

Revisiting quantum BHs from effective LQG

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Loops

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Overview

- LOOP QUANTUM BLACK HOLES
- AOS MODEL
- FAMILY OF SOLUTIONS
- ASYMPTOTIC BEHAVIOR
- SUMMARY

LOOP QUANTUM BLACK HOLES

CLASSICAL

$$\mathcal{H}_{cl}(b, p_b; c, p_c) = -\frac{1}{2G\gamma} \left(2cp_c + \left(b + \frac{\gamma^2}{b} \right) p_b \right)$$

Hamiltonian of the Schwarzschild solution

$$ds^2 = -N^2 dT^2 + \frac{p_b^2}{|p_c| L_o^2} dx^2 + |p_c| d\Omega^2$$

Schwarzschild metric in Kantowski-Sachs form

$$b \rightarrow \frac{\sin(\delta_b b)}{\delta_b} \quad ; \quad c \rightarrow \frac{\sin(\delta_c c)}{\delta_c}$$

EFFECTIVE

$$\mathcal{H}_{eff}(b, p_b, \delta_b; c, p_c, \delta_c) = -\frac{1}{2G\gamma} \left[2 \frac{\sin(\delta_c c) |p_c|}{\delta_c} + \left(\frac{\sin(\delta_b b)}{\delta_b} + \frac{\gamma^2 \delta_b}{\sin(\delta_b b)} \right) p_b \right]$$

Effective loop quantized Schwarzschild solution

δ_i : quantum parameters

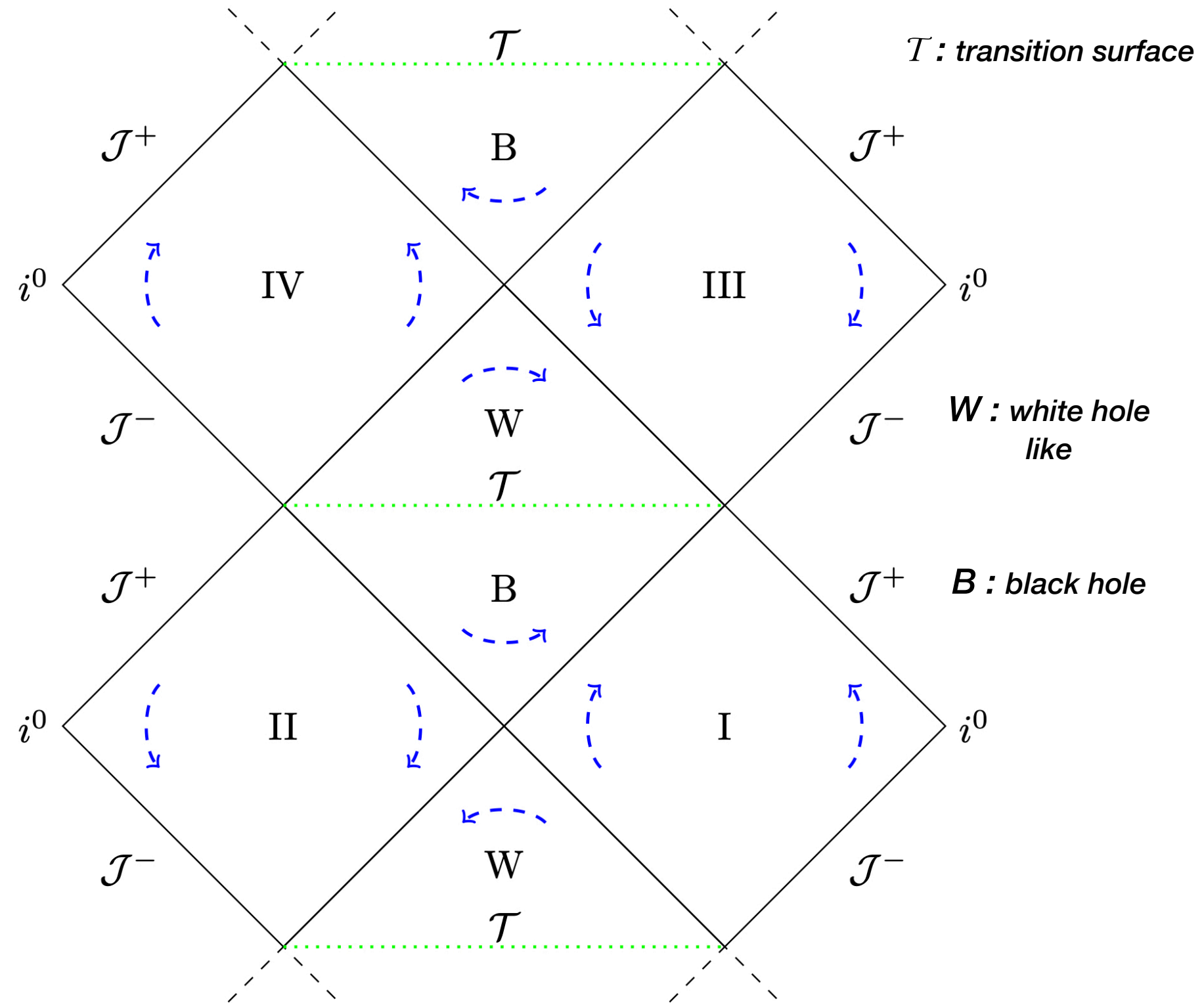
AOS MODEL

AOS: Ashtekar, Olmedo & Singh

$$\delta_b = \left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}} \right)^{\frac{1}{3}}, \quad L_0 \delta_c = \frac{1}{2} \left(\frac{\gamma \Delta^2}{4\pi^2 m} \right)^{\frac{1}{3}}$$

δ_i : Dirac observables

- BH singularity replaced by a transition surface.
- Symmetric spacetime near the transition surface.
- Universal upper bound of invariants at T .
- No mass amplification of the WH.
- Negligible quantum corrections at classical scales.
- Asymptotically flat-like spacetime at infinity.



Penrose diagram of the AOS model

FAMILY OF SOLUTIONS

$$\mathcal{H}_{\text{eff}}(b, p_b, \delta_b; c, p_c, \delta_c)$$

$$(B_o, \delta_b; c_o, p_c^o, \delta_c)$$

The solution space of the effective Hamiltonian has 5 free parameters

$$(\delta_b, \delta_c, c_o, m)$$

Reduces to 4 parameters after matching with classical solution & gauge fixing

$$(\alpha_b, m)$$

Reduces to 2 parameters to ensure negligible quantum corrections at classical scales

FAMILY OF SOLUTIONS

AOS

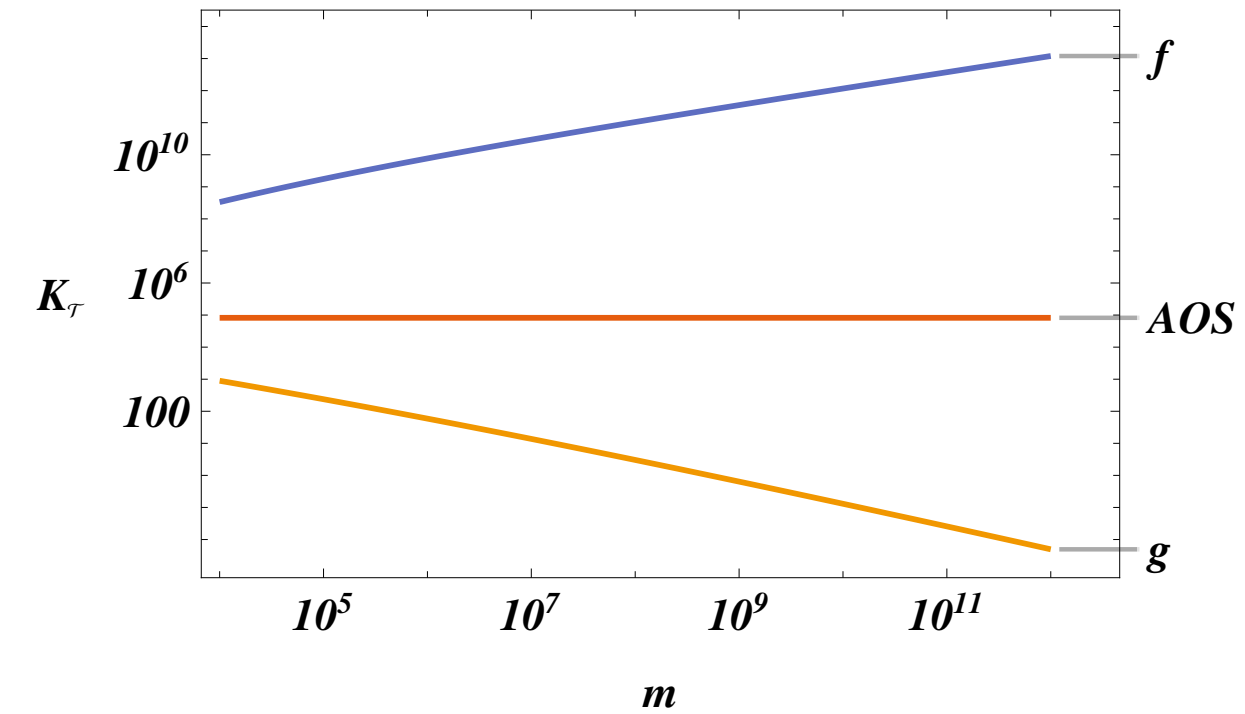
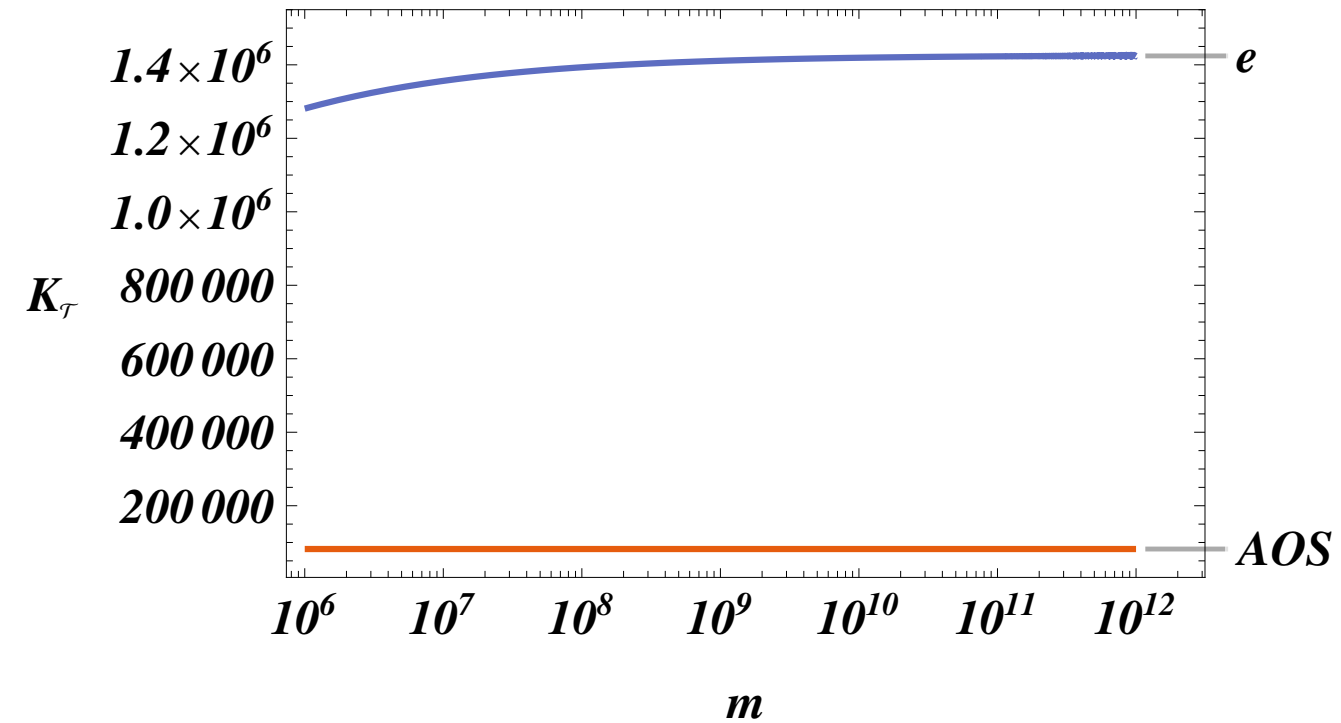
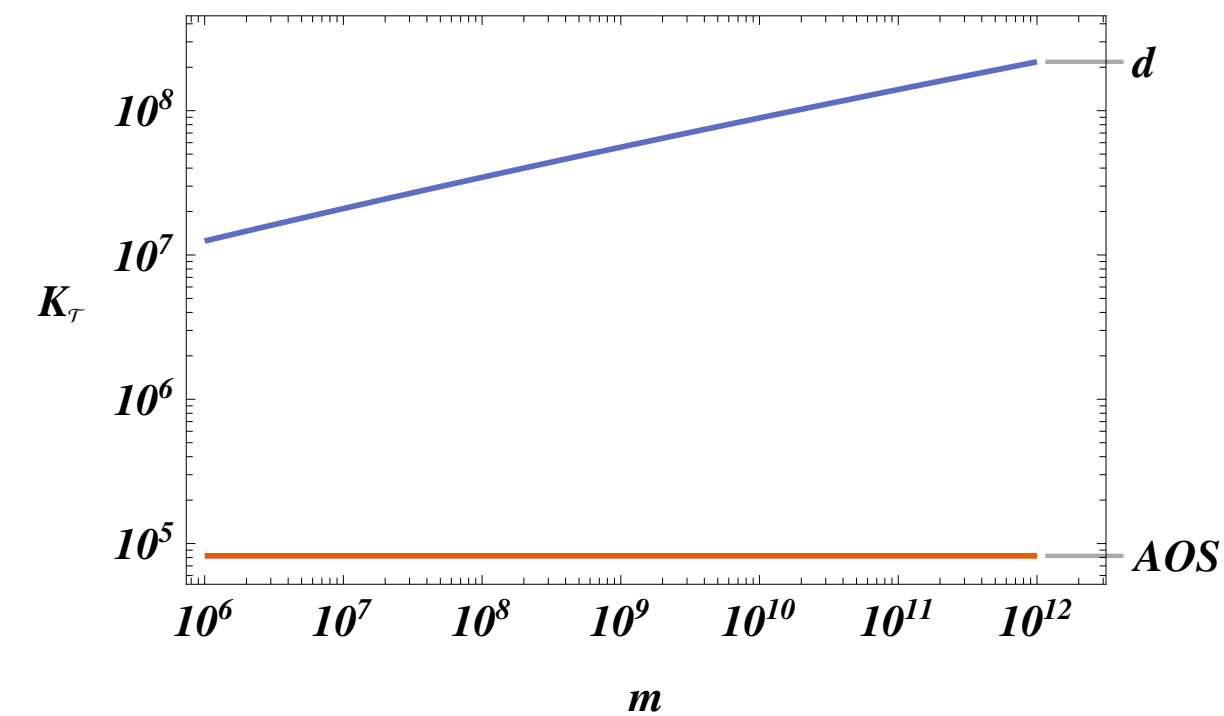
$$\delta_b^{AOS} = \alpha_b^{AOS} \left(\frac{\ell_p}{m} \right)^{1/3}, \quad L_o \delta_c^{AOS} = \alpha_c^{AOS} \left(\frac{\ell_p}{m} \right)^{1/3}$$

$$\alpha_b^{AOS} \simeq 2.52, \quad \alpha_c^{AOS} \simeq 0.27$$

FAMILY

$$\delta_b = \alpha_b \left(\frac{\ell_p}{m} \right)^{1/3}, \quad L_o \delta_c = \alpha_c \left(\frac{\ell_p}{m} \right)^{1/3}$$

FAMILY OF SOLUTIONS



Plots of Kretschmann scalar at transition surface vs mass for various choices in comparison to AOS

$$d: (\delta_b, \delta_c) = \left(\delta_b^{(AOS)}, \frac{1}{2L_o} \left(\frac{\gamma \Delta^{3/2}}{4\pi^2 m} \right)^{1/2} \right)$$

$$e: (\delta_b, \delta_c) = \left(\left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}} \right)^{1/2}, \delta_c^{(AOS)} \right)$$

$$f: (\delta_b, \delta_c) = \left(\left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}} \right)^{1/2}, \frac{1}{2L_o} \left(\frac{\gamma \Delta^{3/2}}{4\pi^2 m} \right)^{1/2} \right)$$

$$g: (\delta_b, \delta_c) = \left(\left(\frac{\sqrt{\Delta}}{\sqrt{2\pi\gamma^2 m}} \right)^{1/4}, \frac{1}{2L_o} \left(\frac{\gamma \Delta^{5/2}}{4\pi^2 m} \right)^{1/4} \right)$$

FAMILY OF SOLUTIONS

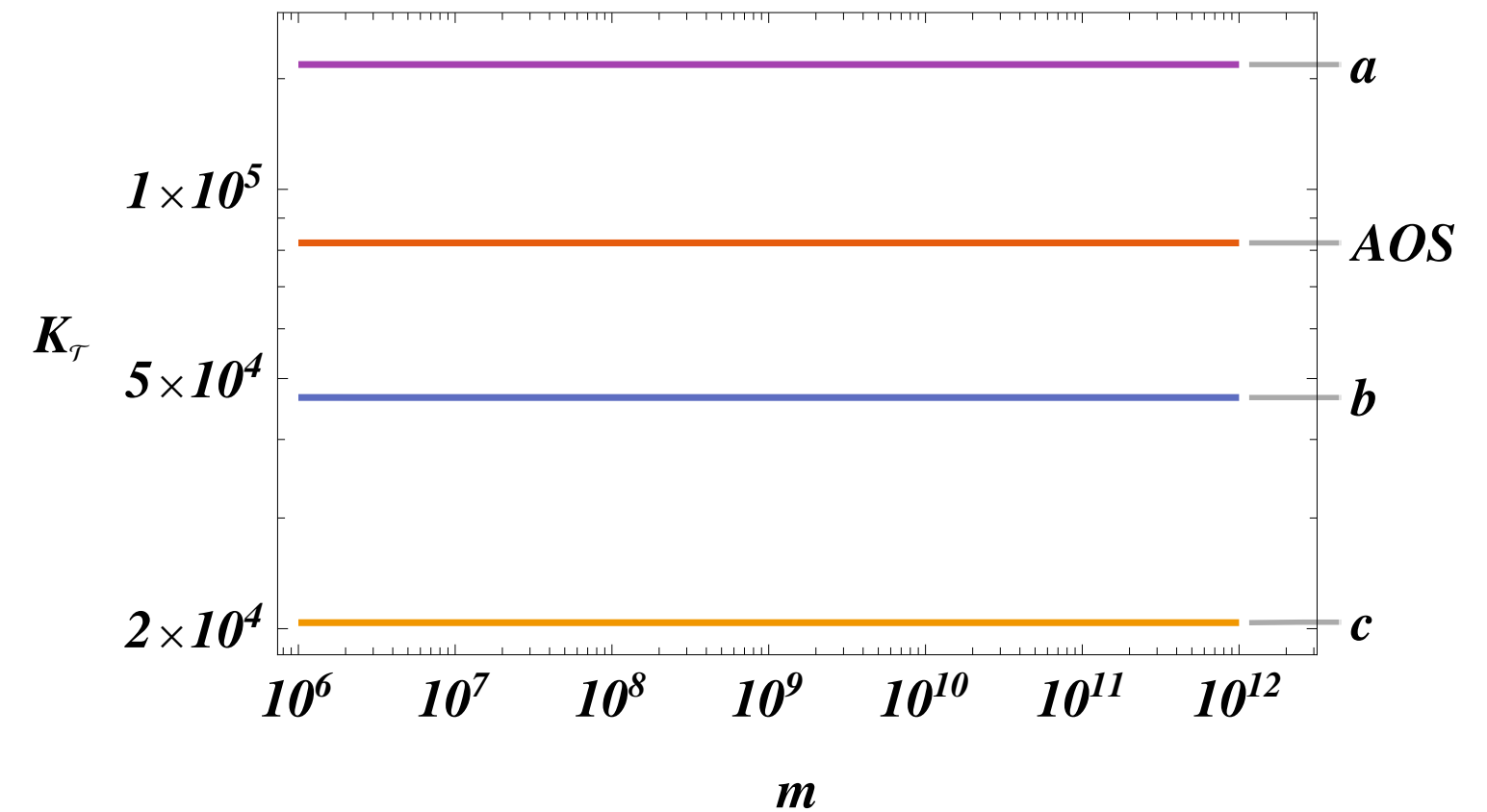
$$\delta_b = \alpha_b \left(\frac{\ell_p}{m} \right)^{1/3}, \quad L_o \delta_c = \alpha_c \left(\frac{\ell_p}{m} \right)^{1/3}$$

&

Demanding radius of BH = WH we get



$$\delta_b = \alpha_b \left(\frac{\ell_p}{m} \right)^{1/3}, \quad L_o \delta_c = \alpha_b^4 \left(\frac{\gamma^3 \ell_p}{2} \right) \left(\frac{\ell_p}{m} \right)^{1/3}$$



Plot of Kretschmann scalar at transition surface vs mass

Few random choices

Line	α_b	α_c
a	1.97	3
b	2	0.5
c	1	1

ASYMPTOTIC BEHAVIOR

- Kretschmann Scalar in the asymptotic limit $r \rightarrow \infty$

$$K = \frac{a_0}{r^4} + \frac{a_1}{r^{4+b_o}} + \frac{a_2}{r^{4+2b_o}} + \frac{a_3}{r^{4+3b_o}} + \frac{a_4}{r^8} + \mathcal{O}\left(\frac{1}{r^{4(1+b_o)}}\right), \quad (b_o \approx 1)$$

- However for a solar mass BH in the AOS and CS model, r^{-4} and r^{-5} terms dominate at scales greater than the size of our observable universe

$$r_{c, M_\odot}^{AOS} > L_{obs}$$

- Spacetime well described by its classical limit for macroscopic BHs

$$K \simeq r^{-6}$$

SUMMARY

- BH singularity replaced by a transition surface
- Antitrapped region (WH) to the future of the transition surface
- Kretschmann scalar of macroscopic BHs decays as r^{-6} within our observable universe
- Family of solutions to the quantum parameters?

Thank you!

Do you have any questions?