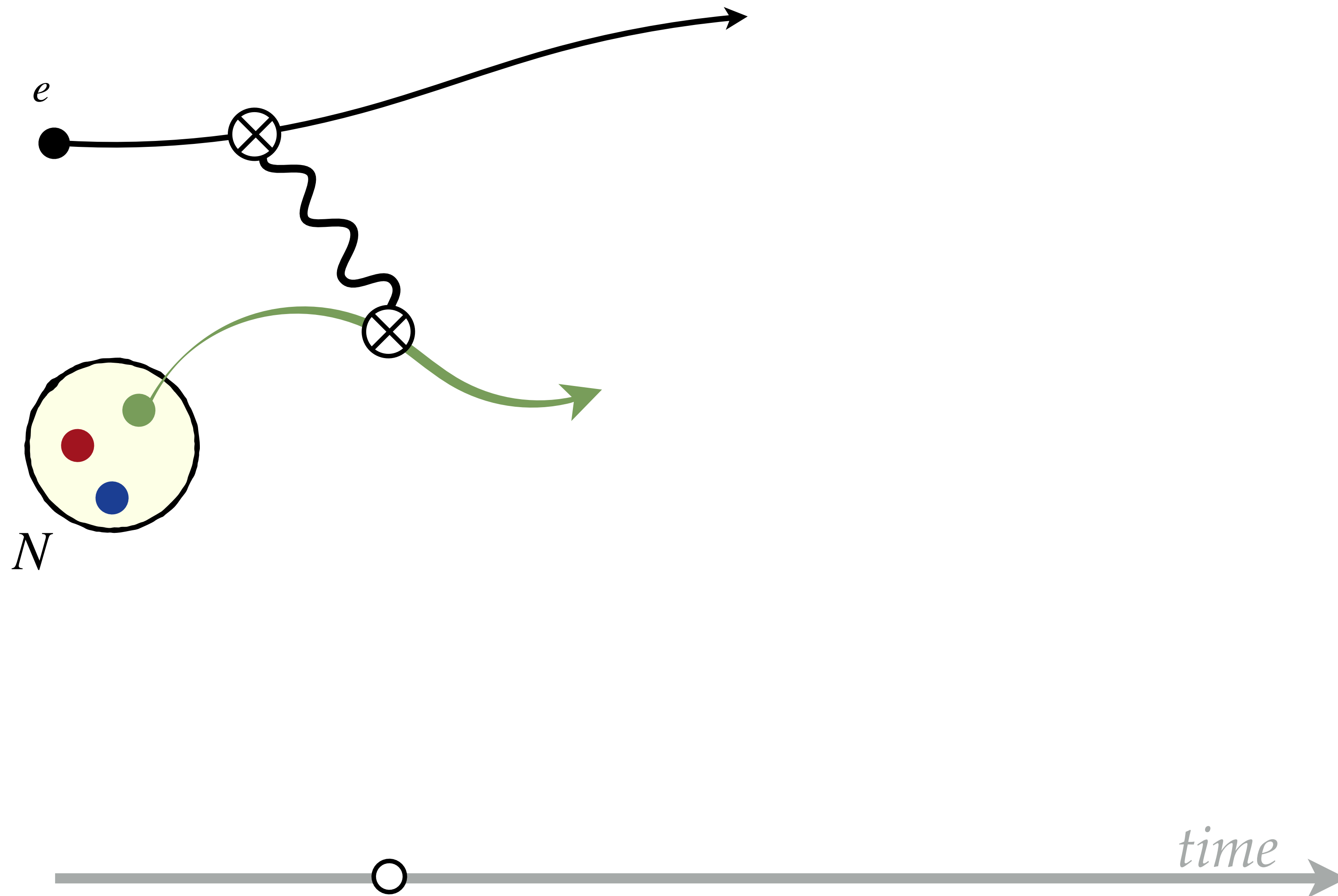
🐦 @RaulBriceno12

Carrillo, Guerrero, Sturzu, RB (to appear)

RB, Guerrero, Hansen, Sturzu (2020)

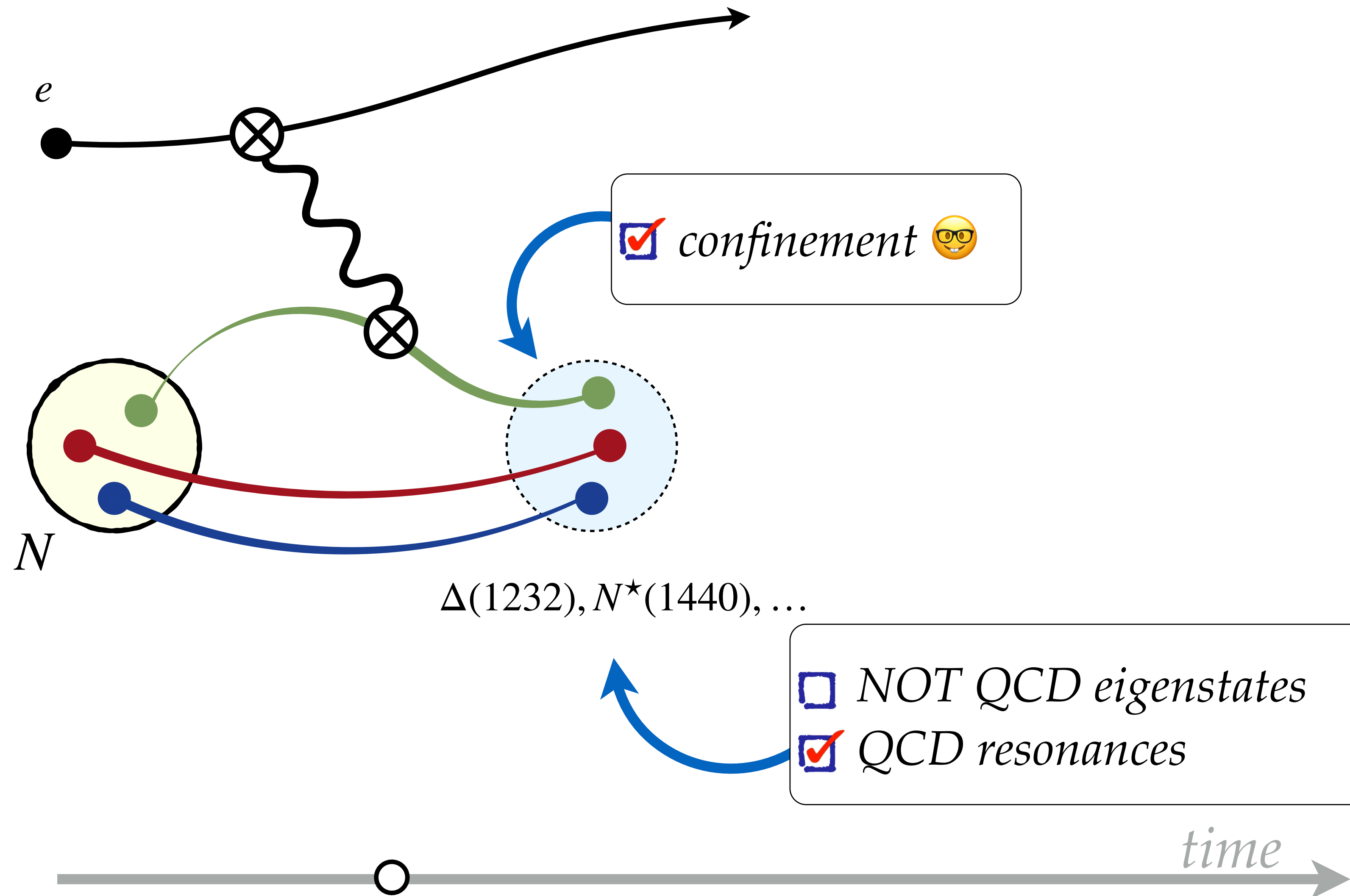
Inclusive reactions *just a subclass*

- Virtual Compton scattering: PDFs, GPDs,...



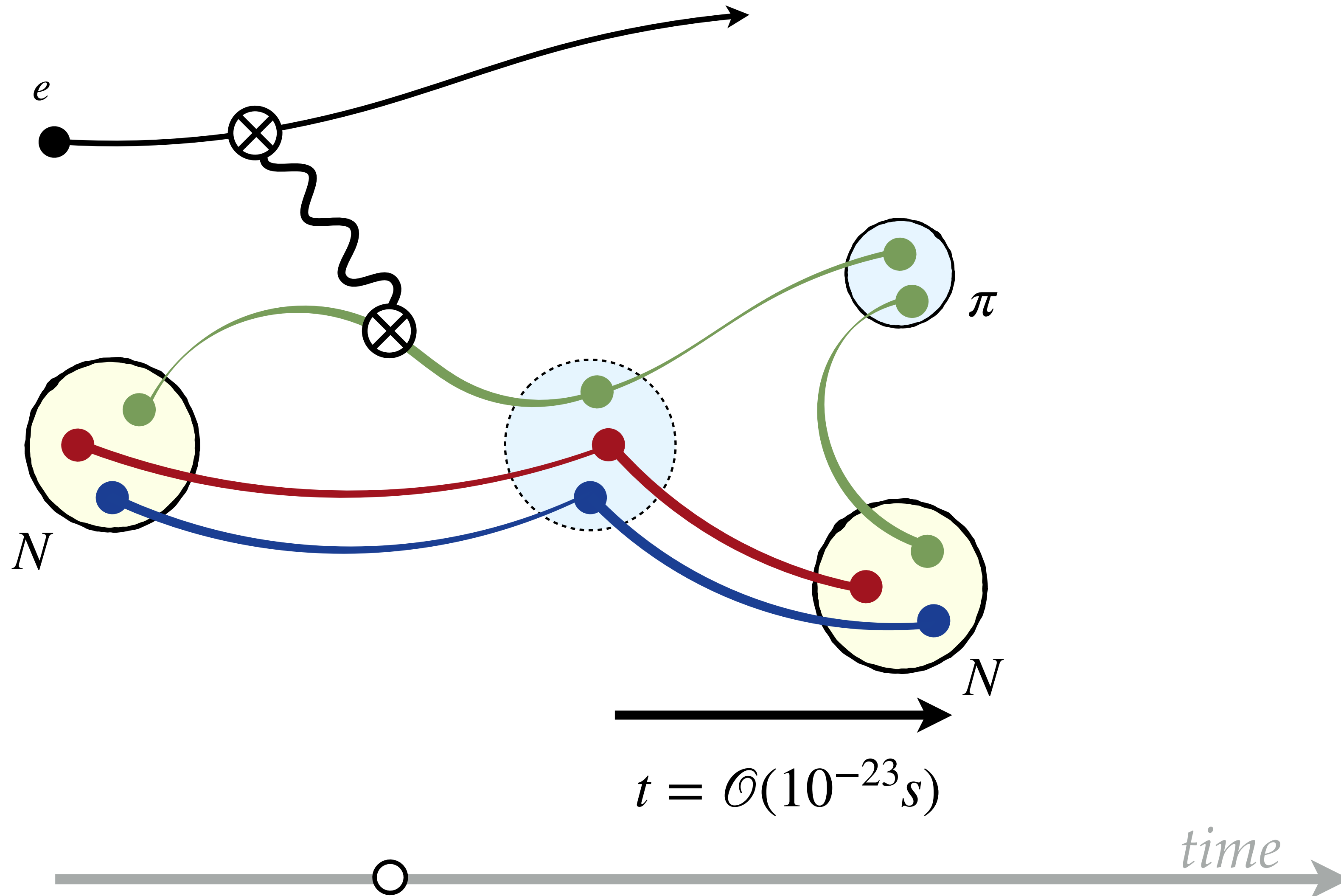
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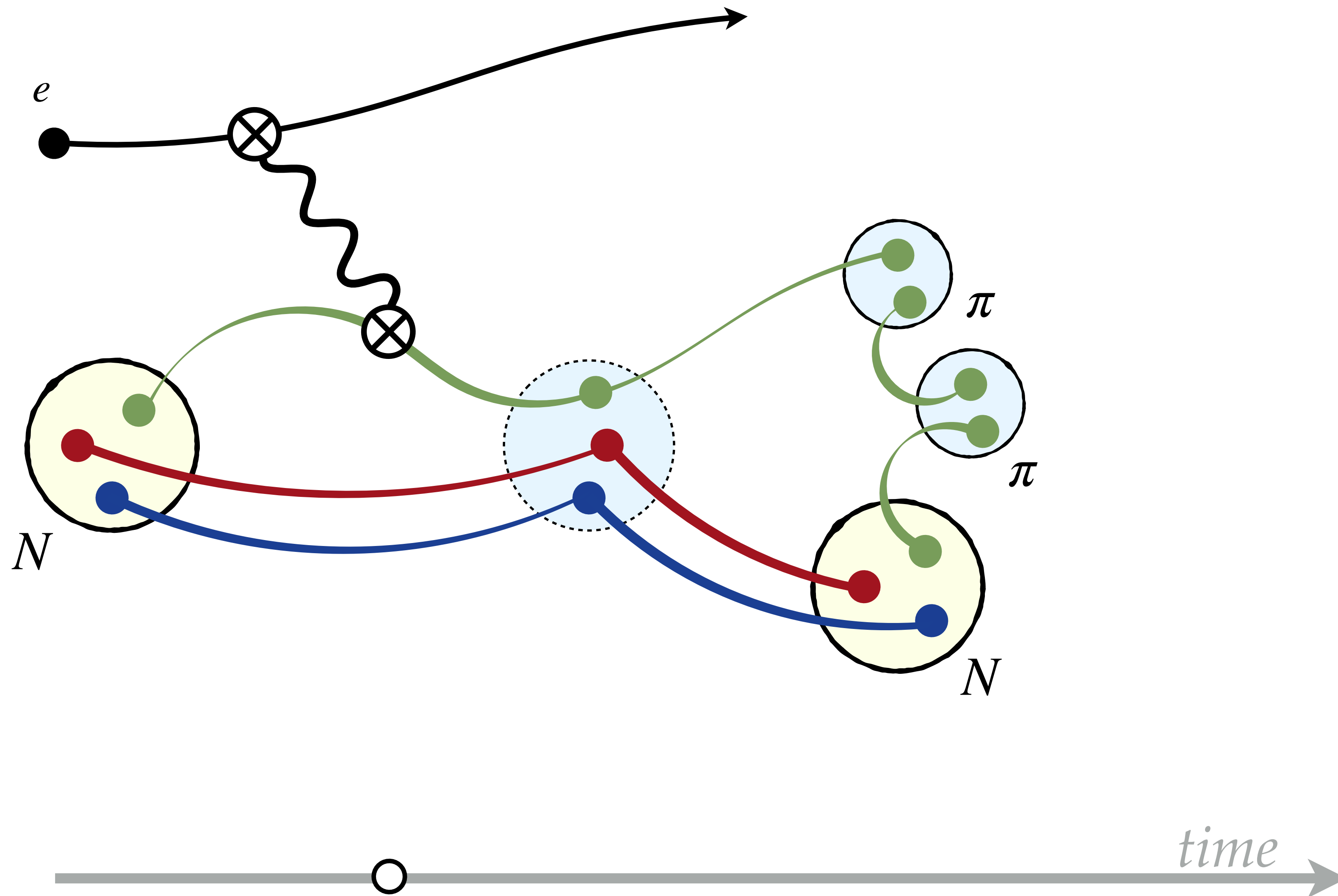
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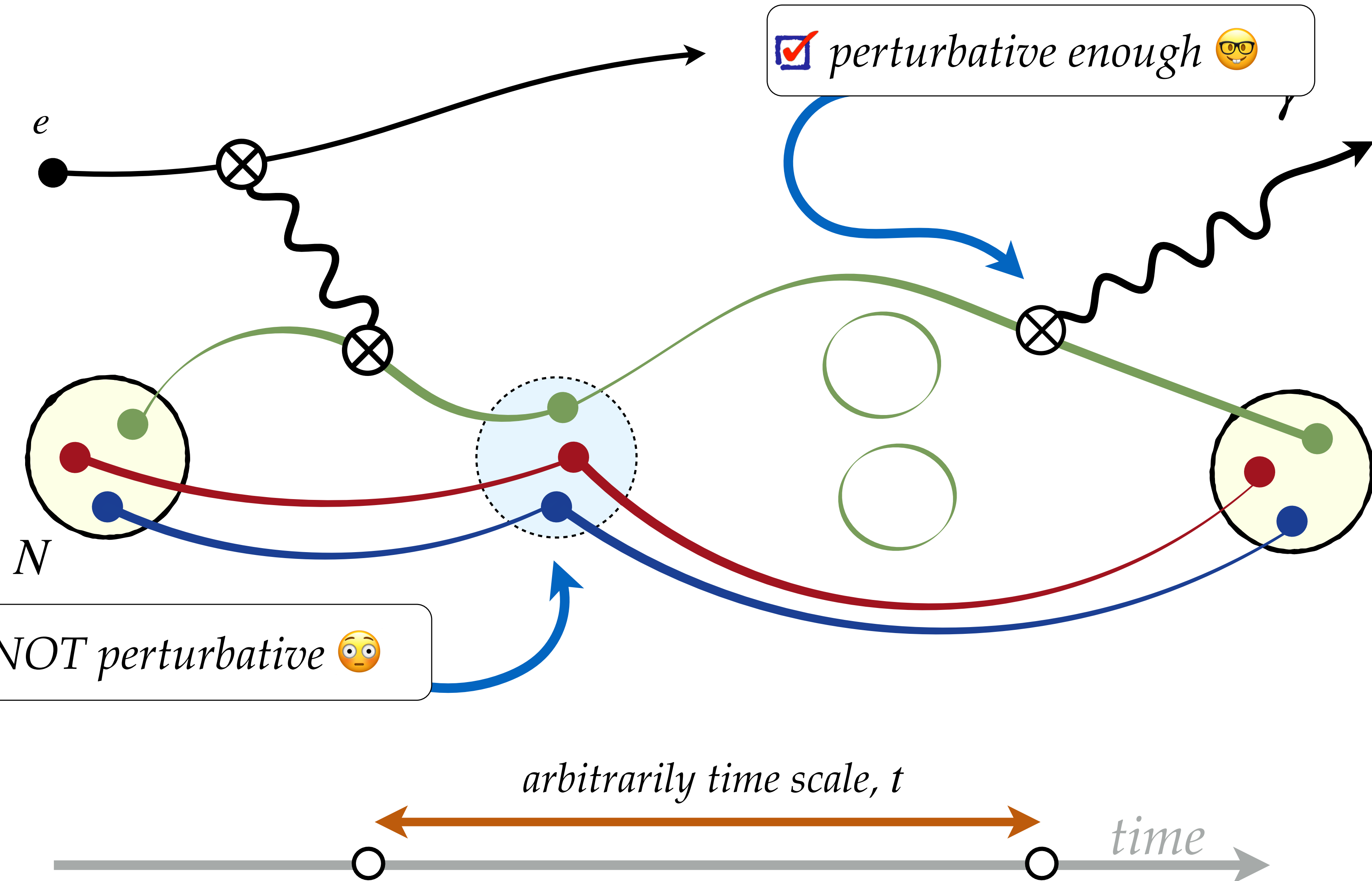
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Inclusive reactions *just a subclass*

☐ Virtual Compton scattering: PDFs, GPDs,...



Inclusive reactions *just a subclass*

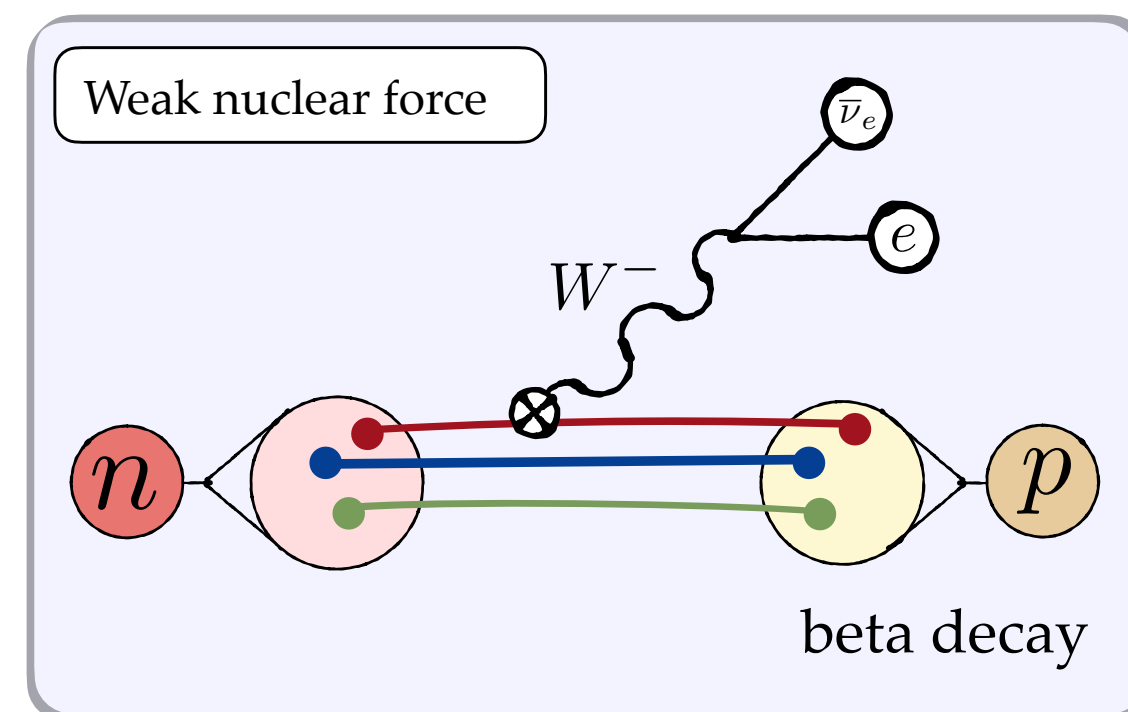
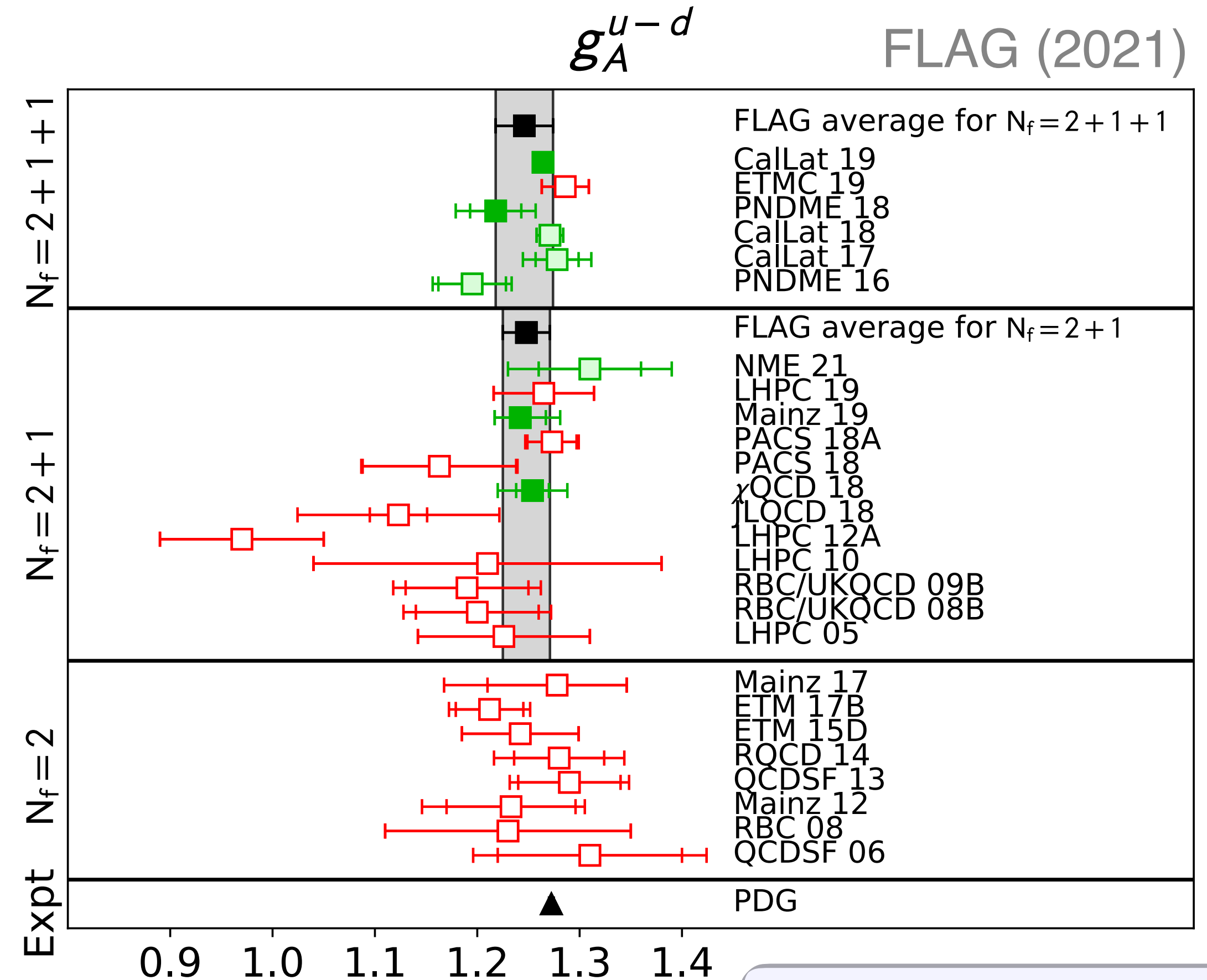
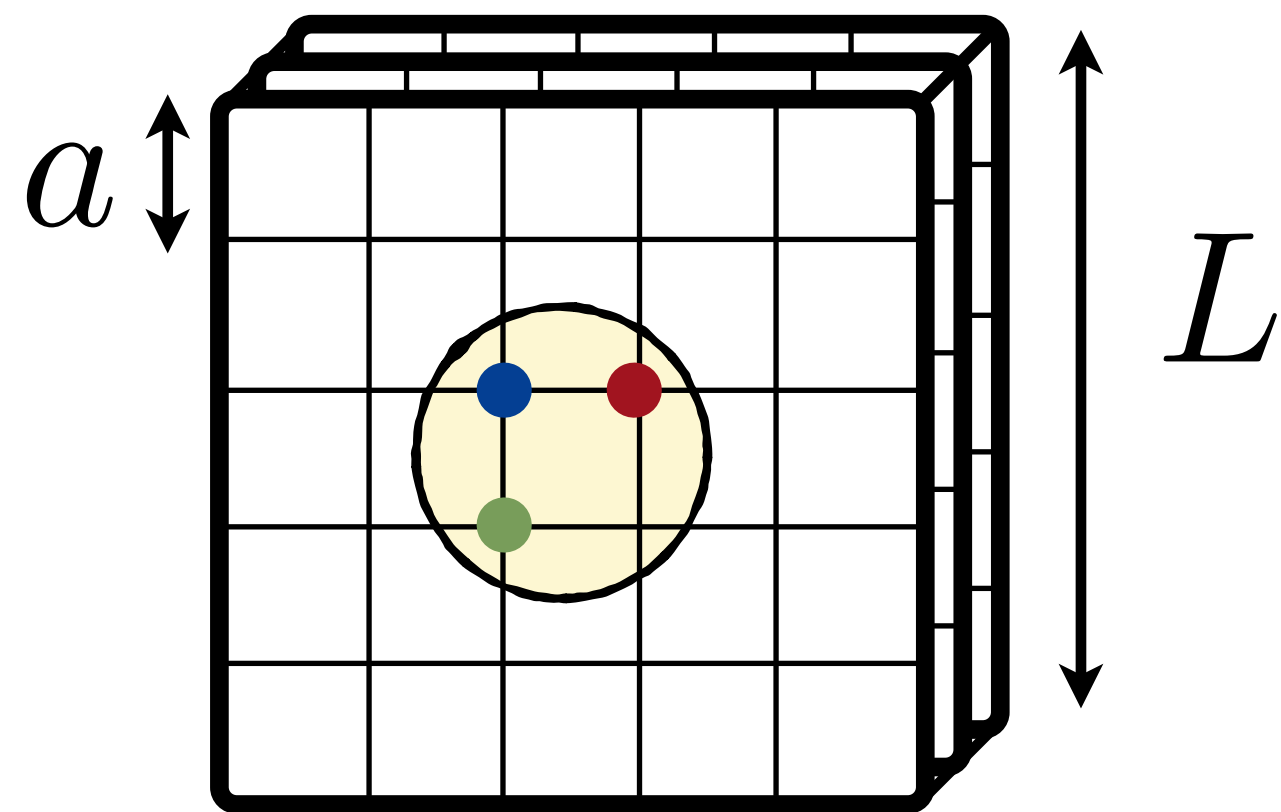
- ❑ Virtual Compton scattering: PDFs, GPDs,...
- ❑ inclusive neutrino-nucleus scattering,
- ❑ double beta decay
- ❑ Glueball structure,
- ❑ Radiative corrections in weak decays
- ❑ ...

All can be defined as: $\mathcal{T} \sim \int d^4x e^{ix \cdot q} \langle n_f | T [\mathcal{J}_{2,M}(t) \mathcal{J}_1(0)] | n_i \rangle_\infty$



lattice QCD *only current non-perturbative QCD tool*

- Euclidean spacetime: $t_M \rightarrow -it_E$
- Monte Carlo sampling
- finite volume: $L \sim 1 - 10$ fm
- lattice spacing: $a \sim 0.03 - 0.1$ fm
- quark masses: $m_q \rightarrow m_q^{\text{phys}}$



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❑ Euclidean spacetime: $t_M \rightarrow -it_E$

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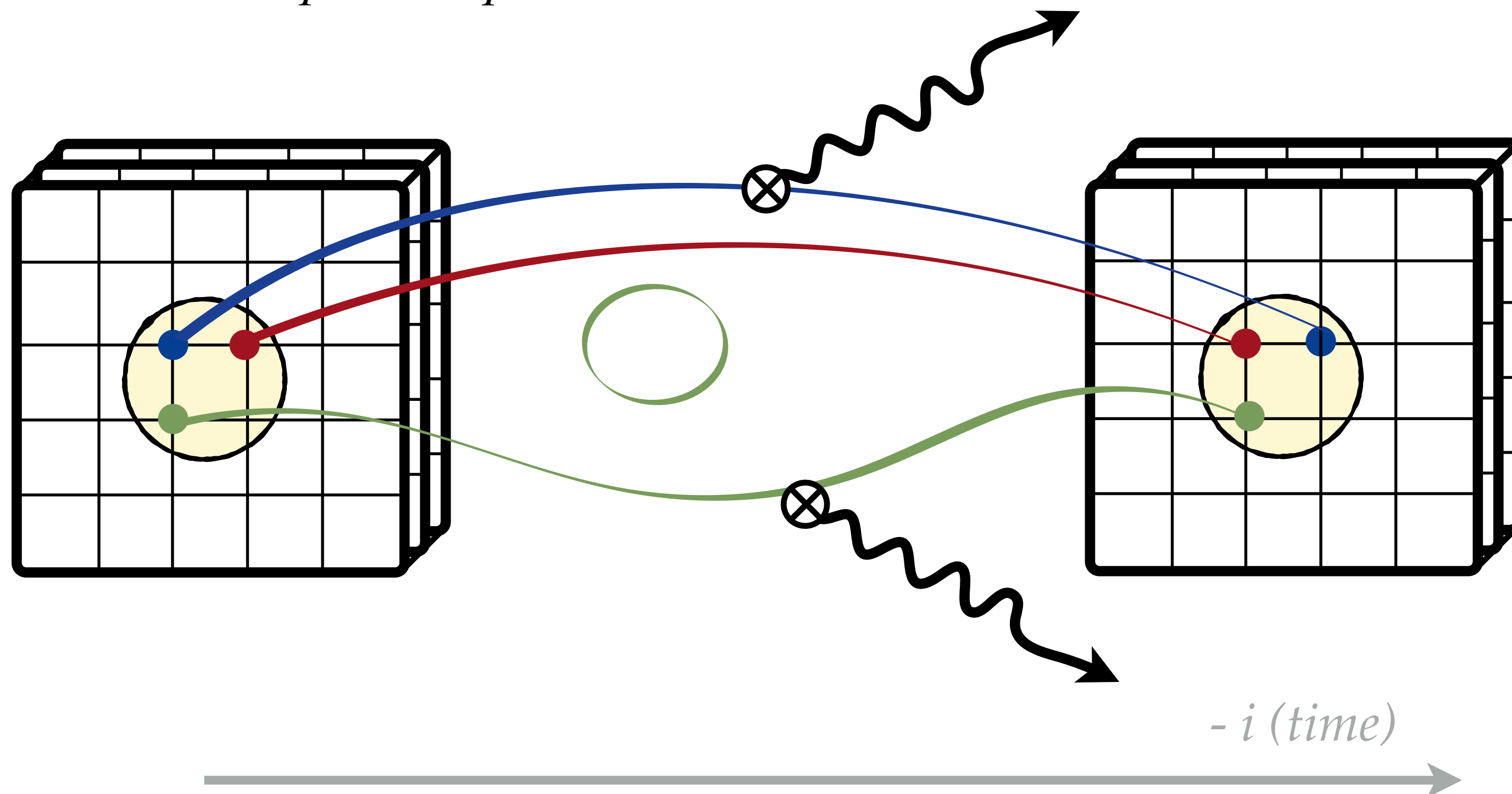
❑ finite volume: $L \sim 1 - 10$ fm

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strongly correlated issues:

*“time evolution operator $\sim e^{-t\hat{H}_L}$
depends on both the time-signature and size
of the volume”*



~~lattice QCD~~ *quantum computers, tensor networks, ...*

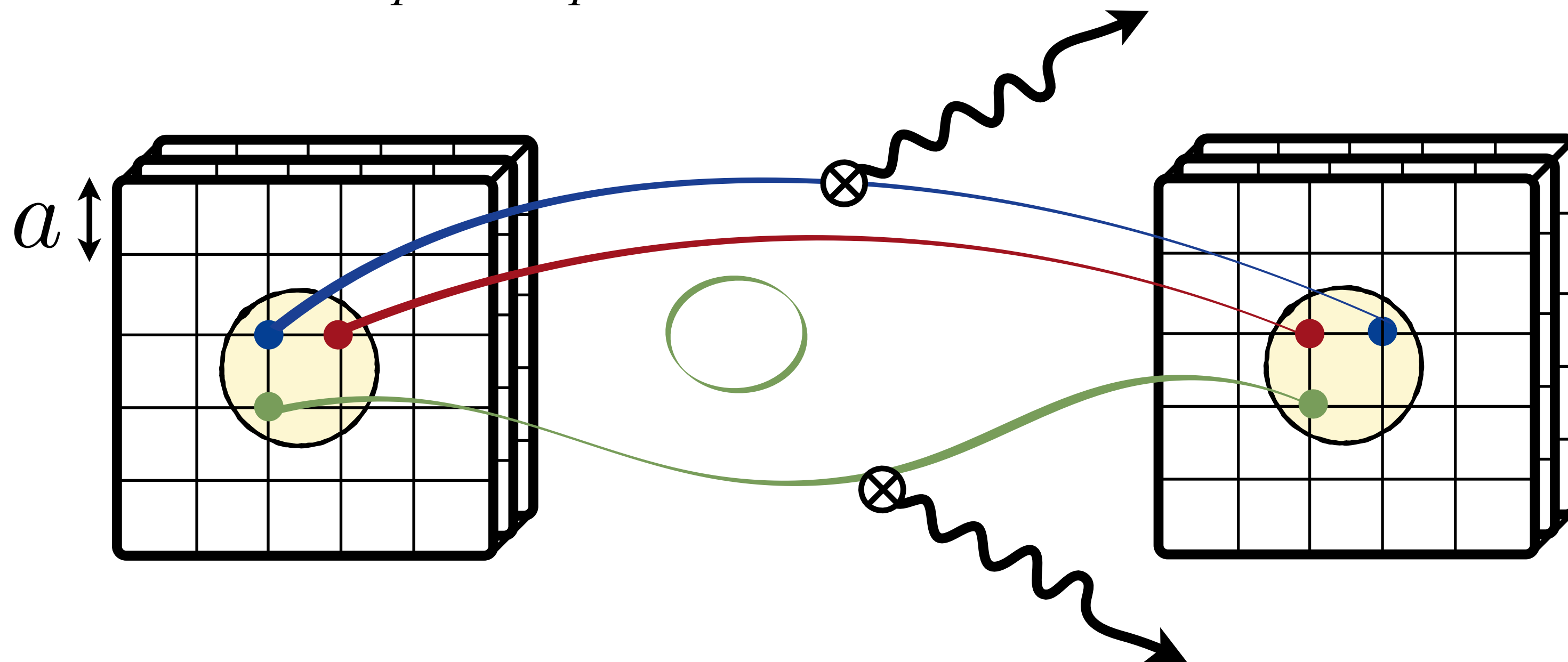
☐ ~~Euclidean spacetime: $t_M \rightarrow -it_E$~~

☒ Monte Carlo sampling

☐ finite volume: $L \sim 1 - 10$ fm

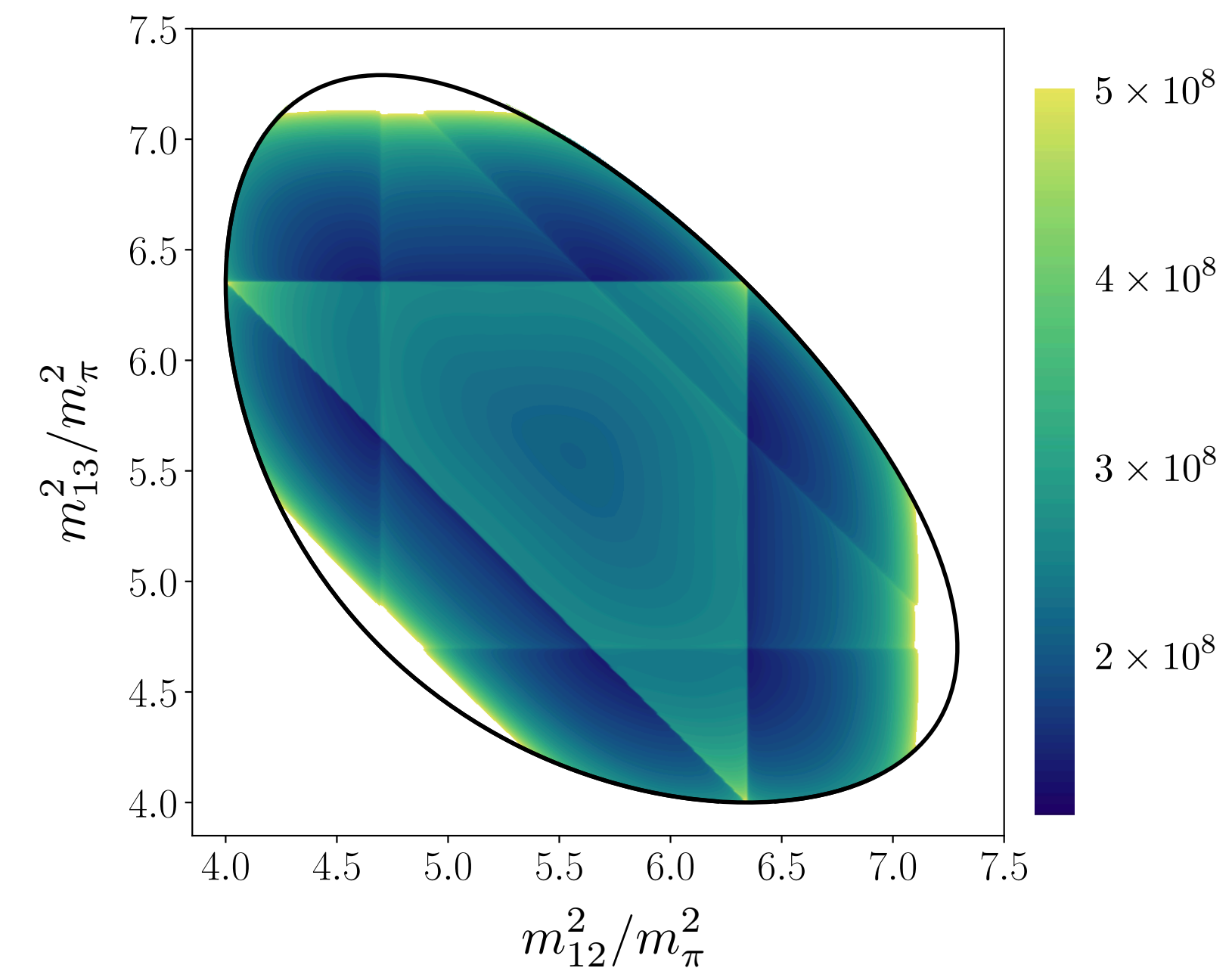
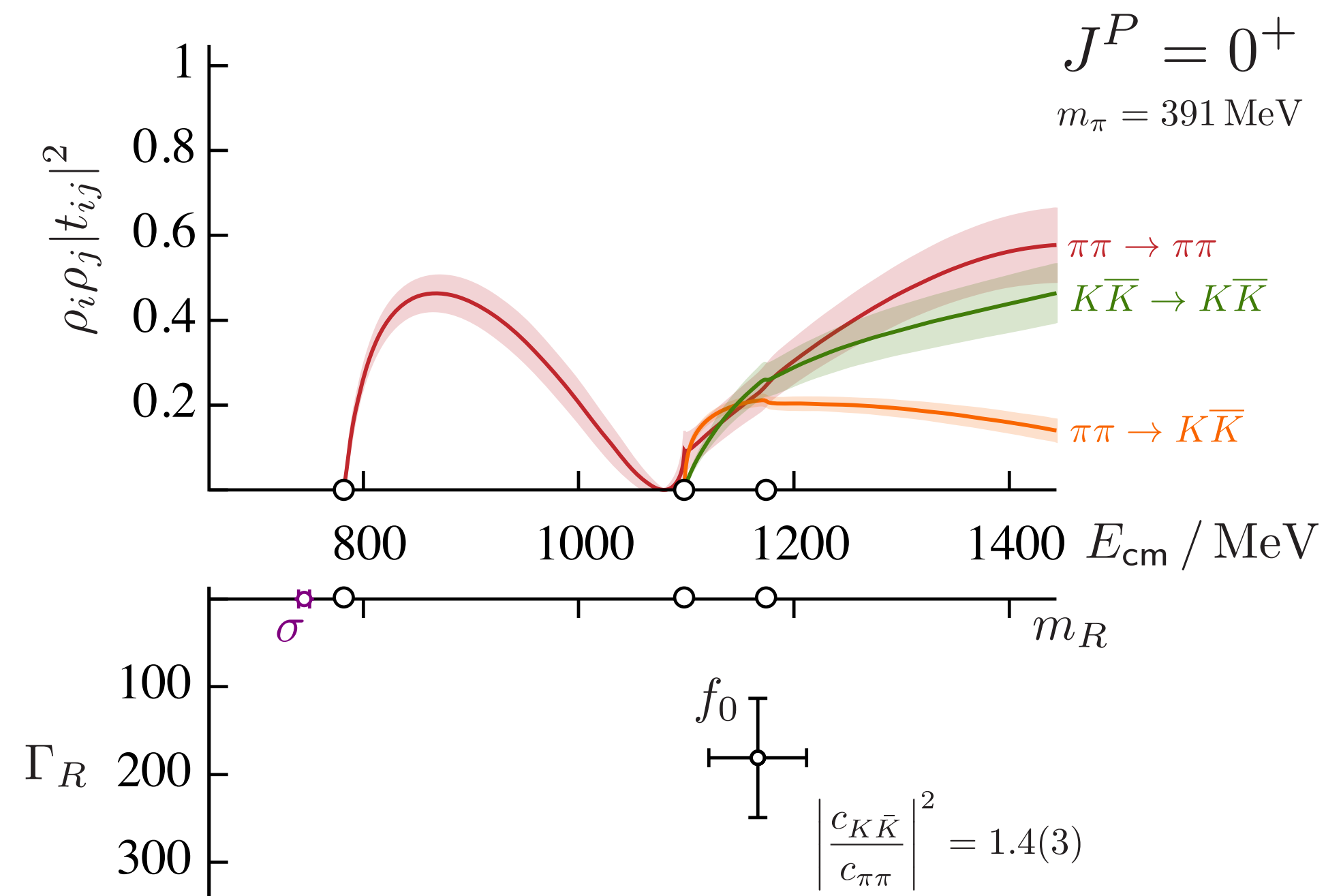
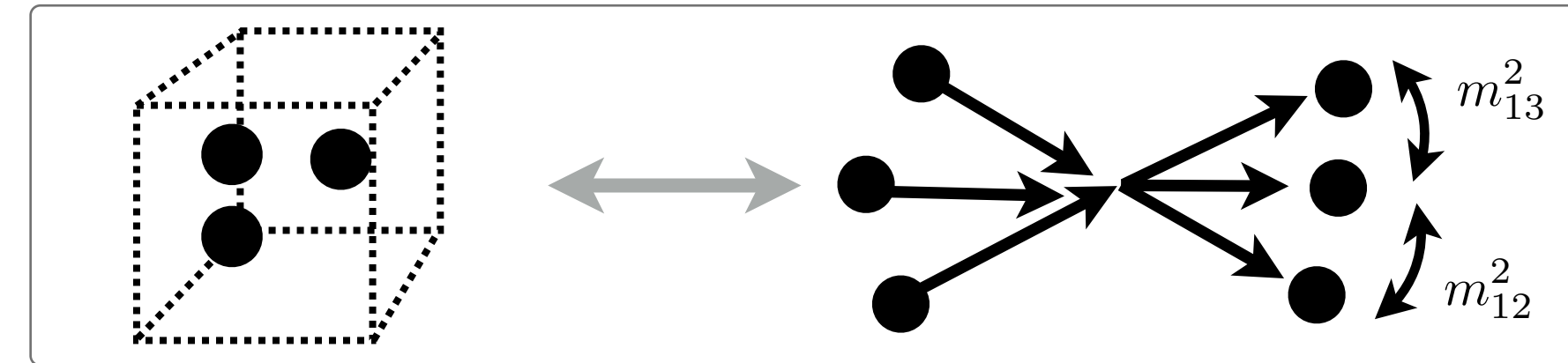
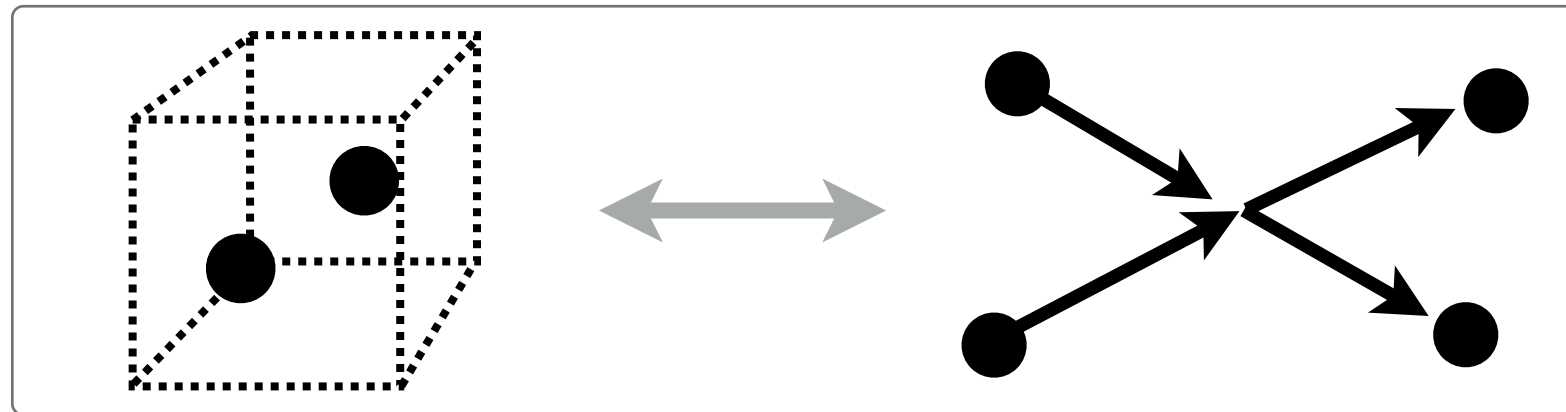
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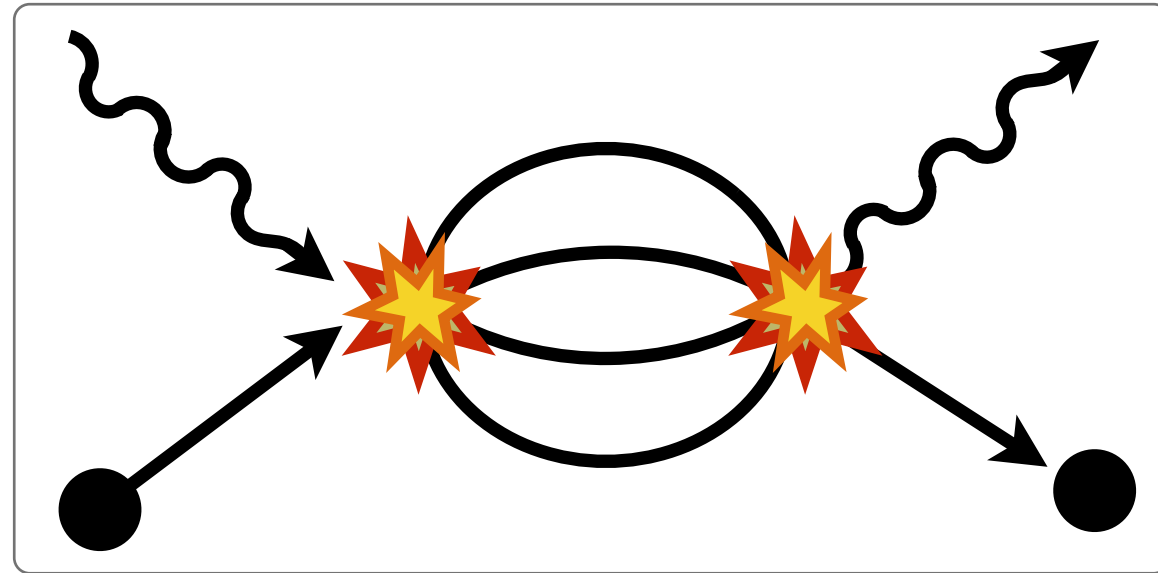
Exclusive vs. inclusive reactions

- ❑ If exclusive and *interesting*
- ❑ After developing increasingly complex formalism...
- ❑ Lattice QCD **will always win**



Exclusive vs. inclusive reactions

- ❑ If exclusive and *interesting*
 - ❑ After developing increasingly complex formalism...
 - ❑ Lattice QCD **will always win**
- ❑ Inclusive reactions, QC methods *may* be needed and worth investigating.



Infinite-volume reactions

- ❑ complex functions,
- ❑ kinematic singularities,
- ❑ due to intermediate on-shell states.

Hamiltonian frameworks

Four-point functions in a finite, Minkowski spacetime

$$\mathcal{T} \sim \int_0^T d^4x e^{it(\omega+i\epsilon)} \langle n_f | \mathcal{J}(t) \mathcal{J}(0) | n_i \rangle_\infty$$

[only considering one time ordering,
introduced ϵ as a regulator]

$$= \sum_n \int_0^T d^4x e^{it(E_f+\omega-E_n+i\epsilon)} \langle n_f | \mathcal{J}(0) | n \rangle \langle n | \mathcal{J}(0) | n_i \rangle_\infty$$

[inserting a complete set of discrete
finite-volume states]

$$\approx \sum_n i \frac{\langle n_f | \mathcal{J}(0) | n \rangle \langle n | \mathcal{J}(0) | n_i \rangle_\infty}{(E_f + \omega - E_n + i\epsilon)}$$

[assuming $\epsilon T \gg 1$]

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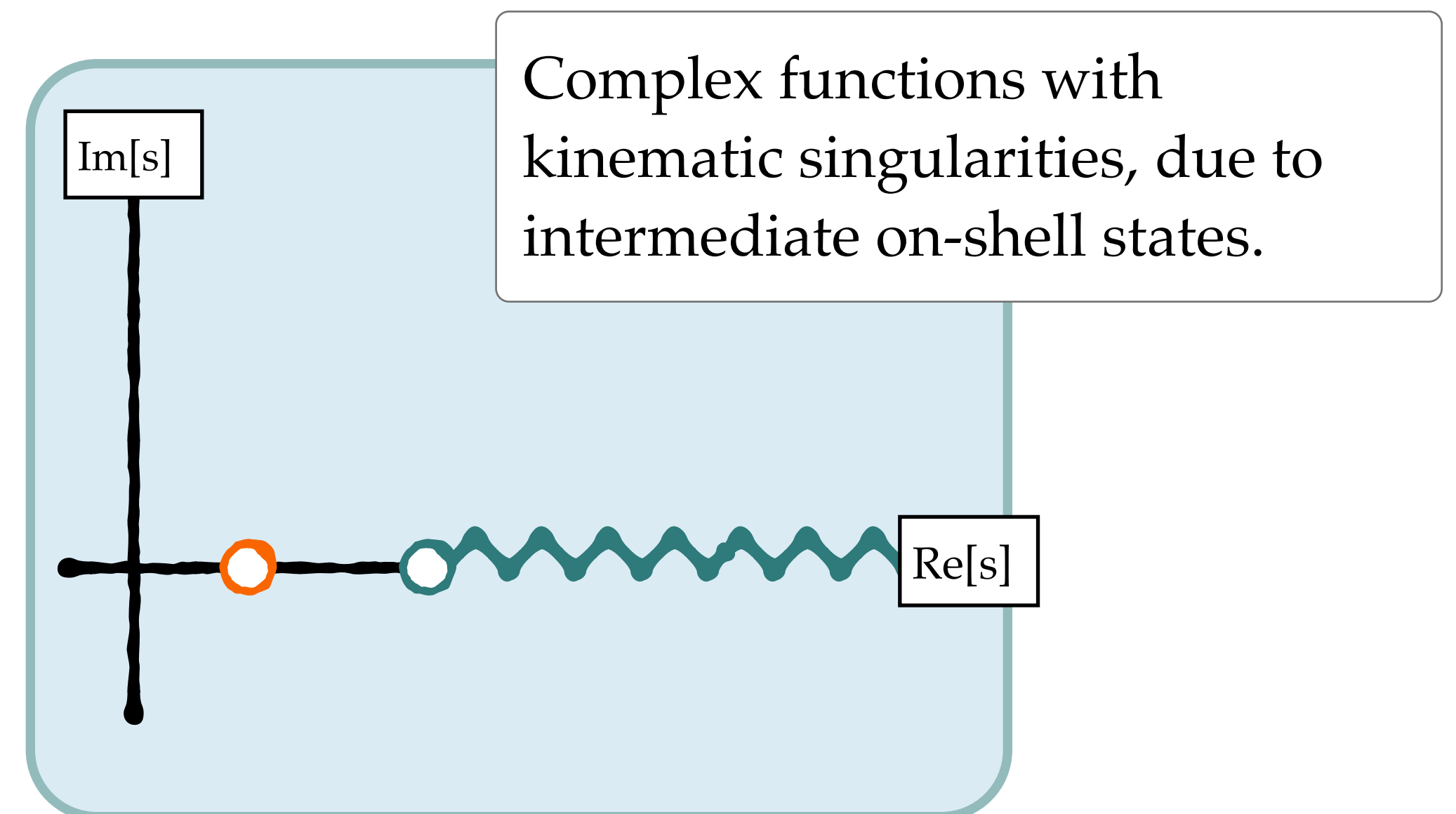
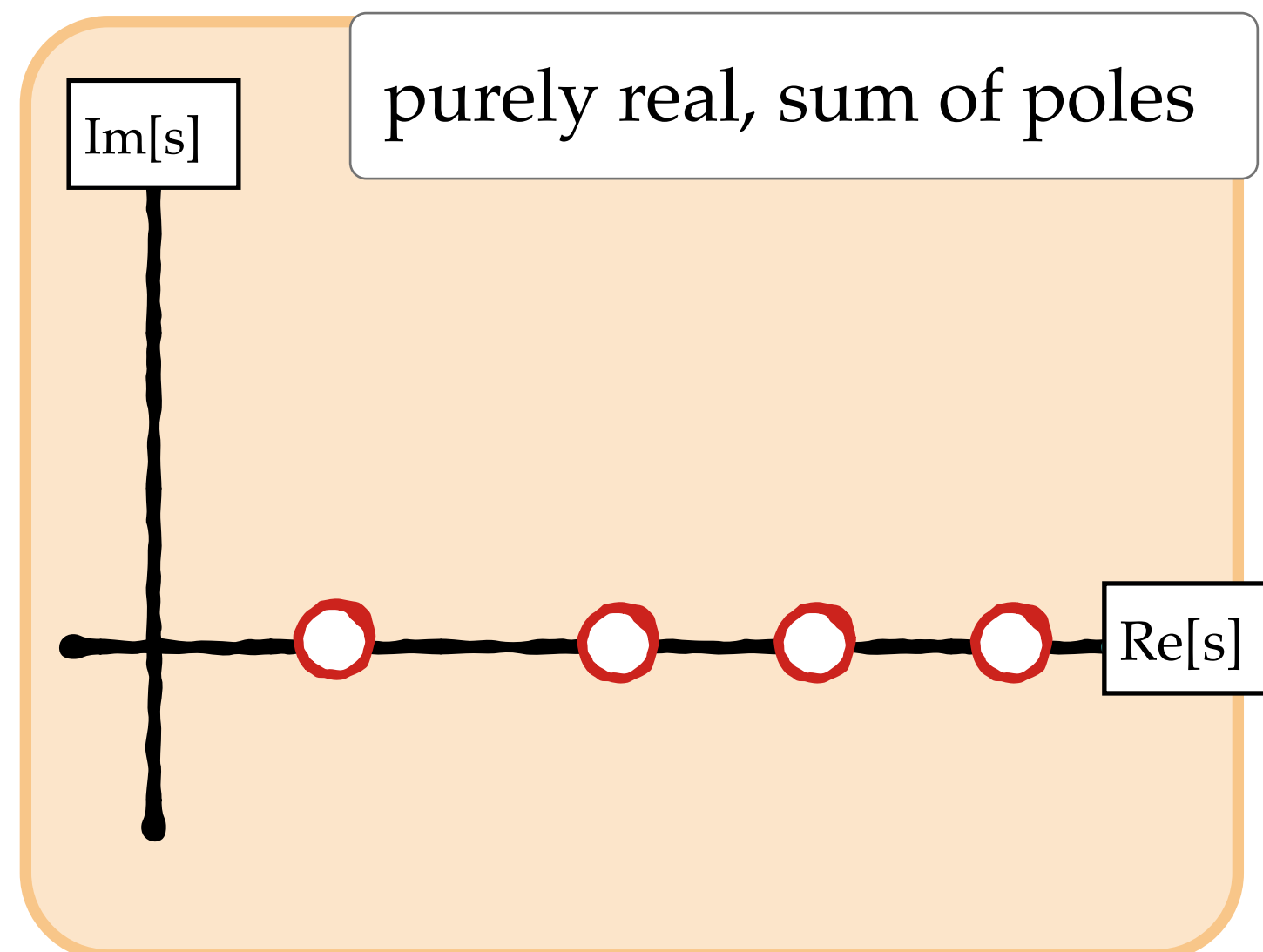
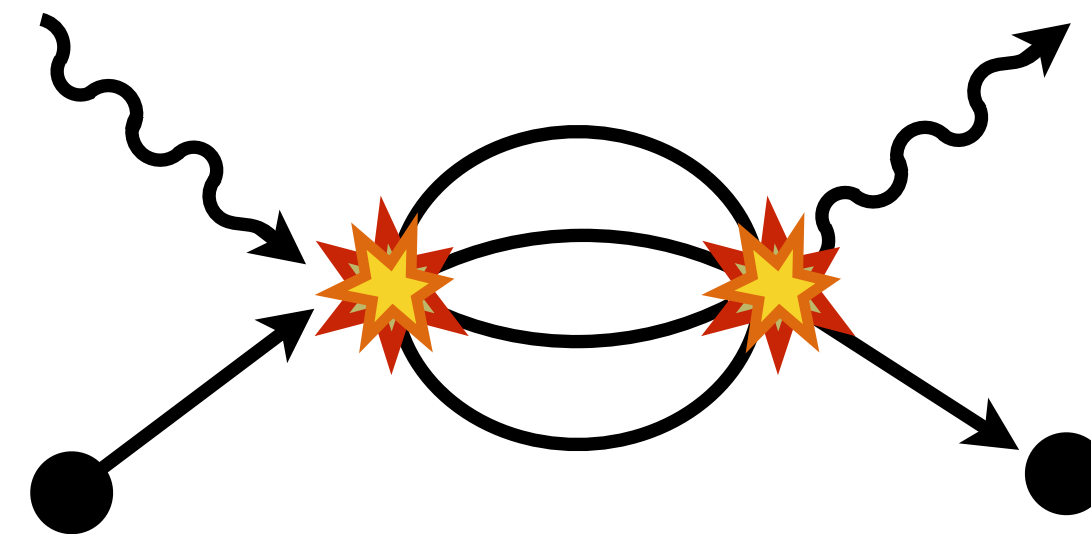
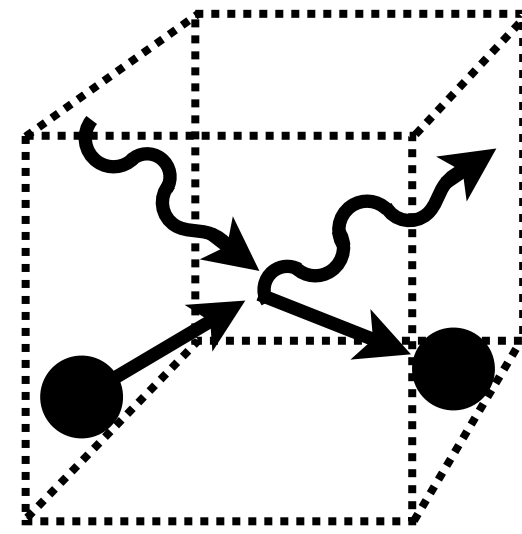
[assuming $\epsilon/E \ll 1$]

Hamiltonian frameworks

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sum over real-valued poles...
clearly unphysical 🤨



Assessing feasibility *in 1+1D*

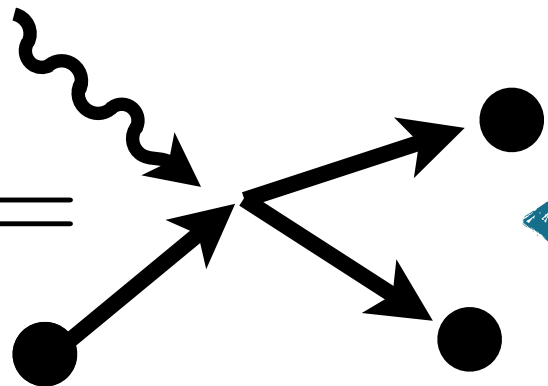
- ❑ (hopefully) interesting and not obviously straightforward.
- ❑ Extrapolation is not even obviously well defined.
- ❑ To explore and test new ideas, we can use existing formalism:
 - ❑ **exact** relationship between finite- and infinite-volume amplitude,
 - ❑ derived using principles of scattering theory.
- ❑ For simplicity: assume scalar currents and current hadrons, and $(2m)^2 < s < (3m)^2$

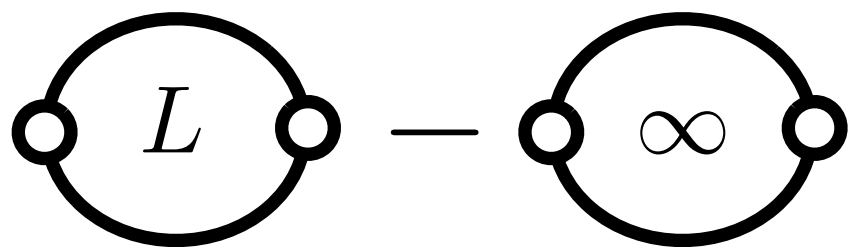
$$\mathcal{T}_L = \mathcal{T} - \mathcal{H}(s, Q_f^2) \frac{1}{F^{-1}(P, L) + \mathcal{M}(s)} \mathcal{H}(s, Q_i^2)$$

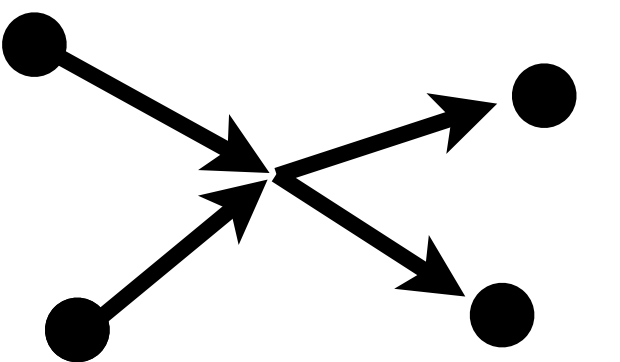
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$i\mathcal{H}(s, Q_f^2) =$


$iF(P, L) =$


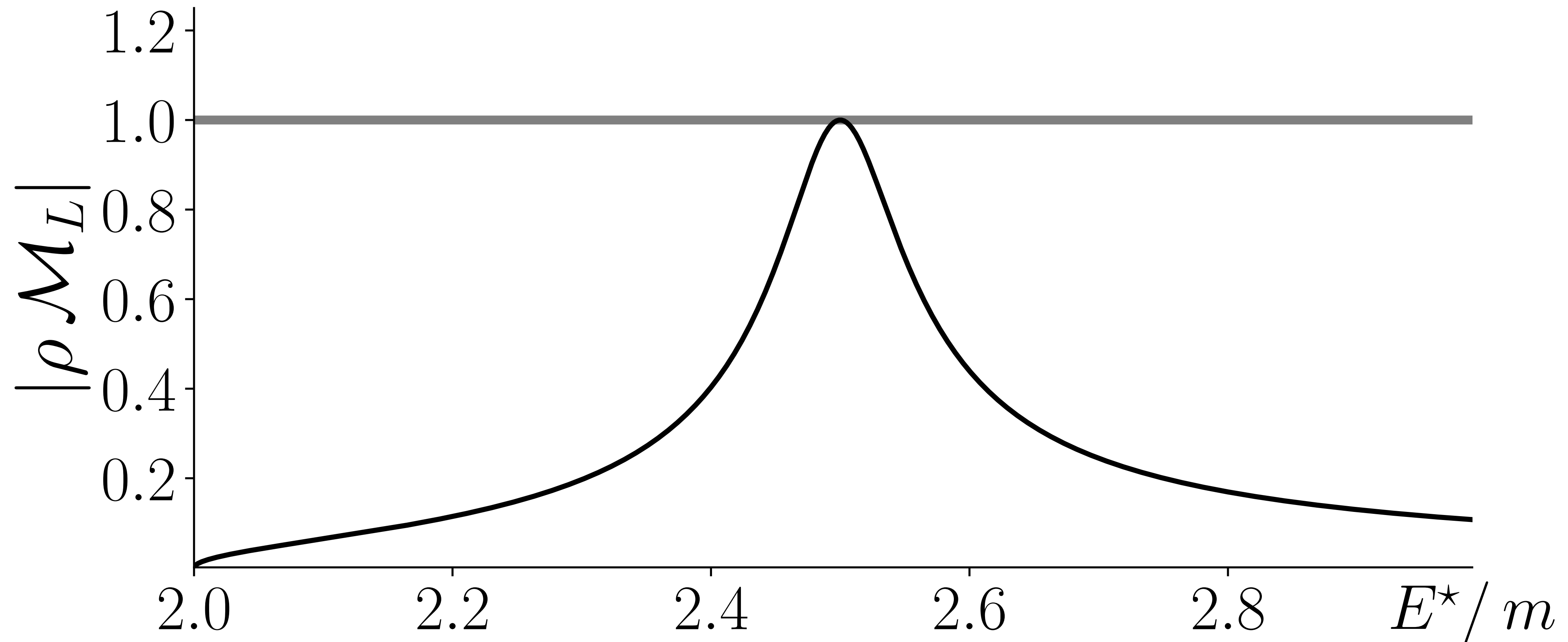
$i\mathcal{M}(s) =$


→ (from $\mathcal{H}(s, Q_f^2)$ to the first vertex)
 → (from $F^{-1}(P, L)$ to the denominator)
 → (from $\mathcal{M}(s)$ to the denominator)

How bad are these effects? *real bad!*

In a finite-volume, the spectrum and all observables depend on the total momentum $P = 2\pi d/L$, where d is discrete.

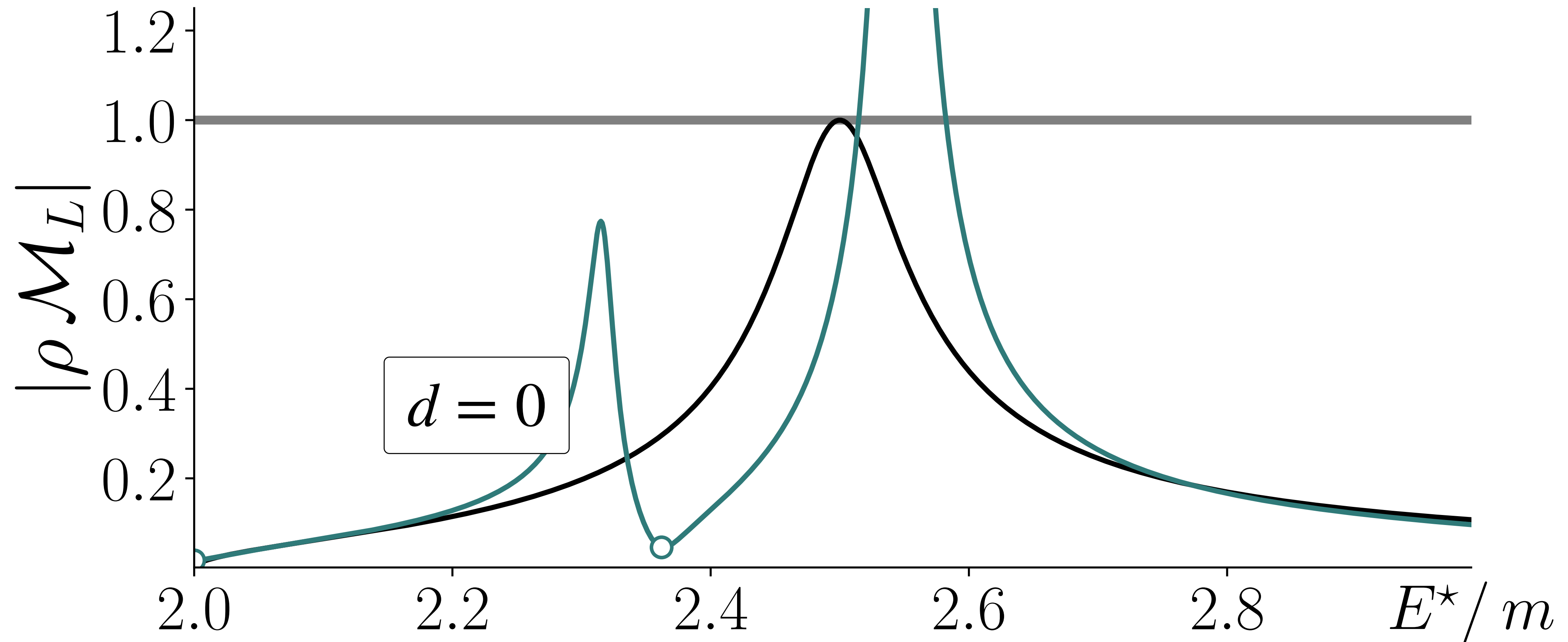
$$\mathcal{M}_L(s) = \mathcal{M}(s) - \mathcal{M}(s) \frac{1}{F^{-1}(P, L) + \mathcal{M}(s)} \mathcal{M}(s)$$



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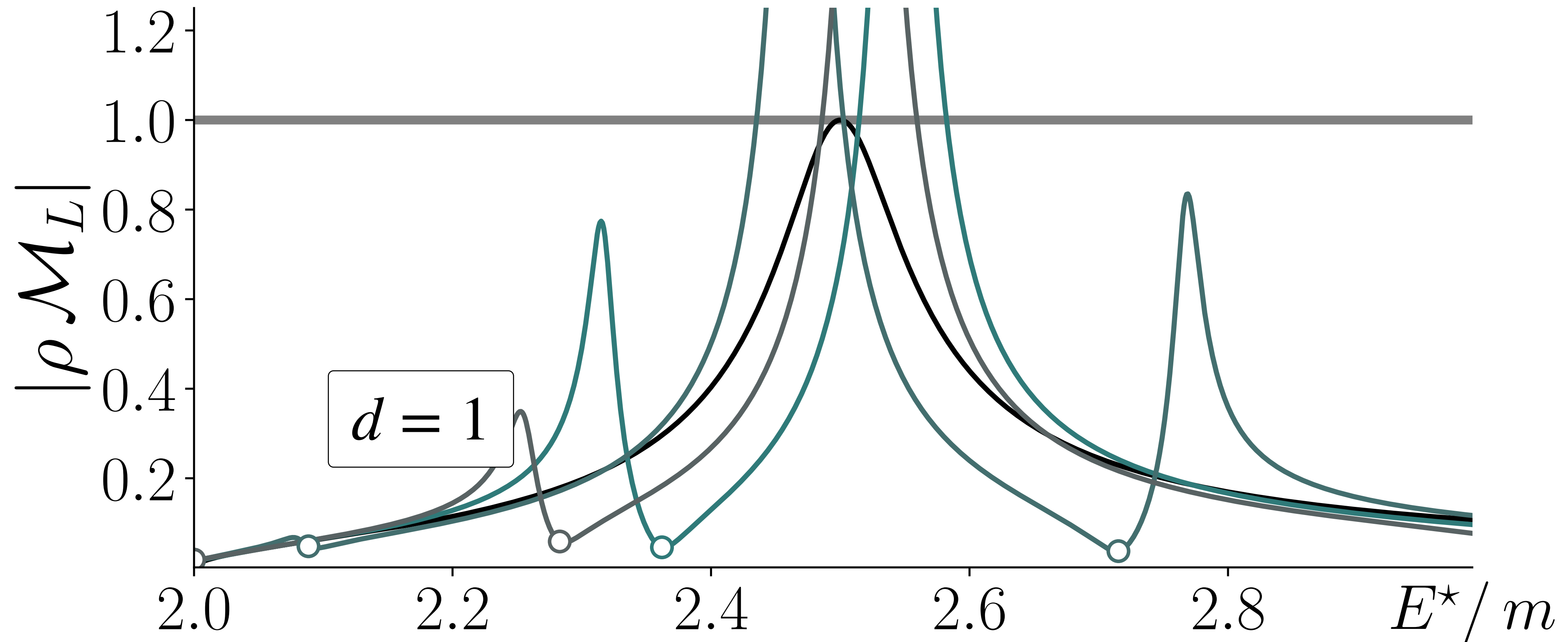
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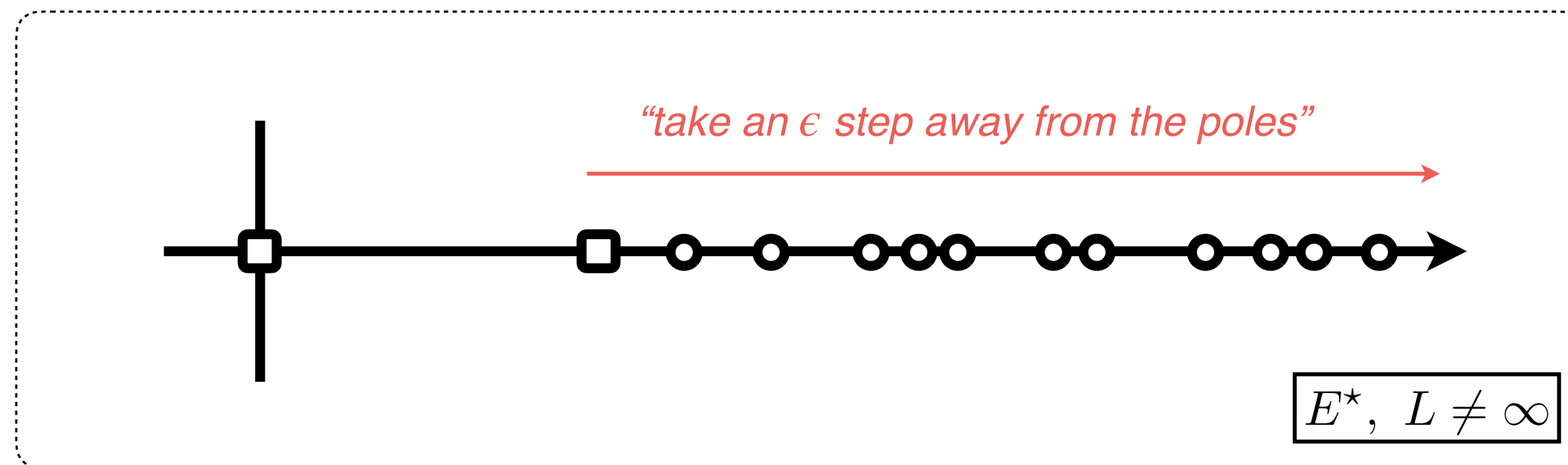
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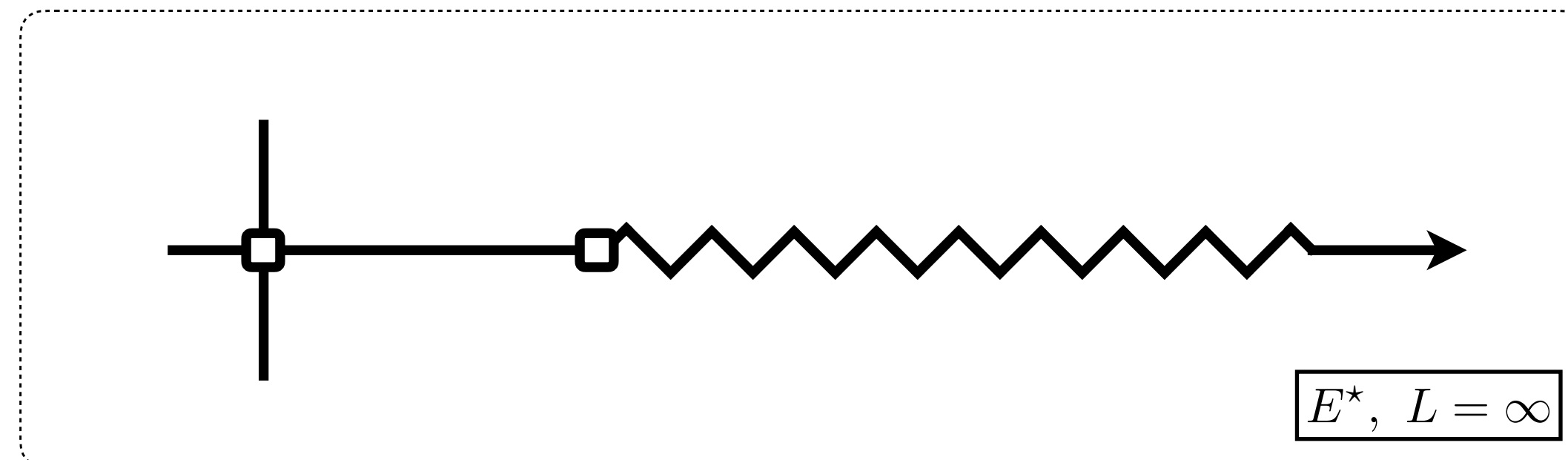
Constructing reliably estimators

- ☑ Determine time-dependent matrix elements [easier said than done 😅]
- ☑ Introduce an $i\epsilon$ by hand

$$\mathcal{T}_L(\epsilon) \sim \int_{-\infty}^{\infty} d\tau e^{iq_0 t - \epsilon|t|} \langle n_f | T[\mathcal{J}_2(t) \mathcal{J}_1(0)] | n_i \rangle_L$$



For large enough L
and small enough ϵ



Constructing reliably estimators

☑ Determine time-dependent matrix elements [easier said than done 😅]

☑ Introduce an $i\epsilon$ by hand [makes sense 😊]

$$\mathcal{T}_L(\epsilon) \sim \int_{-\infty}^{\infty} d\tau e^{iq_0 t - \epsilon|t|} \langle n_f | T[\mathcal{J}_2(t) \mathcal{J}_1(0)] | n_i \rangle_L$$

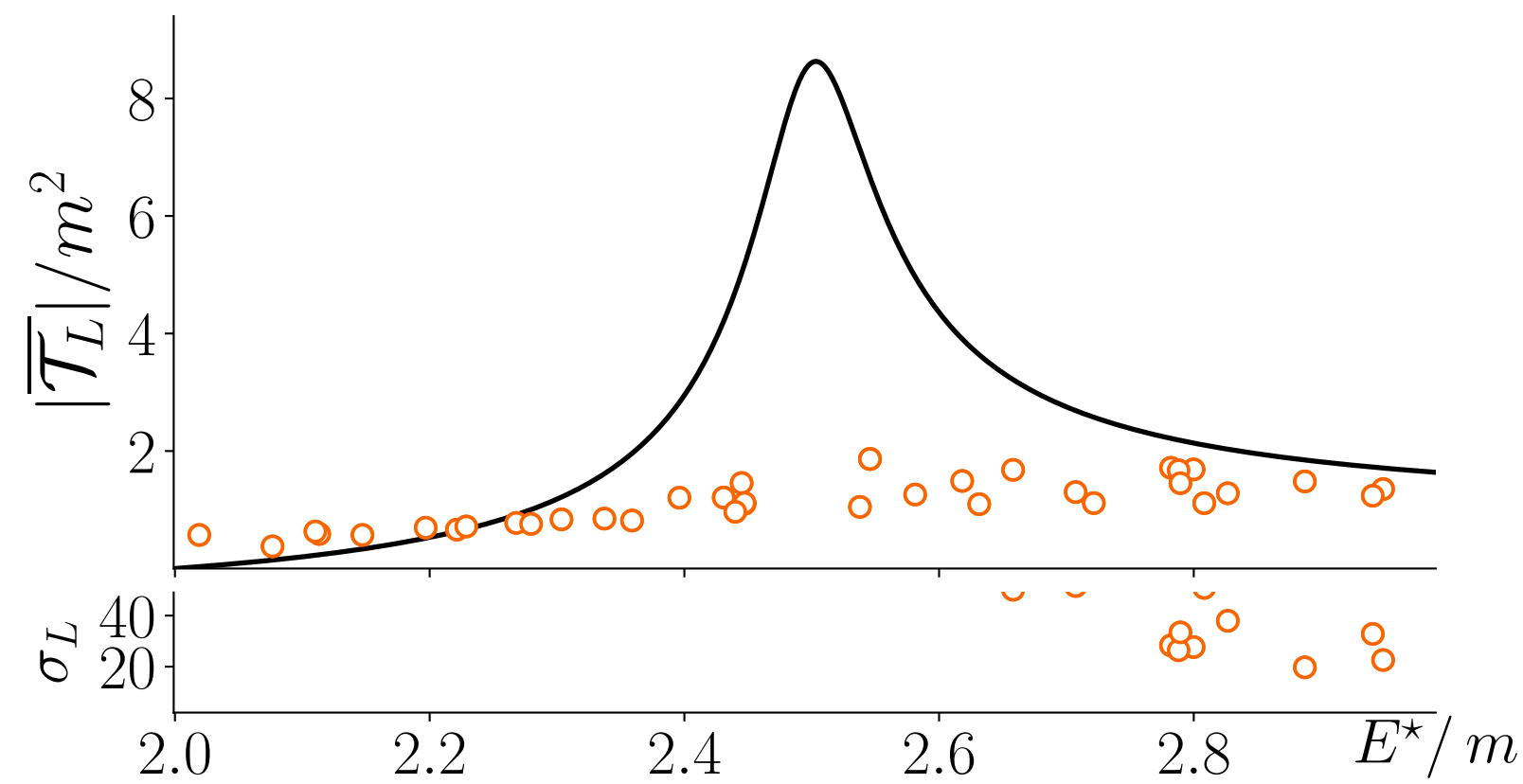
☑ Binning / wave packets [makes sense 😊]

☑ Exploit symmetry: [uh? 🤔]

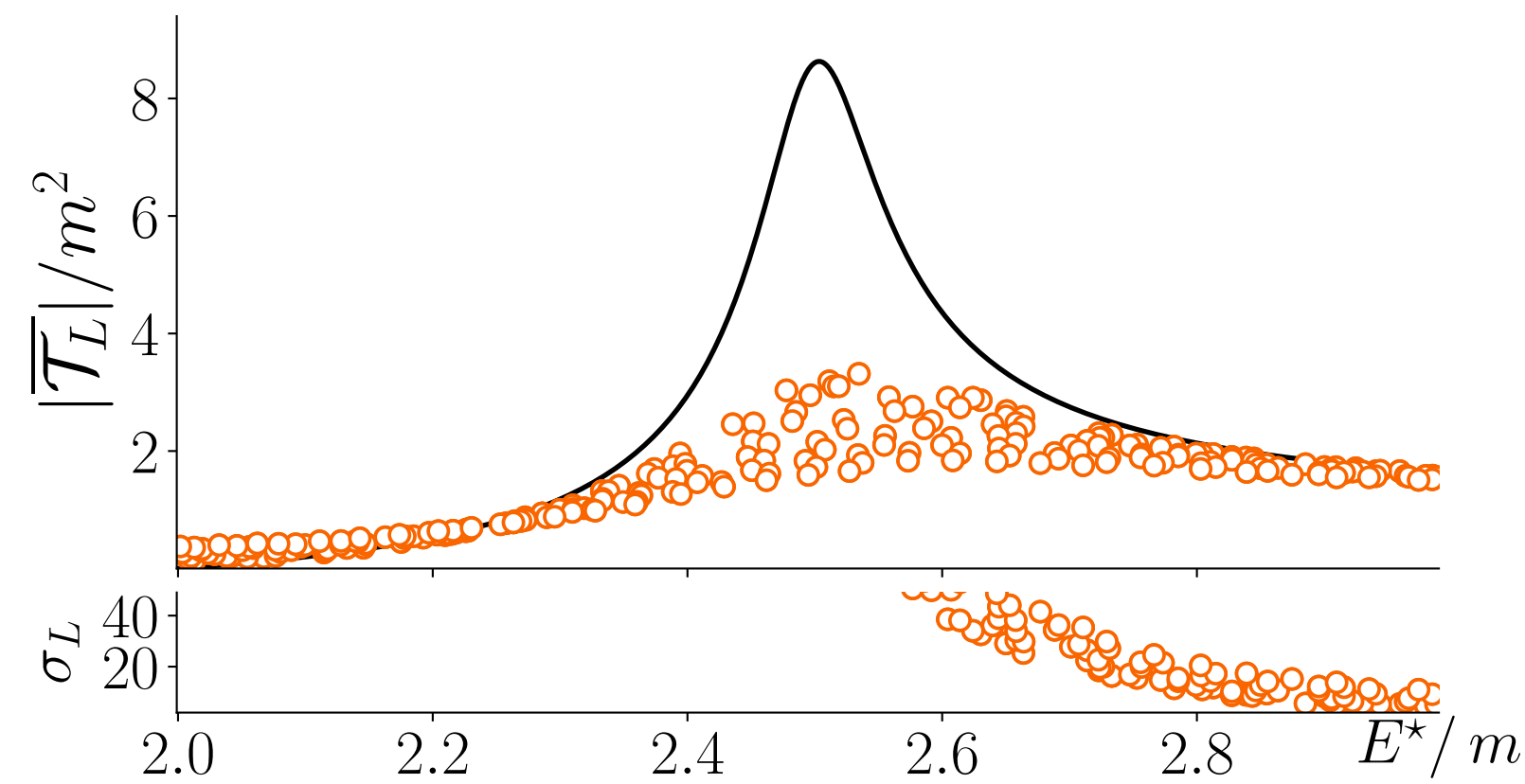
☑ Physical amplitudes only depend on Lorentz scalars.

☑ Boost average

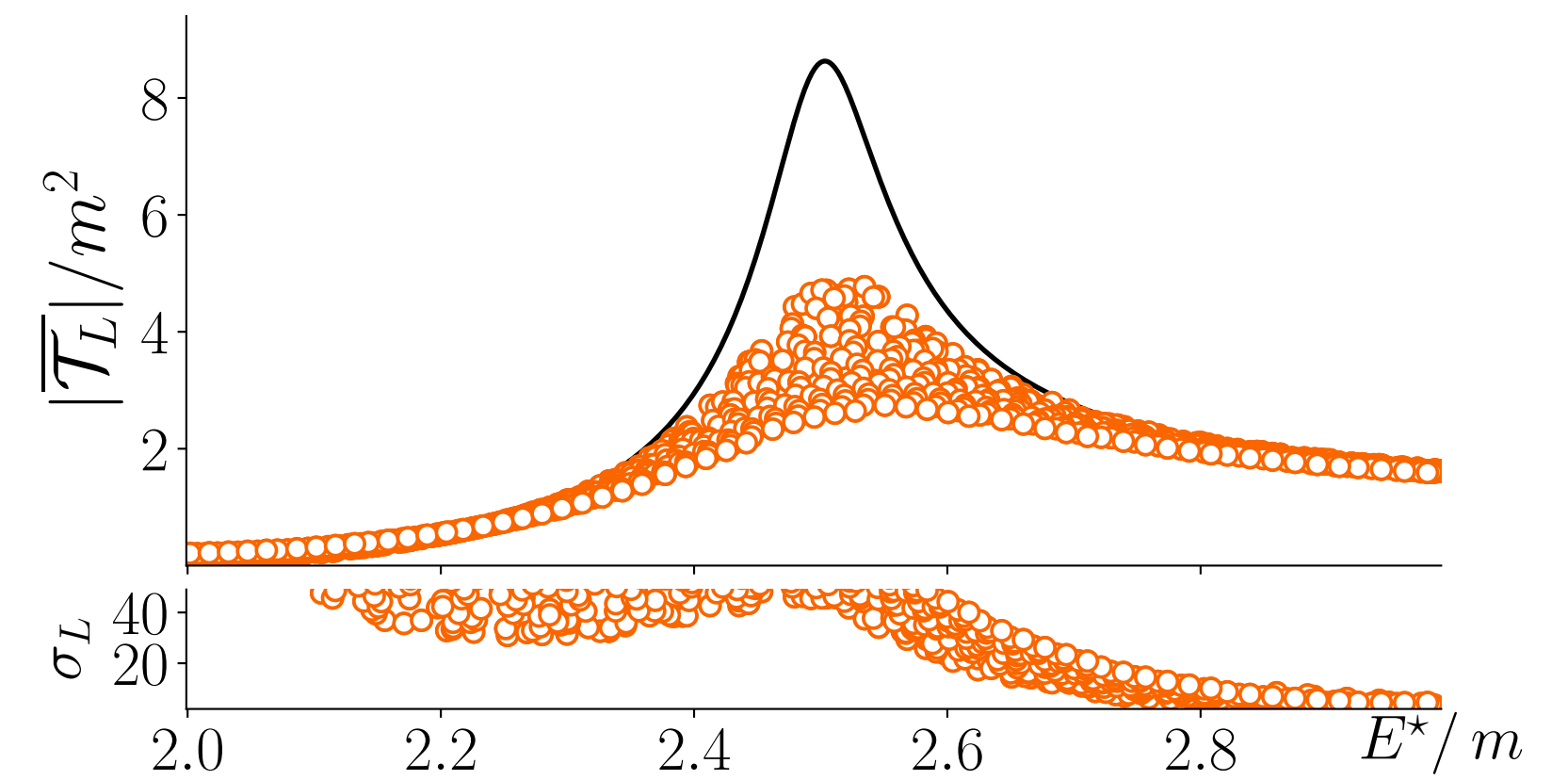
Toy model investigation for \mathcal{T}



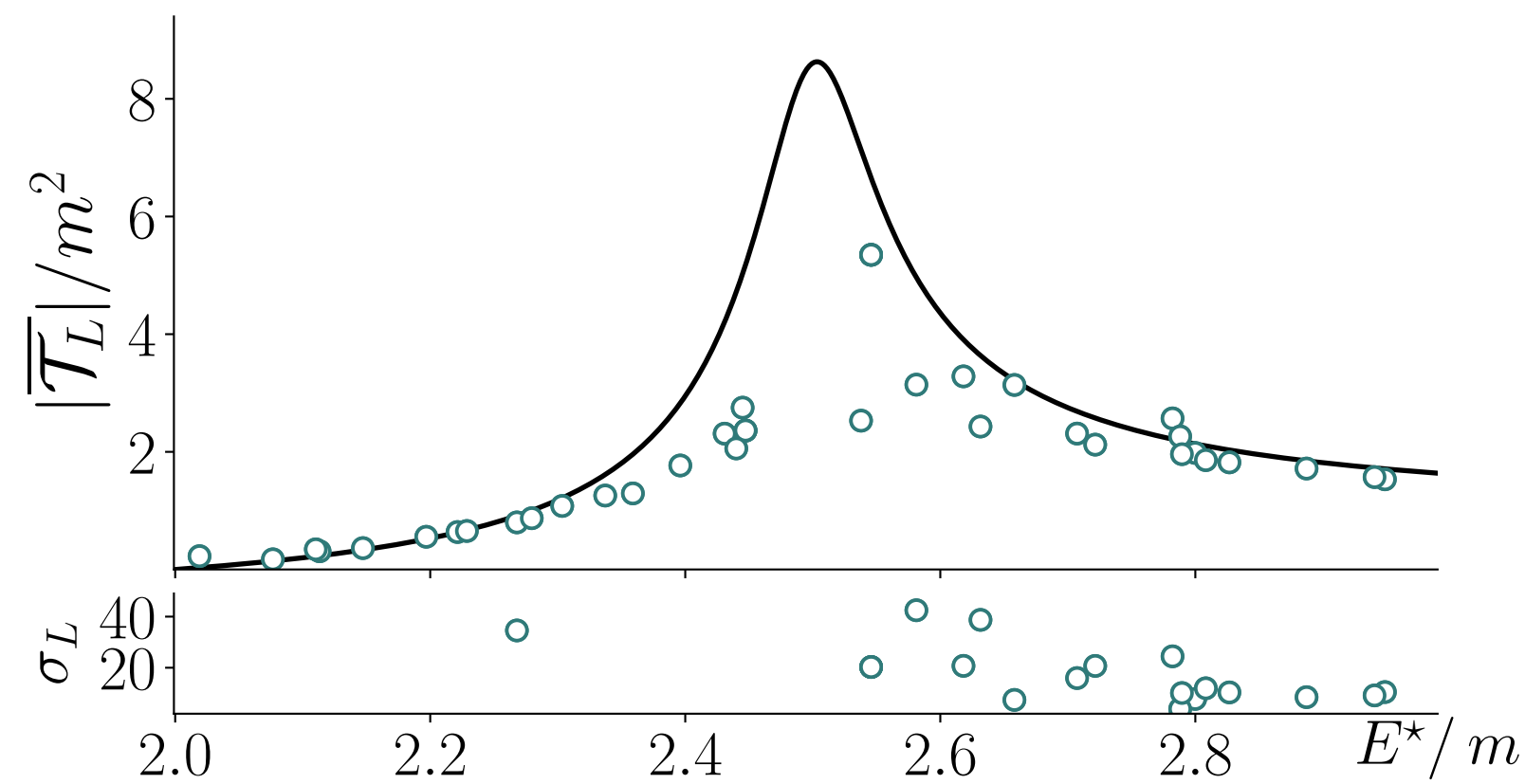
$mL = 20, \epsilon L = 4$



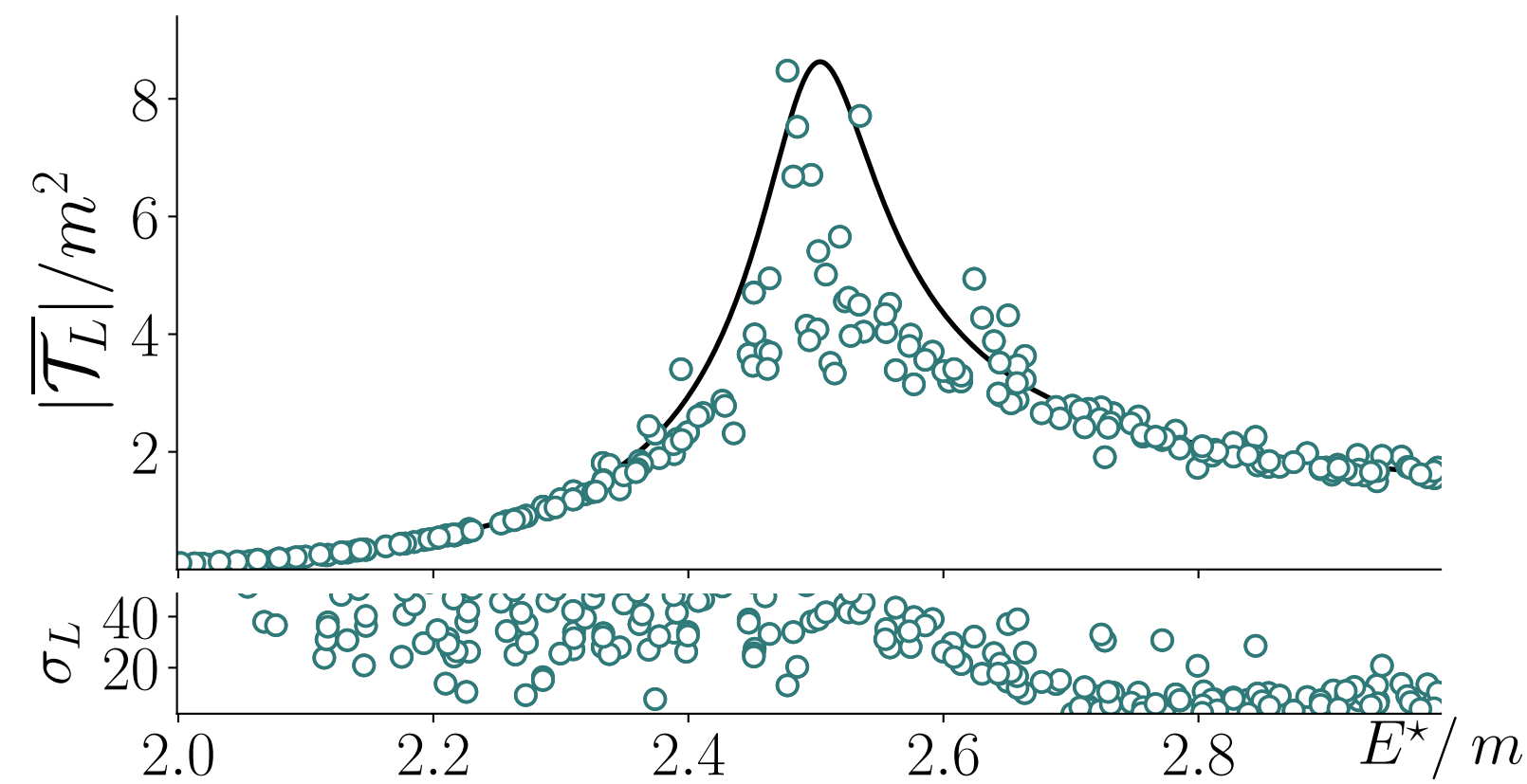
$mL = 50, \epsilon L = 4$



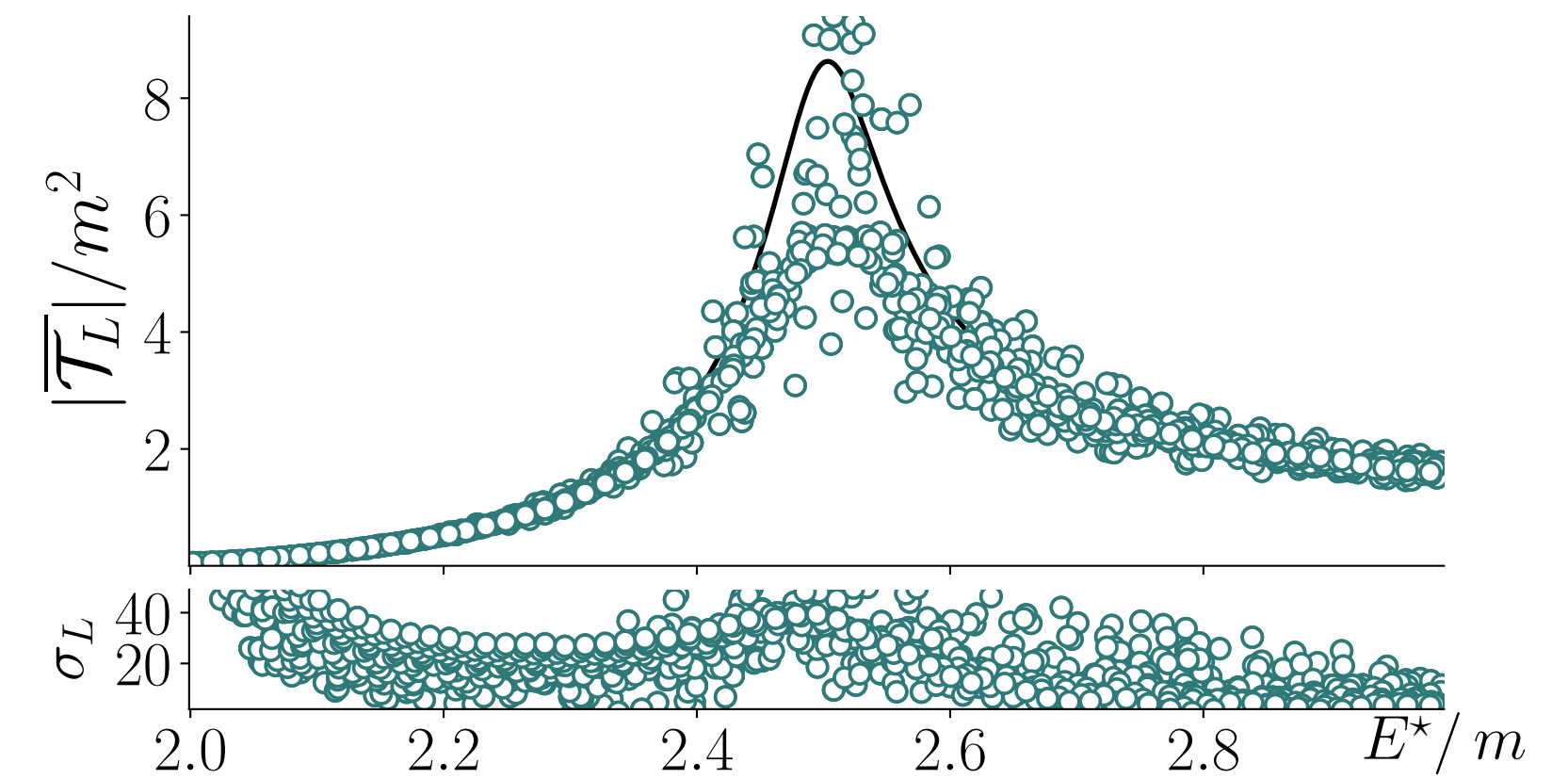
$mL = 100, \epsilon L = 4$



$mL = 20, \epsilon L = 1$



$mL = 50, \epsilon L = 1$



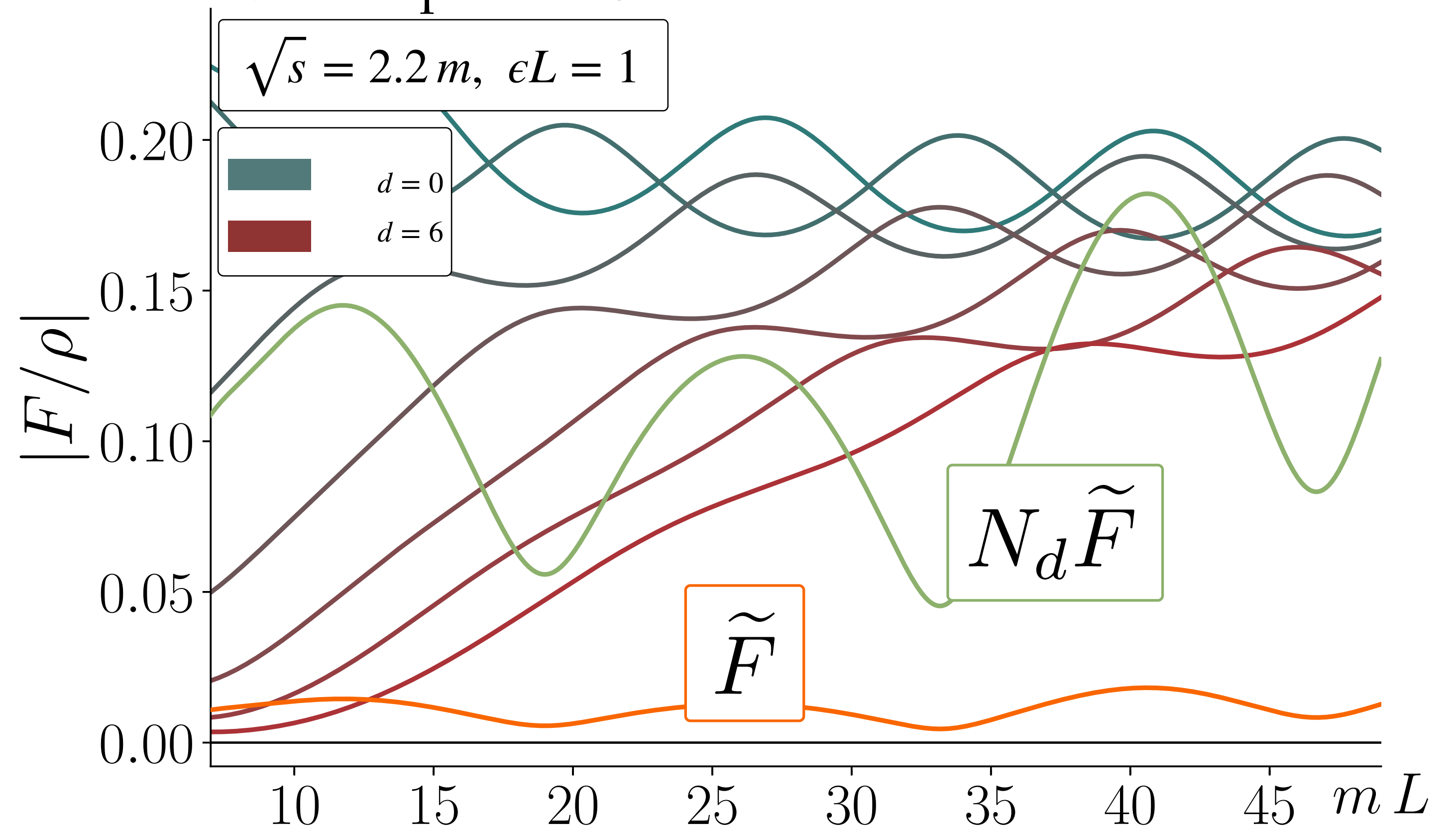
$mL = 100, \epsilon L = 1$

Boost averaging

The volume effects are encoded in $F(P, L)$, which is not a Lorentz scalar.

Asymptotic behavior $F \sim e^{-L\epsilon\alpha_0} (-1)^d$

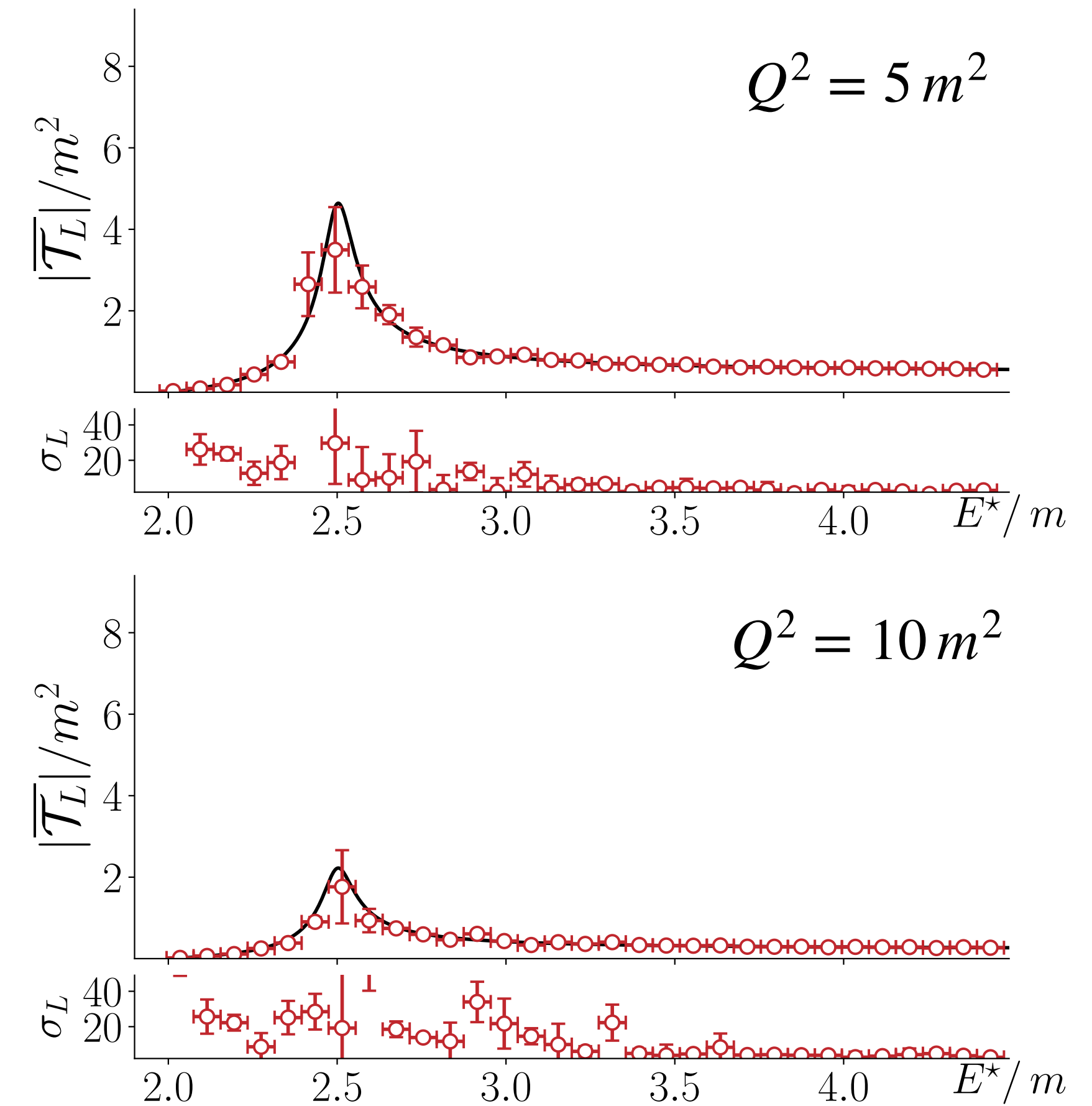
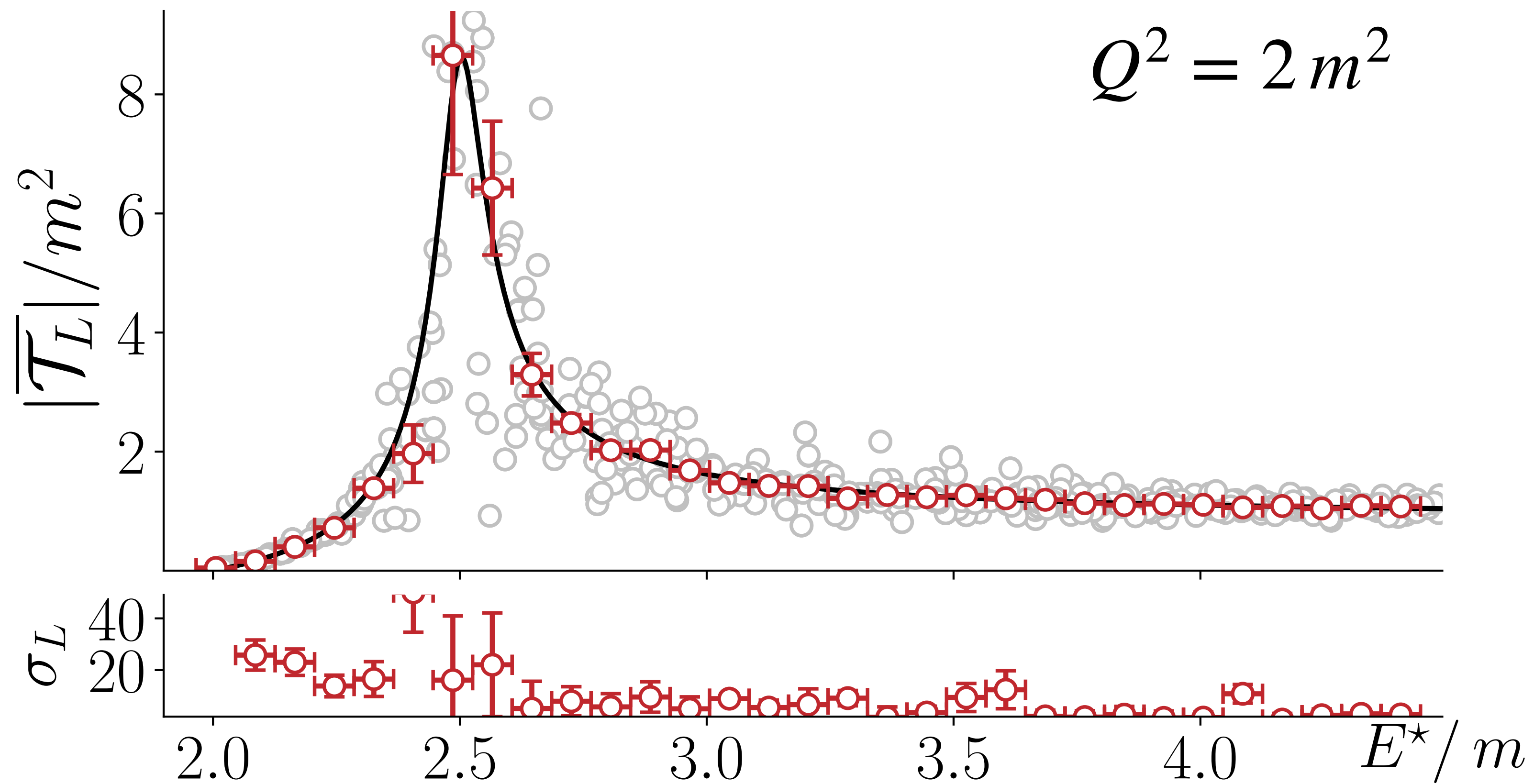
Averaging over different boosts, volume effects are expected to reduce.



$$\mathcal{T}_L = \mathcal{T} - \mathcal{H}(s, Q^2) \frac{1}{F^{-1}(P, L) + \mathcal{M}(s)} \mathcal{H}'(s, Q_{if}^2)$$

Following the recipe

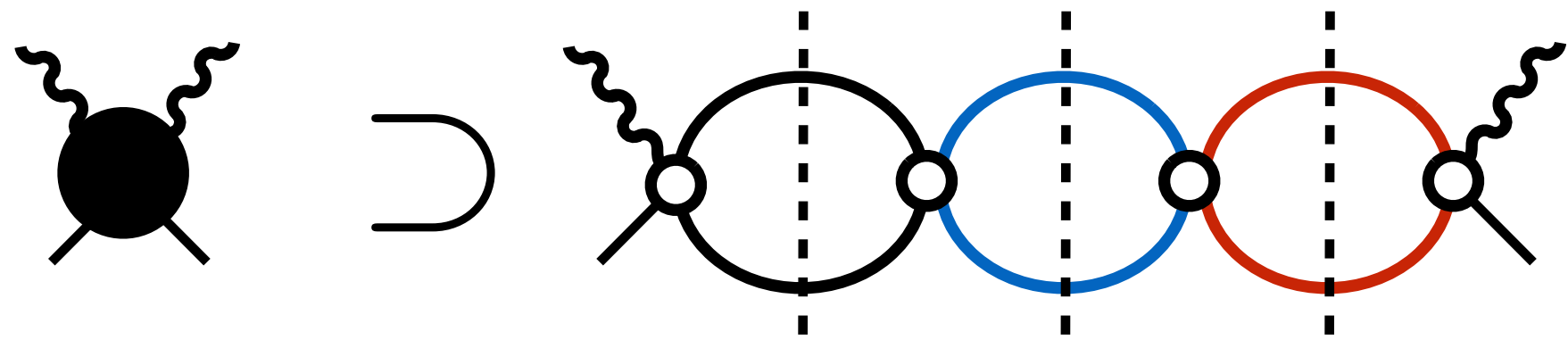
By averaging over $mL = [20,25,30]$ boost with $d \leq mL$, and binning in energy and virtualities.



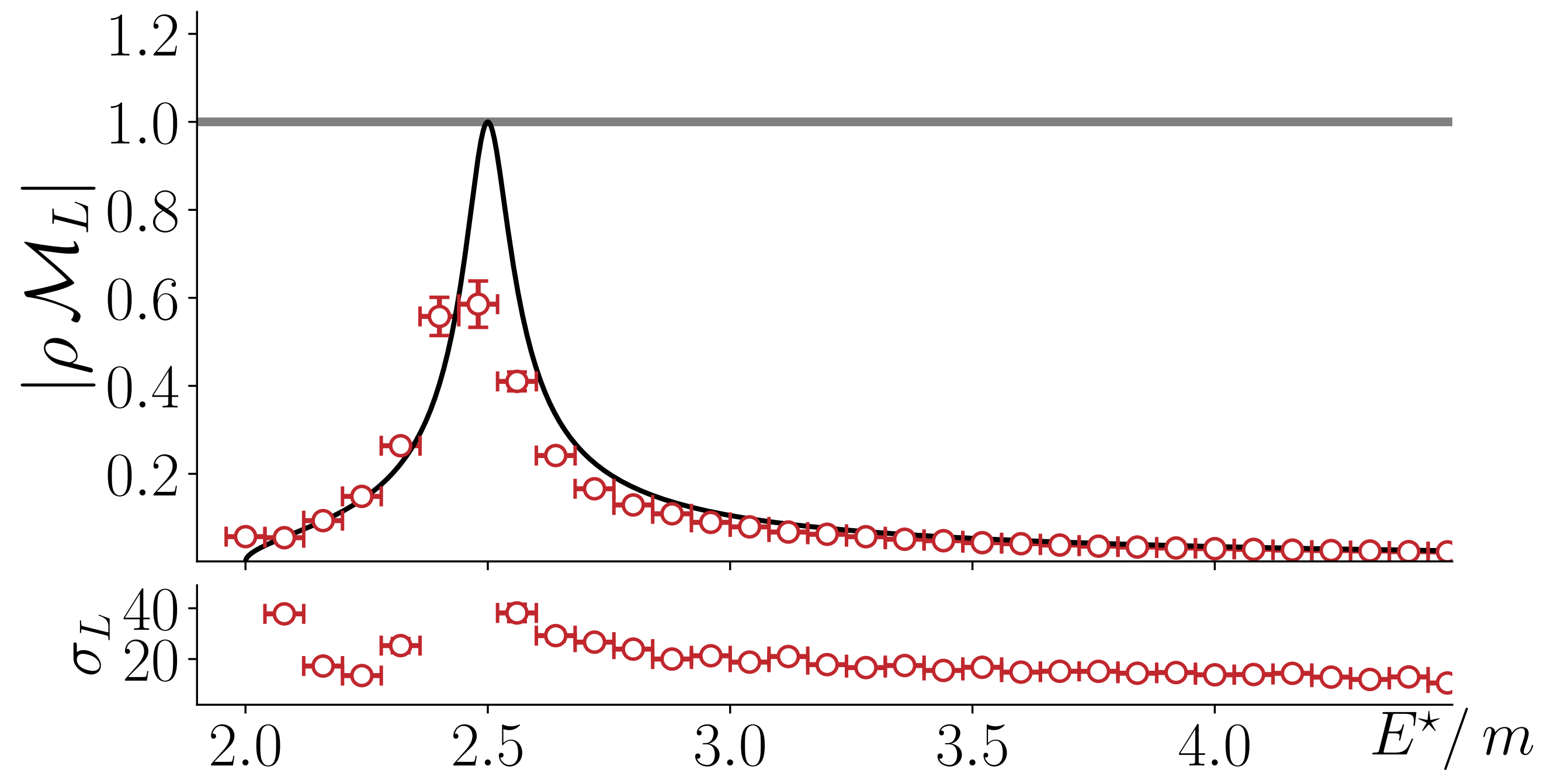
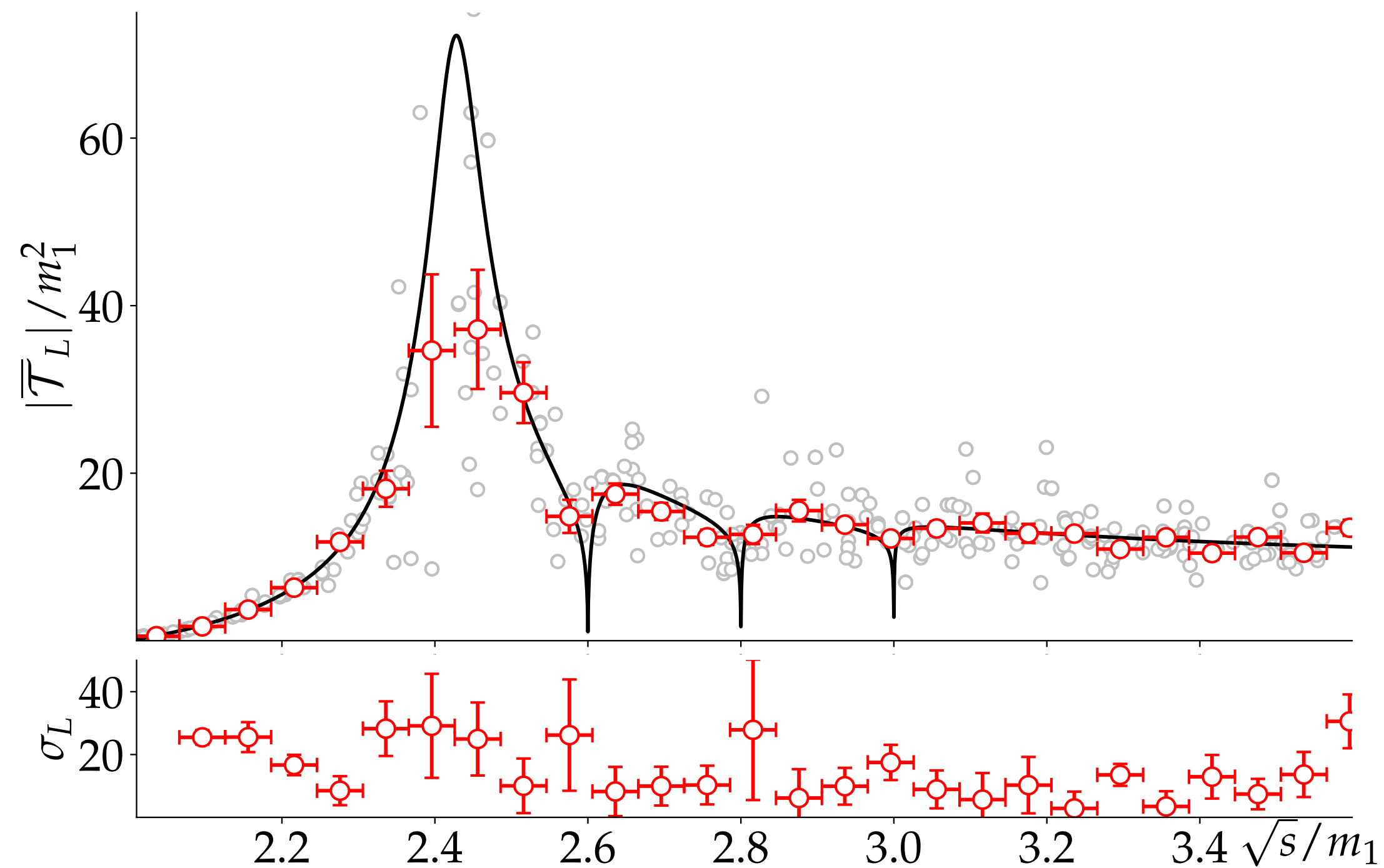
Extensions *other amplitudes*

☑ Checks on arbitrary number of channels

☑ Purely hadronic amplitudes using LSZ



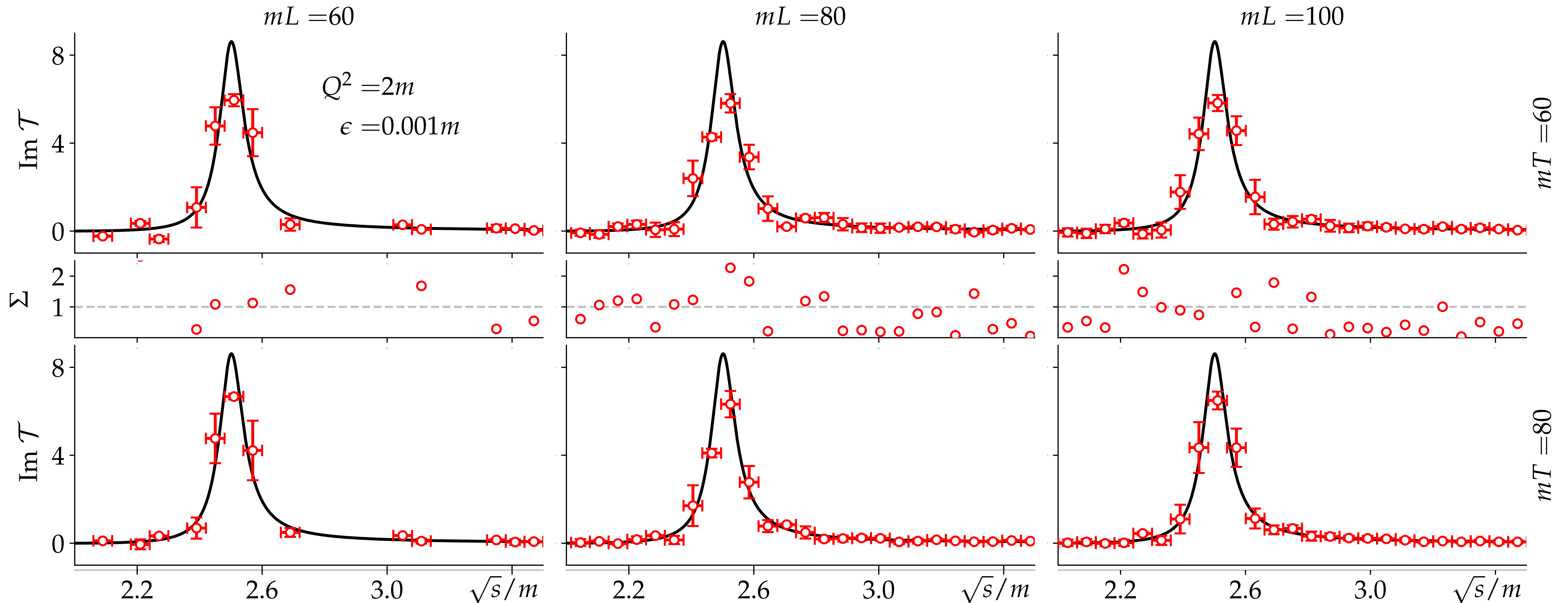
$$\propto \frac{\mathcal{M}(s)}{(q_i^2 - m^2)(q_f^2 - m^2)}$$



Testing finite-time dependence

☑ Using spectral decomposition, we can estimate the order of magnitude of time (T) needed

$$\mathcal{T} \sim \int_0^T d^4x e^{it(\omega+i\epsilon)} \langle n_f | \mathcal{J}(t) \mathcal{J}(0) | n_i \rangle_\infty = \sum_n \int_0^T d^4x e^{it(E_f+\omega-E_n+i\epsilon)} \langle n_f | \mathcal{J}(0) | n \rangle \langle n | \mathcal{J}(0) | n_i \rangle_\infty$$



Take-home message

- ☑ Inclusive scattering observables are not out of the question
- ☑ No sophisticated formalism is needed
 - ☐ existing formalism serves as diagnostic tool
 - ☐ Does life get harder or easier in 3 + 1D?
 - ☐ Test on toy theory [quantum simulation vs. lattice]

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Alex Sturzu
William & Mary



Max Hansen
Edinburgh



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