

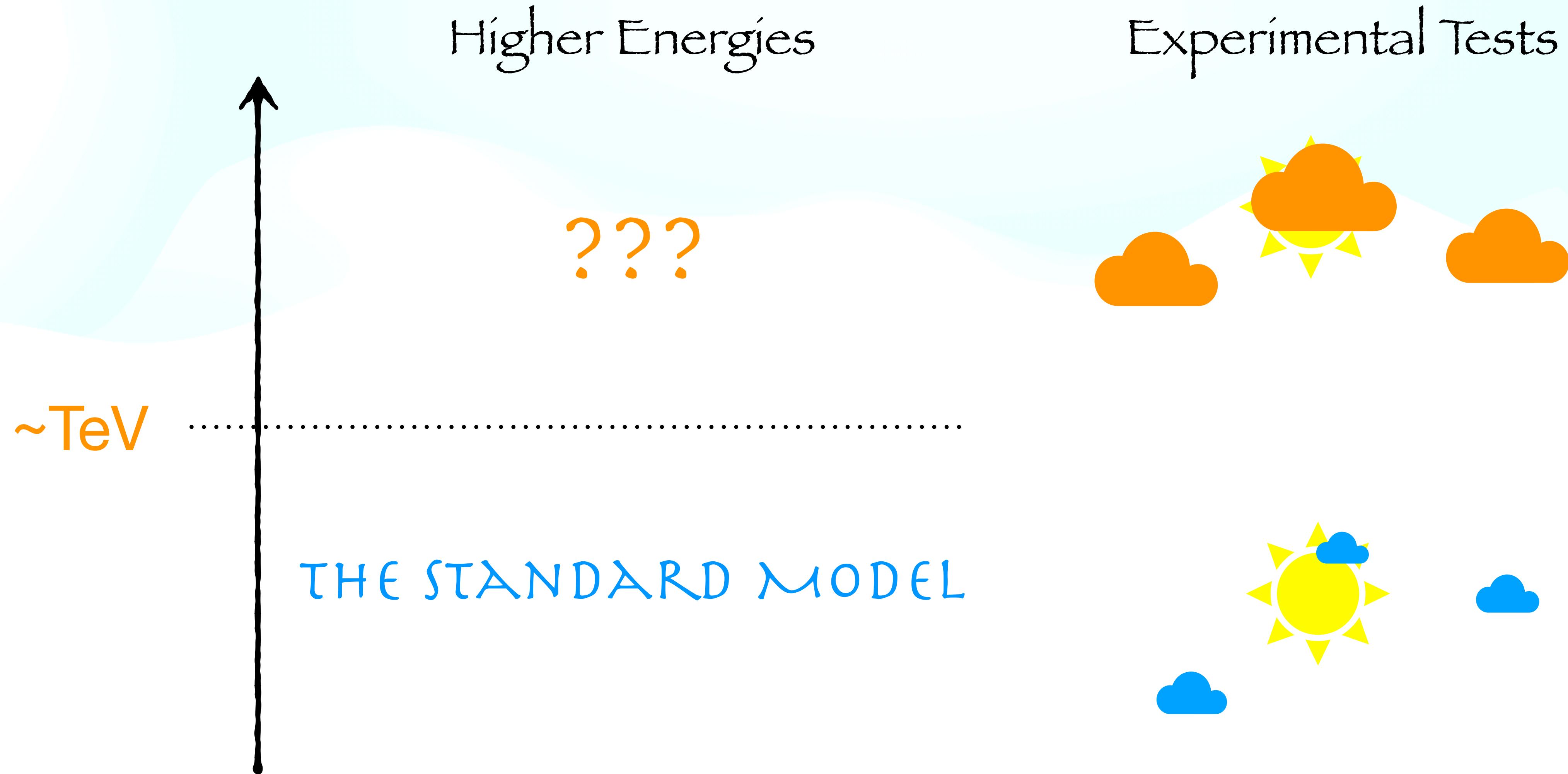
Bubbletrons

Filippo Sala
U. Bologna

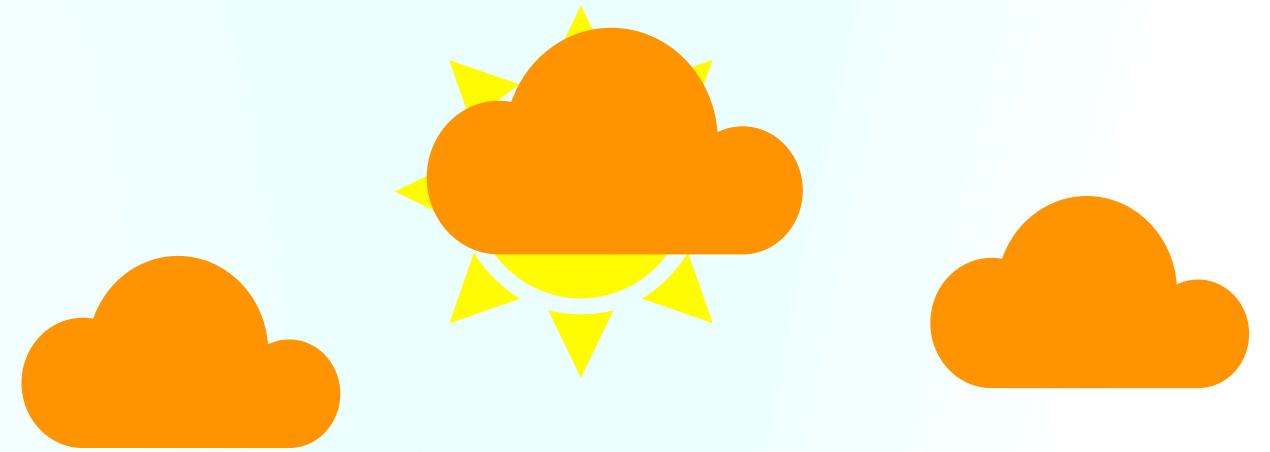
Portorož 14 Apr 2023



A Big Picture

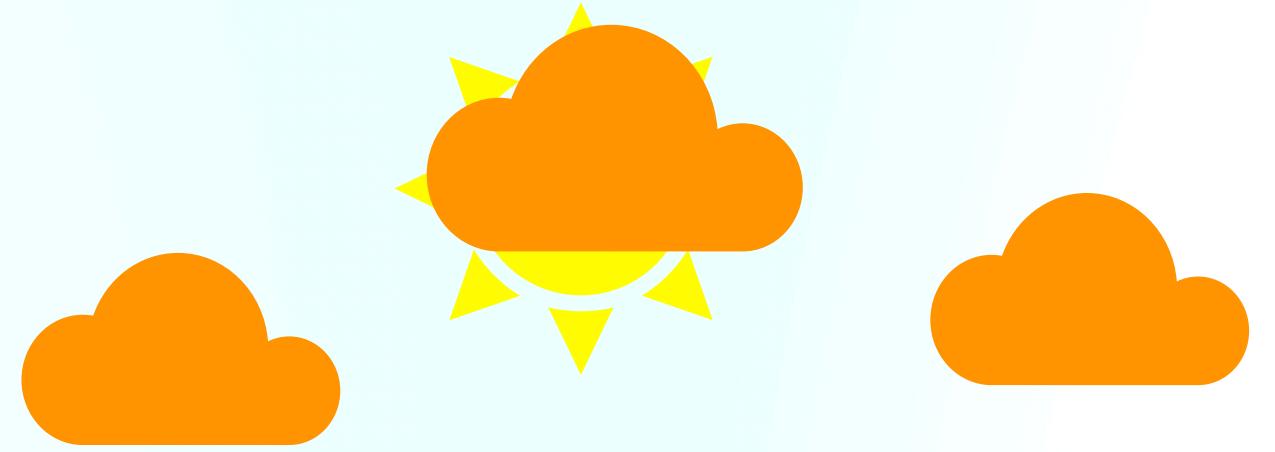


How to access higher energies? (*)



(*) before muon collider or FCC-hh

How to access higher energies? (*)



Look at the Sky!

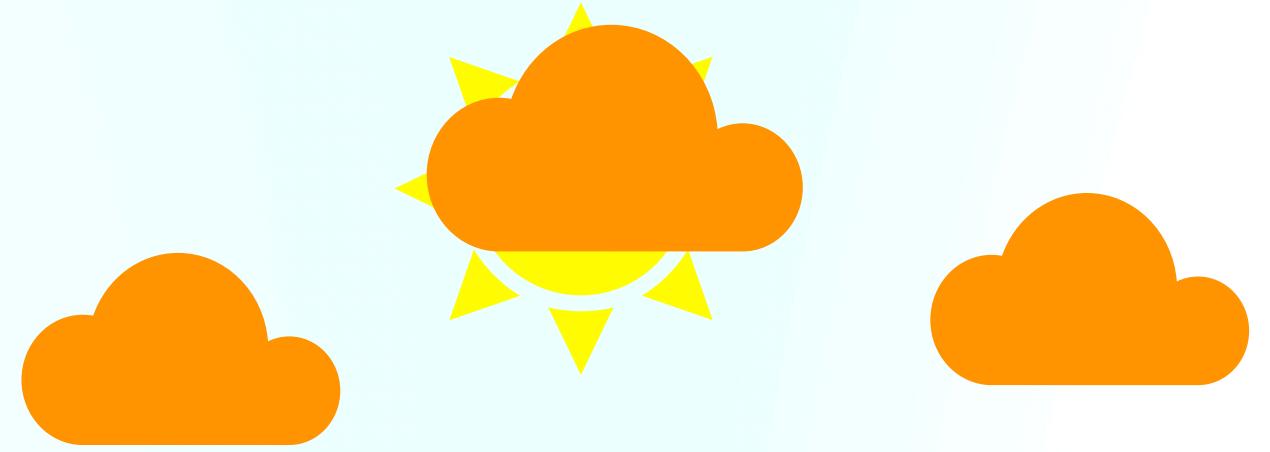
High-Energy Cosmic Rays

Gravitational Waves

The more the merrier

(*) before muon collider or FCC-hh

How to access higher energies? (*)



Look at the Sky!

High-Energy Cosmic Rays

Gravitational Waves

Matter Power Spectrum

Sneaky preview:

Warm EeV Dark Matter from
Weak Scale Phase Transition

(*) before muon collider or FCC-hh

Phase Transitions...

...as Particle Accelerators (“Bubbletrons”)

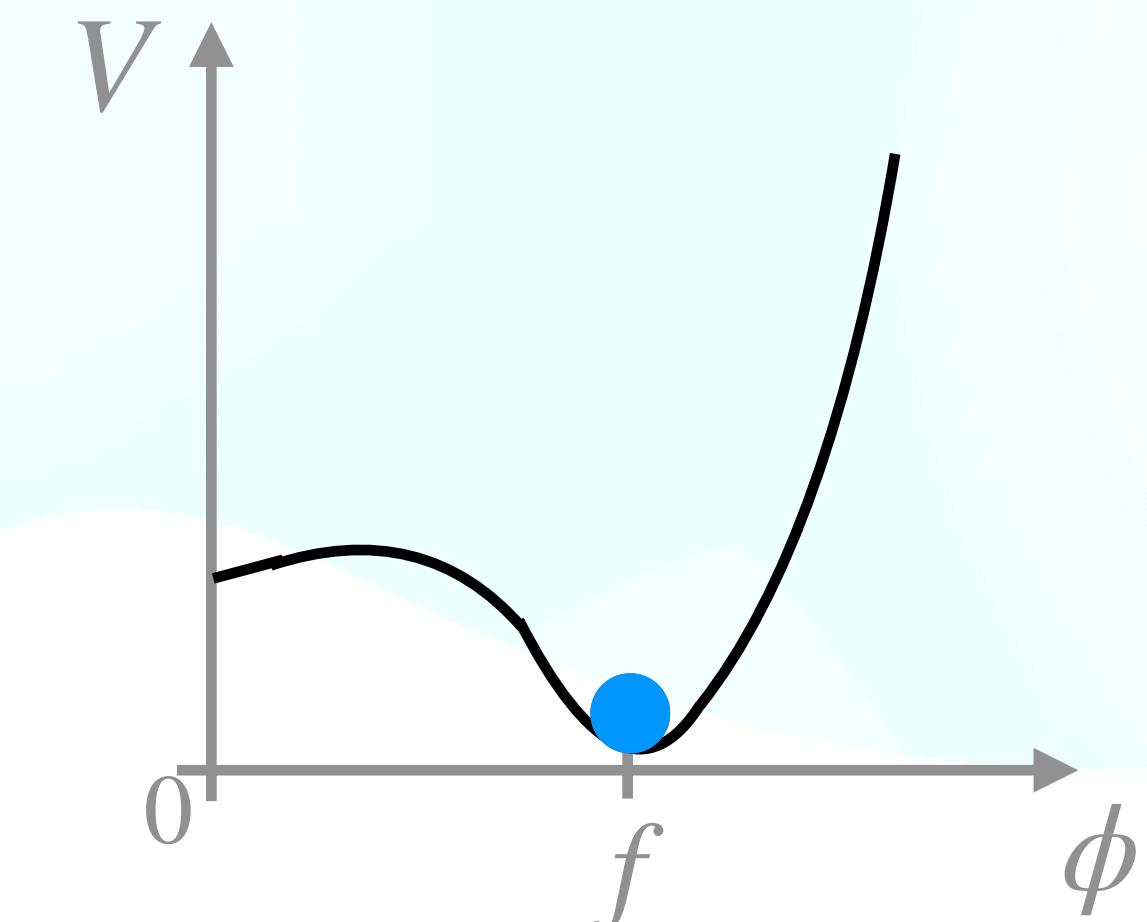
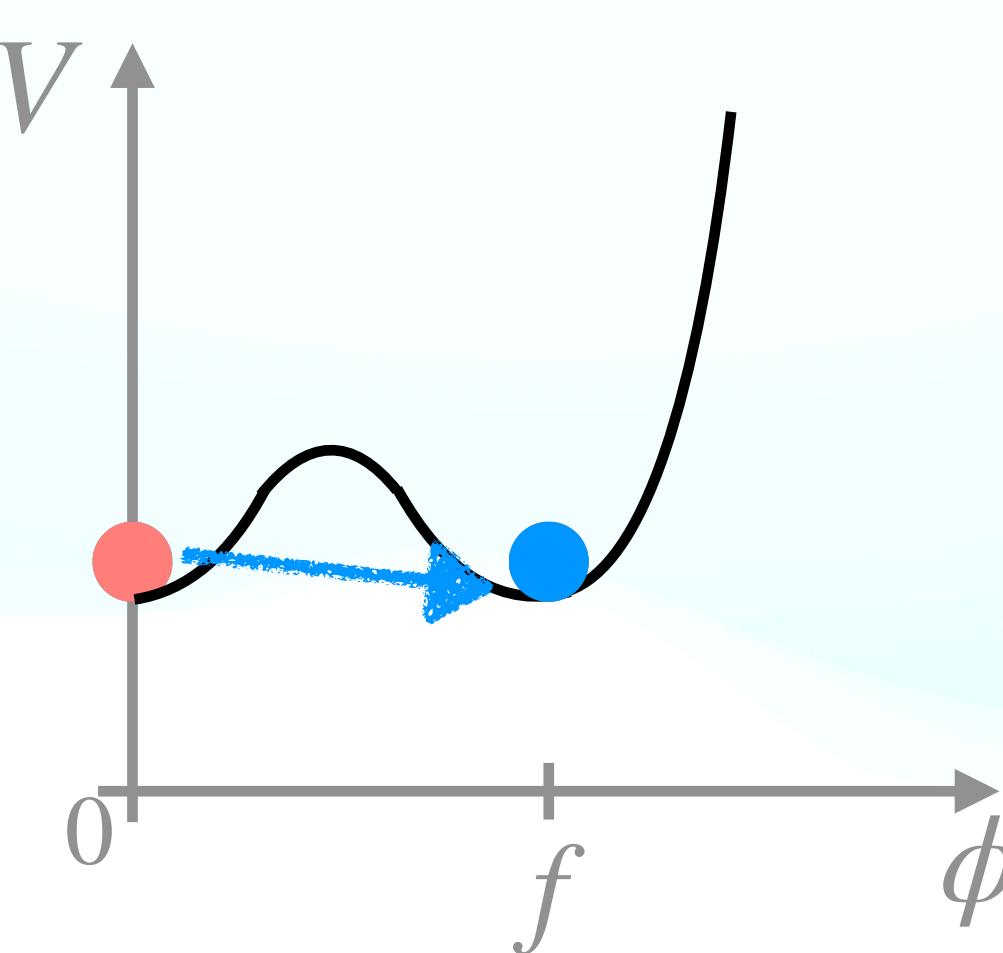
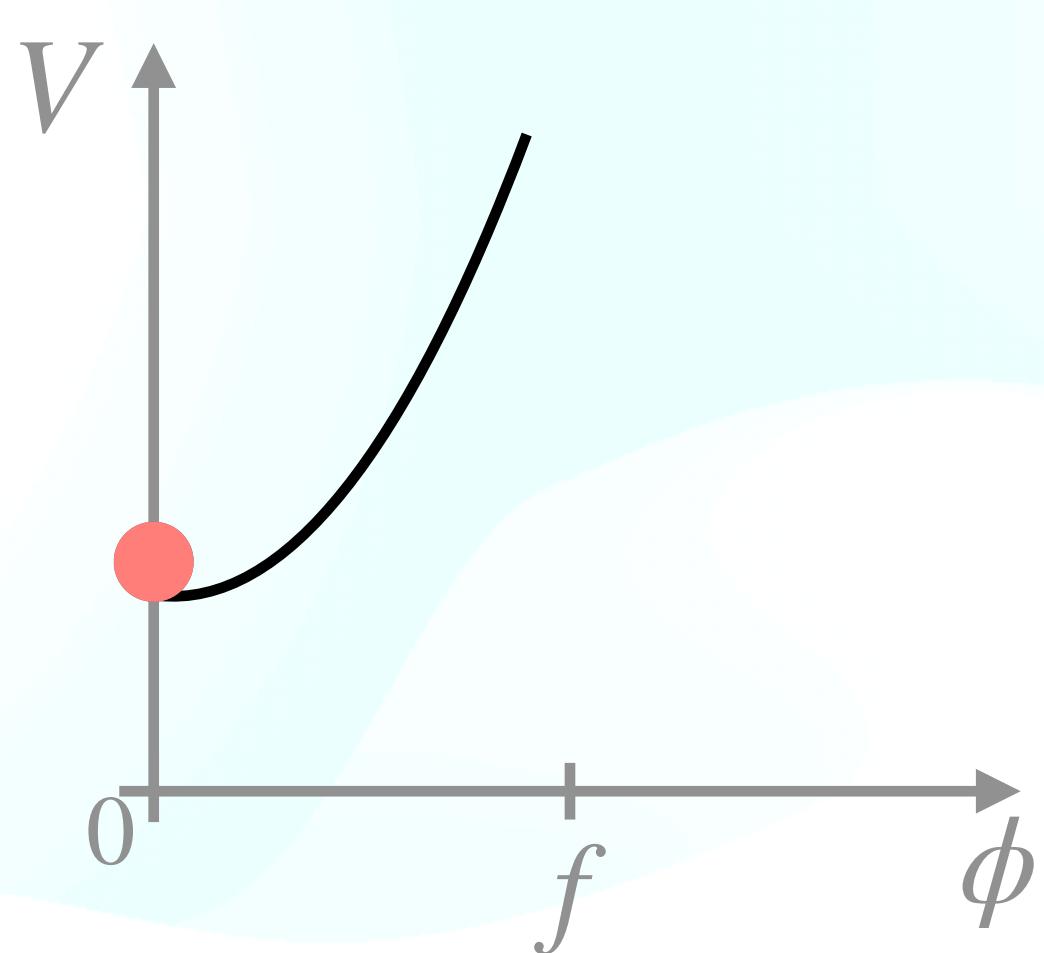
Warm EeV Dark Matter from Weak scale PT

Phase Transitions...

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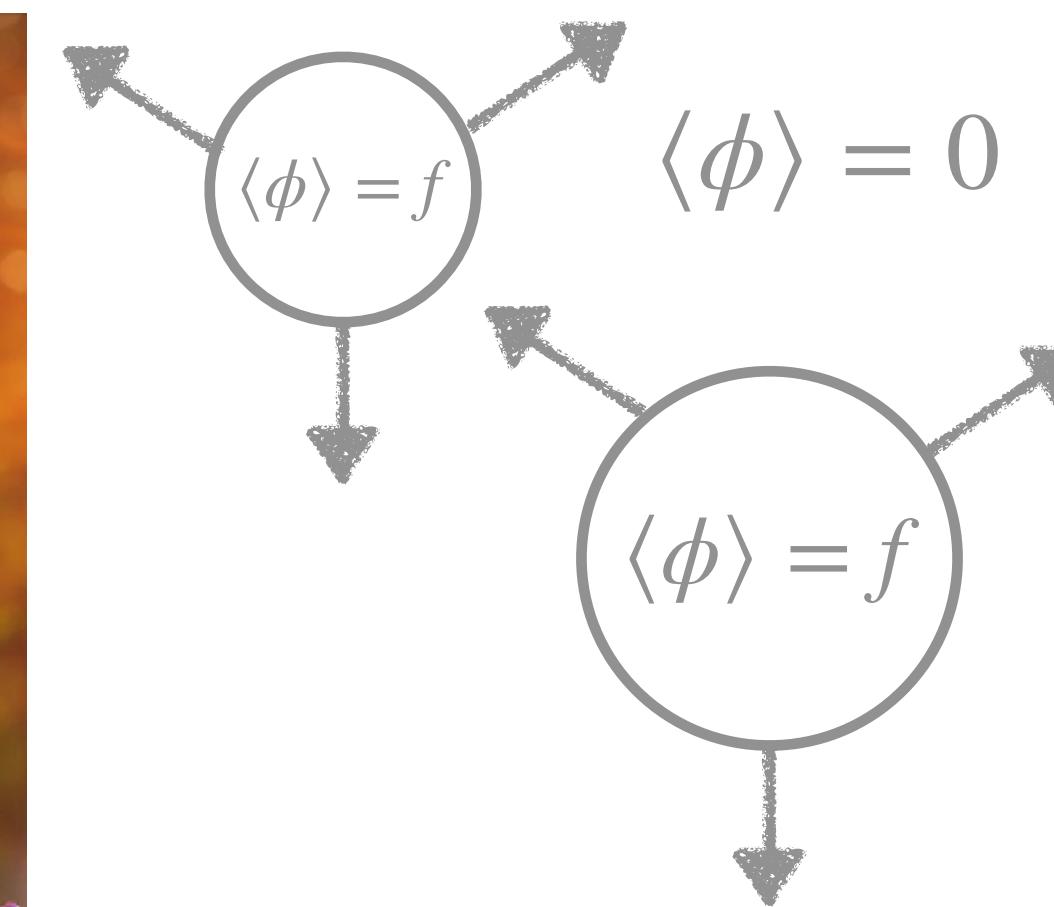
Warm EeV Dark Matter from Weak scale PT

First-order Phase Transitions in Early Universe



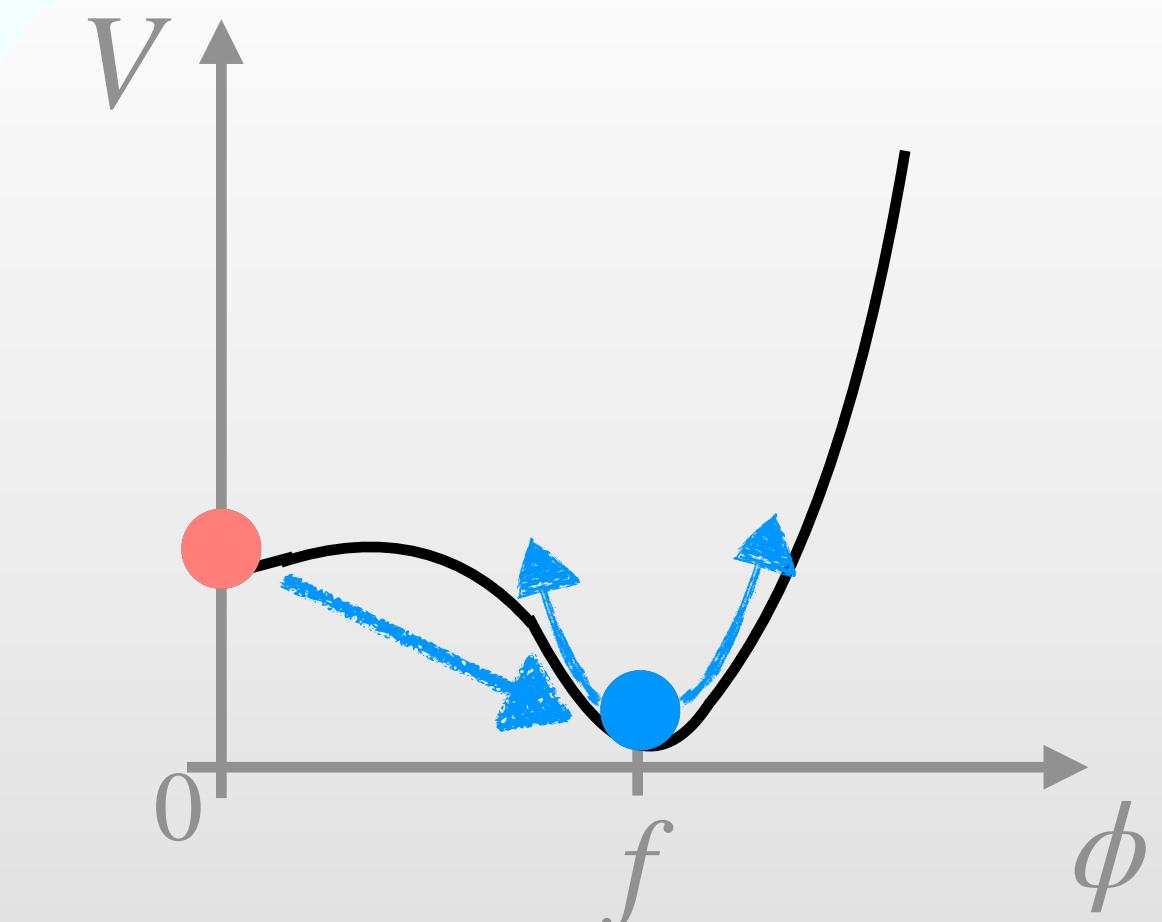
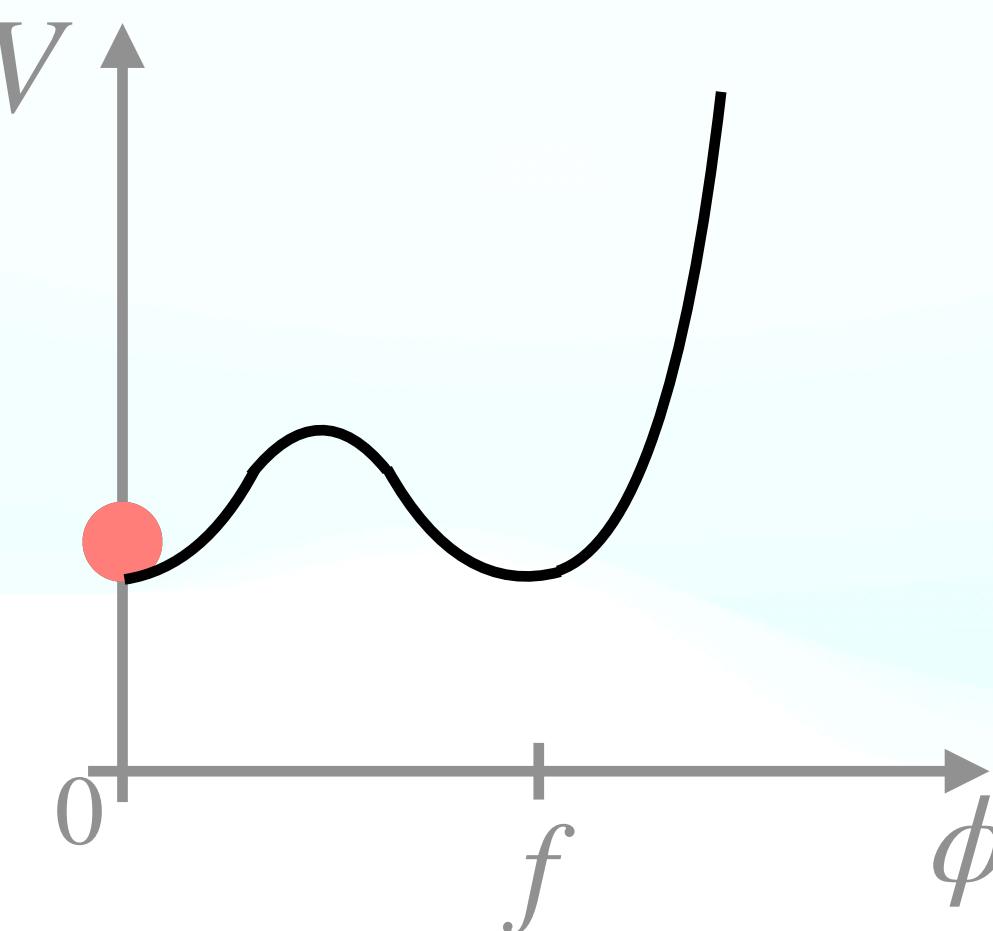
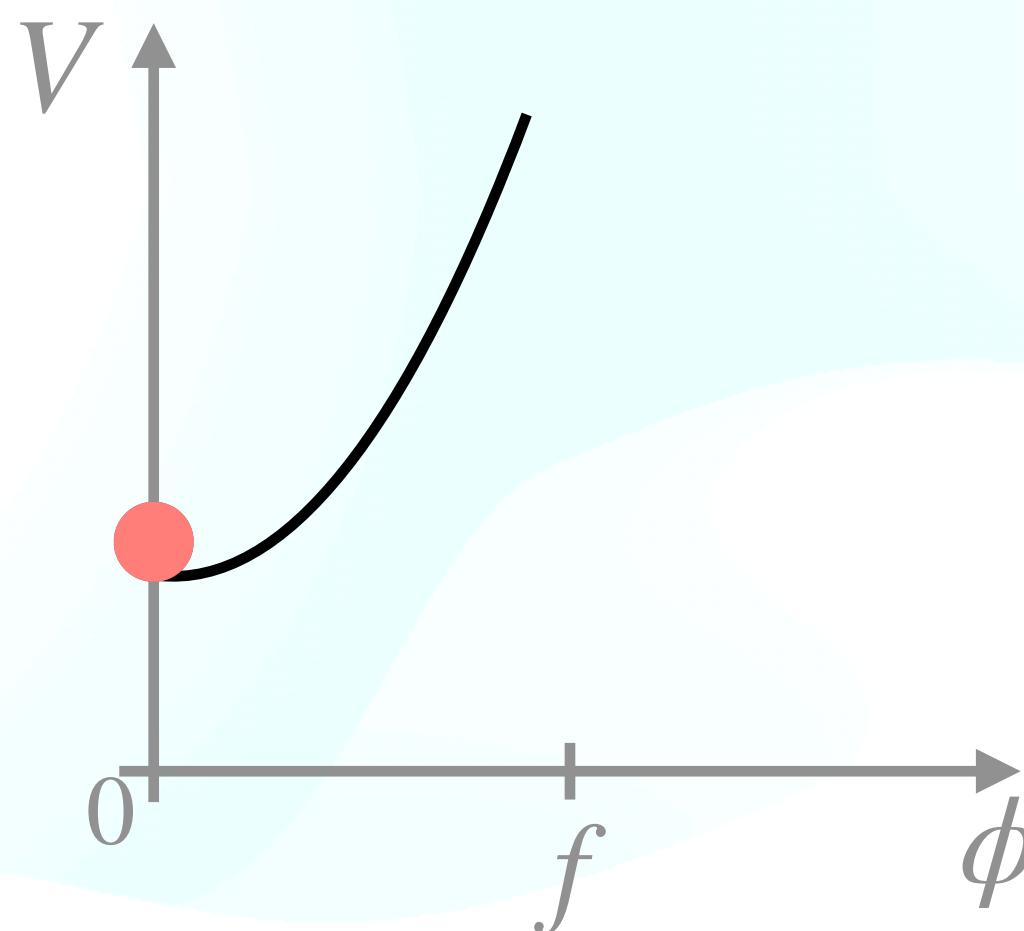
$T_{\text{nuc}} \lesssim f$

Lower Temperature



Gravity Waves!

First-order Phase Transitions in Early Universe



$T_{\text{nuc}} \ll f$

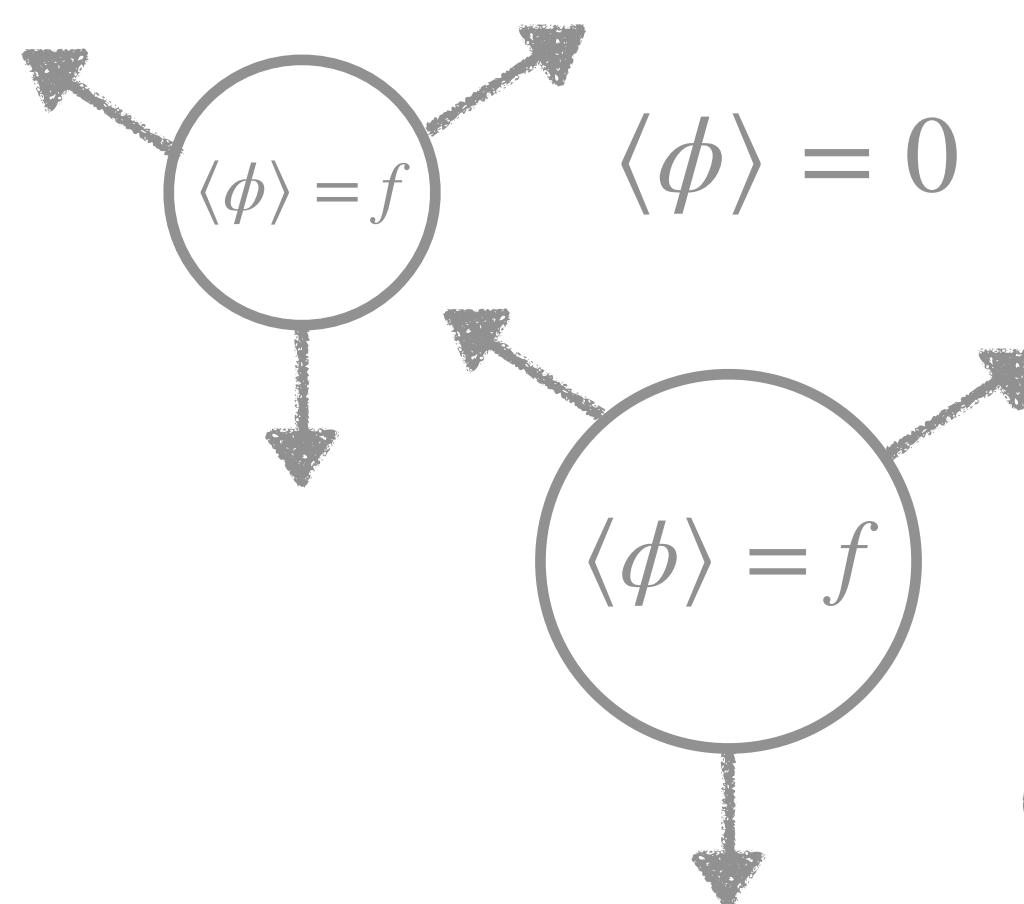
Lower Temperature

IF Phase Transitions here

THEN Supercooling

& oscillations reheat universe to

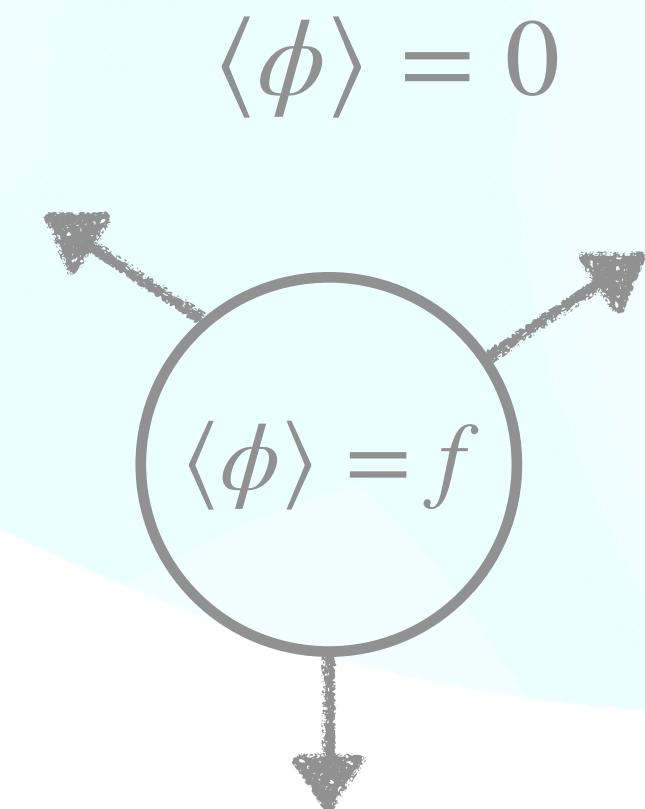
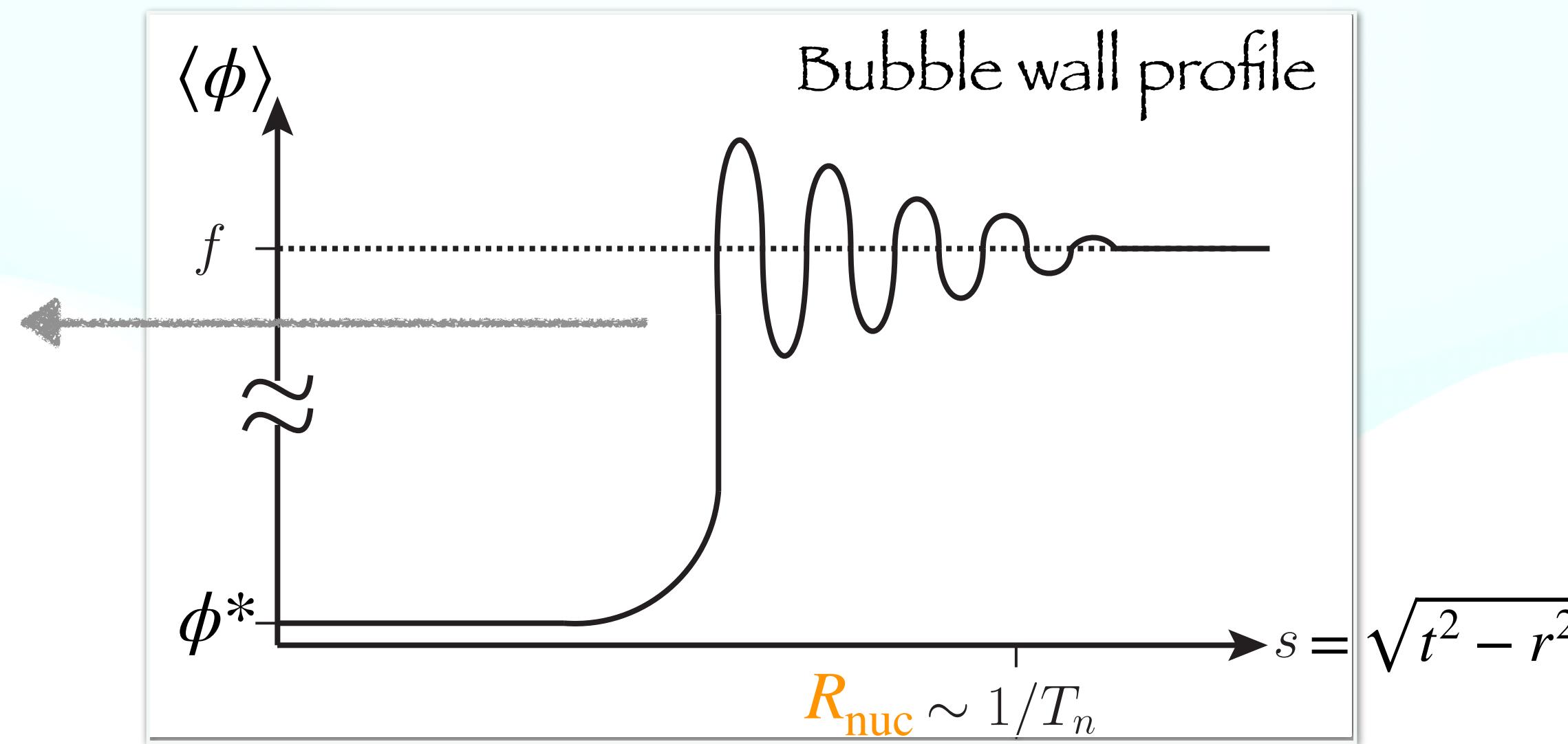
$T_{\text{RH}} \lesssim f$



Gr

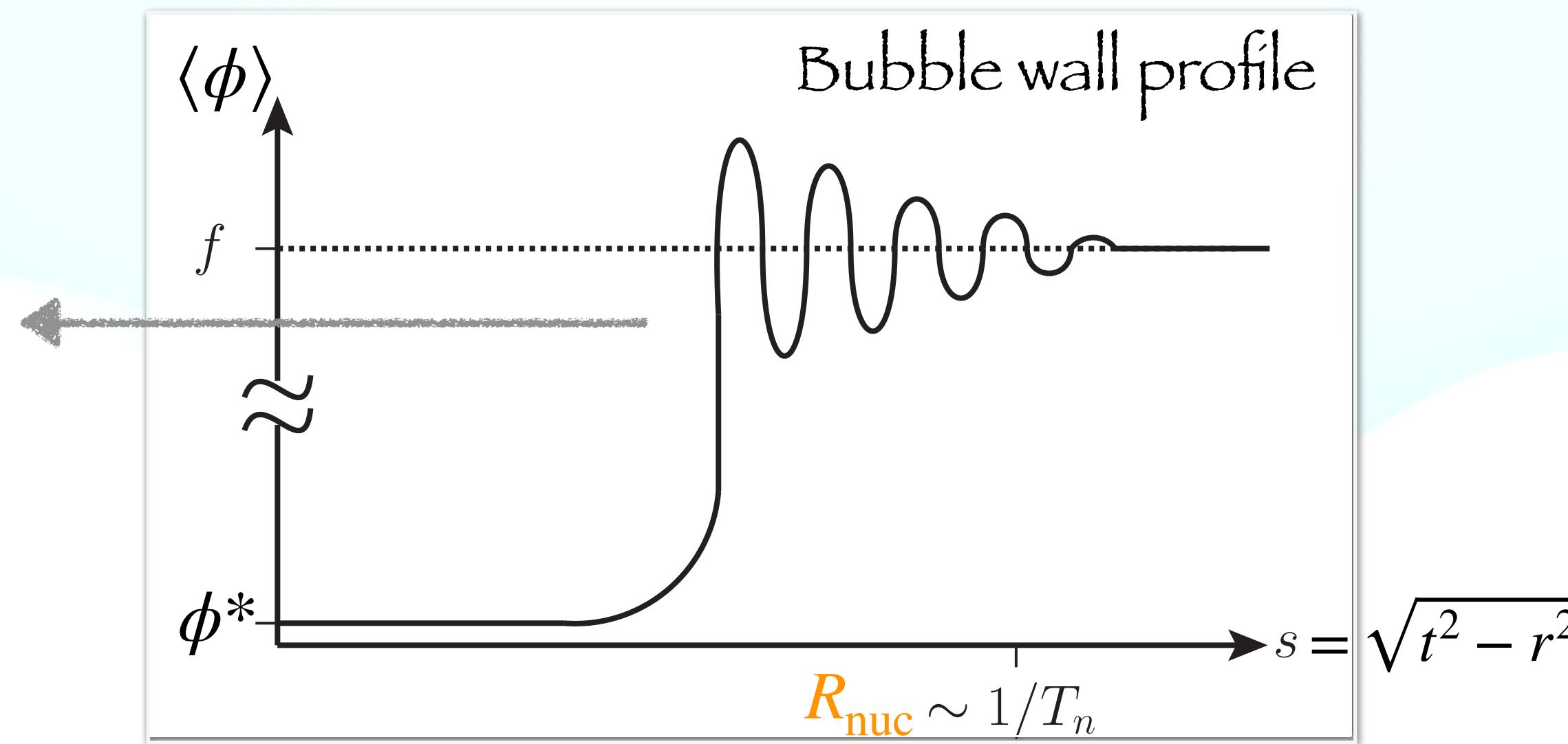
Bubble walls

Wall expands this way
with Lorentz boost γ



Bubble walls

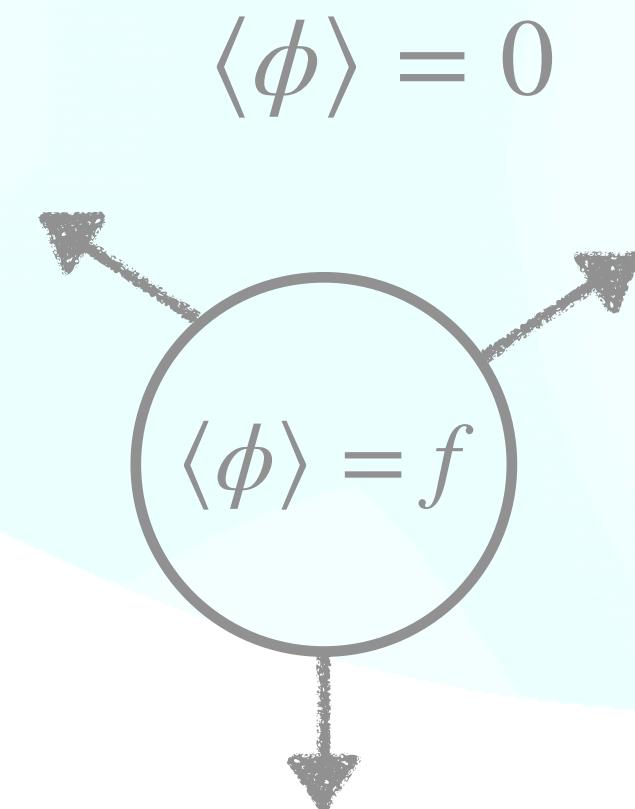
Wall expands this way
with Lorentz boost γ



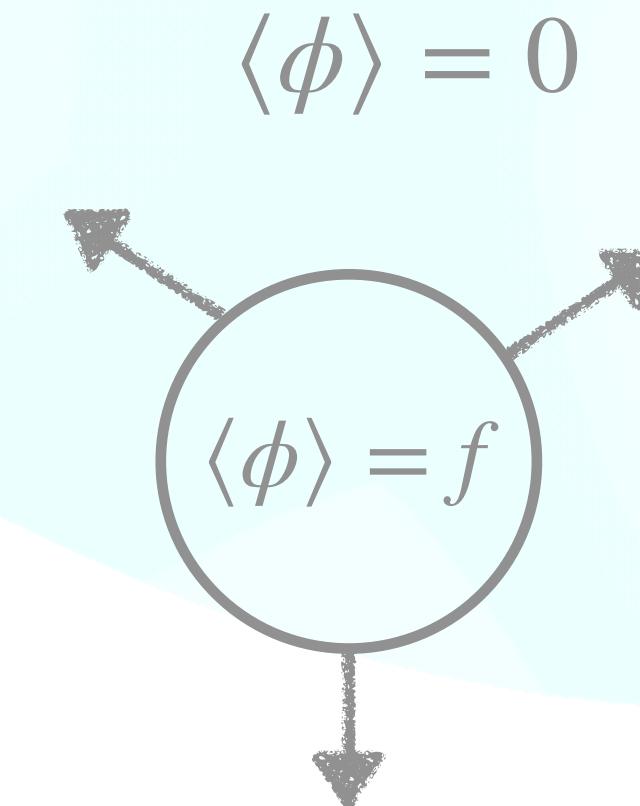
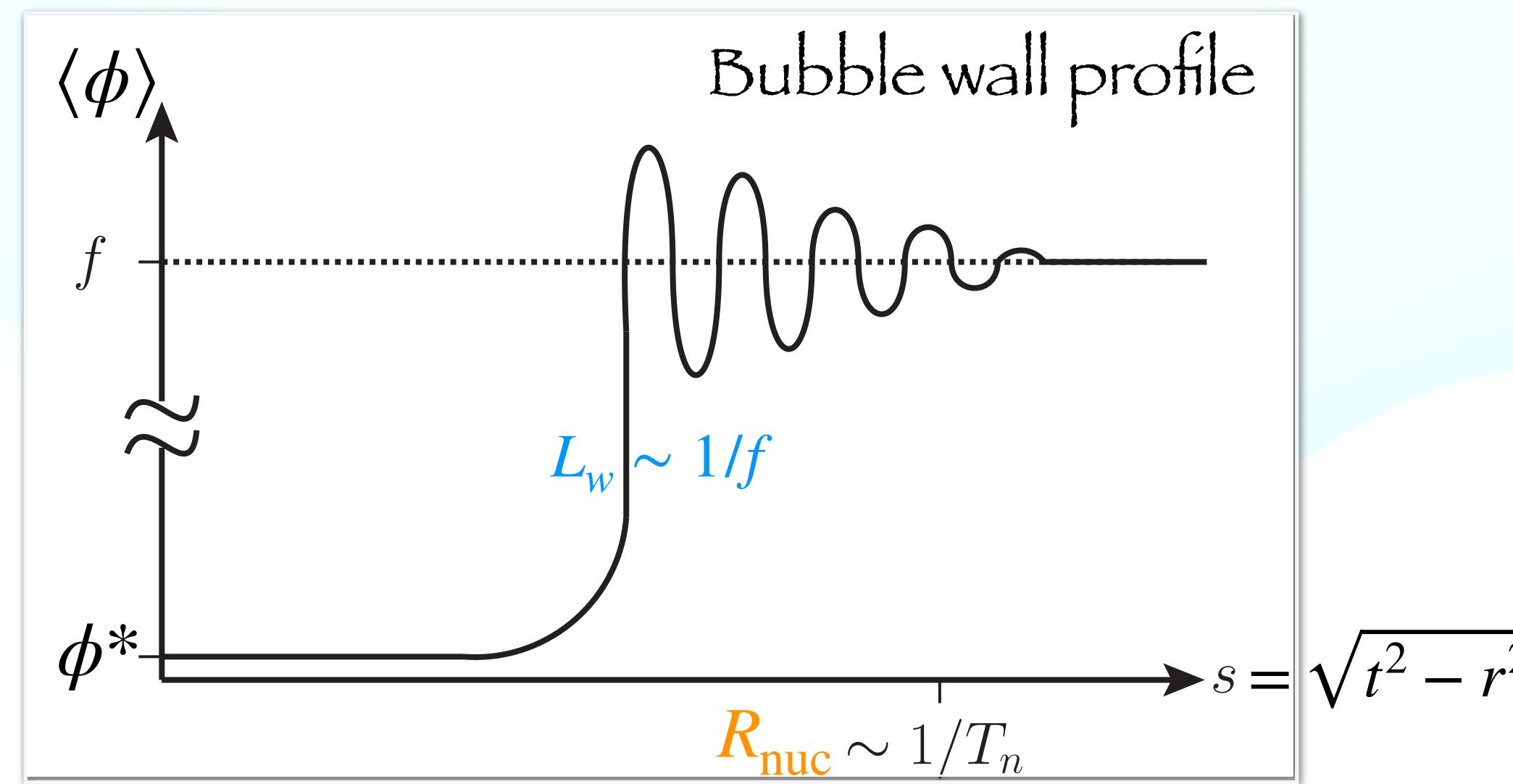
If no pressure then walls run away $\gamma(t) \sim \frac{R_{\text{bubble}}(t)}{R_{\text{nuc}}} \ggg 1$

$$E_{\text{bubble}} \sim R^3 f^4$$

$$E_{\text{wall}} \sim \gamma R^2 \sigma \quad \sigma \sim R_{\text{nuc}} f^4$$



Bubble walls



If no pressure then walls run away $\gamma(t) \sim \frac{R_{\text{bubble}}(t)}{R_{\text{nuc}}} \ggg 1$

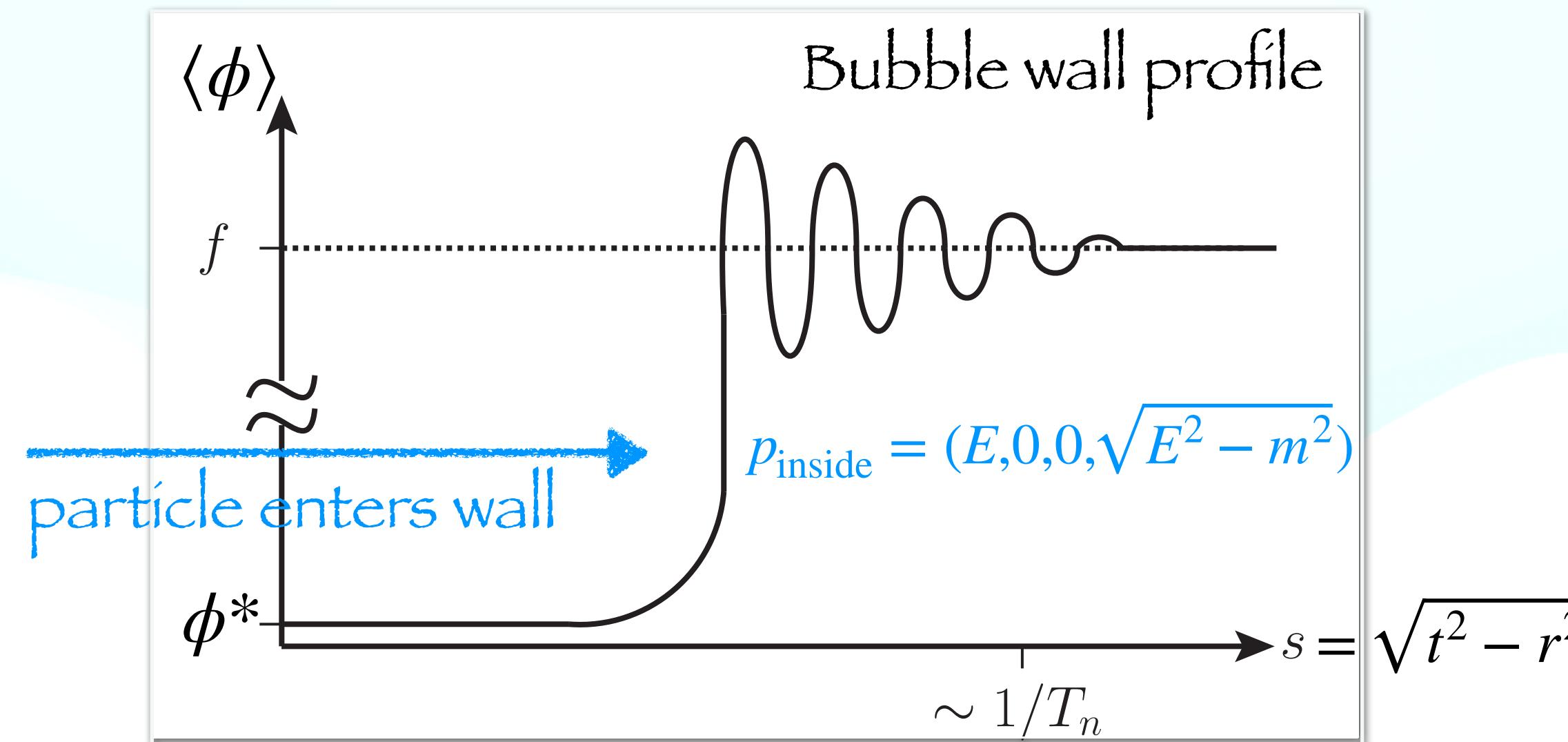
Pressure from plasma effects: Negligible if $T_{\text{nuc}}/f < O(1)$

from $\gamma v/L_w \gg \Gamma_{\text{coll}} \sim \alpha^2 T$

Bubble walls

$$p_{\text{outside}} = (E, 0, 0, E)$$

wall frame : $E \sim \gamma T_{\text{nuc}}$



If no pressure then walls run away $\gamma(t) \sim \frac{R_{\text{bubble}}(t)}{R_{\text{nuc}}} \ggg 1$

Pressure from plasma effects: Negligible if $T_{\text{nuc}}/f < O(1)$

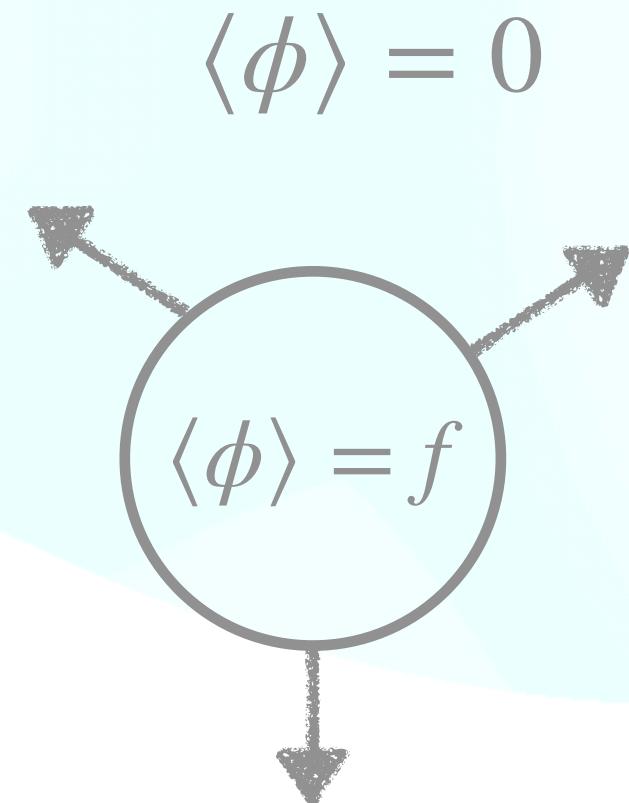
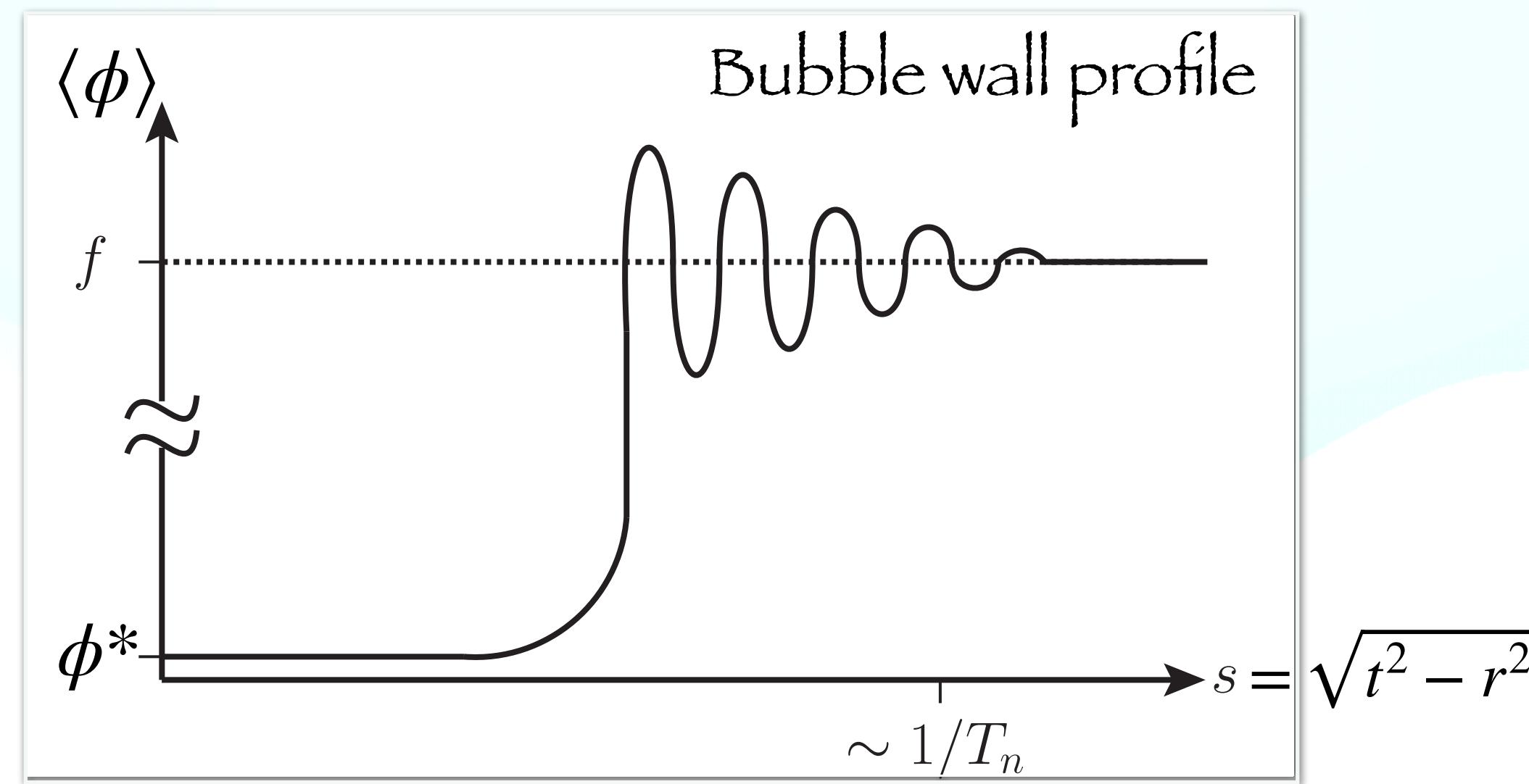
Pressure from single particles: Negligible if $T_{\text{nuc}}/f < O(1)$

$$P_{\text{particle}} = \int \frac{d^3 p}{(2\pi)^3} f(p) \Delta p \sim n \frac{m^2}{E} \sim g^2 f^2 T_{\text{nuc}}^2$$

$$P_{\text{bubble}} = -(V(\phi = f) - V(\phi = 0)) \sim f^4$$

Bodeker Moore 0903.4099

Bubble walls



If no pressure then walls run away $\gamma(t) \sim \frac{R_{\text{bubble}}(t)}{R_{\text{nuc}}} \ggg 1$

Pressure from plasma effects: Negligible if $T_{\text{nuc}}/f < O(1)$

Pressure from single particles: Negligible if $T_{\text{nuc}}/f < O(1)$

Phase Transitions with fast bubble walls

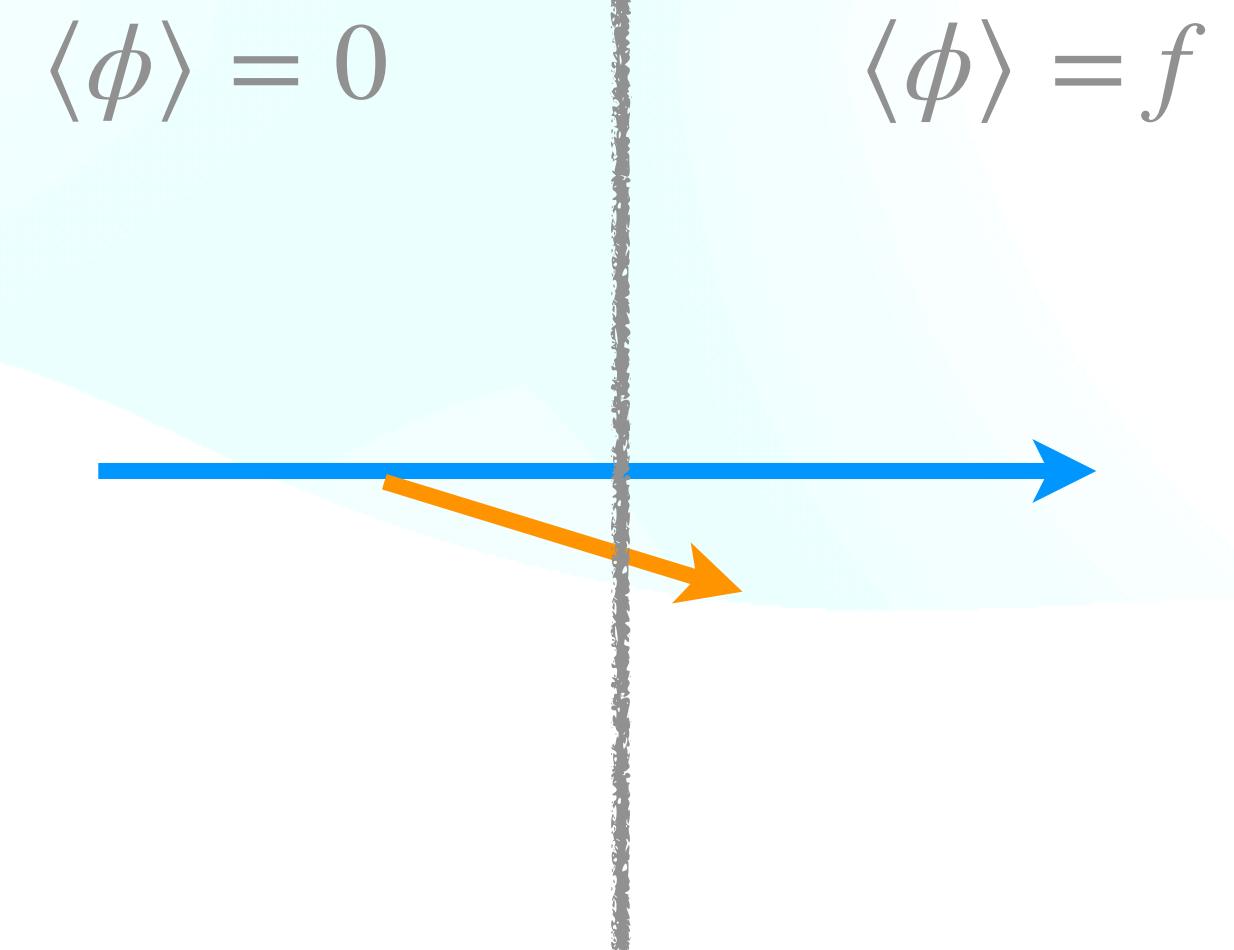
Incoming particles can emit stuff!

Vector emission IR enhanced

$$P_{\text{NLO}} \sim g^3 f T_{\text{nuc}}^3 \gamma \quad P_{\text{bubble}} \sim f^4$$

→ terminal velocity $\gamma \sim \left(\frac{f}{T_{\text{nuc}}}\right)^3$

Bodeker Moore 1703.08215



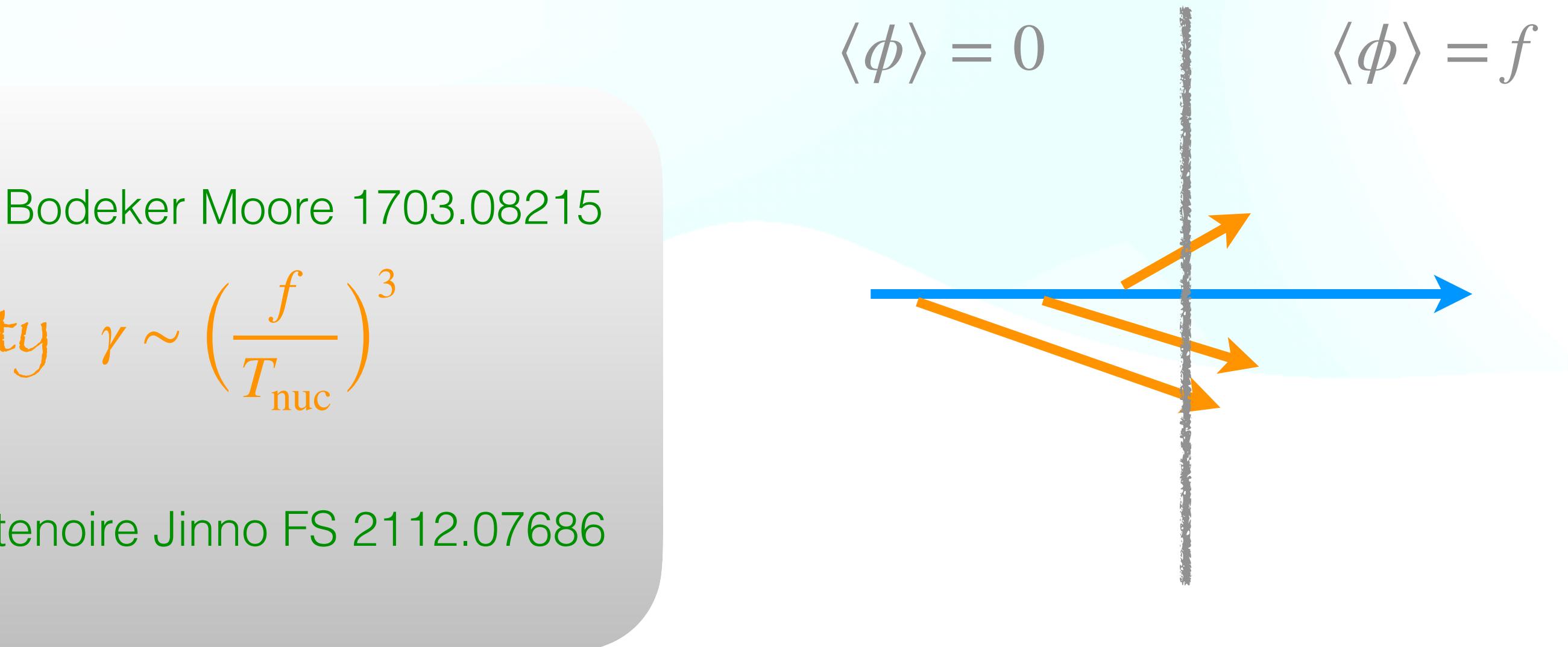
Phase Transitions with fast bubble walls

Incoming particles can emit stuff!

Vector emission IR enhanced

$$P_{\text{NLO}} \sim g^3 f T_{\text{nuc}}^3 \gamma \xrightarrow[P_{\text{bubble}} \sim f^4]{\longrightarrow} \text{terminal velocity } \gamma \sim \left(\frac{f}{T_{\text{nuc}}}\right)^3$$

$$P_{\text{LL}} \simeq g^3 f T_{\text{nuc}}^3 \gamma \times \log \quad \text{when resummed} \quad \text{Gouttenoire Jinno FS 2112.07686}$$



Phase Transitions with fast bubble walls

Incoming particles can emit stuff!

Vector emission IR enhanced

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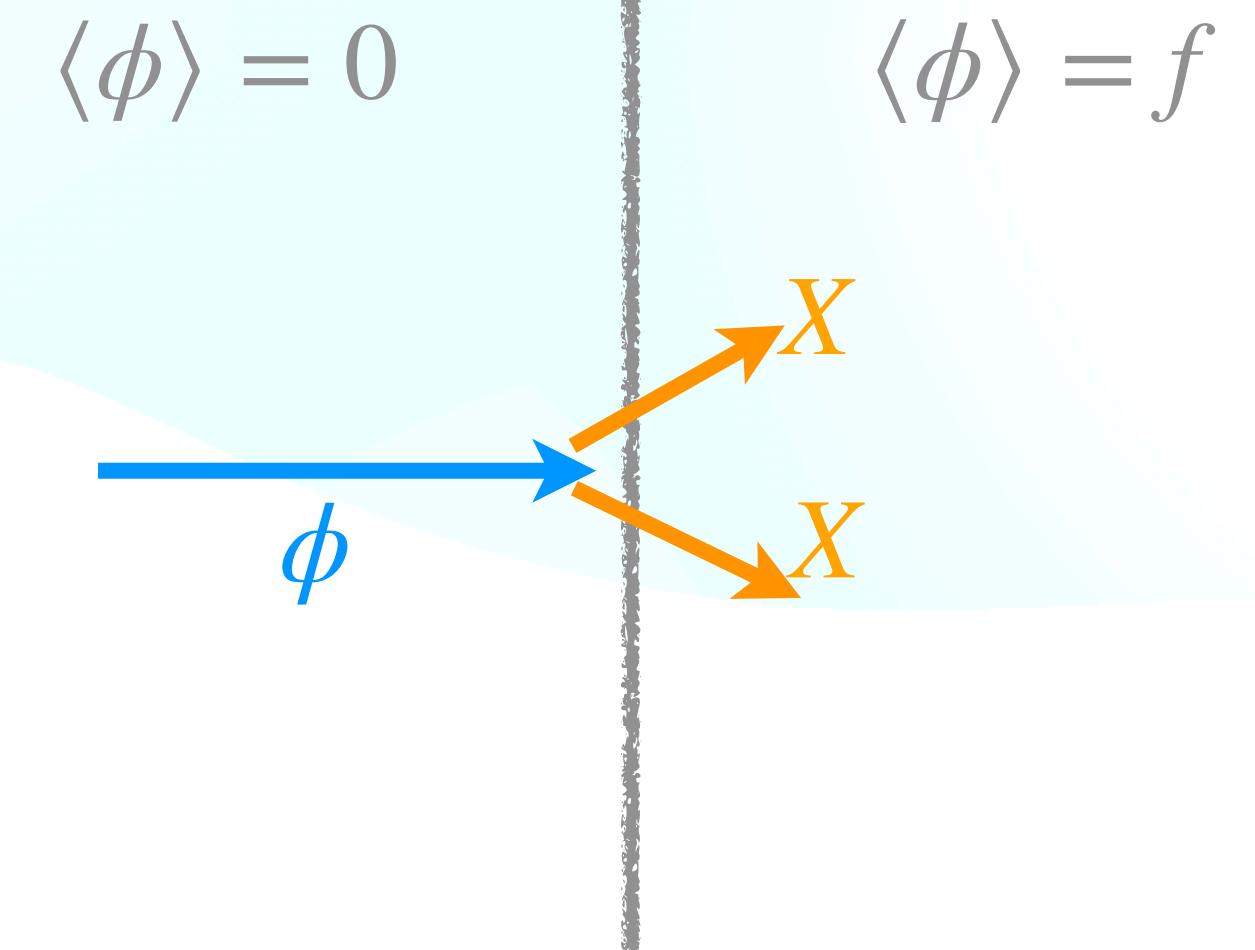
Bodeker Moore 1703.08215

Pressure from heavy states $M \gg f$ produced if $\gamma T_{\text{nuc}} \gg M$

$$\text{E.g. } \mathcal{L} \supset -\frac{M^2}{2} X^2 - \lambda \phi^2 X^2 \quad \text{Prob}(\phi \rightarrow XX) \sim \frac{\lambda^2}{16\pi^2} \frac{f^2}{M^2}$$

$$P_{\text{heavy}} \sim n_\phi \text{Prob}(\phi \rightarrow XX) \frac{M^2}{E} \sim \frac{\lambda}{16\pi^2} \frac{\lambda}{g^2} P_{\text{LO}}$$

Azatov Vanvlasselaer 2010.02590



Phase Transitions...

...as Particle Accelerators (“Bubbletrons”)

Warm EeV Dark Matter from Weak scale PT

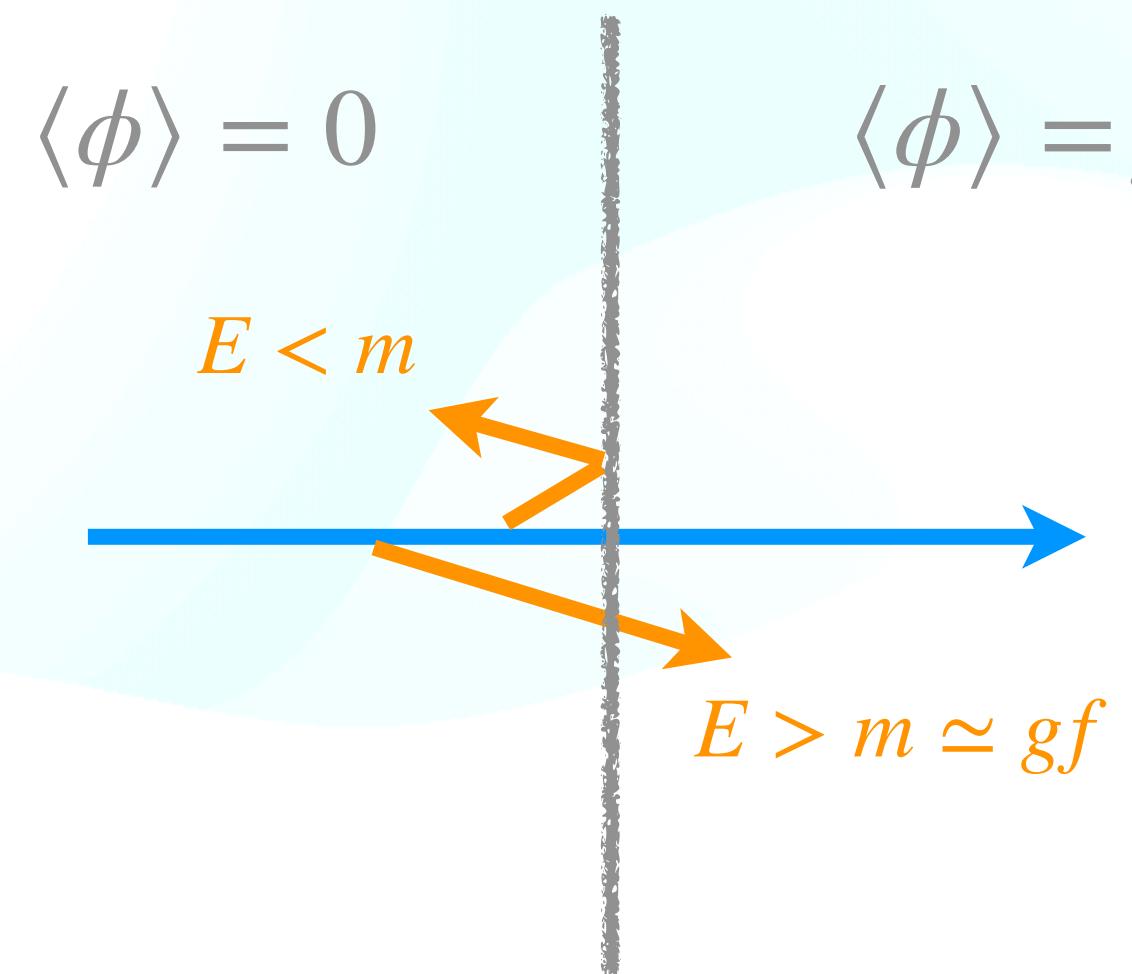
Phase Transitions...

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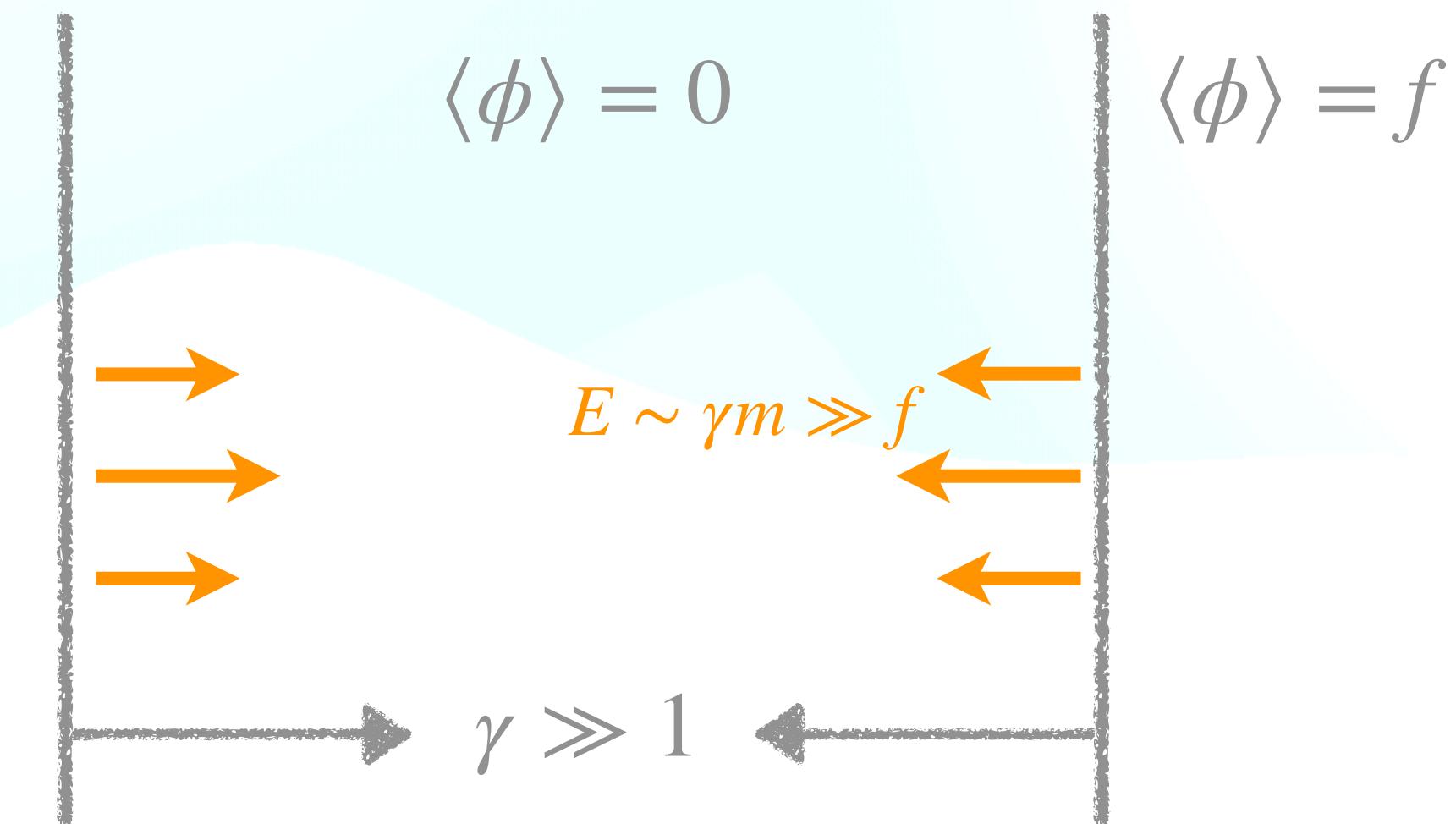
Warm EeV Dark Matter from Weak scale PT

Radiation at the Bubble Walls

wall frame



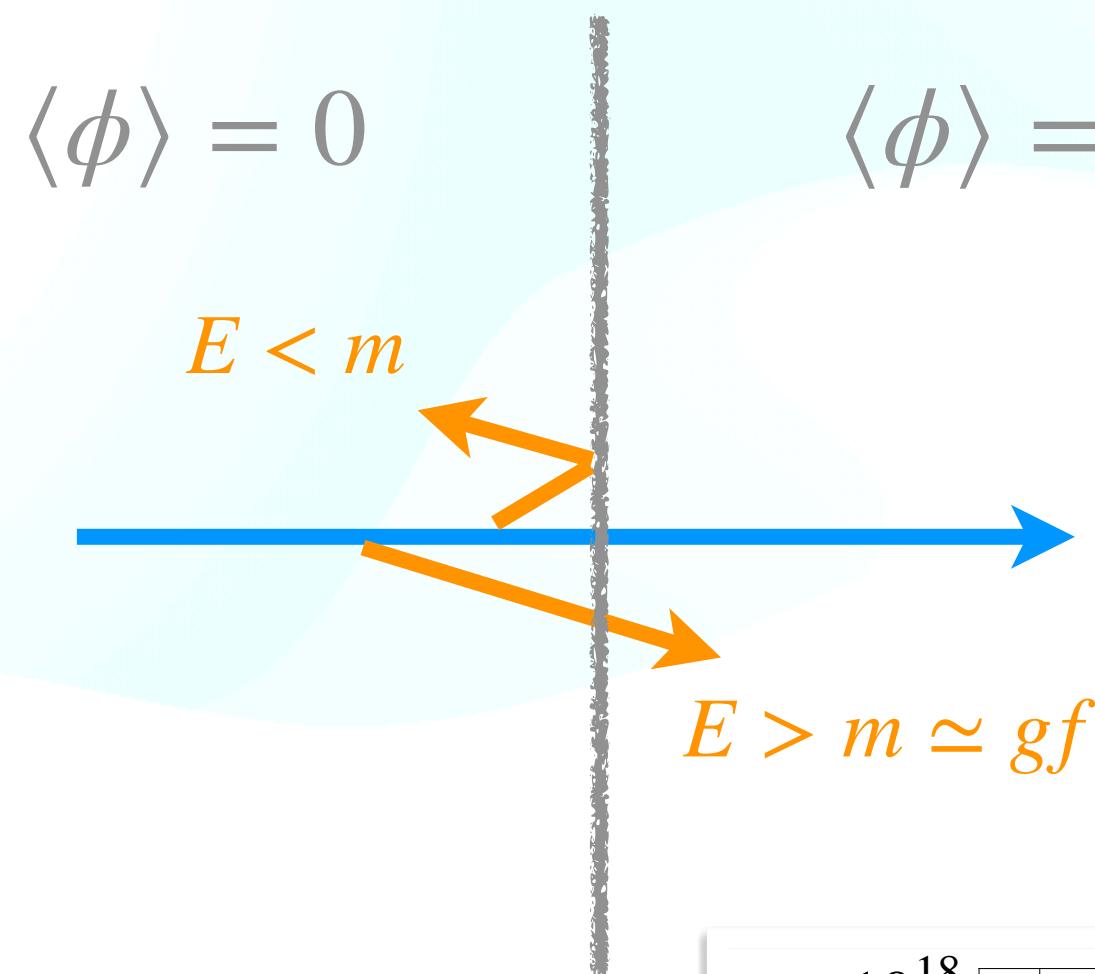
plasma (= CMB) frame



Soft radiation is reflected
fast walls are a particle accelerator!

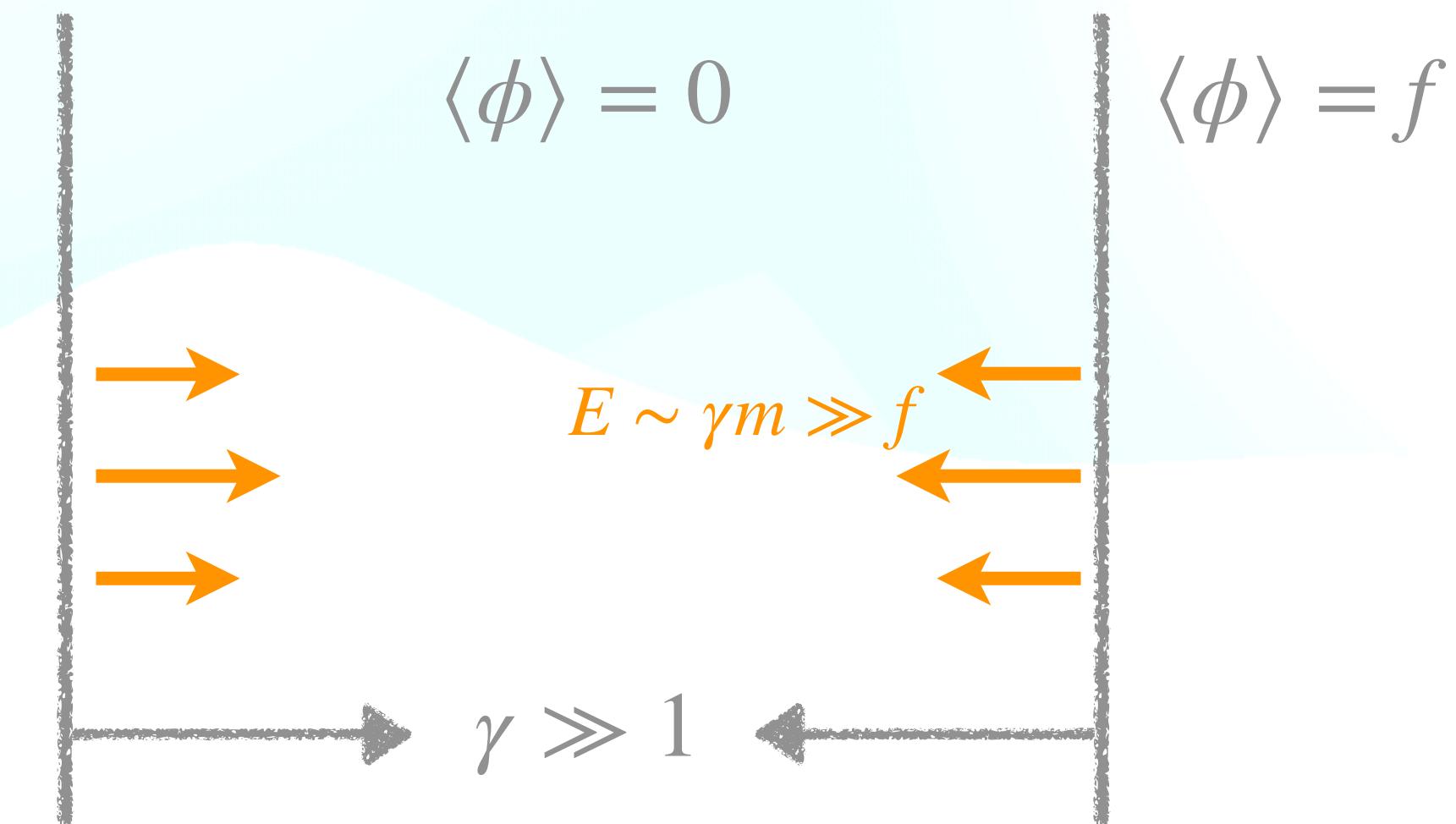
Radiation at the Bubble Walls

wall frame

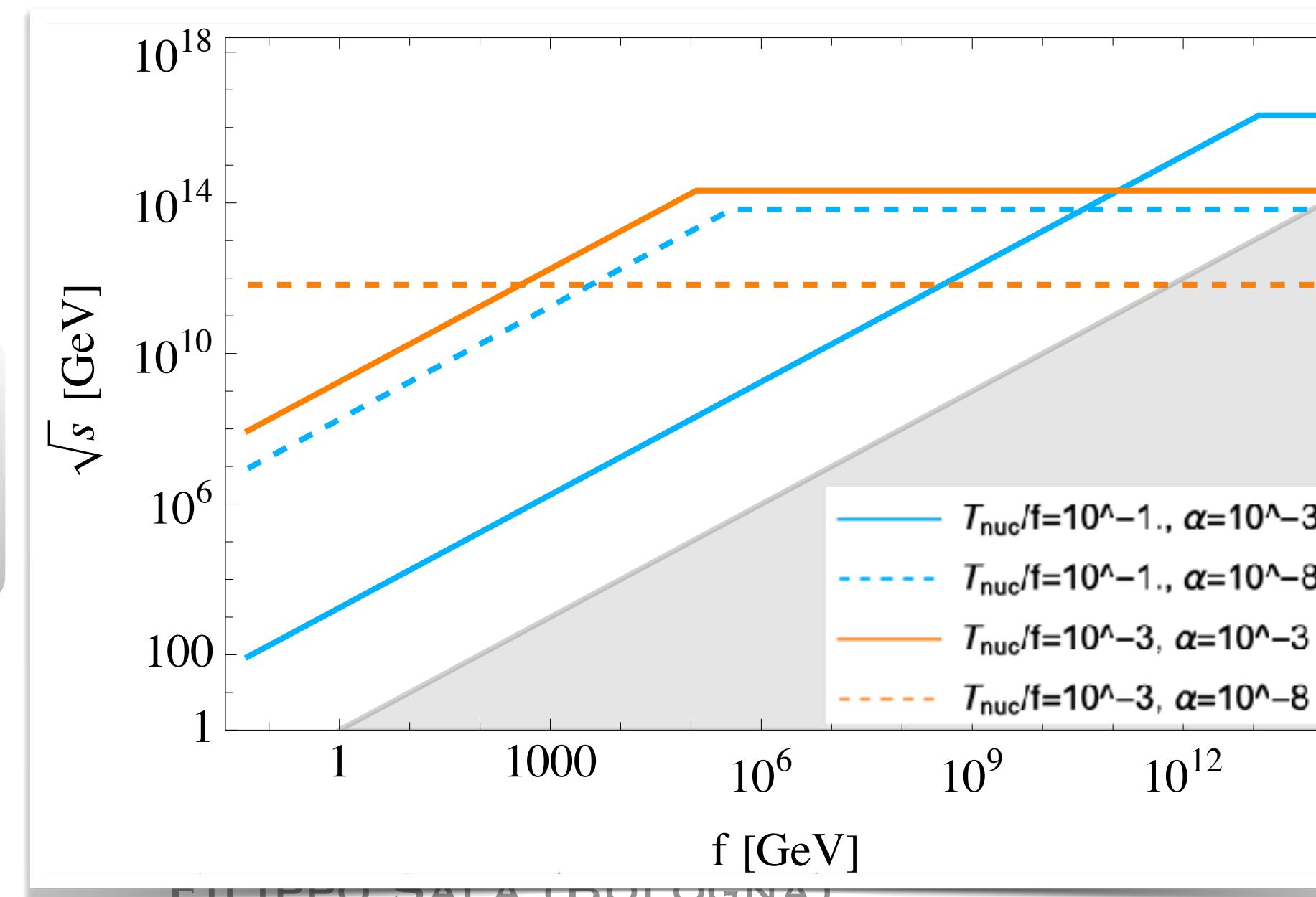


Soft radiation is reflected
fast walls are a particle accelerator!

plasma (= CMB) frame

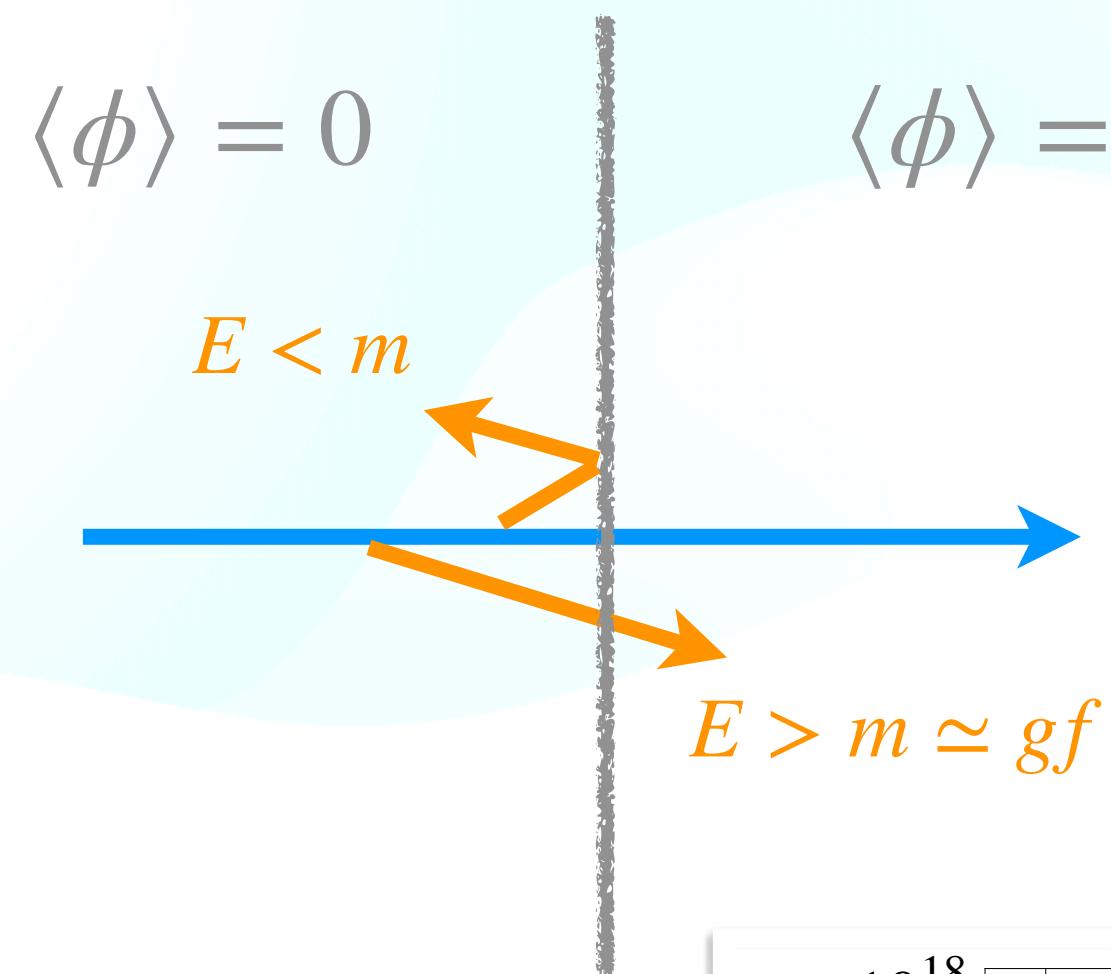


Radiation collide
before walls collide!



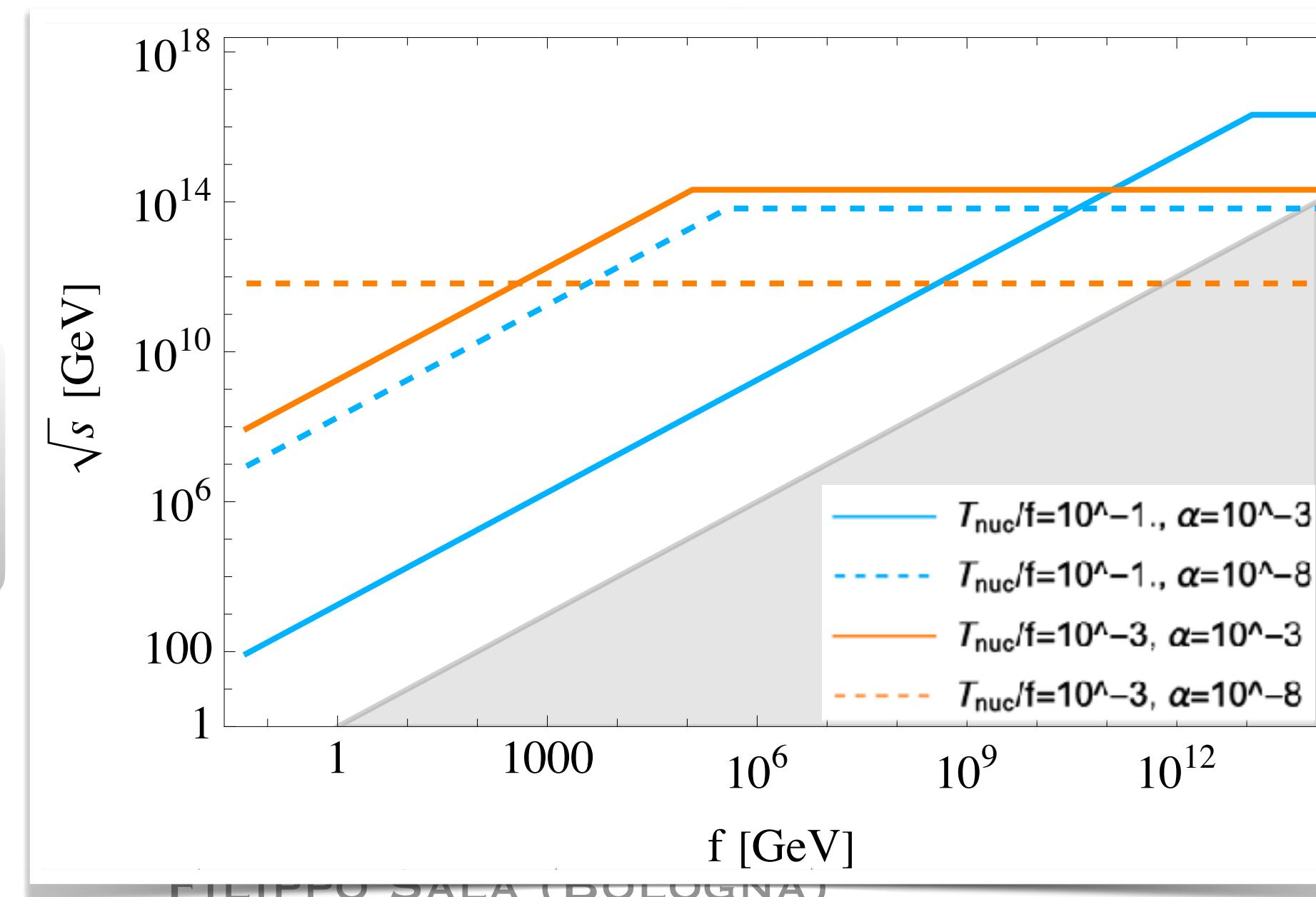
Radiation at the Bubble Walls

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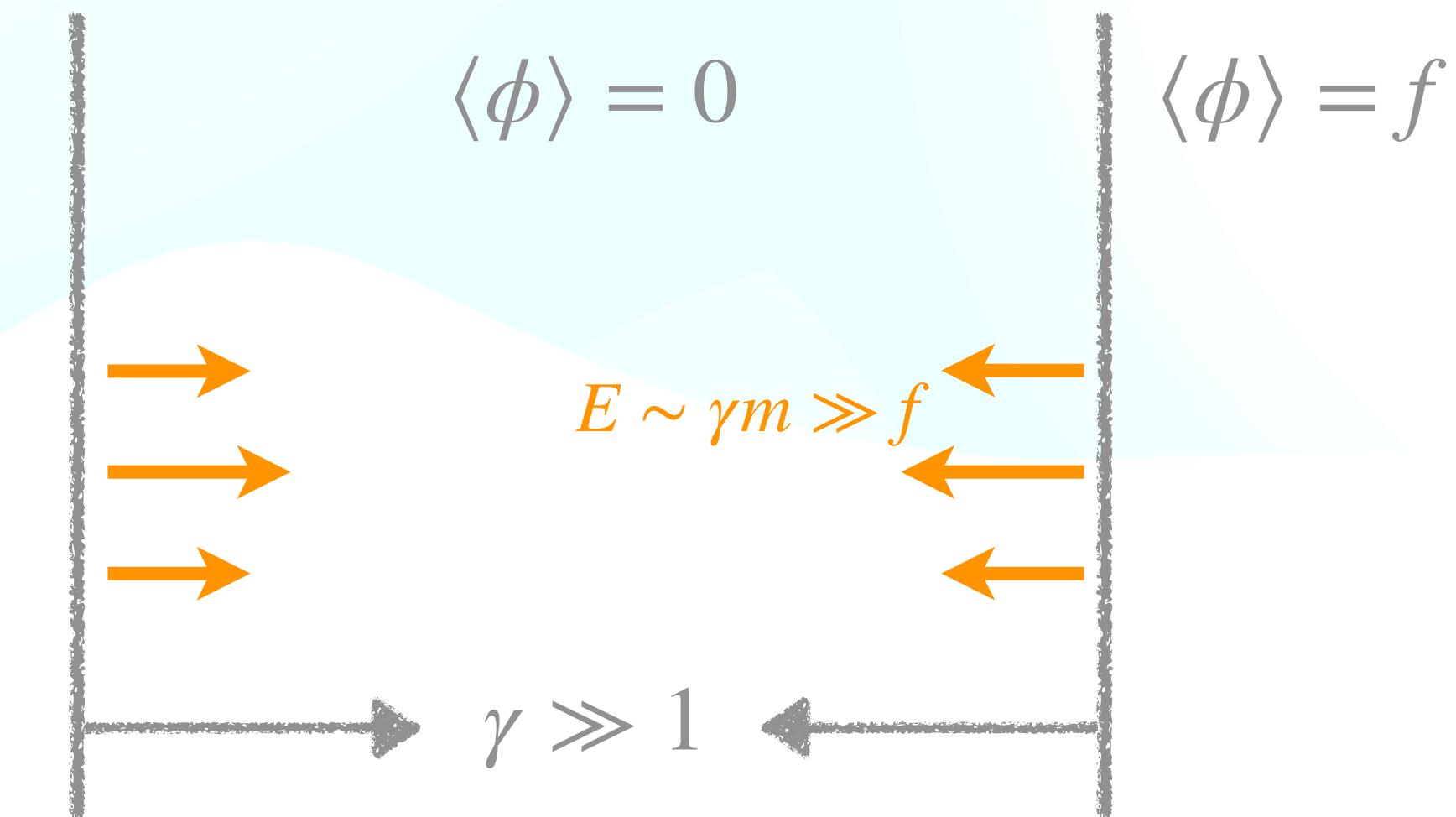


Soft radiation is reflected
fast walls are a **particle accelerator!**

Radiation collide
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plasma (= CMB) frame

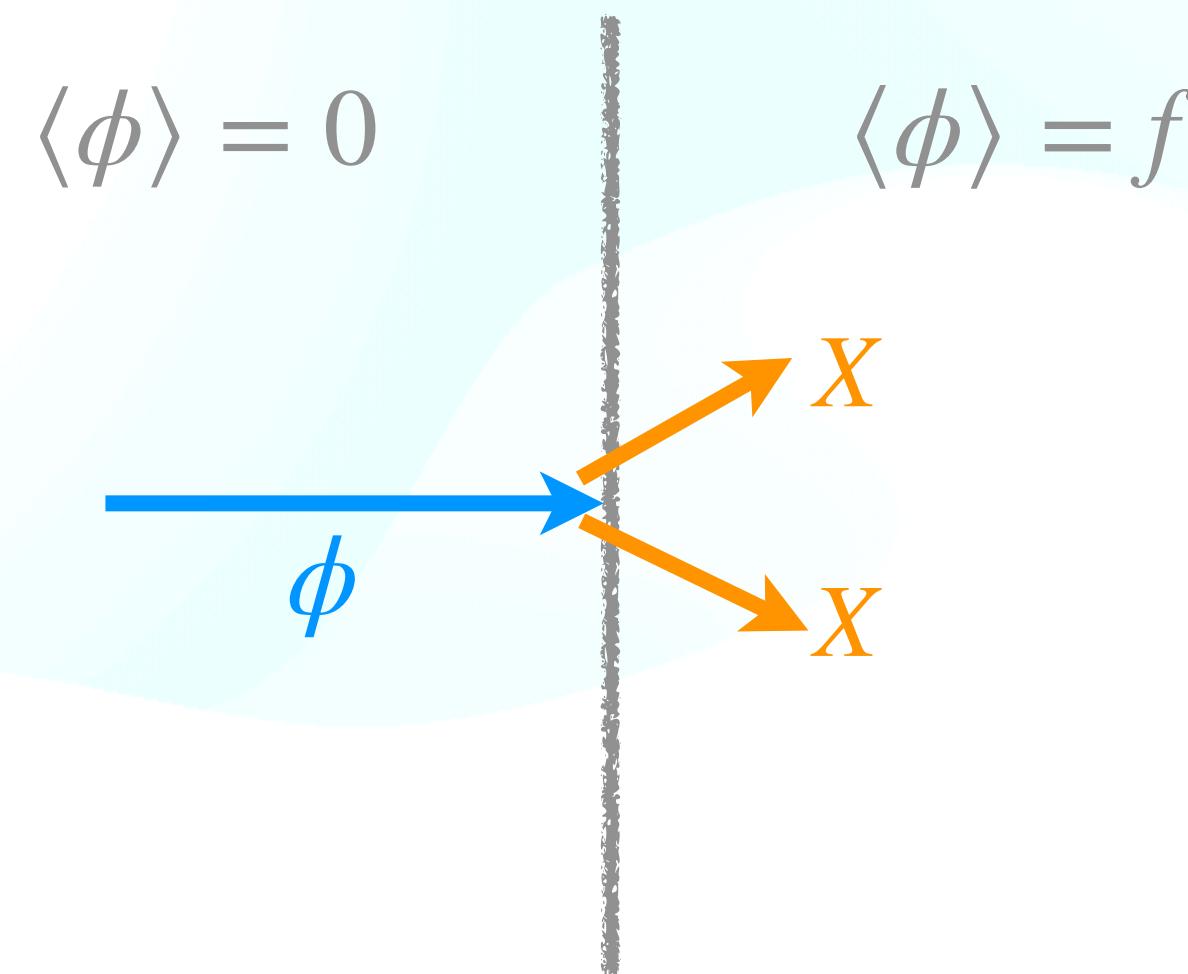


But: ? Shells scatterings with bath particles:
momentum loss: wall catches up,...
number changing: mess

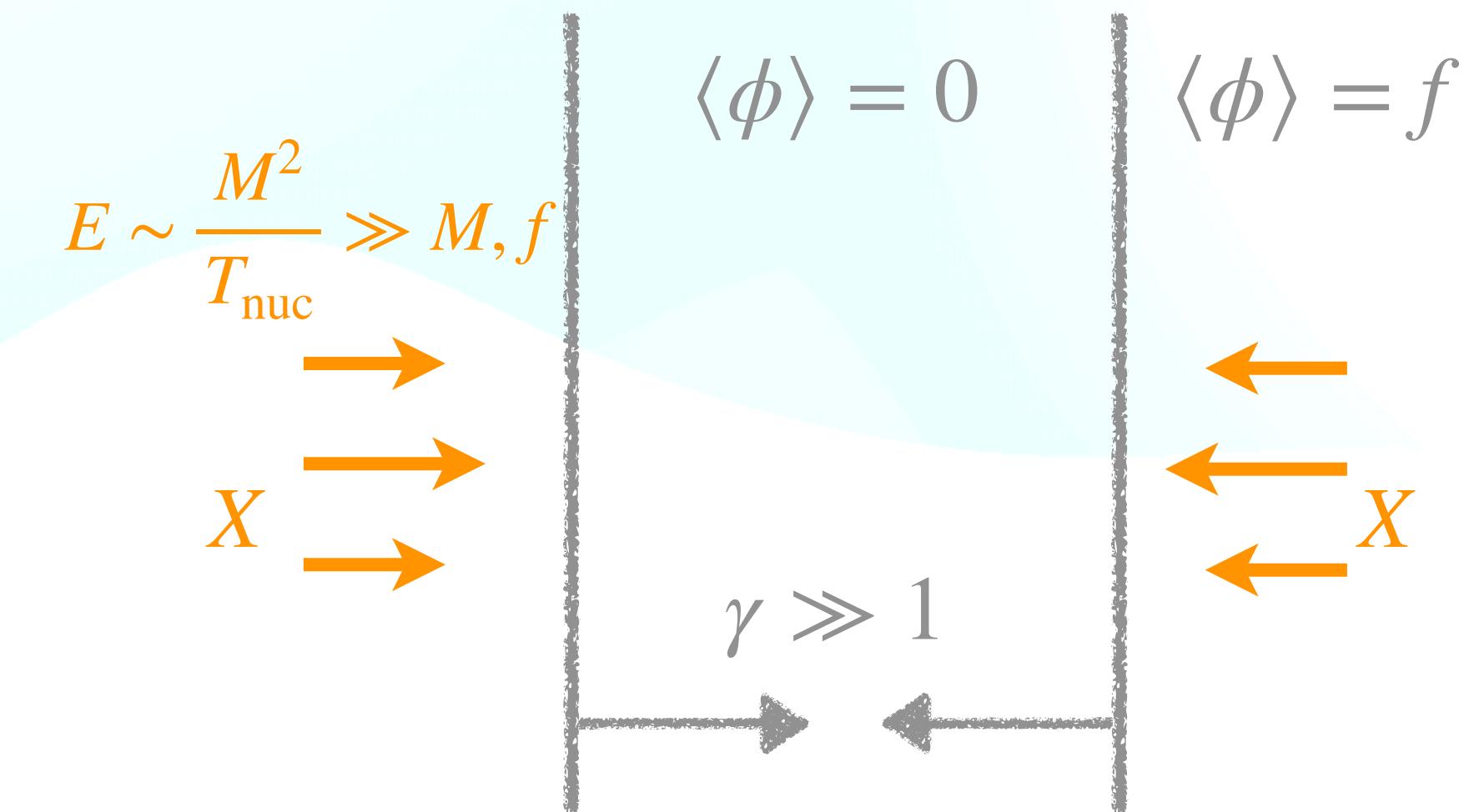
? Shells thermalisation

Heavier Physics

wall frame



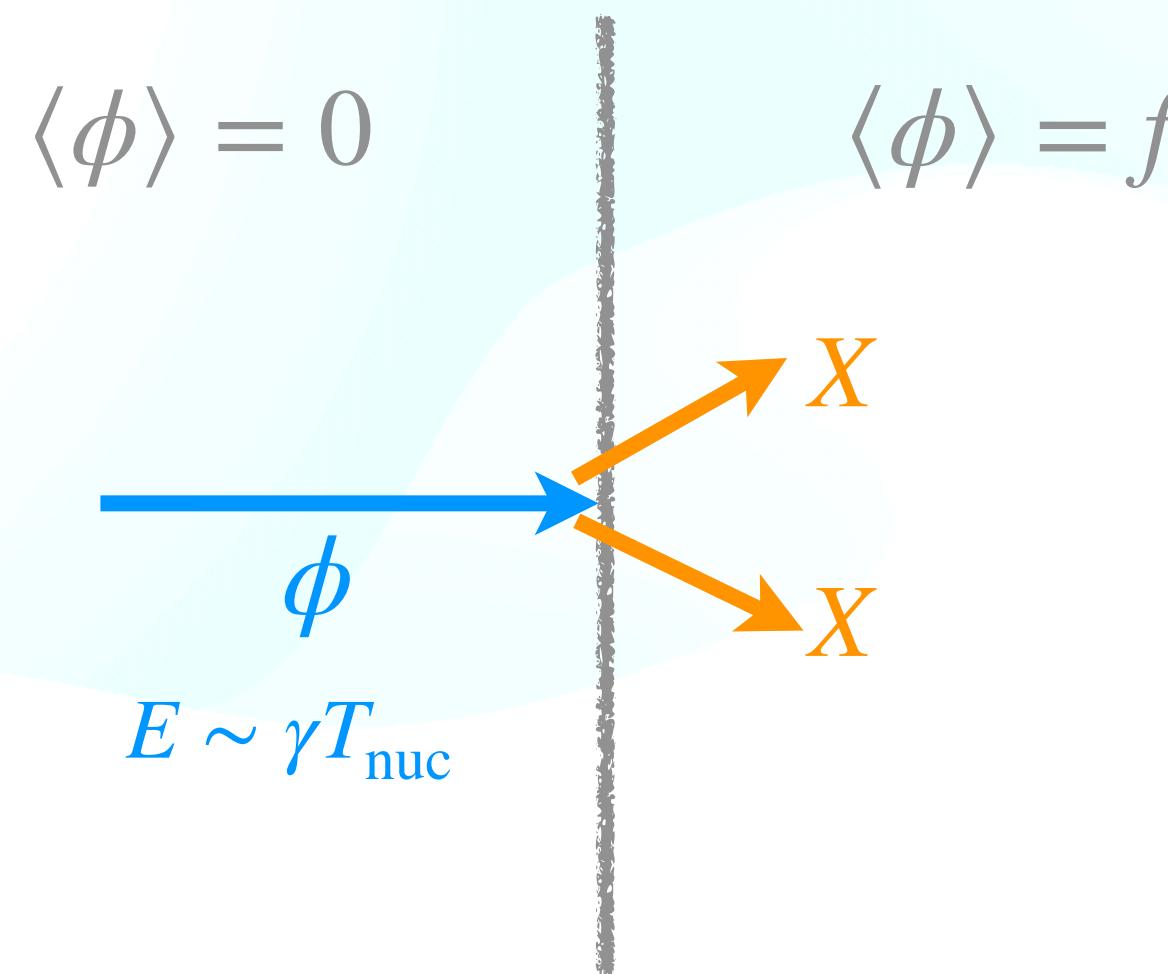
plasma (= CMB) frame



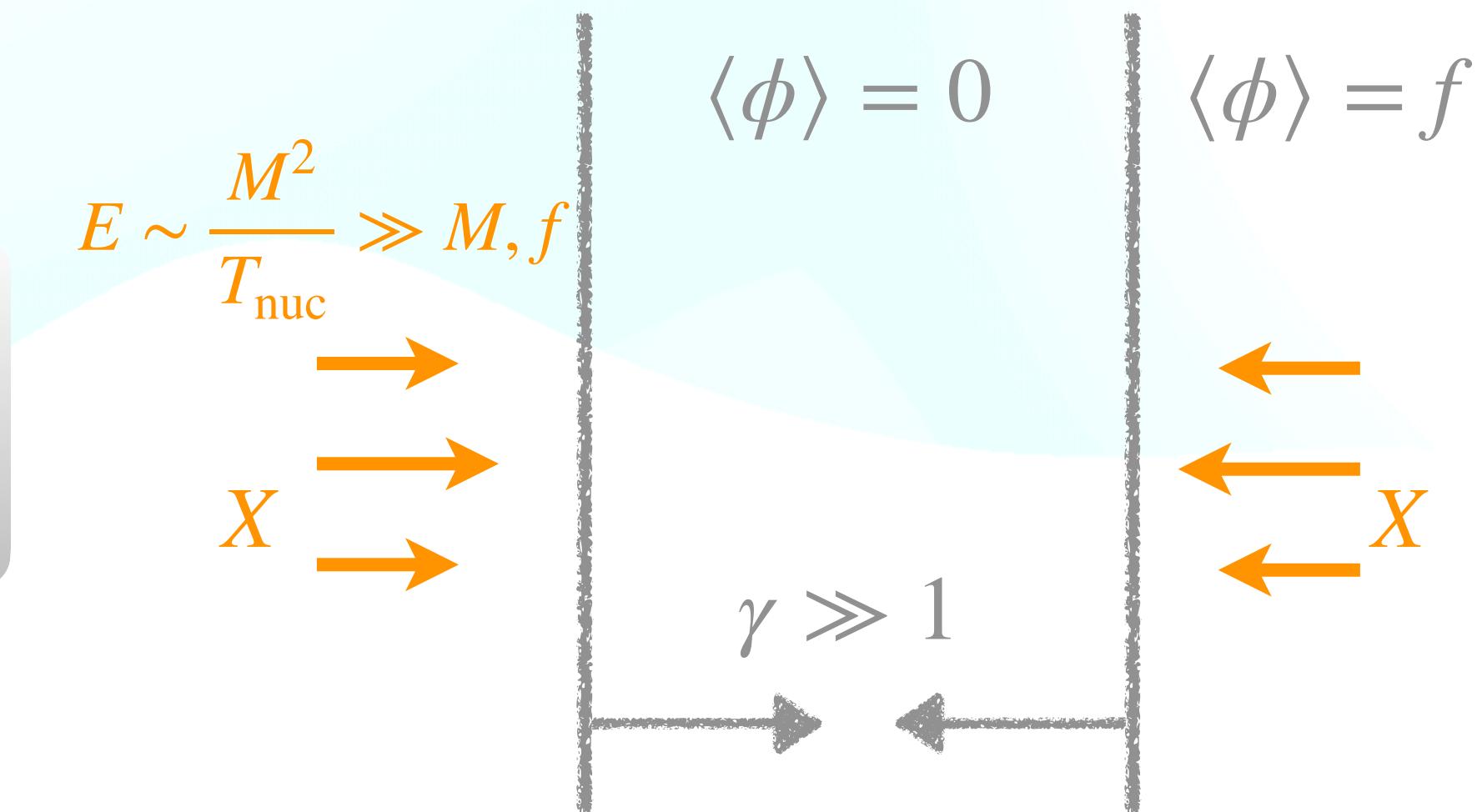
Heavy States follow the walls
If fast, again particle accelerator!

Heavier Physics

wall frame



plasma (= CMB) frame



Heavy States follow the walls
If fast, again particle accelerator!

But: ? Shells scatterings with bath particles

Momentum loss: wall catches up,...

Number changing: mess

? Shells thermalisation

? Shells scattering with wall

Bubbletrons

(Brand by Yann Gouttenoire)



Accelerated Particles

Radiated vectors

Radiated scalars

Heavier states

Dark hadrons

Baldes+ 2007.08440

Wall collisions

Falkowski No 1211.5615

Wall Decays

Gauged PT

Global PT

Shell-wall

...



...

Phase Transitions...

...as Particle Accelerators (“Bubbletrons”)

Warm EeV Dark Matter from Weak scale PT

Based on Baldes Gouttenoire FS 2207.05096 (SciPost Phys.)

Phase Transit

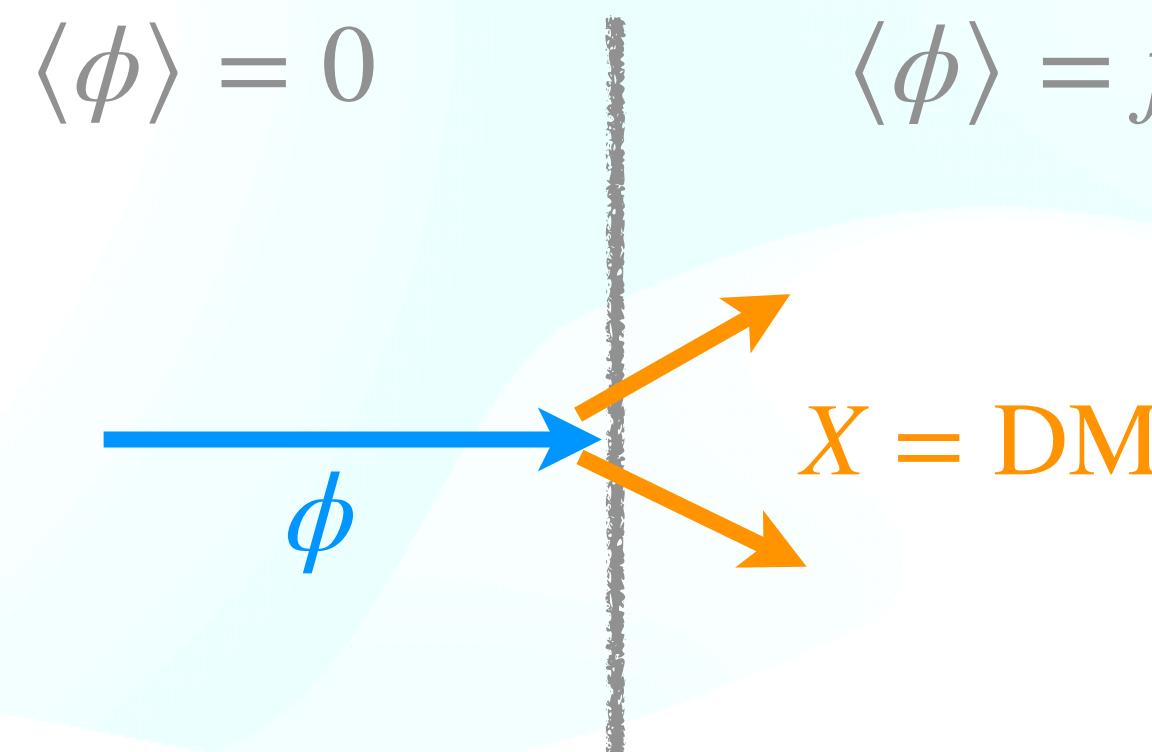
...as Particle



Warm EeV Dark Matter from Weak scale PT

Warm EeV Dark Matter from the Weak Scale

wall frame

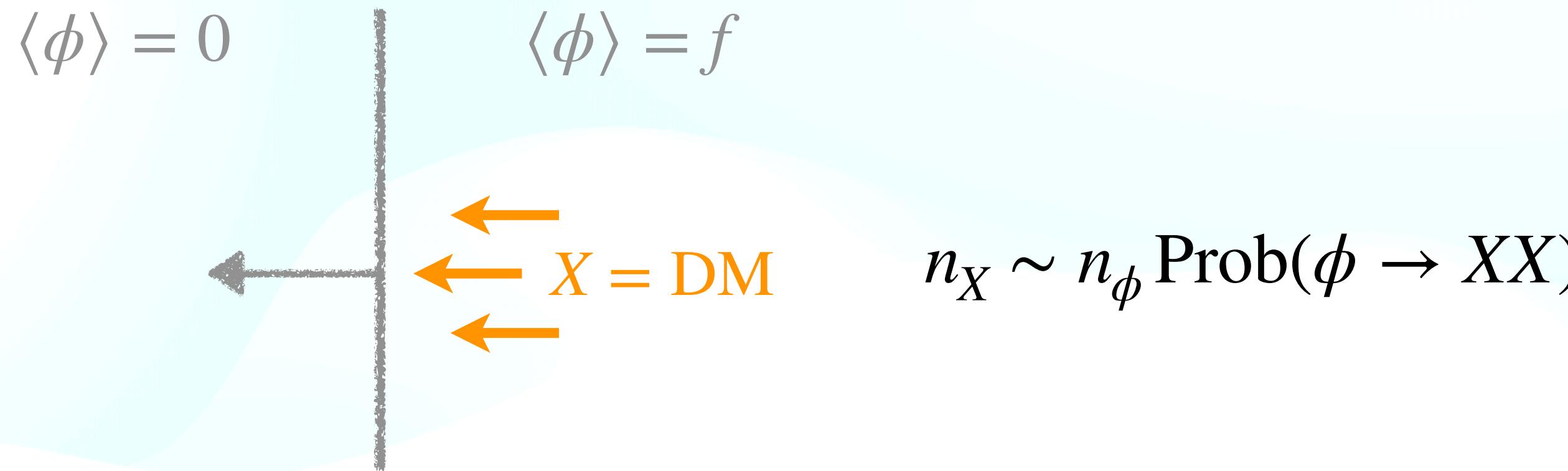


Azatov Vanvlasselaer Yin 2101.05721

$$n_X \sim n_\phi \text{Prob}(\phi \rightarrow XX)$$

Warm EeV Dark Matter from the Weak Scale

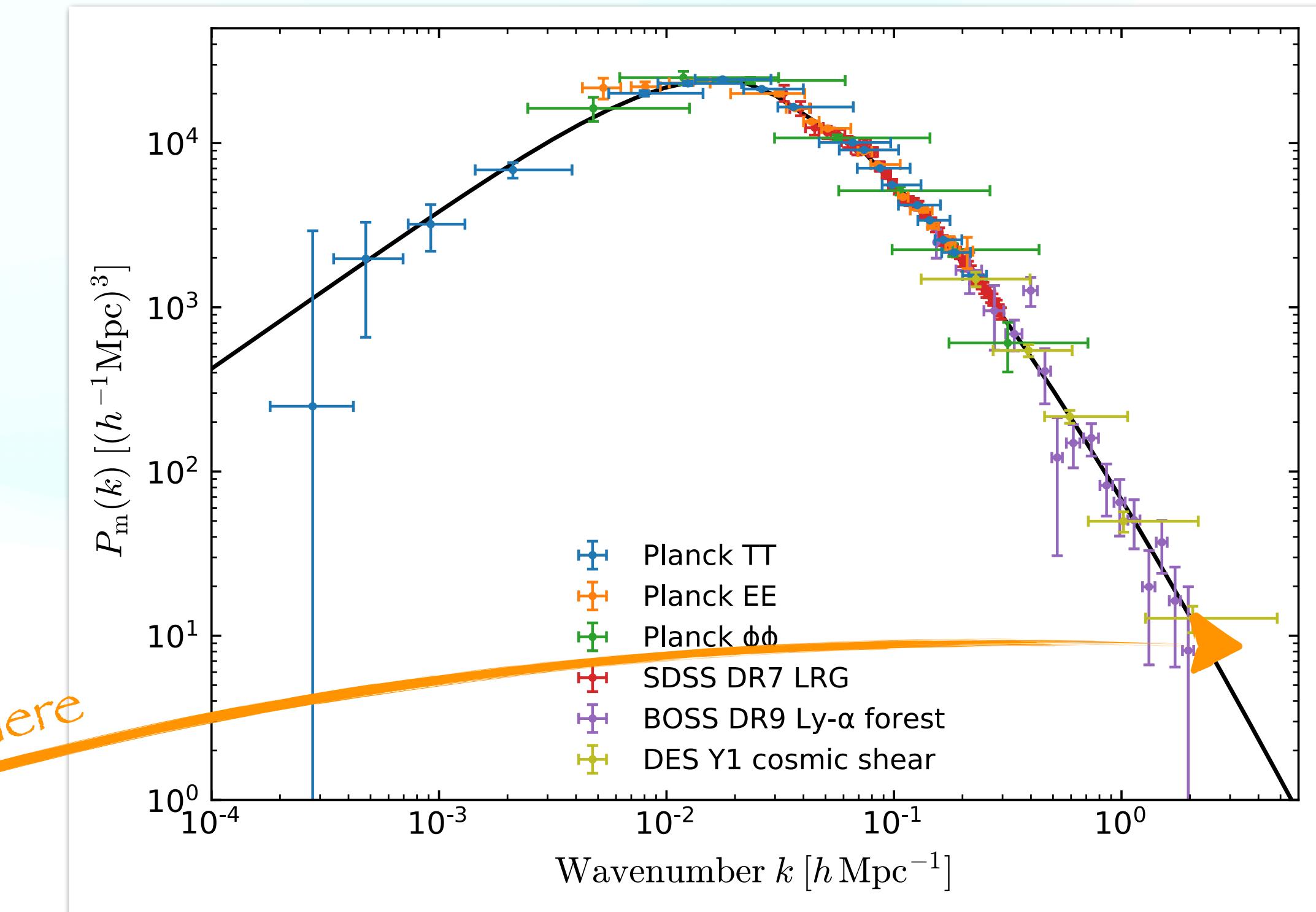
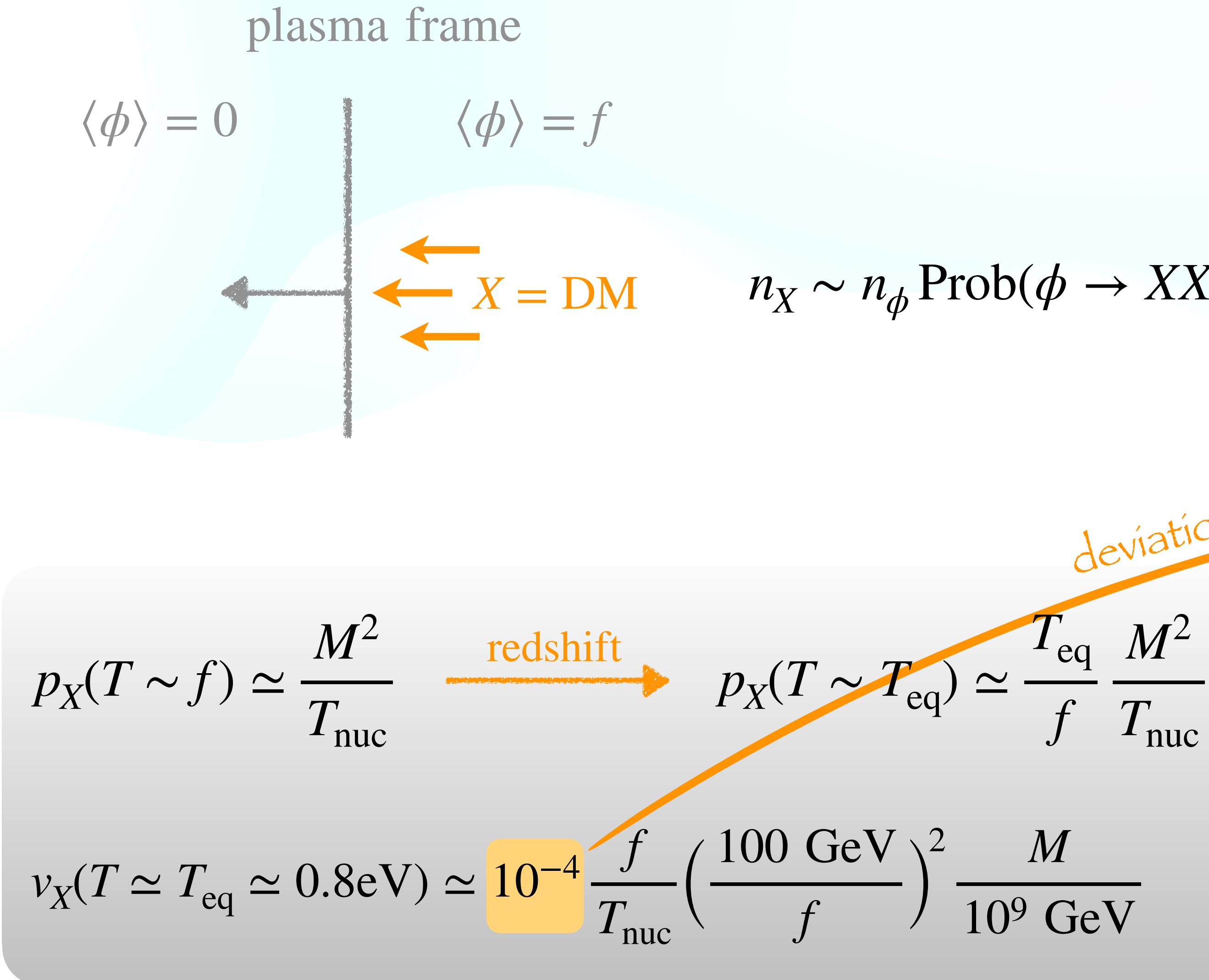
plasma frame



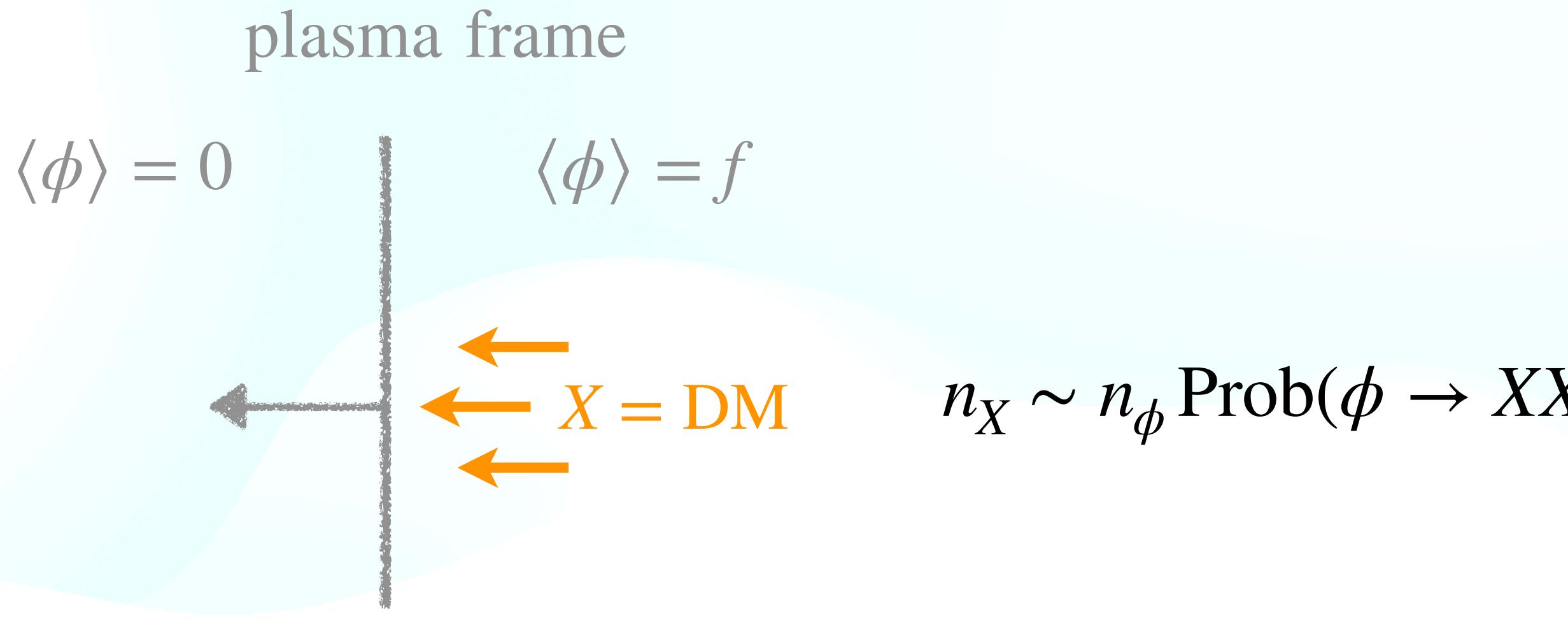
$$p_X(T \sim f) \simeq \frac{M^2}{T_{\text{nuc}}} \xrightarrow{\text{redshift}} p_X(T \sim T_{\text{eq}}) \simeq \frac{T_{\text{eq}}}{f} \frac{M^2}{T_{\text{nuc}}}$$

$$\nu_X(T \simeq T_{\text{eq}} \simeq 0.8 \text{eV}) \simeq 10^{-4} \frac{f}{T_{\text{nuc}}} \left(\frac{100 \text{ GeV}}{f} \right)^2 \frac{M}{10^9 \text{ GeV}}$$

Warm EeV Dark Matter from the Weak Scale



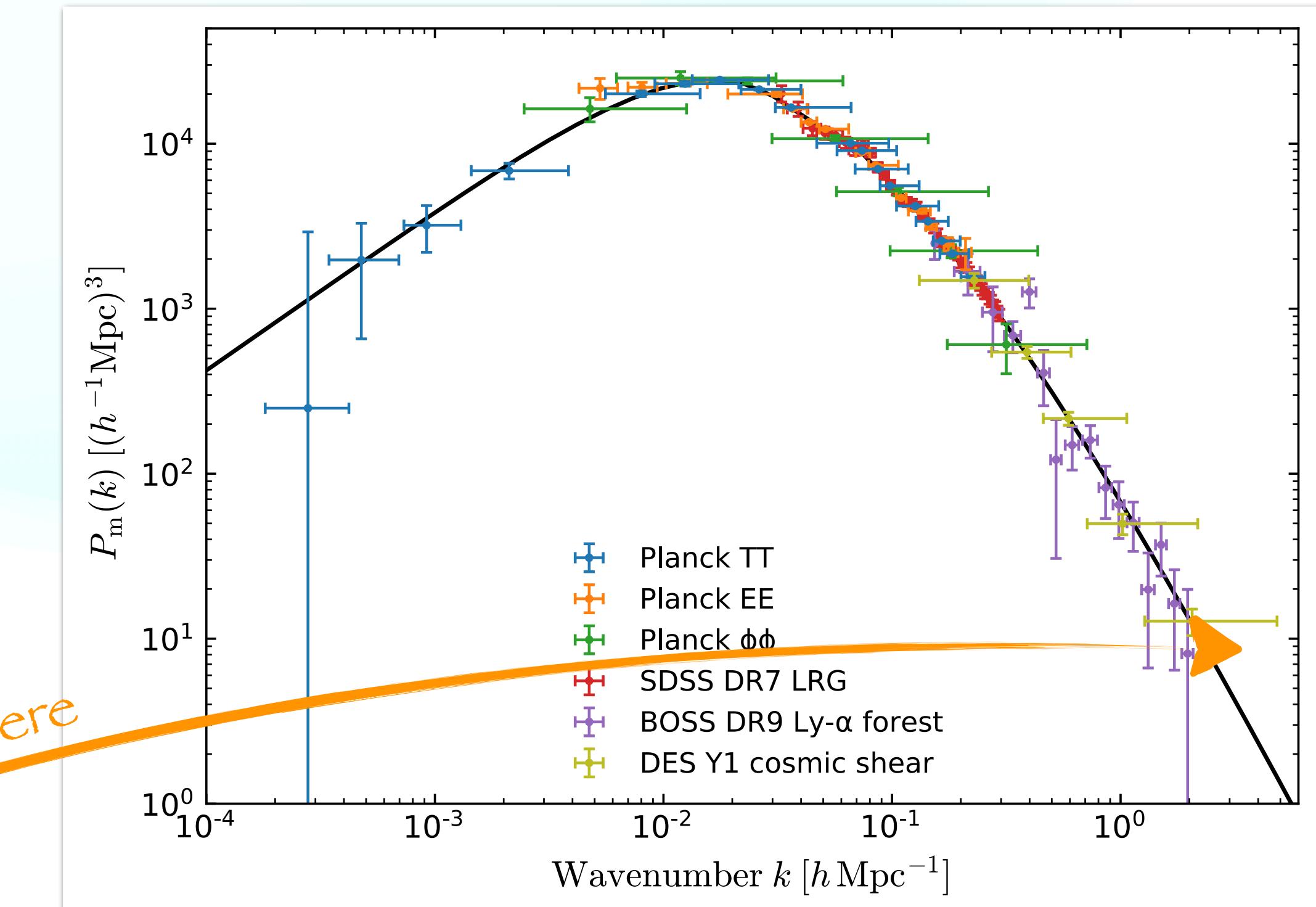
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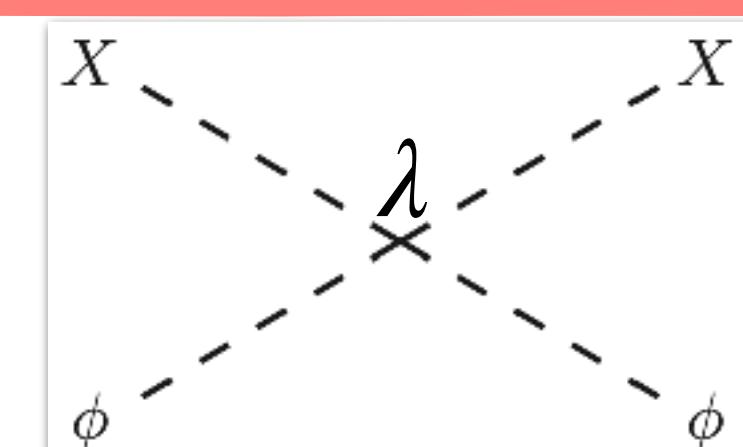
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deviations here

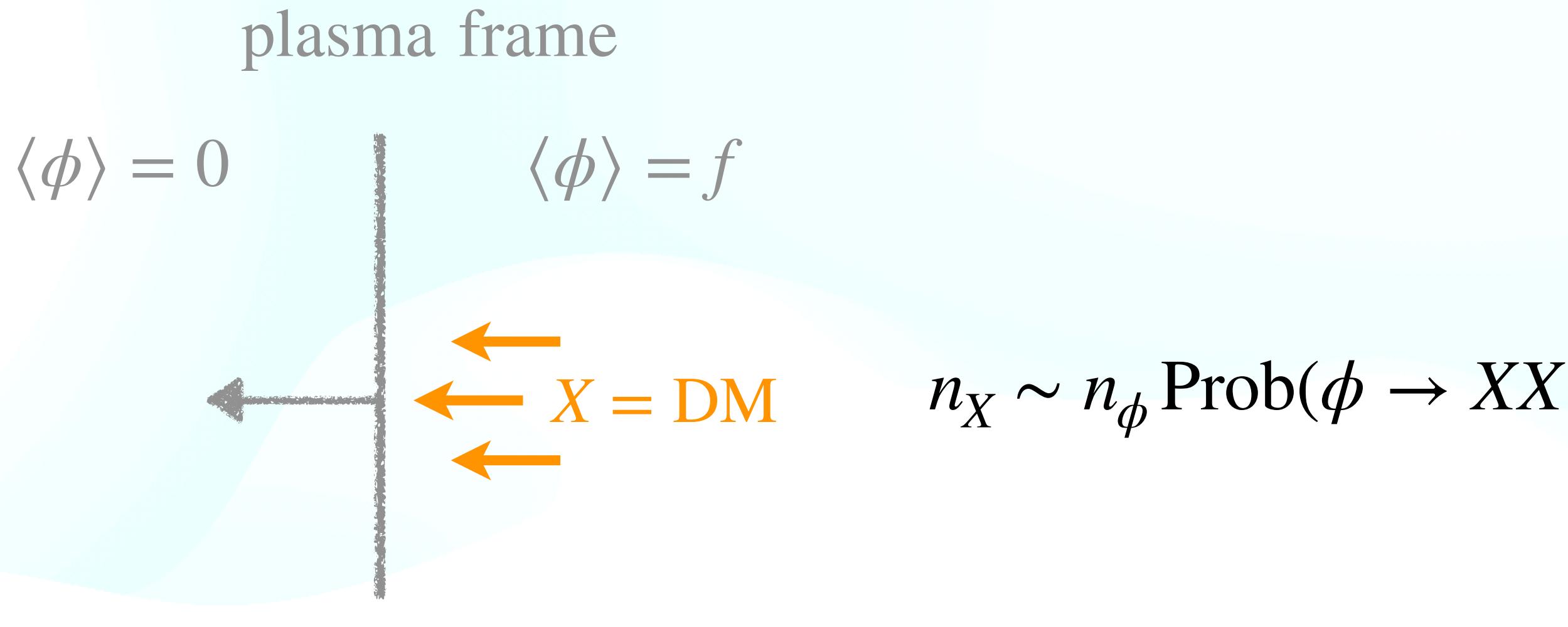


X scatterings with bath



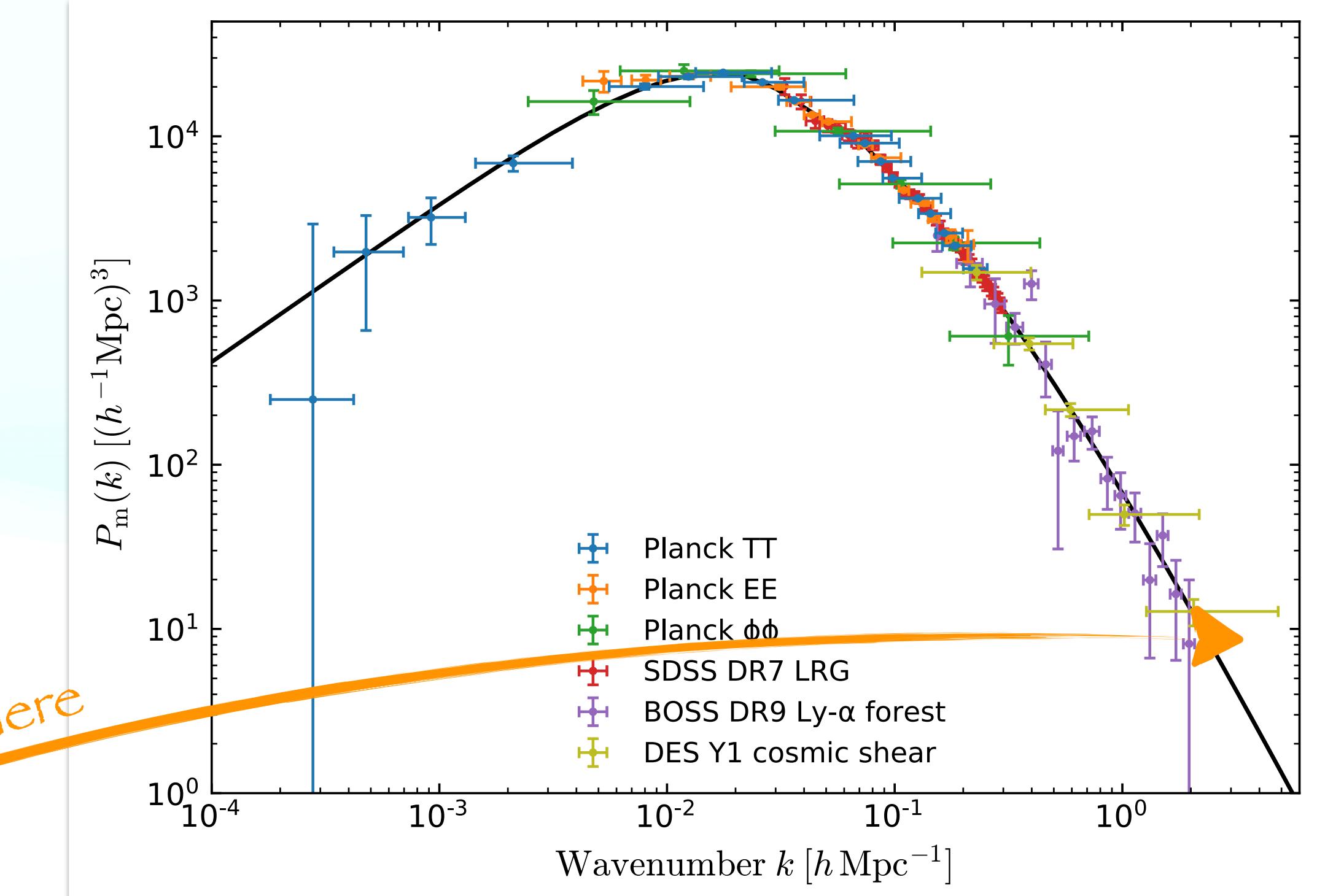
$$\frac{1}{p_X} \frac{dp_X}{dt} < H \Rightarrow \lambda \lesssim \frac{10^9 \text{ GeV}}{M} \left(\frac{f}{T_{\text{nuc}}} \right)^{\frac{1}{4}} \left(\frac{\text{TeV}}{f} \right)^{\frac{1}{2}}$$

Warm EeV Dark Matter from the Weak Scale



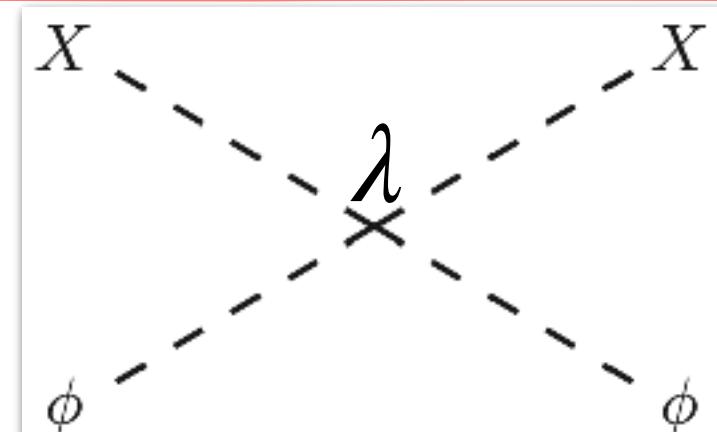
$$p_X(T \sim f) \simeq \frac{M^2}{T_{\text{nuc}}} \xrightarrow{\text{redshift}} p_X(T \sim T_{\text{eq}}) \simeq \frac{T_{\text{eq}}}{f} \frac{M^2}{T_{\text{nuc}}}$$

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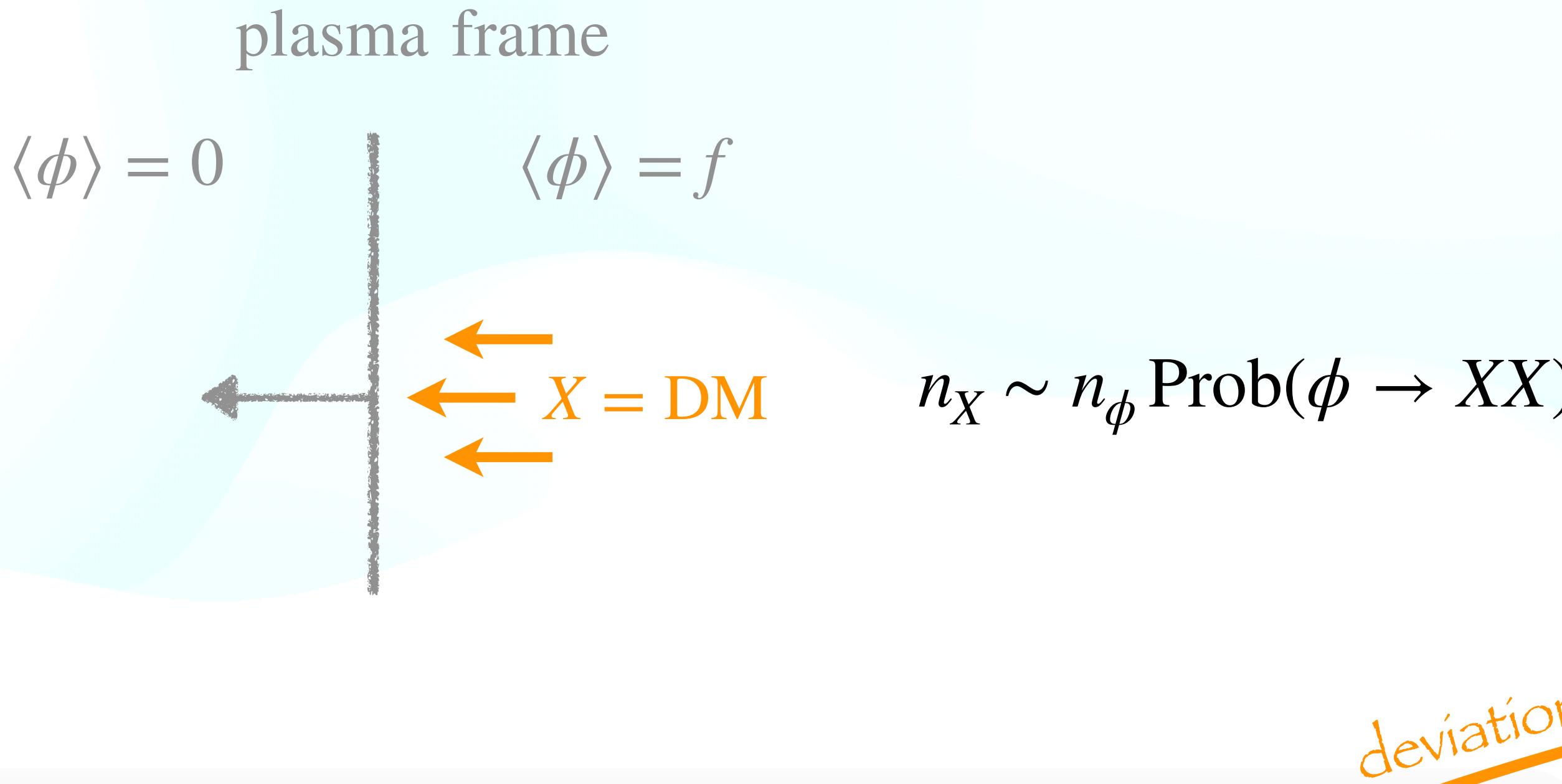


X scatterings with bath

$$\frac{1}{p_X} \frac{dp_X}{dt} < H \Rightarrow \lambda \lesssim \frac{10^9 \text{ GeV}}{M} \left(\frac{f}{T_{\text{nuc}}} \right)^{\frac{1}{4}} \left(\frac{\text{TeV}}{f} \right)^{\frac{1}{2}}$$

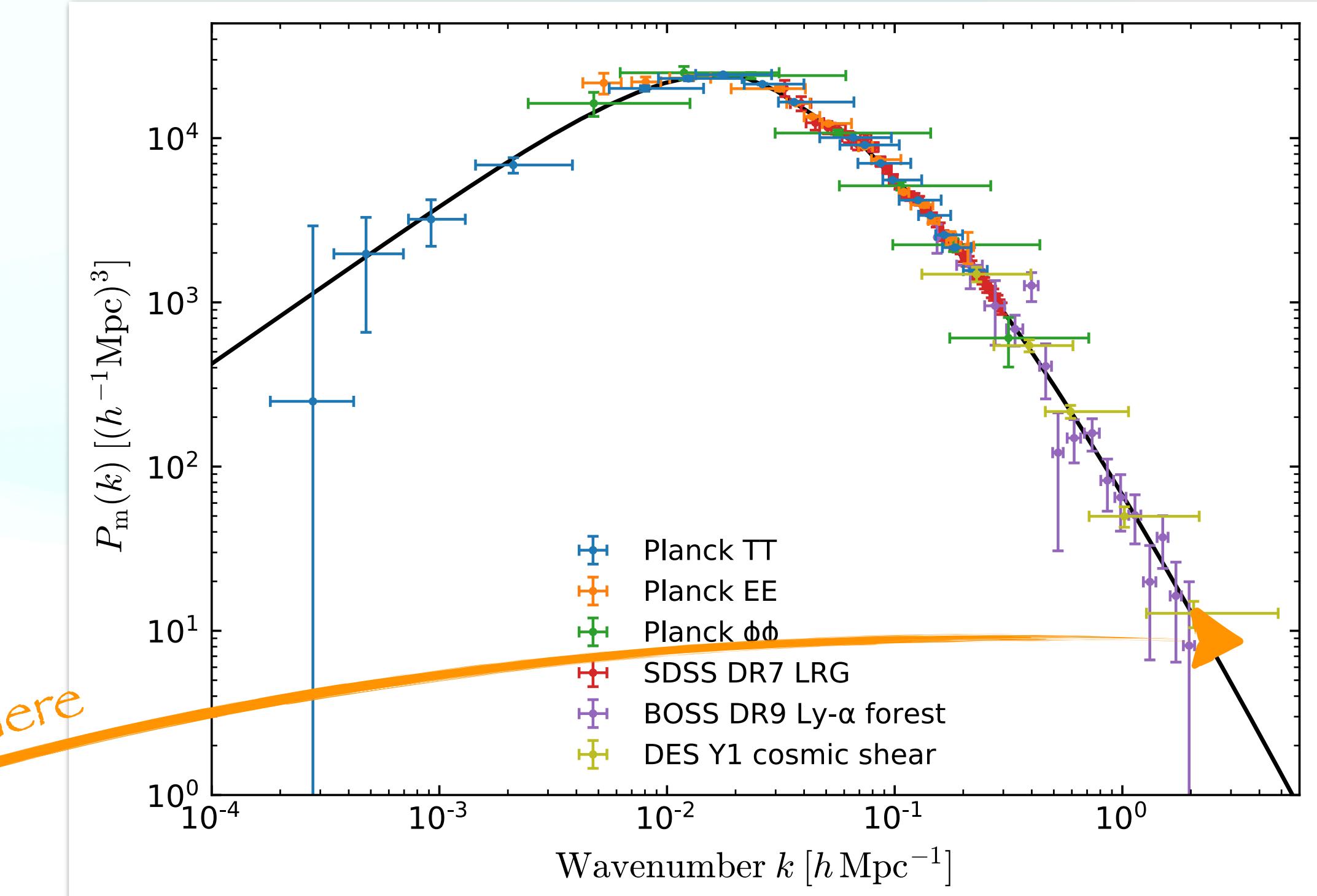


Warm EeV Dark Matter from the Weak Scale



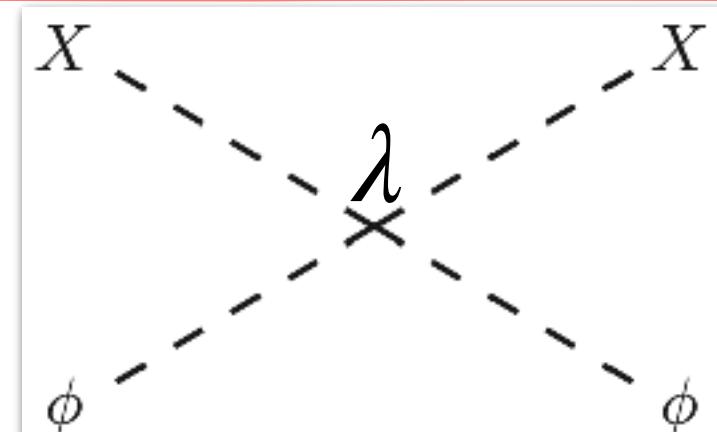
$p_X(T \sim f) \simeq \frac{M^2}{T_{\text{nuc}}}$ redshift $\rightarrow p_X(T \sim T_{\text{eq}}) \simeq \frac{\Gamma_{\text{eq}}}{f} \frac{M^2}{T_{\text{nuc}}}$

$$v_X(T \simeq T_{\text{eq}} \simeq 0.8 \text{ eV}) \simeq 10^{-4} \frac{f}{T_{\text{nuc}}} \left(\frac{100 \text{ GeV}}{f} \right)^2 \frac{M}{10^9 \text{ GeV}}$$

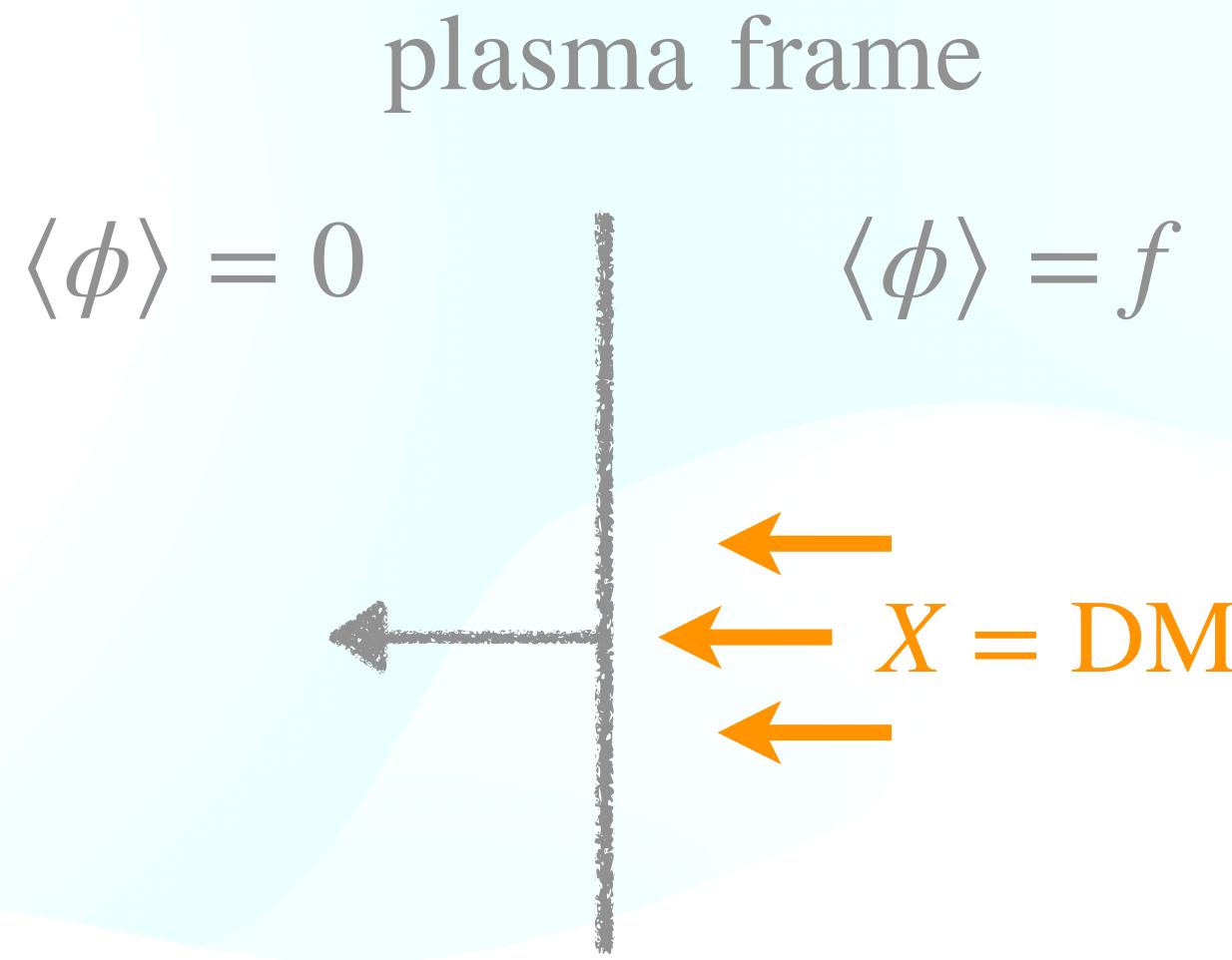


X scatterings with bath

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Warm EeV Dark Matter from the Weak Scale

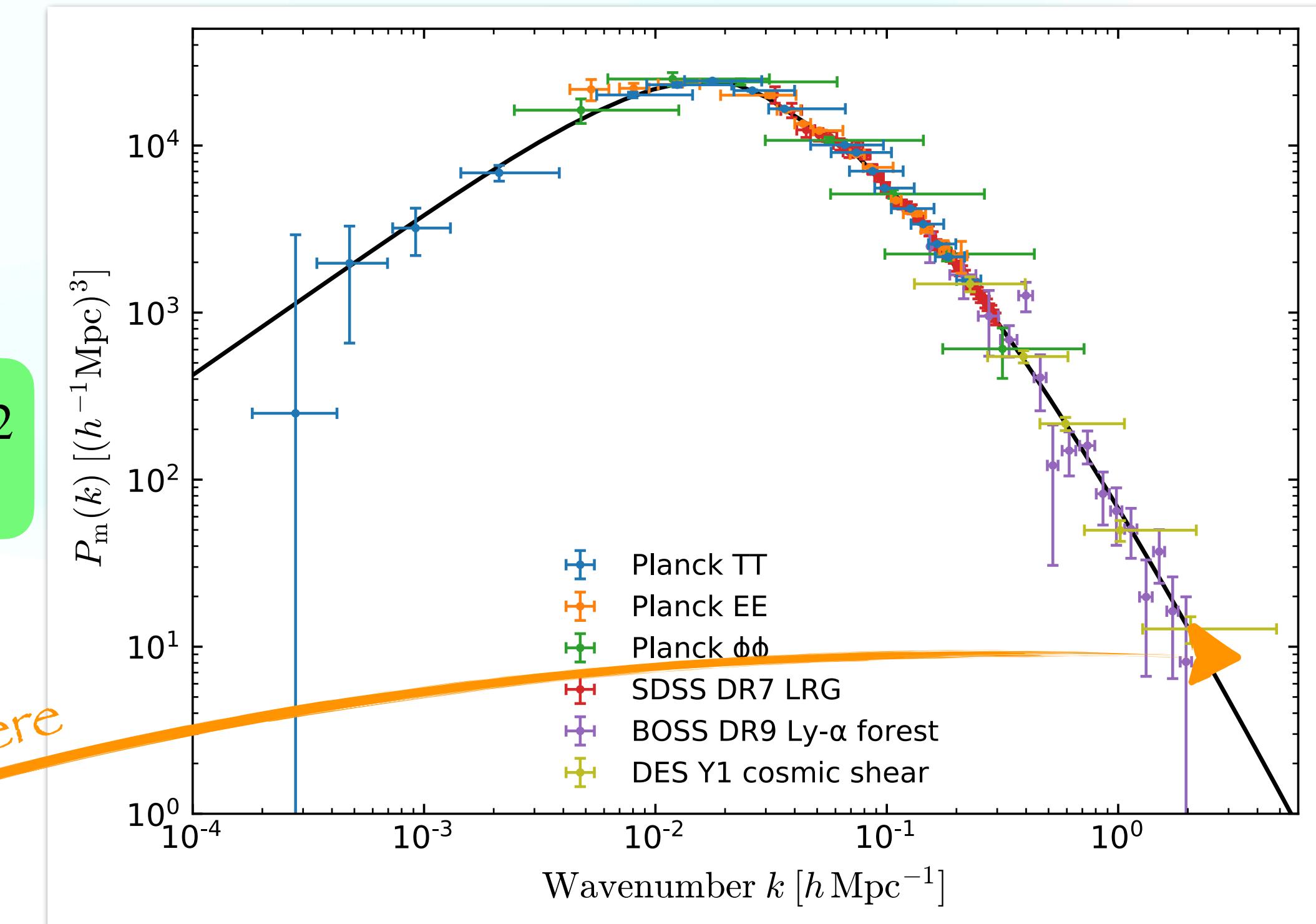


$$n_X \sim n_\phi \text{Prob}(\phi \rightarrow XX) \propto \lambda^2$$

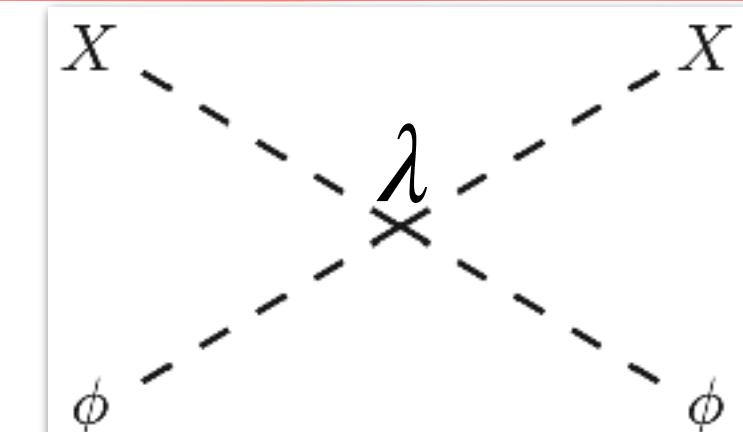
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deviations here

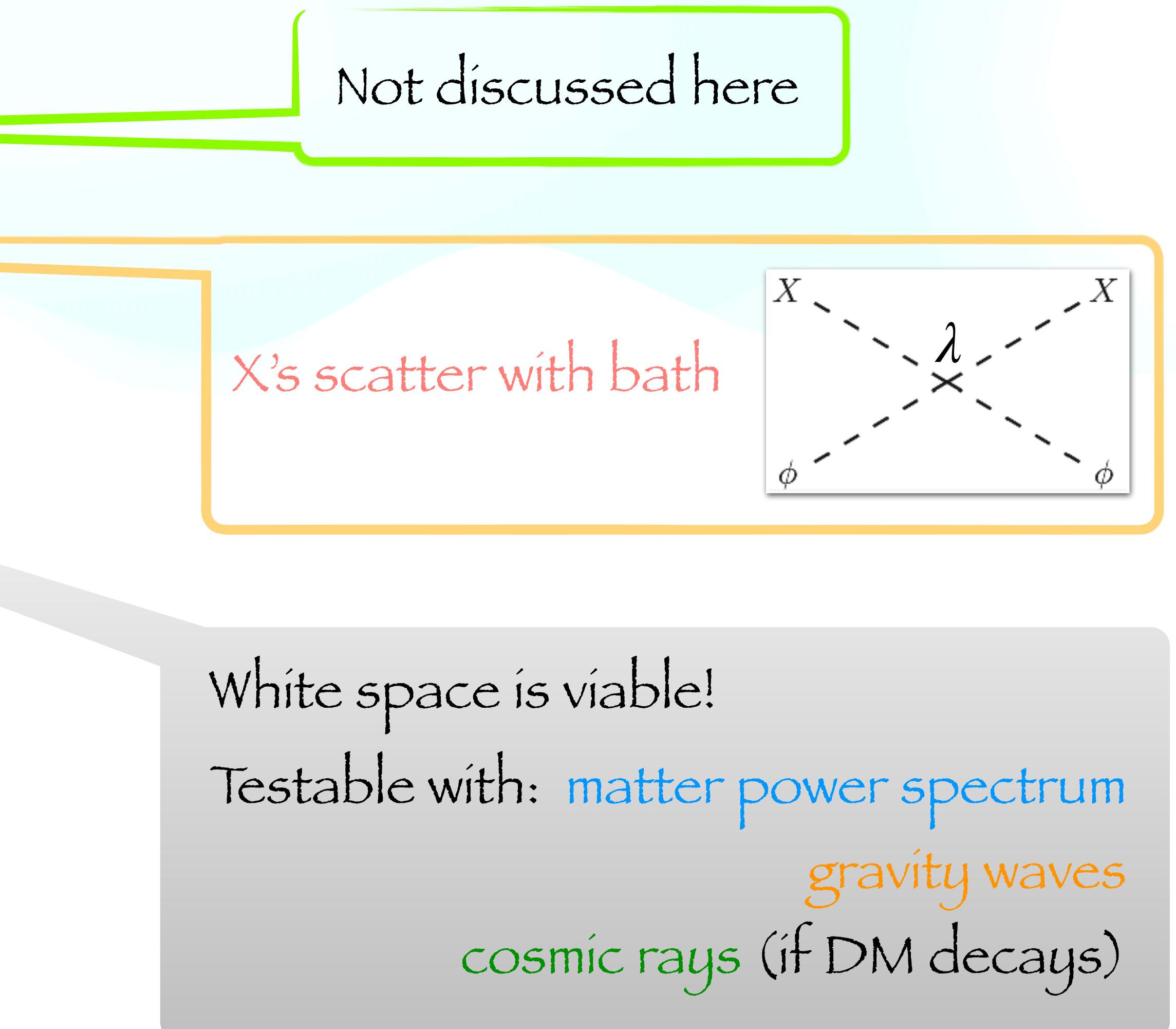
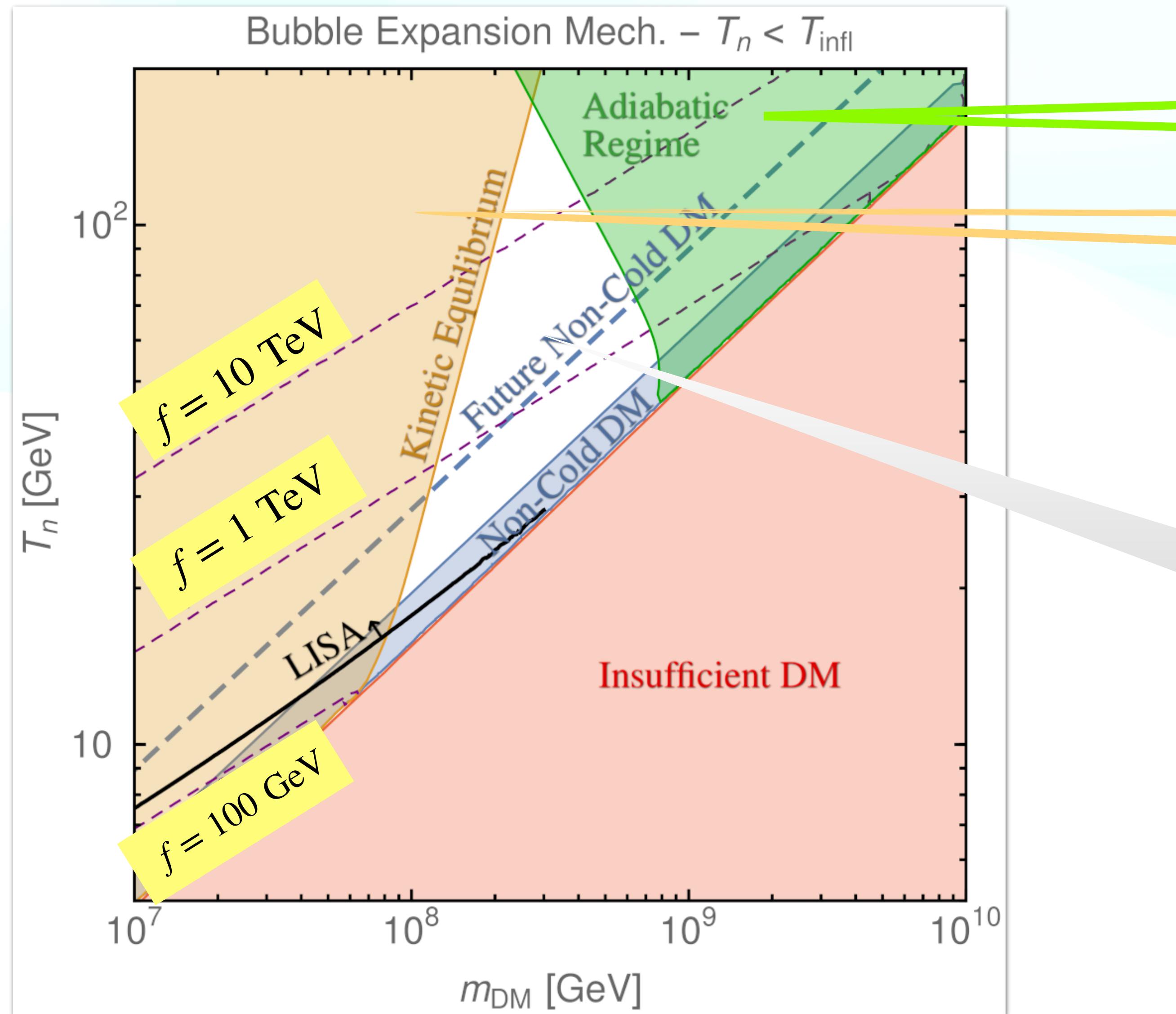


X scatterings with bath

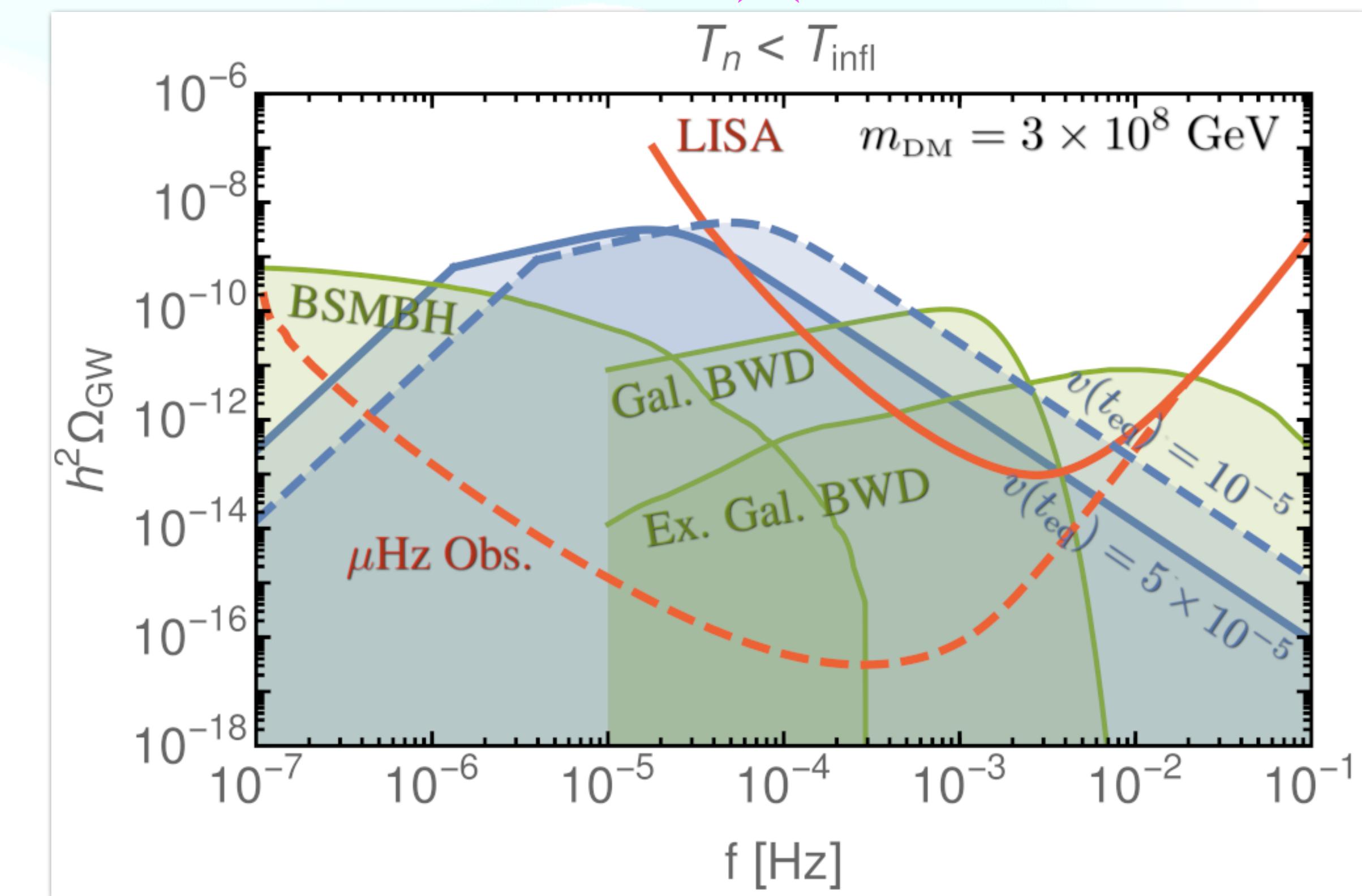
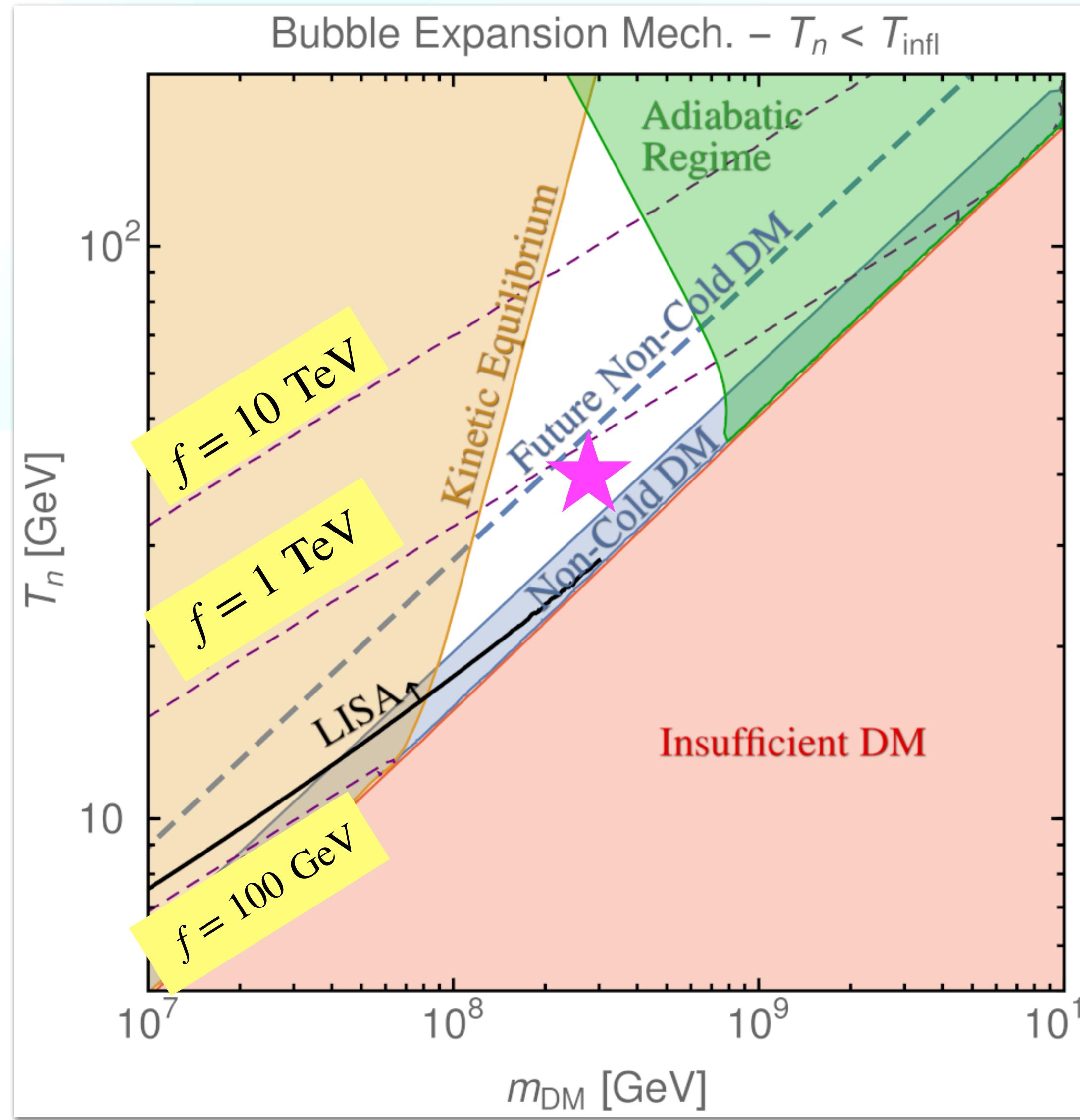


$$\frac{1}{p_X} \frac{dp_X}{dt} < H \Rightarrow \lambda \lesssim \frac{10^9 \text{ GeV}}{M} \left(\frac{f}{T_{\text{nuc}}} \right)^{\frac{1}{4}} \left(\frac{\text{TeV}}{f} \right)^{\frac{1}{2}}$$

Warm EeV Dark Matter from the Weak Scale



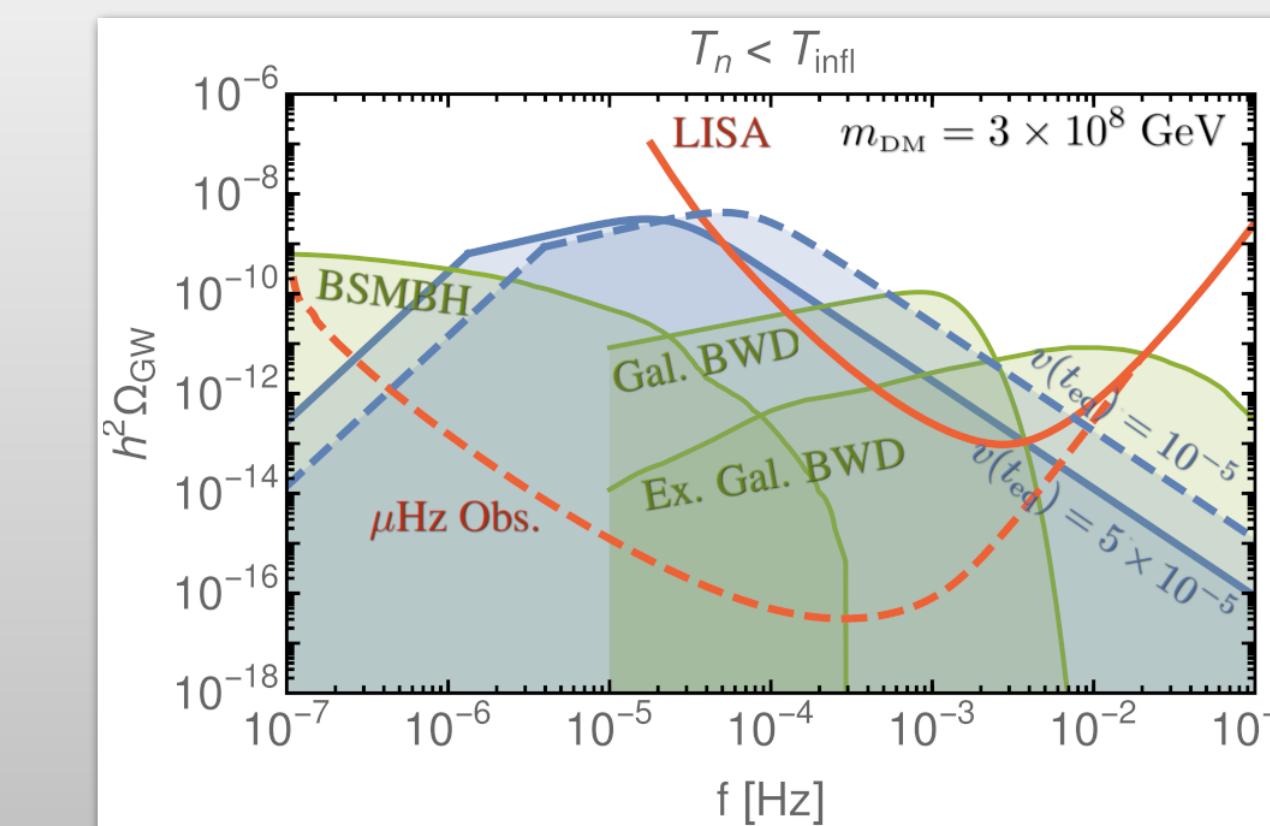
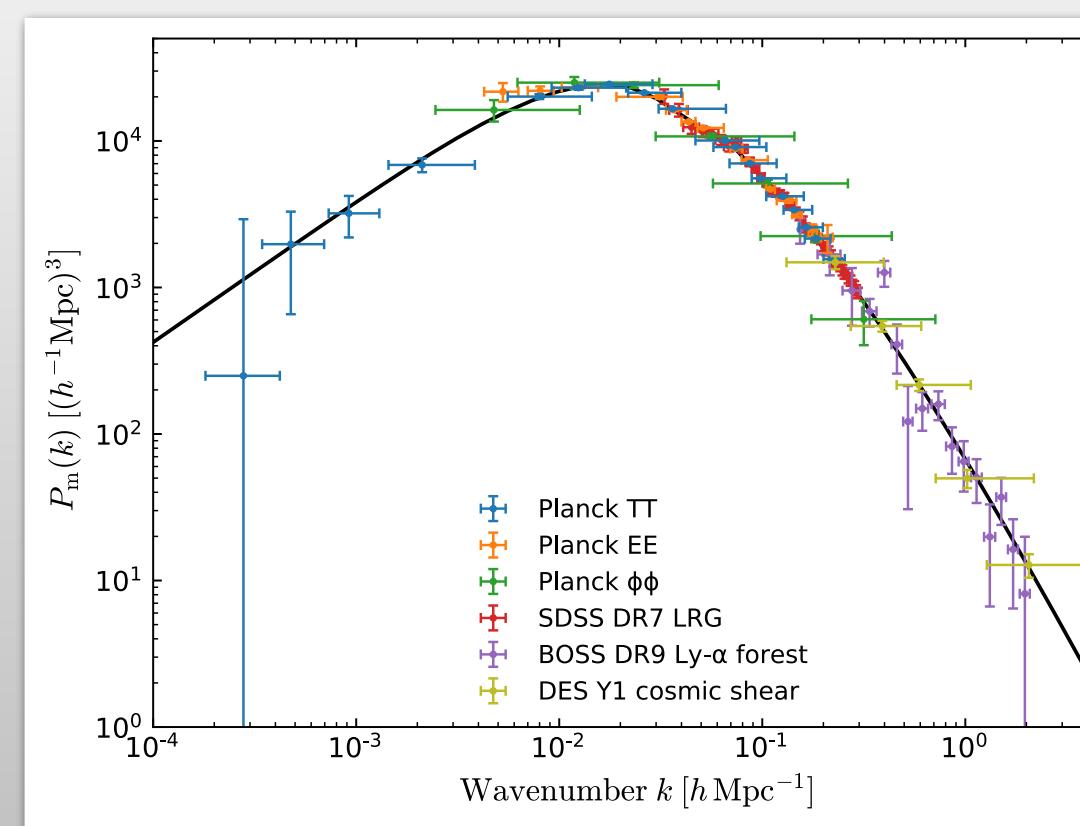
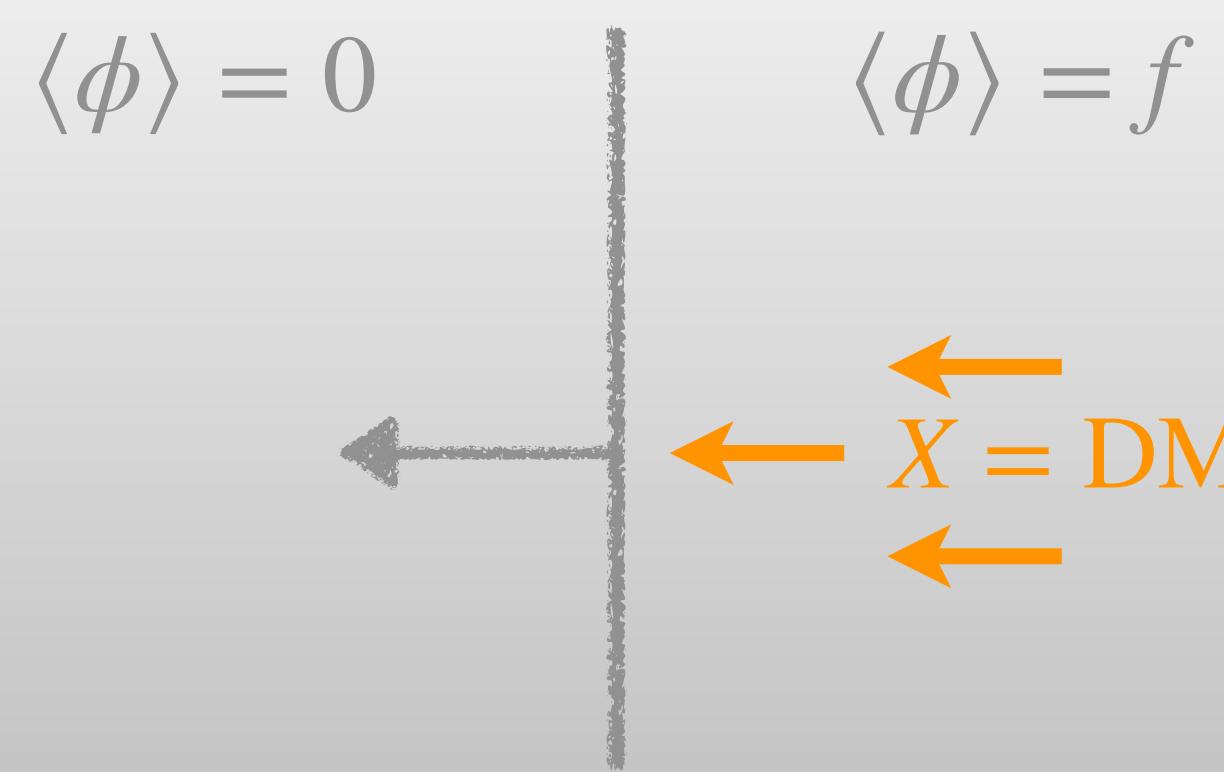
Warm EeV Dark Matter from the Weak Scale



Bubbletrons: Phase transitions with fast bubble walls are particle accelerators

Application: Warm EeV Dark Matter from Weak Scale Phase Transition

Baldes Gouttenoire FS
2207.05096



Still a lot to explore!

E.g. DM beyond GUT scale from PT at much lower scale

Baldes Gouttenoire FS + M. Dichtl in progress

Back up

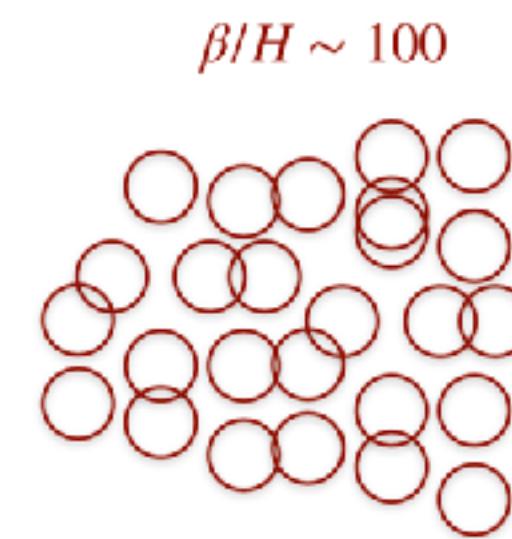
Gravitational Waves from Phase Transition

Randall Servant hep-ph/0607158,...

$$\Omega_{\text{GW}} \propto (H/\beta)^2 \quad \beta^{-1} \sim \text{duration of PT}$$

$$\frac{\beta}{H} \simeq T \frac{dS_4}{dT} \Big|_{T_{\text{nuc}}} \simeq 15 \left(\frac{10}{N_{\text{e-fold}}} \right)^2$$

Standard 1st order PT



Supercooled PT
 $\beta/H \sim 10$

