General covariance and dynamics with a Gauss law

Hassan Mehmood work in collaboration with Vigar Husain (arXiv: 2312.06079)

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Prelude: the Husain-Kuchar (HK) model

Given an $\mathfrak{su}(2)$ -valued triad e_a^i and connection A_a^i ($a \in \{1, ..., 4\}$) on a 4d spacetime M, consider the generally covariant action

$$S = rac{1}{2} \int_M d^4 x \operatorname{Tr}(e \wedge e \wedge F)$$

where $F = dA + A \wedge A$.

Contrast with the 4d Palatini action: $\mathfrak{so}(3,1)$ -valued *tetrads* replaced with $\mathfrak{su}(2)$ -valued *triads*.

Canonical HK

Assuming $M = \mathbb{R} \times \Sigma$, the canonical decomposition of the action is straightforward:

$$S = \int dt \int_{\Sigma} d^3 x (\tilde{E}^a_i \dot{A}^i_a - A^i_0 \tilde{G}_i - (e^i_0 E^a_i) \tilde{C}_a)$$

where $\tilde{E}_{i}^{a} = \det(e)E_{i}^{a} = \tilde{\epsilon}^{0abc}\epsilon_{ijk}e_{b}^{j}e_{c}^{k}$, and

$$\begin{split} \tilde{G}_i &= -D_a \tilde{E}_i^a pprox 0 \quad \mbox{(Gauss)} \\ \tilde{C}_a &= \tilde{E}_i^b F_{ab}^i pprox 0 \quad \mbox{(spatial diffeos)} \end{split}$$

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No Hamiltonian constraint!

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- The theory is non-dynamical: the geometry of Σ does not evolve. But not topological: local degrees of freedom exist.
- ► There's an invertible spatial metric g_{ab} = δ_{ij}eⁱ_ae^j_b, a, b ∈ {1, 2, 3}. Thus interesting three-geometries exist.



Consider the replacement

$$e^i_{a}
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where D is the covariant derivative with connection A and ϕ is an $\mathfrak{su}(2)$ -valued scalar.

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The action now decomposes as

$$S = \int dt d^{3}x (\tilde{E}_{i}^{a} \dot{A}_{a}^{i} + \tilde{p}_{i} \dot{\phi}^{i} - A_{0}^{i} \tilde{G}_{i});$$

$$\tilde{E}_{i}^{a} = \tilde{\epsilon}^{abc} \epsilon_{ijk} D_{b} \phi^{j} D_{c} \phi^{k}, \quad \tilde{p}_{i} = \tilde{\epsilon}^{abc} \epsilon_{ijk} D_{a} \phi^{j} F_{bc}^{k}$$

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with only one constraint, a modified Gauss law with a source:

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No Hamiltonian and diffeomorphism constraints!

But whither the constraints?

The theory is generally covariant. But the first-class constraints of the theory (namely, the Gauss law) generate only SU(2) transformations of the gauge field A. Where does the remaining gauge redundancy go?

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In these theories, for any generator of diffeomorphisms v,

 $\mathcal{L}_{v}A = G$ -transformations + equations-of-motion terms

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where G is the gauge group of the connection A.

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- For instance, canonical quantization via LQG methods yields a Hilbert space of spin network states with a finite number of charges \u03c6 sitting at the vertices.
- Would be interesting to look at the spinfoam and group field theory models of the theory (work currently underway).

Thank you!

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