

CROSS-SECTION CONTINUITY OF ANGULAR MOMENTUM

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THE METRIC NEAR NULL INFINITY

Class of metrics near \mathcal{I}^+ can be written down in $1/r$ expansion ($x^A = (\theta, \phi)$):

$$ds^2 = -du^2 - 2dudr + r^2 h_{AB} dx^A dx^B \\ + \frac{2m}{r} du^2 + r C_{AB} dx^A dx^B + D^B C_{AB} dudx^A \\ + \frac{4}{3r} N_A dudx^A + \dots$$

- Bondi mass aspect $m(u)$
- Shear tensor $C_{AB}(u, x^A)$
- Angular momentum aspect $N_A(u, x^A)$

NB: a 'cross-section' of \mathcal{I}^+ is 2-sphere labelled by u (Bondi frame).

THE SUPERTRANSLATION AMBIGUITY

$$BMS_4 = SO(3,1) \ltimes S$$

Supertranslations are coupled to rotations. Cannot distinguish between ‘rotation’ and ‘rotation + supertranslation’. Cannot recover unique Lorentz subgroup, rather infinite number, each differing by supertranslation.

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AMBIGUITY EXAMPLE

On $u = 0$ cross-section call $X = Y^A \frac{\partial}{\partial x^A}$ ‘rotation’ (Y^A is rotational Killing field on $u = \text{const}$ sphere).

On $u' = 0$ cross-section, where $u' = u - f(x^A)$ and $x'^A = x^A$, symmetry is $X = Y^A \frac{\partial}{\partial x'^A} - Y^A D_A f \frac{\partial}{\partial u'} \rightarrow$ ‘rotation + supertranslation’ $-Y^A D_A f$.

Why? Different foliations!

ANGULAR MOMENTUM AT NULL INFINITY

Immediate consequence: cannot distinguish ‘angular momentum’ and ‘angular momentum + supertranslation charge’. How to proceed?

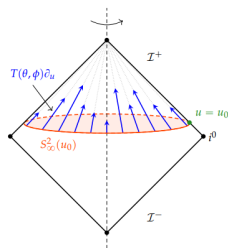


FIGURE: *Compere*, 1801.07064v4.
Supertranslations T act by “wiggling” cross-sections of \mathcal{I} .

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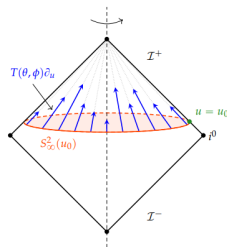


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- 2 Use data (physical fields) on cross-section to determine what should be pure rotation.

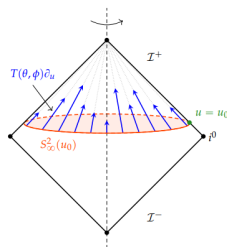


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THE DRAY STREUBEL ANGULAR MOMENTUM

DS angular momentum

$$J^{DS}|_{u=0} = \frac{1}{8\pi} \int Y^A \left(N_A - \frac{1}{4} C_{AB} D_D C^{DB} \right)$$

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All combined into **one-parameter family by Compere and Nichols:**

$$J^{CN}|_{u=0} = \frac{1}{8\pi} \int Y^A \left(N_A - \frac{1}{4} C_{AB} D_D C^{DB} \right) - \frac{\alpha - 1}{32\pi} \int Y^A C_{AB} D_C C^{BC}$$

THE CHEN-WANG-YAO ANGULAR MOMENTUM

Arose from attempts to eliminate supertranslation ambiguity:

$$J^{CWY}|_{u=0} = \frac{1}{8\pi} \int Y^A \left(N_A - \frac{1}{4} C_{AB} D_D C^{DB} \right) + \frac{1}{8\pi} \int m Y^A D_A c$$

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Note: CWY angular momentum does not reduce to Komar formula for arbitrary cross-section in presence of axial Killing field.

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CROSS-SECTION CONTINUITY CONDITION

Let $J : \mathcal{C} \rightarrow \mathbb{R}$ be map from space, \mathcal{C} , of smooth cross-sections of \mathcal{I} into \mathbb{R} . Then J satisfies cross-section continuity condition at $u = 0$ if for any sequence $\{f_n\}$ of smooth functions on sphere such that $f_n \rightarrow 0$ uniformly as $n \rightarrow \infty$ we have $J|_{u=f_n} \rightarrow J|_{u=0}$.

FLUX OF ANGULAR MOMENTUM

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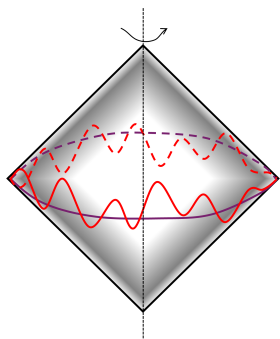
DS angular momentum has well-defined flux and obeys balance law:

$$J^{DS}|_{\mathcal{C}_2} - J^{DS}|_{\mathcal{C}_1} = \int_R \mathcal{F} = -\frac{1}{32\pi} \int N^{AB} \mathcal{L}_Y C_{AB}$$

Balance law is necessary criteria and cross-section continuity can be used as tool to check if satisfied.

RESULTS AND CONCLUSION

Dray-Streubel definition is c.c. since it arises from local flux. So use bootstrapping argument for other definitions: investigate differences $J^{CN} - J^{DS}$ and $J^{CWY} - J^{DS}$.

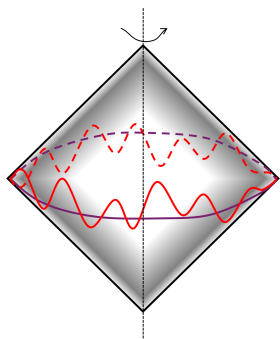


$$\{\mathbf{f}_n\} = \mathbf{F}(\theta) \frac{1}{n} \sin(n\phi)$$

FIGURE: Cross section continuity means $J|_{u=0}$ approaches $J|_{u=f_n}$ as $n \rightarrow \infty$.

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$$\{\mathbf{f}_n\} = \mathbf{F}(\theta) \frac{1}{n} \sin(n\phi)$$

CN one-parameter family only c.c. when $\alpha = 1$. Only DS definition is viable, generalizations are not. CWY definition is also c.c.

FIGURE: Cross section continuity means $J|_{\mathbf{u}=0}$ approaches $J|_{\mathbf{u}=\mathbf{f}_n}$ as $n \rightarrow \infty$.

- Cross-section continuity is natural condition to impose: $u = \text{const}$ cross-sections on \mathcal{I} represent ‘instances of time’.
- Not trivial to see functions on \mathcal{I} satisfy c.c. due to nonlinear transformation laws, unless local flux exists.
- Criterion introduced via angular momentum, but should be useful more broadly to generic functions of cross-sections in dynamical (radiative) eras (c.f. *S. Hollands, R. M. Wald, and V. G. Zhang*).