

DEVELOPMENTS ON XSUITE BEAM-BEAM SIMULATIONS FOR FCC-EE

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Special thanks to:

F. Carlier, M. Hofer, D. Shatilov, K. Oide, D. Zhou

FCCIS Workshop

6th December 2022

Overview

1. xsuite simulation setup for simplified tracking
 - First results
2. FMA
 - FCCee tune footprints
 - Comparison of various codes
 - Slicing algorithms
3. Benchmark against analytical formula ξ
 - Studies of hourglass effect with reduced bunch intensity
4. Strong-strong simulations
 - Tune scan
 - 3D flip-flop
5. Summary

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Simplified tracking simulations with xsuite

- Exploit superperiodicity of machine (2 IP case)

- In code:

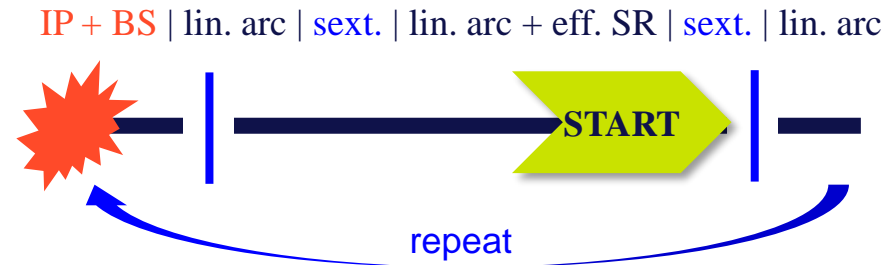
- 1 IP + tracking over half arc with linear transfer matrix

- Arc split into 3 segments

- 2 crab sextupoles between arc segments ($\beta_x=3$ m, $\beta_y=5000$ m)

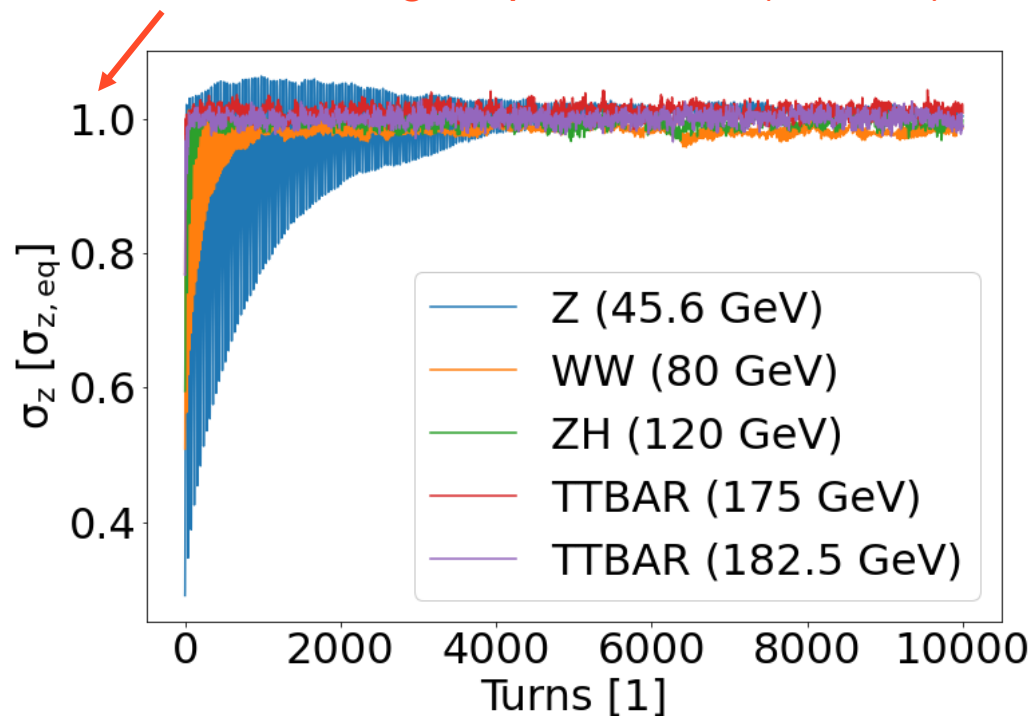
- A «turn» begins in front of the right sextupole:
 - Observation point for coordinates

- Effective radiation (damping+noise) in arc, beamstrahlung in beam-beam
 - No radiation for FMA



Equilibrium bunch length

normalized to design report values (SR+BS)

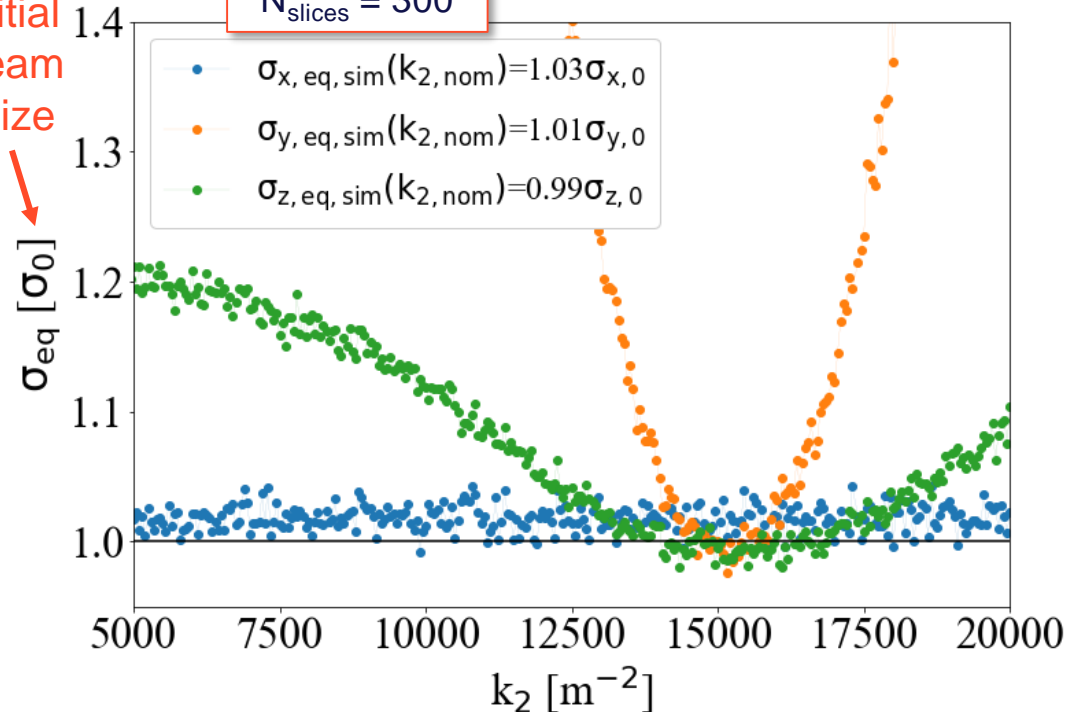


- **Weak-strong** model (1e4 particles)
- Equilibrium bunch length agrees with design report value for all resonances

Crab waist & transverse blowup

FCC-ee Z
 $N_{\text{turns}} = 3e4$
 $N_{\text{particles}} = 1e4$
 $N_{\text{slices}} = 300$

initial
beam
size



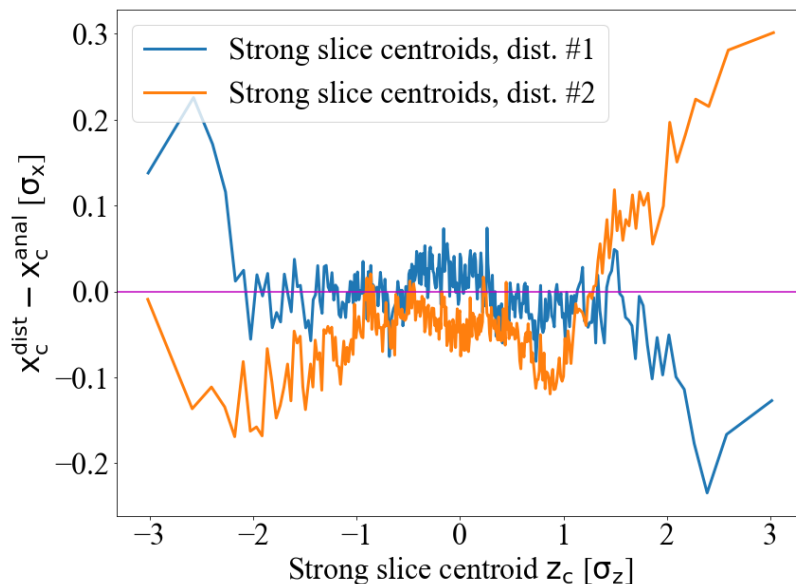
- **Weak-strong** model
- Optimum k_2 close to nominal value ($\sim 0.97 \cdot k_{2, \text{nom}}$ for Z resonance)
- No transverse blowup in optimum setting

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FMA – FCCee Z tune footprint

- In the **weak-strong** model, the strong bunch slice moments can be computed from:
 - Analytical Gaussian
 - Distribution of macroparticles \longrightarrow slice centroid coordinates are noisy
 - Random offset affects tune shift \longrightarrow noise on footprint



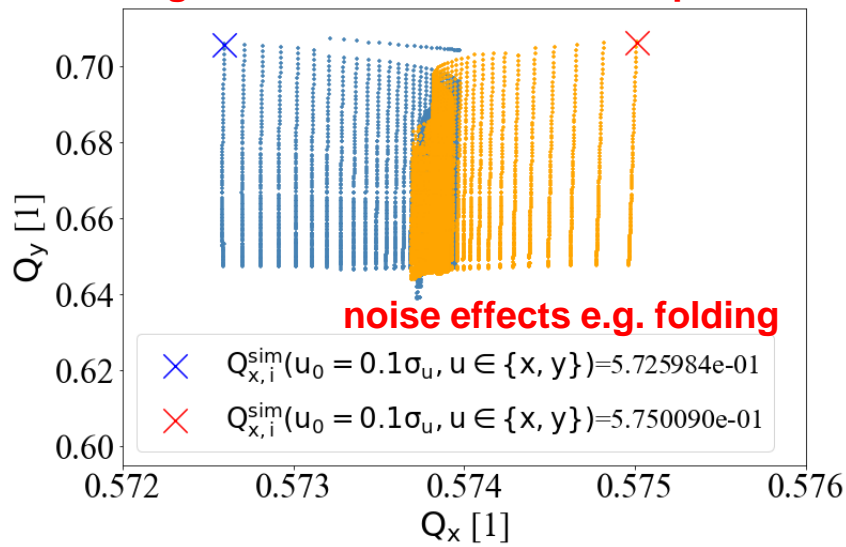
1e6 macroparticles
301 slices (uniform charge)
3333 macroparticles / slice

**compared to slices from analytical
Gaussian (Hirata's slicing)**

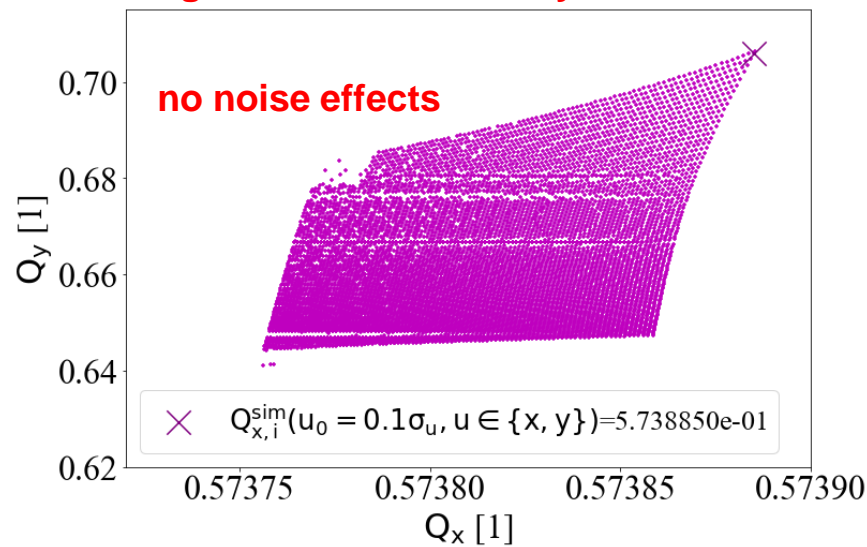
FMA – FCCee Z tune footprint

- Random offset affects tune shift noise on footprint
- Better to use moments from slicing an analytical Gaussian to avoid noise effects
 - Possible in WS but not in SS!

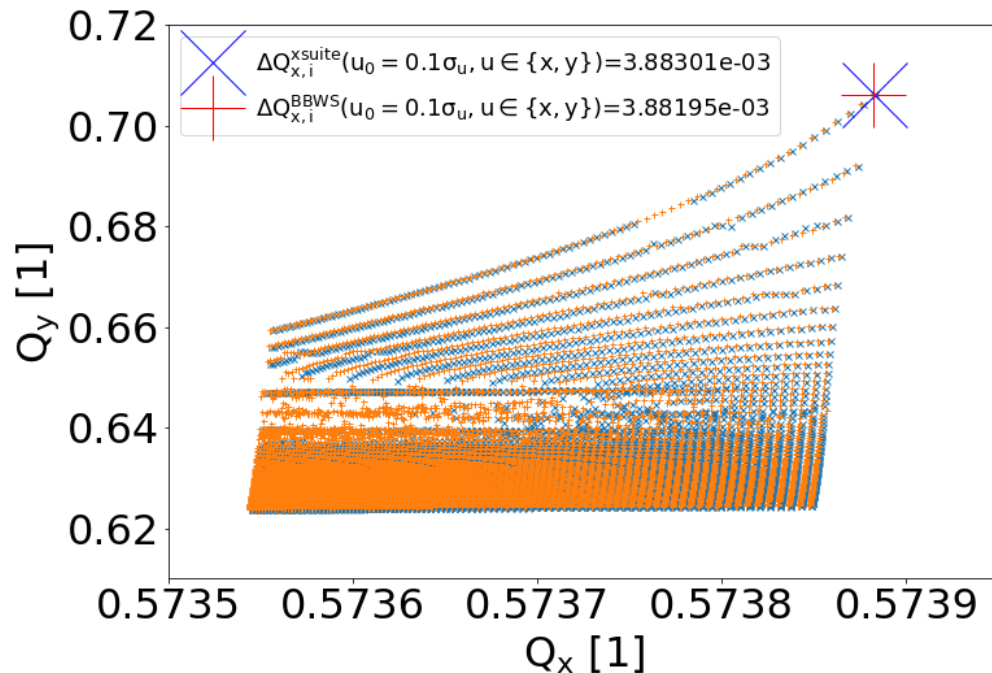
using moments from dist. of 1e6 particles



using moments from analytical Gaussian



FMA - Benchmark of xsuite and BBWS footprint



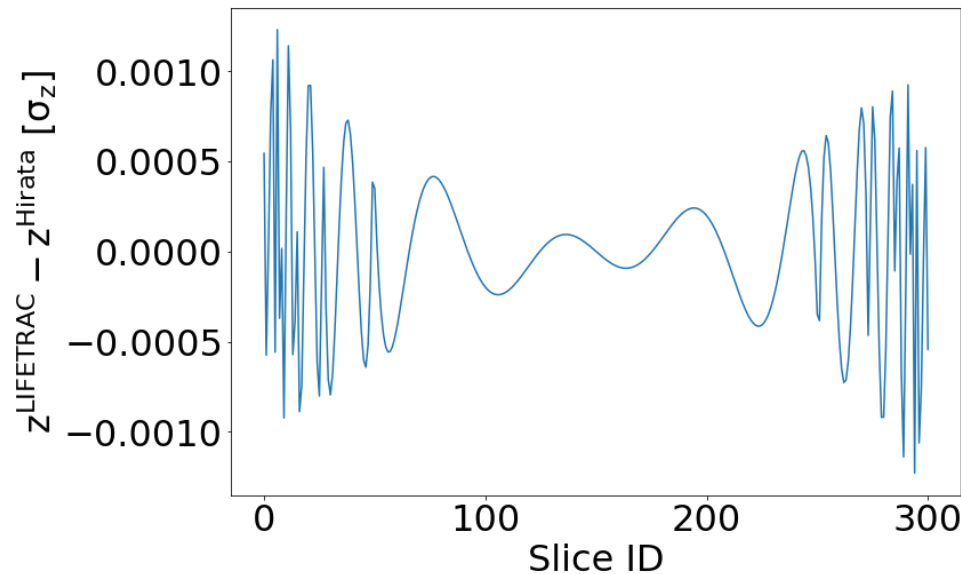
- Good agreement ($1e-6$ accuracy) between BBWS and xsuite
- $1e-4$ discrepancy in ΔQ_x compared to LIFETRAC

Code	ΔQ_x
xsuite / BBWS	$3.88e-3$
LIFETRAC	$3.806e-3$

- xsuite/BBWS both use slicing algorithm from Hirata (SBC) [1]

FMA – slicing algorithms

- Hirata's algorithm has a lower accuracy for slice positions (precomputed table for error function, single precision)
- Comparison of strong slice z centroids, 301 slices, uniform charge



FMA – slicing algorithms

- Switch to slicing algorithms used in LIFETRAC

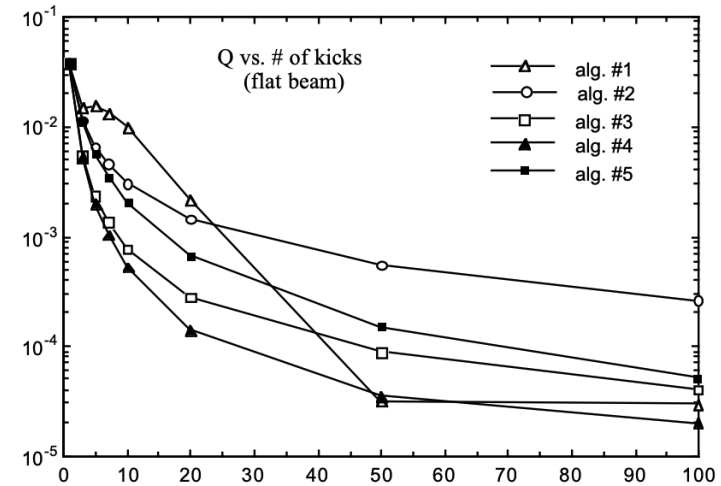
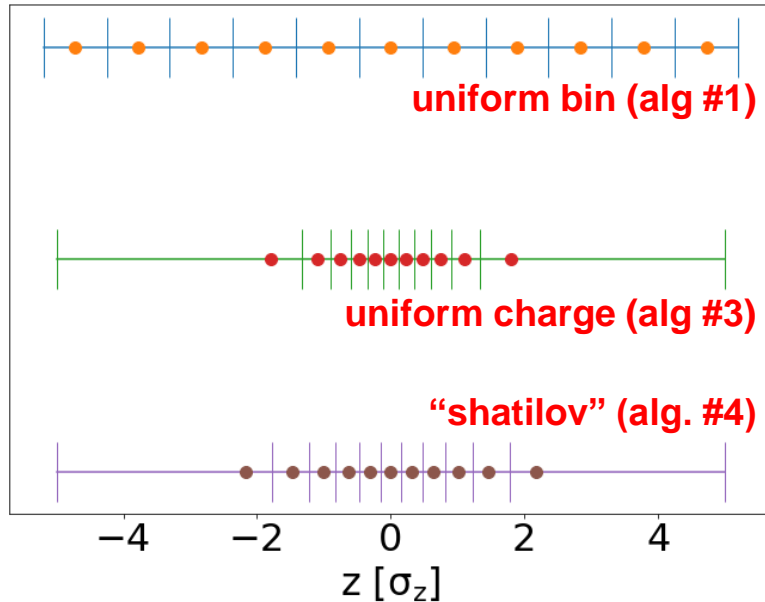
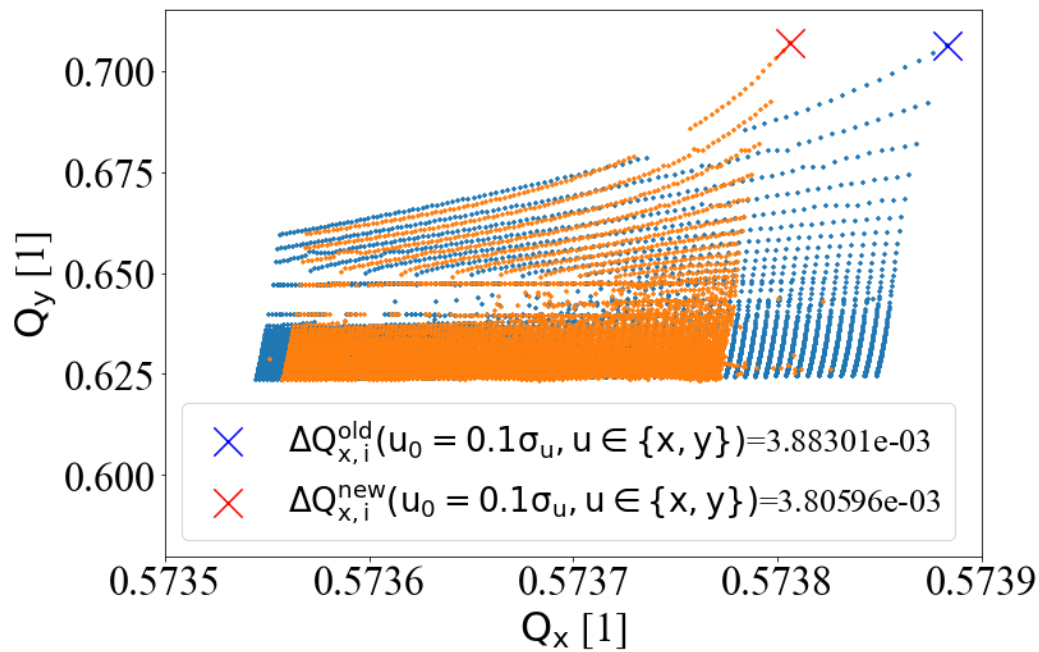


Figure 1: Convergence of the slicing algorithms: Q plotted vs. number of kicks N_s for flat beams (PEP-II-like parameters; see Table 2).

[2]

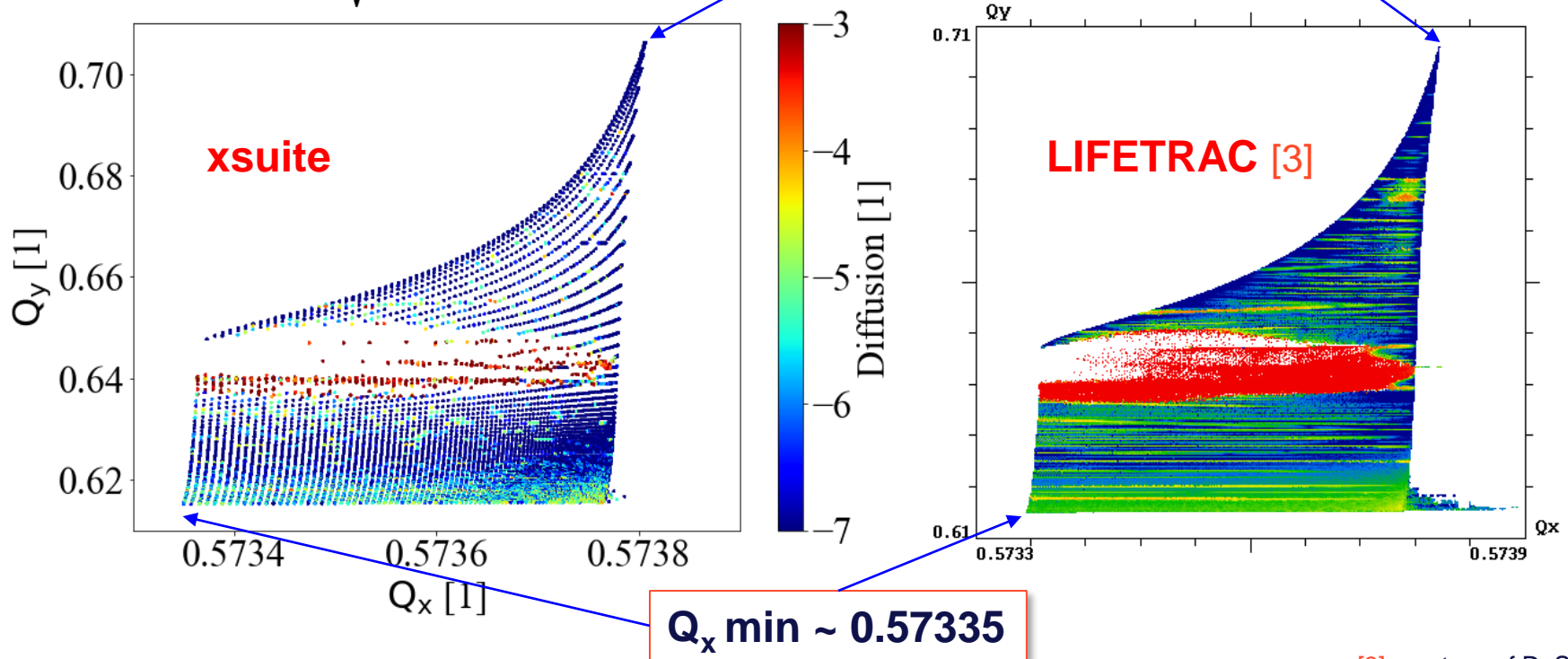
FMA – slicing algorithms

- Compare FCCee Z footprint with old (Hirata's) and new (LIFETRAC) slicing
- All 3 slicing algorithms produce $\Delta Q_x = 3.806e-3$ ($\sim 1e-7$ accuracy)

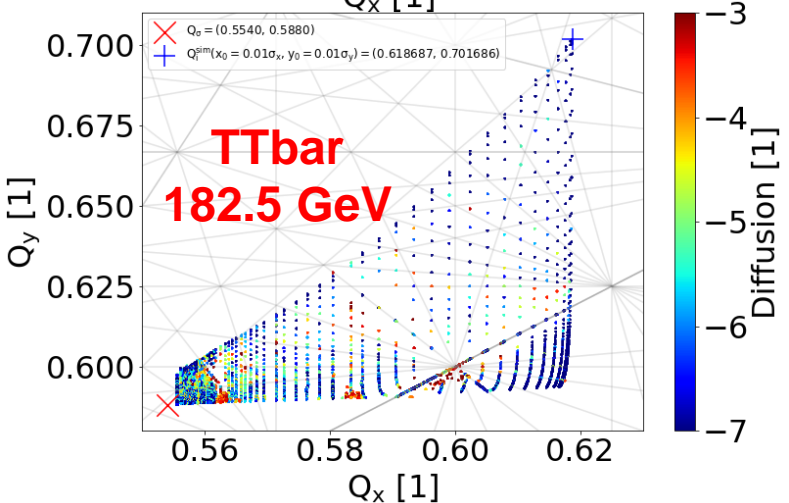
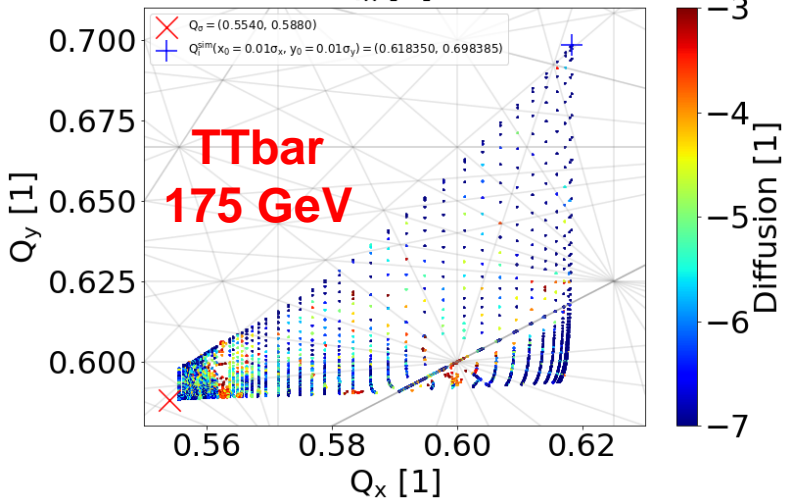
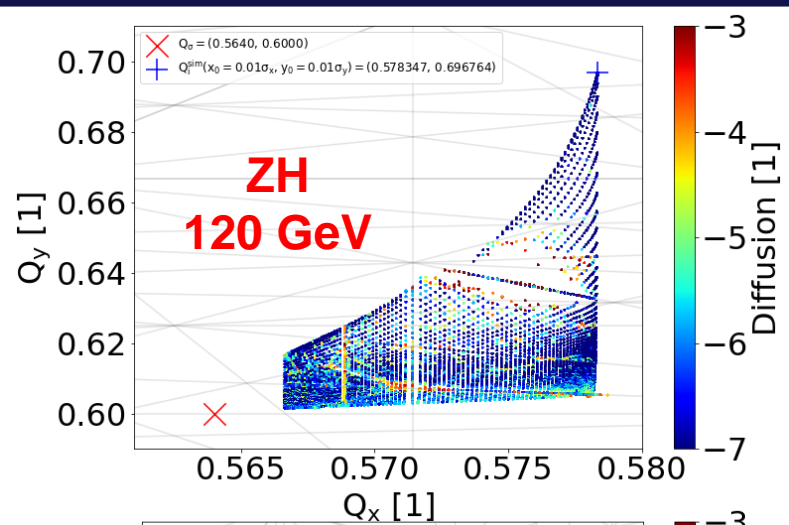
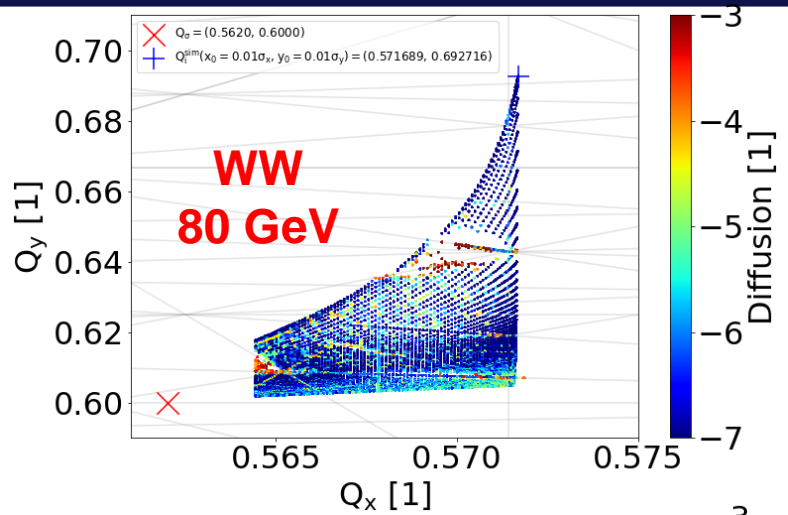


FMA – diffusion

$$D = \log_{10}[\sqrt{\sigma_{Q_{x,i}}^2 + \sigma_{Q_{y,i}}^2}]$$



FMA – Footprint of other FCCee energies



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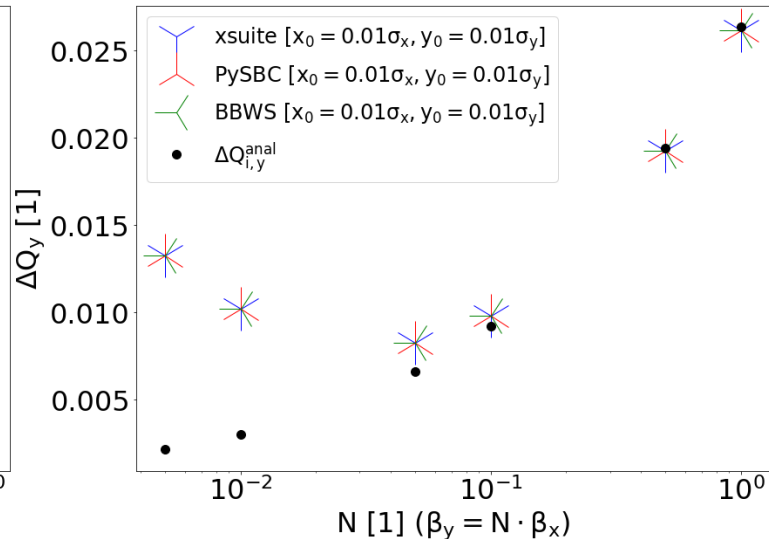
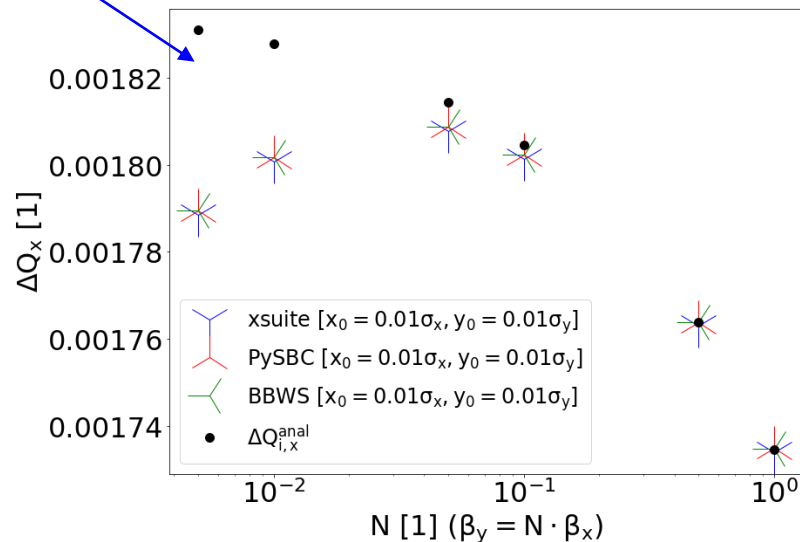
Benchmark against analytical formula ξ

- Head on collision with flat beams ($N_b = 1e8$)
 - All codes agree (xsuite, PySBC, BBWS)
 - Analytical estimate does not account for hourglass effect

$$Q_{x,y}^i = \frac{1}{2\pi} \arccos[\cos(2\pi Q_{x,y}^\sigma) - 2\pi \xi_{x,y} \sin(2\pi Q_{x,y}^\sigma)]$$

$\Phi: 0.0e+00$ [rad]

FCC-ee Z

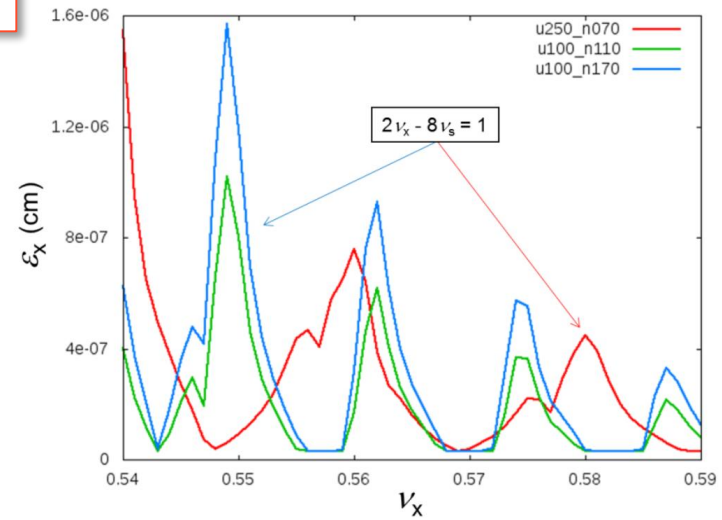
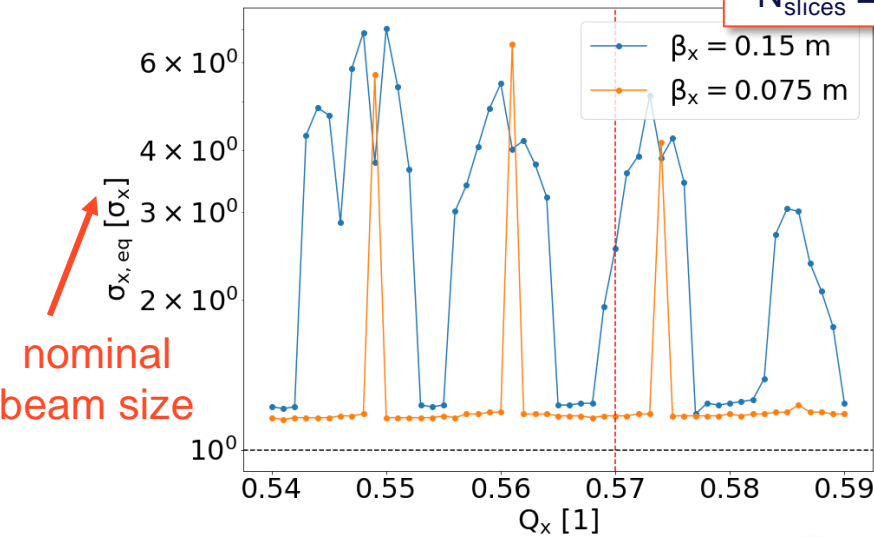


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Tune scan

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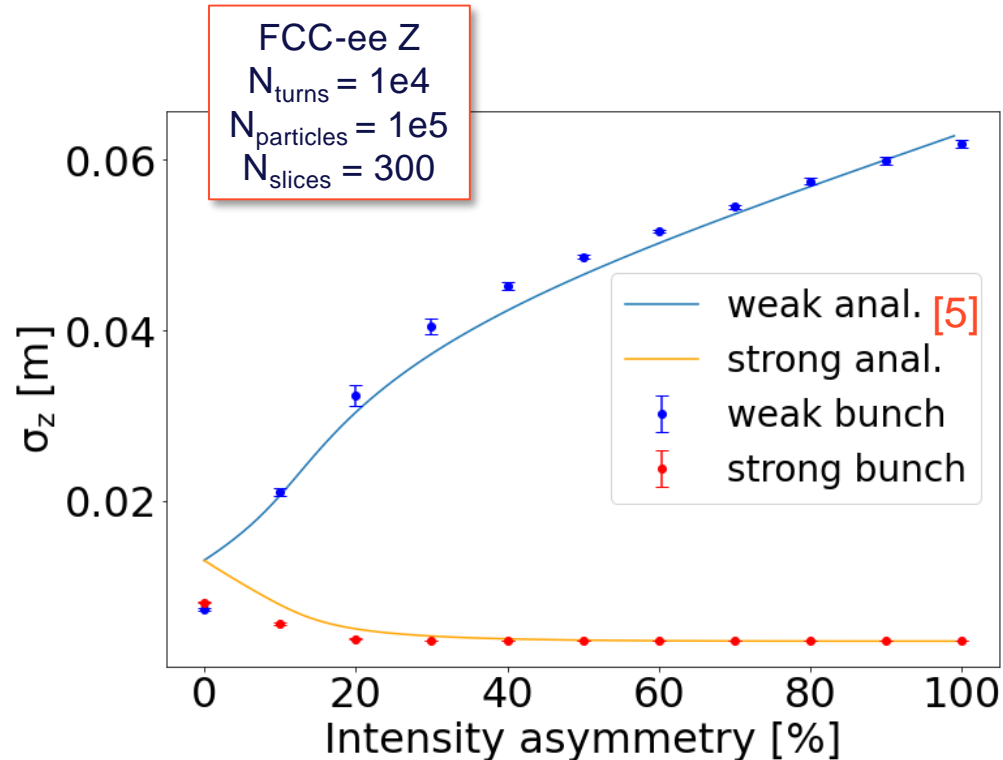


[4]

Figure 7: Growth of ϵ_x due to coherent X-Z instability, as a function of ν_x . Red line corresponds to $U_{RF} = 250 \text{ MV}$, $N_p = 7 \cdot 10^{10}$, green and blue lines – $U_{RF} = 100 \text{ MV}$, $N_p = 1.1 \cdot 10^{11}$ and $1.7 \cdot 10^{11}$.

- Nominal working point is unstable in xsuite
- Reason to be investigated

xsuite benchmark of 3D flip-flop



- First results are promising
- Need more particles and turns
- Improvement of parallelization ongoing
- Symmetric case instability to be understood
- Study ongoing

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Summary

- Successful code benchmarks in weak-strong case
 - xsuite benchmarked against several existing tools, such as GUINEA-PIG, LIFETRAC, PySBC, BBWS
 - Beamstrahlung (feature released on Github [6]), tune footprint
- Ongoing benchmarking and simulations of 3D flip-flop and coherent head-tail instability with the strong-strong model
- Next steps:
 - Link element by element lattice (SAD / MAD-X) to xsuite beam-beam
 - Bhabha scattering

Thank you!

WORK SUPPORTED BY THE SWISS ACCELERATOR RESEARCH AND TECHNOLOGY (CHART)

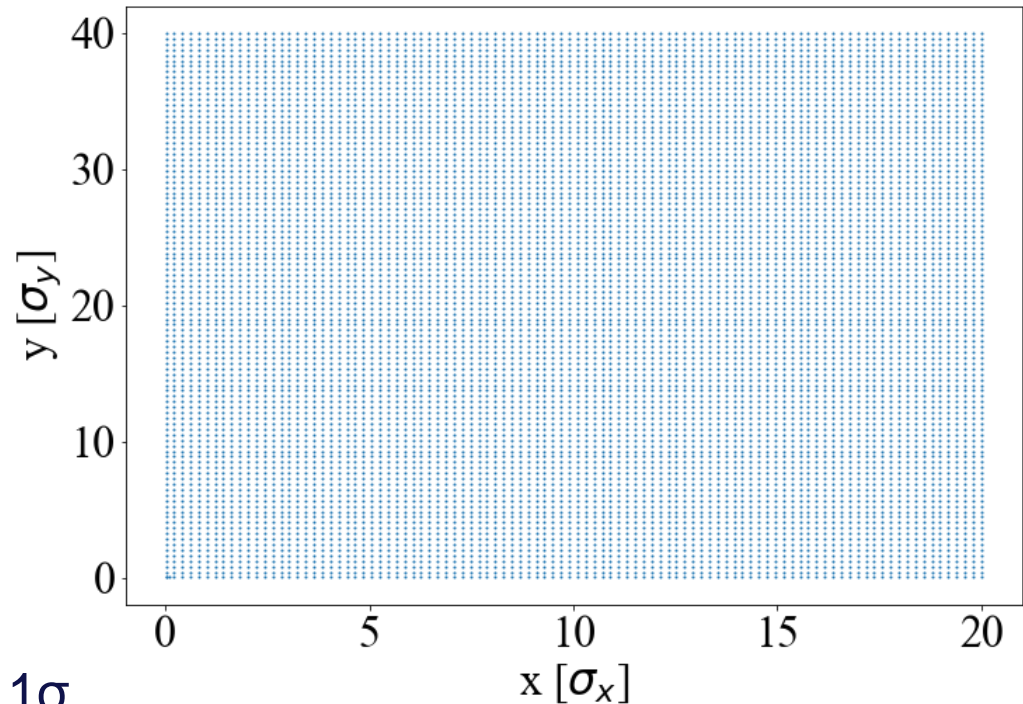
[6] <https://github.com/xsuite/xfields>

BACKUP

FMA – test particle grid

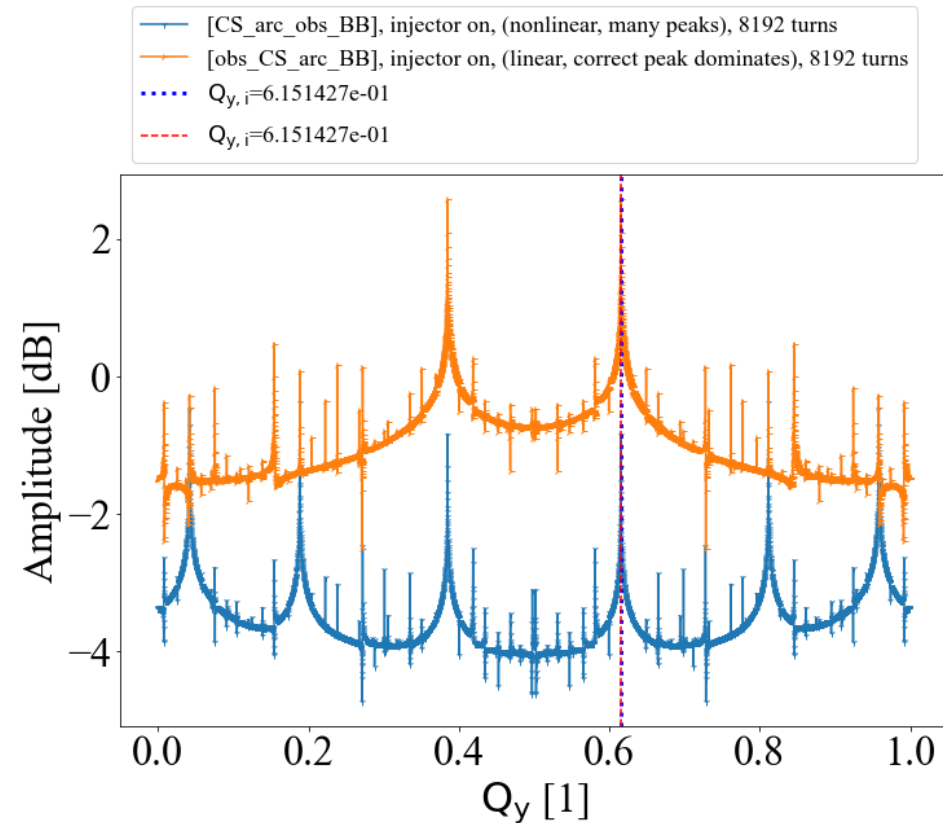
```
n_j = 100 # n_j x n_j particles
j_min = 1e-4
j_max = 10
j_vec_x = np.linspace(0, 2*j_max, n_j)**2/2
j_vec_y = np.linspace(0, 4*j_max, n_j)**2/2
j_vec_x[0] = j_min
j_vec_y[0] = j_min
x_test_vec = np.sqrt(2*j_vec_x)
y_test_vec = np.sqrt(2*j_vec_y)
```

- 10002 particles:
 - 100 x 100 grid
 - Extra test particles: 0.01σ , 0.1σ
- FMA with 8192 turns
- FFT window for diffusion: 4096, stride: 409 (5%), $D = \text{std}$ of 11 tunes
- Observation point before CS: [obs, CS, small arc, IP, small arc, CS, long arc]



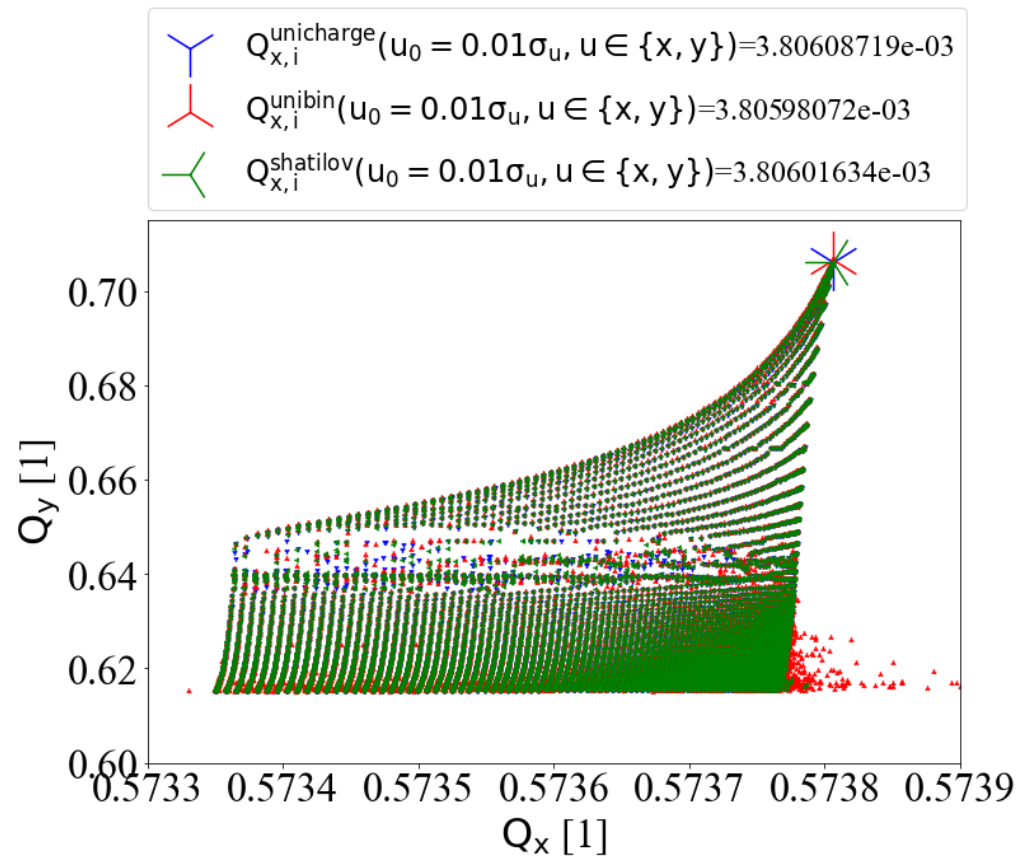
FMA – choice of observation point

- Observation point in linear part of lattice, outside CS (**orange**): spectrum less noisy
- Observation point before/after IP in nonlinear part (**blue**): spectrum more noisy due to CS nonlinearities
- There is no impact on incoh. tune
- But: in the second case for large amplitude particles the peak finder can more easily trigger on wrong peak

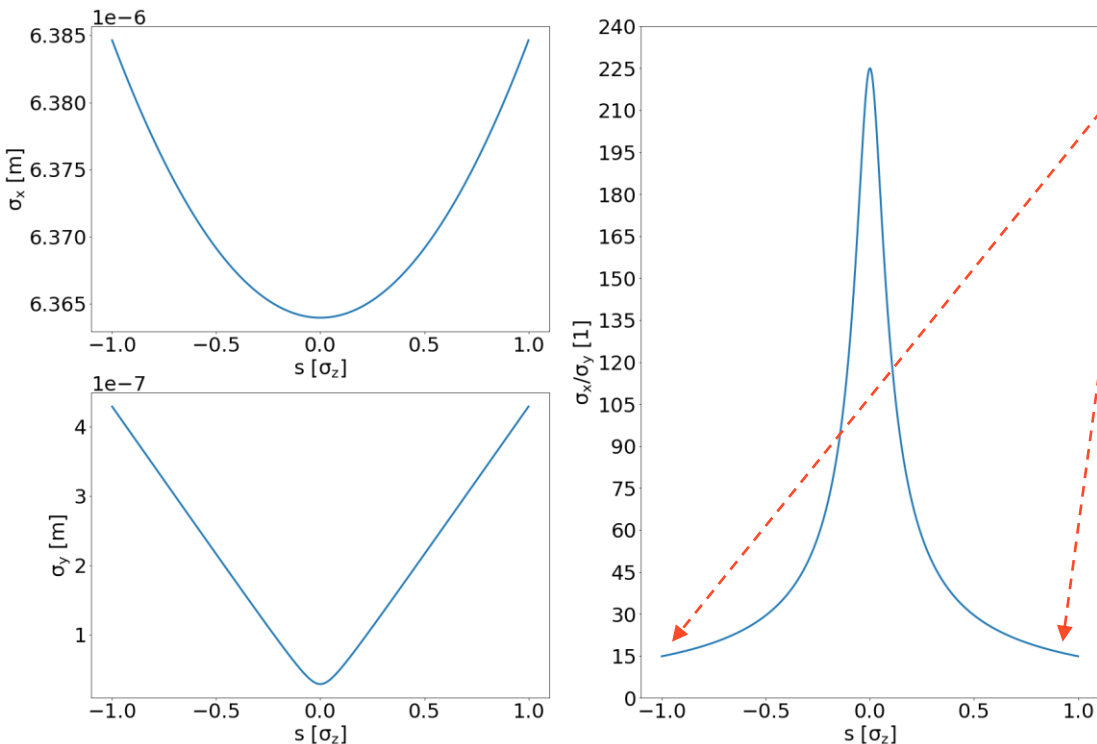


FMA – footprint with all 3 slicing algorithms

- Observation point in linear part of lattice: [obs, CS, arc, BB, arc, CS, arc]
- Injector element before tracking
 - Turn starts with CS
 - Need to match optical functions
- Same settings on slide #12, #13



Benchmark against analytical formula ξ

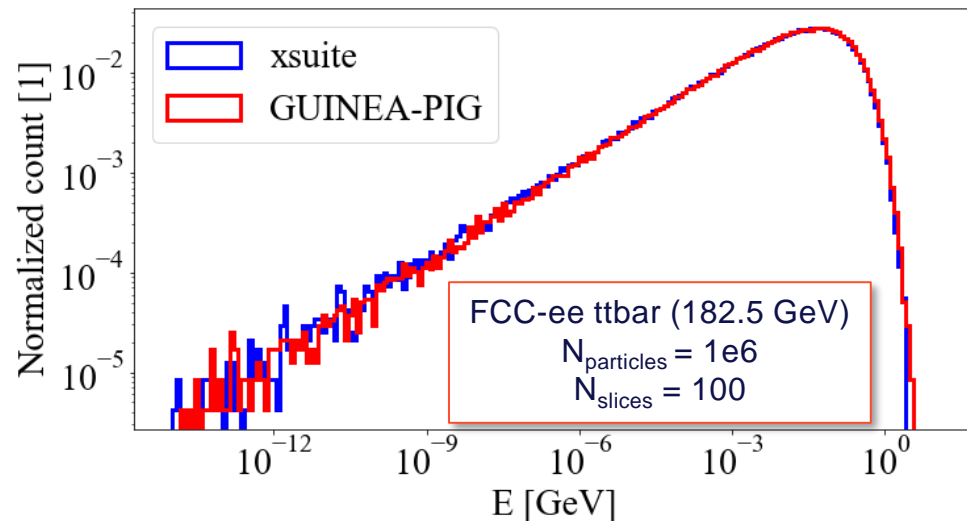


- For FCC-ee Z:
 $\sigma_x/\sigma_y \sim 15$ at $\mathbf{s} = \sigma_z$
- Therefore $\sigma_x \gg \sigma_y$
starts to break
- $\Delta Q_x \sim \xi_x \sim \beta_x / (\sigma_x(\sigma_x + \sigma_y))$
- ΔQ_x still depends on σ_y
- ΔQ_x smaller than w/o
hourglass

Beamstrahlung benchmark

- Benchmark against reference code GUINEA-PIG [1]
 - Beamstrahlung model OK
- xsuite: **weak-strong**
- GUINEA-PIG: **strong-strong**
 - More photons emitted for extremely flat beams (Z)
- Possibility to generate photons for external use (collimation, MDI) [2]
- **TODO:** come up with an efficient model of Bhabha scattering

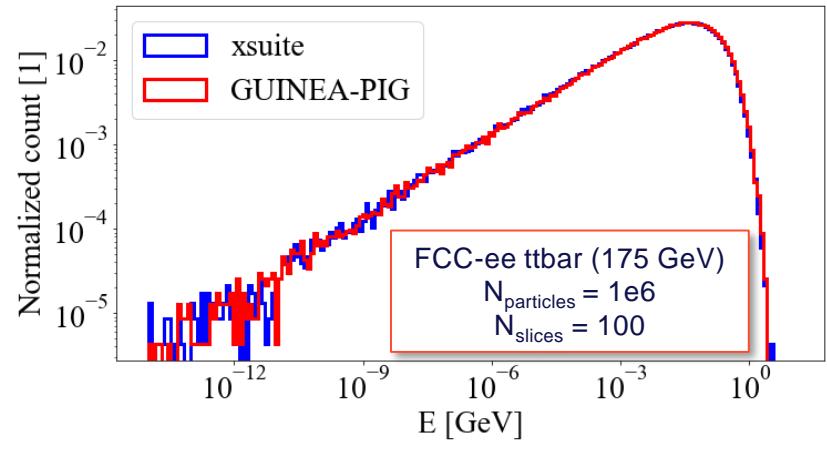
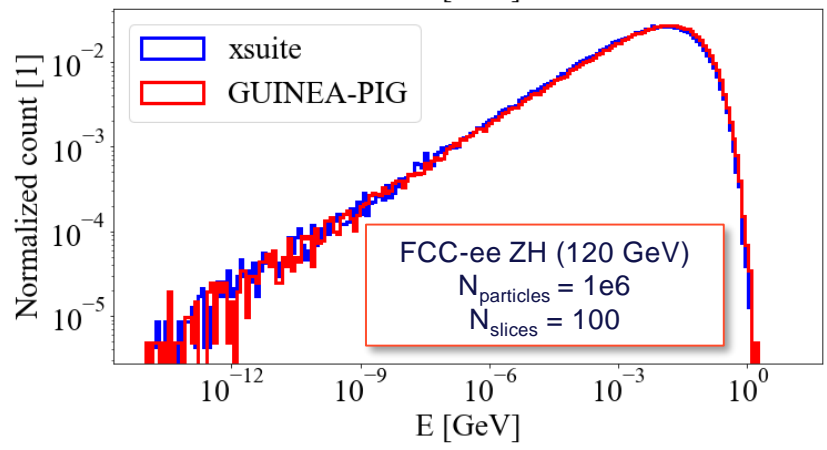
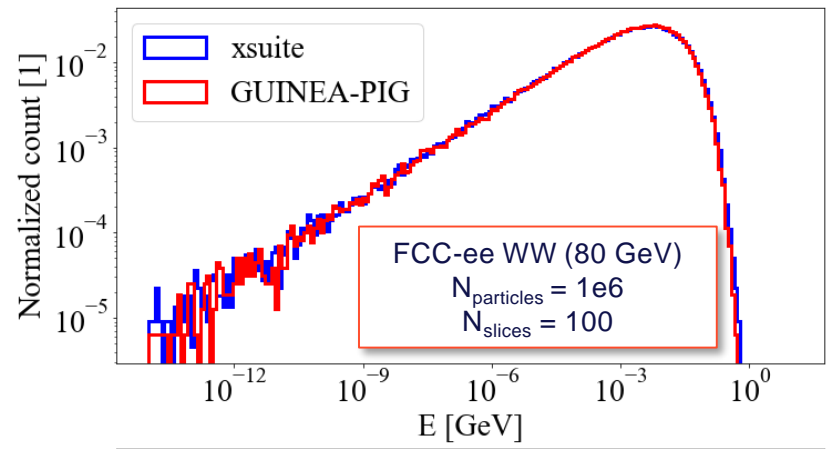
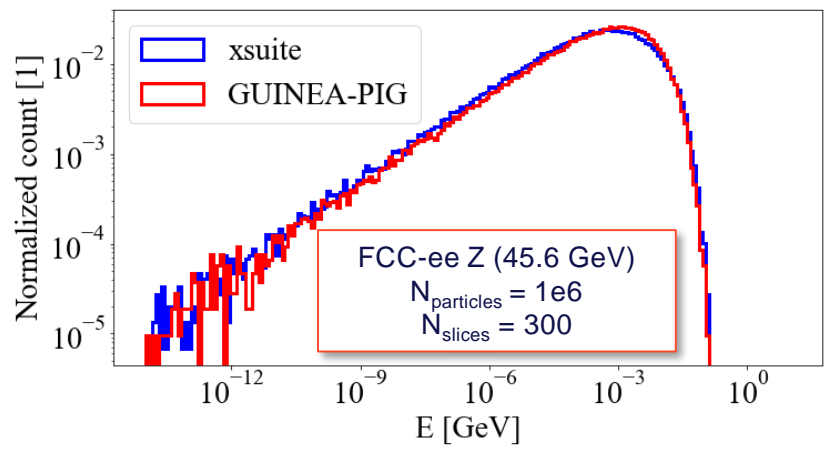
Beamstrahlung photon spectrum / coll.



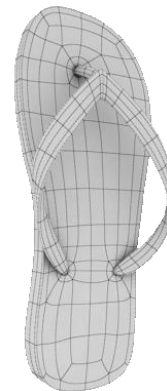
[1] <https://twiki.cern.ch/twiki/bin/view/ABPCComputing/Guinea-Pig>

[2] https://xsuite.readthedocs.io/en/latest/internal_record.html#internal-record-for-elements-used-in-standalone-mode

Beamstrahlung benchmark – other FCCee energies

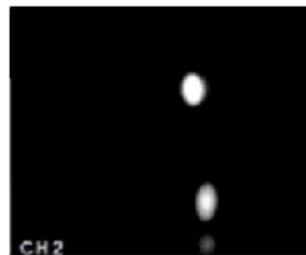


3D flip-flop [1]

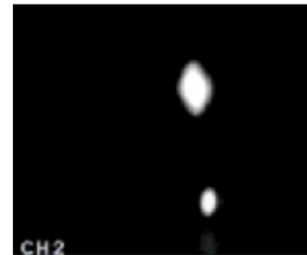


- Occurs only in presence of BS
 1. BS leads to asymmetry in the bunch length σ_z due to asymmetry of bunch currents (which are never exactly the same)
 2. Asymmetry in σ_z leads to amplified synchro-betatron resonances of the longer σ_z bunch
 3. Tune footprint expands and can cross more resonance lines
 4. Increase of both transversal emittances of the longer bunch
 - Increased emittances cause asymmetry in transversal beam size. This asymmetry increases BS and σ_z for the longer bunch but decreases BS and shrinks σ_z for the shorter bunch:
 - σ_z asymmetry increases even more
 5. Repeat from 2.

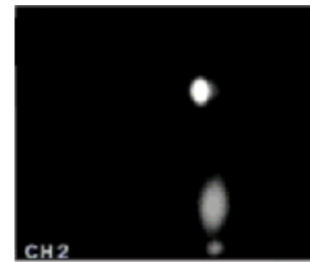
➤ One bunch blows up, the other shrinks



(a) Equal



(b) Blown-up positron beam



(c) Blown-up electron beam

[1] D. Shatilov [<http://www.icfa-bd.org/Newsletter72.pdf>]

[2] X. buffat [<https://e-publishing.cern.ch/index.php/CYRSP/article/view/265/277>]