Machine learning for lattice field theory and back

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Introduction

- \circ past five years or so has seen a rapid rise of applications of machine learning (ML) in fundamental science, particle physics, theoretical physics
- o of course ML has been around for quite some time, especially in experimental particle physics
- o nevertheless, there is an **exponential** increase in activity

year

o find title **learning** on the Inspire data base (highenergy physics)

o exponential growth!

a flavour of ideas and results in theoretical statistical physics biased towards own work and interests in lattice field theory

o classification: order-disorder transition (by now classic application)

o inverse renormalisation group (new application and concepts)

o transfer learning (widely used notion)

o outlook

Based on the following papers:

fruitful collaboration with Dimitrios Bachtis and Biagio Lucini

 \checkmark Extending machine learning classification capabilities with histogram reweighting Phys. Rev. E 102 (2020) 033303 [2004.14341 [cond-mat.stat-mech]]

 \checkmark Mapping distinct phase transitions to a neural network Phys. Rev. E 102 (2020) 053306 [2007.00355 [cond-mat.stat-mech]]

 \checkmark Adding machine learning within Hamiltonians: Renormalization group transformations, symmetry breaking and restoration Phys. Rev. Res. 3 (2021) 013134 [2010.00054 [hep-lat]]

 \checkmark Quantum field-theoretic machine learning Phys. Rev. D 103 (2021) 074510 [2102.09449 [hep-lat]]

 \checkmark Inverse renormalisation group in quantum field theory Phys. Rev. Lett. 128 (2022) 081603 [2107.00466 [hep-lat]] with Dimitrios Bachtis and Biagio Lucini and Francesco di Renzo

Classification of phases of matter

Published: 13 February 2017

Machine learning phases of matter

Juan Carrasquilla ⊠ & Roger G. Melko

Nature Physics 13, 431-434(2017) Cite this article

> 1100 citations since 2017

Classification of phases of matter

- o matter can exist in different phases
- o prototype: 2d Ising model -> ordered/disordered or cold/hot phases
- o task: determine phase a configuration is in, determine critical coupling or temperature

Ordered -- ? -- Disordered

ML excels in pattern finding

- supervised learning problem: use sets of configurations deep in the ordered and in the disordered phase
- input: configurations $\langle -2 \rangle$ $\langle -2 \rangle$ output: ordered/disordered
- "train the ML algorithm", i.e. adjust parameters in the neural network so that it reproduces the correct classification for the training set

• new, unseen configurations -- > determine **probability** to be (dis)ordered

2d Ising model

$$
\circ \, Z = \text{Tr} \, e^{-\beta E} \quad \text{with } E = -\sum_{\langle i j \rangle} s_i s_j \qquad (s_i = \pm 1)
$$

 \circ critical coupling or inverse temperature β_c

 \circ correlation length ξ , magnetic susceptibility χ diverge at transition

 \circ critical exponents $\xi \sim |t|^{-\nu}$ $\chi \sim |t|^{-\gamma}$ reduced temperature

$$
\nu = 1
$$
, $\gamma/\nu = 7/4$, $\beta_c = \frac{1}{2} \ln(1 + \sqrt{2}) \approx 0.440687$

 \circ finite-size scaling $|t| = \left|\frac{\beta_c - \beta_c(L)}{\beta_c}\right| \sim \xi^{-\frac{1}{\nu}} \sim L^{-\frac{1}{\nu}} \sim L^{-\frac{1}{\nu}} \sim \chi \sim L^{\gamma/\nu}$

Probability to be in ordered phase

train NN away from the phase transition: $\beta \leq 0.41$ and $\beta \geq 0.47$

investigate unseen configurations at intermediate β on lattices of different sizes

probability behaves as an approximate order parameter

 \checkmark Extending machine learning classification capabilities with histogram reweighting Phys. Rev. E 102 (2020) 033303 [2004.14341 [cond-mat.stat-mech]]

What's new?

- by now well-established procedure, what can we add?
- interpret output from a NN as an observable in a statistical system
- input: configurations, distributed according to Boltzmann weight
- output: observable, "order parameter" in statistical system

Output of NN as physical observable

- opens up possibility to use "standard" numerical/statistical methods
	- histogram reweighting: extrapolation to other parameter values
- starting from computation at given β_0 : extrapolate to other β values

$$
\langle P \rangle (\beta) = \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum e^{-(\beta - \beta_0)E_i}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.3 \ 0.4 \ 0.4355 \ 0.4365}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4 \ 0.4415 \ 0.4425}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4485}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4485}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4485}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3 \ 0.4415 \ 0.4445 \ 0.4445}} \times \frac{\sum P_i e^{-(\beta - \beta_0)E_i}}{\sum_{\substack{0.3
$$

 \mathcal{L} . The contract of the contract of \mathcal{L}

Critical behaviour from NN observables

• Determine L dependent susceptibility δP and its maximum at $\beta_c(L)$

Extract critical properties from NN observables only

Intermediate summary

- \checkmark proposed to identify NN outputs as observables in statistical physics
- \checkmark introduced histogram reweighting to employ in supervised machine learning
- ü critical properties obtained from a finite-size scaling analysis using quantities derived from NN alone (no need for explicit order parameter, knowledge of symmetries)
- \checkmark quantitative studies of phase transitions based on a synergistic relation between machine learning and statistical mechanics

Transfer learning with histogram reweighting

- Ø NN has learned patterns, or *features*, in 2d Ising model
- \triangleright are these sufficiently universal to predict the structure of phase transitions in other systems?
- \triangleright what about universality class, order of transition, type of degrees of freedom?

 \triangleright apply to q-state Potts model (with $q = 3, ..., 7$), φ^4 scalar field theory

 \checkmark Mapping distinct phase transitions to a neural network

Phys. Rev. E 102 (2020) 053306 [2007.00355 [cond-mat.stat-mech]]

Transfer learning: q-state Potts model

- training on Ising model, not Potts model
- \circ continuous lines using histogram reweighting
- \circ vertical dashed lines indicate expected transition at $\beta_c = \ln (1 + \sqrt{q})$
- $q = 3, 4: 2^{nd}$ order transition, $q = 5, 6, 7: 1^{st}$ order transition

 φ^4 scalar field theory

- reweight in mass parameter, μ^2
- identify regions where phase is clear
- retrain NN using $\mu^2 < -1.0$ and $\mu^2 > -0.9$
- repeat finite-size scaling analysis as in 2d Ising model

$$
\frac{\mu_c^2}{\text{CNN+Reweighting}} \quad \frac{\mu_c^2}{-0.95225(54)} \quad \frac{\nu}{0.99(34)} \quad \frac{\gamma/\nu}{1.78(7)}
$$

- same universality class as 2d Ising model
- critical mass in agreement with results obtained with standard methods (Binder cumulant, susceptibility)

Under the hood: activation functions in NN

mean activation functions in the 64 neurons in the fully connected (FC1) layer of 2d Ising-trained neural network, for:

- o 2d Ising model
- $q = 3$ and $q = 5$ Potts model
- φ φ ⁴ scalar field theory

universal features distinguish ordered and disordered phases, irrespective of e.g. order of transition

Intermediate summary: transfer learning

- \checkmark enables the use of simplistic systems to study complicated models with partially known behaviour
- \checkmark combine with reweighting to scan parameter space and reconstruct effective order parameters
- \checkmark locate (unknown) phase transitions
- \checkmark given this knowledge, train a new NN on configurations of the target system labelled according to previous step
- \checkmark study infinite-volume limit in this new NN to make accurate predictions

Renormalisation Group (RG)

- o standard renormalisation group: coarse-graining, blocking transformation, integrating out degrees of freedom, …
- o Ising model: Kadanoff block spin
- o majority rule
- o reduction of degrees of freedom
- o study critical scaling
- o not invertible: semi-group

Renormalisation group

- o generates flow in parameter space
- o due to repeated blocking: run out of degrees of freedom
- \circ need to start with large system to apply RG step multiple times
- o large systems, close to a transition, suffer from critical slowing down

Inverse renormalisation group

- what if we could invert the RG ?
- o add degrees of freedom, fill in the 'details'
- o inverse flow in parameter space
- \circ can be applied arbitrary number of steps
- o evade critical slowing down

for Ising model: Inverse Monte Carlo Renormalization Group Transformations for Critical Phenomena, D. Ron, R. Swendsen, A. Brandt, Phys. Rev. Lett. 89, 275701 (2002)

How to devise an inverse transformation?

- new degrees of freedom should be introduced
- learn a set of transformations (*transposed convolutions*) to invert a standard RG step
- minimise difference between original and constructed configuration

Inverse renormalisation group

Transposed convolutions

- local transformation
- apply inverse transformations iteratively
- evade critical slowing down
- generate flow in parameter space
- invariance at critical point

Application to φ^4 scalar field theory

- § repeated steps
- locking in on critical point

TABLE I. Values of the critical exponents γ/ν and β/ν . The original system has lattice size $L = 32$ in each dimension and its action has coupling constants $\mu_L^2 = -0.9515$, $\lambda_L = 0.7$, and $\kappa_L = 1$. The rescaled systems are obtained through inverse renormalization group transformations.

L_i/L_i				32/64 32/128 32/256 32/512 64/128 64/256 64/512 128/256 128/512 256/512	
γ/ν β/ν				$1.735(5)$ $1.738(5)$ $1.741(5)$ $1.742(5)$ $1.742(5)$ $1.744(5)$ $1.744(5)$ $1.745(5)$ $1.745(5)$ $1.746(5)$ $0.132(2)$ $0.130(2)$ $0.128(2)$ $0.128(2)$ $0.128(2)$ $0.127(2)$ $0.127(2)$ $0.126(2)$ $0.126(2)$ $0.126(2)$	

GA, Bachtis, Lucini, di Renzo PRL 128 (2022) 081603

Application to φ^4 scalar field theory

- \circ start with lattice of size 32² and apply IRG steps repeatedly
- $32^2 \rightarrow 64^2 \rightarrow 128^2 \rightarrow 256^2 \rightarrow 512^2$
- o IRG flow towards critical point
- o extract critical exponents γ / υ and β / υ from comparison between two volumes
- \circ constructed a large (512²) lattice very close to criticality without critical slowing down

Intermediate summary: inverse RG

 \checkmark flow to critical point without critical slowing down

 \checkmark reach large lattices from easy-to-simulate lattice sizes

 \checkmark relies on 'reliable' blocking step (nontrivial: scalar field majority rule is new)

 \checkmark new concept for continuous field theories

Outlook

 \checkmark machine learning has seen major boost in physical sciences

 \checkmark largely underexplored in statistical/lattice field theory

 \checkmark new concepts introduced

 \checkmark more progress can be found here:

There is great potential to apply machine learning in the area of numerical lattice quantum field theory, but full exploitation of that potential will require new strategies. In this white paper for the Snowmass community planning process, we discuss the unique requirements of machine learning for lattice quantum field theory research and outline what is needed to enable exploration and deployment of this approach in the future.