# Prospects for searches of $b \rightarrow s \nu \bar{\nu}$ decays 

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## Physics motivation

- Considerable interest in the flavour community in $b \rightarrow s \ell^{+} \ell^{-}$and $b \rightarrow c \ell^{-} \bar{\nu}$ transitions
- $b \rightarrow s \nu \bar{\nu}$ transitions are complementary probes ( $\ell^{+}$and $\nu$ share a weak doublet)
- SM predictions are clean:
- Dominant uncertainties from hadronic form-factors and CKM elements
- No long-distance contributions from (in)famous charm loops
- Sensitive to a variety of NP scenarios e.g. $Z^{\prime}$, leptoquarks etc.


$$
\mathcal{H}_{\mathrm{eff}}=-\frac{G_{F}}{\sqrt{2}} V_{t b} V_{t s}^{*} \sum_{i j}\left(C_{L}^{i j} O_{L}^{i j}+C_{R}^{i j} O_{R}^{i j}\right)+h . c
$$

- In the $\mathrm{SM}, C_{L}^{i i}=-6.35(7)$ and $C_{R}^{i j}=0[1,2,3,4]$


## Experimental state-of-the-art

- FCC-ee provides a (possibly unique) opportunity for semileptonic flavour physics
- A beauty/charm factory at the $Z$ gets the best of both LHCb and $B$-factories
- $e^{+}-e^{-}$collision, high production rate, access to high mass states, hermetic detector
- In the $\mathrm{SM} b \rightarrow s \nu \bar{\nu} \mathrm{BF}$ predictions are $\mathcal{O}\left(10^{-5}\right)$
- Not yet seen experimentally
- From the underlying $b \rightarrow s \nu \bar{\nu}$ transition we can study:

| Decay | B-factories | FCC-ee | Current Limit | SM prediction |
| :---: | :---: | :---: | :---: | :---: |
| $B^{+} \rightarrow K^{+} \nu \bar{\nu}$ | $\checkmark$ | $\checkmark$ | $<1.6 \times 10^{-5}$ | $(4.0 \pm 0.5) \times 10^{-6}$ |
| $B^{+} \rightarrow K^{*+} \nu \bar{\nu}$ | $\nu$ | $\checkmark$ | $<4.0 \times 10^{-5}$ | $(9.8 \pm 1.1) \times 10^{-6}$ |
| $B^{0} \rightarrow K_{\mathrm{S}}^{0} \nu \bar{\nu}$ | $\checkmark$ | $\checkmark$ | $<2.6 \times 10^{-5}$ | $(3.7 \pm 0.4) \times 10^{-6}$ |
| $B^{0} \rightarrow K^{* 0} \nu \bar{\nu}$ | $\checkmark$ | $\checkmark$ | $<1.8 \times 10^{-5}$ | $(9.2 \pm 1.0) \times 10^{-6}$ |
| $B_{s}^{0} \rightarrow \phi \nu \bar{\nu}$ | $x$ | $\checkmark$ | $<5.4 \times 10^{-3}$ | $(9.9 \pm 0.7) \times 10^{-6}$ |
| $\Lambda_{b}^{0} \rightarrow \Lambda^{0} \nu \bar{\nu}$ | $x$ | $\checkmark$ | - | - |

- Decays with intermediate vectors are consierably easier experimentally
- single track is hard, final state neutral needs good $K_{\mathrm{S}}^{0} / \Lambda^{0}$ reco
- intermediate scalars are much cleaner for theory
- Decays with intermediate scalars are cleaner for theory
- With 2 neutrinos in the final state, decays are (probably) impossible at the LHC $3 / 28$


## Event topology

We have studied the prospects for $B^{0} \rightarrow K^{* 0} \nu \bar{\nu}$ and $B_{s}^{0} \rightarrow \phi \nu \bar{\nu}$

- Use the thrust axis for $Z^{0} \rightarrow q \bar{q}$ to define event hemispheres
- Due to missing energy in the signal decay the two hemispheres have different energy distributions



## Energy in each hemisphere





## Event-level MVA

- Background sample from inclusive $Z^{0} \rightarrow q \bar{q}, c \bar{c}, b \bar{b}$ using PDG branching fractions
- Input variables are the event energy distributions and vertex information


- Powerful seperation - cut at 0.6 has $>90 \%$ signal efficiency and $\sim 90 \%$ background rejection
- Very similar for the $B_{s}^{0} \rightarrow \phi \nu \bar{\nu}$ mode


## Analysis-level MVA

- Train a second BDT on variables related to the candidate properties:
- Intermediate candidate kinematics
- Intermediate candidate topology
- The nominal $B$-meson energy ( $Z$ mass minus $E_{\text {rec }}$ )
- Use multivariate splines to build efficiency maps across the (BDT1, BDT2) plane
- Then maximise the FOM, $S / \sqrt{S+B}$, as a function of the BDT cuts for a range of BF values

$$
B^{0} \rightarrow K^{* 0} \nu \bar{\nu}
$$




## Signal estimate

- Signal expectation is computed as

$$
S=N_{Z} \mathcal{B}(Z \rightarrow b \bar{b}) 2 f_{B} \mathcal{B}(B \rightarrow Y \nu \bar{\nu}) \mathcal{B}(Y \rightarrow f) \epsilon_{\mathrm{pre}}^{s} \epsilon_{\mathrm{BDTs}}^{s}
$$

- Background expectation computed as

$$
B=\sum_{f \in\{b \bar{b}, c \bar{c}, q \bar{q}\}} N_{Z} \mathcal{B}(Z \rightarrow f) \epsilon_{\text {pre }}^{b} \epsilon_{\mathrm{BDTs}}^{b}
$$

assuming

- $3 \times 10^{12} Z^{0}$ in FCC-ee operation (needs updating to $2 \times 10^{12}$ )
- known / predicted production fractions and branching ratios
- analysis efficiencies

For optimal cuts at the SM prediction:

- Signal efficiency $\sim 10 \%$
- $b \bar{b}$ efficiency $\sim 10^{-5}$
- $c \bar{c}$ efficiency $\sim 10^{-6}$
- $q \bar{q}$ efficiency $\sim 10^{-8}$
- S/B ratio $\sim 1: 20$
- Sensitivity $\sim 0.5 \%$



## $B_{s}^{0} \rightarrow \phi \nu \bar{\nu} \quad$ Efficiency and Sensitivity

For optimal cuts at the SM prediction:

- Signal efficiency $\sim 11 \%$
- $b \bar{b}$ efficiency $\sim 10^{-6}$
- $c \bar{c}$ efficiency $\sim 10^{-8}$
- $q \bar{q}$ efficiency $\sim 10^{-9}$
- S/B ratio $\sim 1: 9$
- Sensitivity $\sim 1.3 \%$
- CEPC at $\sim 1.8 \%$ [5]



## PID requirements of the detector

- For serious flavour analysis at FCC-ee - hadronic PID separation is vital
- Our analysis assumes perfect PID
- Naively investigate this by making random swaps (no momentum dependence)

$$
B^{0} \rightarrow K^{* 0} \nu \bar{\nu}
$$



$$
B_{s}^{0} \rightarrow \phi \nu \bar{\nu}
$$



- $K-\pi$ separation of $2 \sigma$ would have negligible impact on the sensitivity


## Vertexing requirements of the detector

- For serious flavour analysis at FCC-ee - precision vertexing is essential
- Our analysis assumes perfect vertex seeding
- Naively investigate this by making random swaps

$$
B^{0} \rightarrow K^{* 0} \nu \bar{\nu}
$$



$$
B_{s}^{0} \rightarrow \phi \nu \bar{\nu}
$$



- Need $<0.2 \mathrm{~mm}$ resolution to mitigate vertex mis-id
- But this is already above the requirements for vertex precision anyway


## $q^{2}$ distribution and reweighting

- Our simulation uses phase space (PHSP) generation models
- We need to reweight the $q^{2}$ distribution to match the latest theory predictions (from MR and OS)

$$
B^{0} \rightarrow K^{* 0} \nu \bar{\nu}
$$



$$
B_{s}^{0} \rightarrow \phi \nu \bar{\nu}
$$



## Summary and Outlook

- We are now at fairly advanced stages (we have an almost complete paper draft)
- Sensitivity to $b \rightarrow s \nu \bar{\nu}$ BFs of $\mathcal{O}(1 \%)$
- Need to finish the pheno interpretation (sensitivity to Wilson coeffs)
- In parallel considering neutral modes not shown here
- Precise vertexing is vital
- Average flight distance of a $B^{0}$ at FCC-ee is $\sim 3 \mathrm{~mm}$
- Our analysis assumes both the production (PV) and decay (SV) vertices of the $B$ are perfectly seeded
- Need resolution $\mathcal{O}(100-200 \mu \mathrm{~m})$ to mitigate vertex mis-id
- Powerful particle identification is required
- Sensitivity begins to rapidly degrade for separation $<2 \sigma$

Thanks to the authors of Ref. [6] for the inspiration and example codes

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## Back Up

## BACK UP

## Searches at $B$-factories

- Searches at B-factories use $B$-mesons produced via $e^{+} e^{-} \rightarrow \Upsilon(4 S) \rightarrow B^{+} B^{-}$
- Event is tagged either inclusively or using specific hadronic or semileptonic decays of the other $B$.
- Belle II results: $\mathrm{BR}\left(B^{+} \rightarrow K^{+} \nu \bar{\nu}\right)<4.1 \times 10^{-5}$ at $90 \%$ C.L. [arXiv:2104.12624].
- Expect to reach $\sim 10 \%$ precision on $B^{+} / B^{0}$ with $50 \mathrm{ab}^{-1}$ [arXiv:1808.10567]


- FCC-ee is the only foreseen experiment that can improve Belle-II measurement in the (far) future (apart from maybe CEPC)!


## Some places we cheat

## Relevant for detector design

- Use the same vertexing procedure developed for $B_{c}^{+} \rightarrow \tau^{+} \nu_{\tau}$ (see this talk for details) which assumes perfect vertex seeding
$\rightarrow$ implies we will have excellent vertex resolution
- We also truth match the kaon and pion daughters to have the correct mass hypothesis (with the reconstructed momentum)
$\rightarrow$ implies we will have excellent PID
- When we get a bit more advanced it would be nice to understand the impact of relaxing these requirements.
- Also assume the $K^{* 0}$ in the signal mode is pure $K^{*}(892)^{0}$

None of this is particularly relevant for the event level MVA we have trained so far (and show today) but it will be important for the next stage MVA

## Charged energy in each hemisphere

- More discrmination power in the minimum energy hemisphere (signal side) due to missing energy in the signal decay




## Neutral energy in each hemisphere

- More discrmination power in the minimum energy hemisphere (signal side) due to missing energy in the signal decay




## Charged multiplicity in each hemisphere

- More discrmination power in the minimum energy hemisphere (signal side) due to missing energy in the signal decay




## Neutral multiplicity in each hemisphere

- More discrmination power in the minimum energy hemisphere (signal side) due to missing energy in the signal decay




## Stage 1 Inputs

- The total reconstructed energy in each hemisphere,
- The total charged and neutral reconstructed energies of each hemisphere,
- The charged and neutral particle multiplicities in each hemisphere,
- The number of charged tracks used in the reconstruction of the primary vertex,
- The number of reconstructed vertices in the event,
- The number of candidates in the event
- The number of reconstructed vertices in each hemisphere,
- The minimum, maximum and average radial distance of all decay vertices from the primary vertex.


## Stage 1 BDT






## Stage 2 BDT





## Stage 2 Inputs

- The intermediate candidate's reconstructed mass
- The number of intermediate candidates in the event
- The candidate's flight distance and flight distance $\chi^{2}$ from the primary vertex
- The $x, y$ and $z$ components of the reconstructed candidate's momentum
- The scalar momentum of the candidate
- The transverse and longitudinal impact parameter of the candidate
- The minimum, maximum and average transverse and longitudinal impact parameters of all other reconstructed vertices in the event
- The angle between the intermediate candidate and the thrust axis
- The mass of the primary vertex
- The nominal $B$ candidate energy, defined as the $Z$ mass minus all of the reconstructed energy apart from the candidate children


## Backgrounds




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## Spline Drop Off



