

Light Dark World International Forum 2022

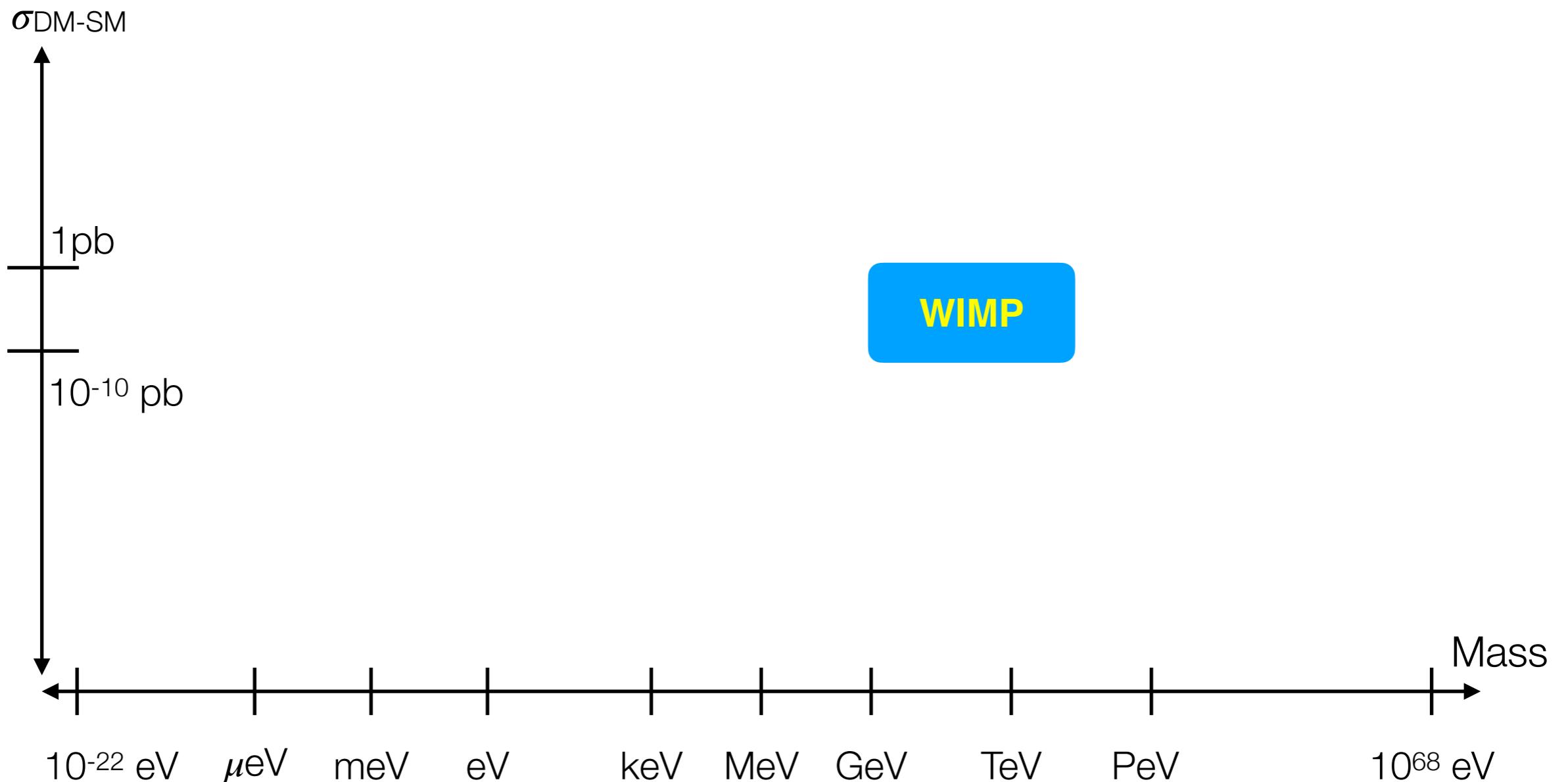
Manifesting hidden dynamics of a sub-component dark matter

Seodong Shin



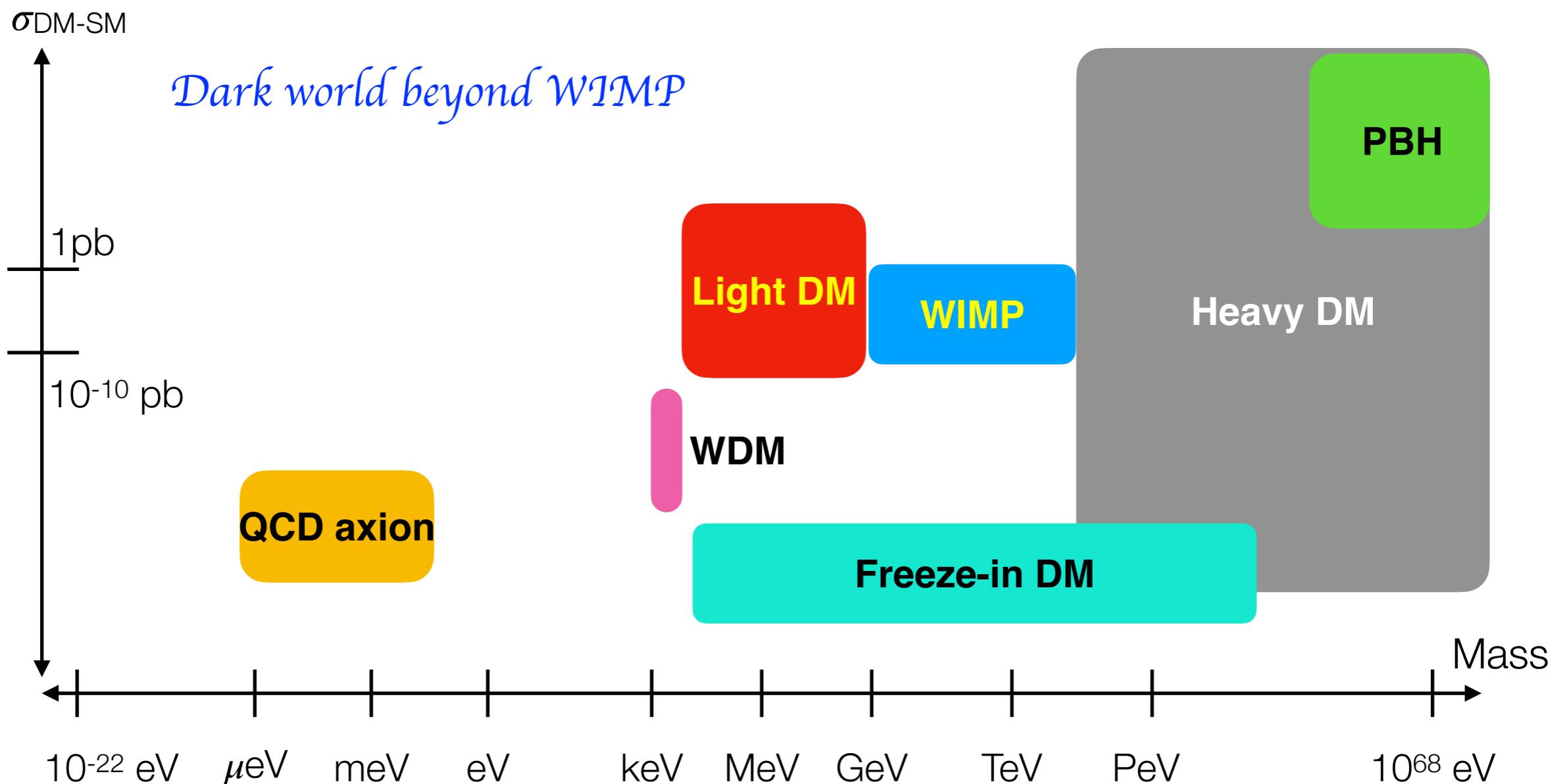
Ayuki Kamada, Hee Jung Kim, Jong-Chul Park, **ss**, JCAP 10, 052 (2022), arXiv: 2111.06808

Dark world beyond WIMP



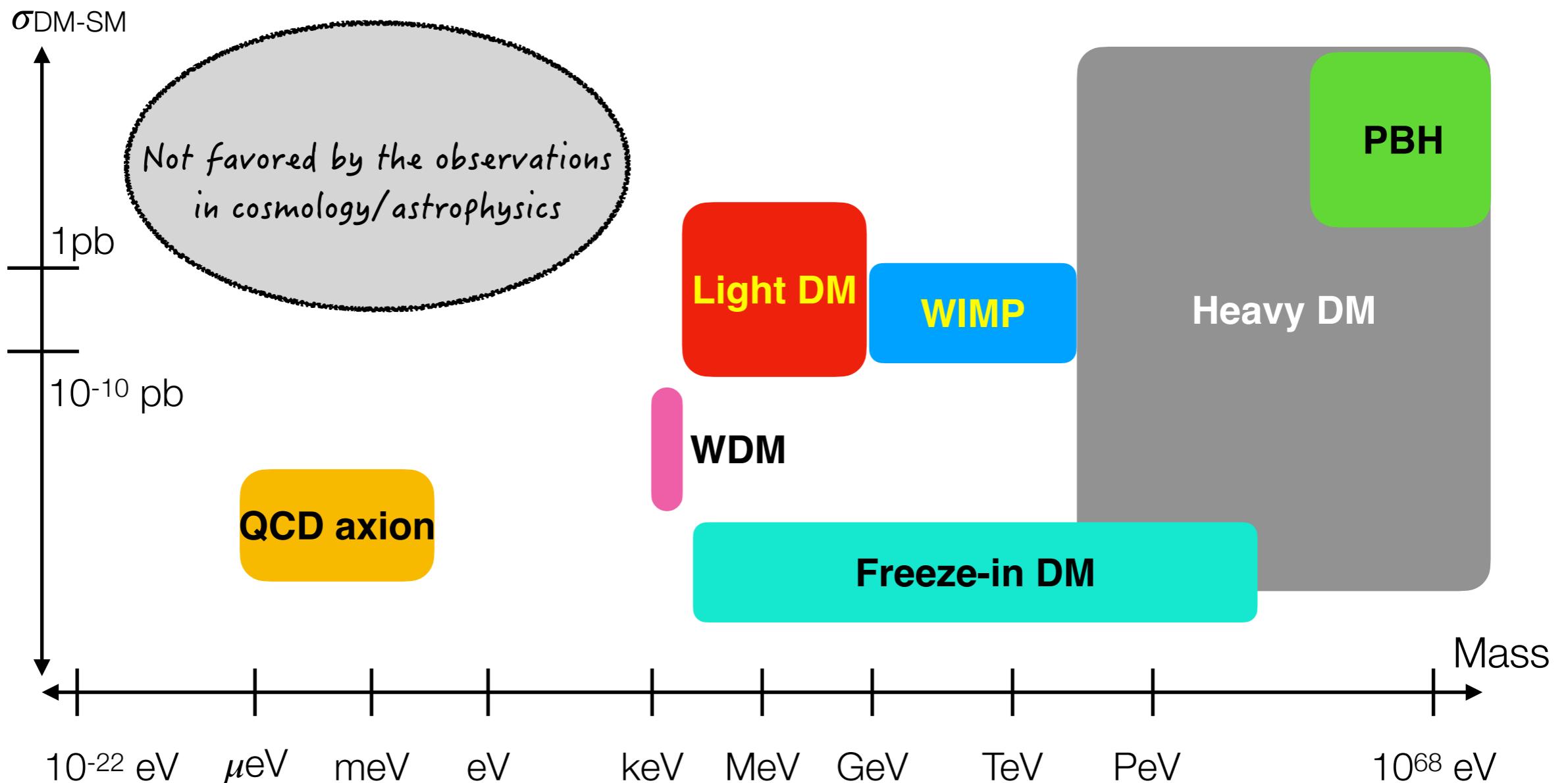
- WIMP: a single species of particles with thermal relic via freeze-out
- Mass in between $1 \text{ GeV} \lesssim m_\chi \lesssim 100 \text{ TeV}$ roughly

Dark world beyond WIMP

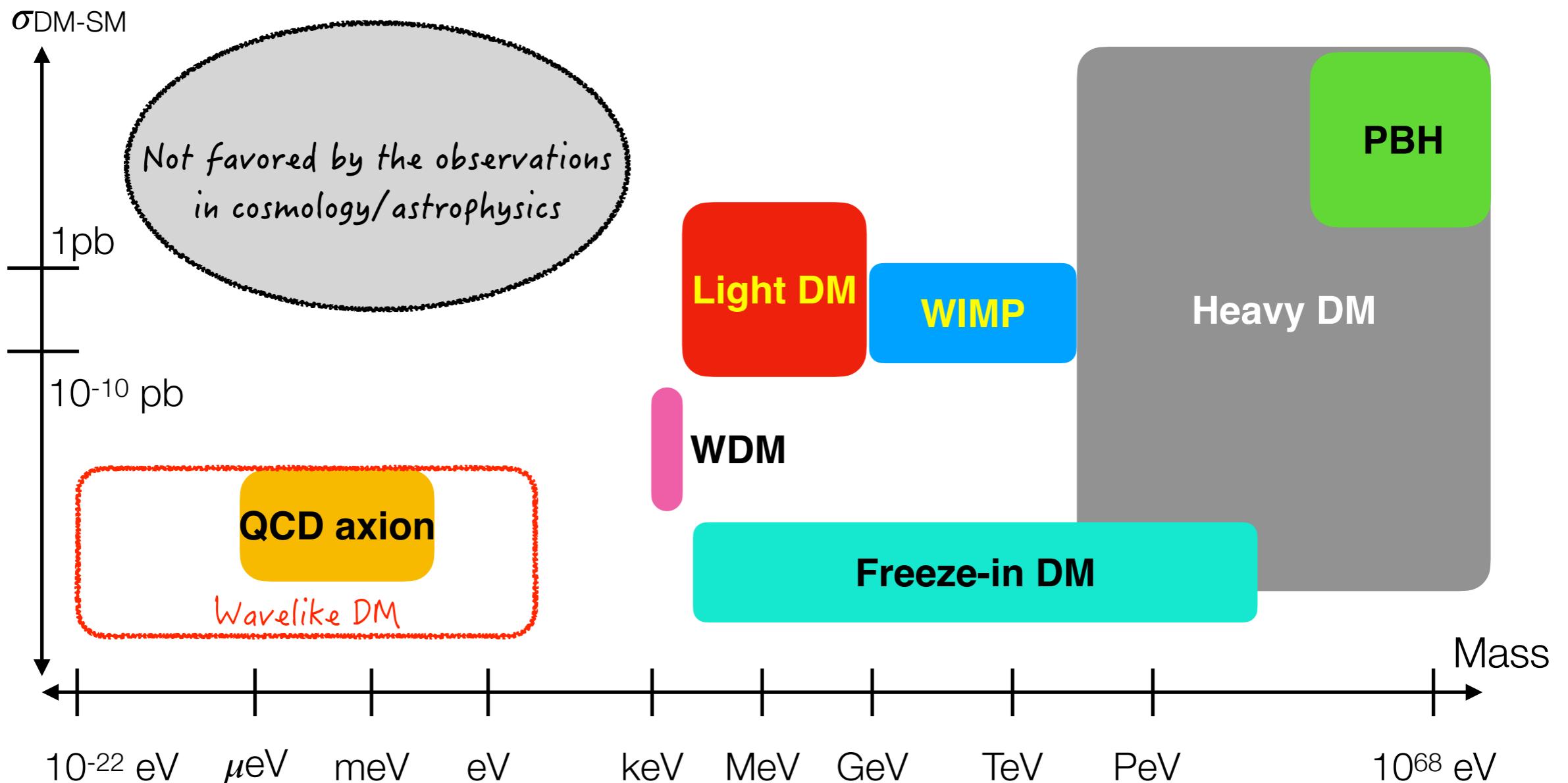


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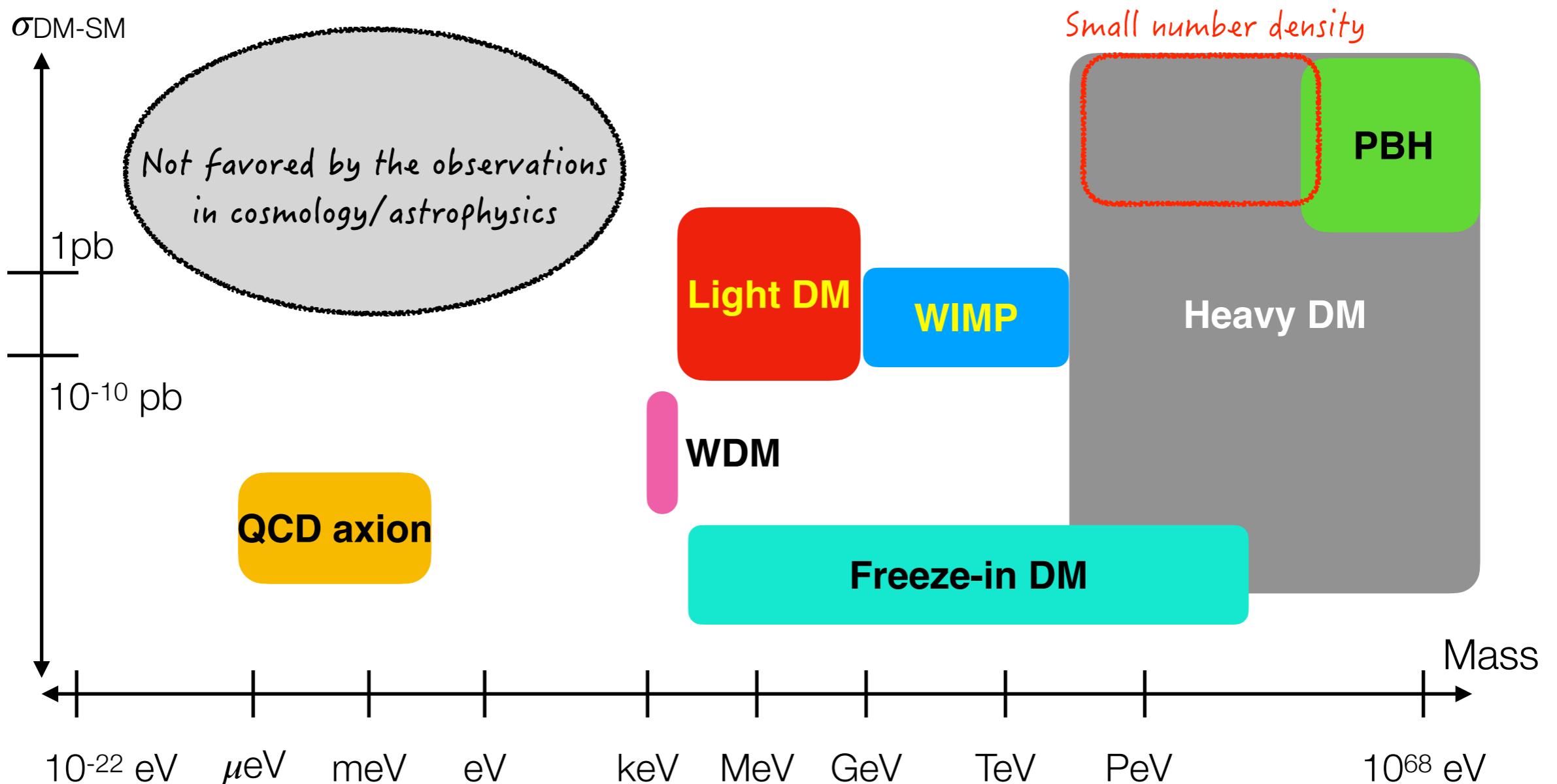


Dark world beyond WIMP



- Wavelike DM: occupation number is large so that the particles are described by classical waves, e.g., axion or ALP
- Experimental probes: axion or ALP search experiments See also Ann. Rev. Nucl. Part. Sci. (2021) by Choi, Im, Shin

Dark world beyond WIMP

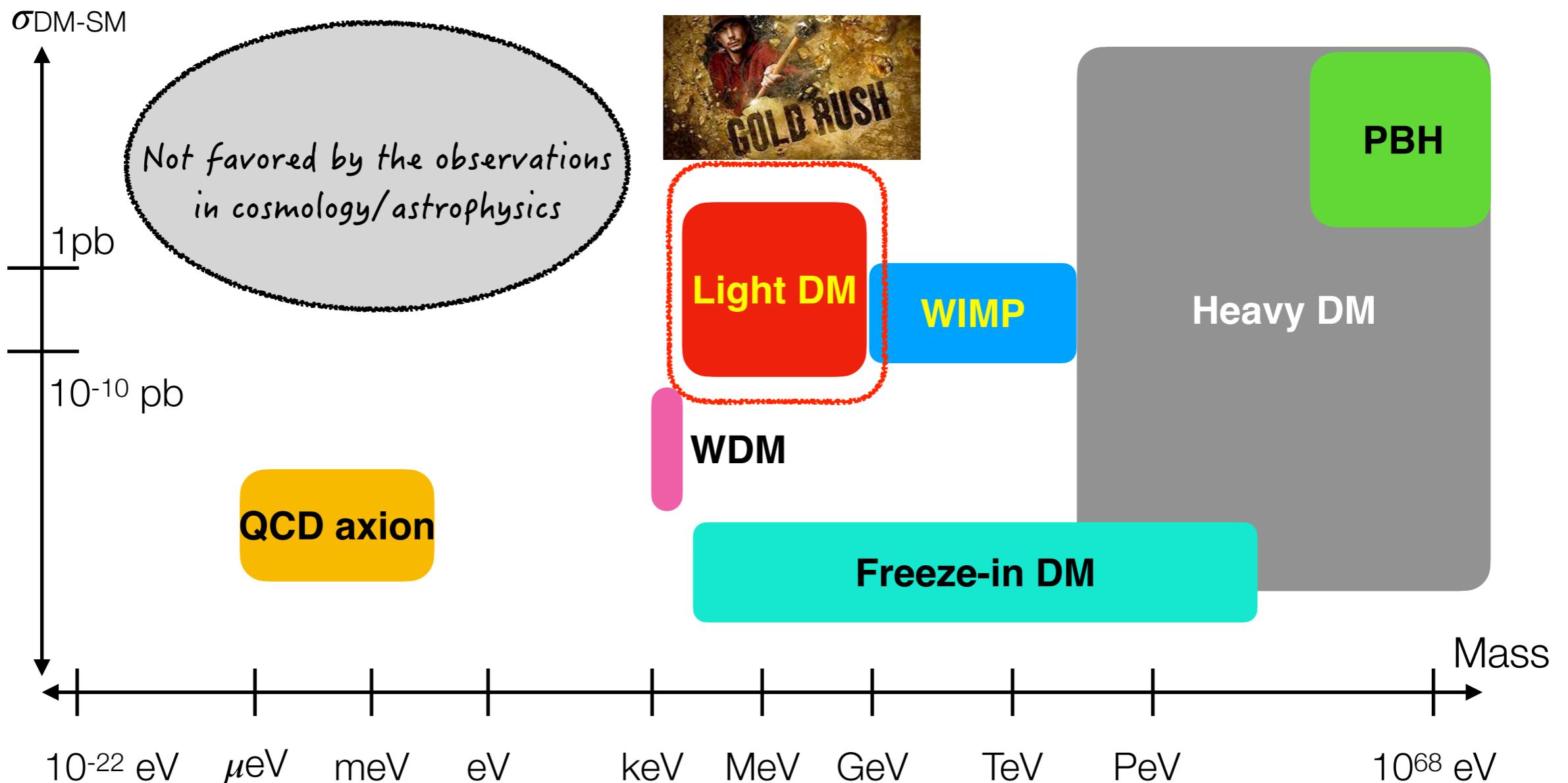


- Heavy DM (**small n**) with large $\sigma_{\text{DM-SM}}$: multiple scattering signals

Bramante et al., PRD 2019, PRD 2018

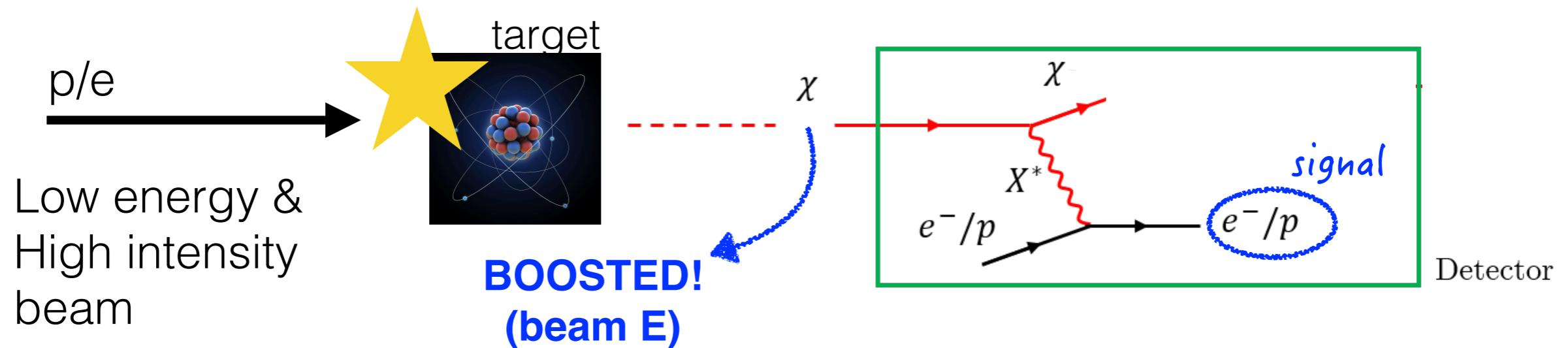
- Primordial Black Hole: cosmo/astro observations

Dark world beyond WIMP

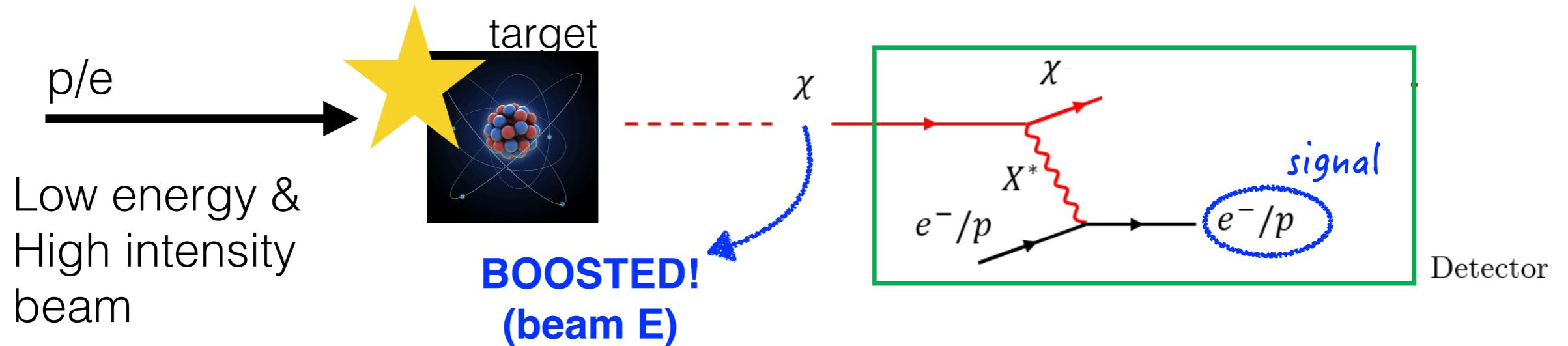


- Light CDM (keV ~ sub GeV): high intensity low energy accelerators, DM - e recoil, ν -experiments, B-factories long-lived particle searches, etc.
New ideas!

Example: intensity frontier



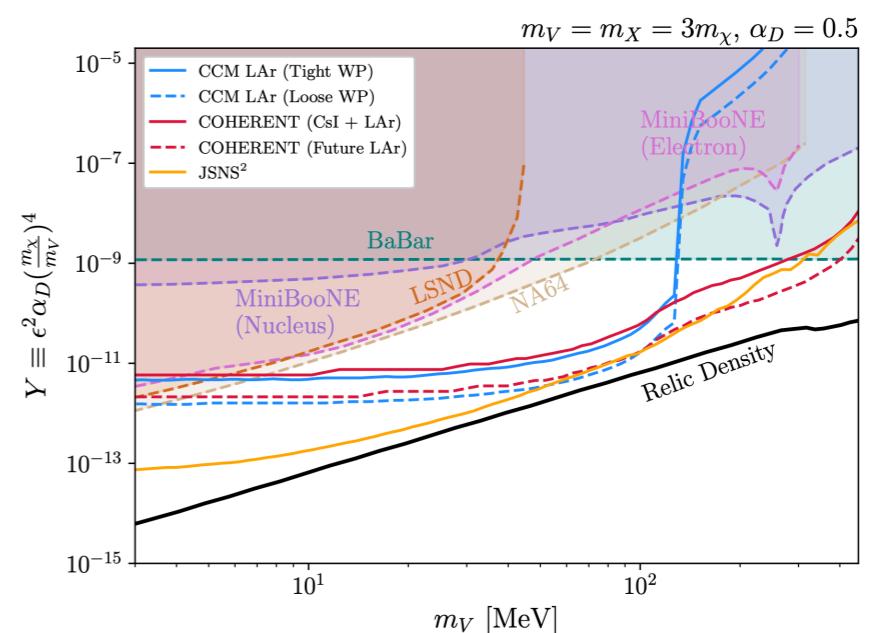
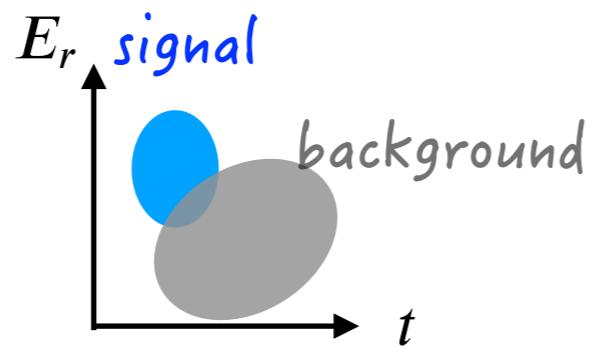
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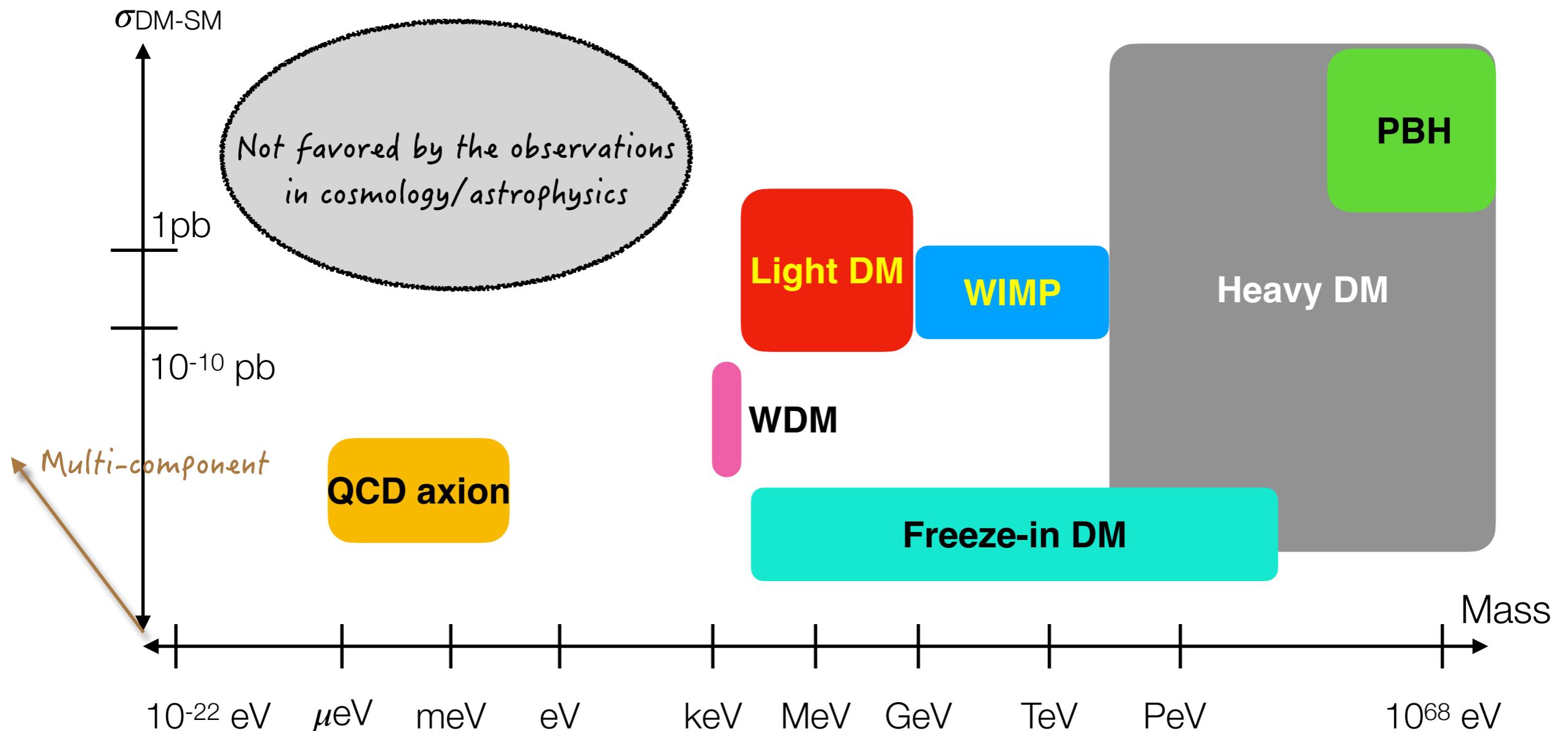
Pulsed beam with high intensity

- Increase the signal sensitivity enormously by using **timing & energy** information together. Dutta, Kim, Liao, Park, **ss**, Strigari, PRL 2020
- Applied to COHERENT, CCM, JSNS², ...

Dutta, Kim, Liao, Park, **ss**, Strigari, Thompson, JHEP 2022



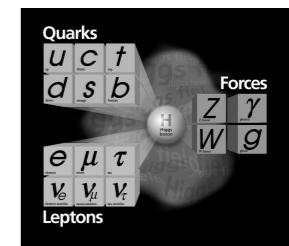
Multi-component dark sector



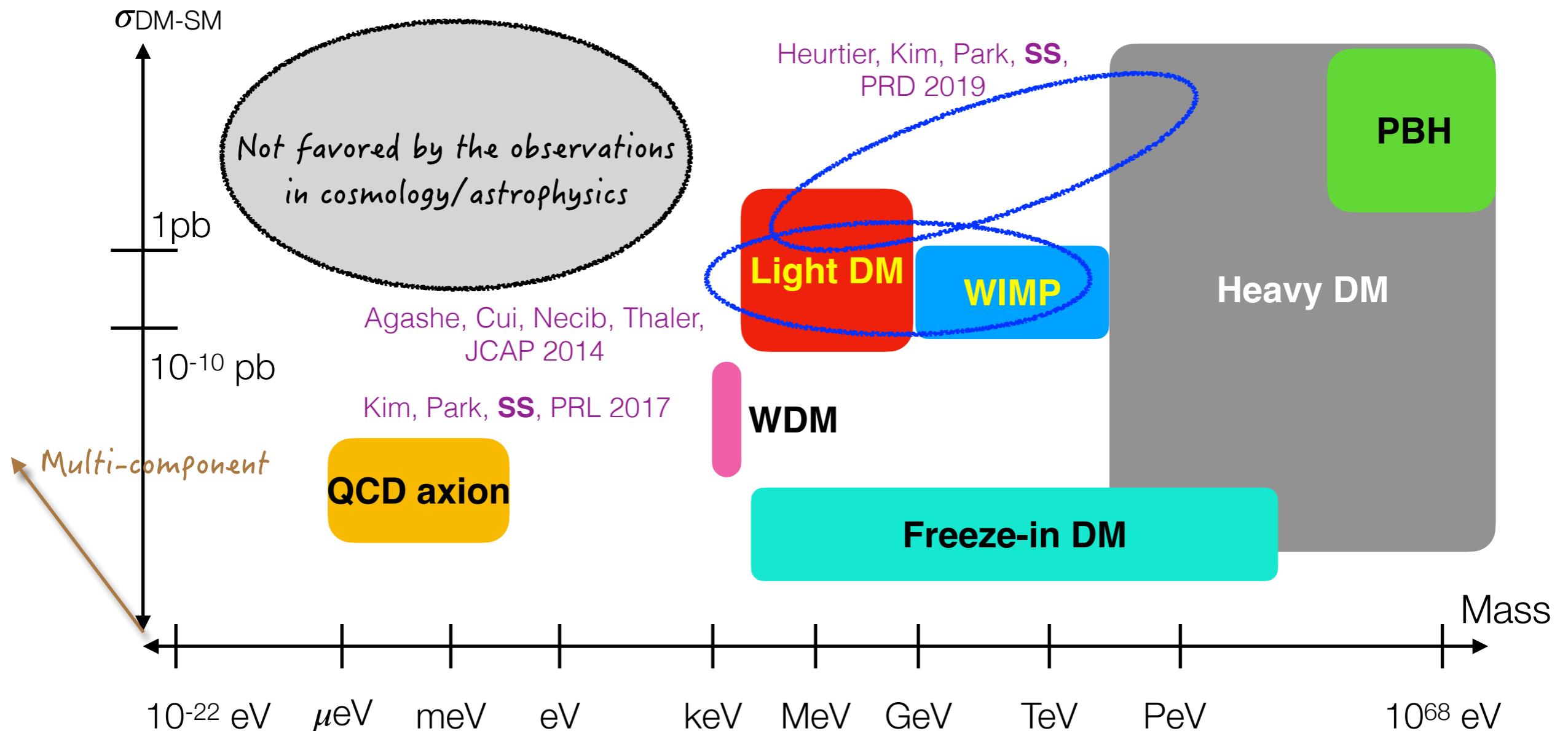
- Non-trivial dark sector structure: e.g., iDM

Smith, Weiner, PRD 2001

- Hybrid models with fruitful experimental signatures: e.g., BDM



Multi-component dark sector



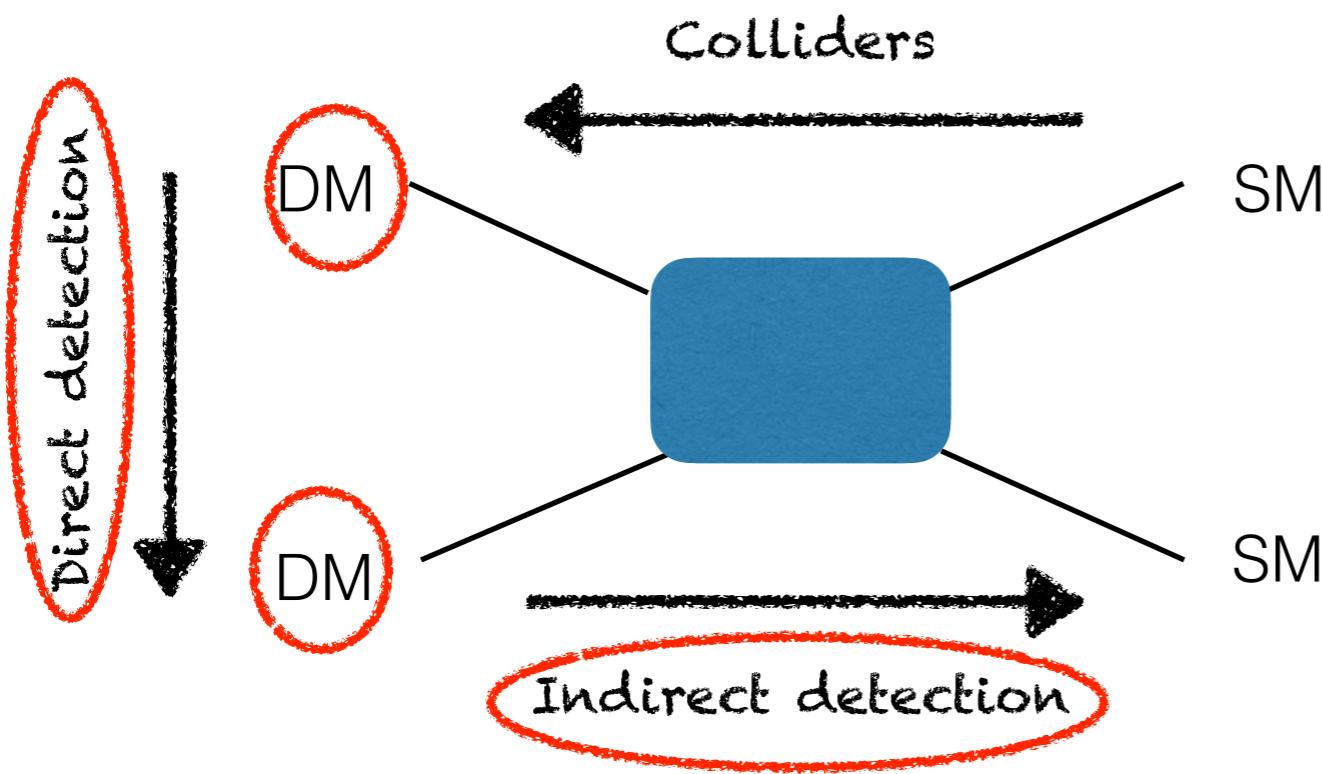
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Sub-dominant component is hidden?

- Conventionally, sub-dominant DM components are thought to be hidden in direct/indirect detection experiments: observables \propto fraction
- The dominant relics might be safe from the experiments so far if it communicates with the SM sector through the sub-dominant relic?
- Question is how the amount of the sub-dominant relic is determined.

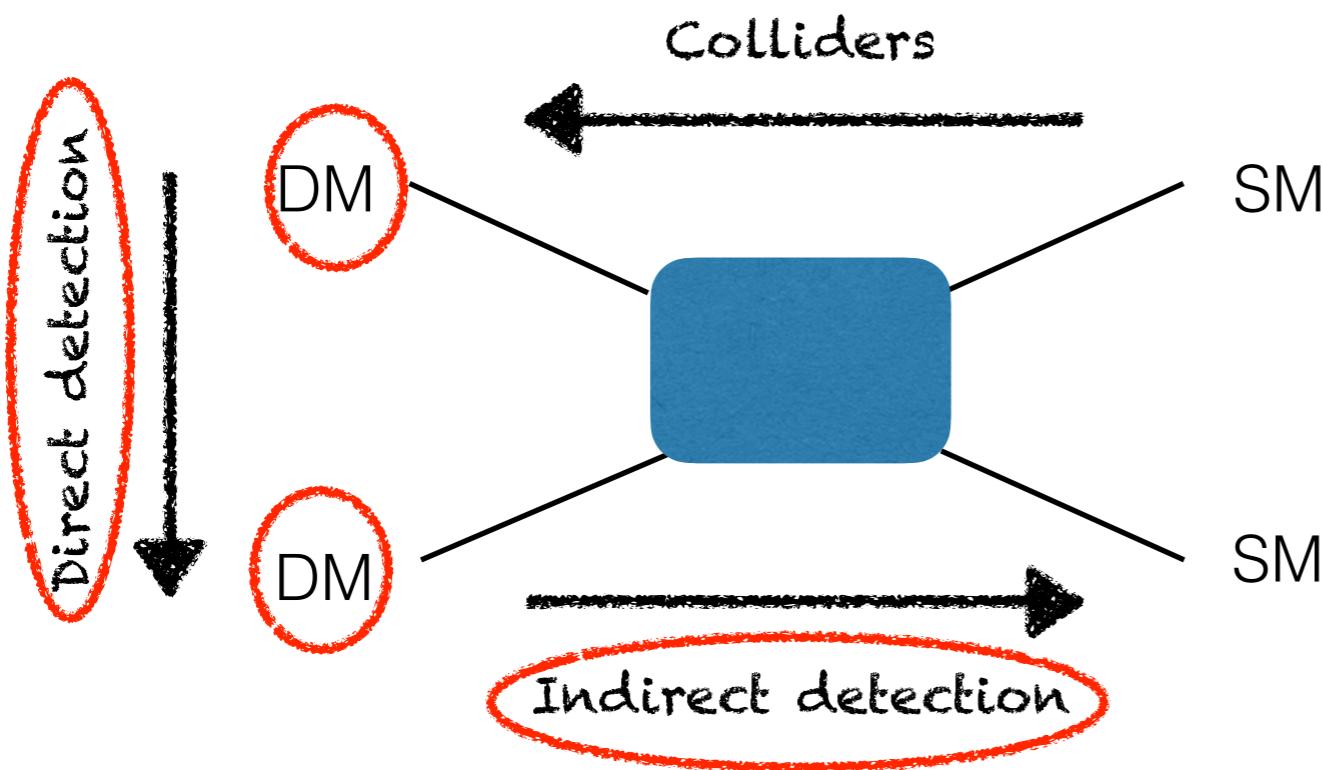
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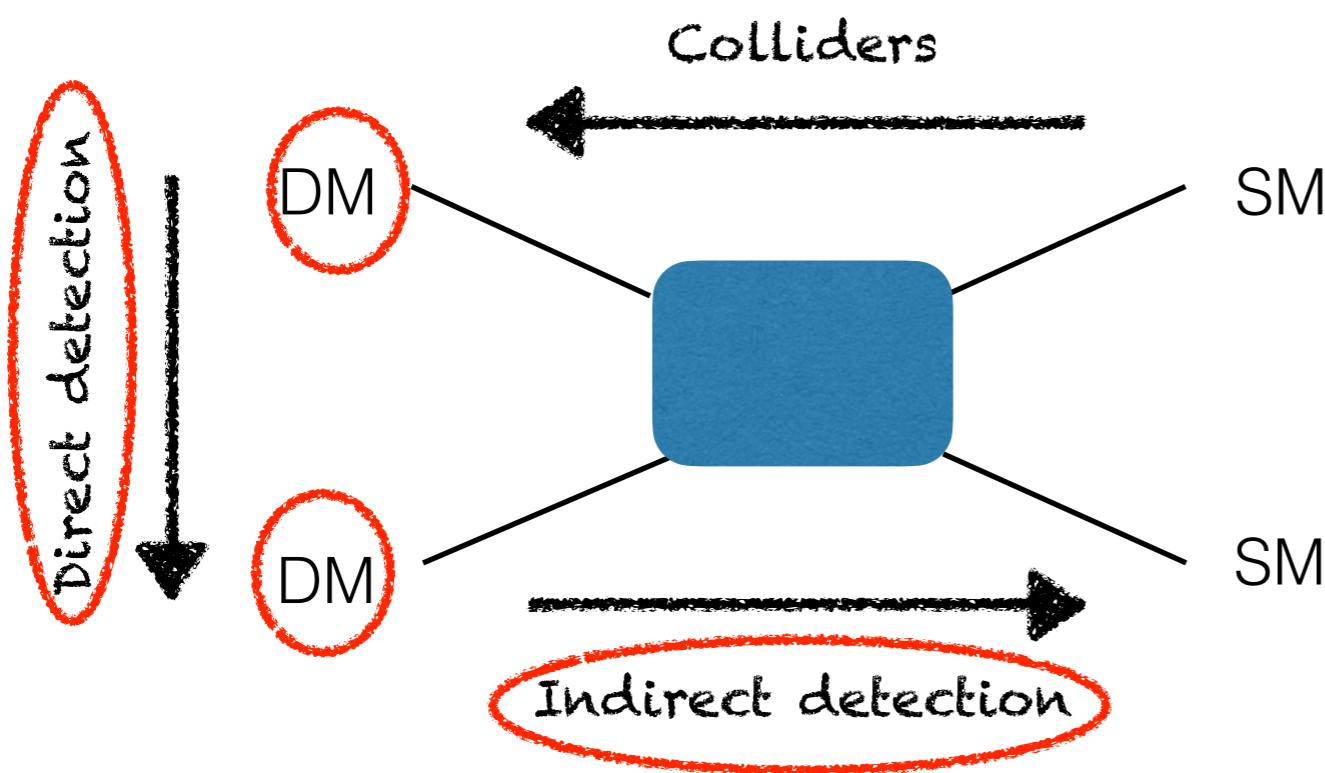
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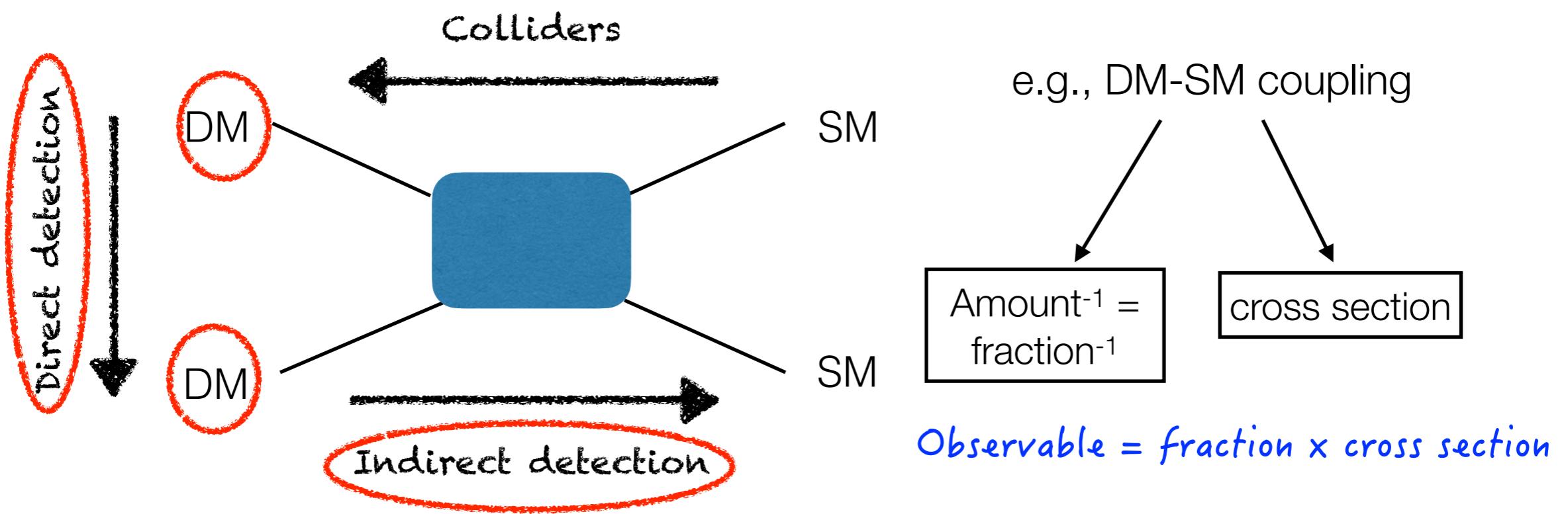
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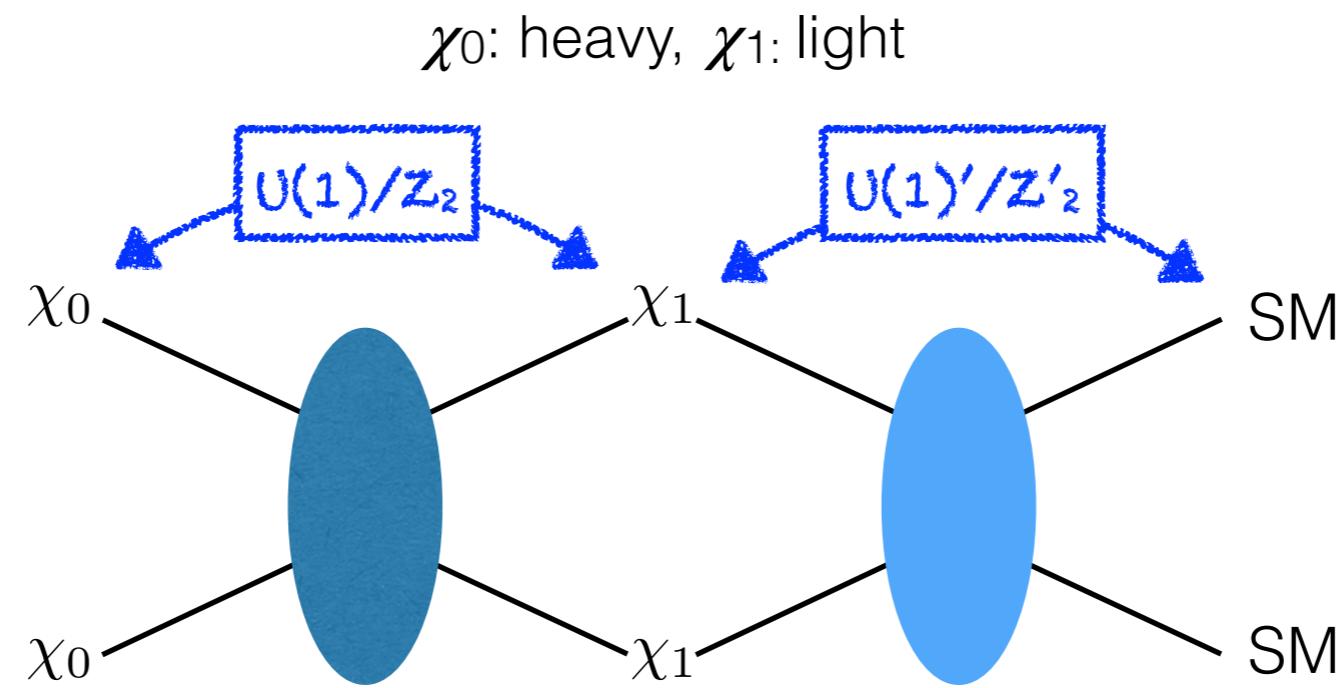


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Reference: Multi-component BDM

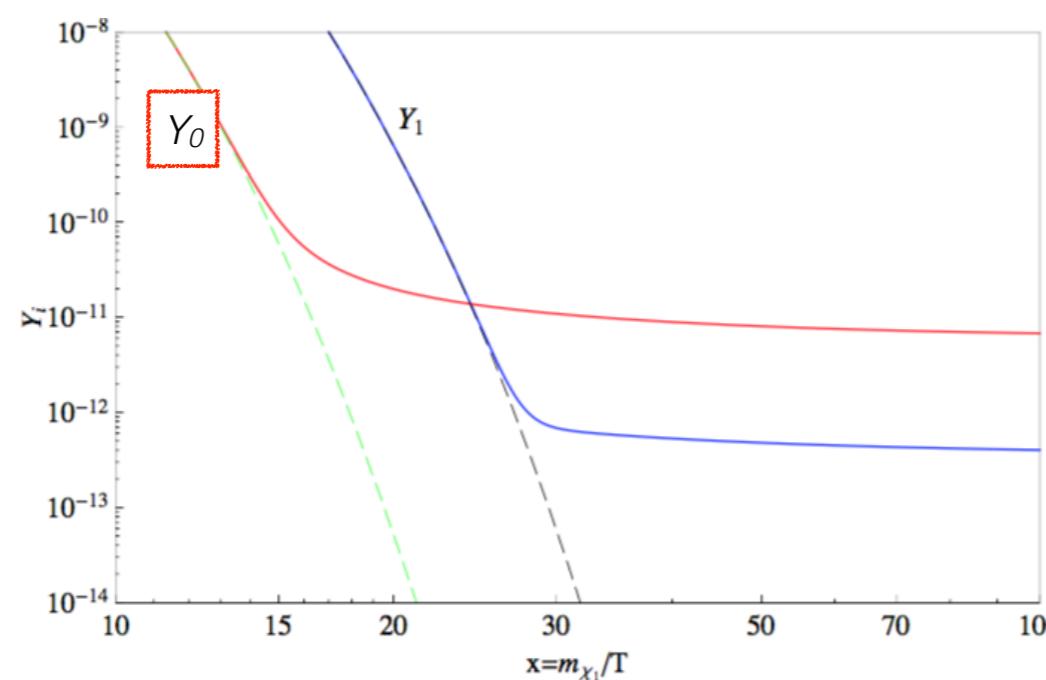
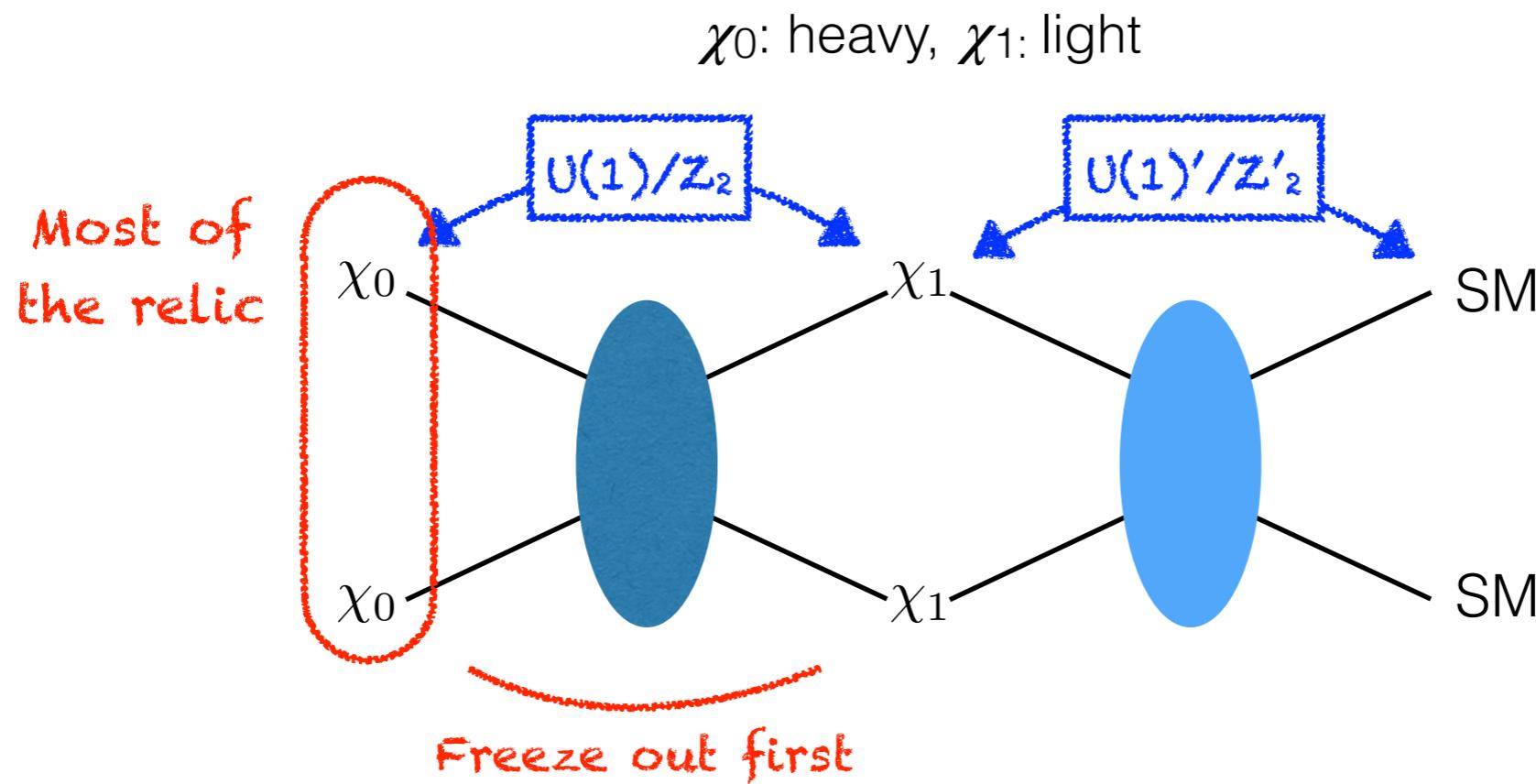


Agashe, Cui, Necib, Thaler, JCAP 2014

Kim, Park , **ss**, PRL 2017

Giudice, Kim, Park , **ss**, PLB 2018

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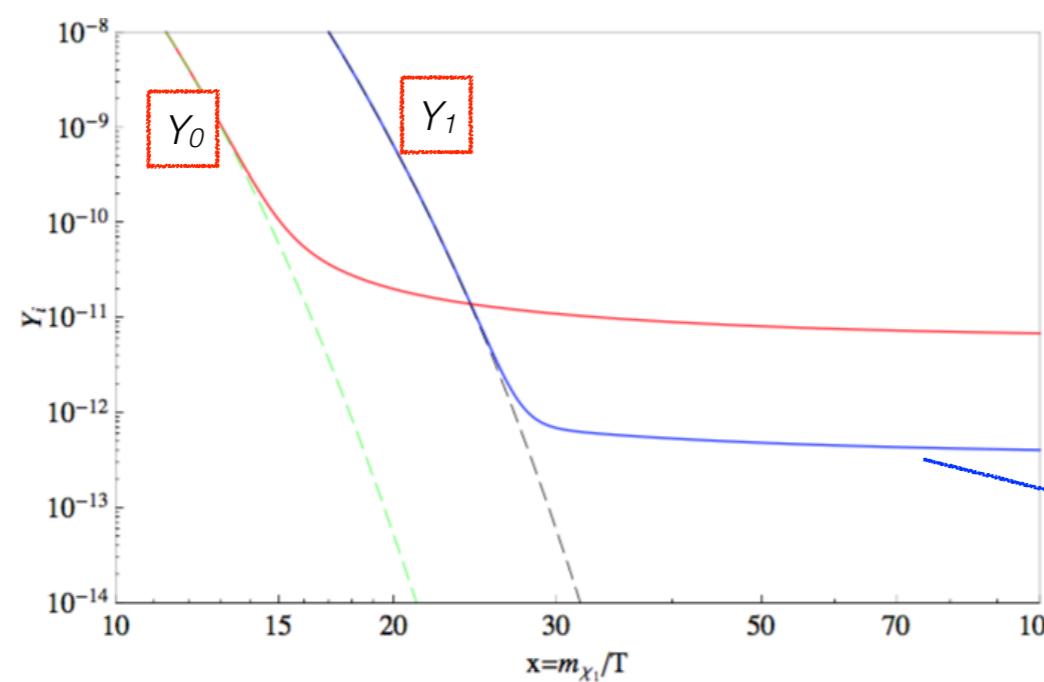
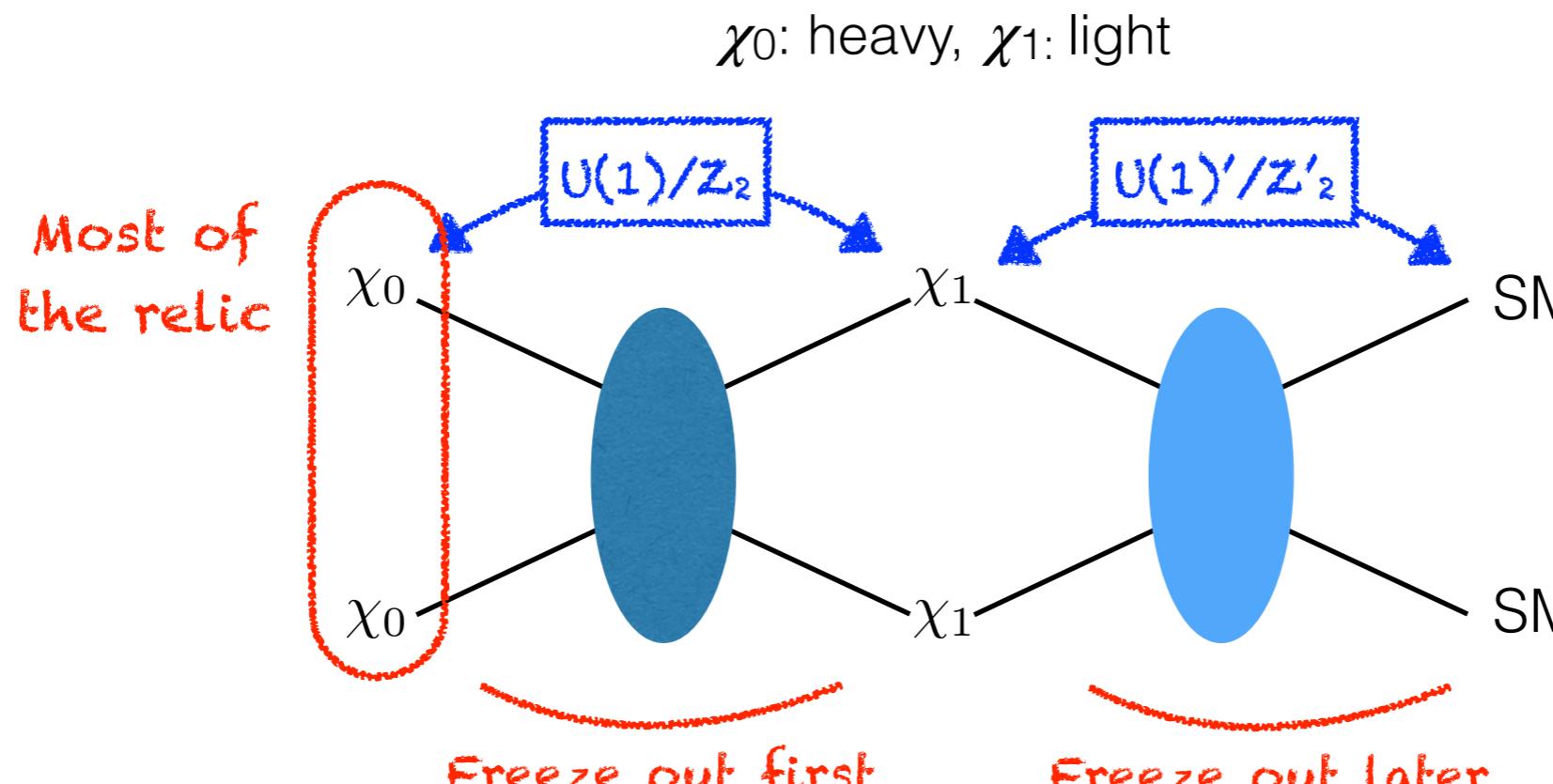


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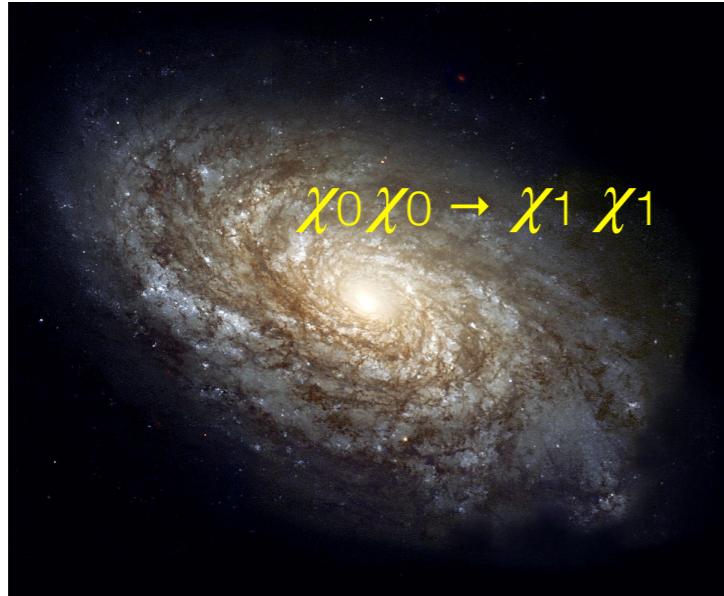
Belanger, Park, JCAP 2012

Assisted freeze-out mechanism

non-relativistic relic χ_1 (negligible)

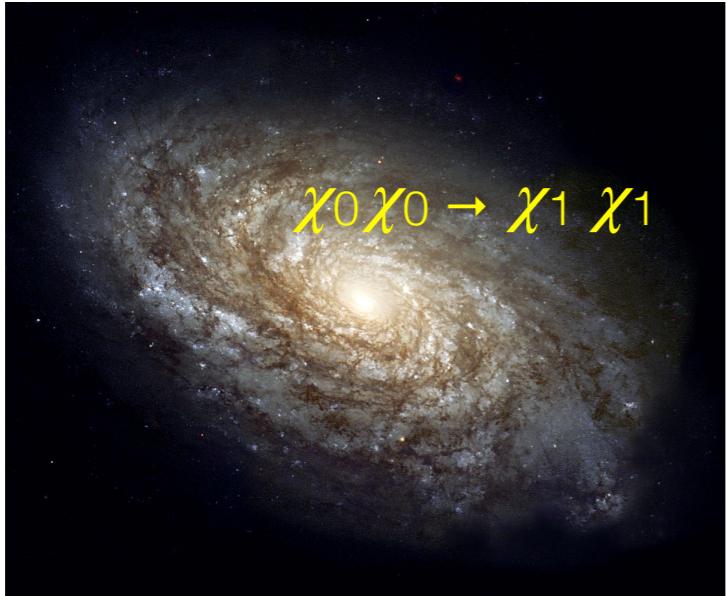
$$Y_0 \gg Y_1$$

Reference: Multi-component BDM



- χ_0 : accumulated
(GC, Sun, dSphs)
- $\chi_0 \chi_0 \rightarrow \chi_1 \chi_1$ (current universe) **relativistic**
※ relic χ_1 is non-relativistic

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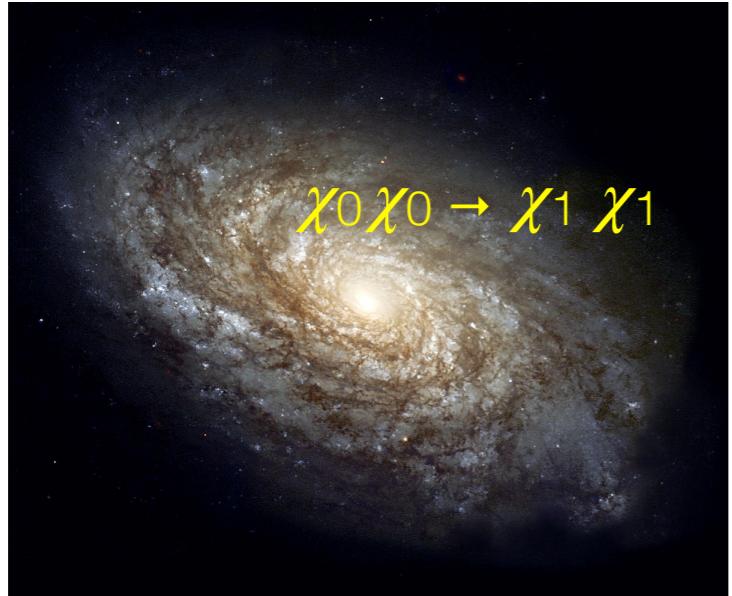


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Observe χ_1 scattering off target with $E_1 > E_{th}$
(indirect detection of χ_0)

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$$\text{Flux of } \chi_1 \simeq 1.6 \times 10^{-8} \text{ cm}^{-2} \text{s}^{-1} \times \left(\frac{\langle \sigma v \rangle_{0 \rightarrow 1}}{5 \times 10^{-26} \text{ cm}^3 \text{s}^{-1}} \right) \times \left(\frac{100 \text{ GeV}}{m_0} \right)^2$$

Assume: NFW

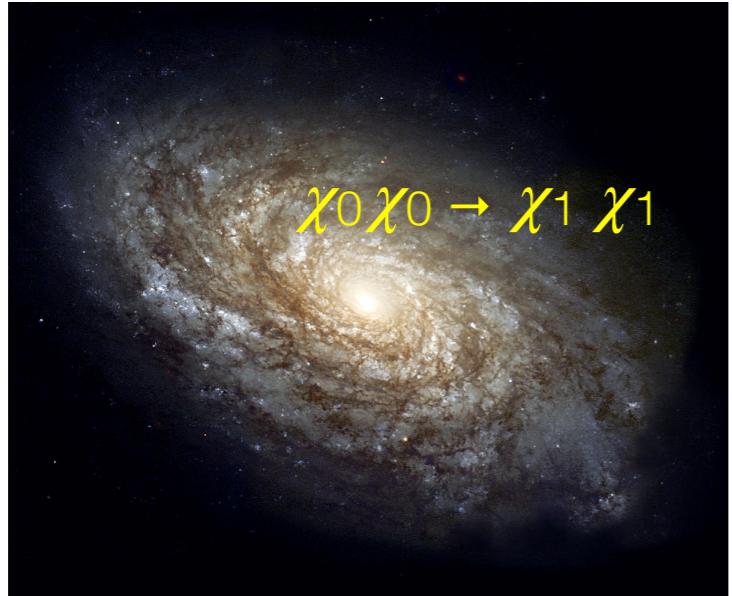


Fixed ~ 1 if **s-wave** annihilation dominates (throughout this work for simplicity)

10,000 times smaller than the flux of atmospheric ν if $m_0 \sim 100$ GeV

Agashe et al.,
JCAP 2014
Kim, Park , **SS**,
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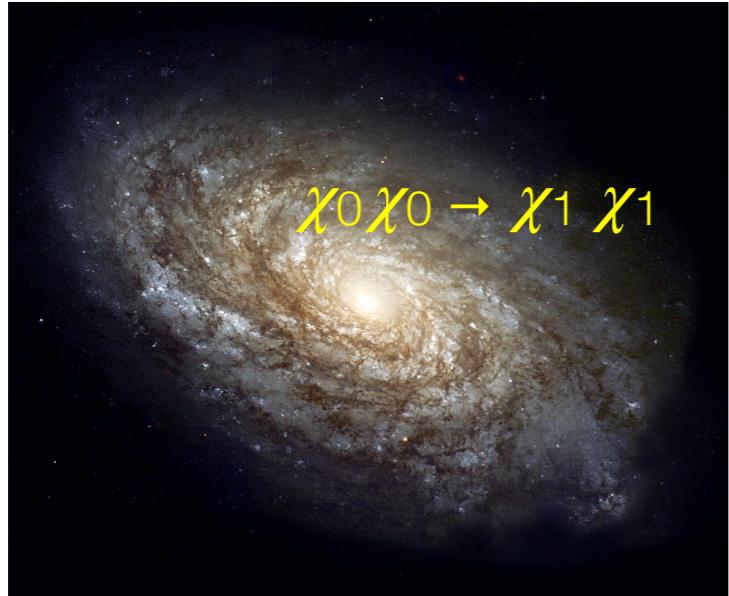
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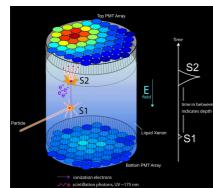
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comparable

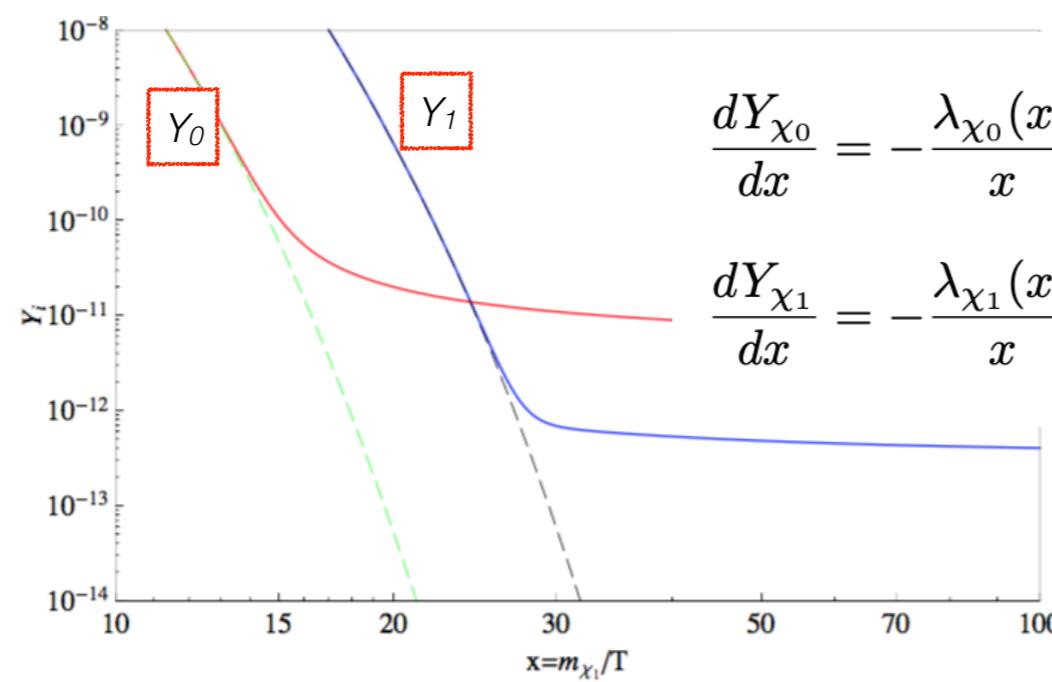
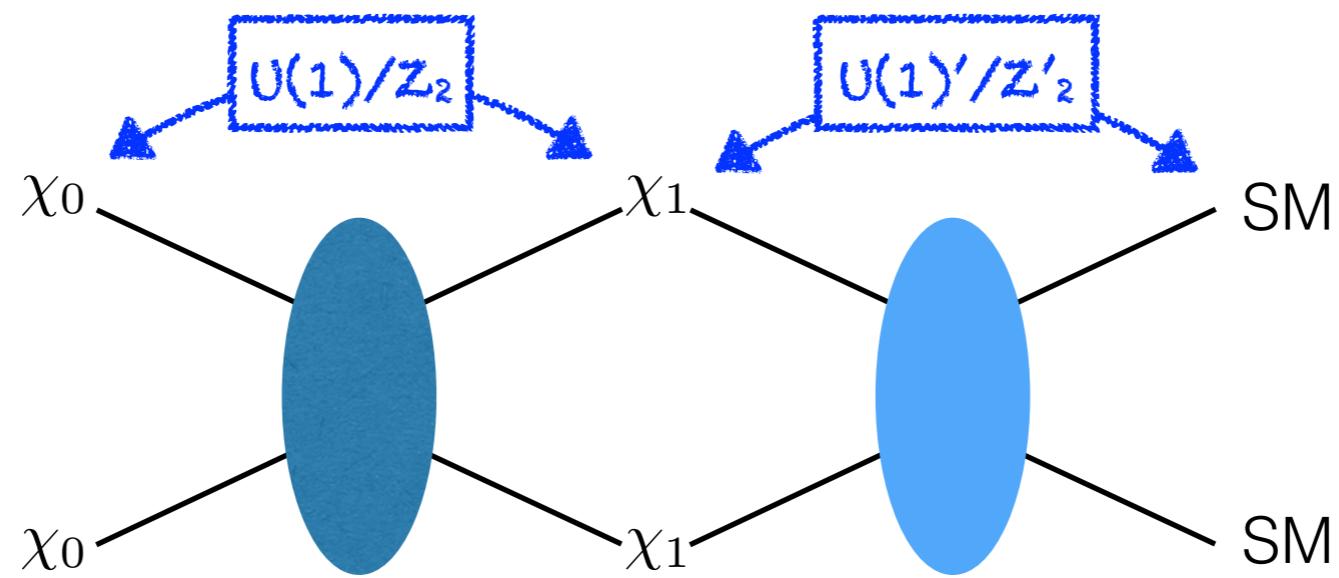
if $m_0 \lesssim 1 \text{ GeV}$

Giudice, Kim, Park , **SS**, PLB 2018

Agashe et al.,
JCAP 2014
Kim, Park , **SS**,
PRL 2017

Structure of $\chi_1 \chi_1 \rightarrow \text{SM}$

χ_0 : heavy (dominant), χ_1 : light (subdominant)



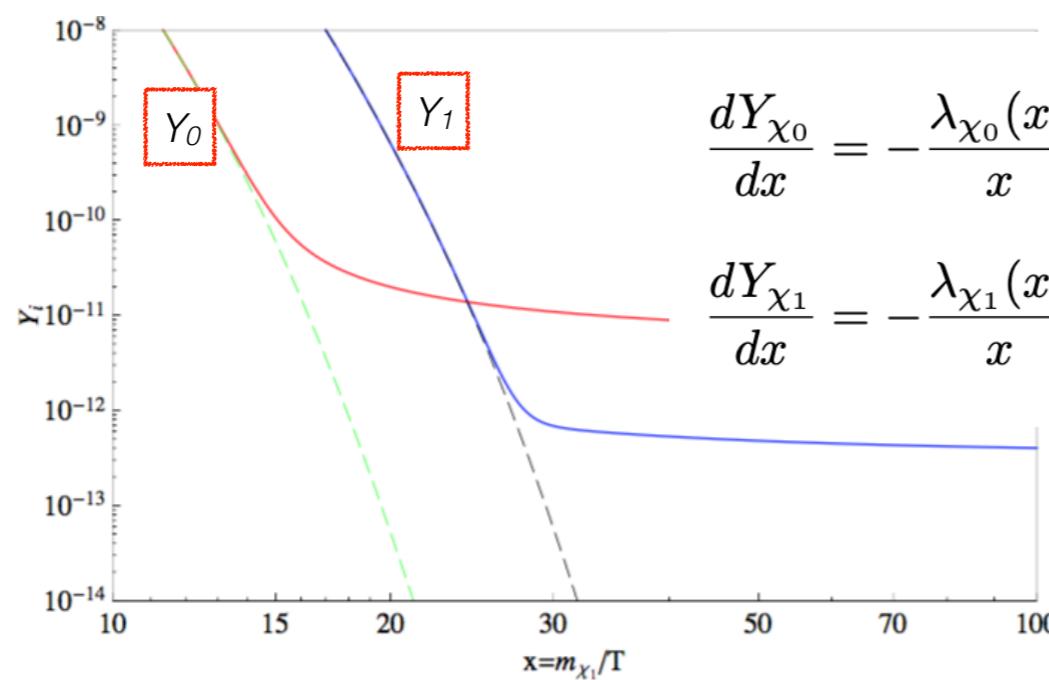
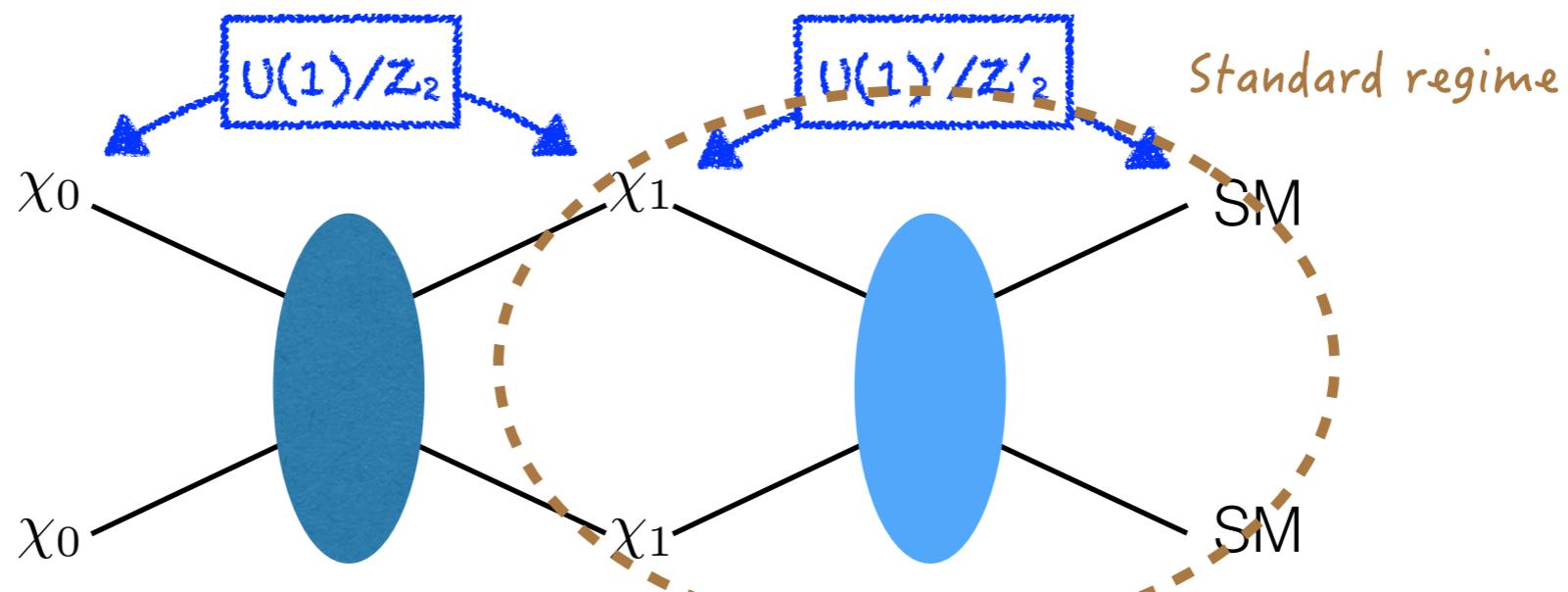
$$\frac{dY_{\chi_0}}{dx} = -\frac{\lambda_{\chi_0}(x)}{x} \left[Y_{\chi_0}^2 - \left(\frac{Y_{\chi_0}^{\text{eq}}(x)}{Y_{\chi_1}^{\text{eq}}(x)} \right)^2 Y_{\chi_1}^2 \right],$$

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$$\lambda_{\chi_i} = s \langle \sigma_i v_{\text{rel}} \rangle / H$$

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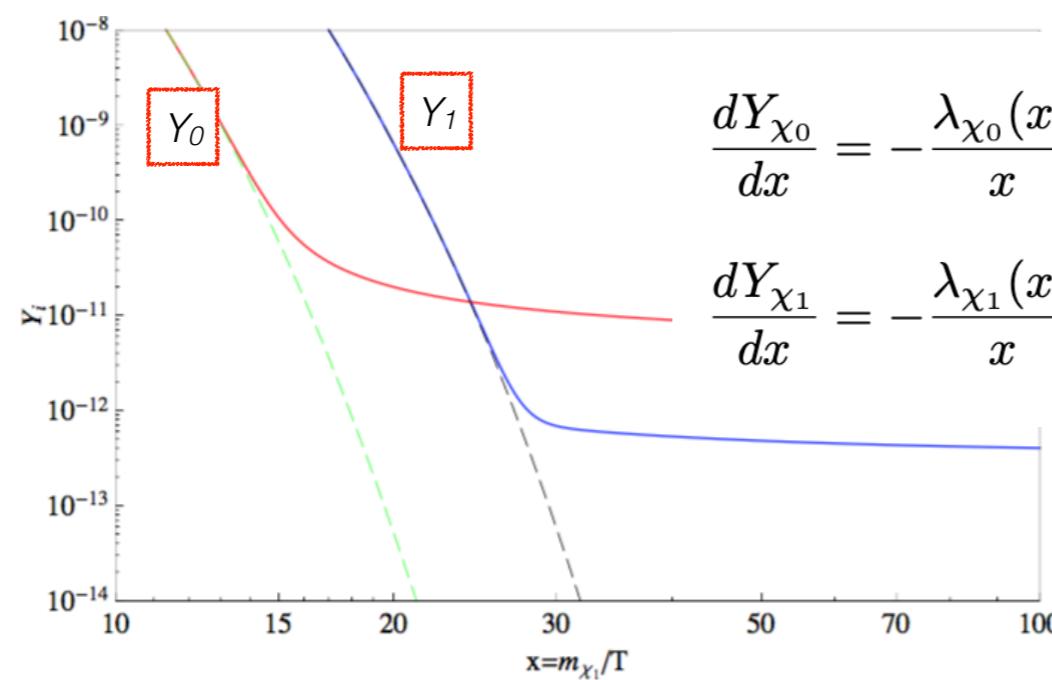
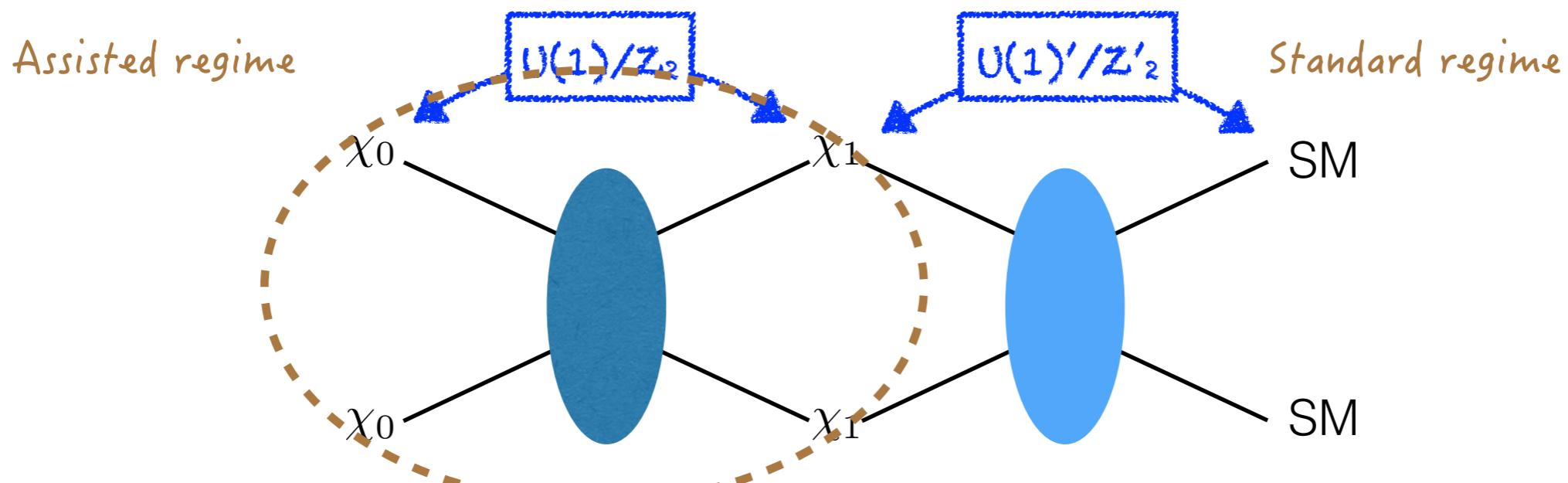
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with SM

$$\lambda_{\chi_i} = s \langle \sigma_i v_{\text{rel}} \rangle / H$$

Structure of $\chi_1 \chi_1 \rightarrow \text{SM}$

χ_0 : heavy (dominant), χ_1 : light (subdominant)



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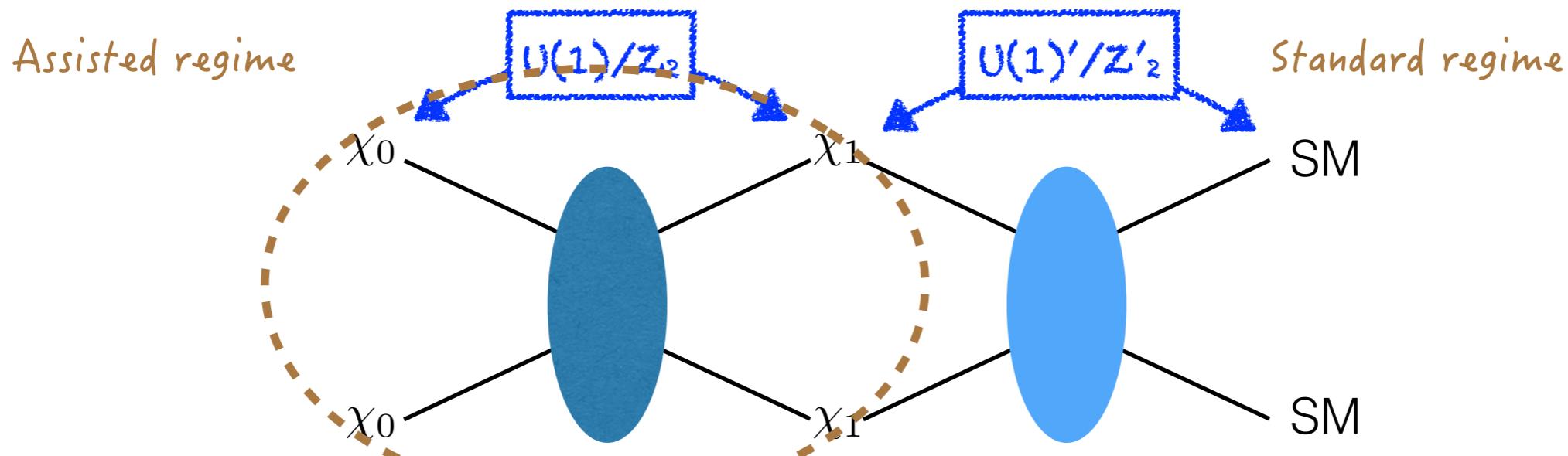
with SM

with heavy DM χ_0

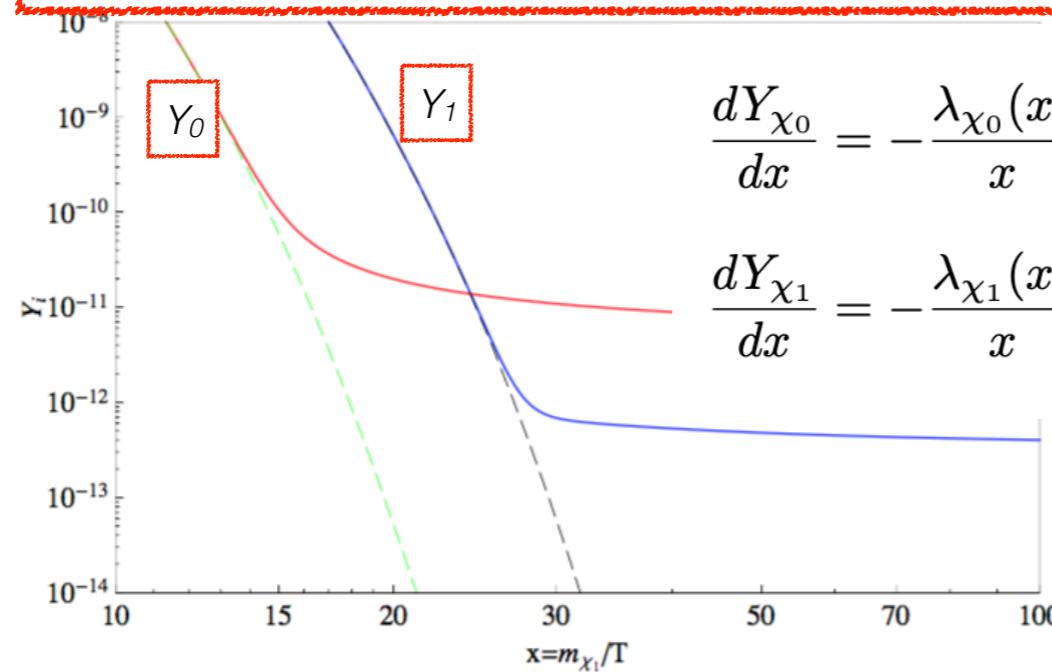
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χ_0 : heavy (dominant), χ_1 : light (subdominant)



Assumption: $\chi_0\chi_0 \rightarrow \chi_1\chi_1$ is *s*-wave & the mediator χ_1 - SM is heavier than χ_1 .



$$\frac{dY_{\chi_0}}{dx} = -\frac{\lambda_{\chi_0}(x)}{x} \left[Y_{\chi_0}^2 - \left(\frac{Y_{\chi_0}^{\text{eq}}(x)}{Y_{\chi_1}^{\text{eq}}(x)} \right)^2 Y_{\chi_1}^2 \right],$$

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with SM

with heavy DM χ_0

$$\lambda_{\chi_i} = s \langle \sigma_i v_{\text{rel}} \rangle / H$$

Structure of $\chi_1 \chi_1 \rightarrow \text{SM}$

After the heavy component χ_0 freezes-out

$$\frac{dY_{\chi_1}}{dx} \simeq -\frac{\lambda_{\chi_1}(x)}{x} \left[Y_{\chi_1}^2 - \underline{(Y_{\chi_1}^{\text{eq}}(x))^2} - Y_{\text{ast.}}^2(x) \right]$$



where $Y_{\text{ast.}}(x) = \sqrt{\frac{\langle \sigma_0 v_{\text{rel}} \rangle}{\langle \sigma_1 v_{\text{rel}} \rangle}} Y_{\chi_0}(x)$ $r_1 = \frac{\Omega_{\chi_1}}{\Omega_{\text{DM,tot}}}$

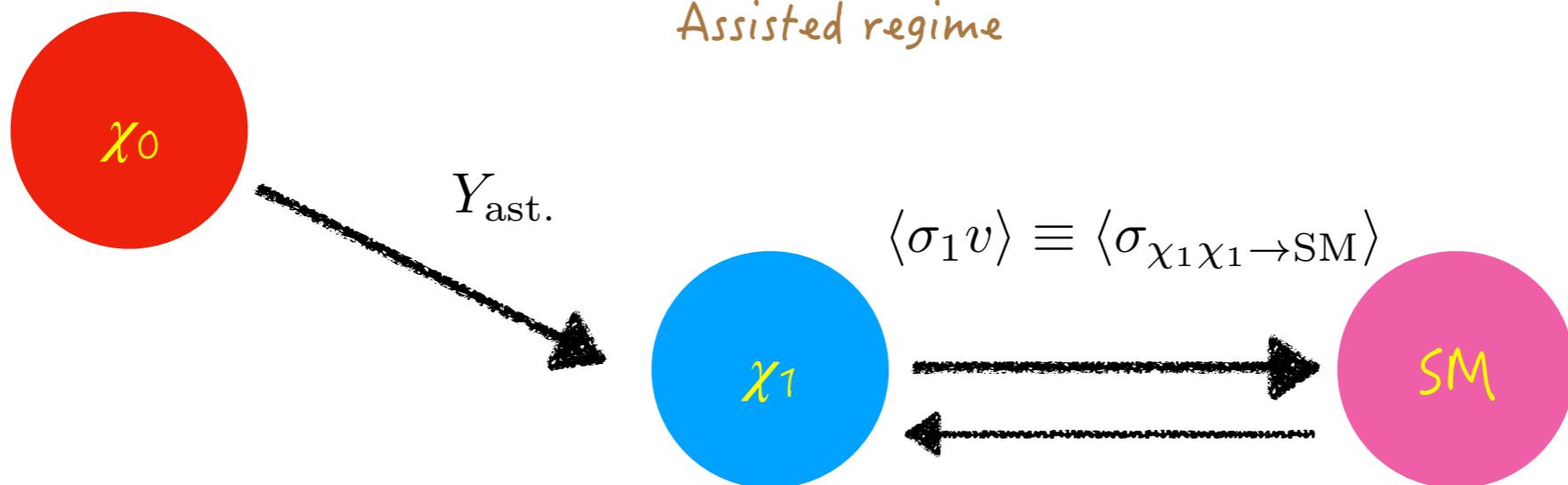
During the decoupling, assume χ_1 is in kinetic equilibrium with the SM

$$(\sigma_{\chi^0-\chi^1} < \sigma_{\chi^1-\text{SM}})$$

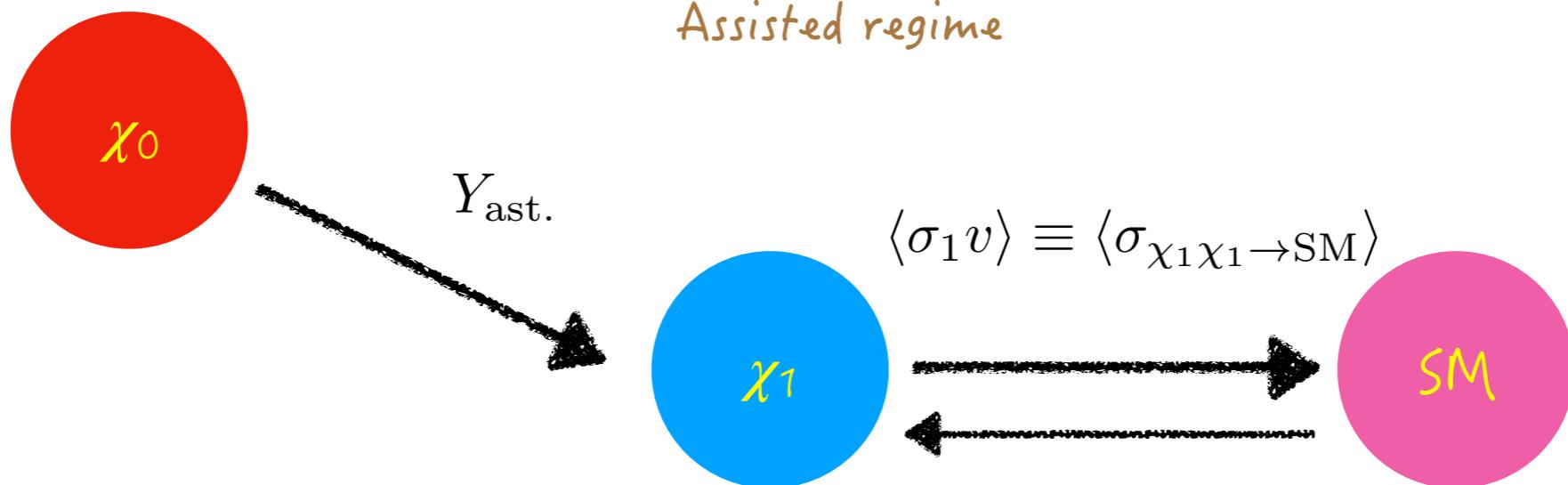
- If $Y_{\text{ast.}}$ is negligible, χ_1 freezes out at $T \sim m_1/20$ as usual. *Standard regime*
- If the fraction of χ_1 is very small, i.e., $r_1 \ll 1$, however, departure from thermal equilibrium is delayed and $Y_{\text{ast.}}$ is **non-negligible** compared to $Y_{\chi_1}^{\text{eq}}$

Assisted regime

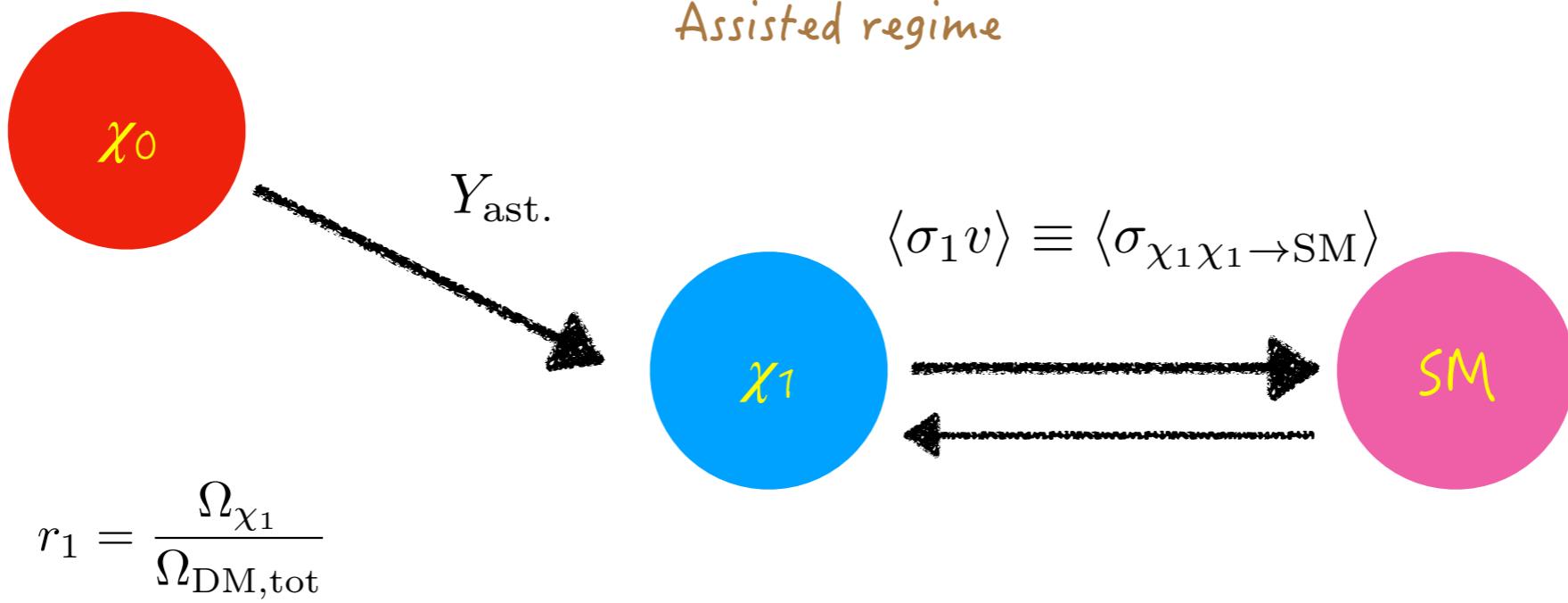
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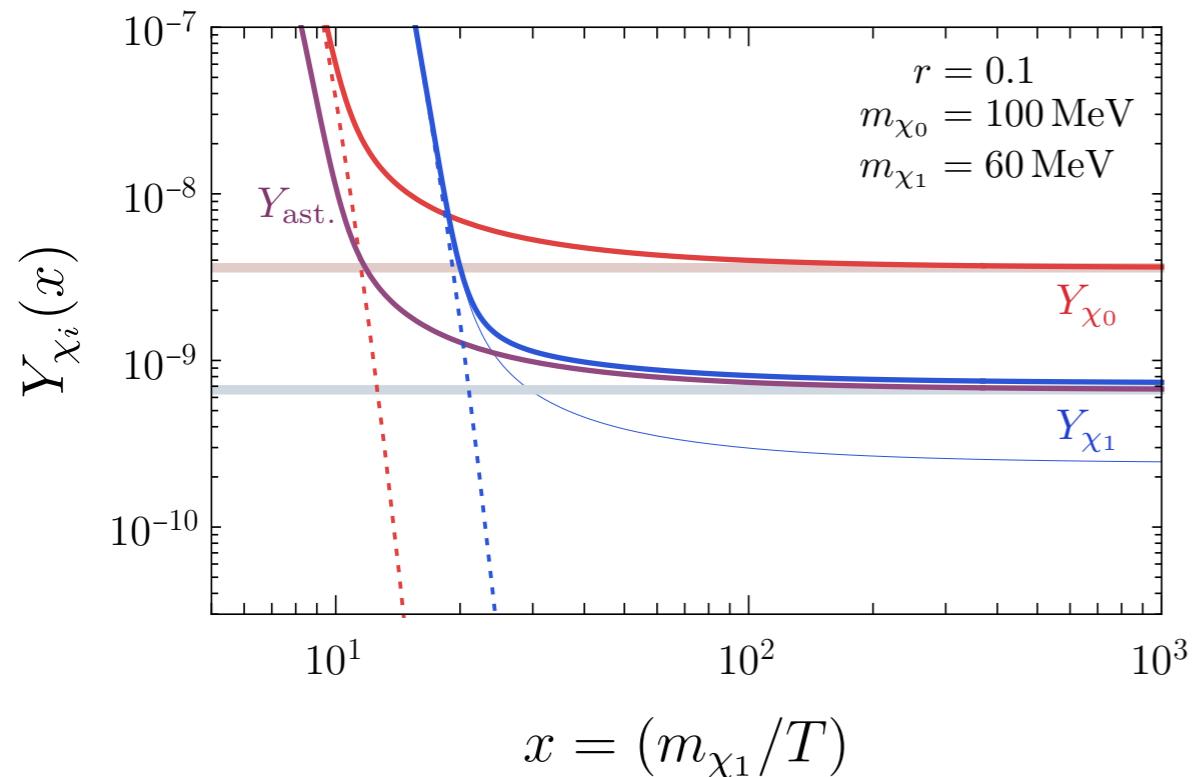
Structure of $\chi_1\chi_1 \rightarrow \text{SM}$



- For a fixed $r_1 \ll 1$, $\chi_1\chi_1 \rightarrow \text{SM}$ should be even larger to deplete the contribution by the residual annihilation $\chi_0\chi_0 \rightarrow \chi_1\chi_1$ ($Y_{\text{ast.}}$).
- We find for s-wave and p-wave, respectively.

Structure of $\chi_1\chi_1 \rightarrow \text{SM}$

When $\chi_1\chi_1 \rightarrow \text{SM}$ is dominated by s-wave



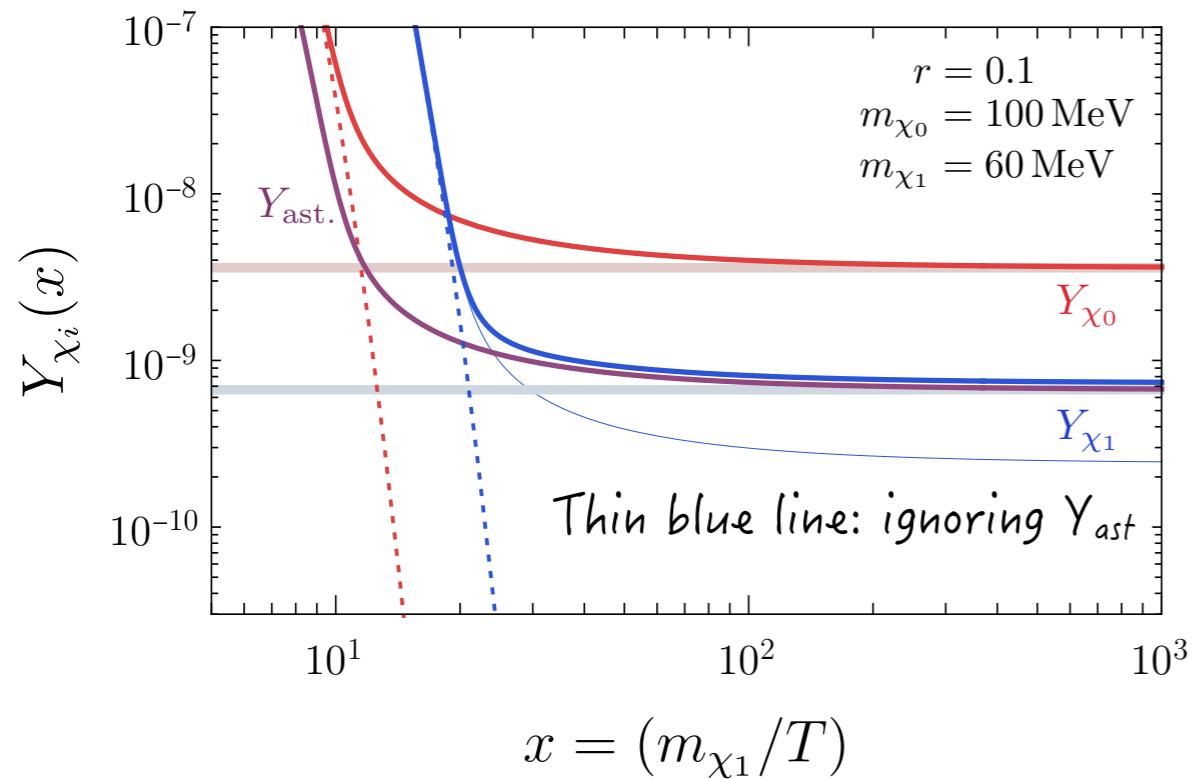
Assisted regime

Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

$$Y_{\text{ast.}}(x) = \sqrt{\frac{\langle \sigma_0 v_{\text{rel}} \rangle}{\langle \sigma_1 v_{\text{rel}} \rangle}} Y_{\chi_0}(x)$$

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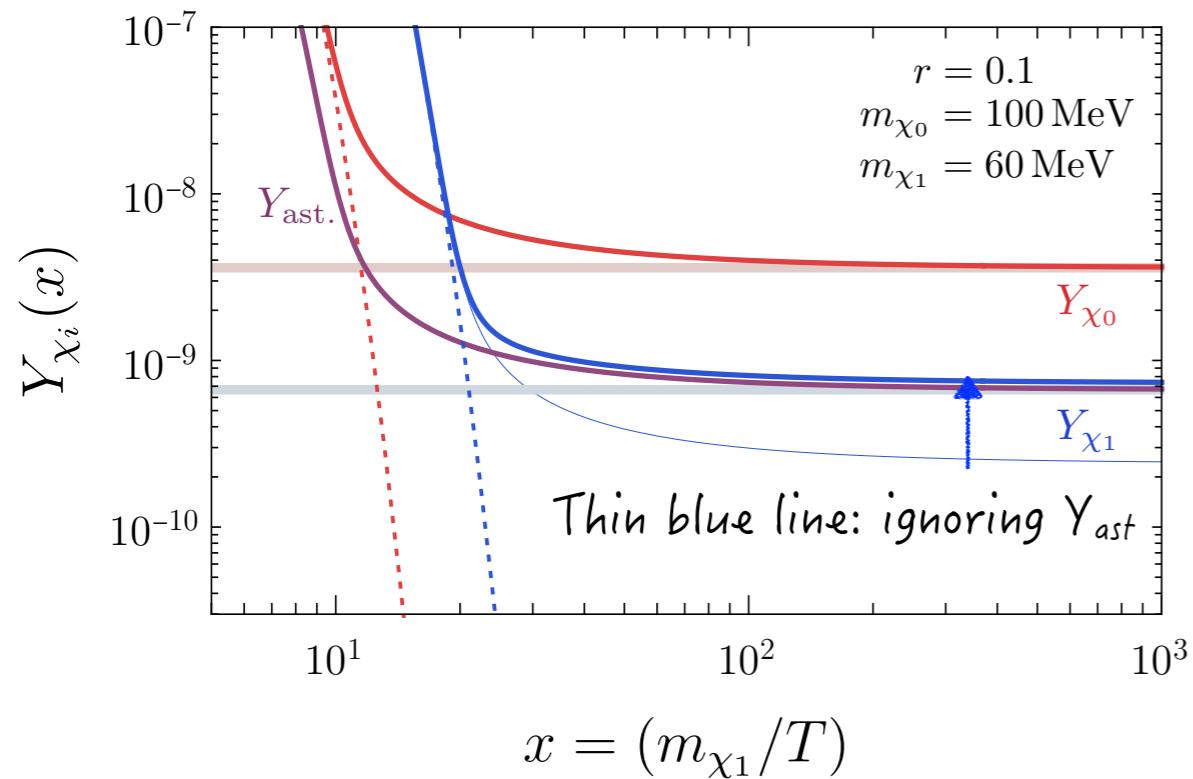
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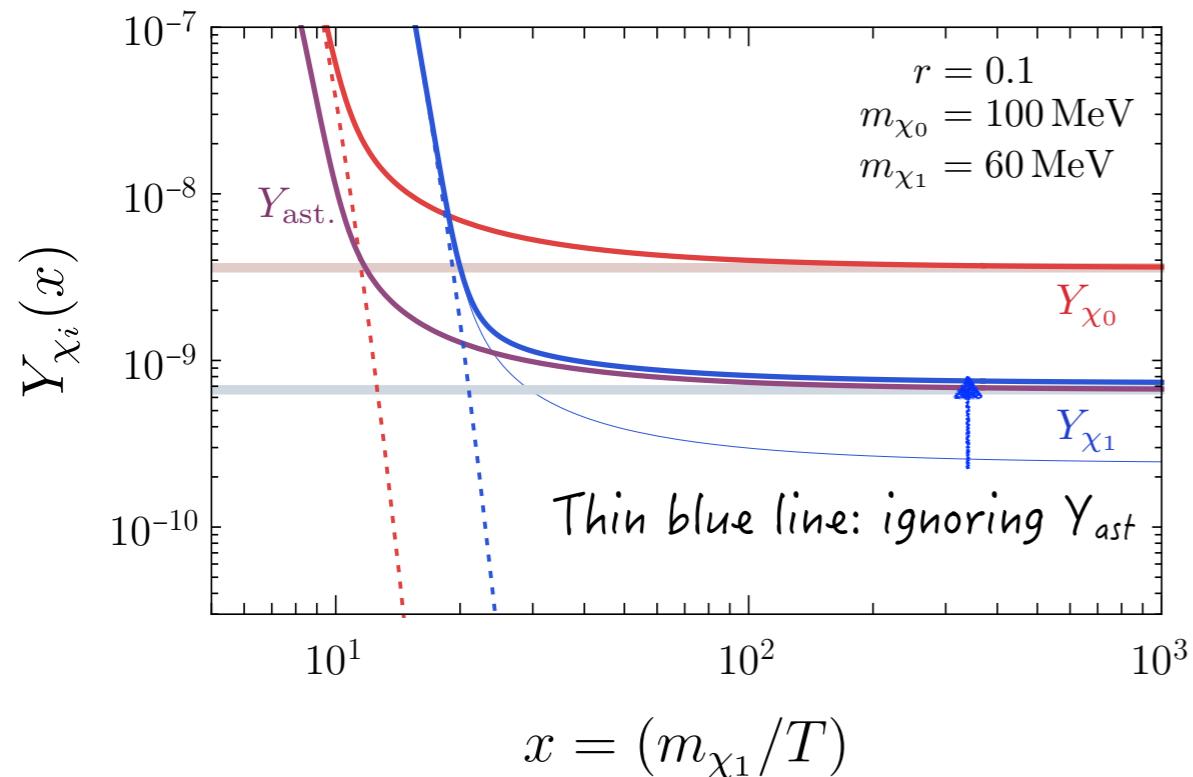
Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

$$r_1 = \frac{\Omega_{\chi_1}}{\Omega_{\text{DM,tot}}} \quad Y_{\text{ast.}}(x) = \sqrt{\frac{\langle \sigma_0 v_{\text{rel}} \rangle}{\langle \sigma_1 v_{\text{rel}} \rangle}} Y_{\chi_0}(x)$$

- For $r_1 \ll 1$, Y_{χ_1} is lifted-up by $Y_{\text{ast.}}$ (follows it when $T \lesssim m_1/30$).

Structure of $\chi_1\chi_1 \rightarrow \text{SM}$

When $\chi_1\chi_1 \rightarrow \text{SM}$ is dominated by s-wave



Assisted regime

Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

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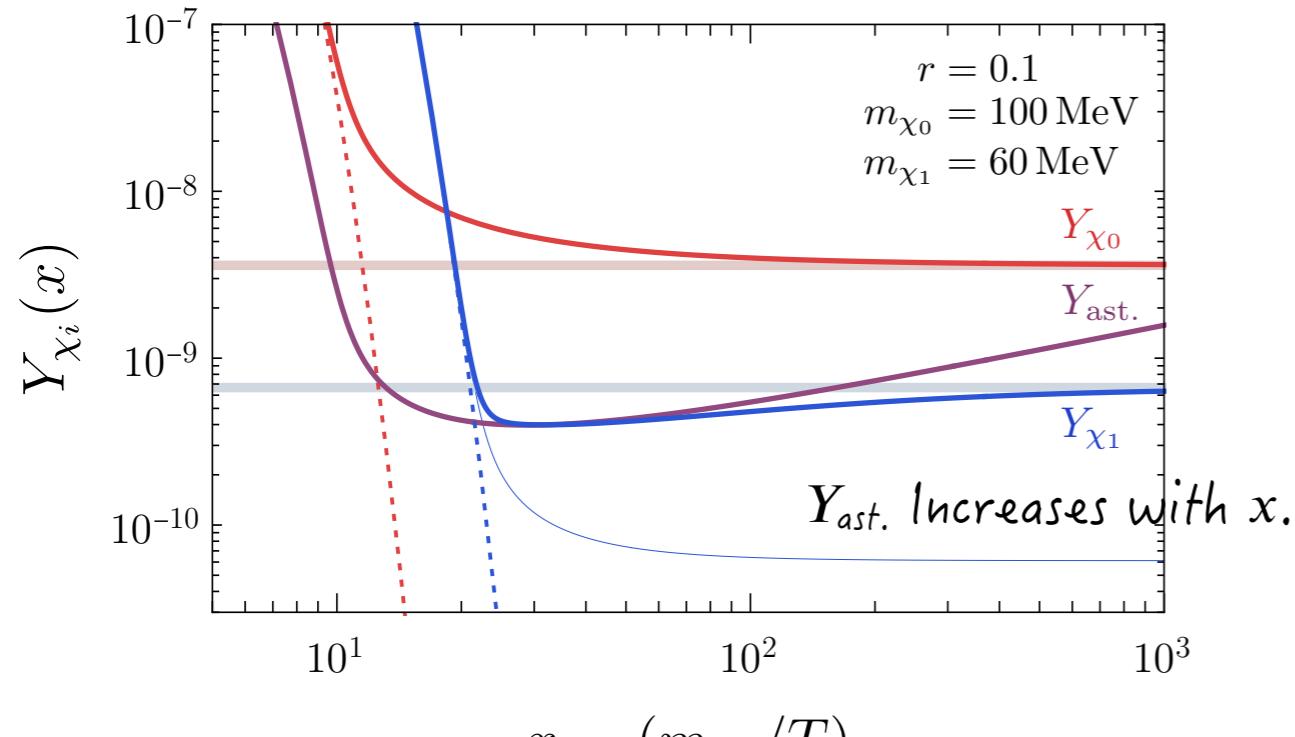
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- The annihilation cross section $\chi_1\chi_1 \rightarrow \text{SM}$ is enhanced by $1/r_1^2$.

Structure of $\chi_1\chi_1 \rightarrow \text{SM}$

When $\chi_1\chi_1 \rightarrow \text{SM}$ is dominated by p-wave Safe from
constraints?

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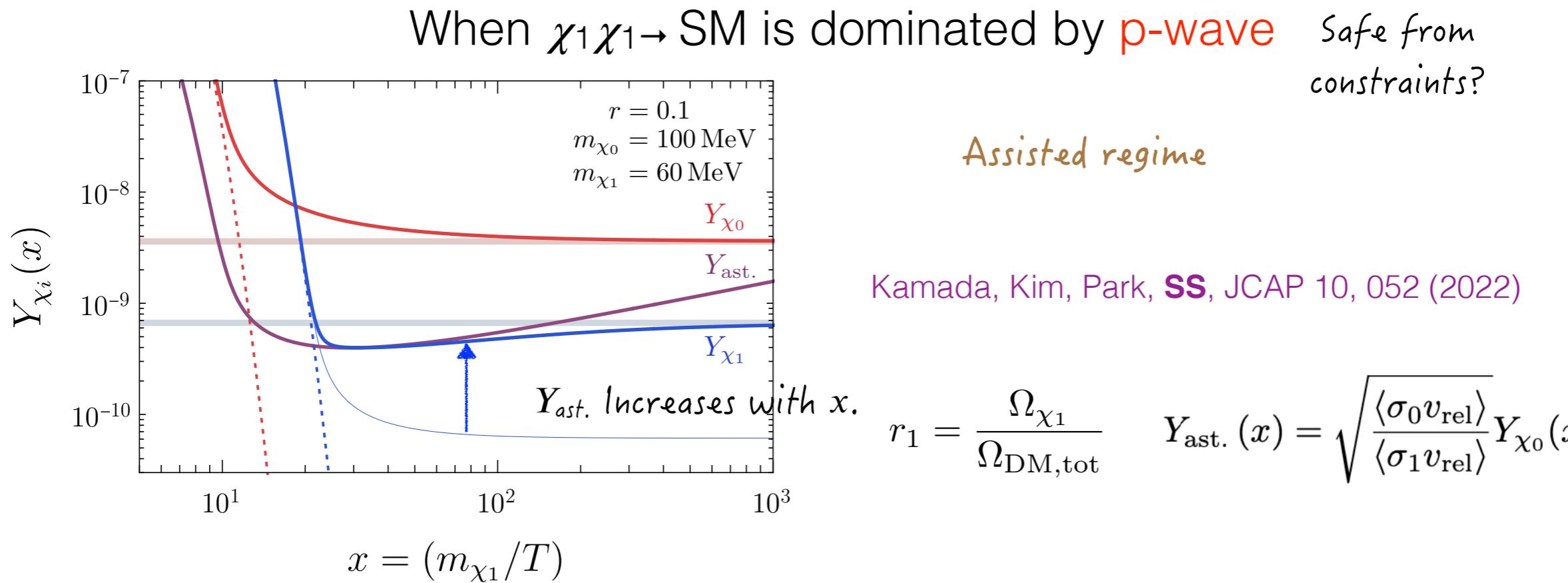
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Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

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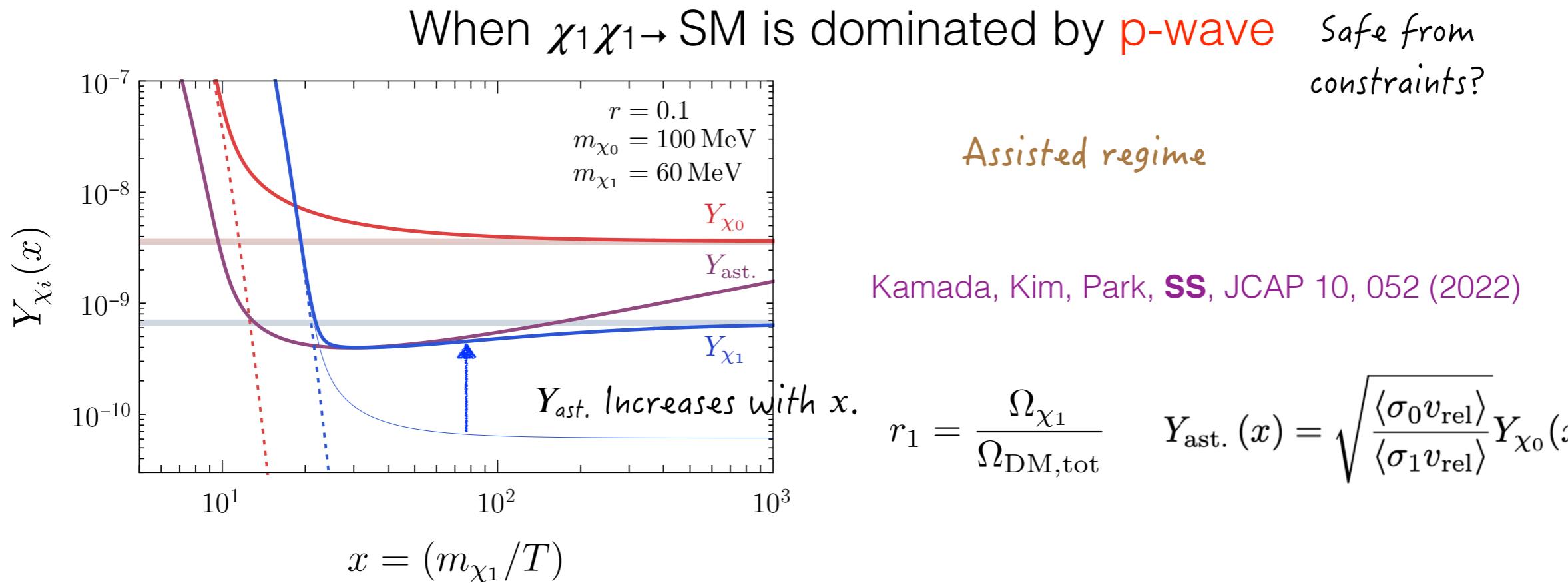
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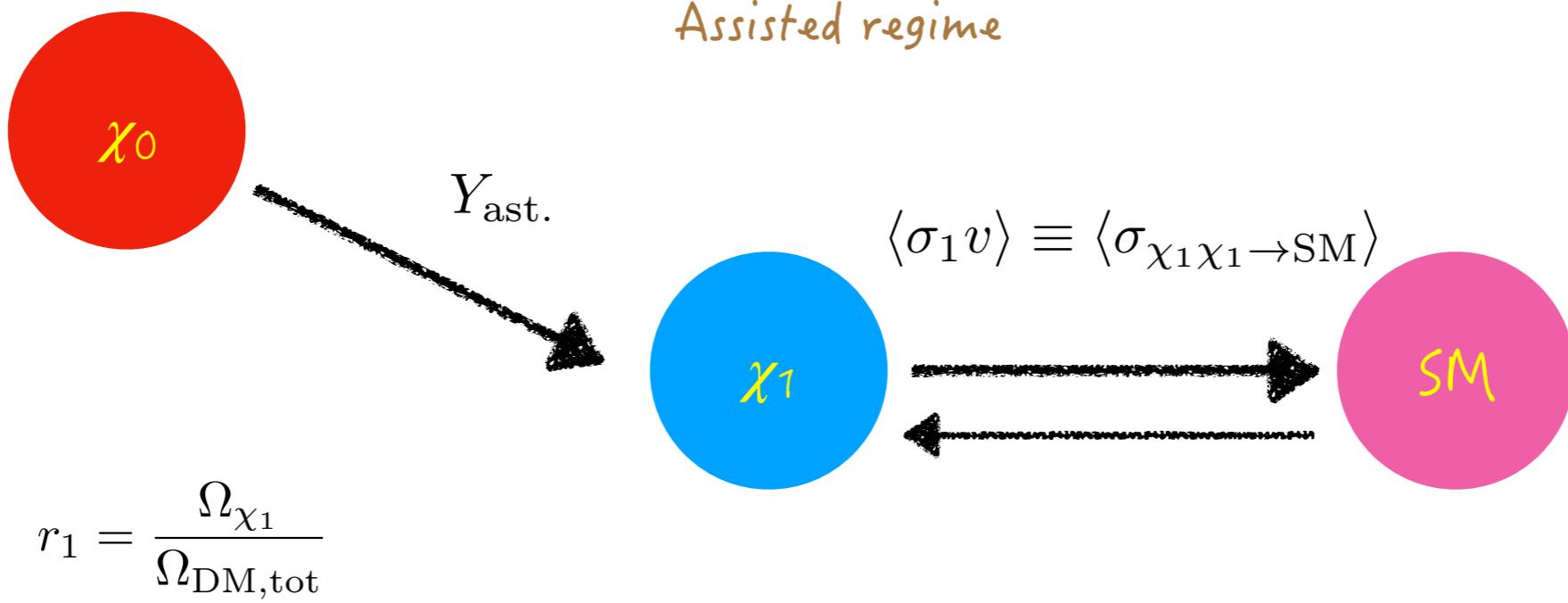
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 (the contribution by p-wave $\chi_1\chi_1 \rightarrow \text{SM}$ gets relatively suppressed.)

Structure of $\chi_1\chi_1 \rightarrow \text{SM}$



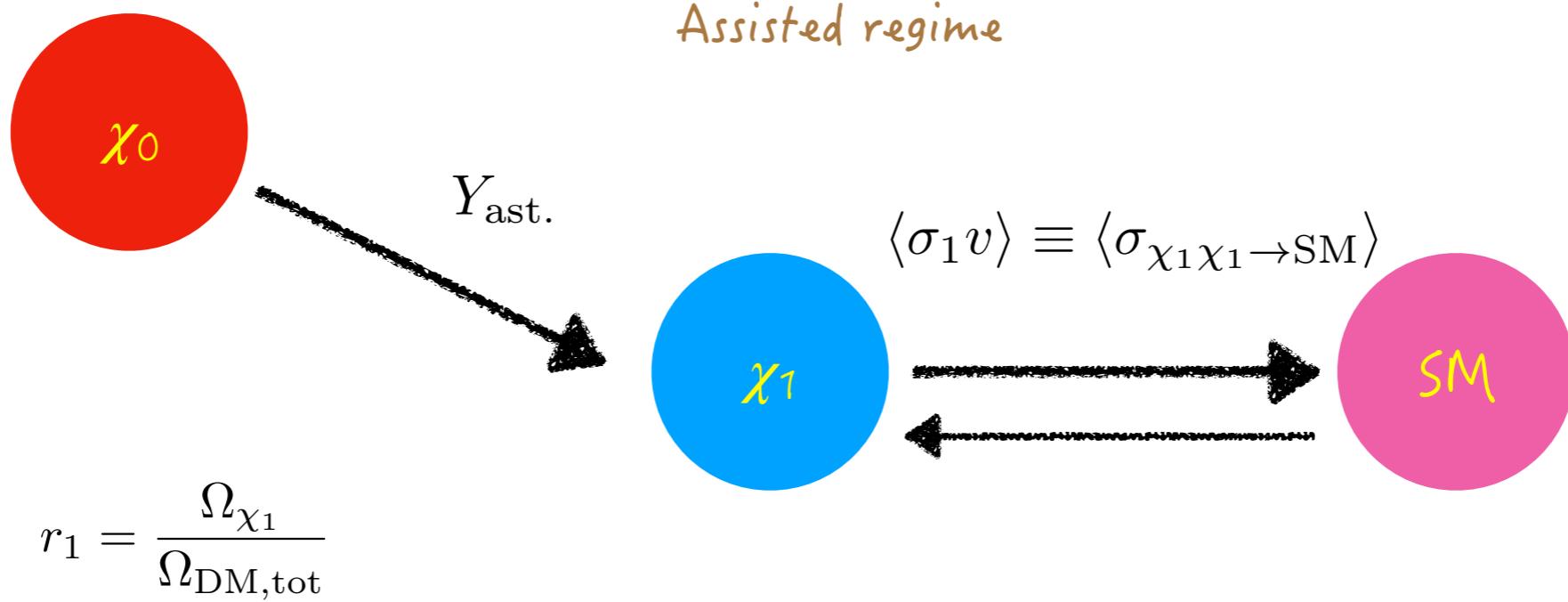
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- The annihilation cross section $\chi_1\chi_1 \rightarrow \text{SM}$ increases as $1/r_1^3$ so the process can be also sensitive to various observables.

Structure of $\chi_1\chi_1 \rightarrow \text{SM}$



- For a fixed $r_1 \ll 1$, $\chi_1\chi_1 \rightarrow \text{SM}$ should be even larger to deplete the contribution by the residual annihilation $\chi_0\chi_0 \rightarrow \chi_1\chi_1$ ($Y_{\text{ast.}}$).
- We find $\langle \sigma_1 v \rangle \propto 1/r_1^2, 1/r_1^3$ for s-wave and p-wave, respectively.

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 - We find $\langle\sigma_1 v\rangle \propto 1/r_1^2, 1/r_1^3$ for s-wave and p-wave, respectively.
- observables $\propto n_{\chi_1}^2 \langle\sigma_1 v\rangle \rightarrow$ No r_1 suppression!
(even enhanced)

Effects of χ_1 to various observables

Sub-component DM can be **not hidden** and $\chi_1\chi_1 \rightarrow \text{SM}$ affect

- Big Bang Nucleosynthesis: photo-dissociation of light elements
e.g., $e\gamma_b \rightarrow e'\gamma'$ changes the ratio of D, ${}^3\text{He}$, ${}^4\text{He}$, ..
- Cosmic microwave background: N_{eff} if χ_1 freeze-out at $T \lesssim T_{\nu,\text{dec}}$,
Energy injection by $\chi_1\chi_1 \rightarrow \text{SM}$ at the recombination epoch
- Diffuse X-rays and γ -rays in the Milky Way
- Direct detection if the crossing symmetry is effective.
(depending on the model)

Effects of χ_1 to various observables

Unprecedented role of a sub-dominant DM component

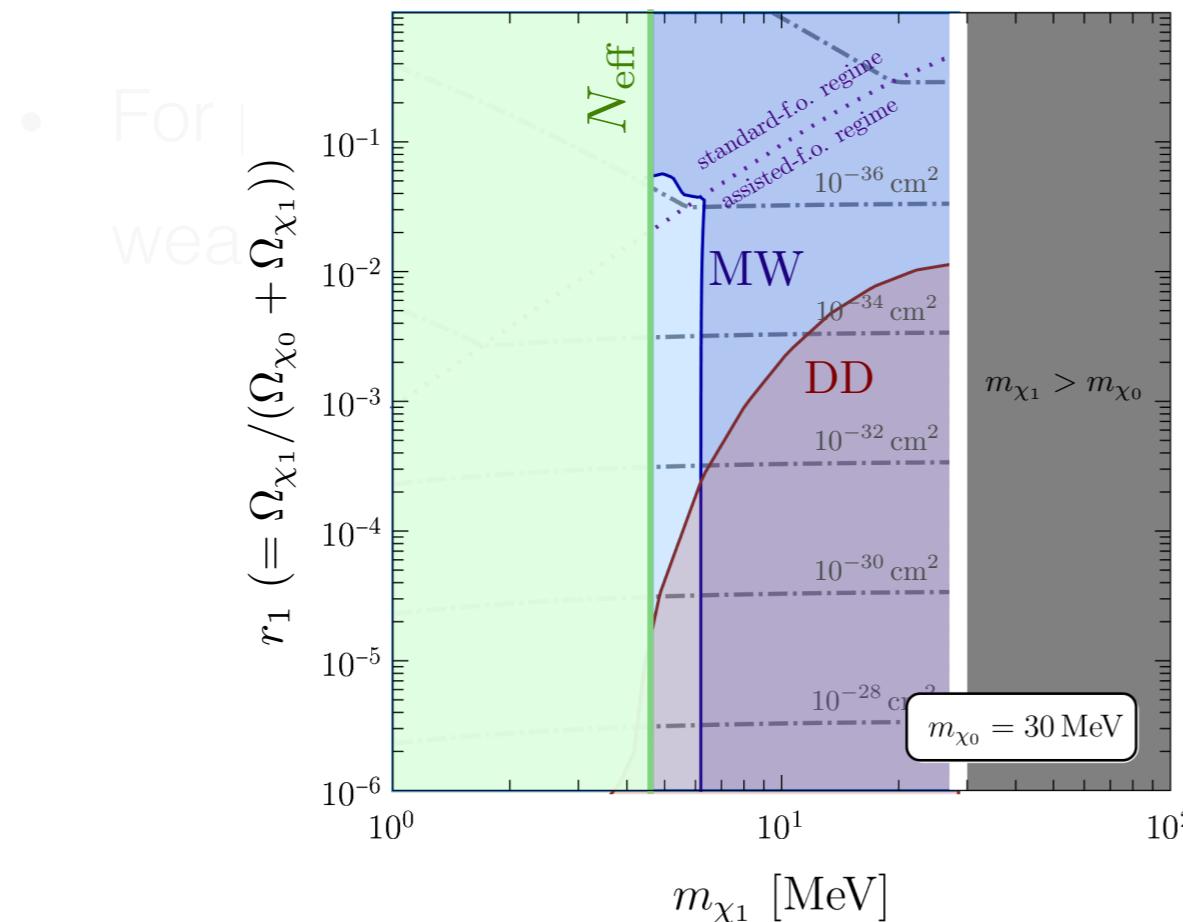
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CMB bound (sky blue): disfavor the whole parameter space

Galactic diffuse X/γ -ray (deep blue)

N_{eff} (green): almost independent of r_1

Direct detection bound (brown): XENON10, 100, DarkSide-50

Effects of χ_1 to various observables

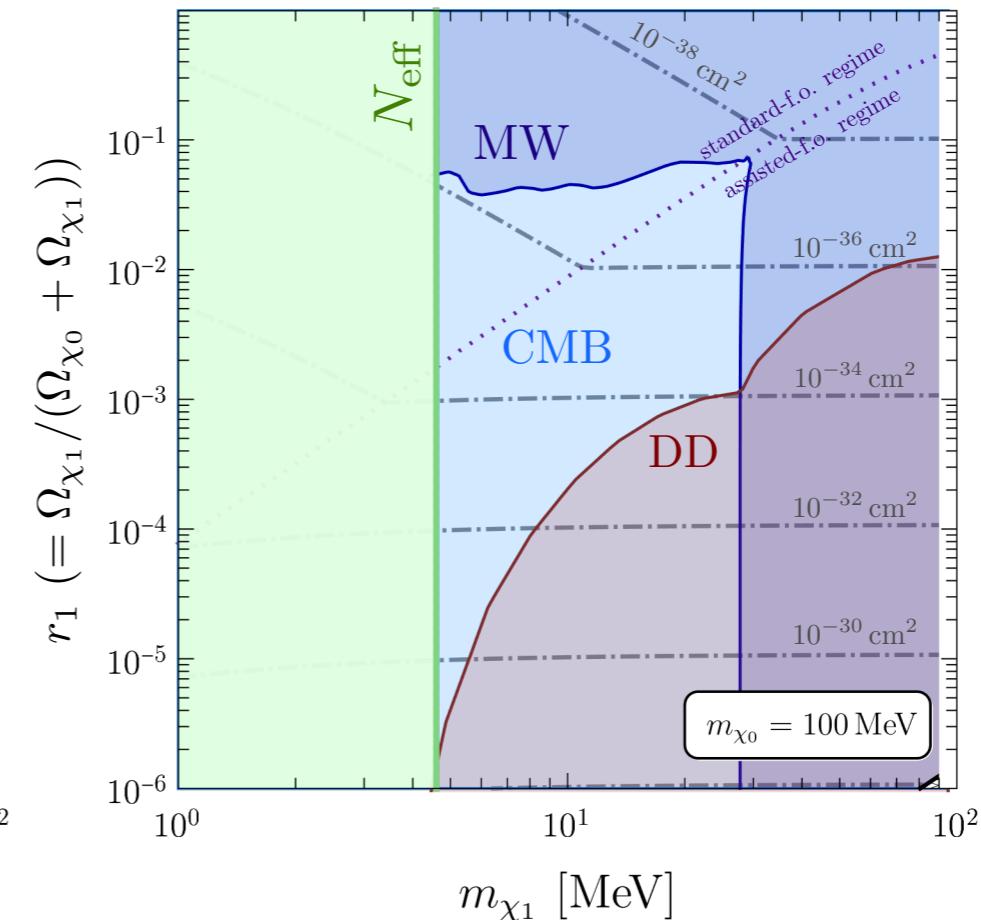
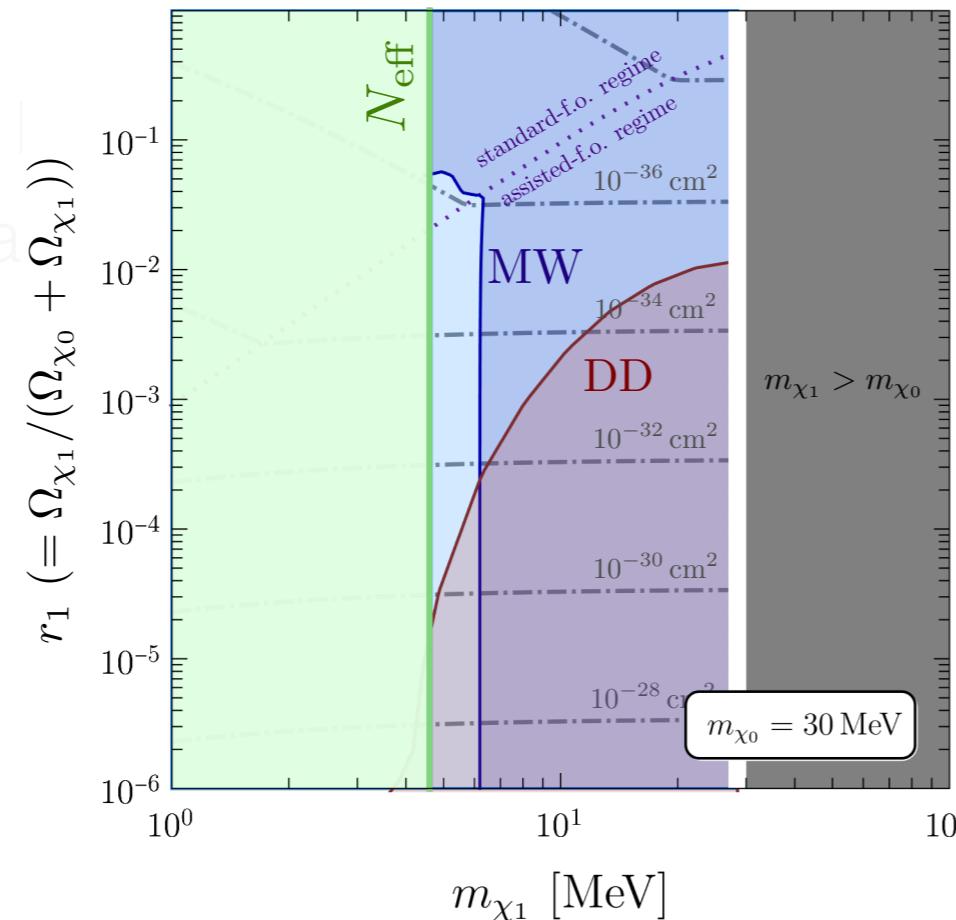
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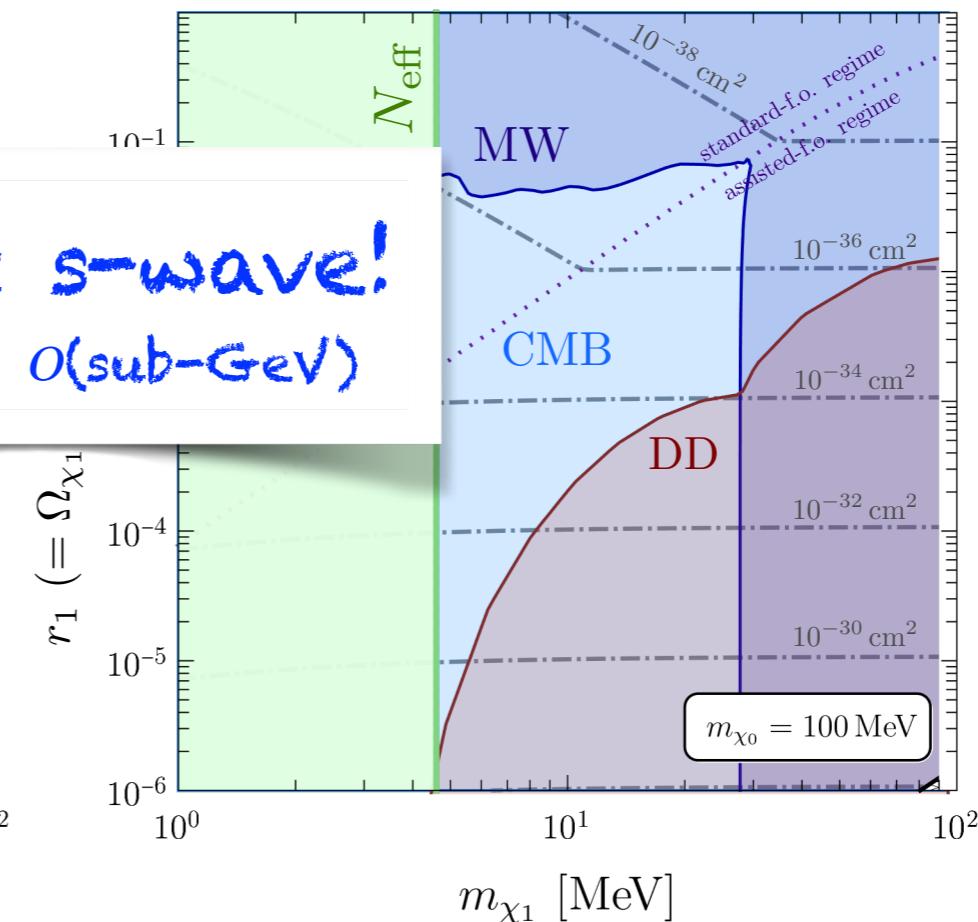
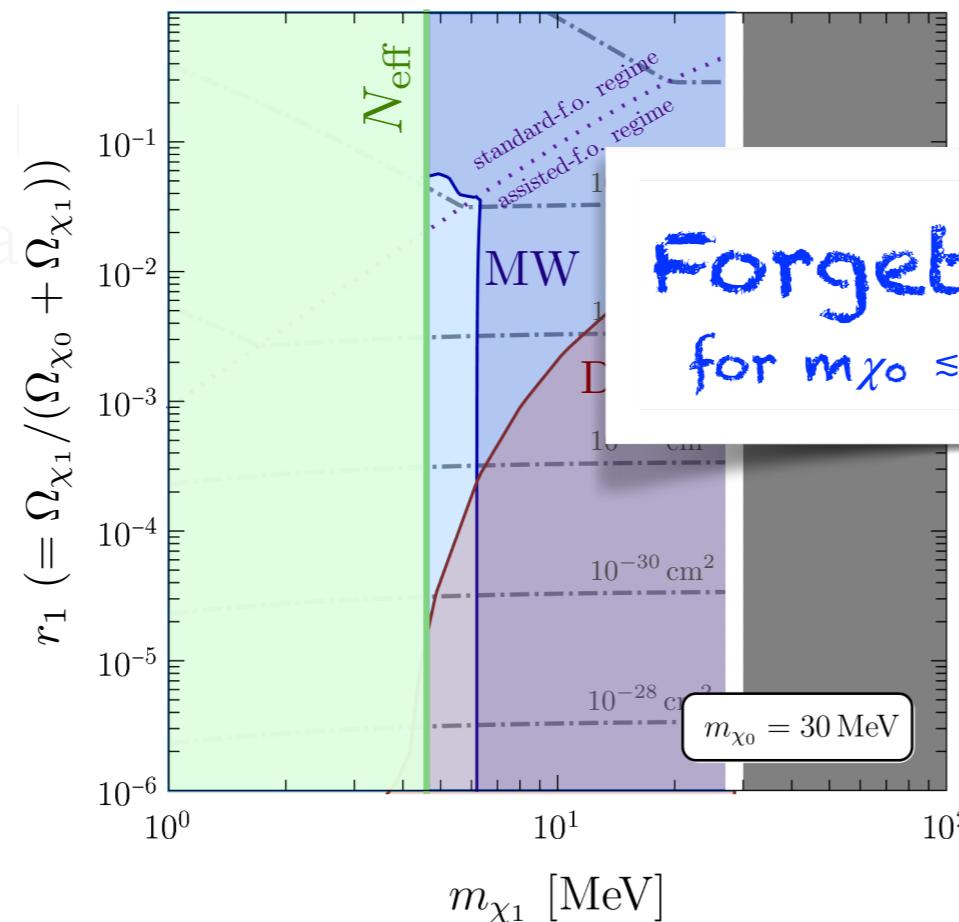
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$\chi_1 - \chi_1$
self-interaction

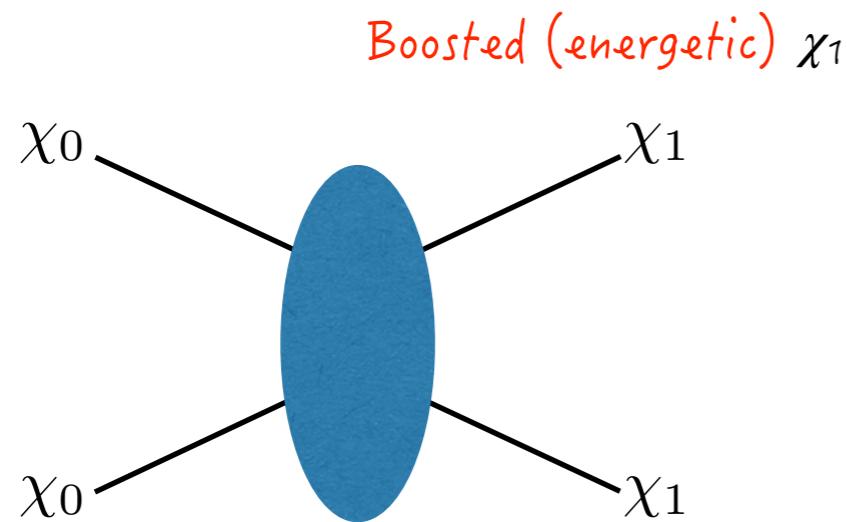
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Self-heating of χ_1

- Self-interacting DM models have been proposed actively recently.
- Self-interactions always exist. The question is how efficient they can transfer energy long after the freeze-out (not effective for WIMP).
- Self-interaction of a subdominant DM χ_1 can be large for the $O(1)$ dark sector coupling.

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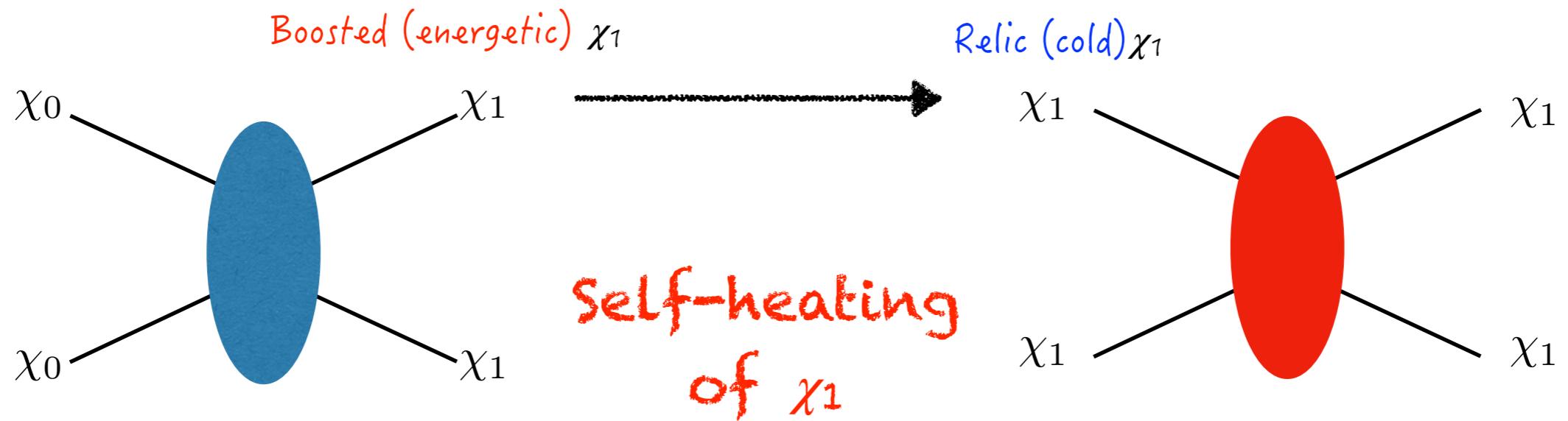
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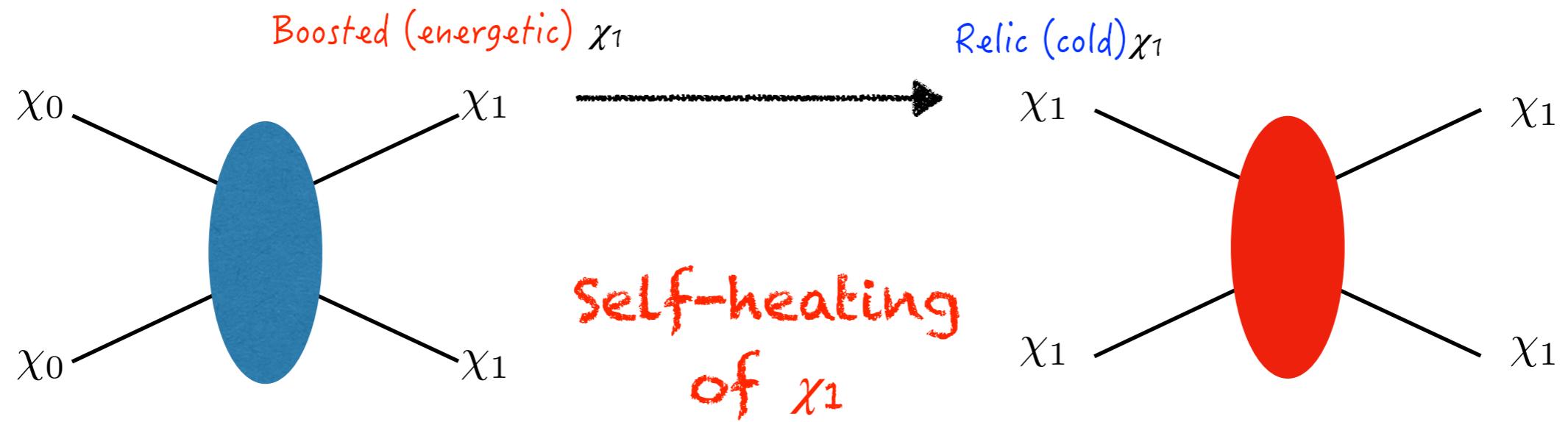
Kamada, Kim, Park, **ss**, JCAP 10, 052 (2022)



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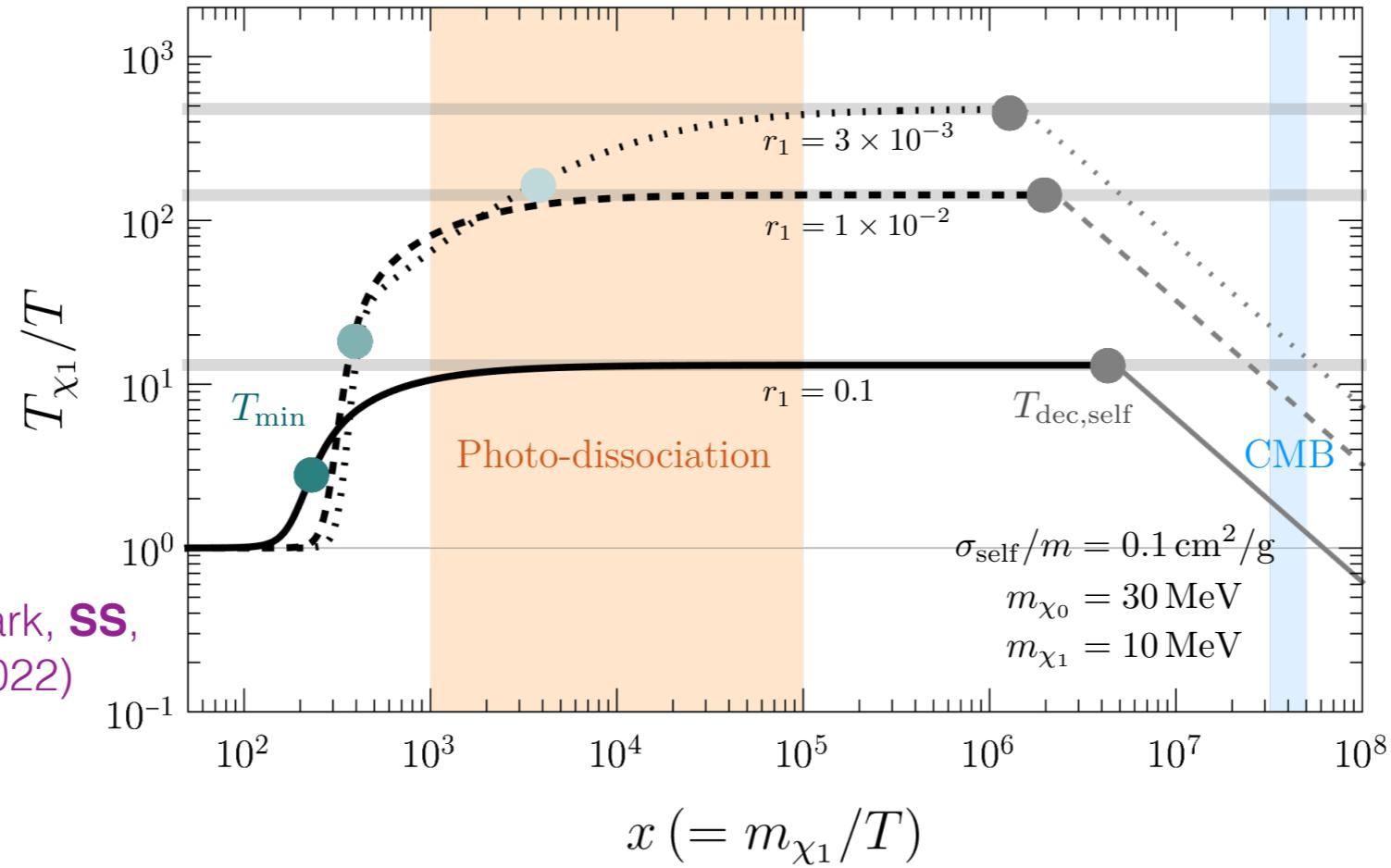
Kamada, Kim, Kim,
Sekiguchi, PRL 2018

Chu, Garcia-Cely, JCAP 2018

Vogelsberger, Zavala,
Schutz, Slatyer, MNRAS 2018

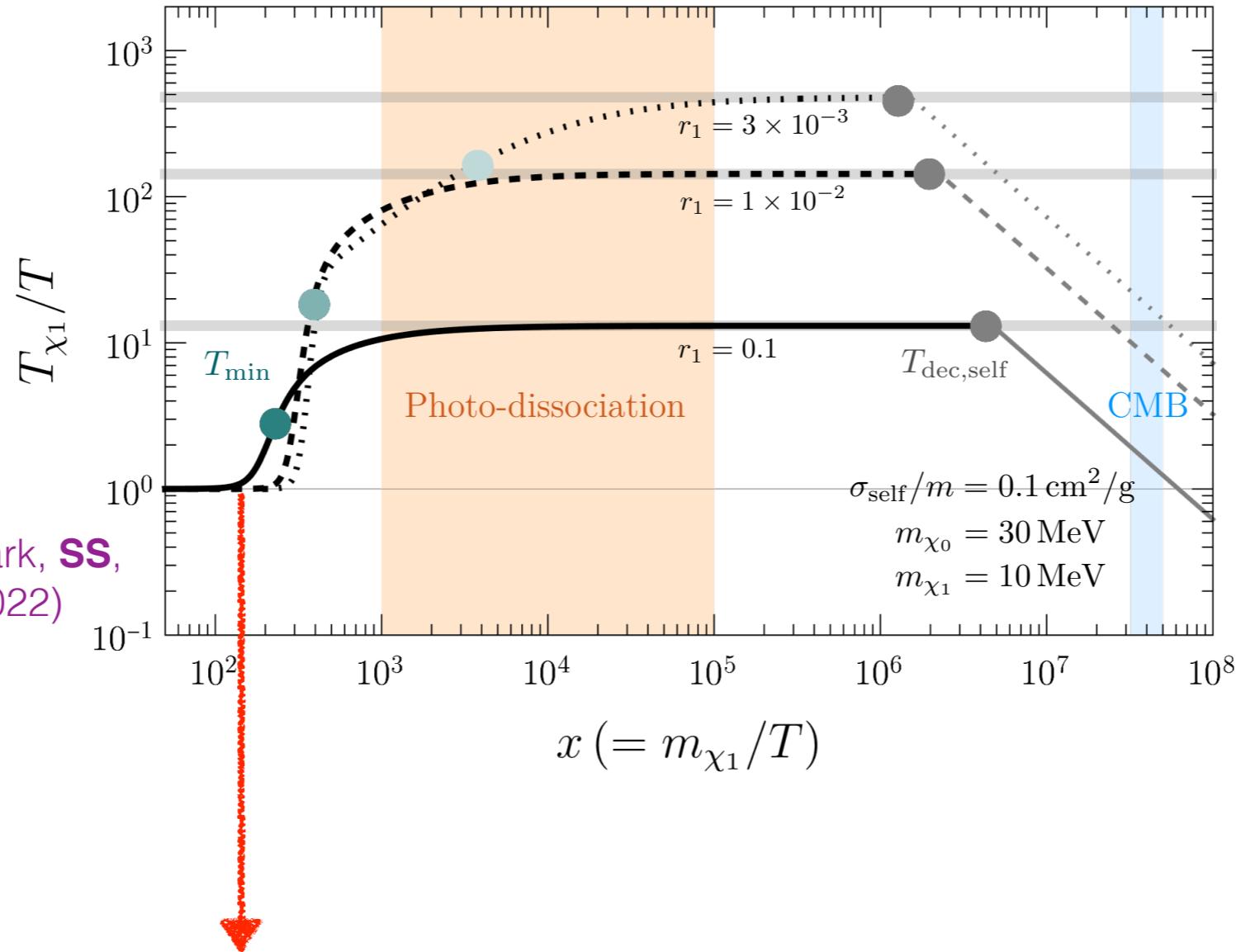
Temperature evolution of χ_1

Kamada, Kim, Park, **SS**,
JCAP 10, 052 (2022)



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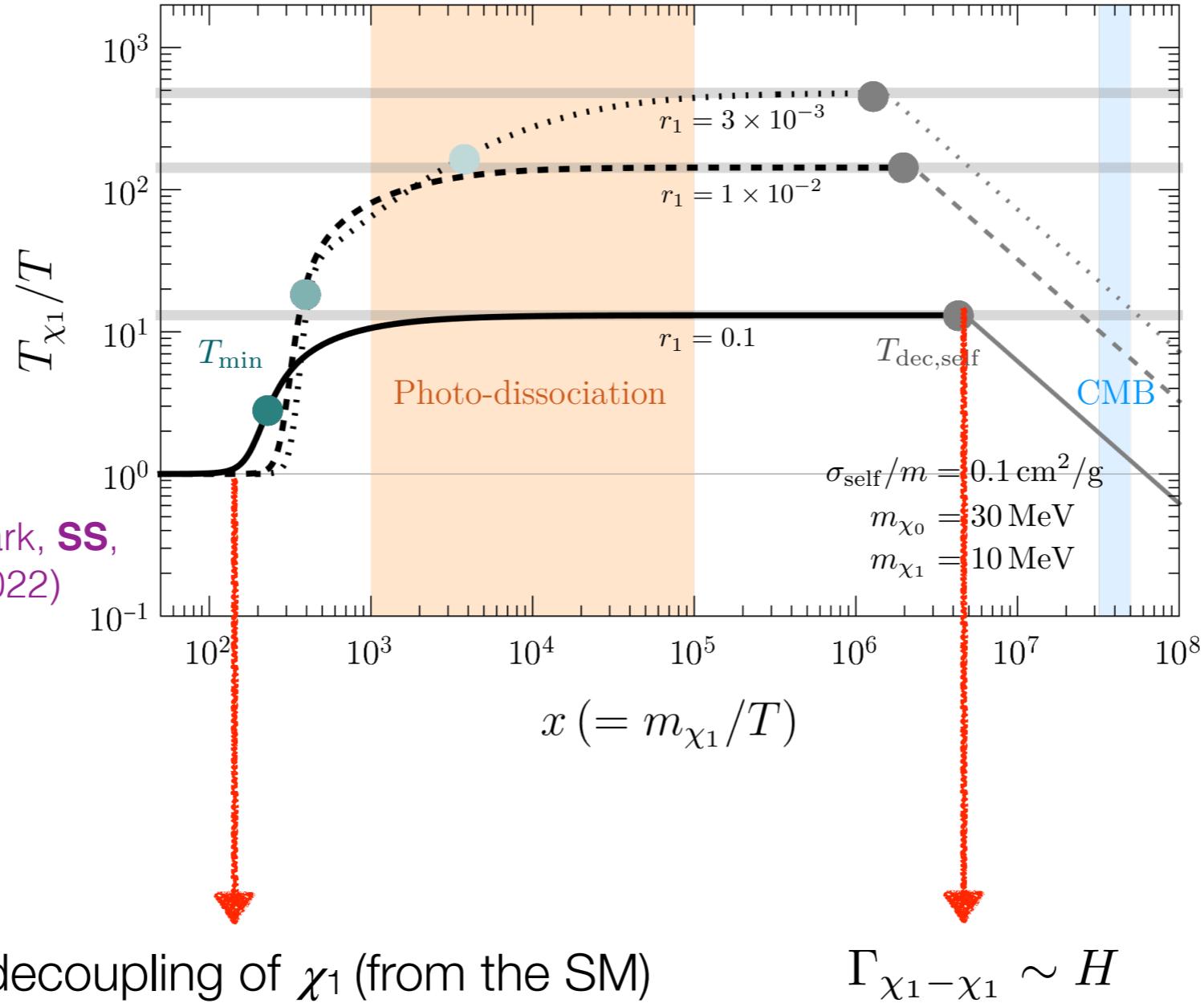
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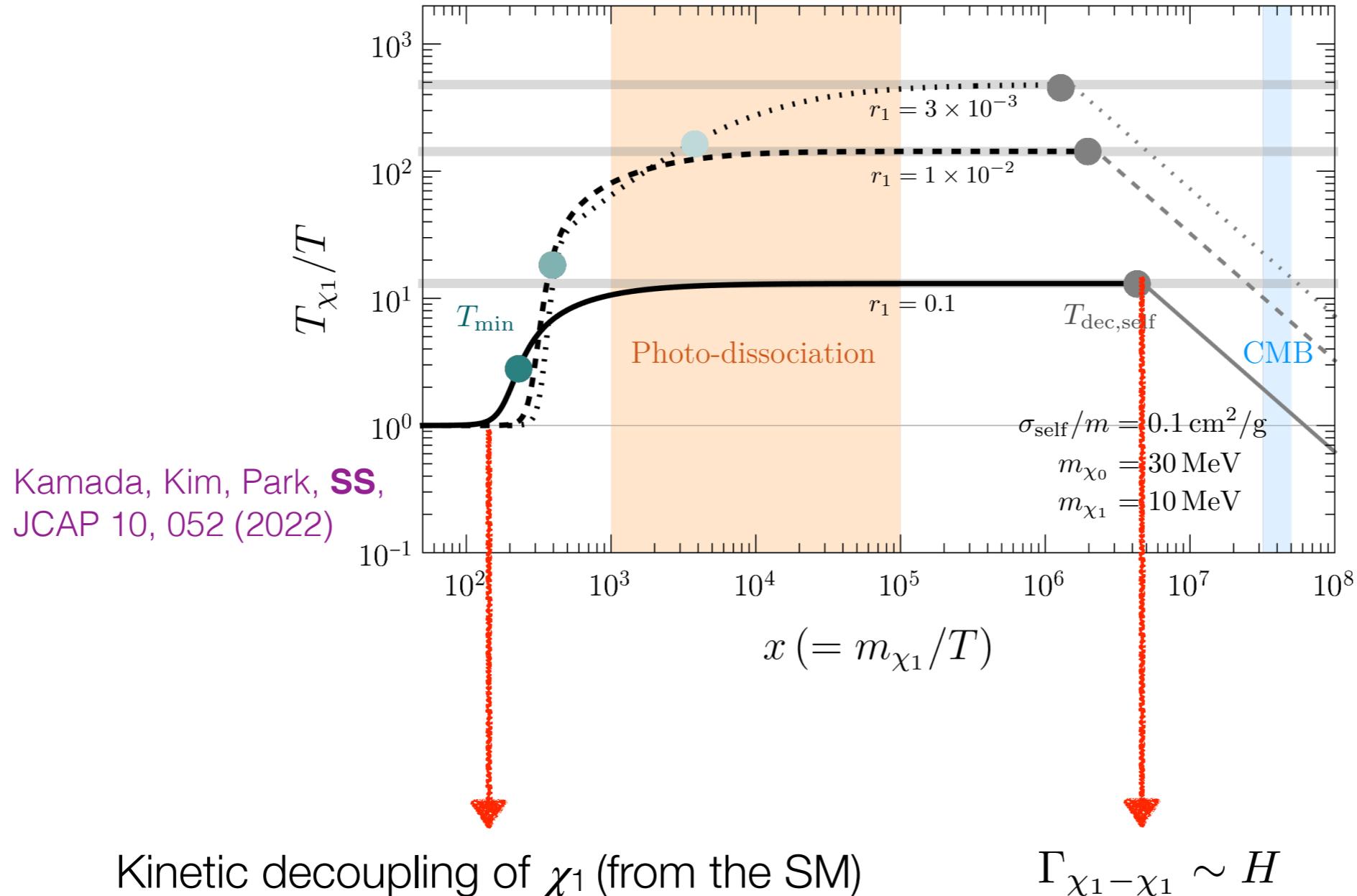
Kinetic decoupling of χ_1 (from the SM)

Temperature evolution of χ_1

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JCAP 10, 052 (2022)



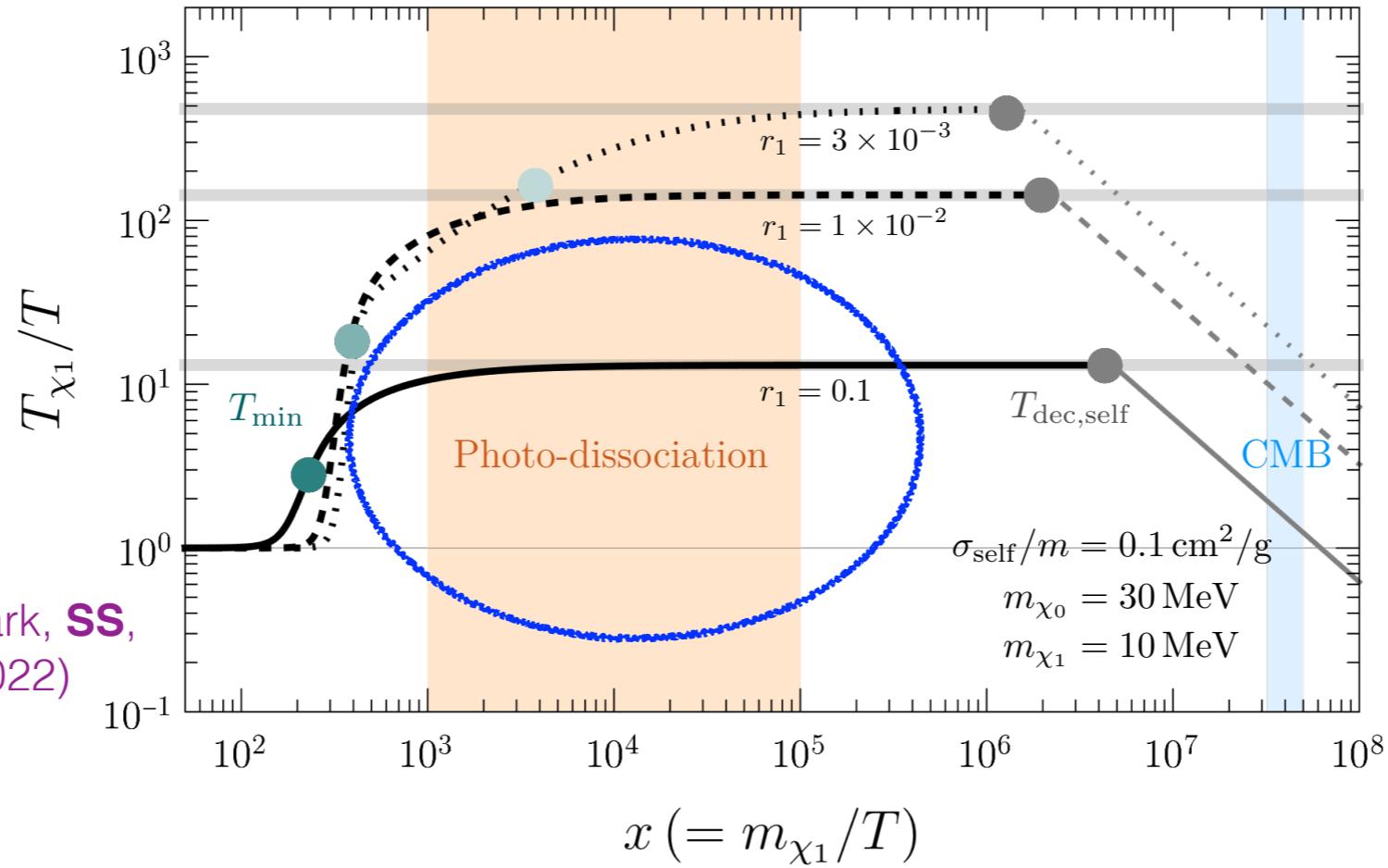
Temperature evolution of χ_1



- If self-heating is efficient even after the kinetic decoupling, the temperature evolution of χ_1 makes it behave like a radiation.

Temperature evolution of χ_1

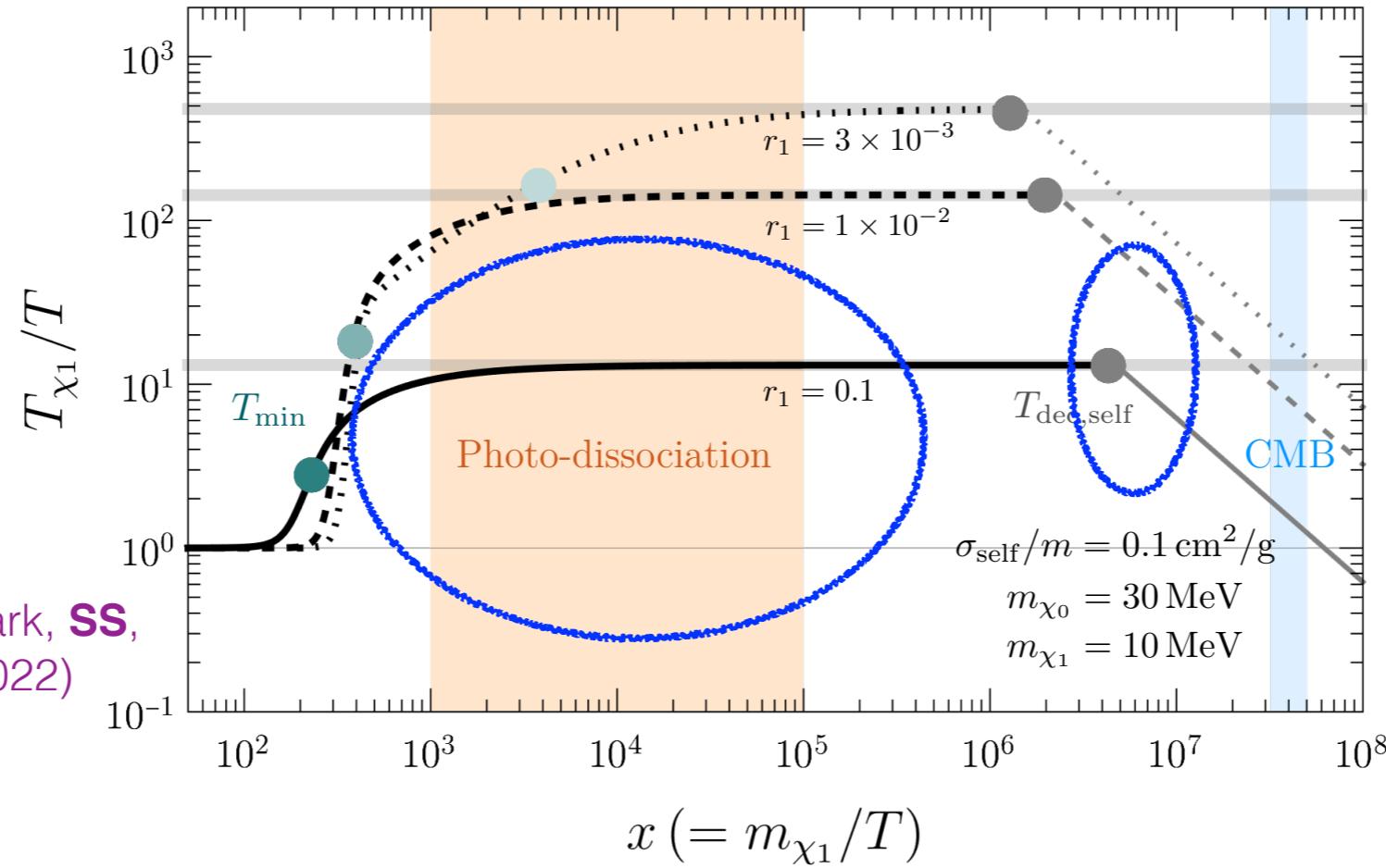
Kamada, Kim, Park, **SS**,
JCAP 10, 052 (2022)



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Temperature evolution of χ_1

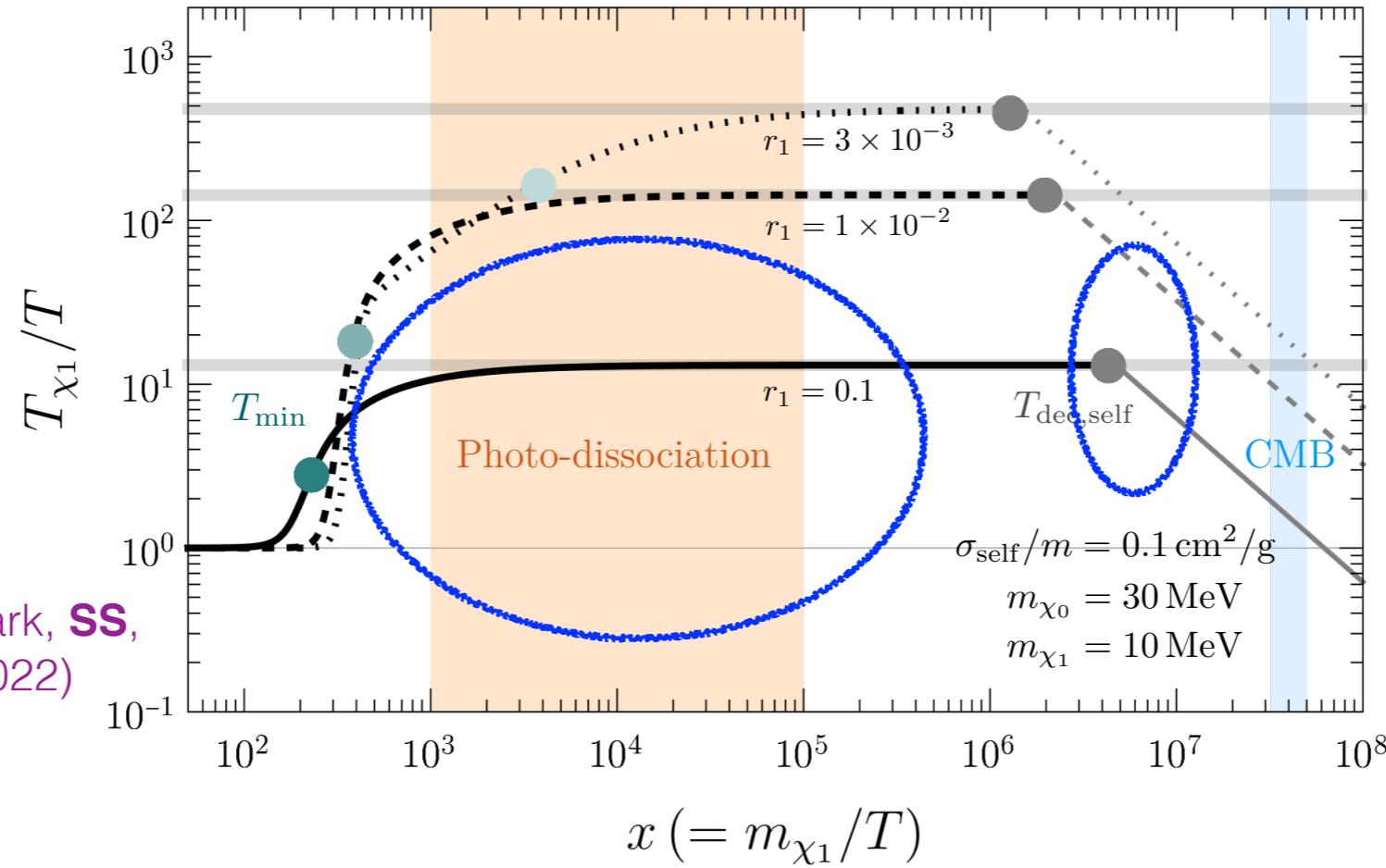
Kamada, Kim, Park, **SS**,
JCAP 10, 052 (2022)



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Temperature evolution of χ_1

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JCAP 10, 052 (2022)



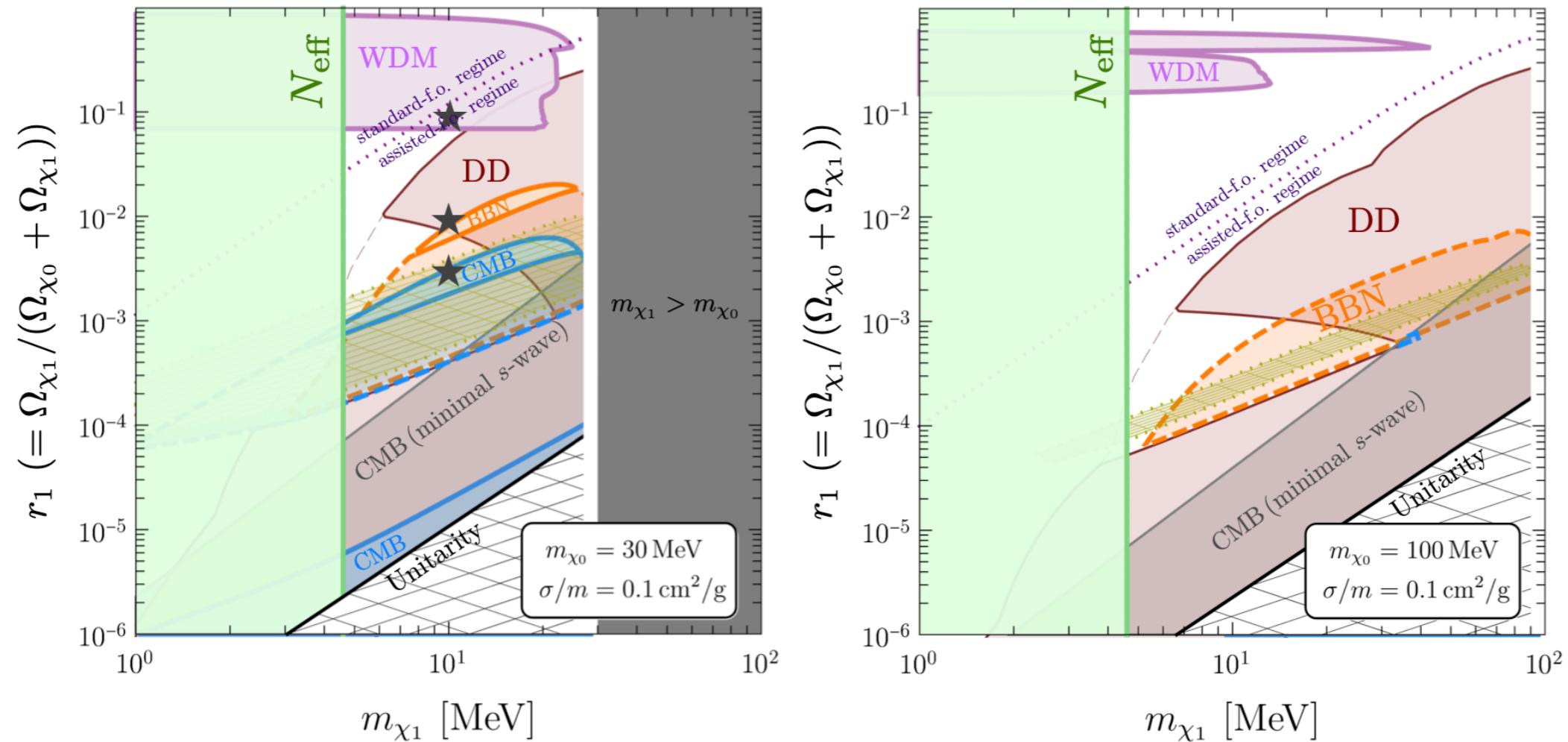
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χ_1 can be **sub-GeV Warm Dark Matter!!**

Lyman- α
of satellites

New bounds due to self-heating

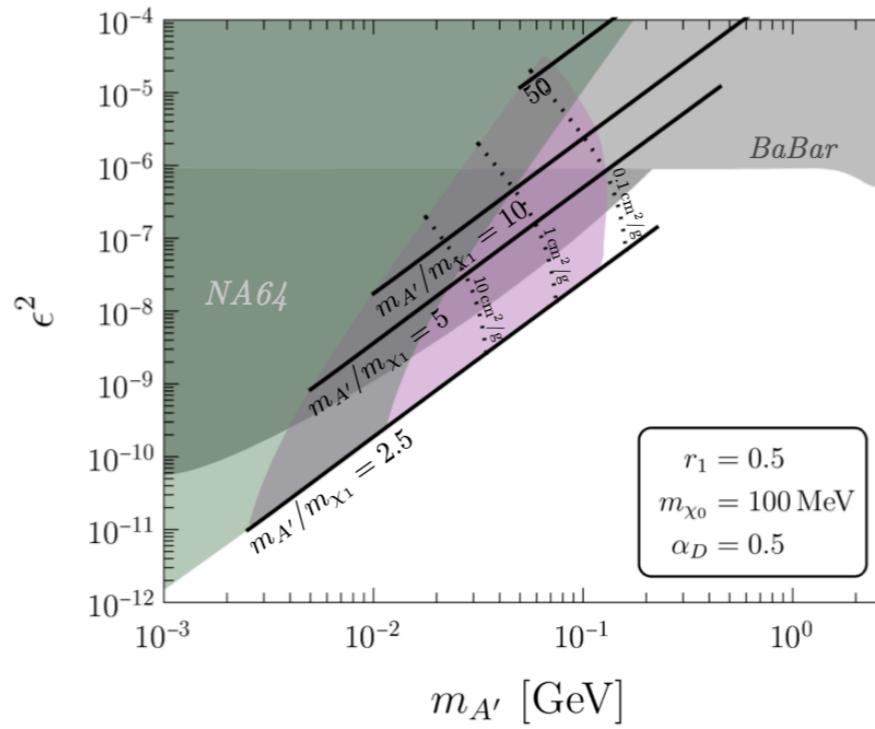
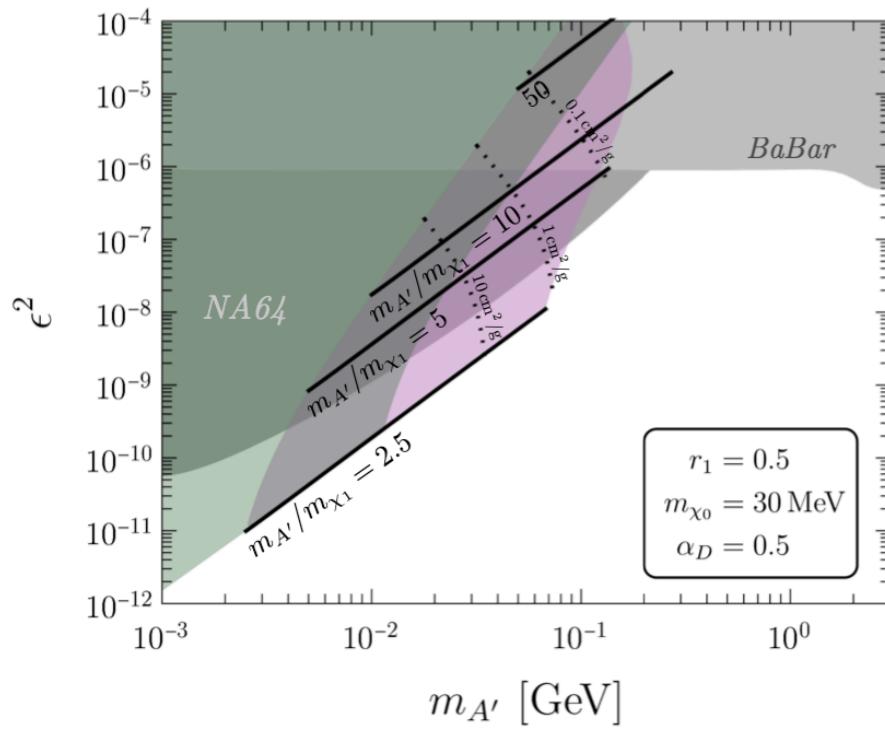
Kamada, Kim, Park, **ss**, JCAP 10, 052 (2022)



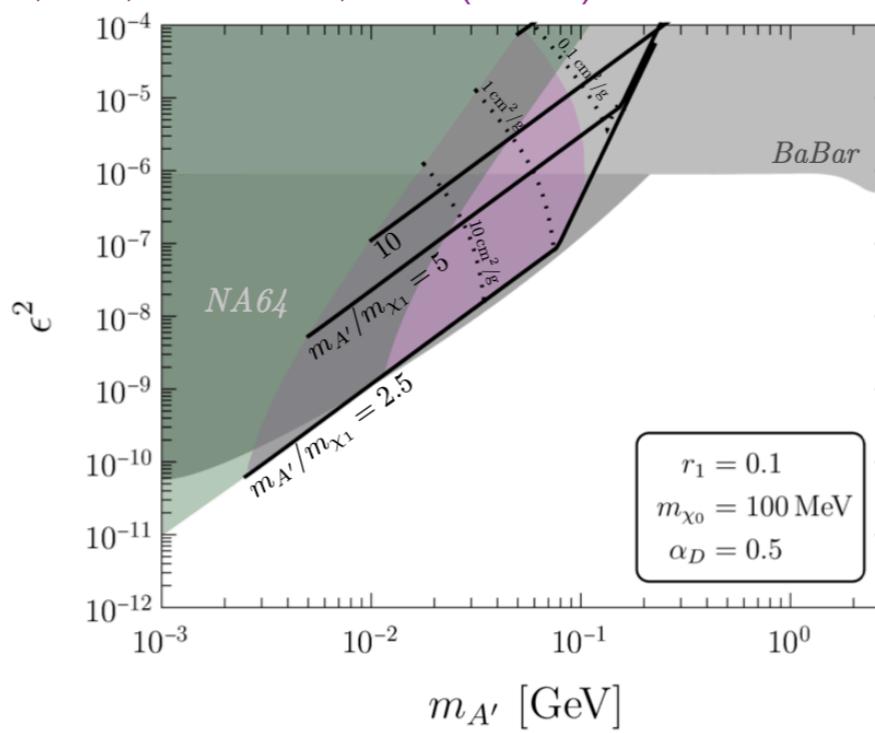
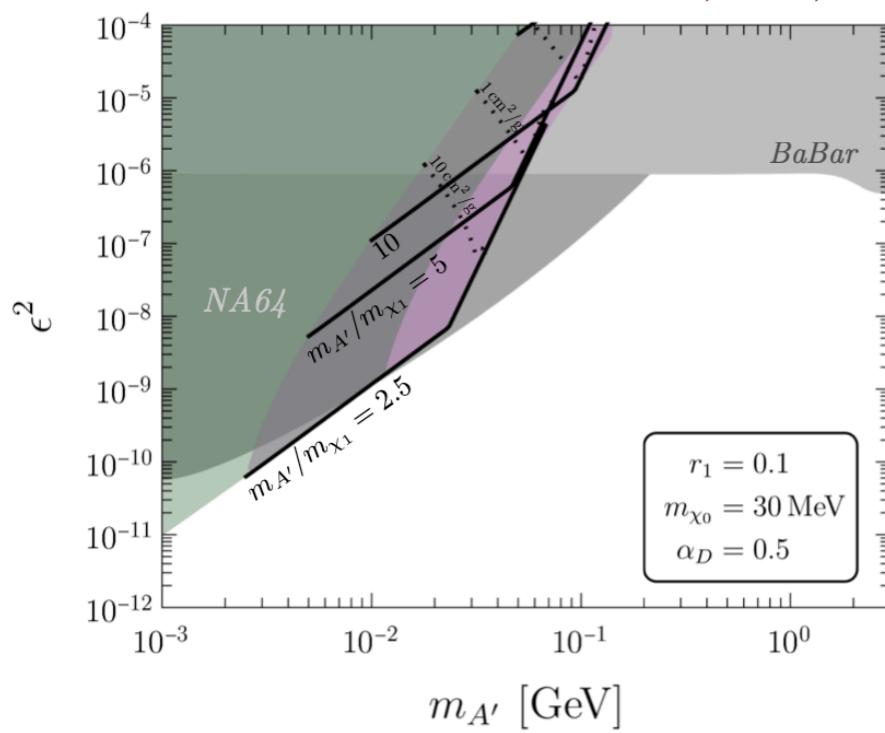
- WDM constraint enters when $r_1 \gtrsim 0.07$ even for $m_{\chi_1} \sim 40$ MeV.
- Direct detection bounds get weaken since n_{χ_1} inside our MW decreases due to the kinetic energy of χ_1
- ★: reference values of r_1 in the temperature evolution (previous slide)

Complementary searches

Light DM can be produced in accelerators with high intensities!



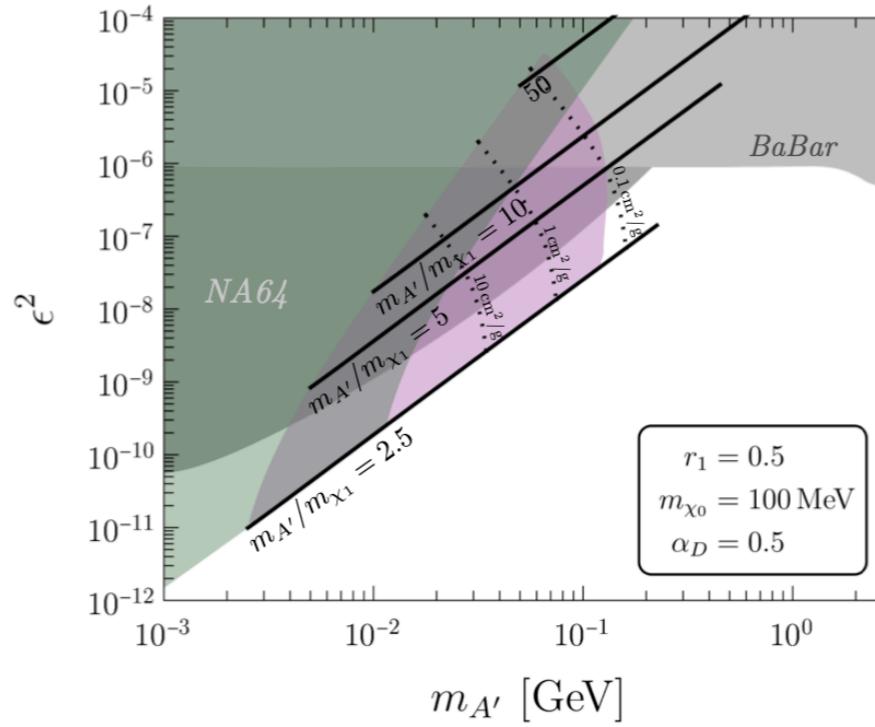
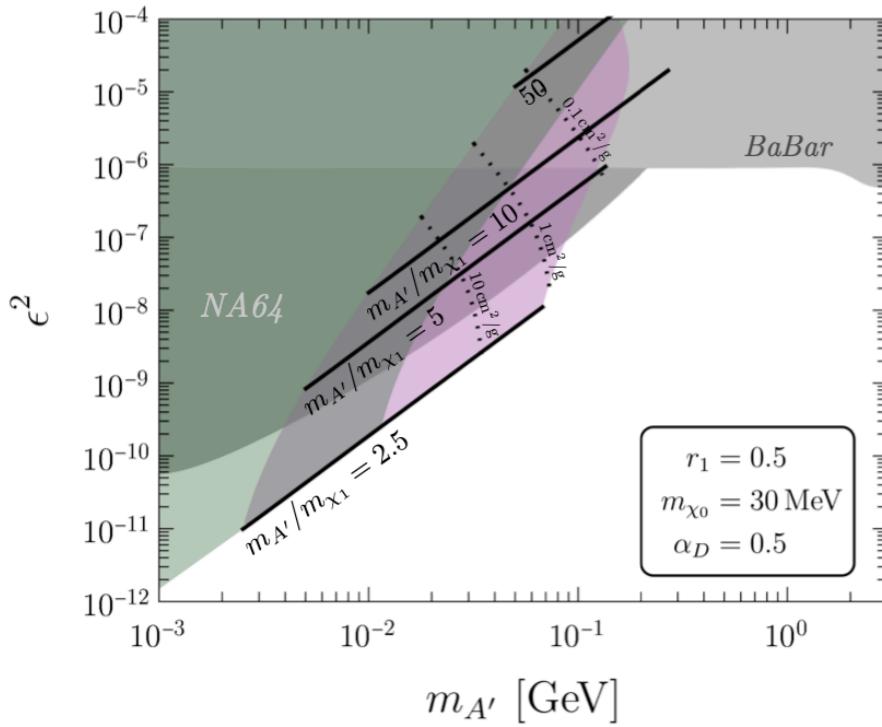
- Reference model: **singlet scalar DM + dark photon (p-wave)**
- Green: N_{eff} ,
Pink: WDM
for $r_1 \gtrsim 0.07$.



Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

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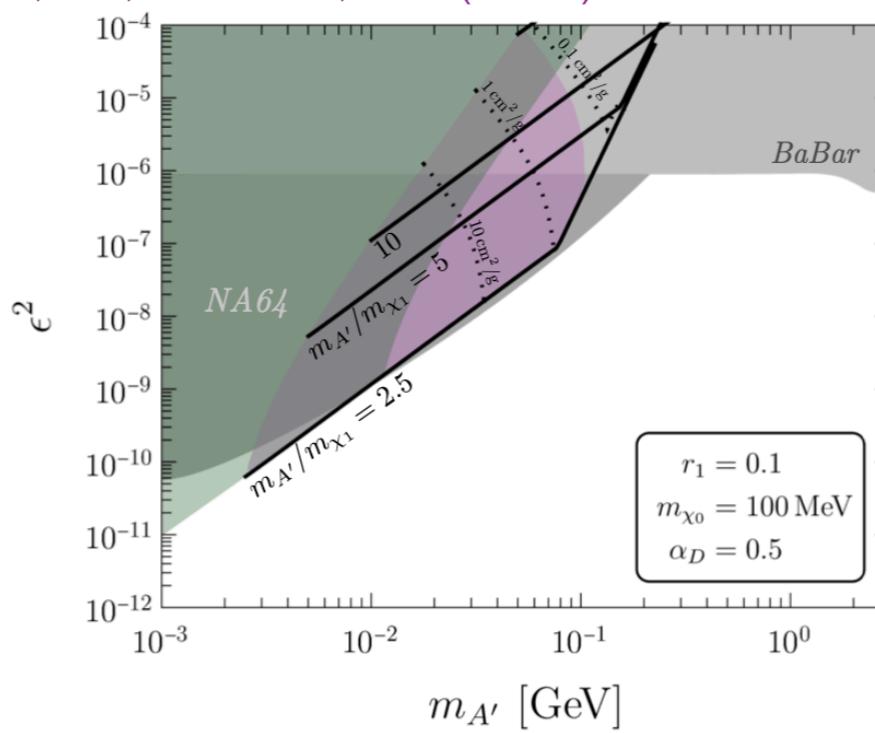
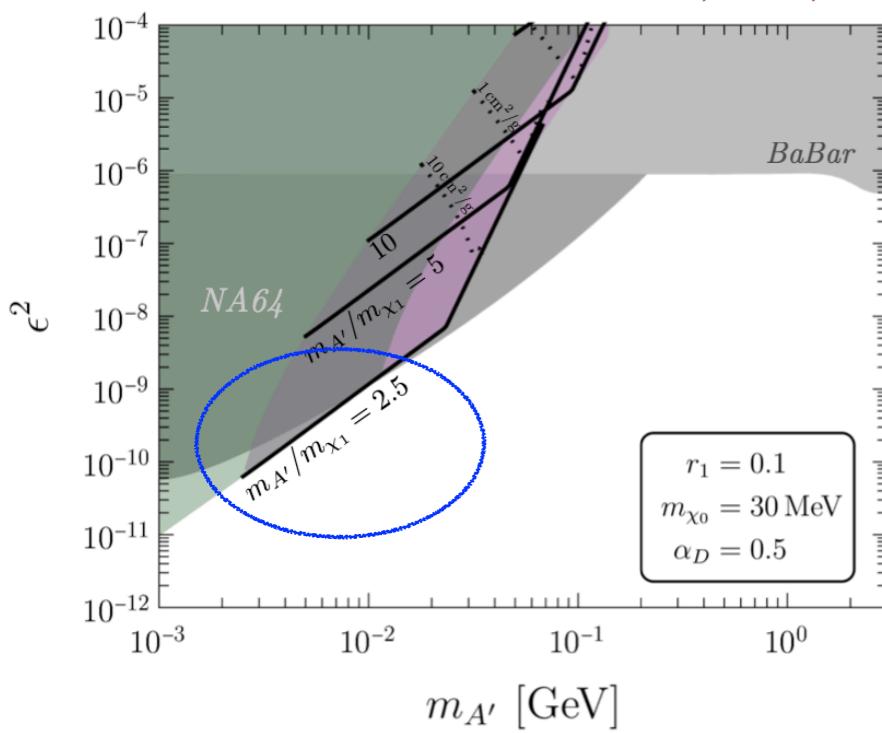
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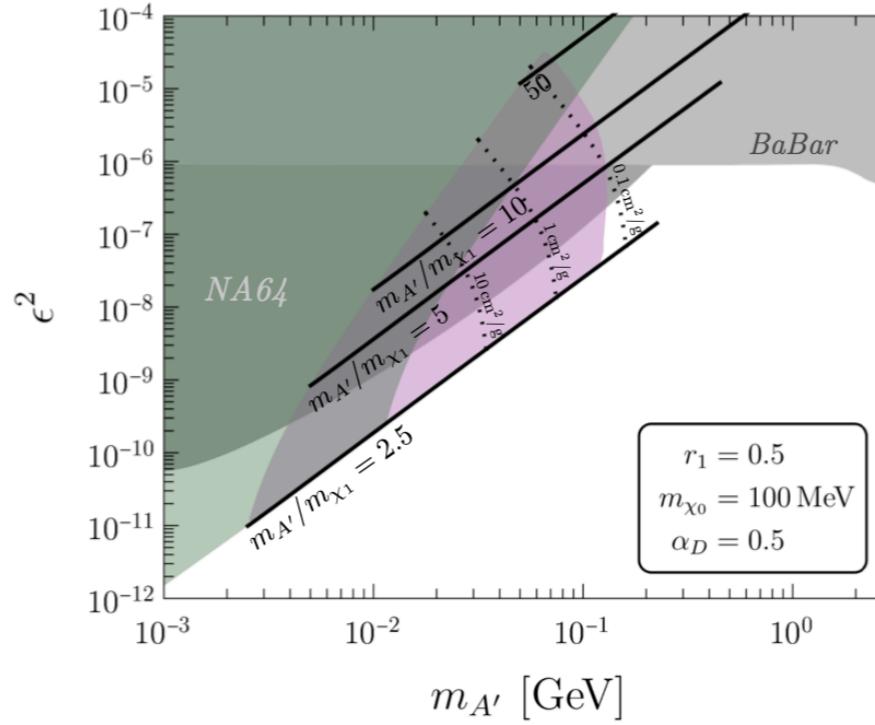
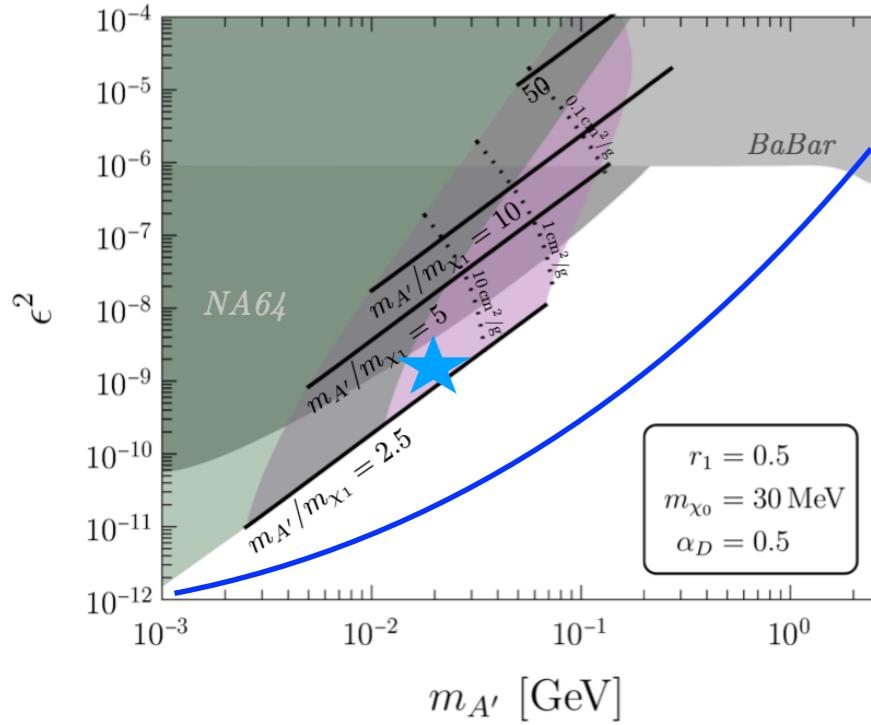
- For $r_1 \lesssim 0.07$, not preferred by the accelerator results.



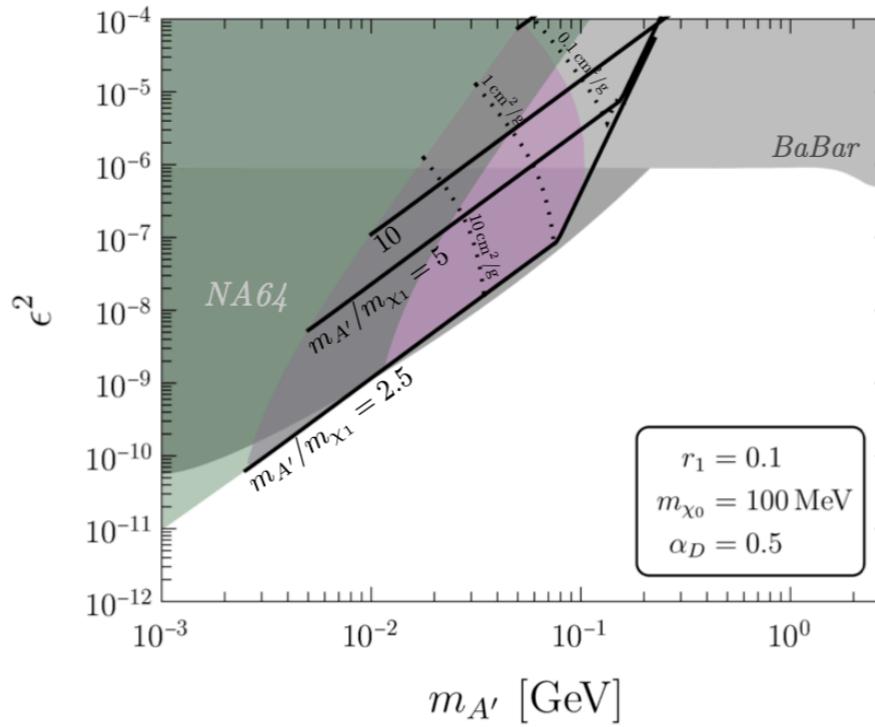
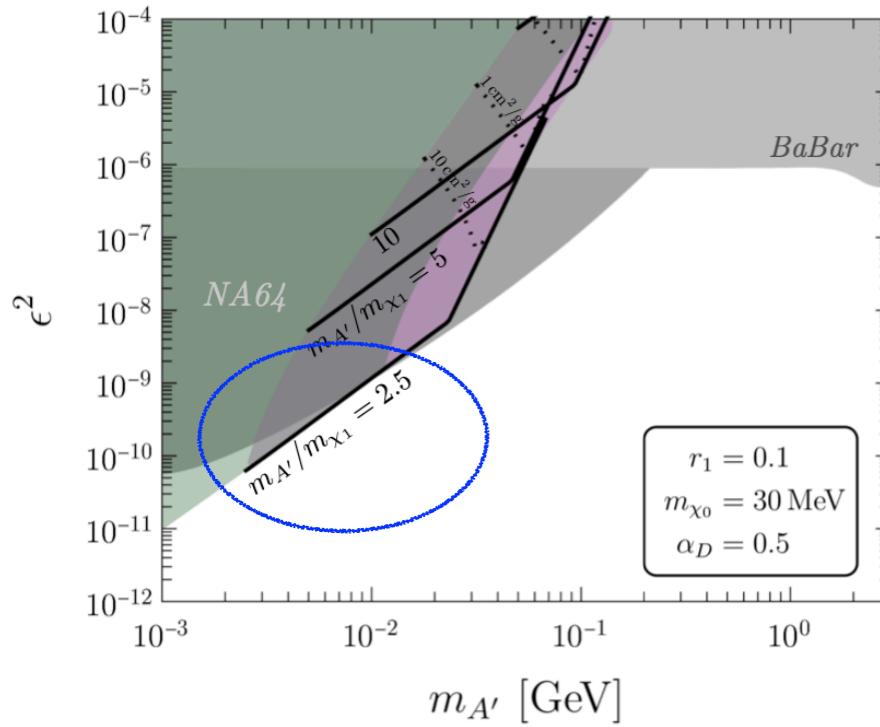
Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)

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Kamada, Kim, Park, **SS**, JCAP 10, 052 (2022)



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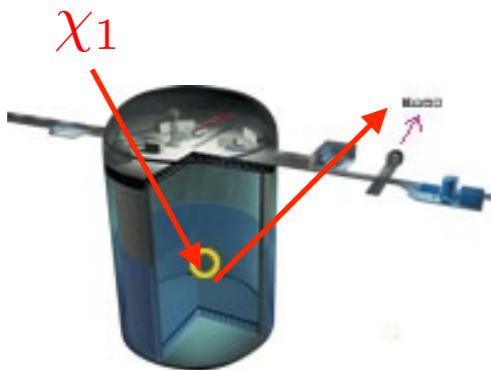
- Future discovery can tell the dark sector details.

Conclusions

- A **sub**-component DM (χ_1) can severely affect the cosmo/astro observables: p-wave χ_1 - SM is preferred but still constrained!.
(Multi-component p-wave scenarios are not always safe.)
- Self-heating naturally arises in a wide range of parameter space and changes the evolution of the temperature of χ_1 after the freeze-out.
- The temperature evolution affects the structure formation of χ_1 :
a **sub-GeV mass Warm Dark Matter** (heavy WDM) for $r_1 \gtrsim 0.07$!
→ This is true even when χ_1 is a dominant component DM.
- Complementary searches in accelerators can give hints on the dark sector details (disfavor $r_1 \lesssim 0.07$ for a reference model).

Backup

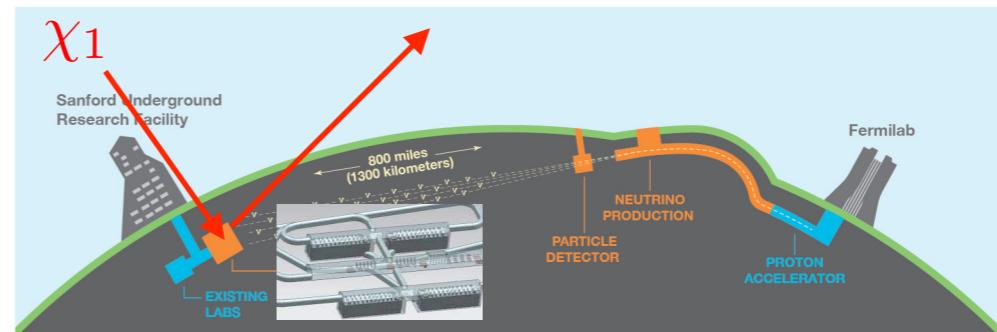
- Neutrino Experiments



PHYSICAL REVIEW LETTERS 120, 221301 (2018)

Editors' Suggestion

Search for Boosted Dark Matter Interacting with Electrons in Super-Kamiokande



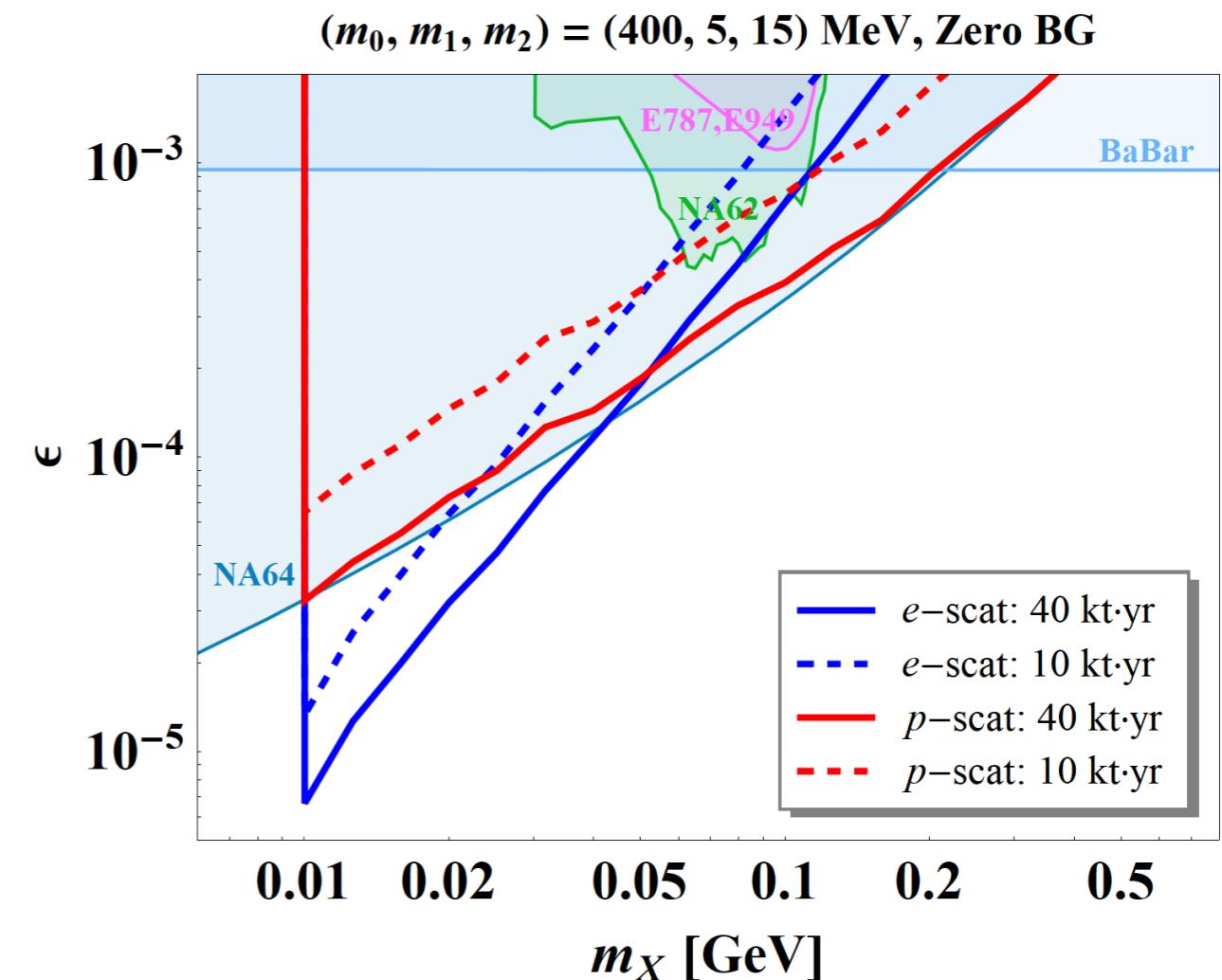
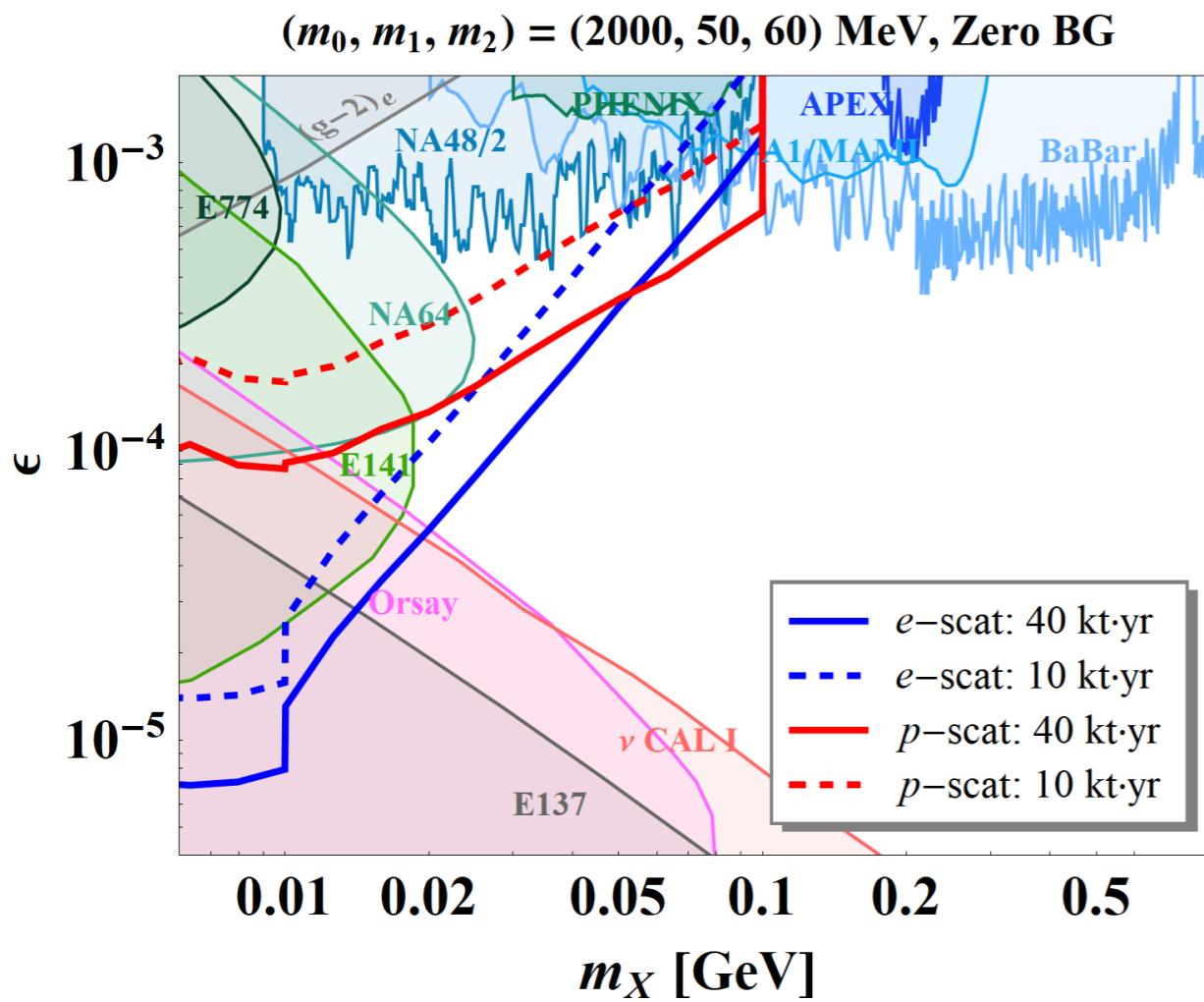
DUNE
DEEP UNDERGROUND
NEUTRINO EXPERIMENT

- 8.8 Dark Matter Probes
8.8.1 Benchmark Dark Matter Models
8.8.2 Search for Low-Mass Dark Mater at the Near Detector
8.8.3 Inelastic Boosted Dark Matter Search at the DUNE FD
8.8.4 Elastic Boosted Dark Matter from the Sun

Kim, Park, SS, PRL 2017

Backup

DUNE



Backup

- Dark Matter direct detection experiments



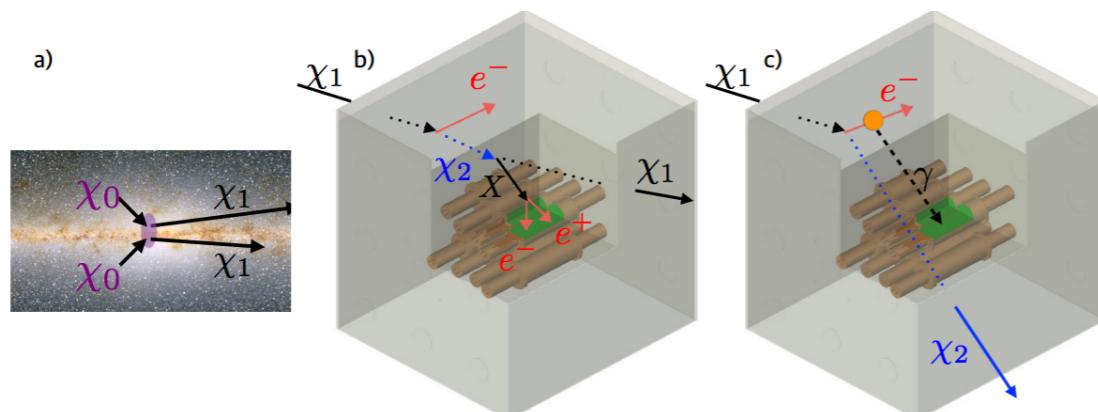
PHYSICAL REVIEW LETTERS 122, 131802 (2019)

Editors' Suggestion

First Direct Search for Inelastic Boosted Dark Matter with COSINE-100

Based on the suggestion in

Giudice, Kim, Park,
SS, PLB 2018

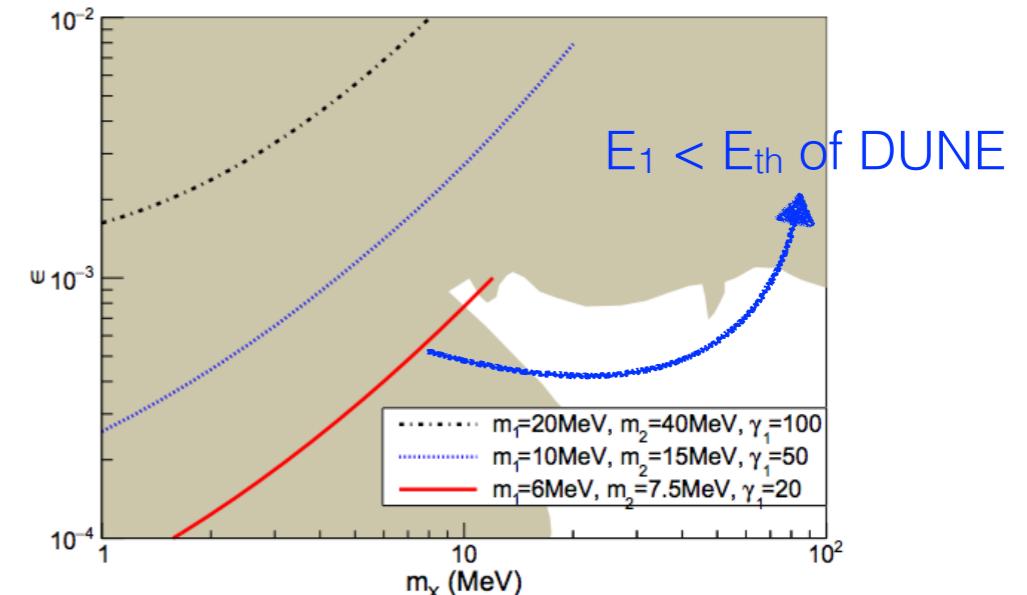


2200L of liquid scintillator
(~ 2 ton)

106kg array of 8 ultra-pure NaI(Tl) crystals
immersed in an active veto detector

Observed: 21 events

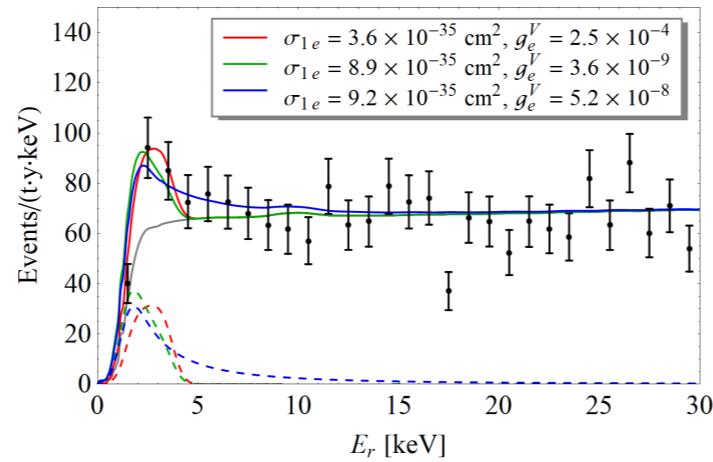
Background expected: 16.4 ± 2.1



Backup

- Dark Matter direct detection experiments

June 2020



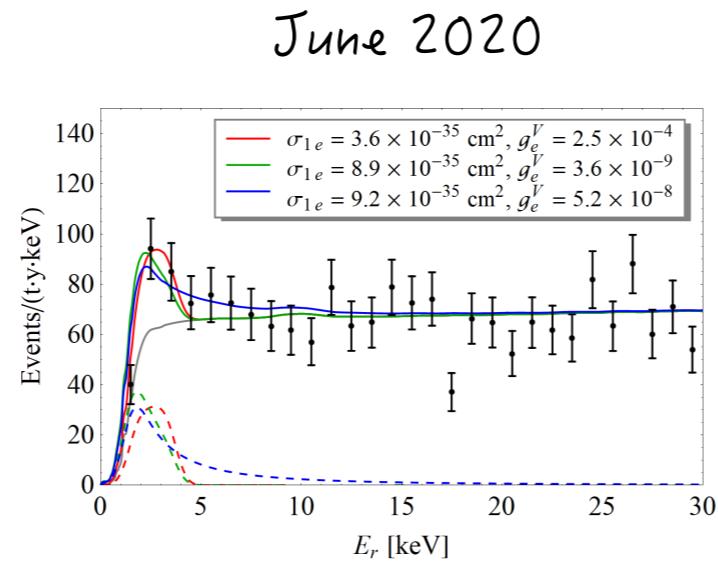
Smoking gun??

Giudice, Kim, Park, **SS**, PLB 2018

Alhazmi, Kim, Kong,
Mohlabeng, Park, **SS**,
JHEP 2021

Backup

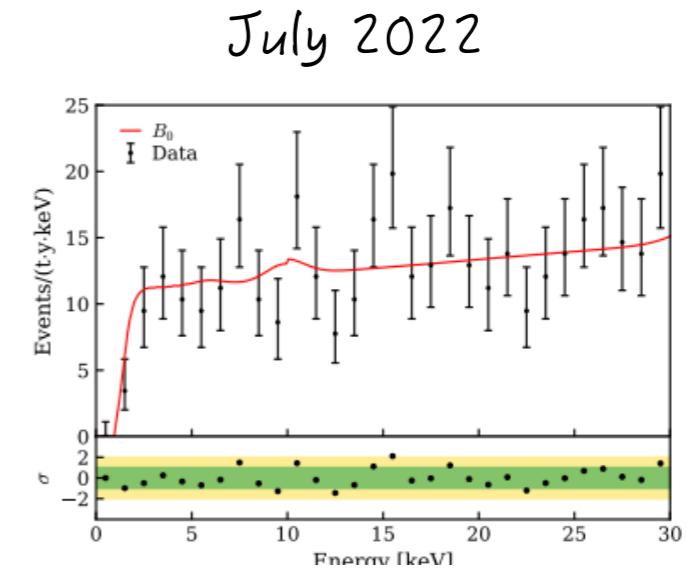
- Dark Matter direct detection experiments



Smoking gun??

Giudice, Kim, Park, **ss**, PLB 2018

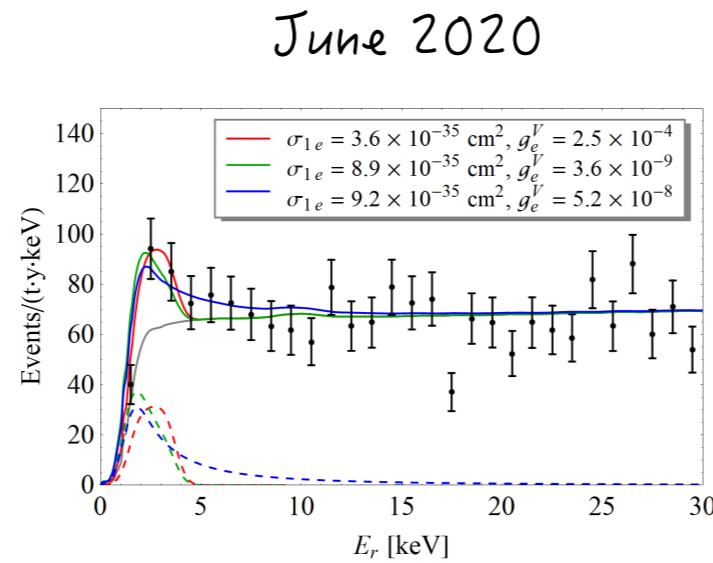
Alhazmi, Kim, Kong,
Mohlabeng, Park, **ss**,
JHEP 2021



Tritium contamination in Xe1T

Backup

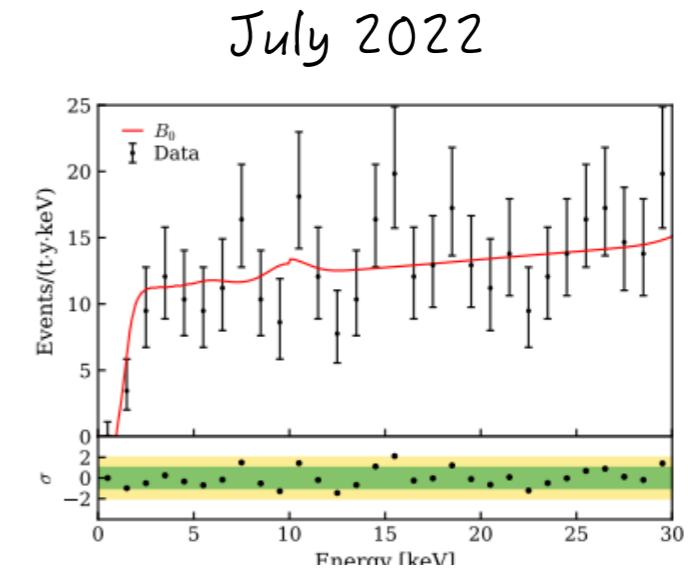
- Dark Matter direct detection experiments



Smoking gun??

Giudice, Kim, Park, **ss**, PLB 2018

Alhazmi, Kim, Kong,
Mohlabeng, Park, **ss**,
JHEP 2021



Tritium contamination in Xe1T

Nevertheless, many powerful DM direct detection experiments are underway and let's see what they can discover!