

DK and related hadronic molecules

Feng-Kun Guo

Institute of Theoretical Physics, Chinese Academy of Sciences

Exotic Hadrons in LHCb, 18 Oct. 2022

L. Liu, K. Orginos, FKG, C. Hanhart, U.-G. Meißner, PRD87(2013)014508;

M. Albaladejo, P. Fernandez-Soler, FKG, J. Nieves, PLB767(2017)465;

M.-L. Du, M. Albaladejo, P. Fernandez-Soler, FKG, C. Hanhart, U.-G. Meißner, J. Nieves, D.-L. Yao, PRD98(2018)094018;

M.-L. Du, FKG, U.-G. Meißner, PRD99(2019)114002;

M.-L. Du, FKG, C. Hanhart, B. Kubis, U.-G. Meißner, PRL126(2021)192001; ...

Positive-parity charmed mesons (1)

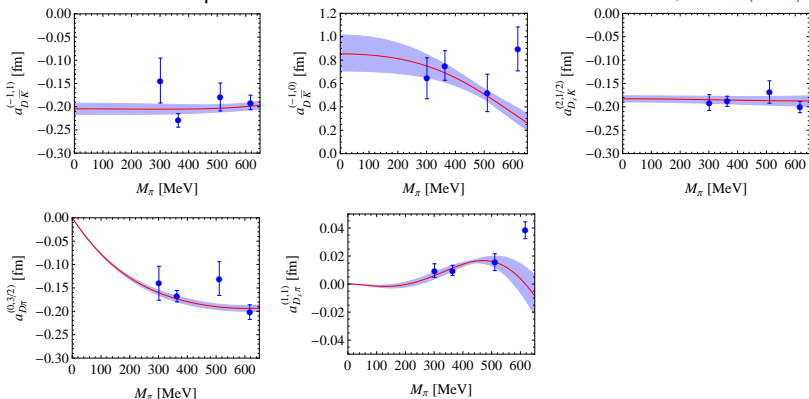
- Hadronic molecular model: $D_{s0}^*(2317)[DK], D_{s1}(2460)[D^*K]$

Barnes, Close, Lipkin(2003); van Beveren, Rupp(2003); Kolomeitsev, Lutz(2004); FKG et al.(2006);



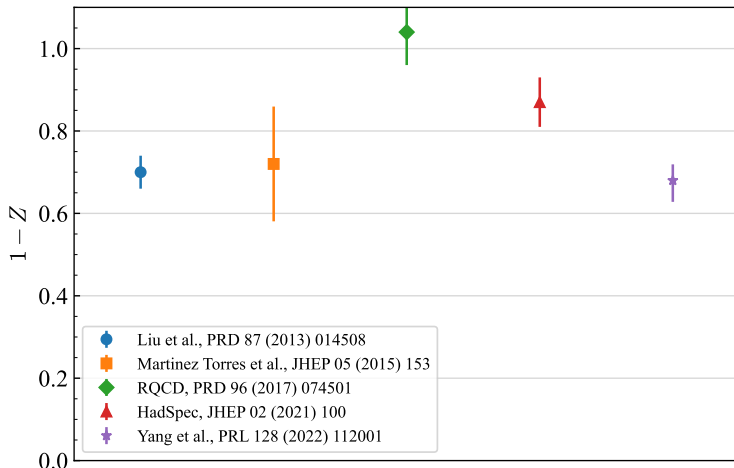
- Unitarized CHPT up to NLO

L.Liu et al., PRD87(2013)014508

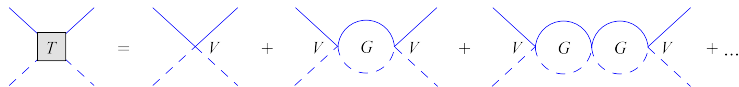


Positive-parity charmed mesons (2): compositeness

- Postdicted mass for $D_{s0}^*(2317)$: 2315_{-28}^{+18} MeV
- Compositeness (DK component) for $D_{s0}^*(2317)$ from (in)direct lattice calculations or analyses of lattice results



Positive-parity charmed mesons (3): bottom-strange



- Natural consequence of HQSS:

similar binding energies $M_D + M_K - M_{D_{s0}^*} \simeq 45 \text{ MeV}$

$M_{D_{s1}(2460)} - M_{D_{s0}^*(2317)} \simeq M_{D^*} - M_D$ is understood

- HQFS: predicting the bottom-partner masses (UCHPT Du et al., PRD98(2018)094018):

$$M_{B_{s0}^*} \simeq M_B + M_K - 45 \text{ MeV} \simeq 5730 \text{ MeV} \quad (5720_{-23}^{+16} \text{ MeV})$$

$$M_{B_{s1}} \simeq M_{B^*} + M_K - 45 \text{ MeV} \simeq 5776 \text{ MeV} \quad (5772_{-21}^{+15} \text{ MeV})$$

to be compared with lattice results for the lowest positive-parity bottom-strange mesons:

Lang, Mohler, Prelovsek, Woloshyn, PLB750(2015)17

$$M_{B_{s0}^*}^{\text{lat.}} = (5.711 \pm 0.013 \pm 0.019) \text{ GeV}$$

$$M_{B_{s1}}^{\text{lat.}} = (5.750 \pm 0.017 \pm 0.019) \text{ GeV}$$

These B_{s0}^* and B_{s1} not found yet; for updated predictions on their radiative/pionic decays, see

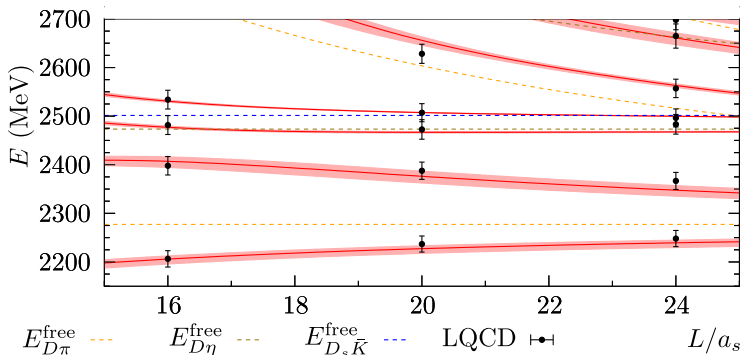
H.-L. Fu, Griebhammer, FKG, Hanhart, Meißner, EPJA58(2022)70

Positive-parity charmed mesons (4): charm-nonstrange

- In a finite volume: $\vec{q} = \frac{2\pi}{L}\vec{n}$, $\vec{n} \in \mathbb{Z}^3$; loop integral $G(s)$: $\int d^3\vec{q} \rightarrow \frac{1}{L^3} \sum_{\vec{q}}$
- Postdicted $I = 1/2$ $D\pi$, $D\eta$, $D_s\bar{K}$ finite volume energy levels in the c.m. frame versus lattice QCD results by [G. Moir *et al.* [HadSpec], JHEP1610(2016)011]

NOT a fit!

M. Albaladejo, P. Fernandez-Soler, FKG, J. Nieves, PLB767(2017)465



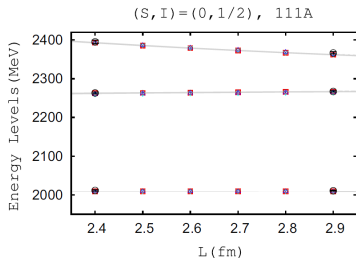
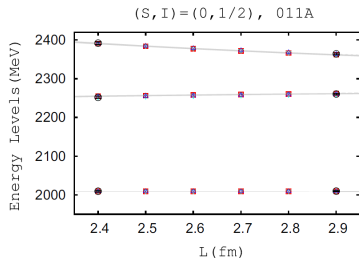
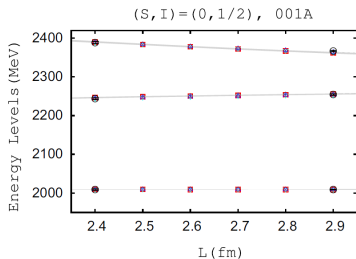
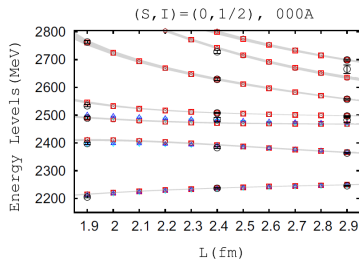
consequence of SU(3) + chiral

Positive-parity charmed mesons (5): charm-nonstrange

A more recent fit to lattice data including the moving frame ones

Z.-H. Guo et al.,

EPJC79(2019)13



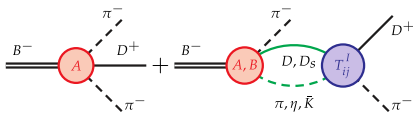
Positive-parity charmed mesons (6): $D_0^*(2300)$ contains TWO states!

- Heavy-nonstrange, two $I = 1/2$ states ($M, \Gamma/2$)

Du et al., PRD98(2018)094018

	Lower (MeV)	Higher (MeV)	PDG2022 (MeV)
D_0^*	$(2105_{-8}^{+6}, 102_{-11}^{+10})$	$(2451_{-26}^{+36}, 134_{-8}^{+7})$	$(2343 \pm 10, 115 \pm 8)$
D_1	$(2247_{-6}^{+5}, 107_{-10}^{+11})$	$(2555_{-30}^{+47}, 203_{-9}^{+8})$	$(2412 \pm 9, 157 \pm 15)$
B_0^*	$(5535_{-11}^{+9}, 113_{-17}^{+15})$	$(5852_{-19}^{+16}, 36 \pm 5)$	—
B_1	$(5584_{-11}^{+9}, 119_{-17}^{+14})$	$(5912_{-18}^{+15}, 42_{-4}^{+5})$	—

With the same LECs, the LHCb data for a set of B decays are well described



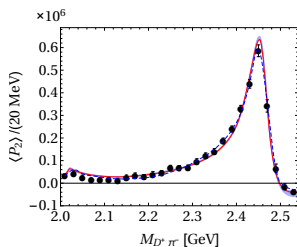
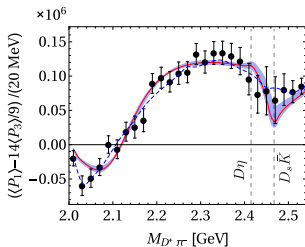
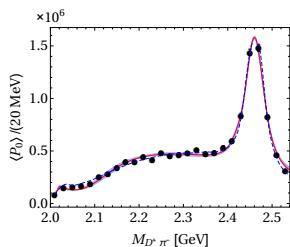
- S -wave: FSI, two new parameters
- P, D -wave: BWs from the LHCb fit

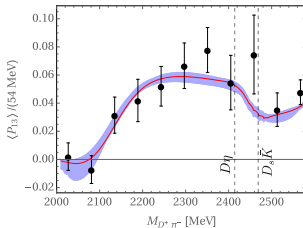
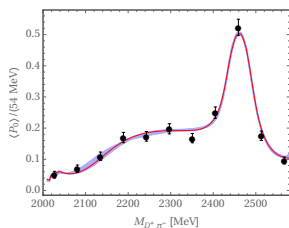
Angular moments for $B^- \rightarrow D^+ \pi^- \pi^-$:

$$\langle P_0 \rangle \propto |\mathcal{A}_0|^2 + |\mathcal{A}_1|^2 + |\mathcal{A}_2|^2, \quad \langle P_2 \rangle \propto \frac{2}{5} |\mathcal{A}_1|^2 + \frac{2}{7} |\mathcal{A}_2|^2 + \frac{2}{\sqrt{5}} |\mathcal{A}_0| |\mathcal{A}_2| \cos(\delta_2 - \delta_0),$$

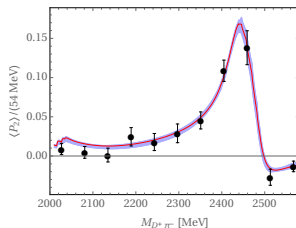
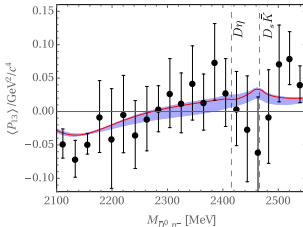
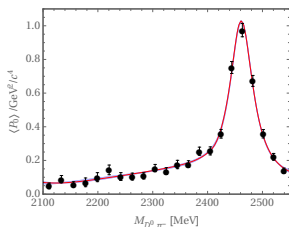
$$\langle P_{13} \rangle \equiv \langle P_1 \rangle - \frac{14}{9} \langle P_3 \rangle \propto \frac{2}{\sqrt{3}} |\mathcal{A}_0| |\mathcal{A}_1| \cos(\delta_1 - \delta_0)$$

Data: LHCb, PRD94(2016)072001

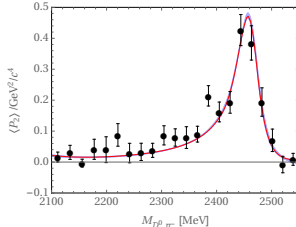


Fit to data of $B^- \rightarrow D^+ \pi^- K^-$ 

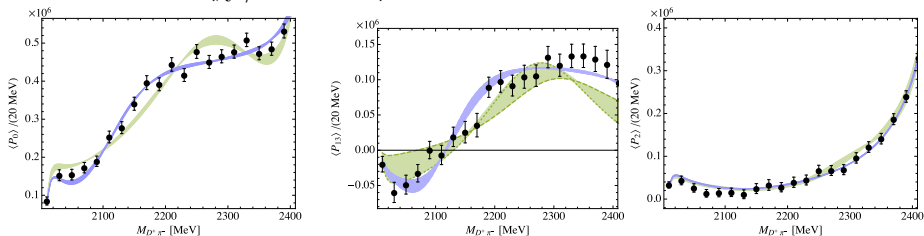
Data: LHCb, PRD91(2015)092002

Fit to data of $B^0 \rightarrow \bar{D}^0 \pi^- \pi^+$ 

Data: LHCb, PRD92(2015)032002

and also $B^0 \rightarrow \bar{D}^0 \pi^- K^+$, $B_s^0 \rightarrow \bar{D}^0 K^- \pi^+$

- Fits with the Khuri-Treiman equation taking into account three-body unitarity: using S -wave $D\pi$ scattering phase from UCHPT ($\chi^2/\text{d.o.f.} = 1.2$) and from BW ($\chi^2/\text{d.o.f.} = 2.0$)



- The LHCb data are well described with UCHPT amplitude with two D_0^* states; the lower one has a mass about 2.1 GeV

- Support from lattice results:

D_0^* with $M \approx 2.2$ GeV and $\Gamma \approx 0.4$ GeV obtained using $M_\pi \approx 239$ MeV L. Gayer et al. [HadSpec], JHEP07(2021)123

$M = (2.12 \pm 0.03)$ MeV with $M_\pi \approx 266$ MeV

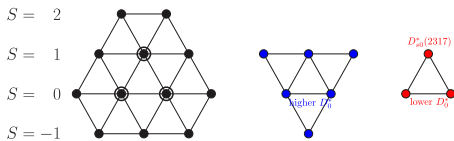
Mohler, Prelovsek, Woloshyn,

PRD87(2013)034501; pole given in J. Bulava et al., arXiv:2203.03230

Positive-parity charmed mesons (9): SU(3) structure

- SU(3) irreps: $\bar{\mathbf{3}} \otimes \mathbf{8} = \bar{\mathbf{15}} \oplus \mathbf{6} \oplus \bar{\mathbf{3}}$

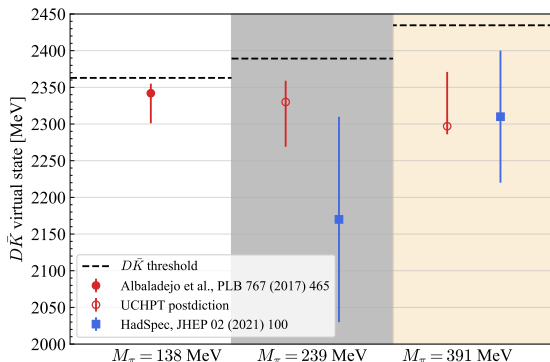
Albaladejo et al., PLB767(2017)465



$(S, I) = (-1, 0)$: virtual state at 2342^{+13}_{-41} MeV at the physical mass

- HadSpec also found a virtual state

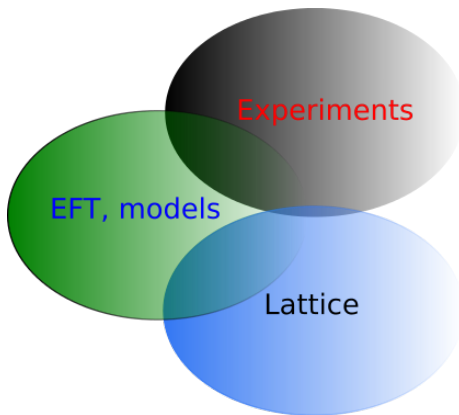
HadSpec, arXiv:2008.06432



Summary and speculations

- Very strong evidence supporting D_{s0}^* (2317) and D_{s1} (2460) to be molecules
- ☞ By the same leading order interaction, there should be a whole family of kaonic bound states FKG, Meißner, PRD84(2012)014013
 D_{s1} (2860) as D_1K with $J^P = 1^-$; D_2K molecule with 2^- , $M \simeq 2.91$ GeV;
 $\Xi_{cc}\bar{K}$ molecule; $B_{1,2}\bar{K}$ molecules ($M \simeq 6.15, 6.17$ GeV); ...
- Kaons are pseudo-Goldstone bosons, the interaction receives chiral suppression. If they form hadronic molecules, there must be many other molecules. Du et al., PRD98(2018)094018
- $D^{(*)}\bar{K}$ form $I = 0$ virtual states (there should also be $D_{1,2}\bar{K}$ virtual states similarly); it's reasonable to expect $D^{(*)}\bar{K}^*$ would have bound states.
 $D^*\bar{K}^*$ molecule: $X(2900)$ found by LHCb?
Then there must be partners:
spin partner $D\bar{K}^*$: ~ 2760 MeV
bottom partners BK^* : ~ 6175 MeV; B^*K^* : ~ 6220 MeV
Then there might be deeper bound $D^{(*)}K^*/B^{(*)}\bar{K}^*$ isoscalar states.

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Thank you !

Backup slides

Compositeness (1)

Weinberg, PR137(1965); Baru *et al.*, PLB586(2004); ...

see also, e.g., Weinberg's books: QFT Vol.I, Lectures on QM

Model-independent result for S -wave loosely bound composite states:

Consider a system with Hamiltonian

$$\mathcal{H} = \mathcal{H}_0 + V$$

\mathcal{H}_0 : free Hamiltonian, V : interaction potential

- **Compositeness:**

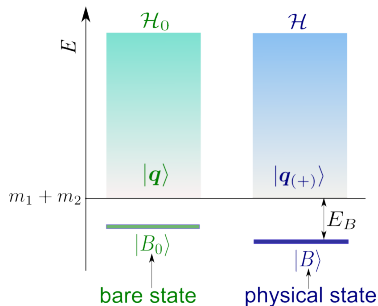
the probability of finding the physical state $|B\rangle$ in the 2-body continuum $|\mathbf{q}\rangle$

$$1 - Z = \int \frac{d^3\mathbf{q}}{(2\pi)^3} |\langle \mathbf{q} | B \rangle|^2$$

- $Z = |\langle B_0 | B \rangle|^2$, $0 \leq (1 - Z) \leq 1$

- ☞ $Z = 0$: pure composite state

- ☞ $Z = 1$: pure elementary state



Compositeness : $1 - Z = \int \frac{d^3 \mathbf{q}}{(2\pi)^3} |\langle \mathbf{q} | B \rangle|^2$

- Schrödinger equation

$$(\mathcal{H}_0 + V)|B\rangle = -E_B|B\rangle$$

multiplying by $\langle \mathbf{q} |$ and using $\mathcal{H}_0|\mathbf{q}\rangle = \frac{\mathbf{q}^2}{2\mu}|\mathbf{q}\rangle$

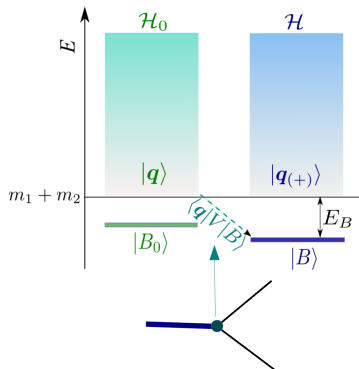
$$\langle \mathbf{q} | B \rangle = -\frac{\langle \mathbf{q} | V | B \rangle}{E_B + \mathbf{q}^2 / (2\mu)}$$

- S*-wave, small binding energy so that $R = 1/\sqrt{2\mu E_B} \gg r$, r : range of forces

$$\langle \mathbf{q} | V | B \rangle = g_{\text{NR}} [1 + \mathcal{O}(r/R)]$$

- Compositeness:

$$1 - Z = \int \frac{d^3 \mathbf{q}}{(2\pi)^3} \frac{g_{\text{NR}}^2}{[E_B + \mathbf{q}^2 / (2\mu)]^2} \left[1 + \mathcal{O}\left(\frac{r}{R}\right) \right] = \frac{\mu^2 g_{\text{NR}}^2}{2\pi \sqrt{2\mu E_B}} \left[1 + \mathcal{O}\left(\frac{r}{R}\right) \right]$$



$X(3872)$ [$\chi_{c1}(3872)$]

- discovered in $B^\pm \rightarrow K^\pm J/\psi \pi \pi$

Belle, PRL91(2003)262001

$$M_{D^0} + M_{D^{*0}} - M_X = (0.00 \pm 0.18) \text{ MeV}$$

- $\Gamma < 1.2 \text{ MeV}$
- $J^{PC} = 1^{++}$, S -wave coupling to $D\bar{D}^*$
- Observed in the $D^0\bar{D}^{*0}$ mode as well

Belle, PRD84(2011)052004

LHCb PRL110(2013)222001

BaBar, PRD77(2008)011102

- Large coupling to $D^0\bar{D}^{*0}$:

$$\mathcal{B}(X \rightarrow D^0\bar{D}^{*0}) > 30\%$$

PDG2022

- Large isospin breaking:

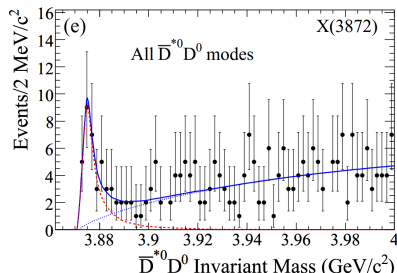
$$\frac{\mathcal{B}(X \rightarrow \omega J/\psi)}{\mathcal{B}(X \rightarrow \pi^+\pi^- J/\psi)} = 0.8 \pm 0.3$$

- $1^{++} D\bar{D}^*$ bound state around 3.87 GeV

predicted by Törnqvist

ZPC61(1994)525

- The large coupling to $D^0\bar{D}^{*0} + c.c.$ implies that $X(3872)$ must be an extended object with the longest-distance component given by $D^0\bar{D}^{*0} + c.c.$



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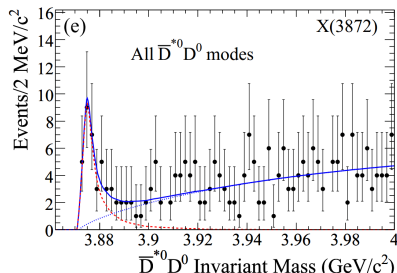
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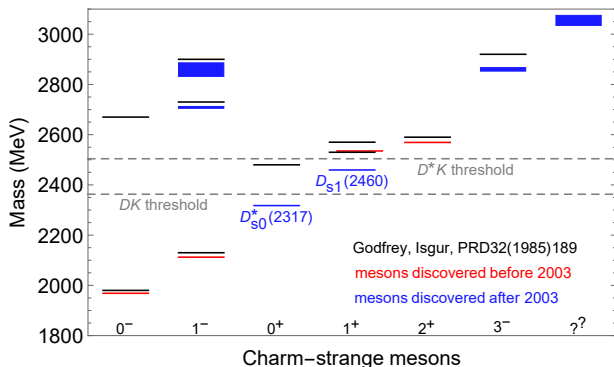
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Charm-strange mesons: D_{s0}^* (2317) and D_{s1} (2460)

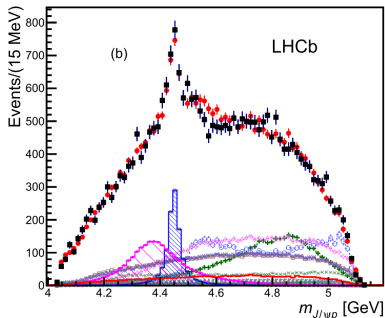
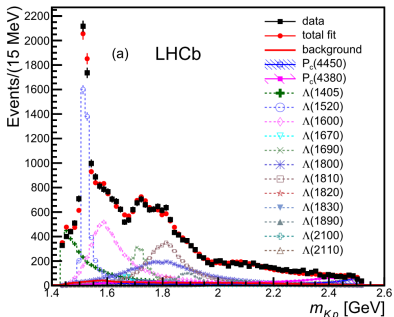
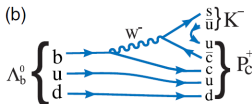
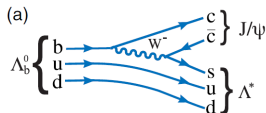


- D_{s0}^* (2317): BaBar (2003)
 $J^P = 0^+$, $\Gamma < 3.8$ MeV
- D_{s1} (2460): CLEO (2003)
 $J^P = 1^+$, $\Gamma < 3.5$ MeV
- no isospin partner observed, tiny widths
 $\Rightarrow I = 0$

- Mass problem: Why are D_{s0}^* (2317) and D_{s1} (2460) so light?
- Naturalness problem: Why $\underbrace{M_{D_{s1}(2460)} - M_{D_{s0}^*(2317)}}_{(141.8 \pm 0.8) \text{ MeV}} \simeq \underbrace{M_{D^{*+}} - M_{D^+}}_{(140.67 \pm 0.08) \text{ MeV}} ?$

Discovered in $\Lambda_b^0 \rightarrow J/\psi p K^-$

LHCb, PRL115(2015)072001 [arXiv:1507.03414]



Two Breit–Wigner resonances, quantum numbers not fixed:

$$M_1 = (4380 \pm 8 \pm 29) \text{ MeV},$$

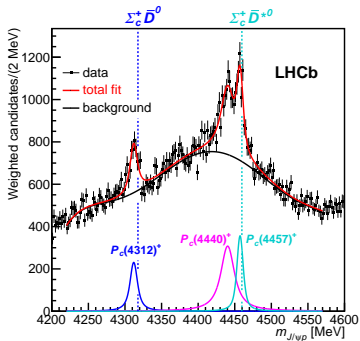
$$\Gamma_1 = (205 \pm 18 \pm 86) \text{ MeV},$$

$$M_2 = (4449.8 \pm 1.7 \pm 2.5) \text{ MeV},$$

$$\Gamma_2 = (39 \pm 5 \pm 19) \text{ MeV}.$$

The 2019 update of LHCb's P_c : three narrow states

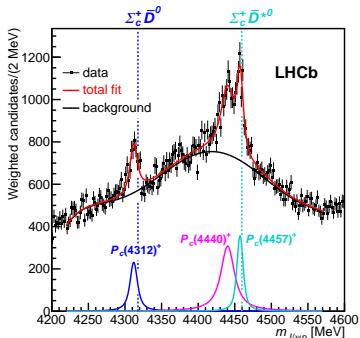
LHCb, PRL122(2019)222001



State	M [MeV]	Γ [MeV]
$P_c(4312)^+$	$4311.9 \pm 0.7^{+6.8}_{-0.6}$	$9.8 \pm 2.7^{+3.7}_{-4.5}$
$P_c(4440)^+$	$4440.3 \pm 1.3^{+4.1}_{-4.7}$	$20.6 \pm 4.9^{+8.7}_{-10.1}$
$P_c(4457)^+$	$4457.3 \pm 0.6^{+4.1}_{-1.7}$	$6.4 \pm 2.0^{+5.7}_{-1.9}$

The 2019 update of LHCb's P_c : three narrow states

LHCb, PRL122(2019)222001



$\bar{D}^{(*)}\Sigma_c$ molecules predicted:

prediction of narrow N^* and Λ^* resonances with hidden charm above 4 GeV, J.-J. Wu, R. Molina,

E. Oset, B.-S. Zou, PRL105(2010)232001

Other predictions: W.L.Wang et al.('11); Z.C.Yang et al.('12); Xiao, Nieves, Oset('13); Karliner, Rosner('15)

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$P_c(4457)^+$	$4457.3 \pm 0.6^{+4.1}_{-1.7}$	$6.4 \pm 2.0^{+5.7}_{-1.9}$

J^P	Λ	$\bar{D}\Sigma_c$ PB system	$\bar{D}^*\Sigma_c$ VB system
		$M - i\Gamma/2$	$M - i\Gamma/2$
$\frac{1}{2}^-$	650	–	–
	800	–	$4462.178 - 0.002i$
	1200	$4318.964 - 0.362i$	$4459.513 - 0.417i$
	1500	$4314.531 - 1.448i$	$4454.088 - 1.662i$
	2000	$4301.115 - 5.835i$	$4438.277 - 7.115i$
$\frac{3}{2}^-$	650	–	–
	800	–	$4462.178 - 0.002i$
	1200	–	$4459.507 - 0.420i$
	1500	–	$4454.057 - 1.681i$
	2000	–	$4438.039 - 7.268i$

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