

# A Logarithmically Accurate Resummation In C++

Parton Showers and Resummation 2023, 7 June 2023

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# A Logarithmically Accurate Resummation In C++

- Event simulation factorised into

- Hard Process
- Parton Shower

- Underlying event

- Hadronisation

See also talks by  
Basem El-Menoufi,  
Christian Preuss

- QED radiation

- Hadron Deepsys

See also talks by  
Silvia Ferrario Ravasio,  
Alexander Karlberg

## This Talk:

### Why?

- parton showers resum large logs  $\sim$  NLL, but open questions on actual accuracy

- starting work towards NNLL/NLO evolution → probably better resolve this first

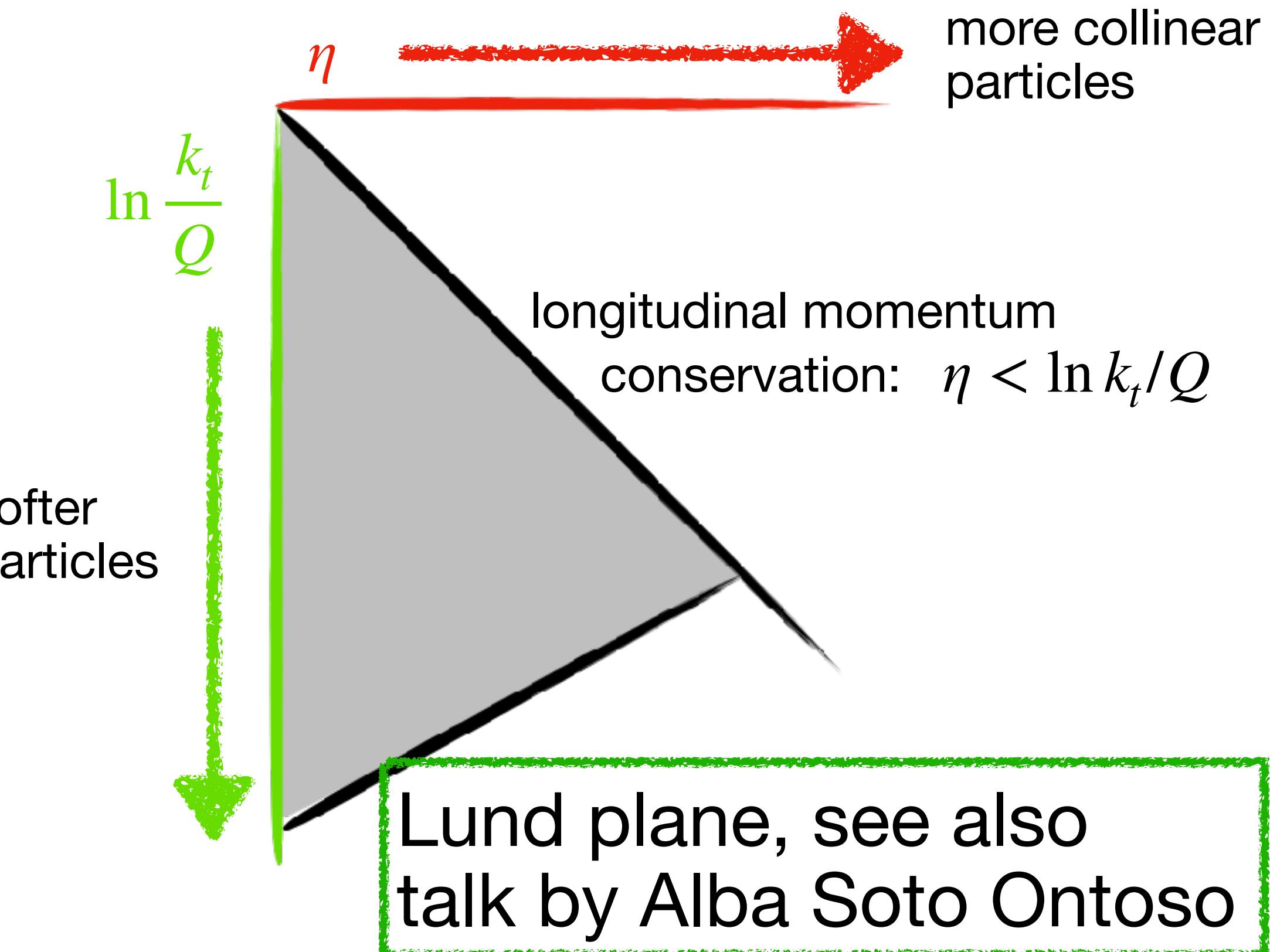
- recent formal discussion → current dipole showers need reworking

[Dasgupta, Dreyer, Hamilton, Monni, Salam '18]

# parton showers - Cliff notes version

- no-emission probability (sudakov factor)
- splitting kernels  $P(z)$  captures soft and collinear limits of matrix elements
- fill phase space ordered in evolution variable  $(k_t, \theta, q^2, \dots) \Rightarrow$  here  $k_t$  ordered shower
- generate new final state after emission according to recoil scheme

$$\sim \exp \left[ - \int_{t_0}^{t_1} \frac{dk_t}{k_t} dz \frac{\alpha_S}{2\pi} P(z) \right]$$



# splitting of Eikonal

Starting point: eikonal



$$\frac{p_i p_k}{(p_i p_j)(p_j p_k)} = \frac{1}{E_j^2} \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})} \equiv \frac{W_{ik,j}}{E_j^2}$$

naive implementation leads to soft double counting need to split into  $ij$  and  $kj$  collinear terms

[Marchesini, Webber '88]

Option 1:

$$W_{ik,j} = \tilde{W}_{ik,j}^i + \tilde{W}_{ki,j}^k , \quad \text{where} \quad \tilde{W}_{ik,j}^i = \frac{1}{2} \left( \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(1 - \cos \theta_{jk})} + \frac{1}{1 - \cos \theta_{ij}} - \frac{1}{1 - \cos \theta_{jk}} \right)$$

- e.g. Angular ordered shower, downside: problems with NGLs

Option 2: follow [Catani, Seymour '97]

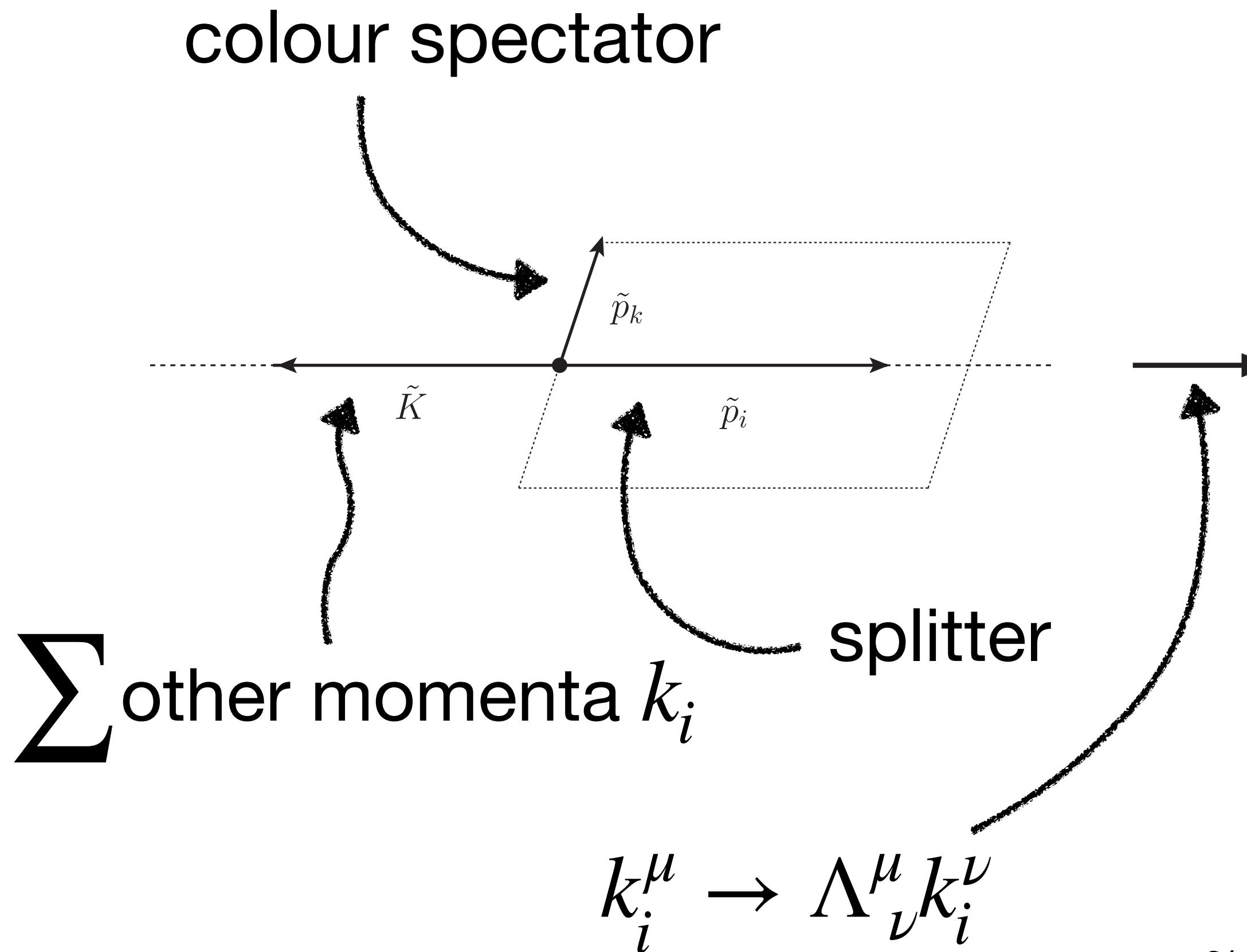
$$W_{ik,j} = \bar{W}_{ik,j}^i + \bar{W}_{ki,j}^k , \quad \text{where} \quad \bar{W}_{ik,j}^i = \frac{1 - \cos \theta_{ik}}{(1 - \cos \theta_{ij})(2 - \cos \theta_{ij} - \cos \theta_{jk})}$$

- full phase space coverage, splitting functions remain positive definite

Note related ideas in [Forshaw, Holguin, Plätzer '20]

# kinematics - global recoil scheme

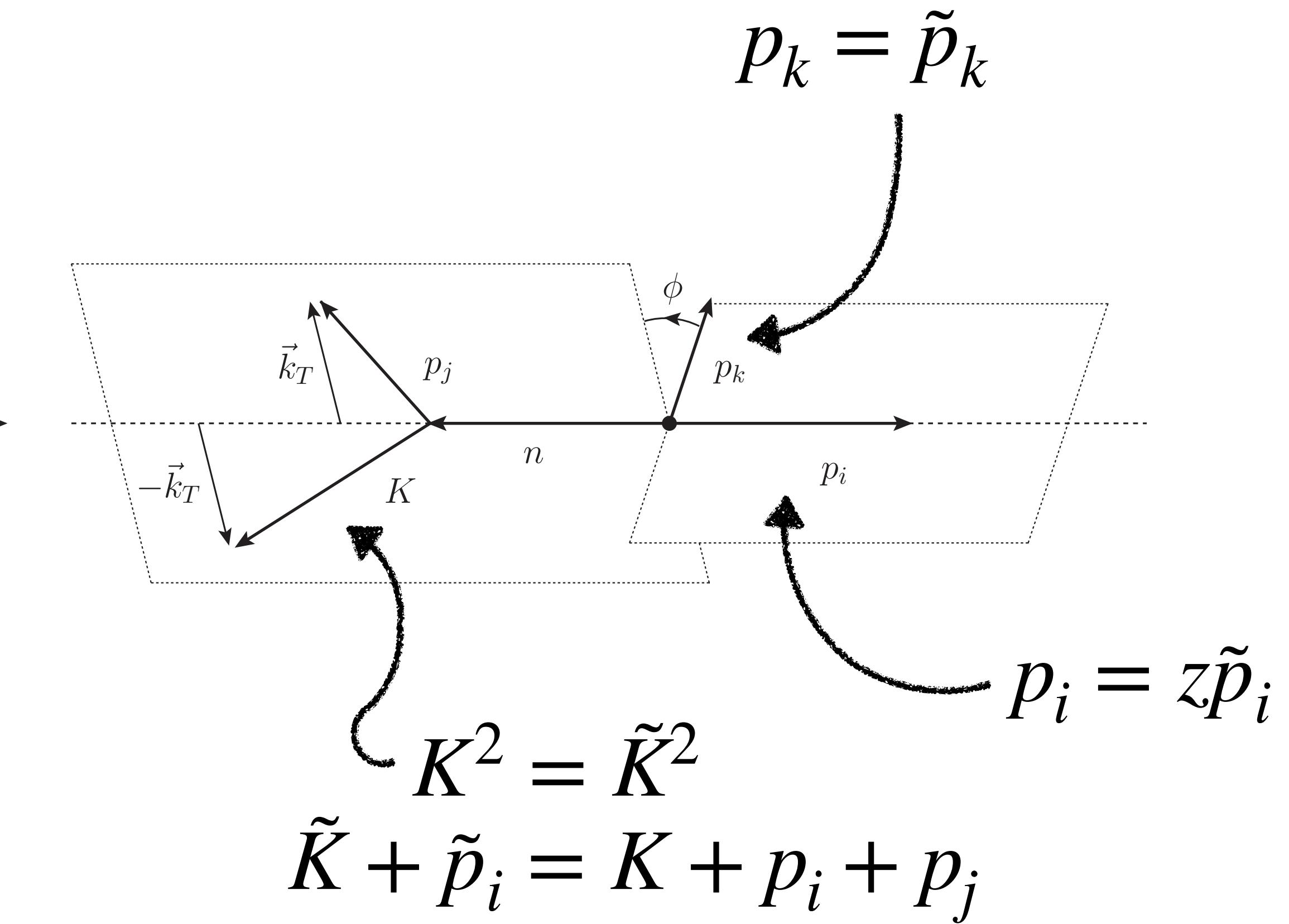
- Before splitting:



[Catani, Seymour '97]

$$\Lambda_\nu^\mu = g_\nu^\mu - \frac{(K + \tilde{K})^\mu (K + \tilde{K})_\nu}{K \cdot \tilde{K} + \tilde{K}^2} + 2 \frac{K^\mu \tilde{K}_\nu}{\tilde{K}^2} \rightarrow \Lambda_\nu^\mu \tilde{K}^\nu = K^\mu$$

- After splitting:



# effect of recoil on accuracy - multiple emissions

- QCD coherence  $\rightarrow$  factorised emissions
- observables dependence correlated  $\rightarrow$  how to extract NLL without additional information?
- method from [Banfi, Salam, Zanderighi '05]: need explicit soft-collinear limit\*:

$$k_t^\rho = k_t \rho$$

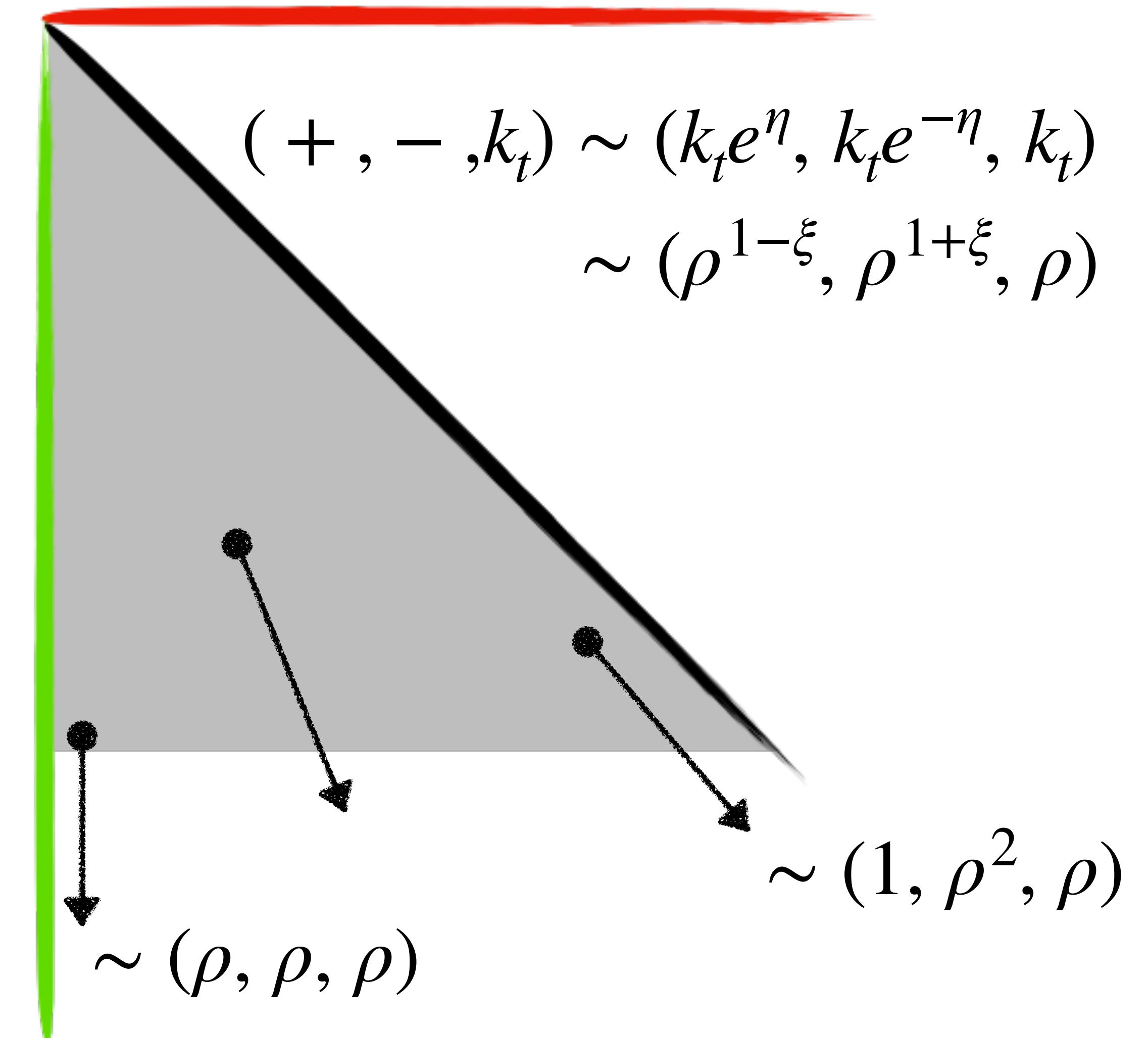
$$\eta^\rho = \eta - \xi \ln \rho$$

and assume

$$V(k_i^\rho) = \rho V(k_i)$$

$$\xi = \frac{\eta}{\eta_{\max}}$$

$\rightarrow$  numerically evaluate integrals in this limit



\* again assume  $V(k_t, \eta) \sim k_t/Q$  for brevity

# effect of recoil on accuracy

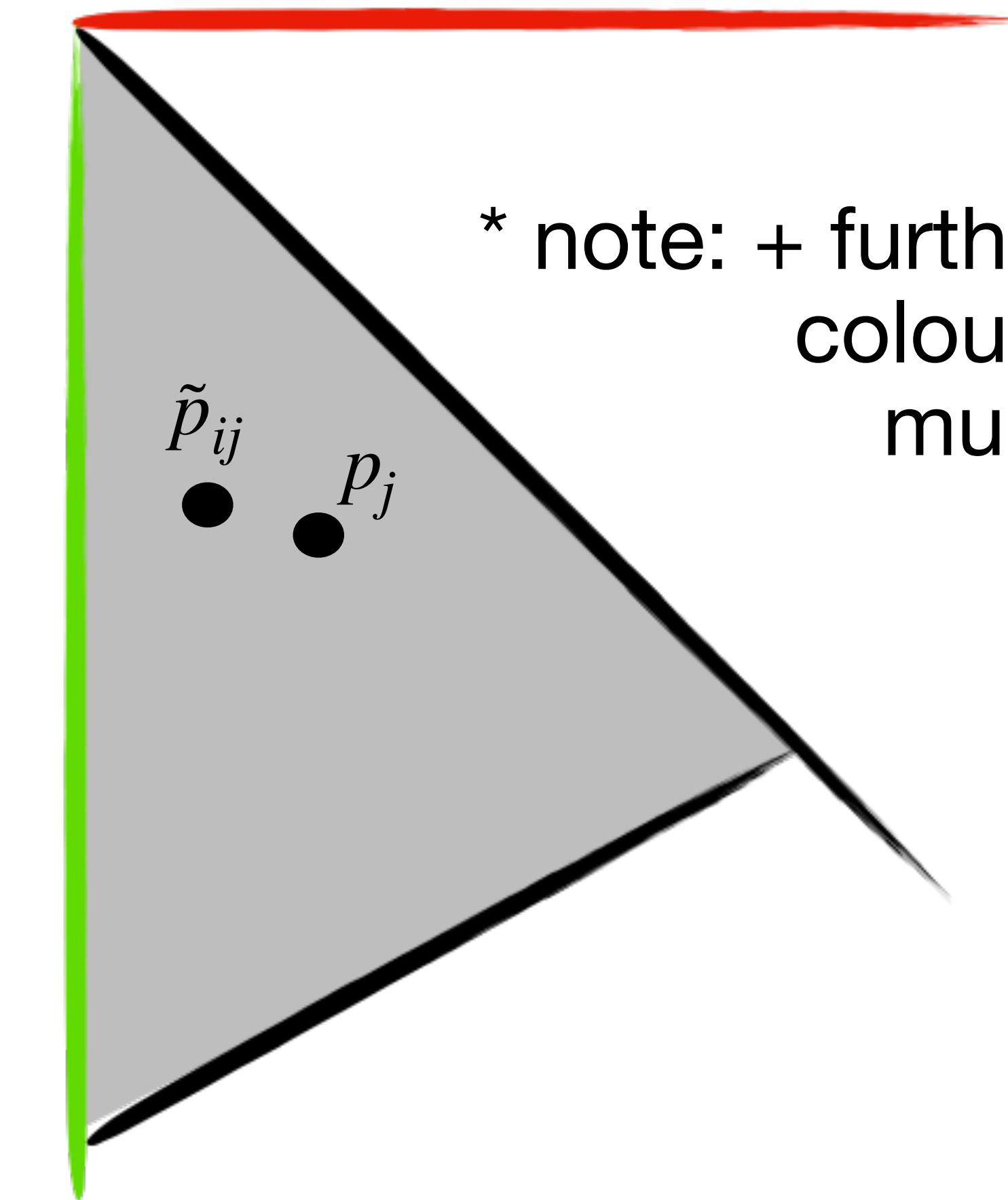
- question: do recoil effects indeed vanish in soft limit (i.e.  $\rho \rightarrow 0$ )?\*

[Dasgupta,Dreyer,Hamilton,Monni,Salam '18]

- consider situation where we first emit  $\tilde{p}_{ij}$  from  $p_a, p_b$ , then emit  $p_j$ ,  
 $\tilde{p}_{ij} \rightarrow p_i, p_j$
- transverse momentum of  $p_i$  will be  
 $\sim k_t^{ij} + k_t^j$

$$\Rightarrow \frac{\Delta k_t^{ij}}{k_t^{ij}} \rightarrow \frac{\rho k_t^j}{\rho k_t^{ij}} = \mathcal{O}(1)$$

$$p_i = z\tilde{p}_{ij} + (1 - z)y\tilde{p}_k + k_\perp$$
$$p_j = (1 - z)\tilde{p}_{ij} + zy\tilde{p}_k - k_\perp$$
$$p_k = (1 - y)\tilde{p}_k .$$



\* note: + further problems for colour assignment in multiple emissions

# analytic proof of accuracy

$$\Lambda_{\nu}^{\mu}(K, \tilde{K}) = g_{\nu}^{\mu} + \tilde{K}^{\mu} A_{\nu} + X^{\mu} B_{\nu} \xrightarrow{\text{vanishes in soft limit}}$$

work out  $\rho \rightarrow 0$  limit:  $A^{\nu} \xrightarrow{\rho \rightarrow 0} 2 \frac{\tilde{K}X}{\tilde{K}^2} \frac{\tilde{K}^{\nu}}{\tilde{K}^2} - \frac{X^{\nu}}{\tilde{K}^2}$ , and  $B^{\nu} \xrightarrow{\rho \rightarrow 0} \frac{\tilde{K}^{\nu}}{\tilde{K}^2}$

apply to soft momentum  $p_l$ :

$$\frac{\Delta p_l^{0,3}}{p_l^{0,3}} \sim \rho^{1-\max(\xi_i, \xi_j)}$$

$$\frac{\Delta p_l^{1,2}}{p_l^{1,2}} \sim \rho^{(1-\xi_l)(\max(\xi_i, \xi_j) - \xi_l)}$$

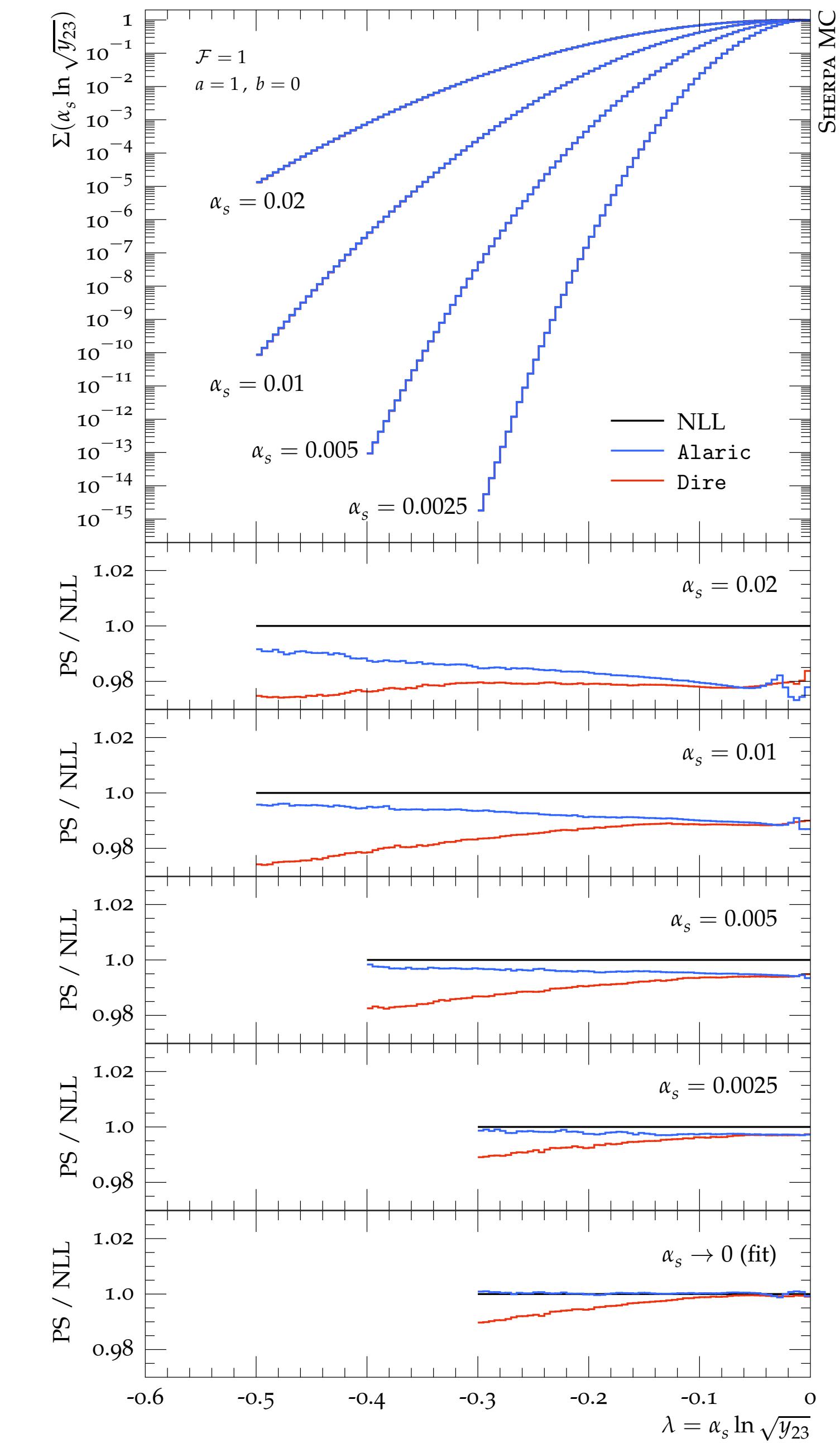
compare to  $\frac{\Delta k_t}{k_t} \sim \mathcal{O}(1)$  from local dipole scheme

# numerical validation I

- Limit  $\alpha_s \rightarrow 0$  with  $\lambda = \alpha_s L = \text{const.}$  of

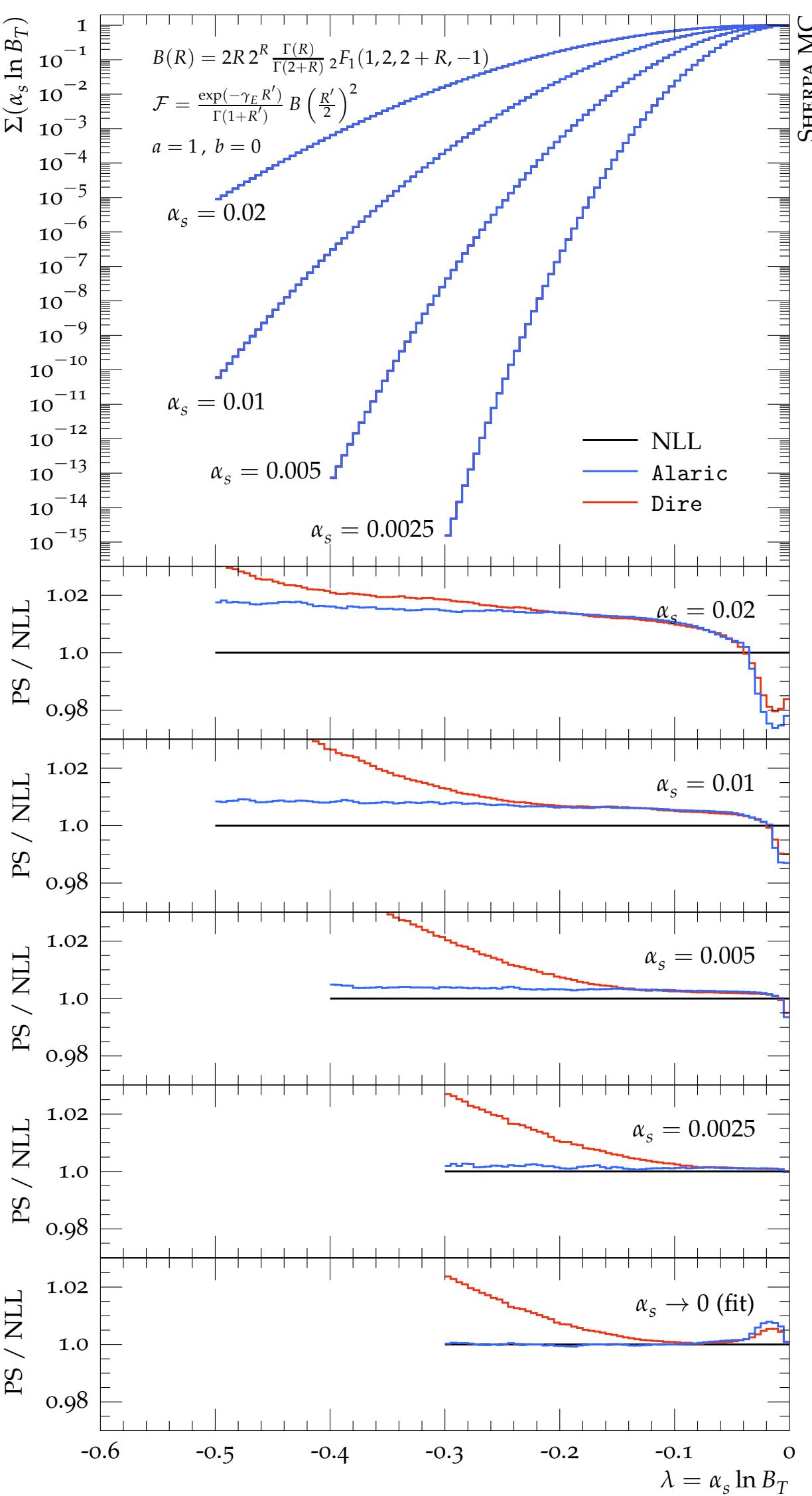
$$\begin{aligned} \frac{\sum \text{Shower}}{\sum \text{NLL}} &\sim \exp \left( f_{\text{Shower}}^{LL} - L g_1(\alpha_s^n L^n) \right) \\ &\times \exp \left( f_{\text{Shower}}^{NLL} - g_2(\alpha_s^n L^n) \right) \\ &\times \exp \left( \mathcal{O}(\alpha_s^{n+1} L^n) \right) \\ \rightarrow 1 & \quad \text{if shower reproduces LL, NLL logs} \end{aligned}$$

- Observable: jet resolution  $y_{23}$  in Cambridge jet measure,  $\mathcal{F} = 1 \rightarrow$  only largest emission matters, check that additional shower emissions vanish

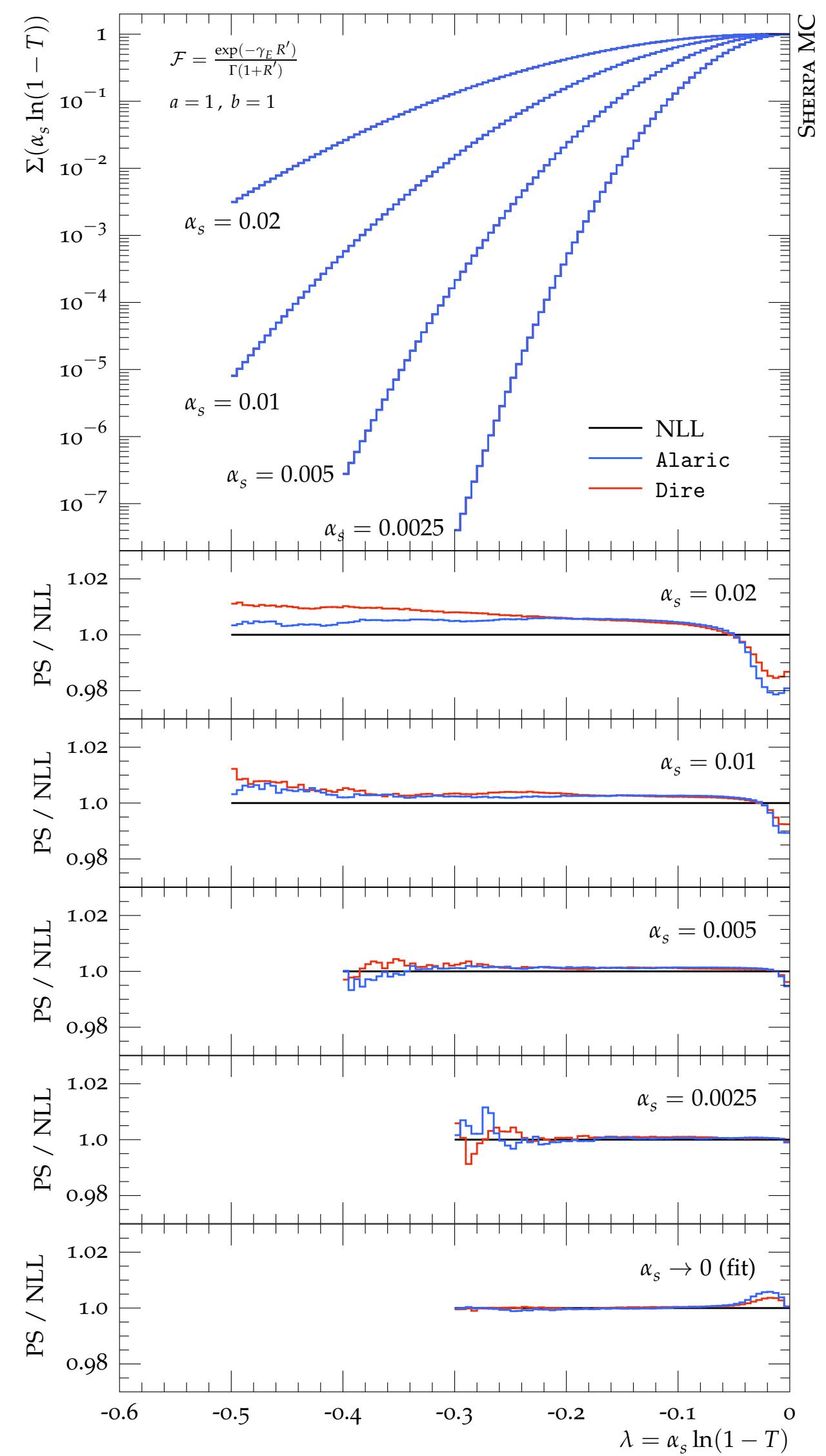


# numerical validation II

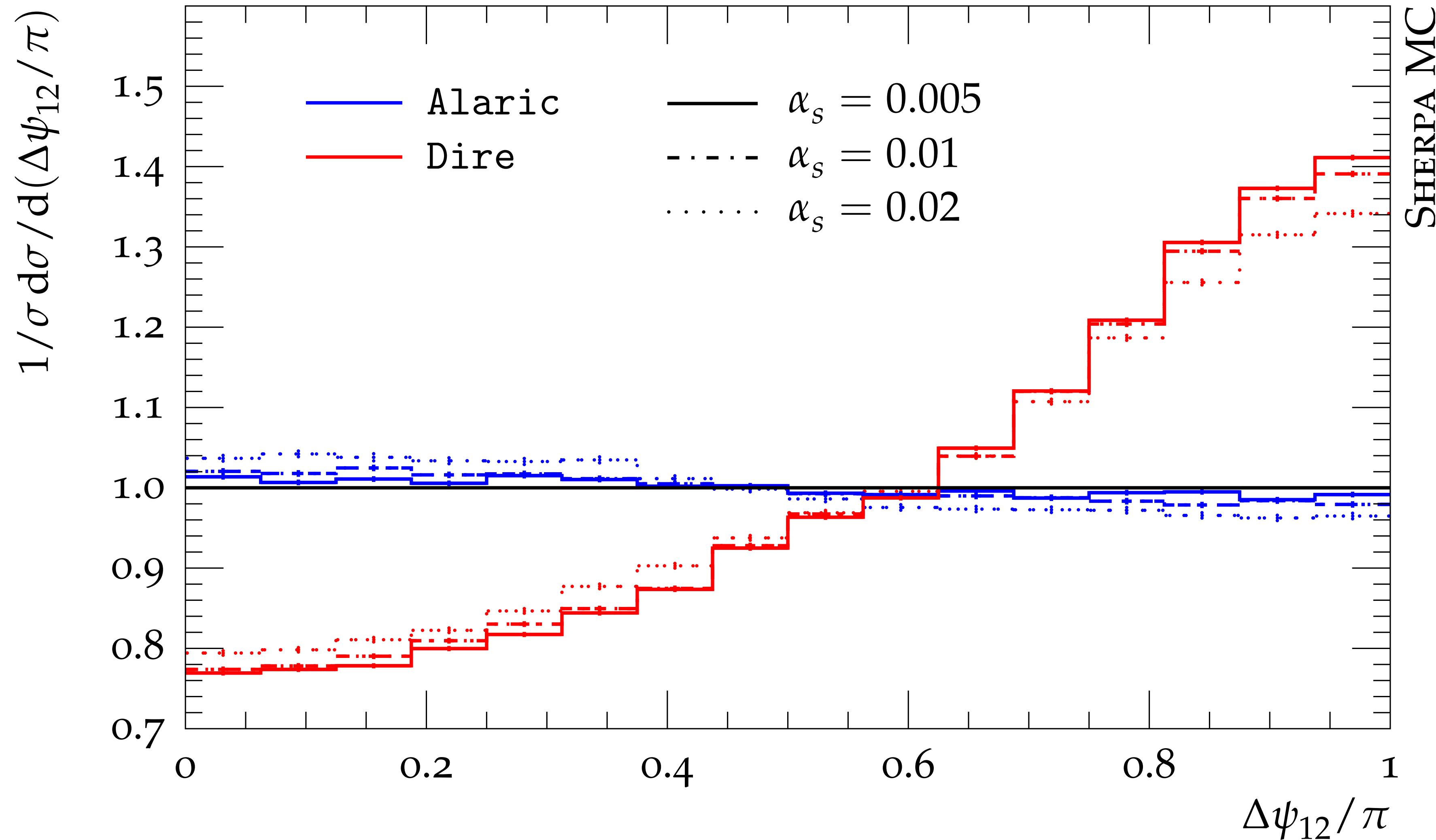
- total broadening  $B_T = B_L + B_R$
- scaling  $k_t$  like, similar to  $y_{23}$
- but non-trivial  $\mathcal{F}$  function



- thrust  $\tau = 1 - t$
- scaling like virtuality  $k_t e^{-\eta}$
- standard function  $\mathcal{F} = \frac{\exp(-\gamma_E R)}{\Gamma(1 + R')}$
- no evidence for NLL violation even for standard showers



# numerical validation III

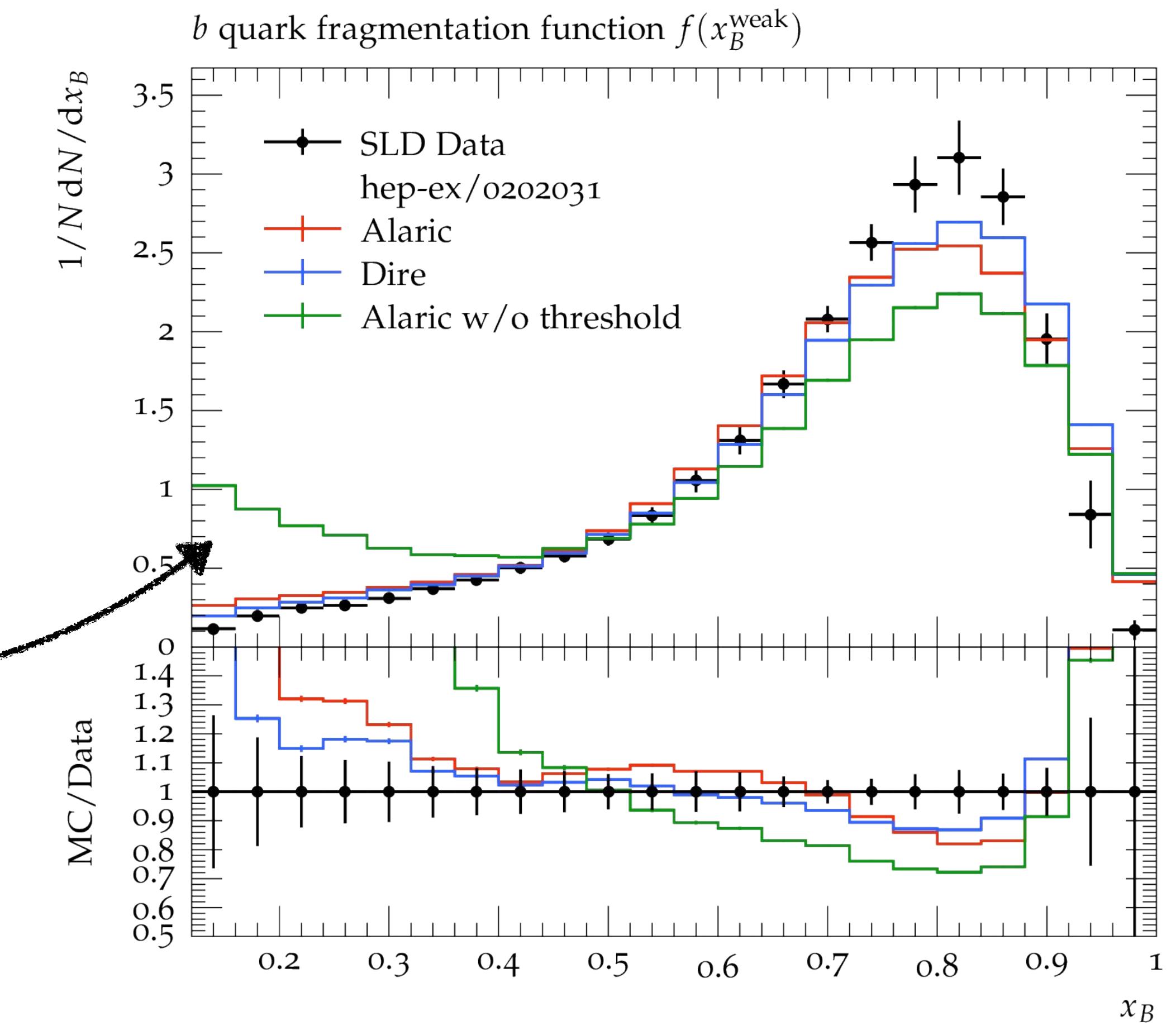


SHERPA MC

See also talks by  
Gregory Soyez

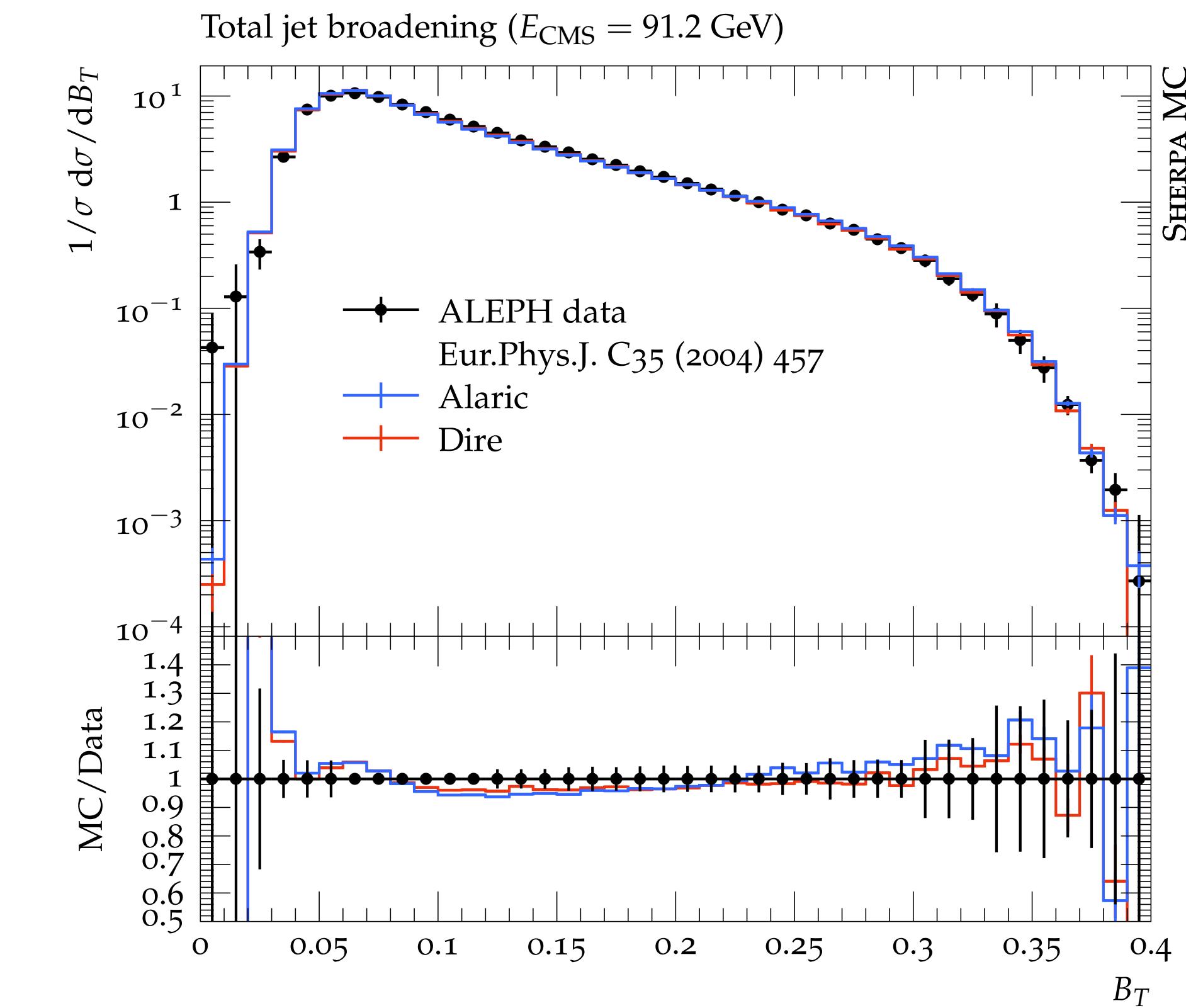
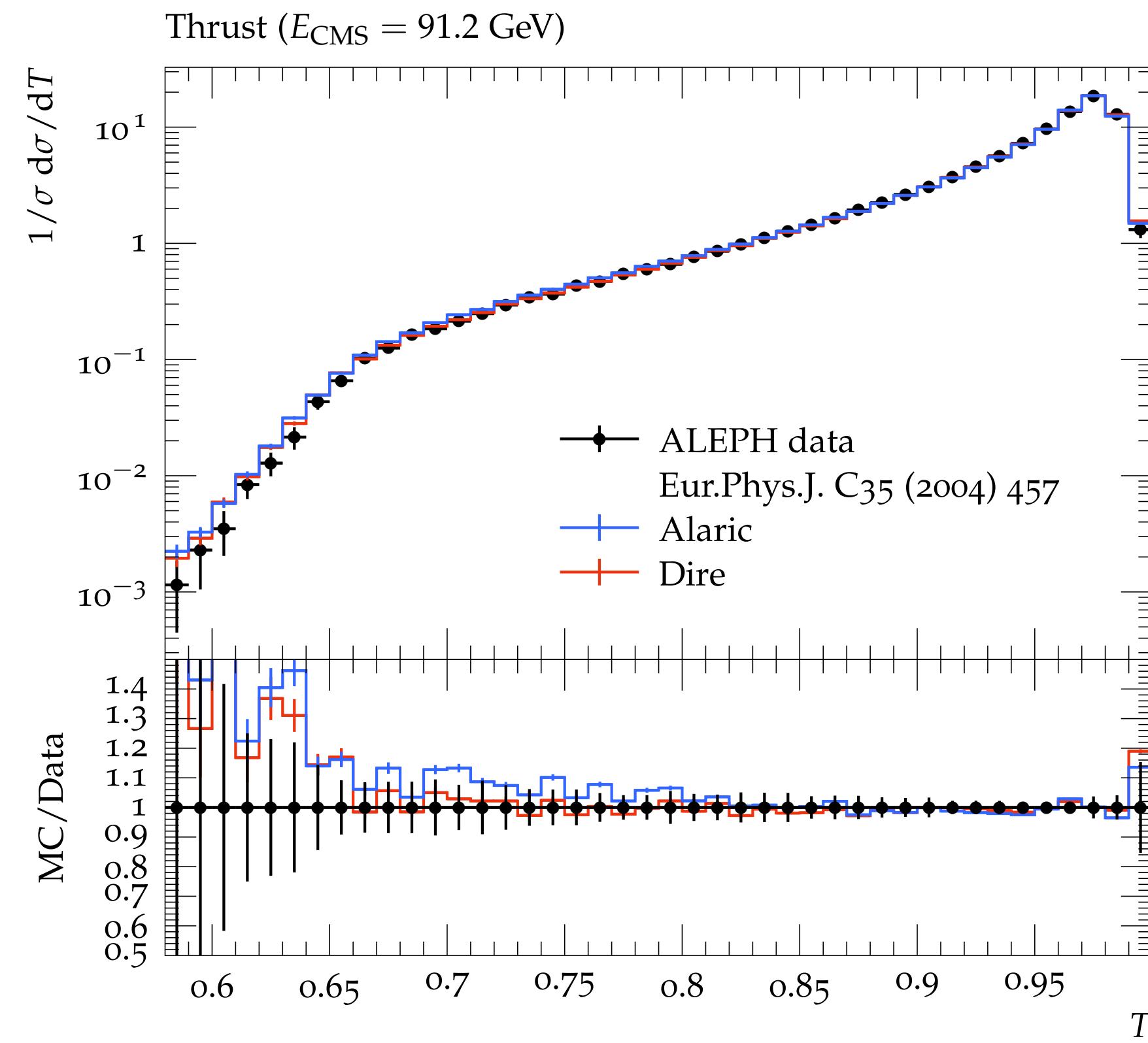
# pheno, details and b fragmentation

- first caveat: no quark masses implemented yet
- problem for cluster hadronisation → use Lund model via Pythia
- + need flavour threshold for  $g \rightarrow b\bar{b}/g \rightarrow c\bar{c}$  splittings
- Dire parton shower as implemented in Sherpa as reference, Lund model tuned for Alaric  $\sigma = 0.3 \text{ GeV}, a = 0.4, b = 0.36 \text{ GeV}^{-2}$  and for Dire  $\sigma = 0.3 \text{ GeV}, a = 0.4, b = 0.46 \text{ GeV}^{-2}$



$$x_B \sim \frac{E_{B-\text{Hadron}}}{E_{\text{tot.}}/2}$$

# pheno, LEP observables



Thrust:

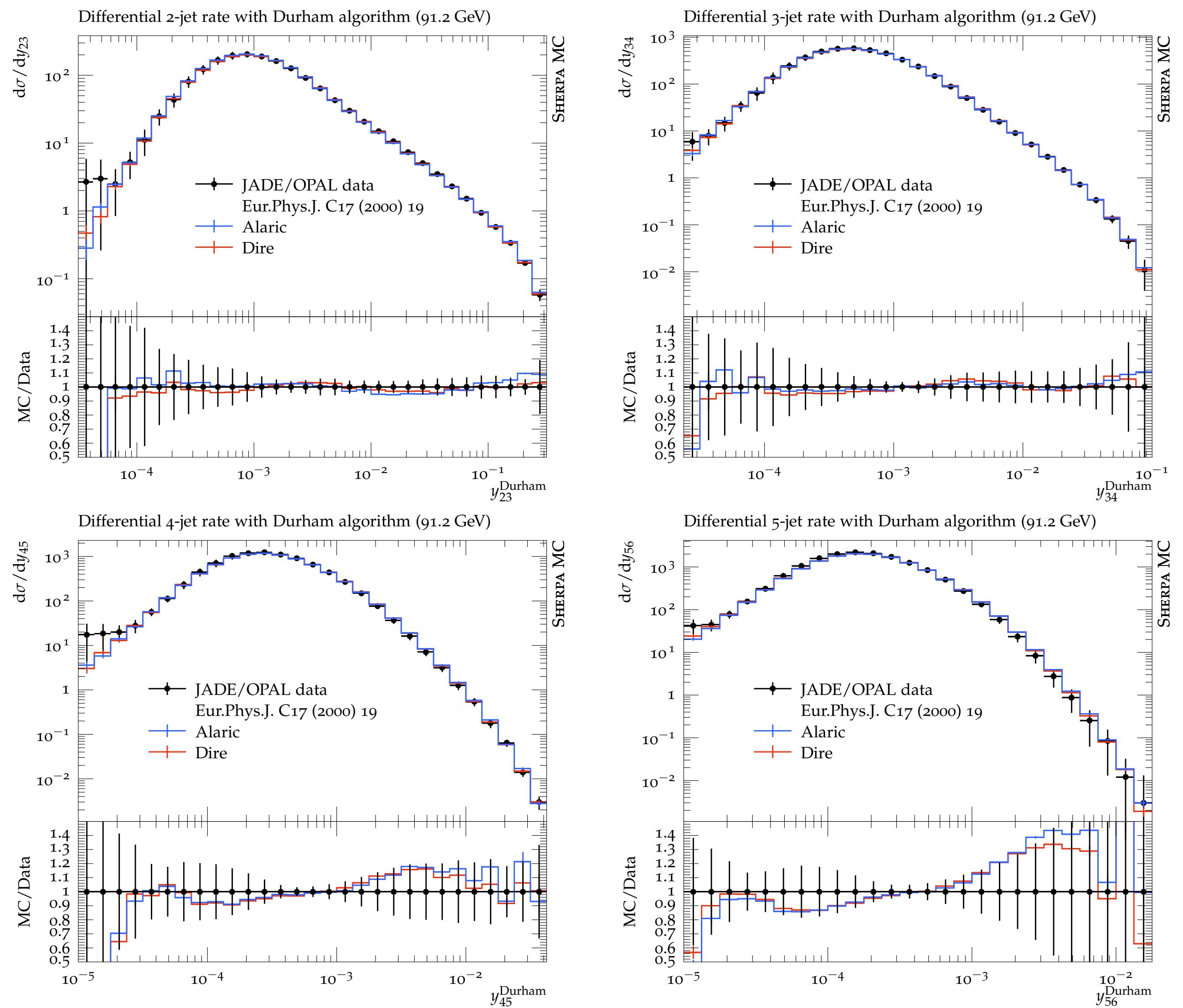
- Note this is  $T$ , not  $1-T$ : soft physics is to the right
- Note there is no matching, relevant for small  $T$

Total Broadening:

- soft physics is left hand side
- some deviations from data, but similar to Dire

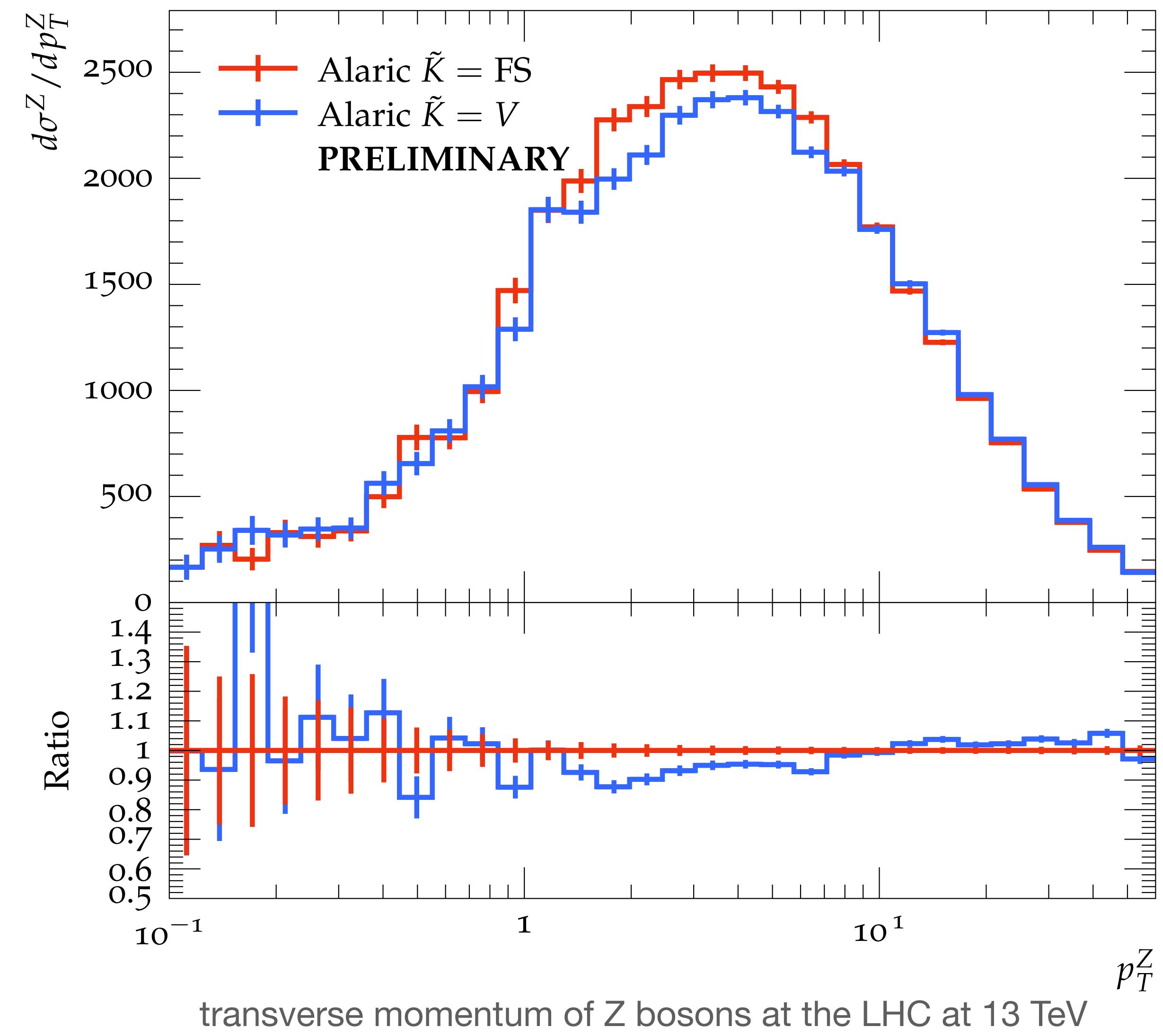
# pheno, LEP observables

- Durham resolution scales  
 $y_{n,n+1} \sim k_t^2/Q^2$
- higher Born multiplicities → sensitivity to multiple emissions increased
- again, note no matching/merging involved



# Alaric initial state shower (outlook)

- Formalism presented in [Herren, Höche, Krauss, DR, Schönherr '22] general and applicable to initial state evolution
- practical considerations:
  - precise definition of evolution variable
  - PDFs, clear in principle, but more choices to make
  - distribution of recoil (i.e. definition of  $\tilde{K}$ )
    - e.g. Drell-Yan process, could be EW boson, or full final state (or ...?)



# A Logarithmically Accurate Resummation In C++

- NLL resummation in CAESAR formalism as definition and validation of parton shower accuracy
- New parton shower Alaric
  - partial fractioning of eikonal → positive definite splitting function with full phase space coverage
  - global kinematics scheme enables analytic proof of NLL accuracy + numerical validation
  - included in Sherpa framework and first pheno results