

# ATLAS recent EFT results

*Eleonora Rossi*

on behalf of the ATLAS Collaboration

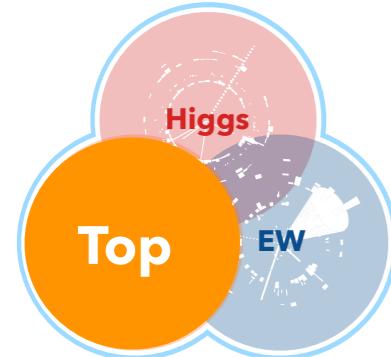


5th General Meeting of the  
LHC EFT Working Group  
23/05/2022



# Latest results

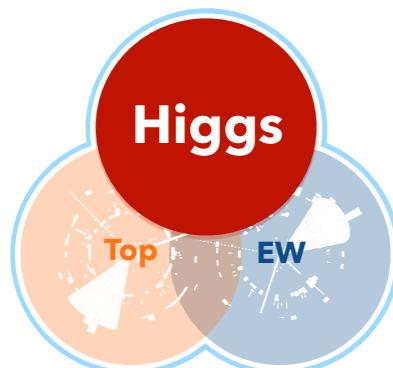
Summary of new results that came out after the 4th LHCEFT General meeting on 23rd May 2022: [Talk](#)



Sketch from Rahul inspired by Ken Mimasu

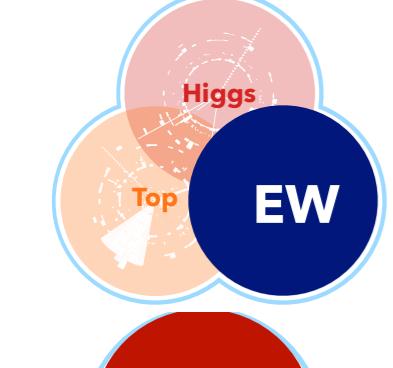
- $tH(\tau\tau)q$  FCNC: [arXiv:2208.11415](https://arxiv.org/abs/2208.11415) (TopFCNC UFO)
- $t\bar{t}$  charge (rapidity) asymmetry: [arXiv:2208.12095](https://arxiv.org/abs/2208.12095) (SMEFT@NLO)

[Top-quark LHC measurements in the SMEFT](#)



- $H \rightarrow \gamma\gamma$ : STXS measurements:  
[CERN-EP-2022-094](https://cds.cern.ch/record/2683422) (SMEFTsim + SMEFT@NLO)

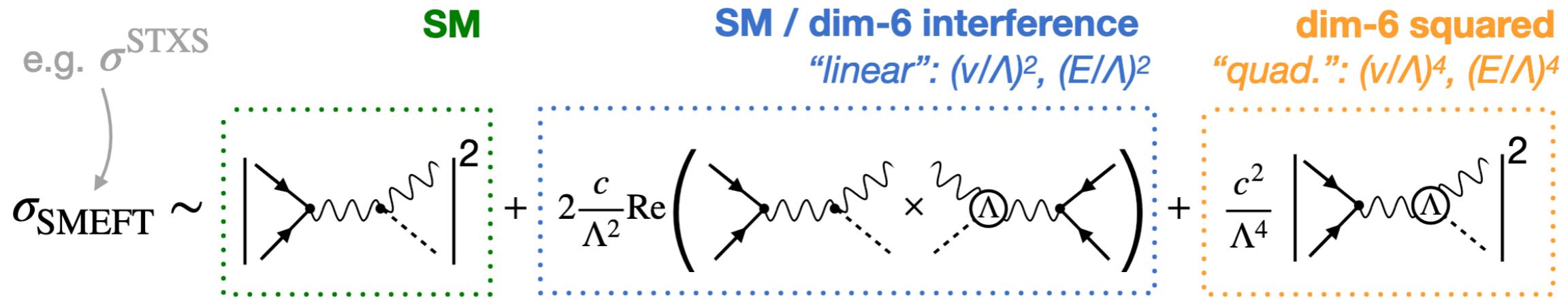
- Electroweak  $Z(\nu\bar{\nu})\gamma jj$  production and limits on anomalous quartic gauge couplings: [arXiv:2208.12741](https://arxiv.org/abs/2208.12741)



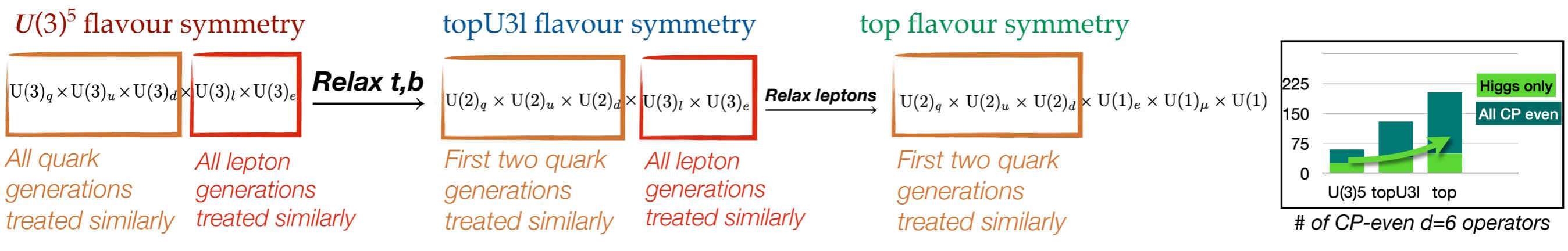
- First ATLAS Global combination (Higgs+EW+EWPO):

[ATL-PHYS-PUB-2022-037 - \(SMEFTsim + SMEFT@NLO\)](https://cds.cern.ch/record/2683422)

# EFT: theoretical framework



- Warsaw basis, assuming  $\Lambda = 1$  TeV.
- SMEFTsim + SMEFT@NLO + TopFCNC.
- Results are usually provided for **linear** model (+ **linear-quadratic** models).
- SMEFTsim: different flavour symmetries used to reduce the number of Wilson coefficients.  
 (" **$U(3)^5$  flavour symmetry**", "**topU3I**") : ( $H \rightarrow \gamma\gamma$ , ATLAS Global combination)



# $tqH(\tau\tau)$ FCNC

[arXiv:2208.11415](https://arxiv.org/abs/2208.11415)

Top

Higgs

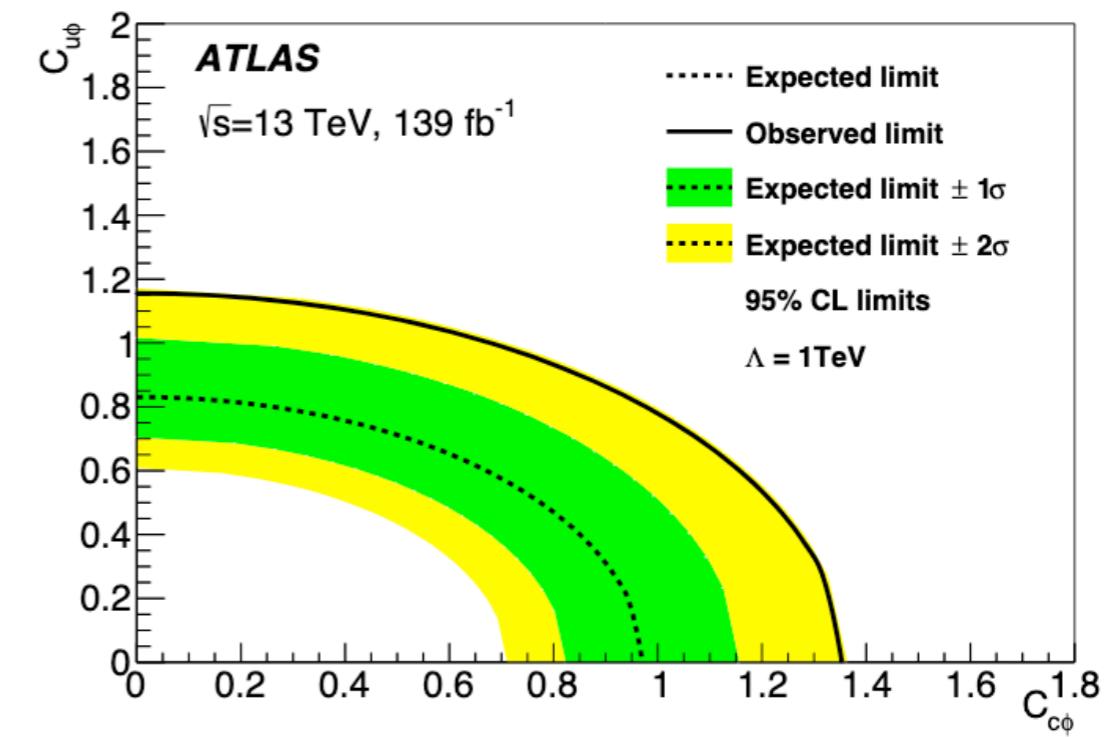
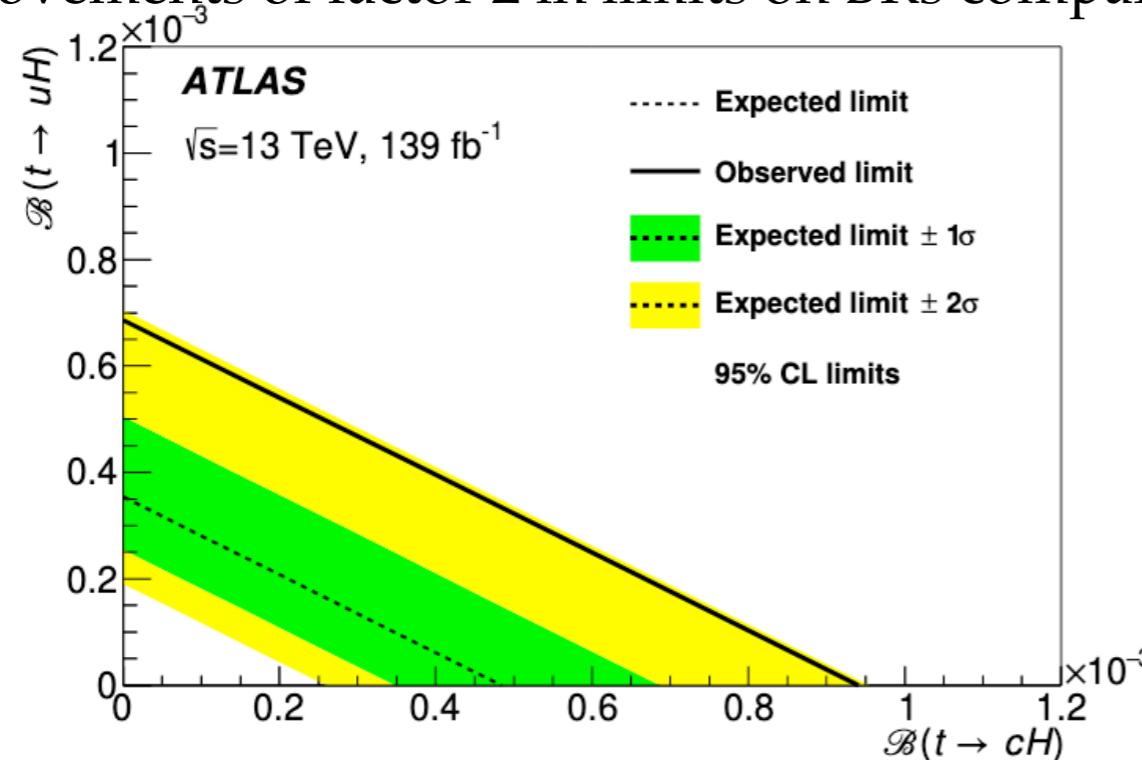
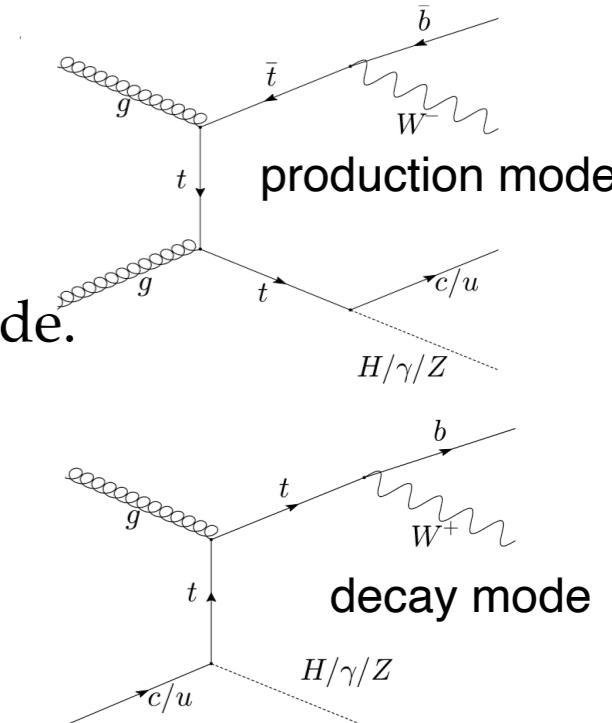
EW



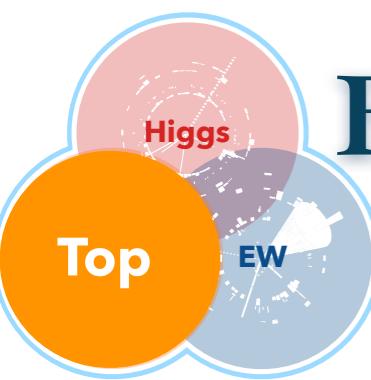
- Search for  $tqH$  vertex focusing on  $H \rightarrow \tau\tau$  in this analysis, where the remaining top decays into a W boson and a b quark - Production & decay mode.
- BDT trained to separate signal and background processes.
- BDT output used in a profile-likelihood fit to extract limits on the FCNC signal processes.
- The FCNC coupling parametrised using dim-6 operators:

$$\mathcal{L}_{EFT} = \frac{C_{u\phi}^{i3}}{\Lambda^2} (\phi^\dagger \phi) (\bar{q}_i t) \tilde{\phi} + \frac{C_{u\phi}^{3i}}{\Lambda^2} (\phi^\dagger \phi) (\bar{t}_i q) \tilde{\phi}$$

- The coefficient  $c_{q\phi}$  can be extracted according to the Madgraph calculation using TopFCNC UFO.
- Improvements of factor 2 in limits on BRs compared to partial Run2 analysis.



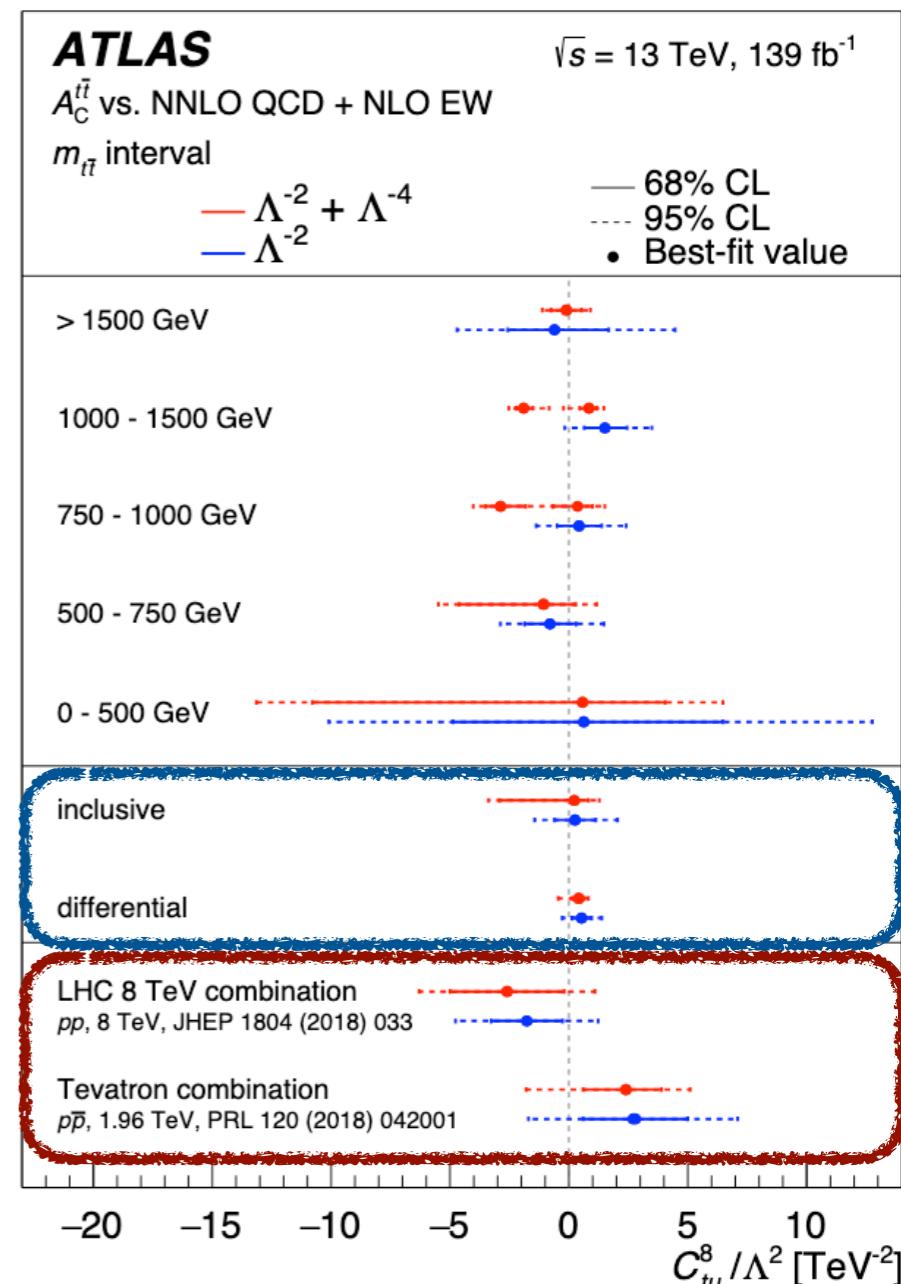
# Evidence for the charge asymmetry in $pp \rightarrow t\bar{t}$



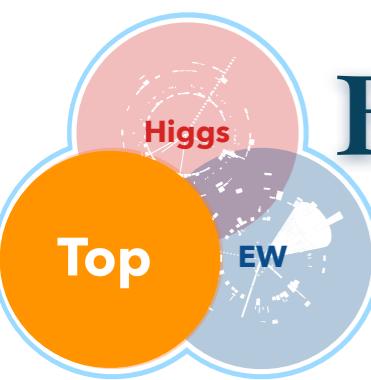
[arXiv:2208.12095](https://arxiv.org/abs/2208.12095)

- Inclusive and differential full Run2 measurements of the top–antitop ( $t\bar{t}$ ) charge asymmetry  $A_c^{t\bar{t}}$  and the leptonic asymmetry  $A_c^{\ell\bar{\ell}}$
- Differential measurements are performed as a function of the invariant mass, transverse momentum and longitudinal boost of the  $t\bar{t}$  system.
- Combined results are interpreted in the SMEFT framework.
- **14 four-fermion operators** + 1 operator for top–gluon interaction.
  
- Large improvement comparing with LHC 8TeV / Tevatron results.
- Interplay between sensitivity, which increases rapidly at higher  $m_{t\bar{t}}$ , and uncertainty, which grows from 0.2%–0.3% in the lowest mass bin to 2.9% in the highest bin.
- For the linear fit, the tightest limit is obtained in the mass bin from 1 to 1.5 TeV.
- Constraint from the differential  $m_{t\bar{t}}$  measurement more than a factor 2 stronger than the one from inclusive measurement (increase in sensitivity with higher  $m_{t\bar{t}}$ ).

The combined inclusive charge asymmetry is measured to be  $A_c^{t\bar{t}} = 0.0068 \pm 0.0015$ , which differs from zero by **4.7 standard deviations**.

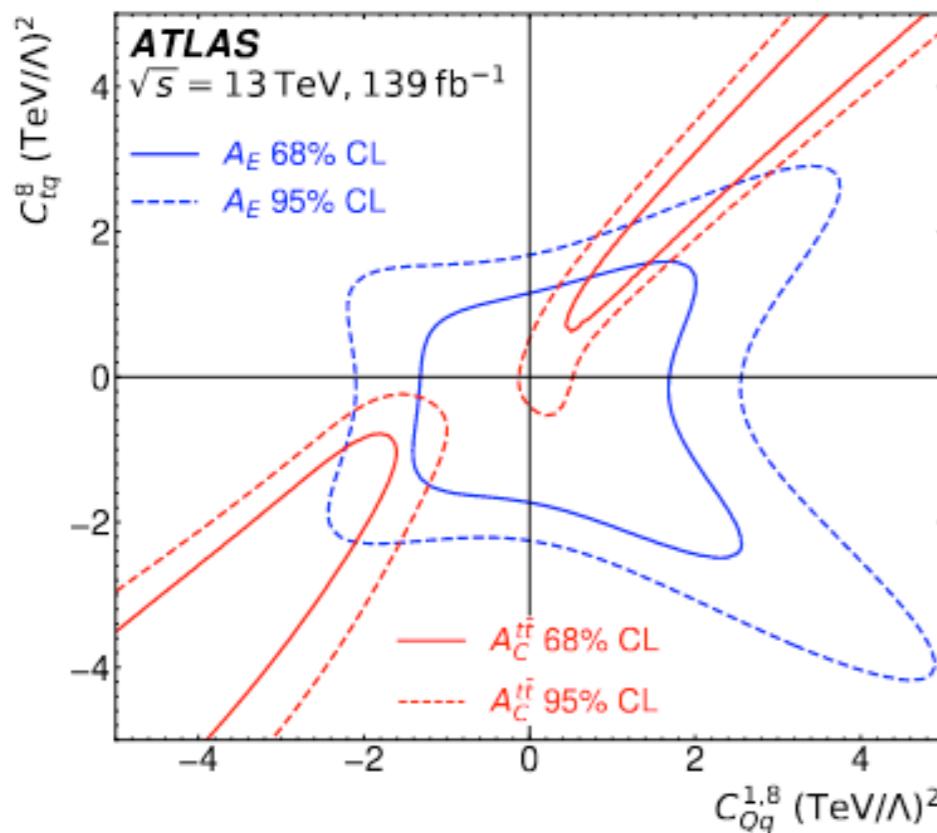


# Evidence for the charge asymmetry in $pp \rightarrow t\bar{t}$



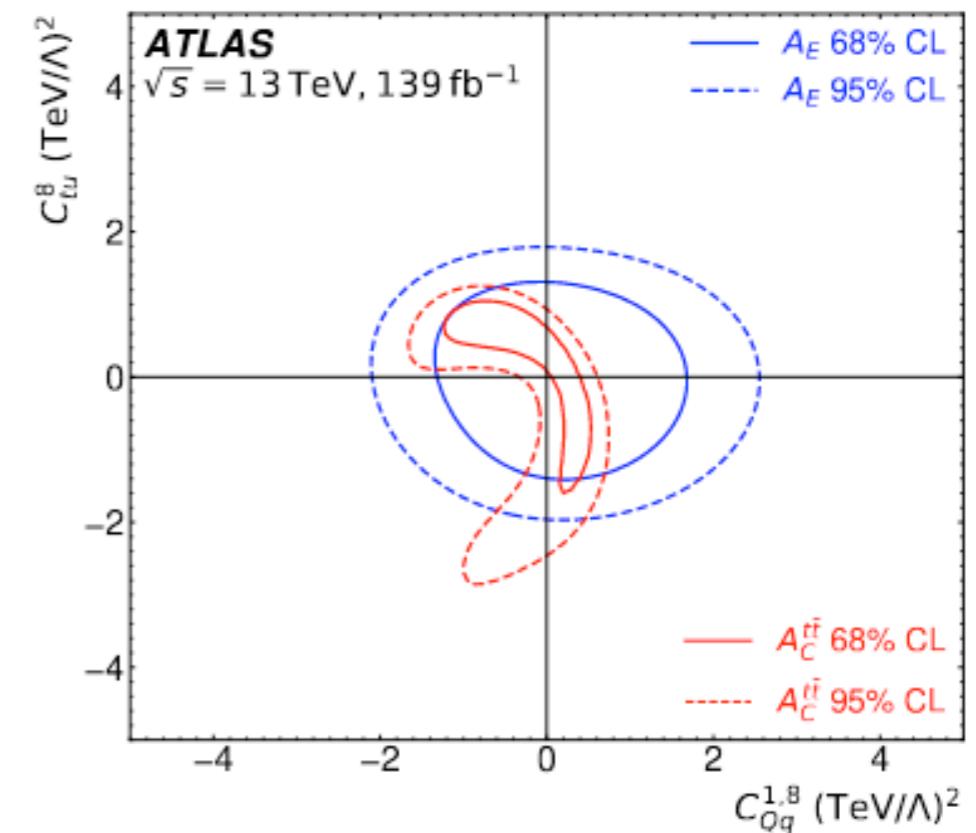
[arxiv:2208.12095](https://arxiv.org/abs/2208.12095)

- Due to the extra jet in  $t\bar{t}j$  production, the QCD structure of the energy asymmetry is not the same as for the charge asymmetry in  $t\bar{t}$  production, so the two asymmetries probe different directions in chiral and colour space.
- For colour-octet operators with the same chirality scenarios the shapes of the bounds look very different: the charge asymmetry (dashed/solid red lines) leaves a blind direction which is broken by the energy asymmetry (dashed/solid blue lines) due to operator interference with the QCD amplitude.

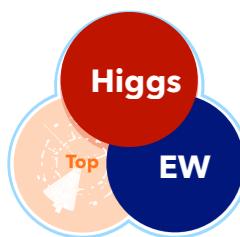


blind direction for  
rapidity  
asymmetry broken  
using  
energy asymmetry  
measurement

Bounds on color-octet  
operators

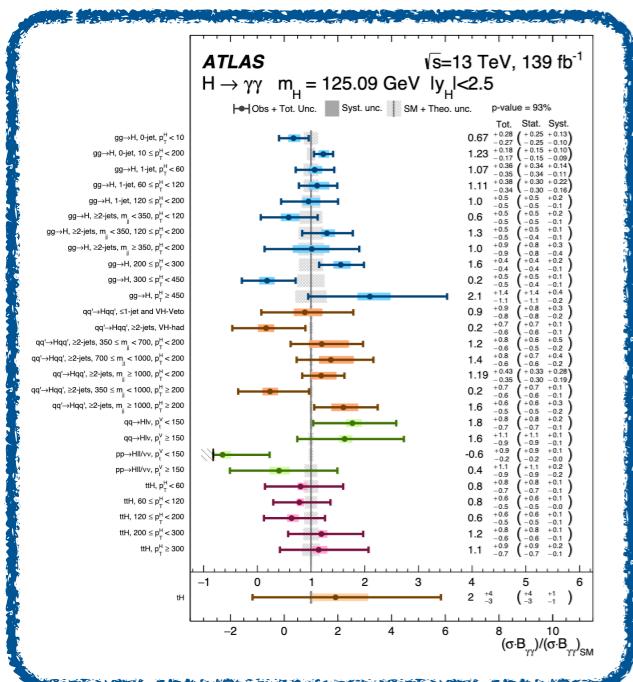


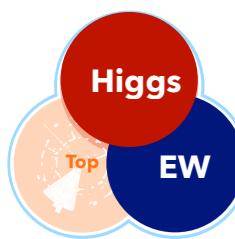
other 2D scans in backup



# Higgs + EW: interpretation workflow

STXS x BR  
Likelihood/Differential  
cross-section  
measurements





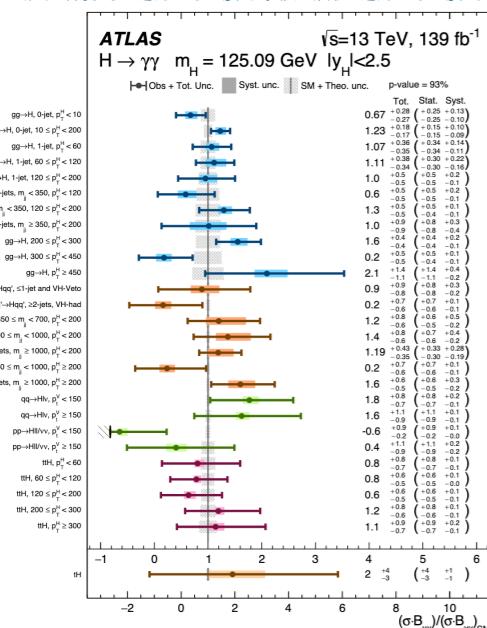
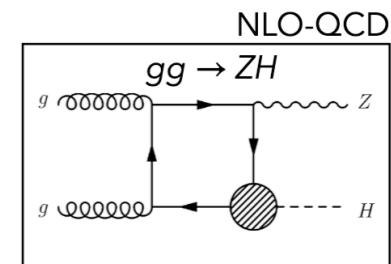
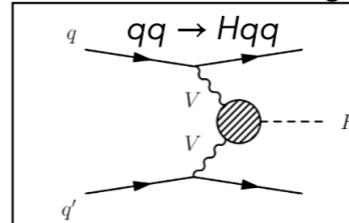
# Higgs + EW: interpretation workflow

**STXS X BR**  
Likelihood/Differential cross-section measurements

re-parameterisation

EFT Likelihood

Tree-level insertion, eg,

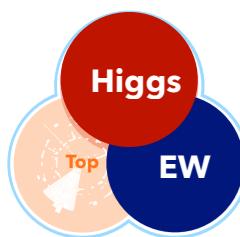


## Parameterisation

- Parameterise EFT dependence at the level of STXS and BR, differential cross sections.
- Analytical Expression in term of the Wilson Coefficients.

- Dimension-six operator effects are calculated:
  - at tree level using SMEFTsim 3.0.
  - for loop-induced processes (e.g  $ggH$ ,  $ggZH$ ): SMEFTatNLO.
  - Analytic formulas for  $H \rightarrow \gamma\gamma$  including NLO EW corrections and LEP observables.
- Acceptance corrections to account for kinematic differences between SM and SMEFT in Higgs boson decays.
- Effects of width changes of intermediate particles (“propagator corrections”) included.



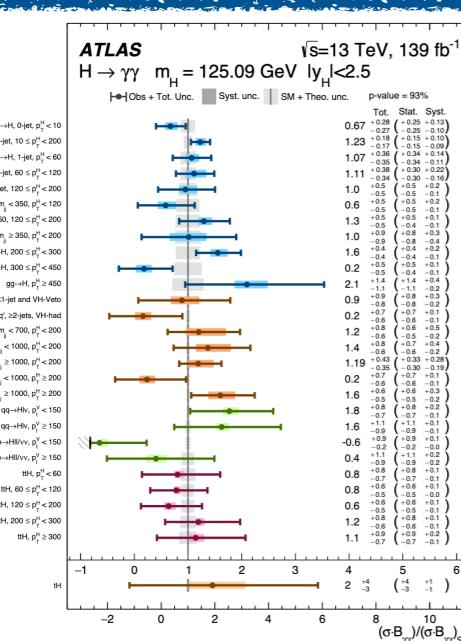


# Higgs + EW: interpretation workflow

STXS x BR  
Likelihood/Differential cross-section measurements

re-parameterisation

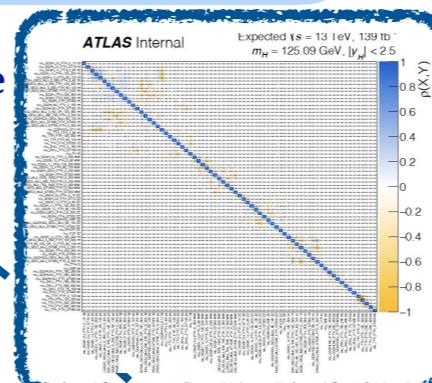
EFT Likelihood



## Parameterisation

- Parameterise EFT dependence at the level of STXS and BR, differential cross sections.
- Analytical Expression in term of the Wilson Coefficients.

covariance matrix



## Sensitivity studies

- Removing degenerate directions in the likelihood.
- Keep only the sensitivity directions for final fits.



# Higgs + EW: interpretation workflow

STXS x BR  
Likelihood/Differential cross-section measurements

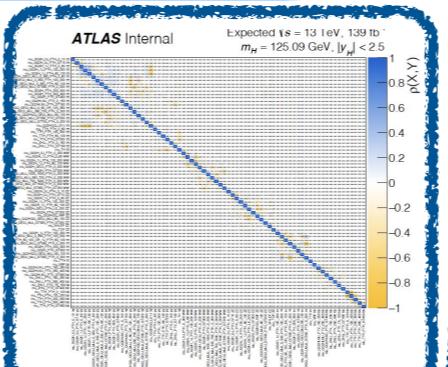
re-parameterisation

EFT Likelihood

sensitive directions

Re-functioned EFT Likelihood

covariance matrix



Sensitivity studies

- Removing degenerate directions in the likelihood.
- Keep only the sensitivity directions for final fits.

Parameterisation

- Parameterise EFT dependence at the level of STXS and BR, differential cross sections.
- Analytical Expression in term of the Wilson Coefficients.

**SMEFT<sub>sim</sub>**  
v3.0



**SMEFT@NLO**

Fit

Results

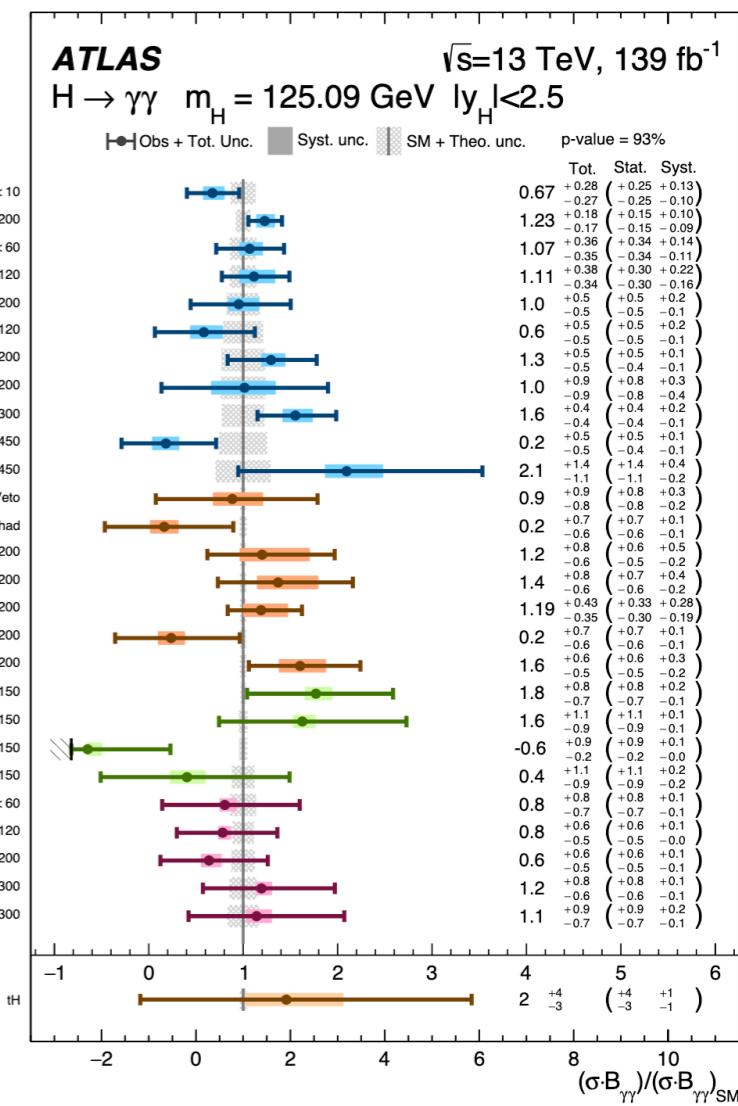
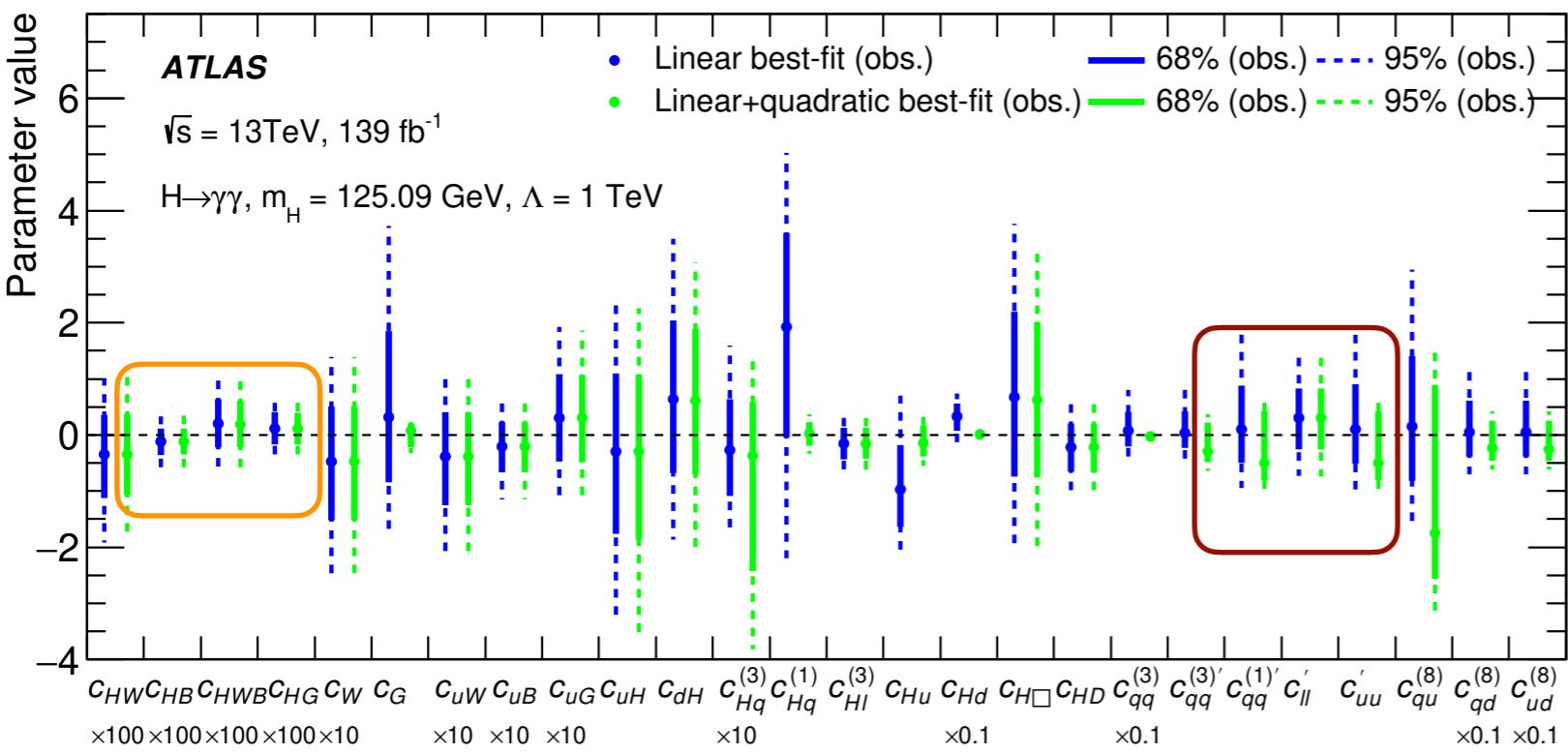
# $H \rightarrow \gamma\gamma$ : individual WCs constraints

Decay channel	Target Production Modes	$\mathcal{L} [\text{fb}^{-1}]$
$H \rightarrow \gamma\gamma$	ggF, VBF, WH, ZH, $t\bar{t}H$ , $tH$	139

CERN-EP-2022-094

- Constraints on individual WCs (one at the time).
- All values compatible with the SM within measurement uncertainties.
- When the sensitivity on one parameter is driven by inclusive event yields, linear and linear+quadratic parameterisations provide **similar results**.
- Operators with sensitivity to high- $p_T^H$  bins show **smaller confidence intervals** for the linear+quadratic case.

## SMEFTsim: " $U(3)^5$ flavour symmetry"



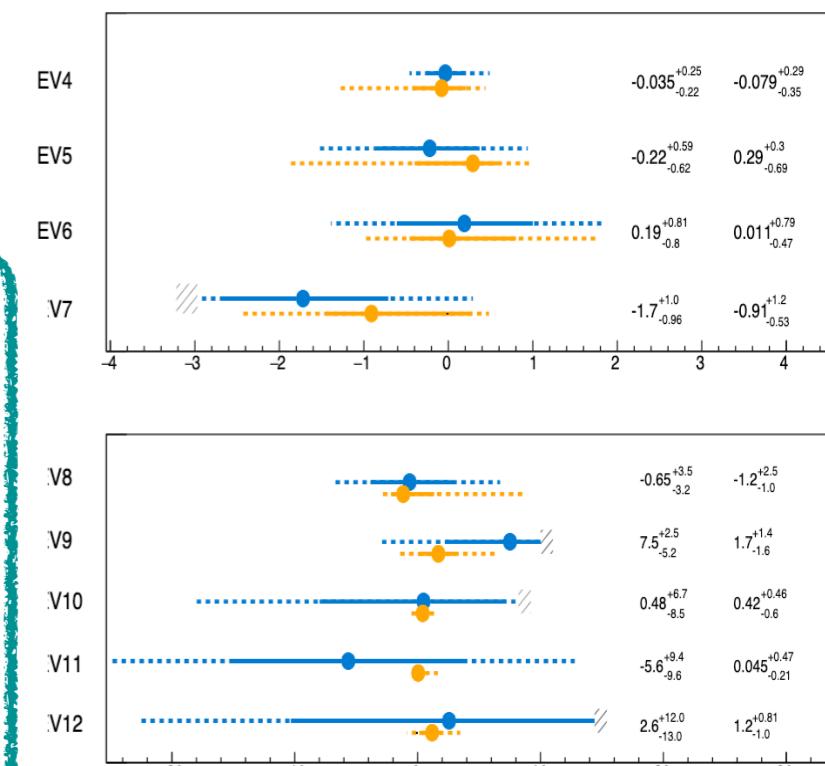
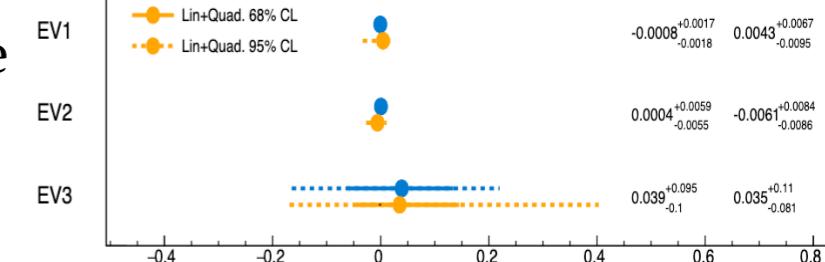
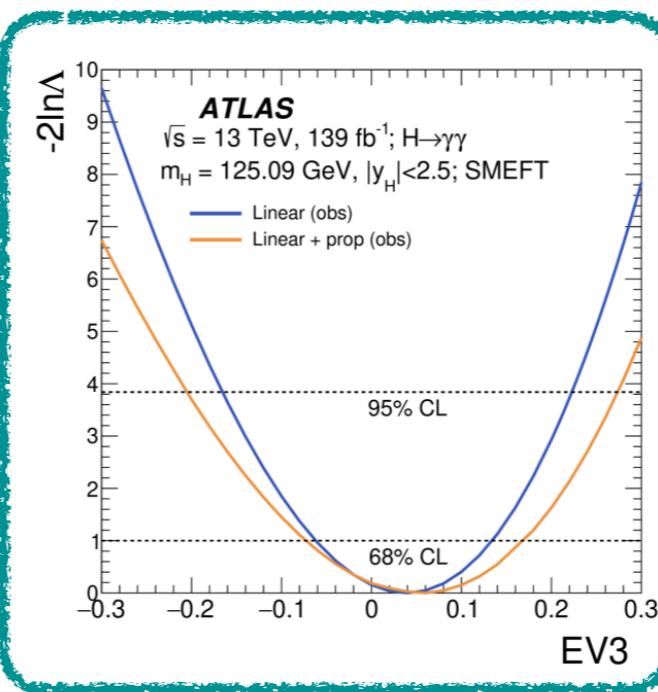
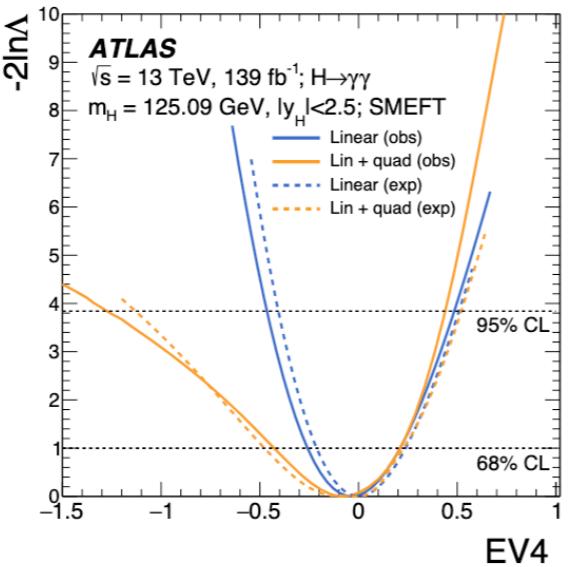
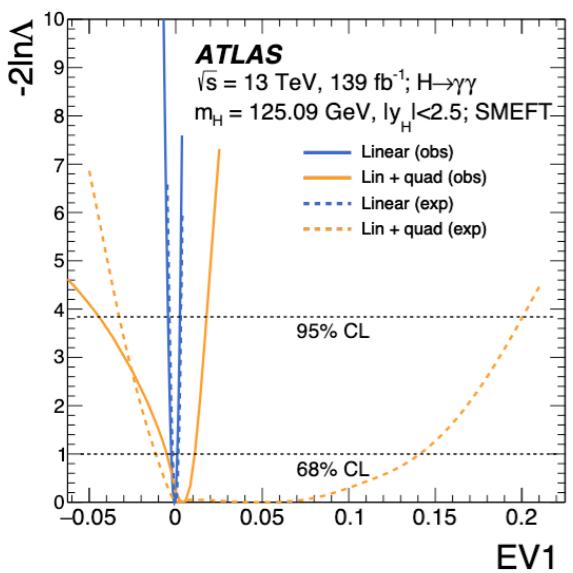
# Higgs

# $H \rightarrow \gamma\gamma$ : constraints on sensitive directions

Decay channel	Target Production Modes	$\mathcal{L} [\text{fb}^{-1}]$
$H \rightarrow \gamma\gamma$	ggF, VBF, WH, ZH, $t\bar{t}H$ , $tH$	139

CERN-EP-2022-094

- Cannot constraints simultaneously all the WCs the STXS measurements are affected by (34).
- Operators grouped according to the eigenvectors of the information matrix of the measurement.
- Inclusion of **propagator corrections** has been tested.
- Both linear and linear+quadratic results are provided.
- All results are in agreement with SM expectations.



10 directions from Higgs  
2020 combination  
(3 channels)

# First ATLAS Global combination

Decay channel	Target Production Modes	$\mathcal{L} [\text{fb}^{-1}]$
$H \rightarrow \gamma\gamma$	ggF, VBF, WH, ZH, $t\bar{t}H$ , $tH$	139
$H \rightarrow ZZ^*$	ggF, VBF, WH, ZH, $t\bar{t}H(4\ell)$	139
$H \rightarrow WW^*$	ggF, VBF	139
$H \rightarrow \tau\tau$	ggF, VBF, WH, ZH, $t\bar{t}H(\tau_{\text{had}}\tau_{\text{had}})$	139
	WH, ZH	139
$H \rightarrow b\bar{b}$	VBF	126
	$t\bar{t}H$	139

- **ATLAS Higgs boson data (2021 combination)**
  - **Higgs boson production and decay combined measurements in STXS bins**
- Higgs Combination

Process	Important phase space requirements	Observable	$\mathcal{L} [\text{fb}^{-1}]$
$pp \rightarrow e^\pm \nu \mu^\mp \nu$	$m_{\ell\ell} > 55 \text{ GeV}$ , $p_T^{\text{jet}} < 35 \text{ GeV}$	$p_T^{\text{lead. lep.}}$	36
$pp \rightarrow \ell^\pm \nu \ell^\pm \ell^-$	$m_{\ell\ell} \in (81, 101) \text{ GeV}$	$m_T^{WZ}$	36
$pp \rightarrow \ell^+ \ell^- \ell^+ \ell^-$	$m_{4\ell} > 180 \text{ GeV}$	$m_{Z2}$	139
$pp \rightarrow \ell^+ \ell^- jj$	$m_{jj} > 1000 \text{ GeV}$ , $m_{\ell\ell} \in (81, 101) \text{ GeV}$	$\Delta\phi_{jj}$	139

WW,WZ,4l, Z+2jets combination

- **ATLAS electroweak data**
- **Differential cross-section measurements for diboson and Z production via VBF**

Observable	Measurement	Prediction	Ratio
$\Gamma_Z$ [MeV]	$2495.2 \pm 2.3$	$2495.7 \pm 1$	$0.9998 \pm 0.0010$
$R_\ell^0$	$20.767 \pm 0.025$	$20.758 \pm 0.008$	$1.0004 \pm 0.0013$
$R_c^0$	$0.1721 \pm 0.0030$	$0.17223 \pm 0.00003$	$0.999 \pm 0.017$
$R_b^0$	$0.21629 \pm 0.00066$	$0.21586 \pm 0.00003$	$1.0020 \pm 0.0031$
$A_{\ell,\ell}^{0,\ell}$	$0.0171 \pm 0.0010$	$0.01718 \pm 0.00037$	$0.995 \pm 0.062$
$A_{c,c}^{0,c}$	$0.0707 \pm 0.0035$	$0.0758 \pm 0.0012$	$0.932 \pm 0.048$
$A_{b,b}^{0,b}$	$0.0992 \pm 0.0016$	$0.1062 \pm 0.0016$	$0.935 \pm 0.021$
$\sigma_{\text{had}}^0$ [pb]	$41488 \pm 6$	$41489 \pm 5$	$0.99998 \pm 0.00019$

Precision Electroweak Measurements on the Z Resonance

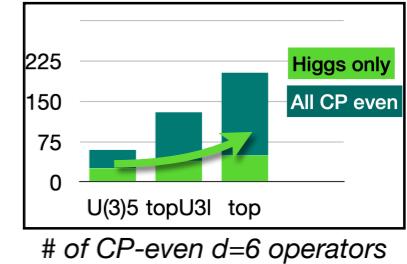
- **Electroweak precision observables measured at LEP and SLC**
- **Eight pseudo observables describing the physics at the Z-pole are interpreted.**
- **Measurement probed with high sensitivity O(1 - 0.01 %)**

# First ATLAS Global combination

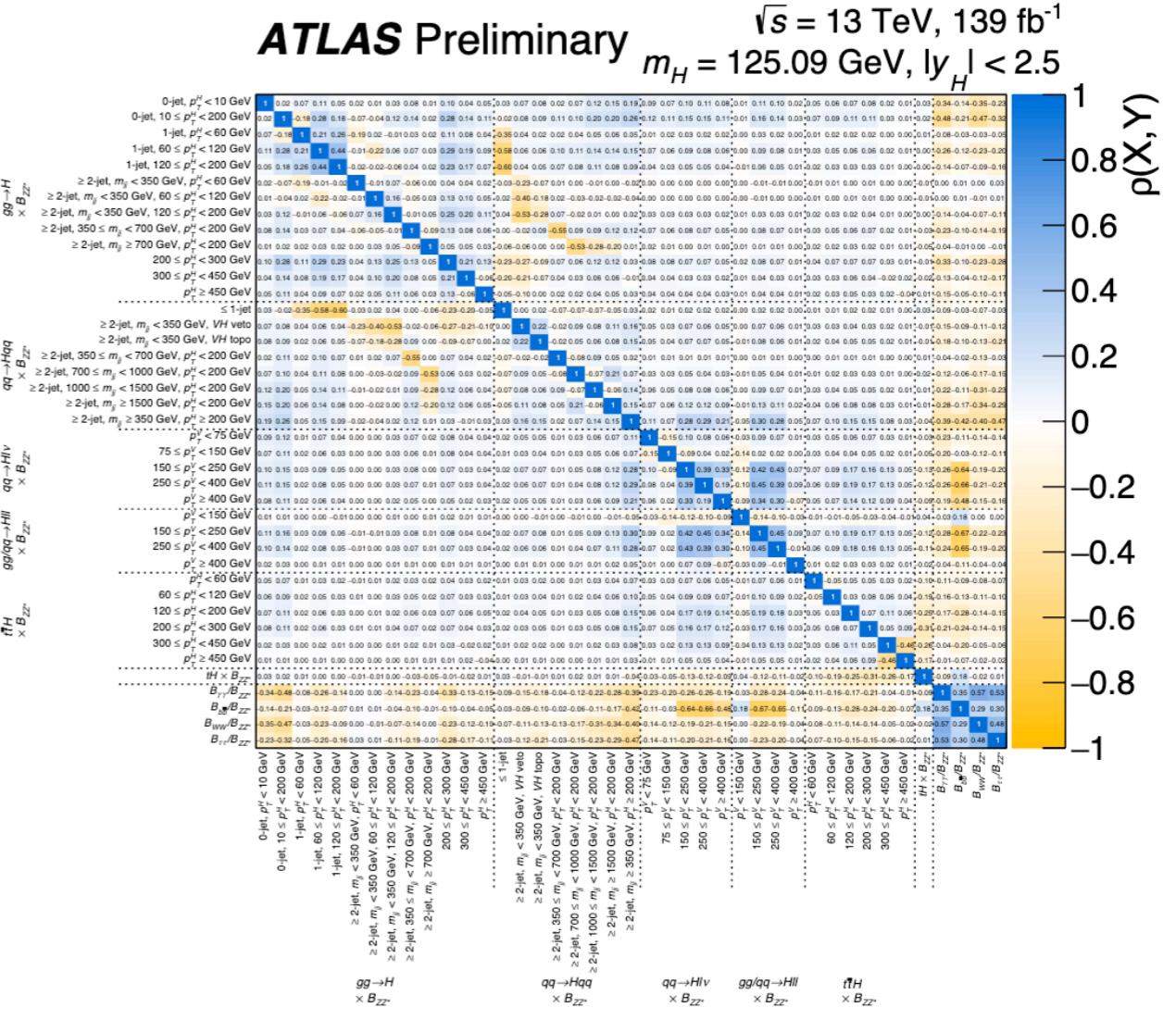
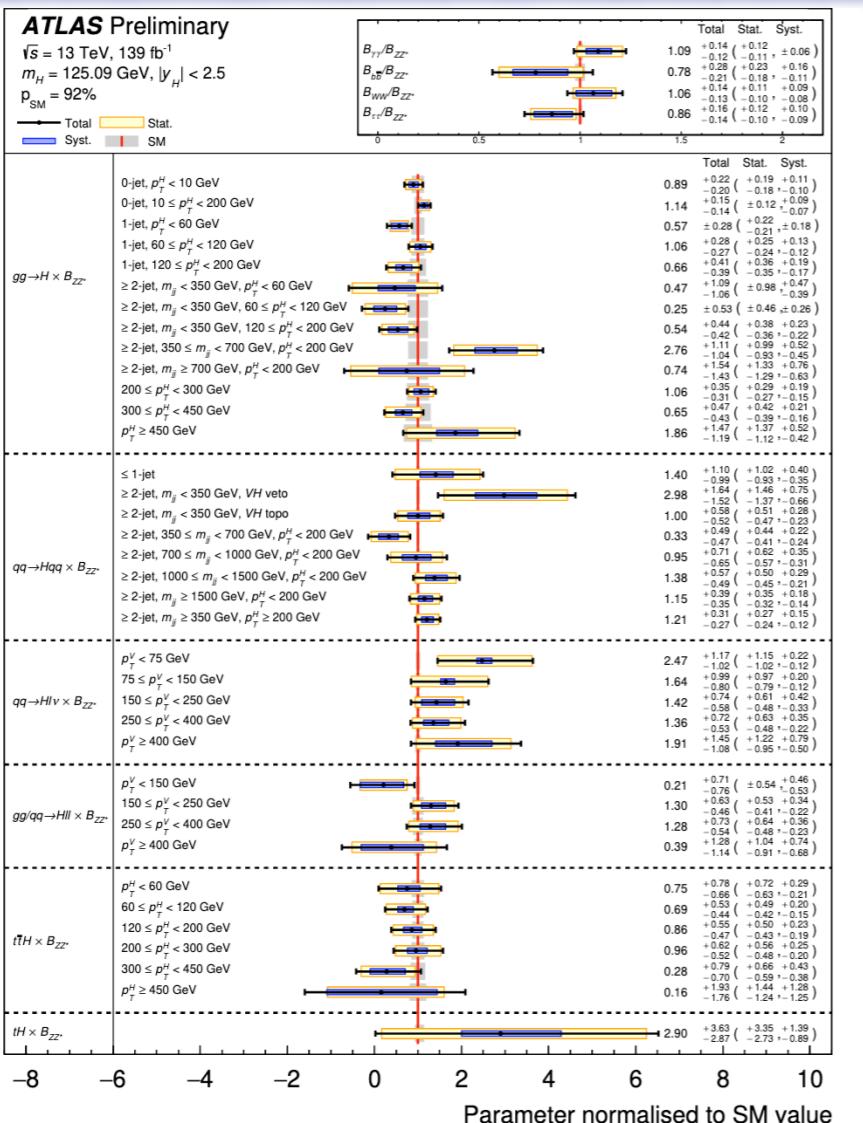
## Higgs Combination

Decay channel	Target Production Modes	$\mathcal{L} [\text{fb}^{-1}]$
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- ATLAS Higgs boson data (2021 combination)
- Higgs boson production and decay combined measurements in STXS bins



SMEFTsim: “topU31” flavour symmetry”



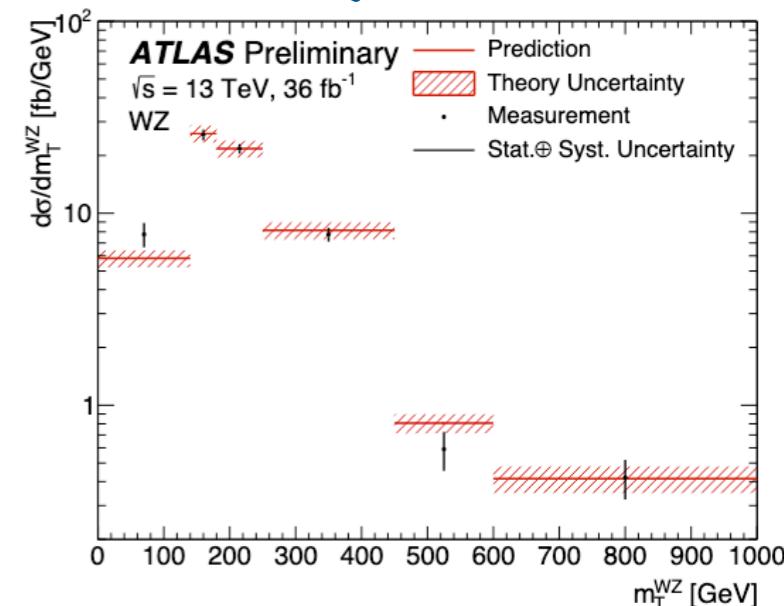
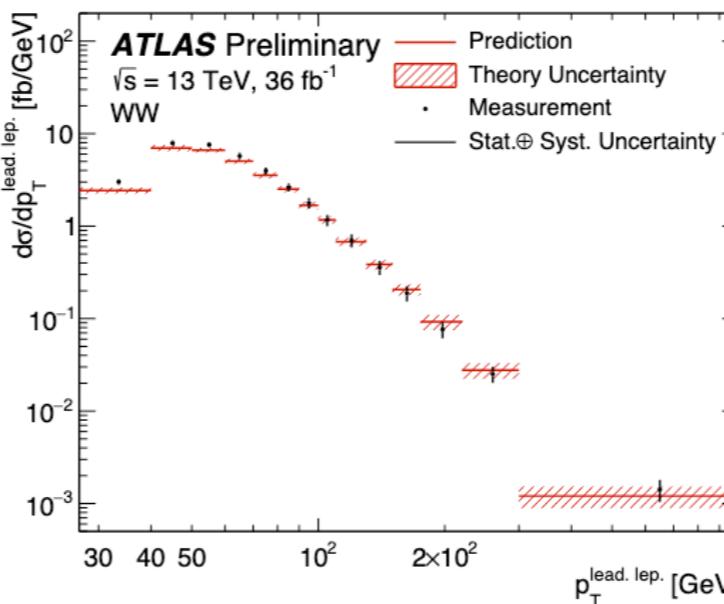
# First ATLAS Global combination

WW,WZ,4l, Z+2jets combination

$$L(x|c, \theta) = \frac{1}{\sqrt{(2\pi)^n \det(V)}} \exp\left(-\frac{1}{2} \Delta x^\top(c, \theta) V^{-1} \Delta x(c, \theta)\right)$$

$$\times \prod_i^{n_{\text{theo syst}}} f_i(\theta_{\text{theo syst},i}) \times \prod_i^{n_{\text{exp syst}}} f_i(\theta_{\text{exp syst},i}).$$

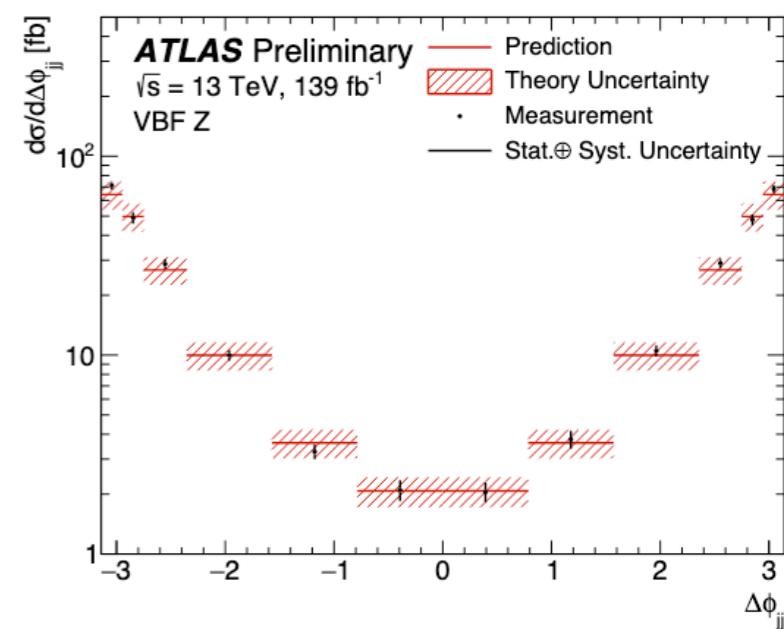
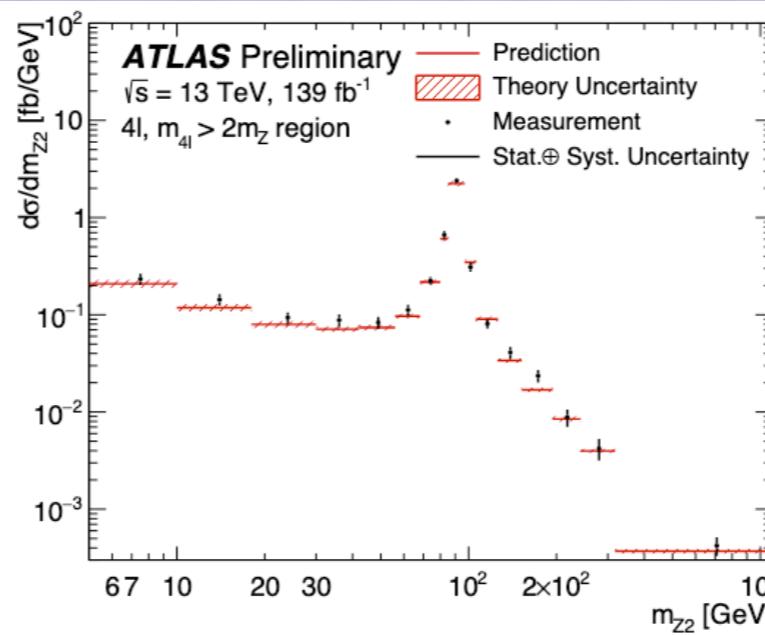
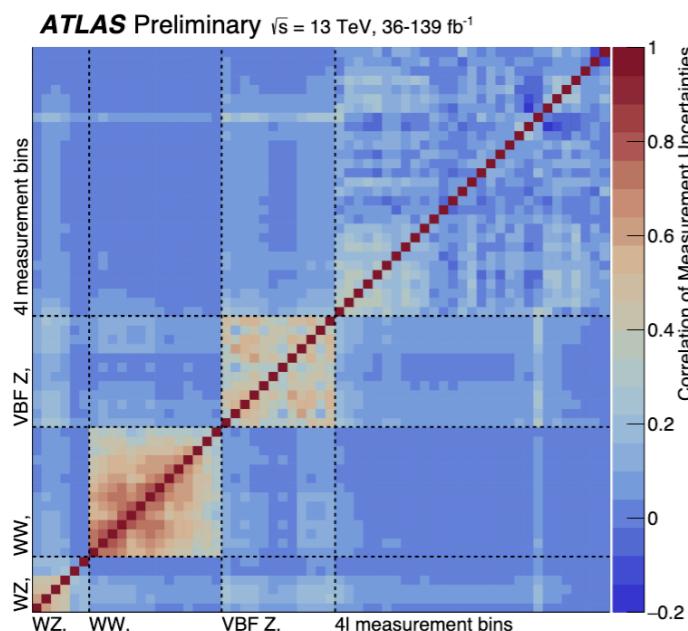
Multivariate gaussian



Process	Important phase space requirements	Observable	$\mathcal{L}$ [fb $^{-1}$ ]
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$pp \rightarrow \ell^\pm \nu \ell^\pm \ell^-$	$m_{\ell\ell} \in (81, 101)$ GeV	$m_T^{WZ}$	36
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- ATLAS electroweak data
- Differential cross-section measurements for diboson and Z production via VBF

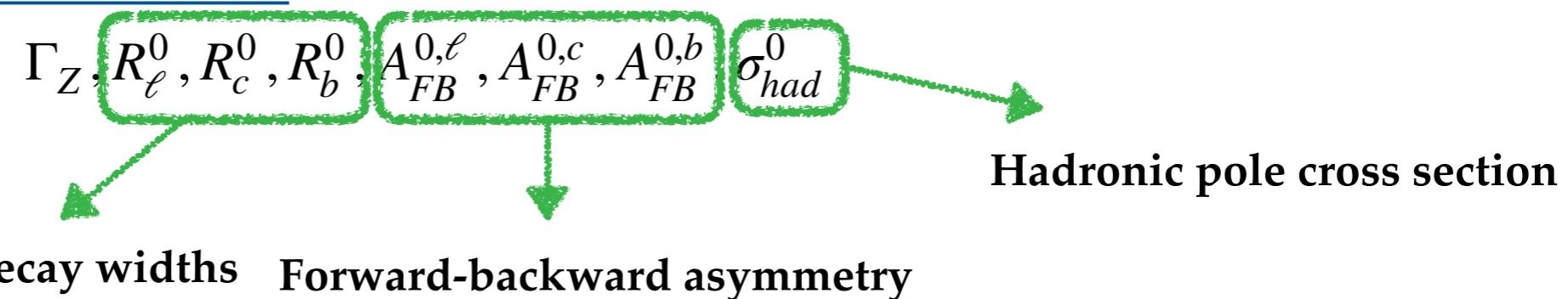
SMEFTsim: “topU3l” flavour symmetry”



# First ATLAS Global combination

## Precision Electroweak Measurements

### on the Z Resonance



- Tight limit provided by LEP-> only sensitive to a limited number of parameters.
- Parametrisation of EW pole observables only in the linear approximations:
  - Two different fit setups: Higgs+EW and Higgs+EW+EWPO
- The likelihood is modelled as a multivariate Gaussian, both theoretical and experimental uncertainties are included in the covariance matrix.

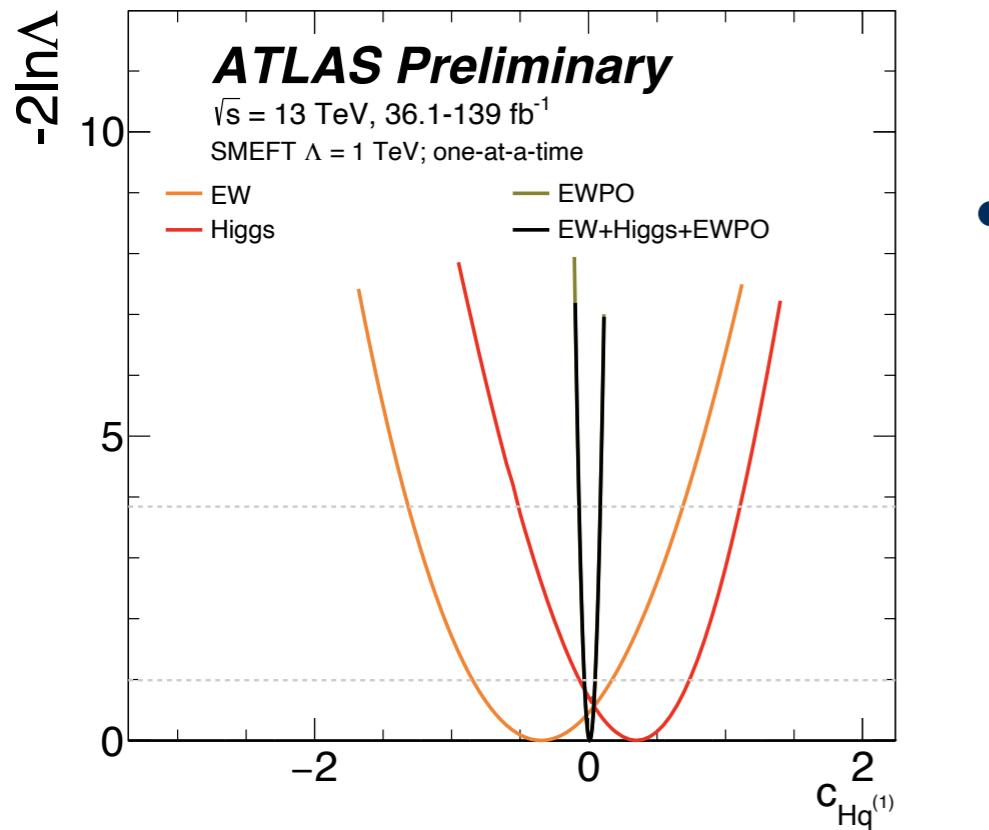
## EWPD in the SMEFT to dimension eight

Observable	Measurement	Prediction	Ratio
$\Gamma_Z$ [MeV]	$2495.2 \pm 2.3$	$2495.7 \pm 1$	$0.9998 \pm 0.0010$
$R_\ell^0$	$20.767 \pm 0.025$	$20.758 \pm 0.008$	$1.0004 \pm 0.0013$
$R_c^0$	$0.1721 \pm 0.0030$	$0.17223 \pm 0.00003$	$0.999 \pm 0.017$
$R_b^0$	$0.21629 \pm 0.00066$	$0.21586 \pm 0.00003$	$1.0020 \pm 0.0031$
$A_{FB}^{0,\ell}$	$0.0171 \pm 0.0010$	$0.01718 \pm 0.00037$	$0.995 \pm 0.062$
$A_{FB}^{0,c}$	$0.0707 \pm 0.0035$	$0.0758 \pm 0.0012$	$0.932 \pm 0.048$
$A_{FB}^{0,b}$	$0.0992 \pm 0.0016$	$0.1062 \pm 0.0016$	$0.935 \pm 0.021$
$\sigma_{had}^0$ [pb]	$41488 \pm 6$	$41489 \pm 5$	$0.99998 \pm 0.00019$

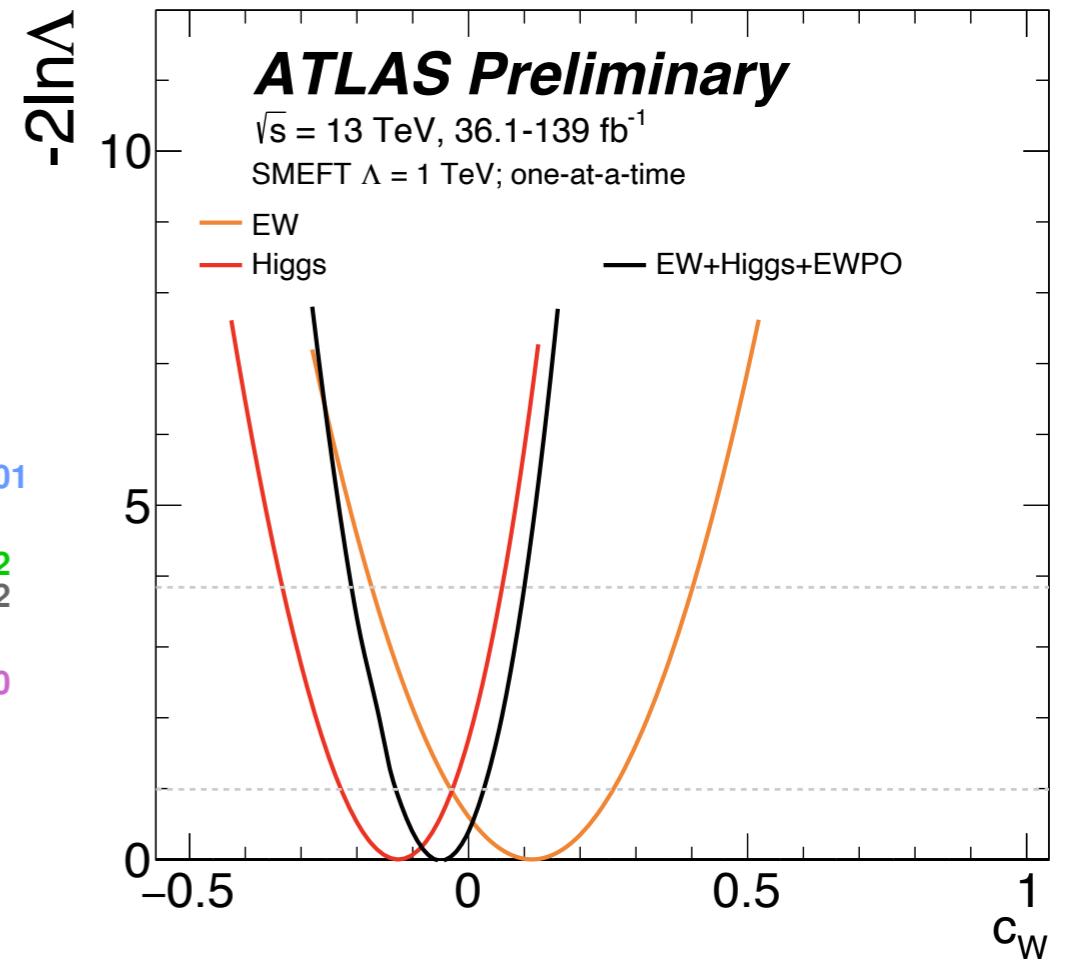
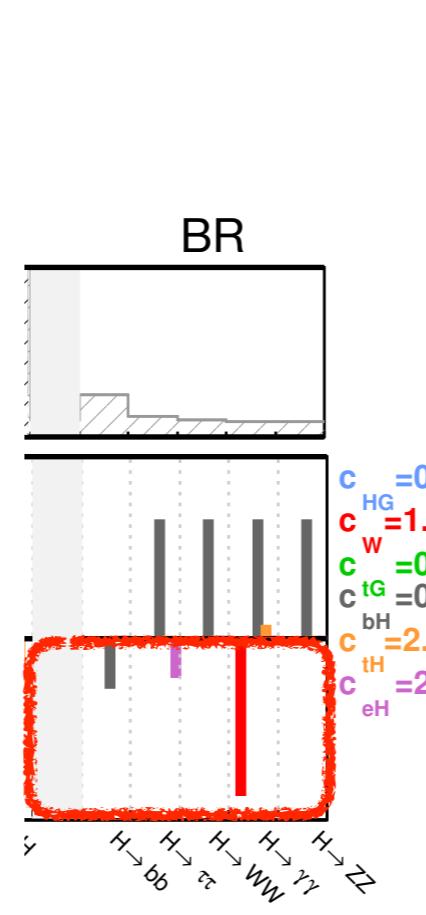
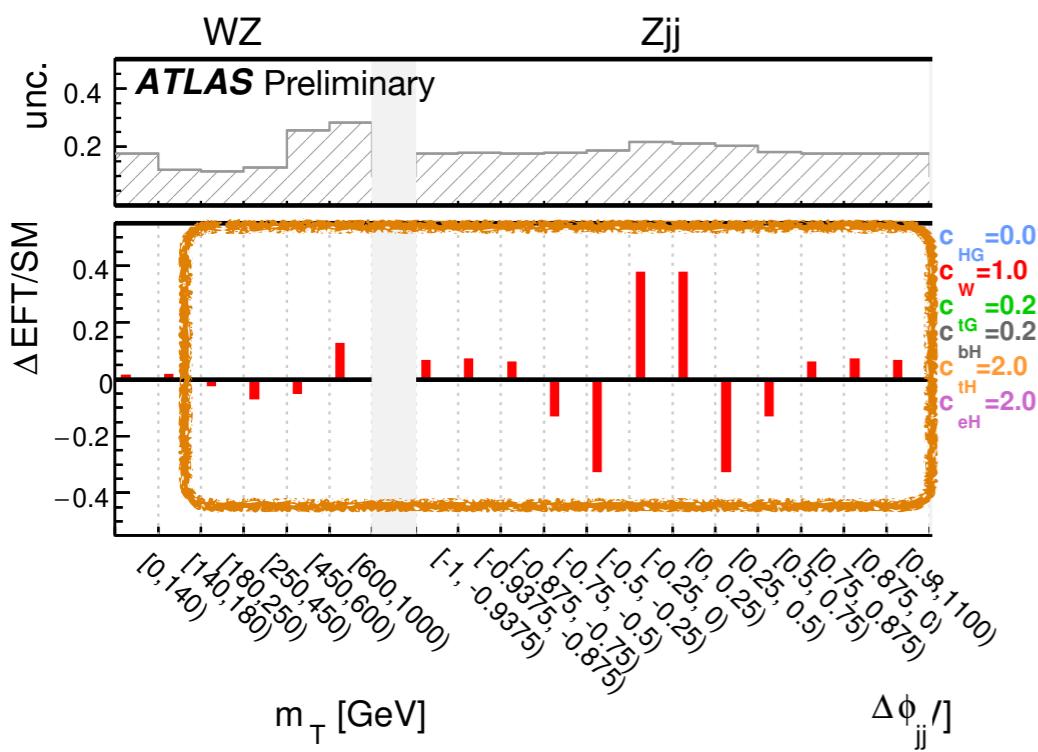
- Electroweak precision observables measured at LEP and SLC
- Eight pseudo observables describing the physics at the Z-pole are interpreted.
- Measurement probed with high sensitivity O(1 - 0.01 %)

# ATLAS Global combination: one at a time

[ATL-PHYS-PUB-2022-037](#)



- One parameter at a time scans to compare sensitivity to an operator across the 3 measurement groups;
  - all remaining Wilson coefficients fixed to zero;
  - correlations between operators are neglected.



# ATLAS Global combination

## HIGGS+EW

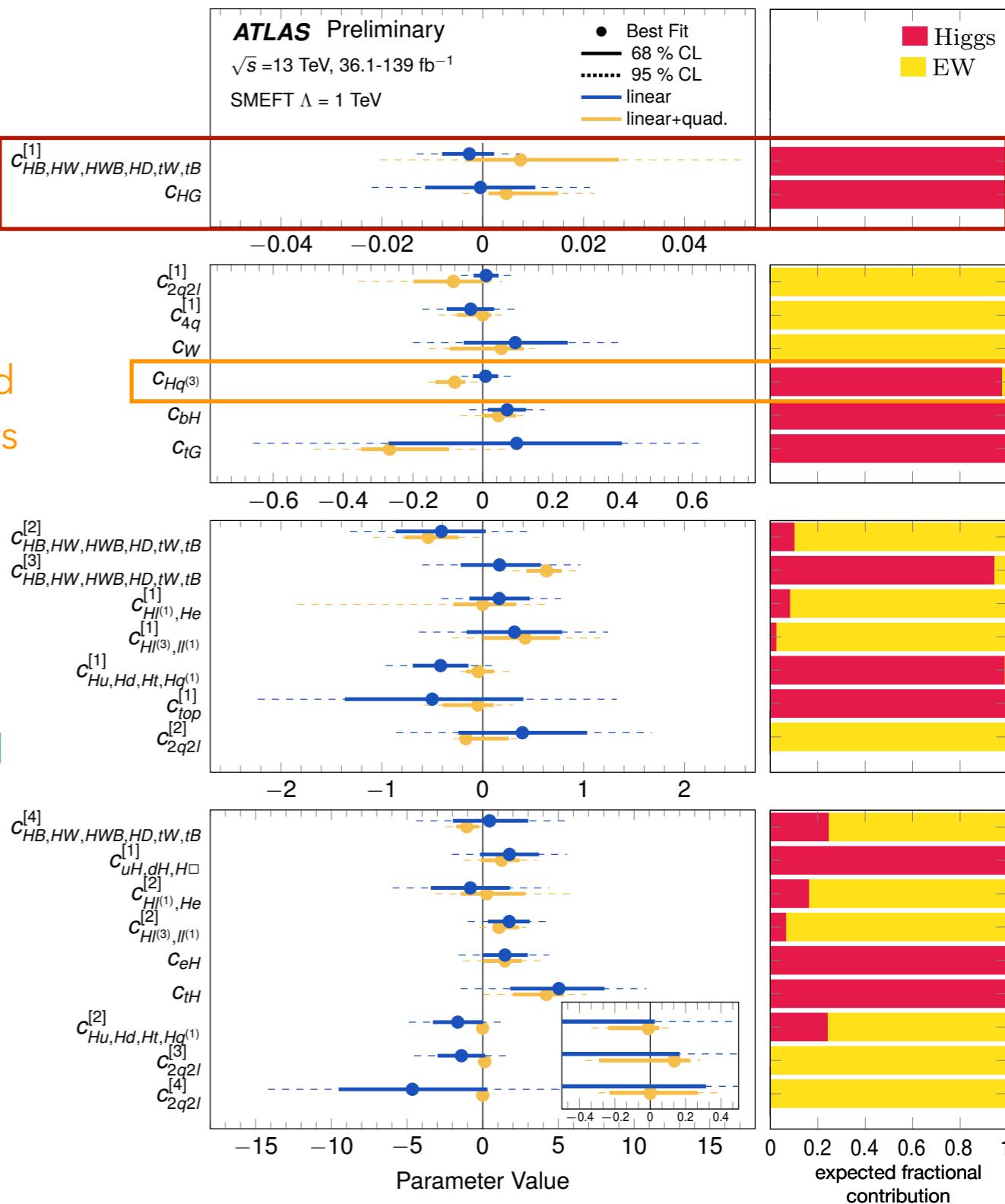
- Constraining 7 individual and 17 linear combinations of Wilson coefficients.
- Data overlap across datasets checked -> remove from the combination whenever relevant.
- PCA to identify sensitive directions-> a modified basis of linear combinations of WCs is defined.
- The fit uses sensitivity eigenvectors instead of original WCs.

Most stringent constraints

Constrained by both diboson and VH measurements

Weakly constrained fit directions-> quadratic contributions are often large; validity of the obtained constraints is questionable

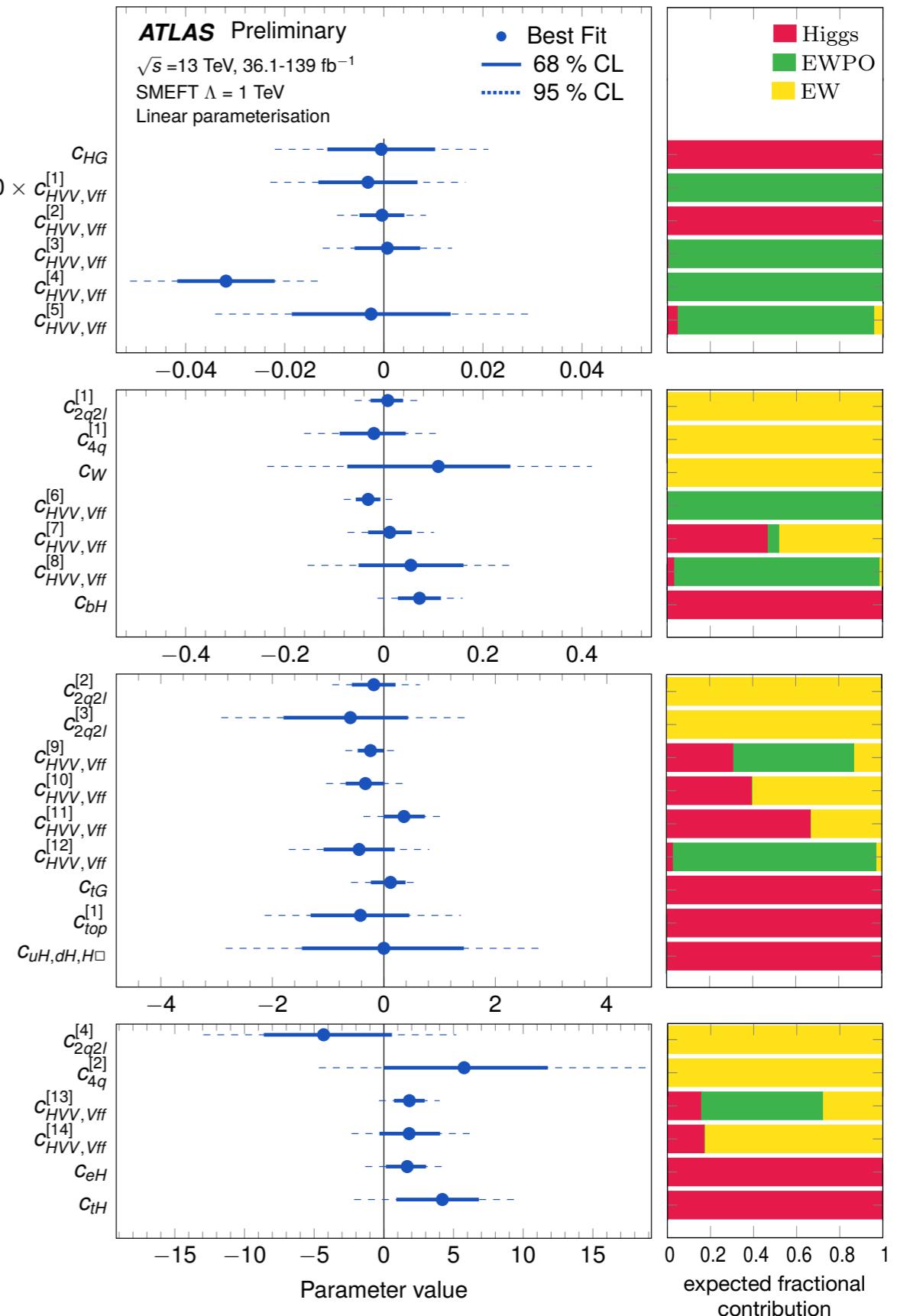
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# ATLAS Global combination

## HIGGS+EW+EWPO

- Constraining 6 individual and 22 linear combinations of Wilson coefficients.
- Several constraints driven by both ATLAS and LEP/SLD.
- Complementary information.
- Linear fits agree with the SM expectation for most fitted parameters, except for:
  - $c_{HVV,Vff}^{[4]}$  → excess driven by a well-known discrepancy in  $A_{FB}^{0,b}$  from the SM expectation.

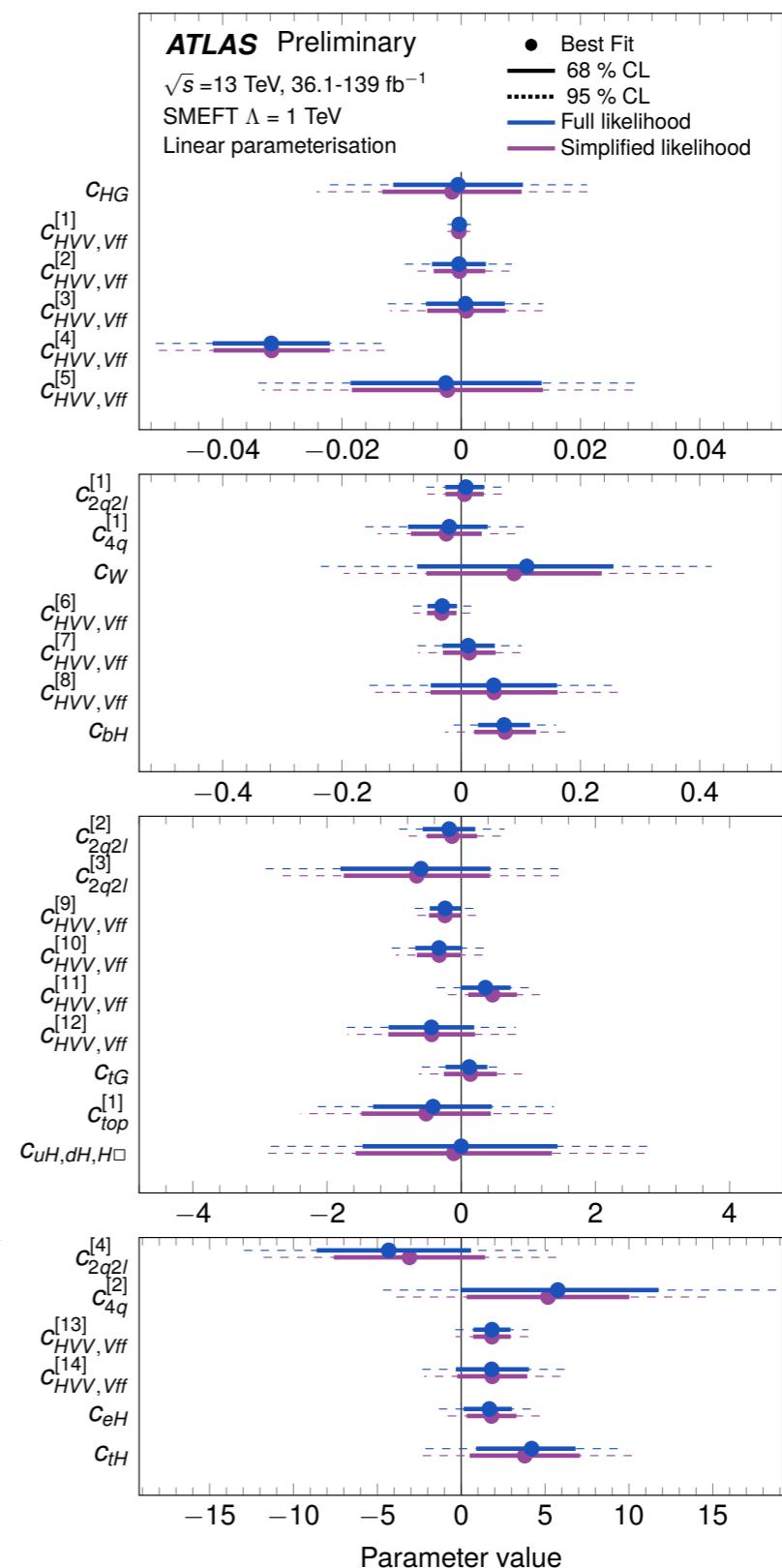


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# ATLAS Global combination: simplified likelihood

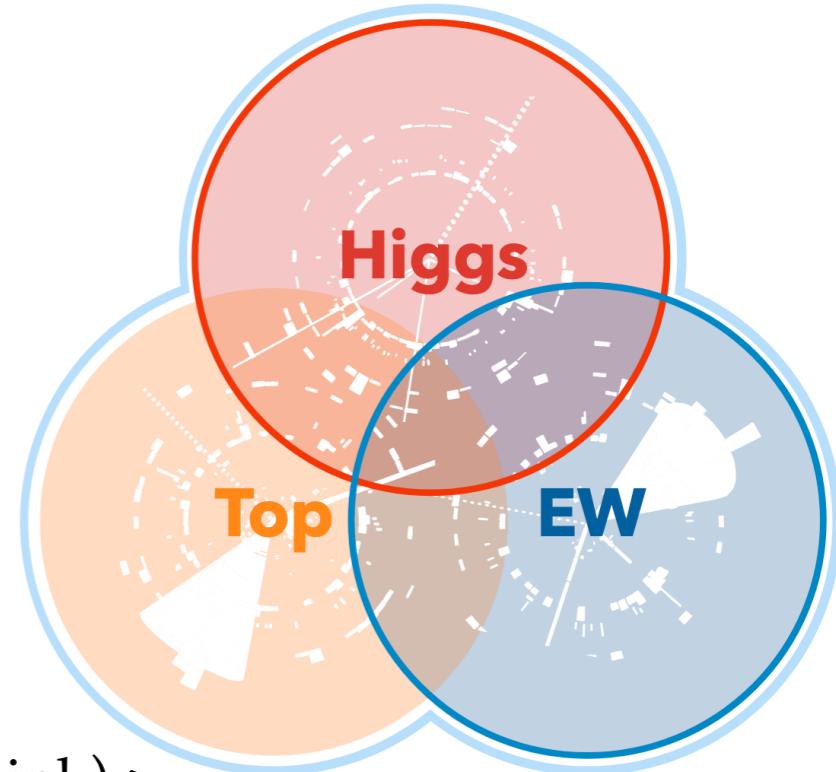
ATL-PHYS-PUB-2022-037

- Simplified likelihood model:
  - format to deliver results for re-interpretation;
  - signal strength modifier + correlation matrix.
- Results from the full likelihood fit compared to those using a simplified likelihood following a multi-variate Gaussian approach:
  - minimal differences between the two methods;
  - the simplified model is nuisance parameter free, as the effect of all uncertainties is encoded in the covariance matrix-> computationally inexpensive.
- Signal strength modifiers + correlation matrix available, preparing shared parameterisation.



# What's next

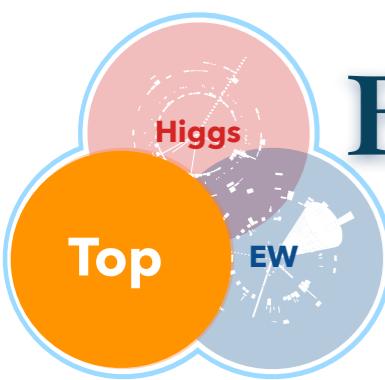
- First global ATLAS EFT interpretation available, also providing a simplified likelihood model for re-interpretation.
  - Well established framework used to perform the ATLAS Global combination.
  - Combination with additional Top + EW analyses ongoing-> provide complementary sensitivity.
- Ongoing Higgs interpretation of STXS results from Nature Paper ([Link](#))-> improvement expected coming from the inclusion of new inputs such as:
  - boosted-Hbb,  $H \rightarrow Z\gamma$ ,  $H \rightarrow cc\bar{c}$ ,  $H \rightarrow \mu\mu$ .
  - "Top" flavour scheme to fully exploit all the channels included in the combination.
- Ongoing studies on the inclusion of dimension 8 terms and mapping EFT and UV theories.
- ATLAS + CMS: ongoing exercise to include few channels and test the combination using consistent parameterisations and assumptions-> developing tools for EFT parameterisation (further details in [Fabian's talk](#))



*Stay tuned for many interesting results to come!!*

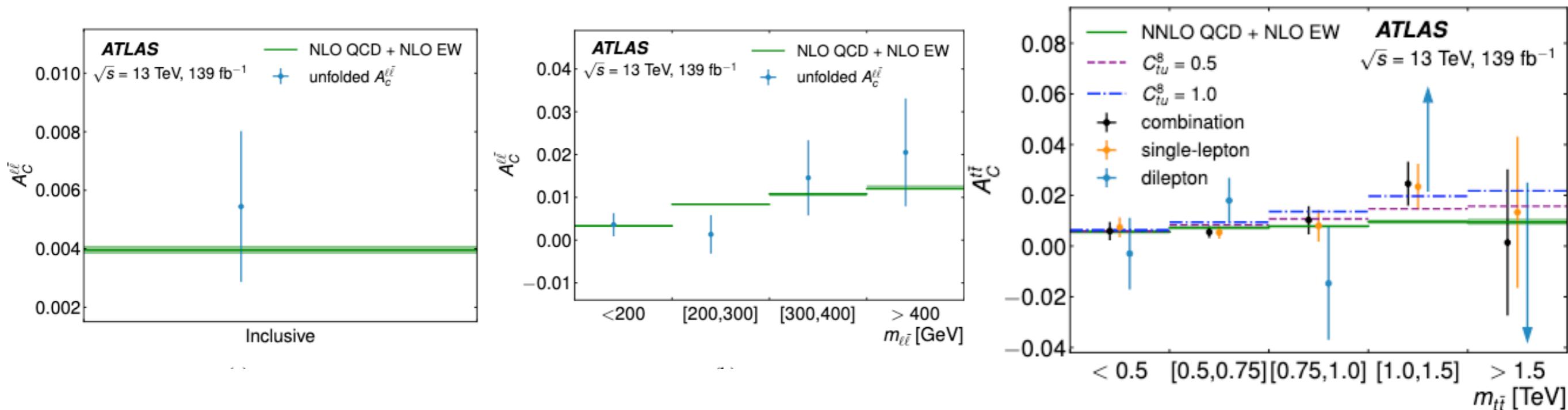
Thanks for your  
attention!

# Evidence for the charge asymmetry in $pp \rightarrow t\bar{t}$



[arxiv:2208.12095](https://arxiv.org/abs/2208.12095)

- Inclusive and differential full Run2 measurements of the top–antitop ( $t\bar{t}$ ) charge asymmetry  $A_c^{t\bar{t}}$  and the leptonic asymmetry  $A_c^{\ell\bar{\ell}}$
- Differential measurements are performed as a function of the invariant mass, transverse momentum and longitudinal boost of the  $t\bar{t}$  system.
- Combined results are interpreted in the SMEFT framework.
- **14 four-fermion operators** + 1 operator for top–gluon interaction.



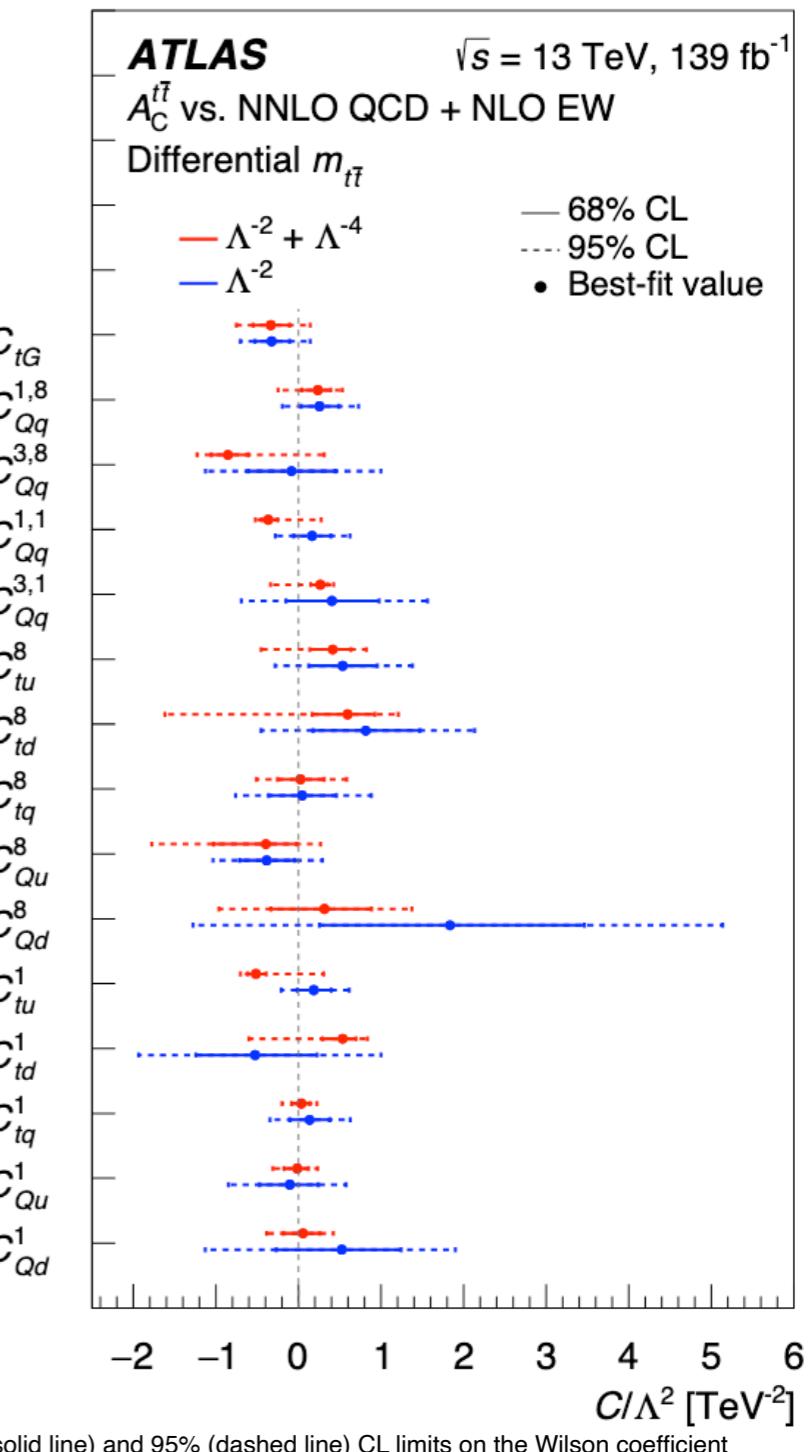
The combined inclusive charge asymmetry is measured to be  $A_c^{t\bar{t}} = 0.0068 \pm 0.0015$ , which differs from zero by **4.7 standard deviations**.

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- Combined results are interpreted in the SMEFT framework.
- **14 four-fermion operators** + 1 operator for top–gluon interaction.
  
- Large improvement comparing with LHC 8TeV / Tevatron results.
- Interplay between sensitivity, which increases rapidly at higher  $m_{t\bar{t}}$ , and uncertainty, which grows from 0.2%–0.3% in the lowest mass bin to 2.9% in the highest bin.
- For the linear fit, the tightest limit is obtained in the mass bin from 1 to 1.5 TeV.
- Constraint from the differential  $m_{t\bar{t}}$  measurement more than a factor 2 stronger than the one from inclusive measurement (increase in sensitivity with higher  $m_{t\bar{t}}$ ).

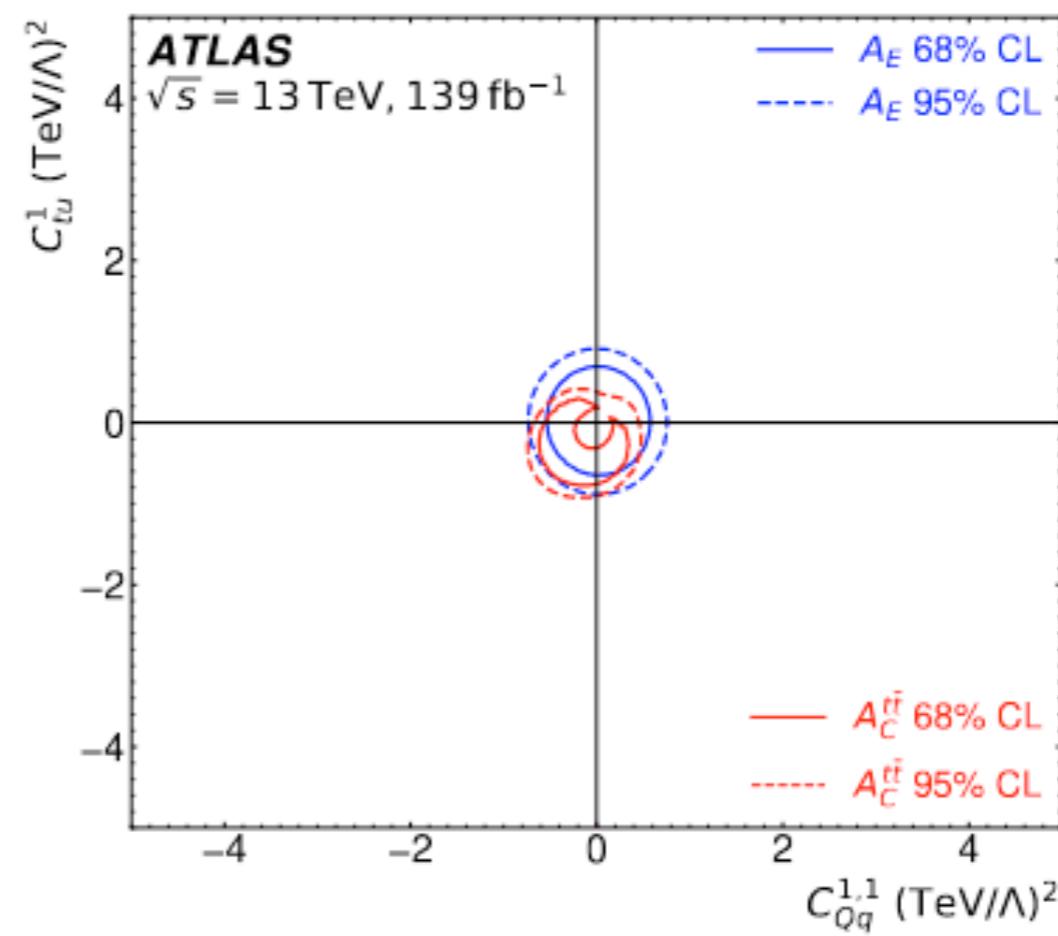
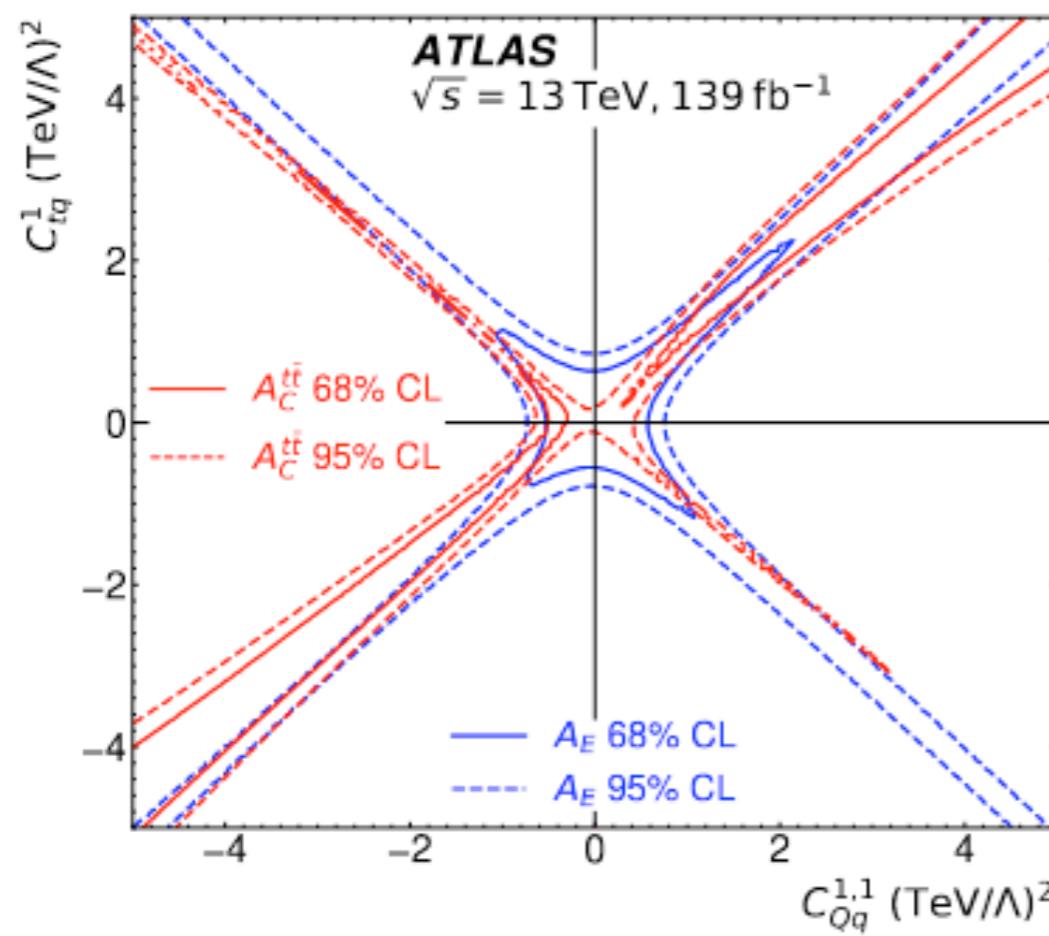
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# Evidence for the charge asymmetry in $pp \rightarrow t\bar{t}$

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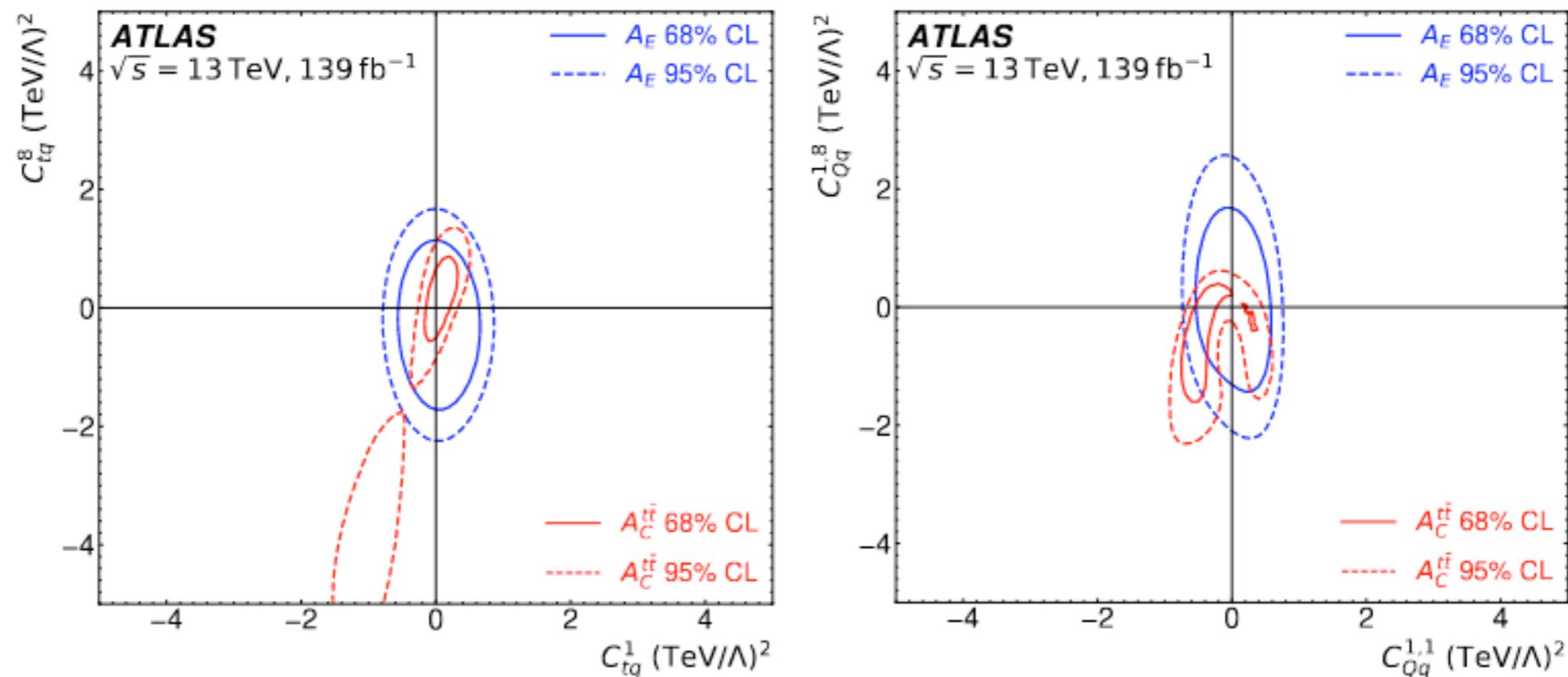
- Due to the extra jet in  $t\bar{t}j$  production, the QCD structure of the energy asymmetry is not the same as for the charge asymmetry in  $t\bar{t}$  production, so the two asymmetries probe different directions in chiral and colour space.
- For colour-singlet operators with different quark chiralities (top row), the two asymmetries probe similar areas in the parameter space.



# Evidence for the charge asymmetry in $pp \rightarrow t\bar{t}$

[arxiv:2208.12095](https://arxiv.org/abs/2208.12095)

- Due to the extra jet in  $t\bar{t}j$  production, the QCD structure of the energy asymmetry is not the same as for the charge asymmetry in  $t\bar{t}$  production, so the two asymmetries probe different directions in chiral and colour space.
- The bottom row shows. Here, the different shapes of the bounds are due to the different colour-singlet and colour-octet contributions to  $t\bar{t}$  and  $t\bar{t}j$  production, which is probed with high sensitivity by the asymmetries.

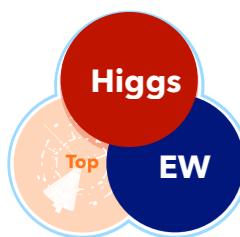


colour-singlet versus colour-octet operators with the same quark chiralities

# New SMEFTsim v3.0

	general		U35		MFV		top		topU31	
	all	<del>CP</del>	all	<del>CP</del>	all	<del>CP</del>	all	<del>CP</del>	all	<del>CP</del>
$\mathcal{L}_6^{(1)}$	4	2	4	2	2	-	4	2	4	2
$\mathcal{L}_6^{(2,3)}$	3	-	3	-	3	-	3	-	3	-
$\mathcal{L}_6^{(4)}$	8	4	8	4	4	-	8	4	8	4
$\mathcal{L}_6^{(5)}$	54	27	6	3	7	-	14	7	10	5
$\mathcal{L}_6^{(6)}$	144	72	16	8	20	-	36	18	28	14
$\mathcal{L}_6^{(7)}$	81	30	9	1	14	-	21	2	15	2
$\mathcal{L}_6^{(8a)}$	297	126	8	-	10	-	31	-	16	-
$\mathcal{L}_6^{(8b)}$	450	195	9	-	19	-	40	2	27	2
$\mathcal{L}_6^{(8c)}$	648	288	8	-	28	-	54	4	31	4
$\mathcal{L}_6^{(8d)}$	810	405	14	7	13	-	64	32	40	20
tot	2499	1149	85	25	120	-	275	71	182	53

From Ilaria's talk

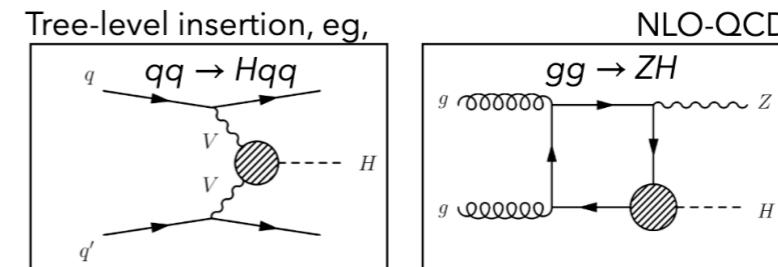


# SMEFT parameterisation



The impact of dim-6 CP-even operators is estimated using both MC truth and analytical predictions for all the Wilson coefficients that have numerically relevant contributions (62).

- Dimension-six operator effects are calculated:
  - at tree level using SMEFTsim 3.0.
  - for processes that are loop-induced in the SM, thus  $ggH$  and  $ggZH$  production, Higgs boson decays into gluons -> SMEFTatNLO.
  - Analytic formulas for  $H \rightarrow \gamma\gamma$  including NLO EW corrections and LEP observables.
- Theory uncertainties on SM predictions, no additional uncertainties on SMEFT.
- Acceptance corrections to account for kinematic differences between SM and SMEFT in Higgs boson decays on both **linear** and **linear+quadratic** terms.
- Effects of width changes of intermediate particles (“propagator corrections”) included.



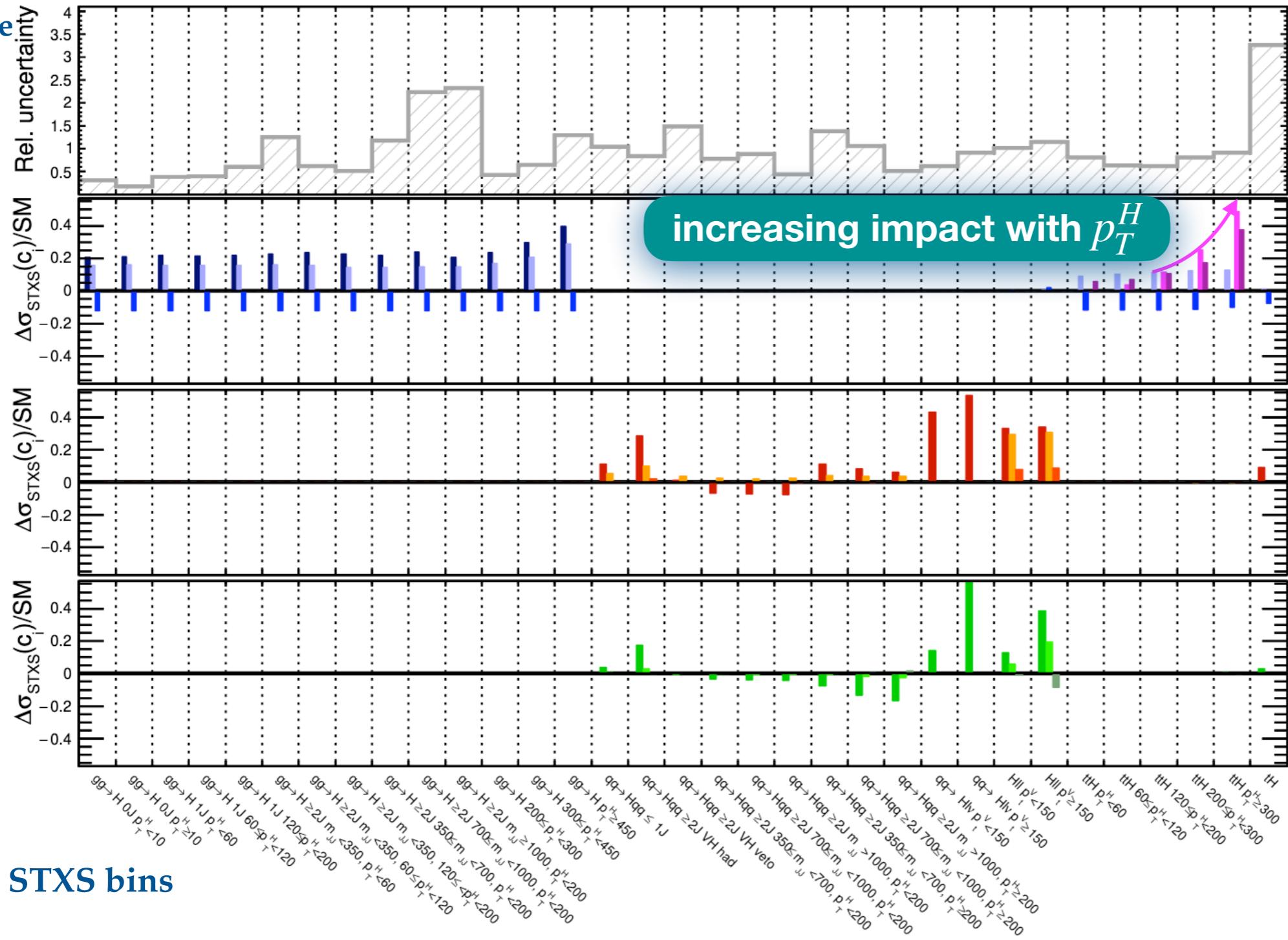
# SMEFT interpretation workflow

## Parameterisation

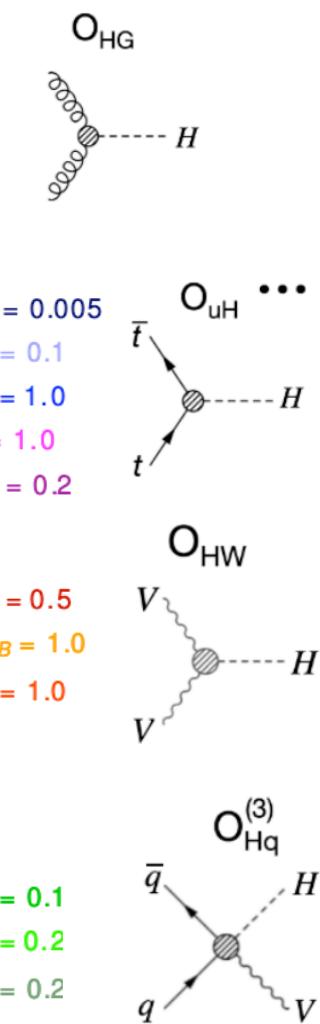
CERN-EP-2022-094

**ATLAS Simulation**  $\sqrt{s}=13 \text{ TeV } 139\text{fb}^{-1}$   $H \rightarrow \gamma\gamma, m_H = 125.09 \text{ GeV}, \Lambda = 1 \text{ TeV}$

Relative Impact w.r.t SM



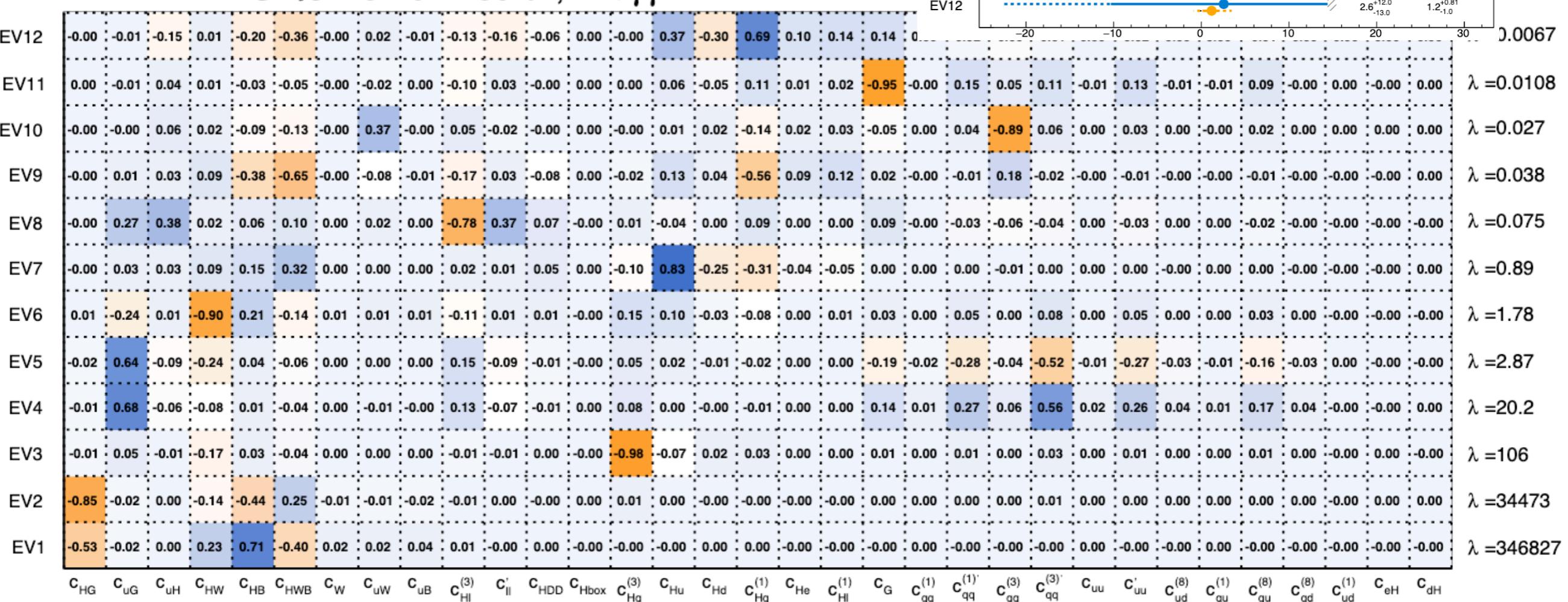
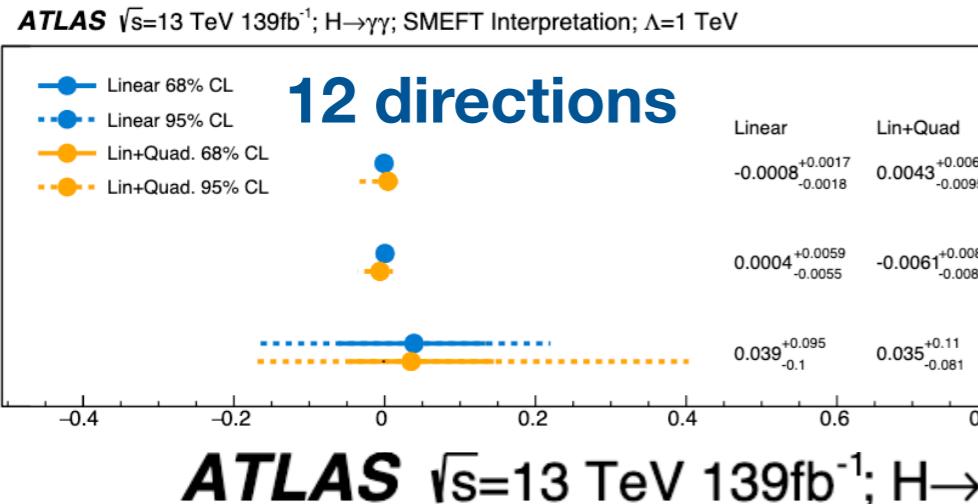
STXS bins



# $H \rightarrow \gamma\gamma$ : constraints on sensitive directions

CERN-EP-2022-094

- Operators grouped according to the eigenvectors of the information matrix of the measurement.
- All results are in agreement with SM expectations.



# Electroweak precision observables

- Electroweak precision observables (EWPO) measured e.g. at LEP and SLC, including 8 observables in our fit.
- Constraints obtained from the interpretation of these observables are typically more precise than LHC constraints but only a limited number of directions in parameter space can be constrained.
- The tight EWPO constraints provided on operators affecting weak-boson–fermion couplings allow to **disentangle** their effect from those affecting only Higgs or weak-boson self couplings, which cannot be constrained with  $Z$ -pole data.
- The precision observables agree very well with the SM expectation, with the exception of  $A_{FB}^{0,b}$  and  $A_{FB}^{0,c}$  for which deviations of more than three and more than one standard deviations, respectively, are found.
- Linear only parameterisation.

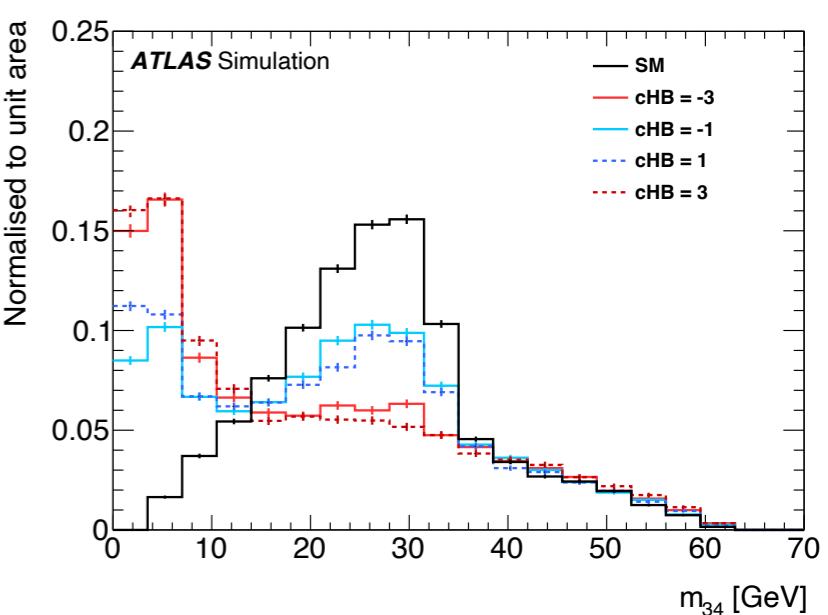
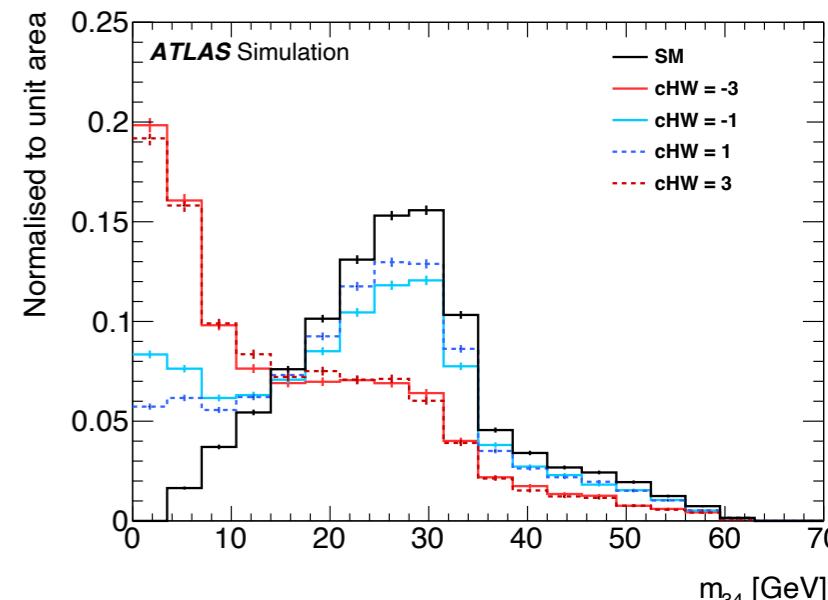
## Example of linear parameterisation for EWPO observables

$$\frac{\Gamma_Z}{\Gamma_{Z,\text{SM}}} = 1 + 0.059c_{ll}^{(1)} - 0.004c_{Hb} - 0.004c_{HD} - 0.007c_{Hd} - 0.011c_{Hl}^{(1)} - 0.081c_{Hl}^{(3)} - 0.011c_{He} \\ + 0.007c_{HQ}^{(1)} + 0.021c_{HQ}^{(1)} + 0.075c_{HQ}^{(3)} + 0.021c_{HQ}^{(3)} + 0.014c_{Hu} + 0.027c_{HWB}$$

Observable	Measurement	Prediction	Ratio
$\Gamma_Z$ [MeV]	$2495.2 \pm 2.3$	$2495.7 \pm 1$	$0.9998 \pm 0.0010$
$R_\ell^0$	$20.767 \pm 0.025$	$20.758 \pm 0.008$	$1.0004 \pm 0.0013$
$R_c^0$	$0.1721 \pm 0.003$	$0.17223 \pm 0.00003$	$0.999 \pm 0.017$
$R_b^0$	$0.21629 \pm 0.00066$	$0.21586 \pm 0.00003$	$1.0020 \pm 0.0031$
$A_{FB}^{0,\ell}$	$0.0171 \pm 0.0010$	$0.01718 \pm 0.00037$	$0.995 \pm 0.062$
$A_{FB}^{0,c}$	$0.0707 \pm 0.0035$	$0.07583 \pm 0.00117$	$0.932 \pm 0.048$
$A_{FB}^{0,b}$	$0.0992 \pm 0.0016$	$0.10615 \pm 0.00162$	$0.935 \pm 0.021$
$\sigma_{\text{had}}^0$ [pb]	$41488 \pm 6$	$41489 \pm 5$	$0.99998 \pm 0.00019$

# ATLAS Global combination: acceptance

- EFT parameterisation is affected by analysis level selections used to reconstruct SM Higgs.
- Acceptance effect can be considerable when the shape of EFT distribution is different from the SM distribution.
- Assuming SM-like acceptance not always possible-> modifications for  $H \rightarrow ZZ$  (ggH production - valid for all) and  $H \rightarrow WW$  are considered.
- Example: shape of the invariant mass distribution of the secondary lepton pair in  $H \rightarrow 4l$  decays for various values of the Wilson coefficients  $c_{HW}$  and  $c_{HB}$ .
- Acceptance parameterisation is a ratio of polynomials, linearised using Taylor expansion.



For one coupling, the acceptance can be approximated using a Taylor expansion:

$$\frac{A^{BSM}}{A^{SM}} = \frac{1 + a_{fid}c_i + b_{fid}c_i^2}{1 + a_{tot}c_i + b_{tot}c_i^2}$$

$$\approx 1 + (a_{fid} - a_{tot})c_i + (b_{fid} - b_{tot} - a_{fid}a_{tot} + a_{tot}^2)c_i^2 + \dots$$

# ATLAS Global combination: statistical combination

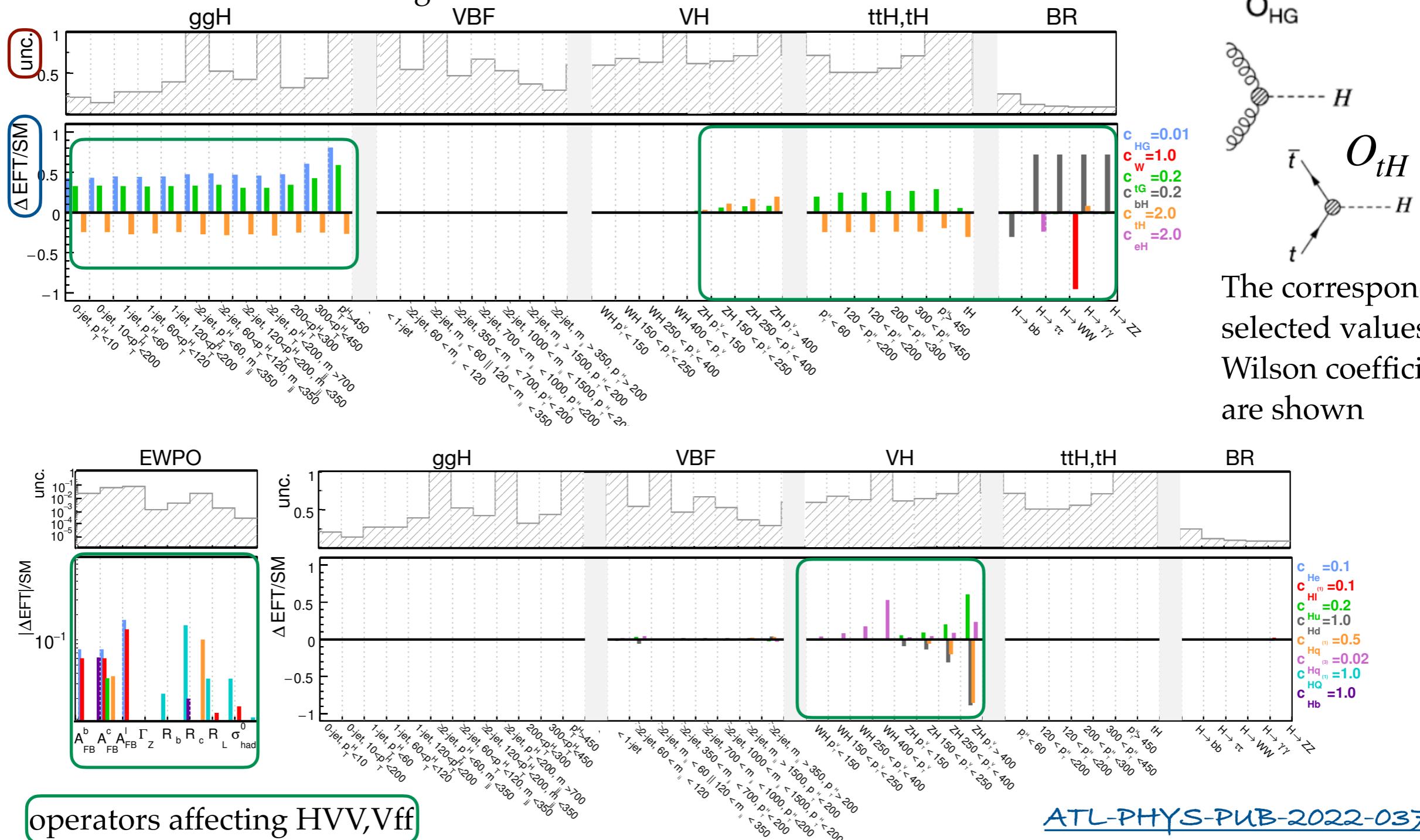
- Overlapping categories:
  - regions of the inclusive  $4\ell$  analysis that target  $m_{4\ell} < 180 \text{ GeV}$  are excluded (small impact on SMEFT).
  - the 0-jet  $WW$  control region from HWW is excluded;  $WW$  normalisation is correlated with  $WW$  signal normalisation.
- A multivariate Gaussian model is used for the interpretation of both LHC EW measurements and EWPO.
- Systematic uncertainties modelled with common nuisance parameter:
  - for unfolded SM measurements: experimental nuisance parameter shift unfolded results;
  - same nuisance parameter shift reco-level prediction for Higgs measurements.
- For EWPO the model contains no nuisance parameters and both theoretical and experimental uncertainties are included in the covariance matrix.
- Limits on WCs extracted using combined likelihood (product of individual likelihood).

Correlated Uncertainty Source	Parameters
Luminosity (correlated part)	1
Luminosity 2015/16	1
Luminosity 2017/18	1
Pile-up modelling	1
Pile-up jet suppression	1
Jet energy scale (pile-up modelling)	3
Jet energy scale $\eta$ -inter-calibration	1
Jet energy resolution	12
B-tagging efficiency ( $WW$ and $H \rightarrow WW^*$ )	1
$WW$ modelling ( $WW$ and $H \rightarrow WW^*$ )	2

# ATLAS Global combination

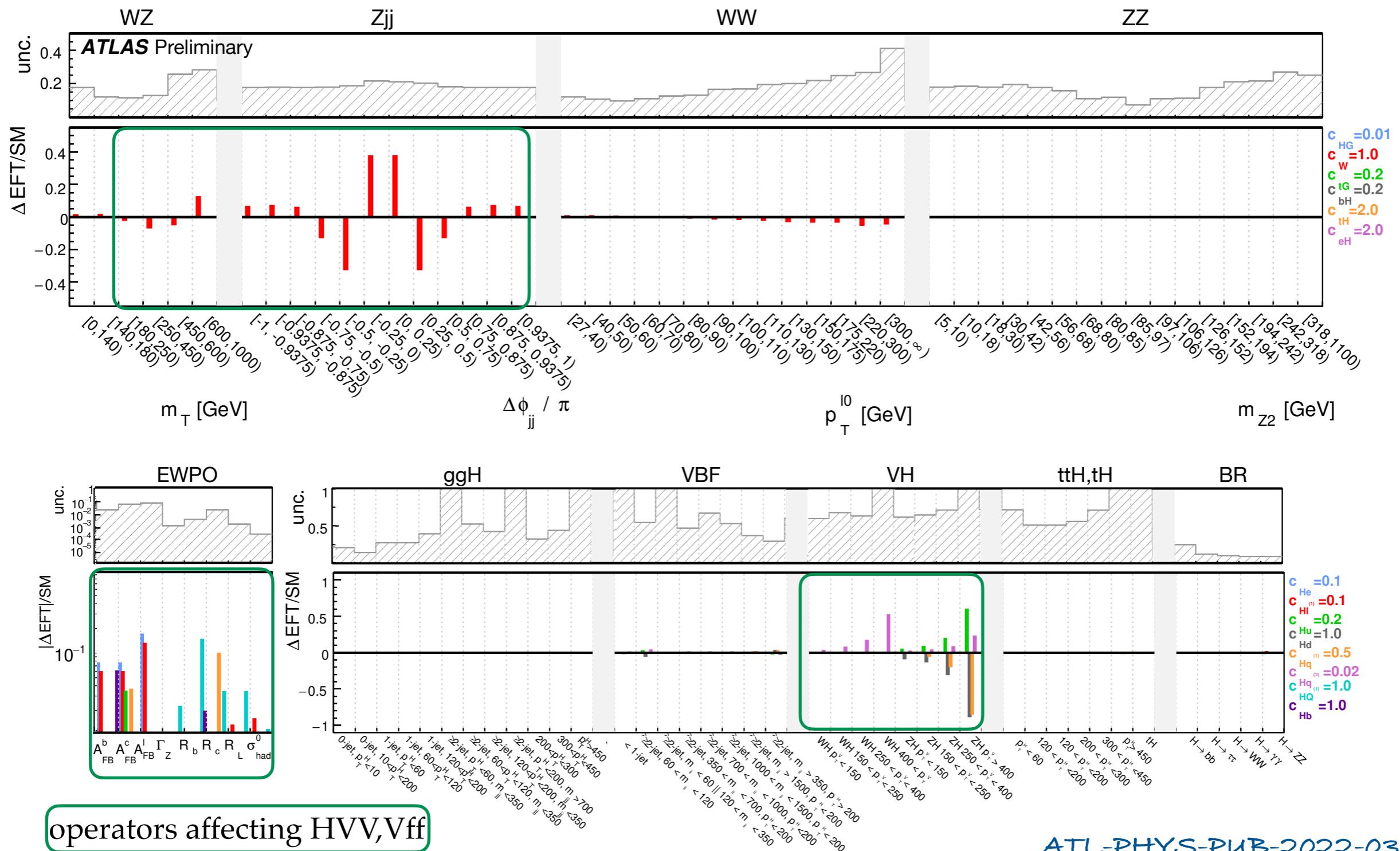
Impact of linear SMEFT parameterisation shown for bins along with corresponding measurement uncertainty

- Relative impact of linear SMEFT terms with Wilson coefficients  $c_{HG}$ ,  $c_W$ ,  $c_{tG}$ ,  $c_{bH}$ ,  $c_{tH}$ , and  $c_{eH}$  on the Higgs STXS cross sections and branching ratios.



# ATLAS Global combination

- Additional sensitivity coming from EW measurements and EWPO, e.g. cW that cannot be disentangled using just  $H \rightarrow \gamma\gamma$  decay.

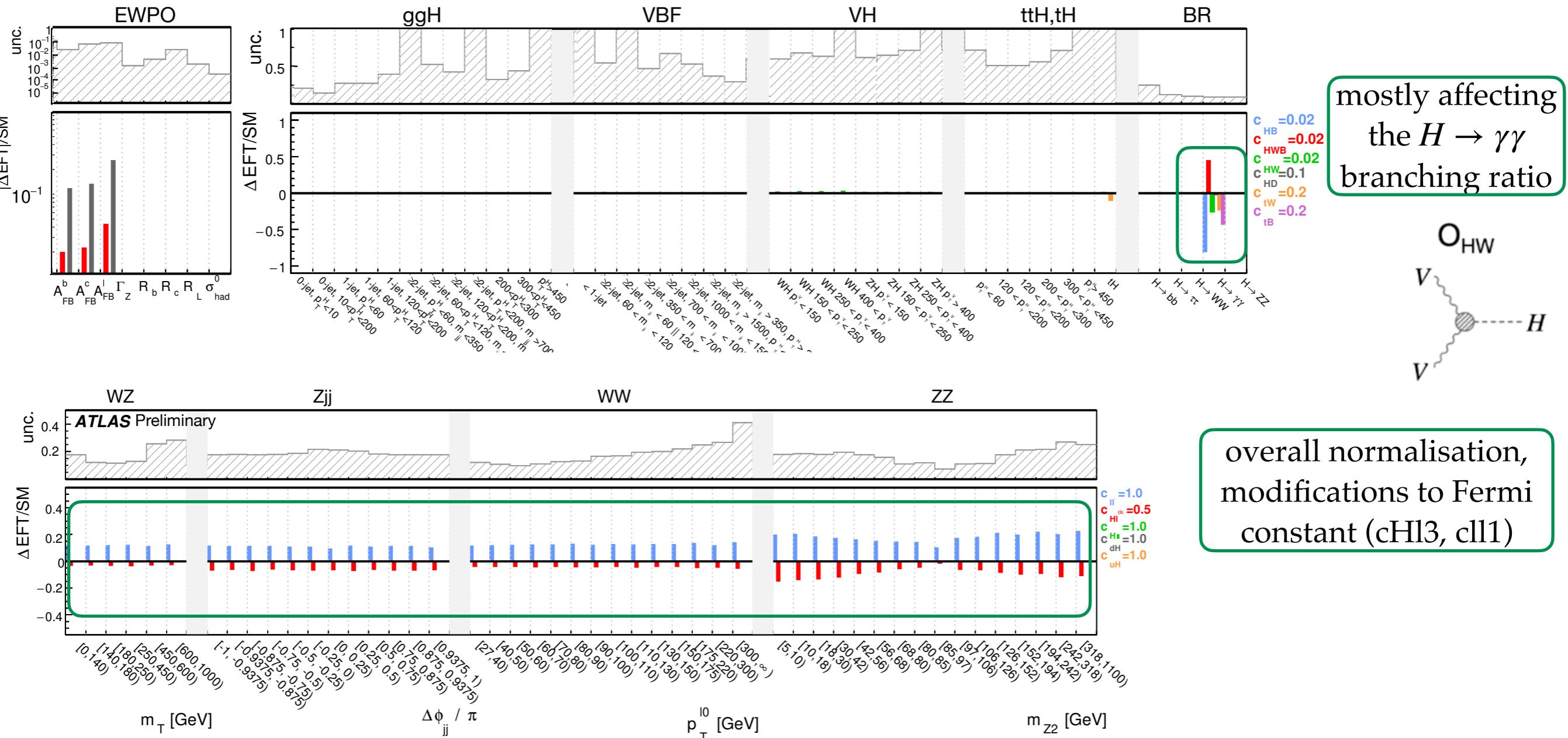


operators affecting HVV, Vff

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# ATLAS Global combination

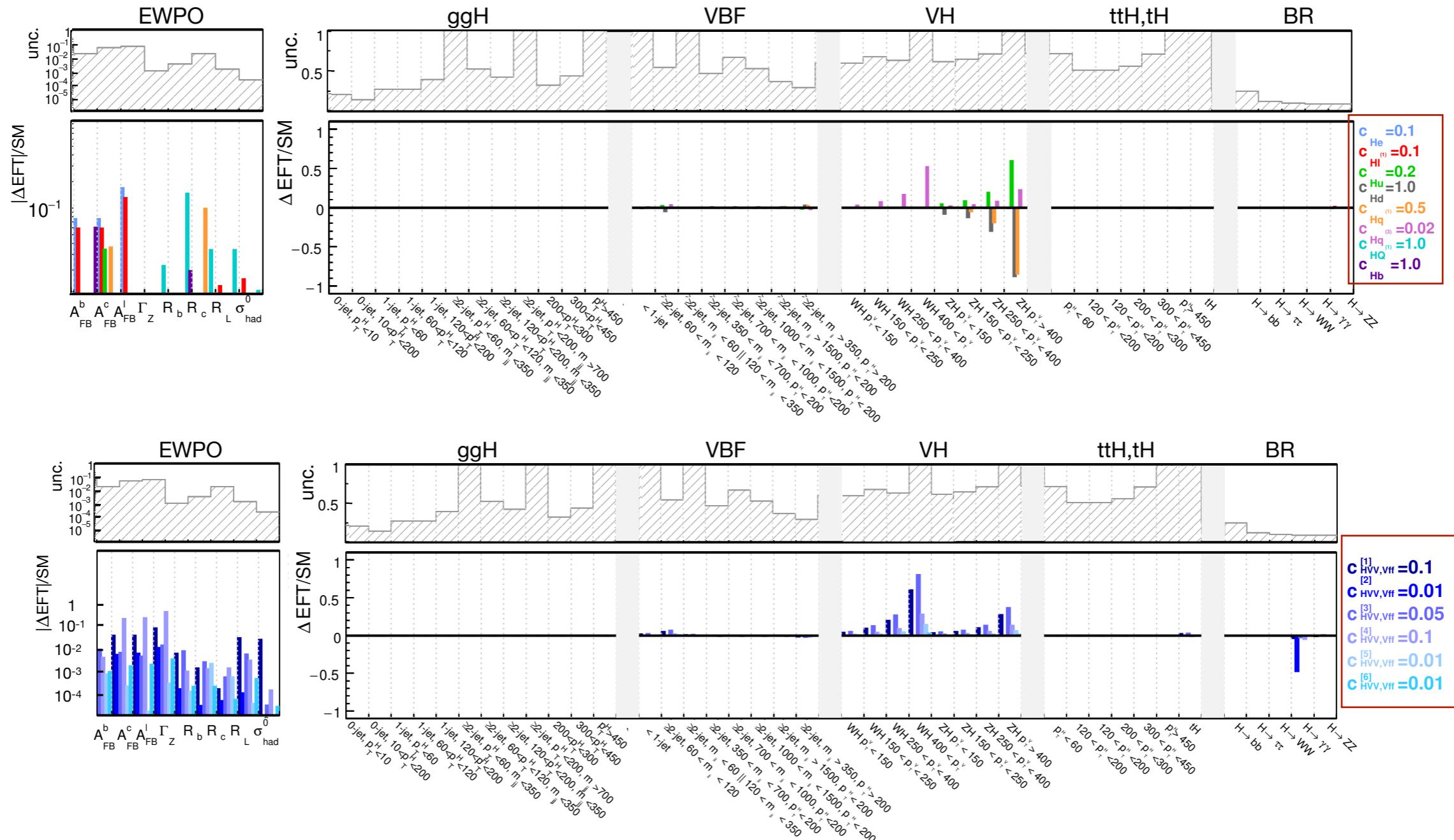
Impact of linear SMEFT parameterisation shown for bins along with corresponding measurement uncertainty



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# ATLAS Global combination

- SMEFT impact on measurements shown in Warsaw basis and fit basis-> allows to understand the impact of the different fit directions on measurements.



# ATLAS Global combination: sensitivity studies

- Principal Component Analysis to reduce the dimensionality of the fit. [ATL-PHYS-PUB-2022-037](#)
- A simultaneous measurement of the signal strengths in each measurement bin is performed using a maximum likelihood fit exploiting a signal strength parameter  $\mu_b$ :

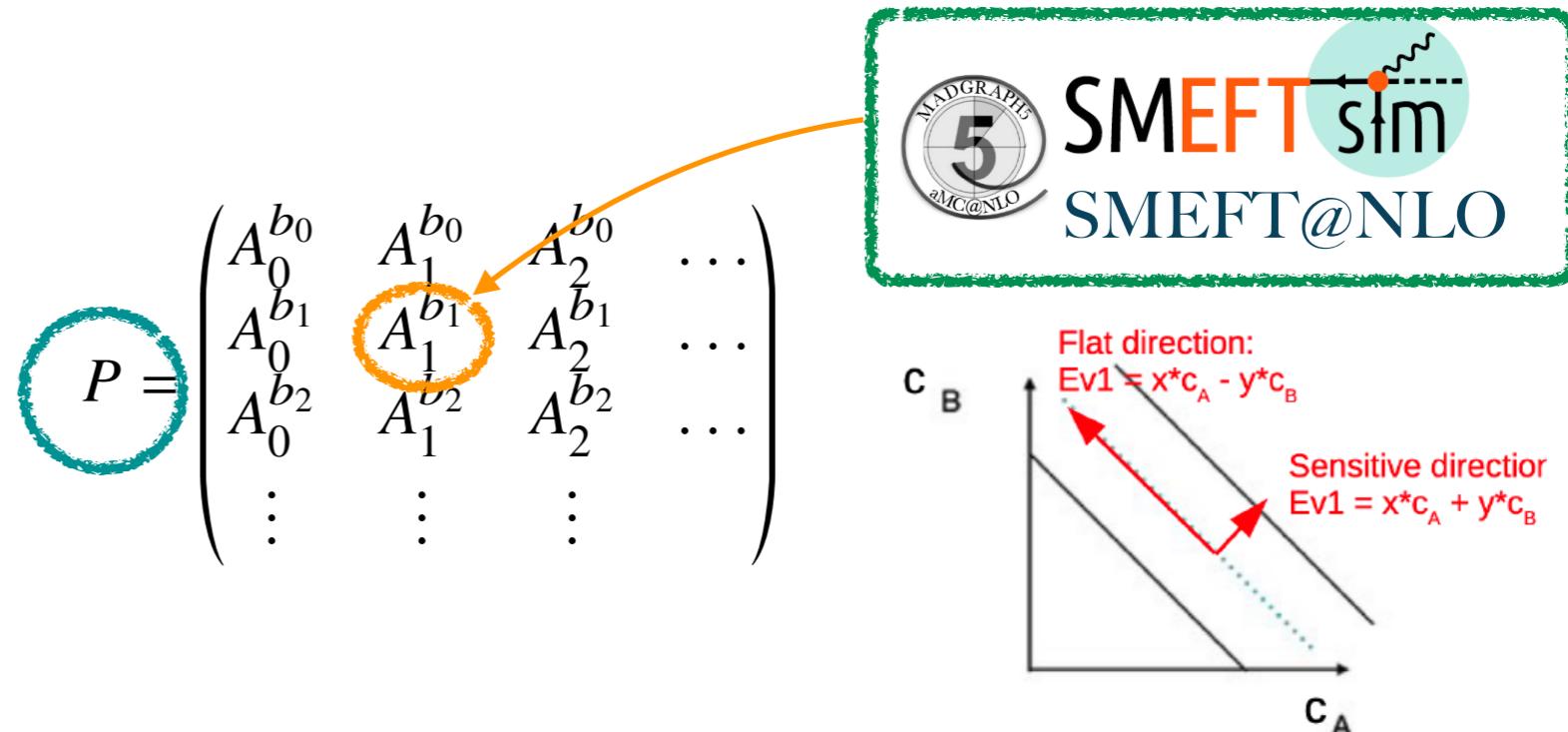
$$1 + \sum_i A_{bi} c_i + \sum_i B_{bi} c_i^2 + \sum_{i < j} C_{bij} c_i c_j \rightarrow \mu_b$$

- The Hessian matrix  $H_\mu$  at the minimum of the negative log likelihood of the  $\mu_b$  measurement is reparameterised in terms of Wilson coefficients using the linear parameterisation matrix  $P$ .
- $H_{SMEFT}$ , can be obtained from the Hessian matrix  $H_\mu$  in  $\mu_b$  space:

$$H_{SMEFT} = P^T H_\mu P$$

- $P$ : matrix that gives the parametrisation of the observables as a function of the Wilson coefficients;
- $A_i^{b_j}$  : factors obtained from the simulation.

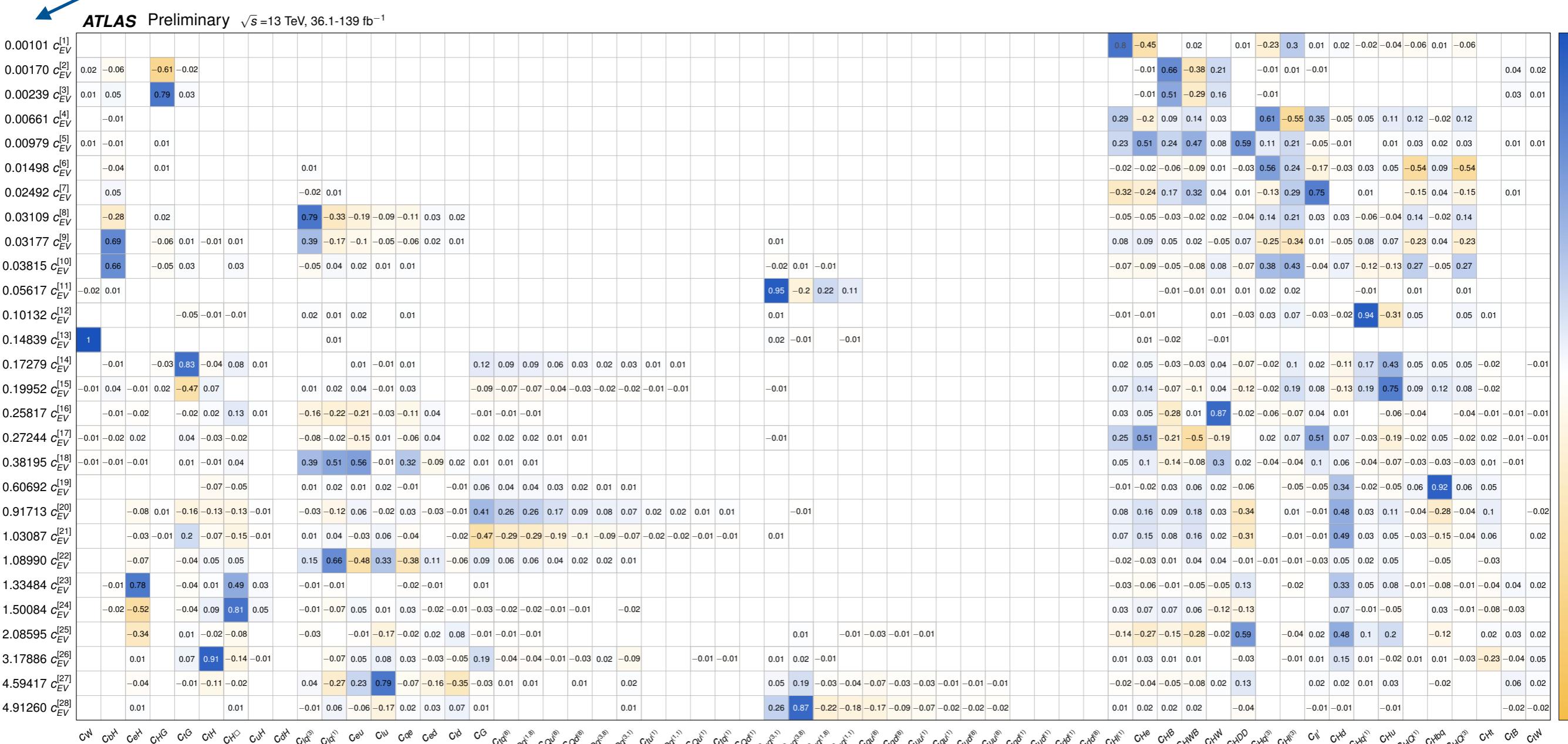
$$P = \begin{pmatrix} A_0^{b_0} & A_1^{b_0} & A_2^{b_0} & \dots \\ A_0^{b_1} & A_1^{b_1} & A_2^{b_1} & \dots \\ A_0^{b_2} & A_1^{b_2} & A_2^{b_2} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$$



# ATLAS Global combination: sensitivity studies

- PCA considering all operators: directions ordered by increasing uncertainties, keeping  $\sigma < 5$ ;
- Wilson coefficients expected to be at most order 1, new physics scale  $\Lambda$  expected to be at least 1 TeV -> directions with  $\sigma > 5$  have very little impact on the measurement.

## Eigenvectors from PCA, corresponding eigenvalues

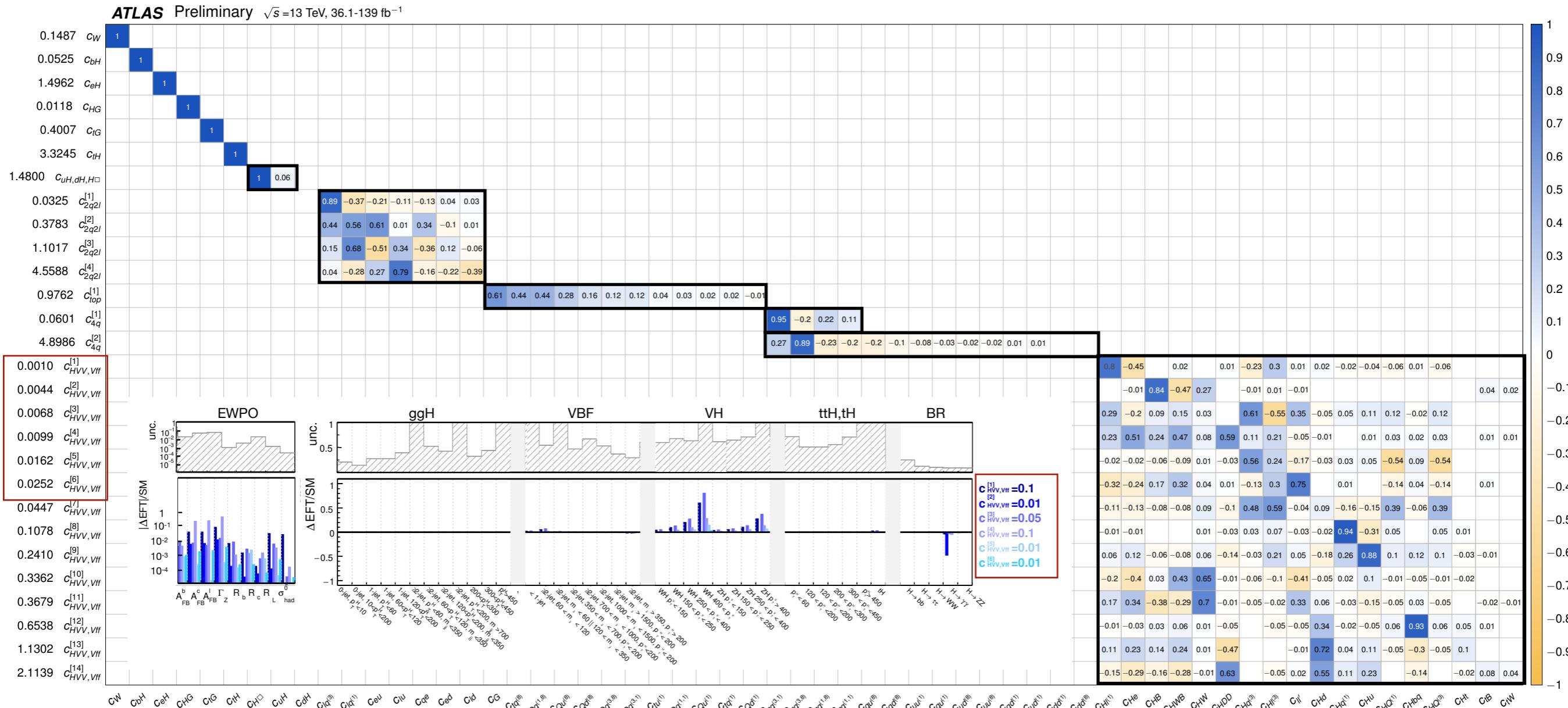


ATL-PHYS-PUB-2022-037 Warsaw Basis, Wilson coefficients

# ATLAS Global combination: sensitivity studies

ATL-PHYS-PUB-2022-037

- PCA considering all operators: directions ordered by increasing uncertainties, keeping  $\sigma < 5$ ;
- Fit basis defined by grouping operators of similar physics impact together.



**Warsaw Basis, Wilson coefficients**