

# Tensions between terrestrial and cosmological neutrino mass determinations

## ALPS 2023, Obergurgl, Austria, 30 March 2023



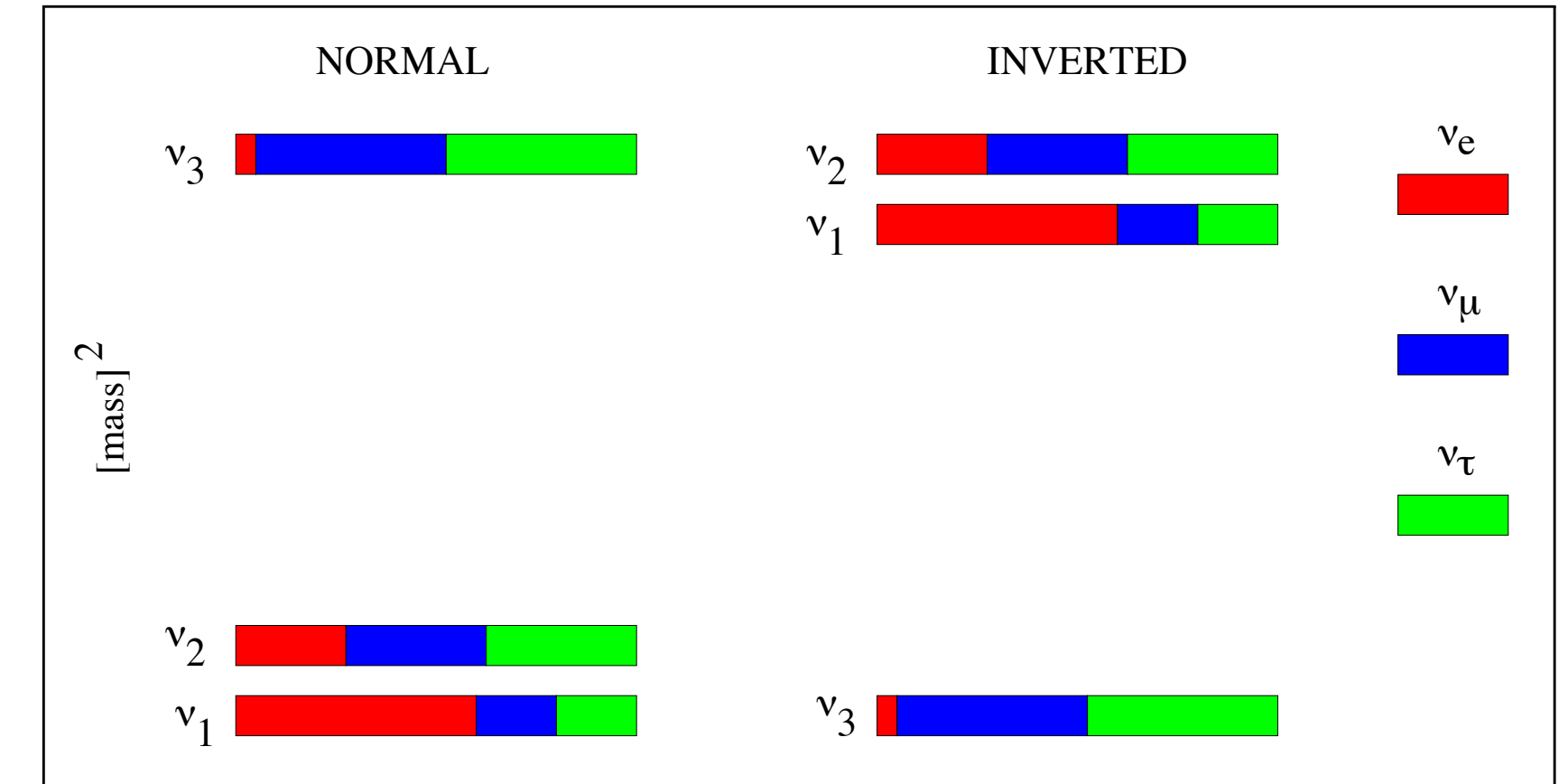
Thomas Schwetz  
Karlsruhe Institute of Technology, Institute for Astroparticle Physics



# Neutrino masses

## Neutrino oscillations:

- $|m_3^2 - m_1^2| \approx (2.5 \pm 0.03) \times 10^{-3} \text{ eV}^2$
- $m_2^2 - m_1^2 = (7.42 \pm 0.21) \times 10^{-5} \text{ eV}^2$



## Absolute mass determinations:

- beta-decay spectrum(KATRIN)
- neutrinoless double-beta decay (assuming Majorana neutrinos)
- cosmology

$$m_\beta = \sqrt{\sum_i |U_{ei}|^2 m_i^2} < 0.8 \text{ eV}$$

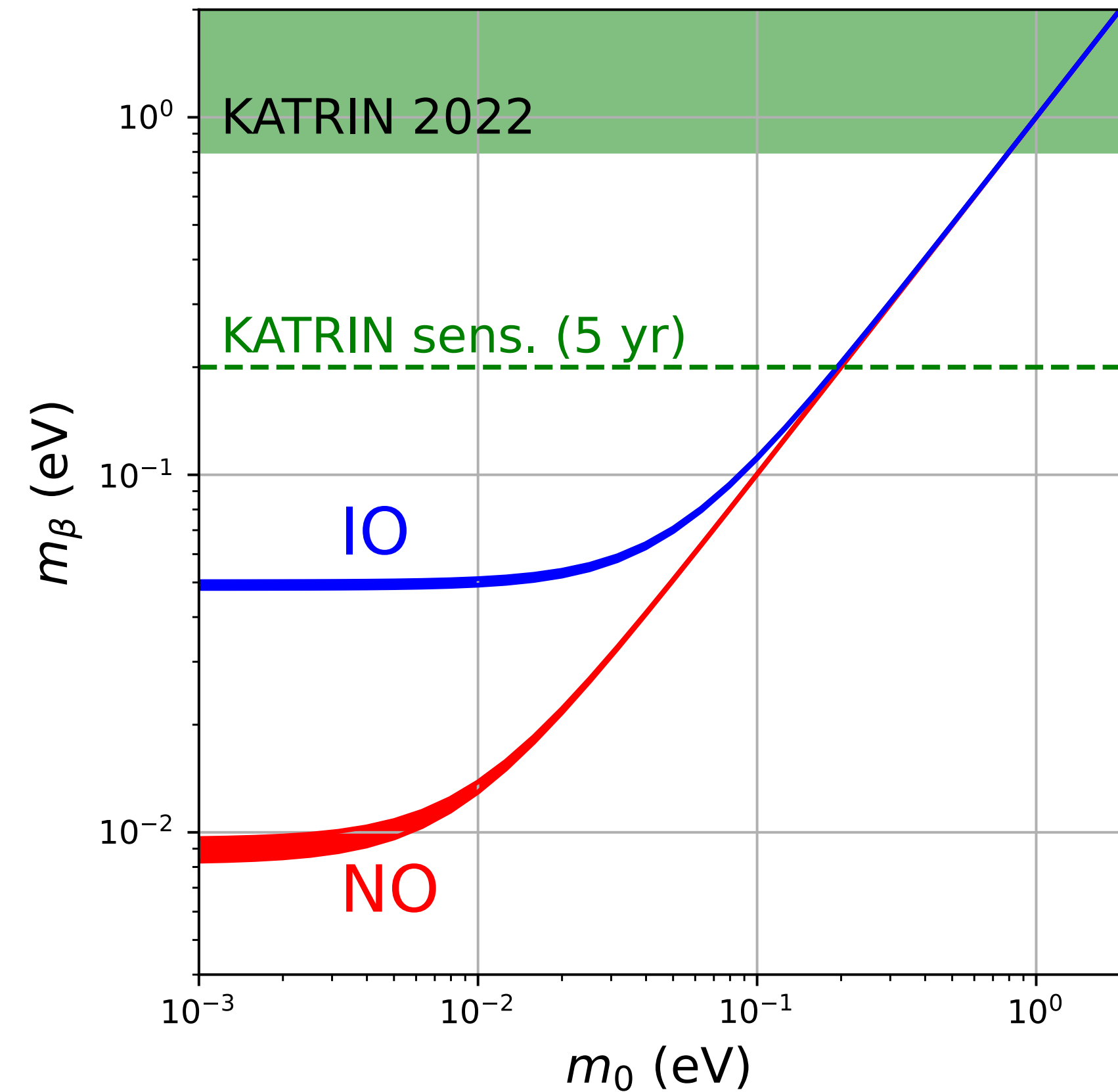
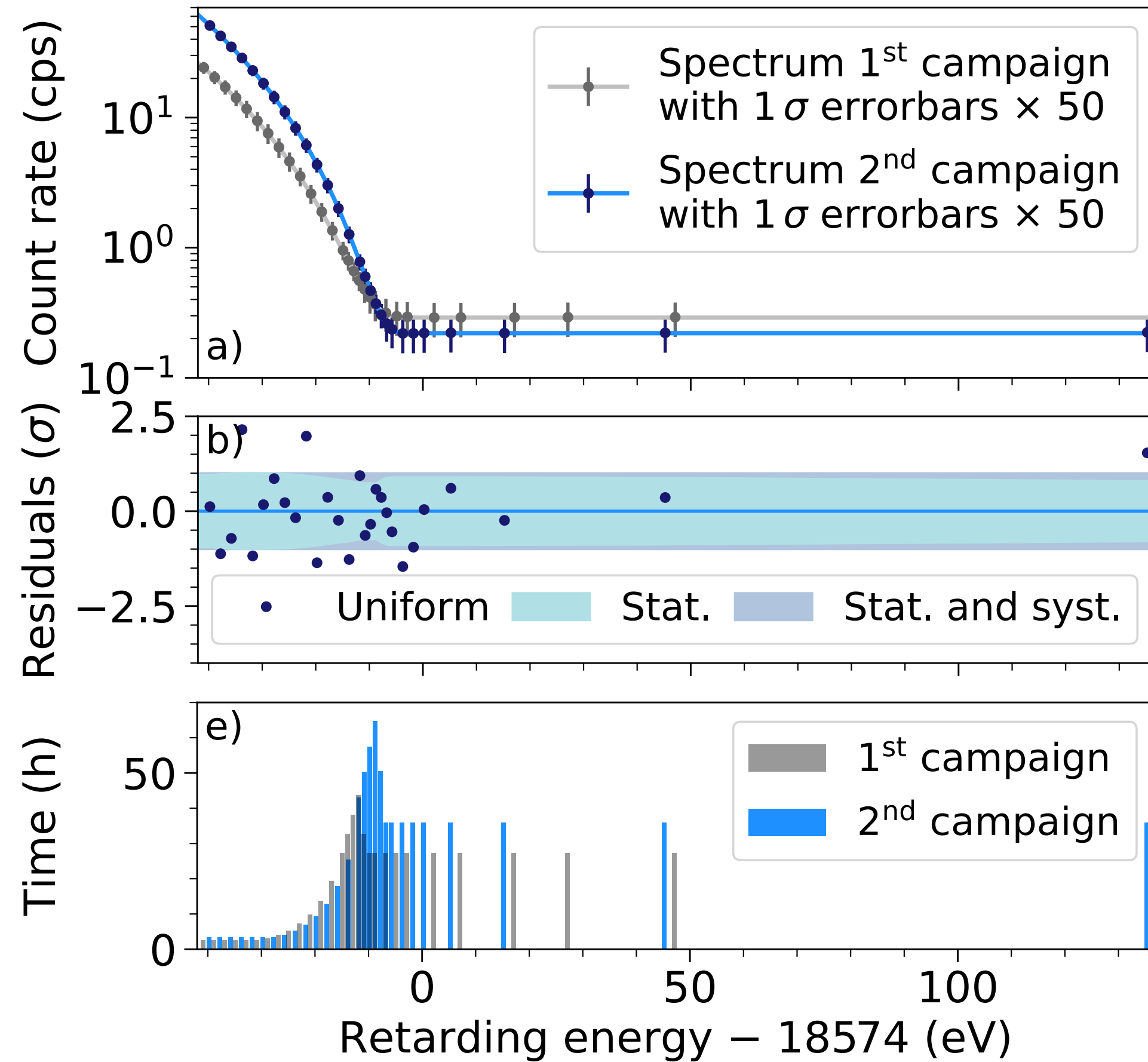
$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right| \lesssim 0.07 \text{ eV}$$

$$\sum_i m_i \lesssim 0.1 \text{ eV}$$

# Beta decay spectrum — KATRIN

KATRIN, Nature Phys. 18 (2022) 160 [2105.08533]

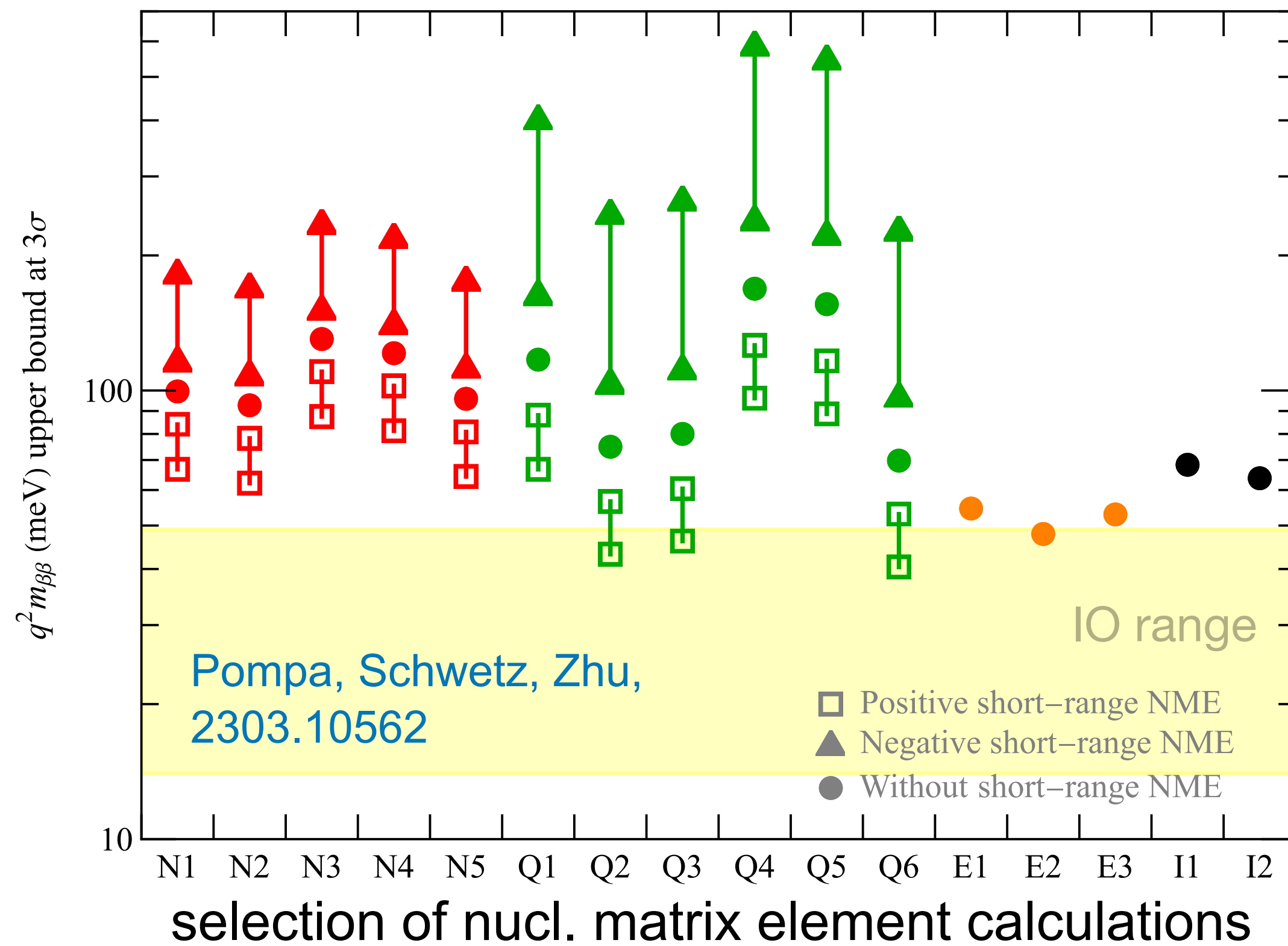
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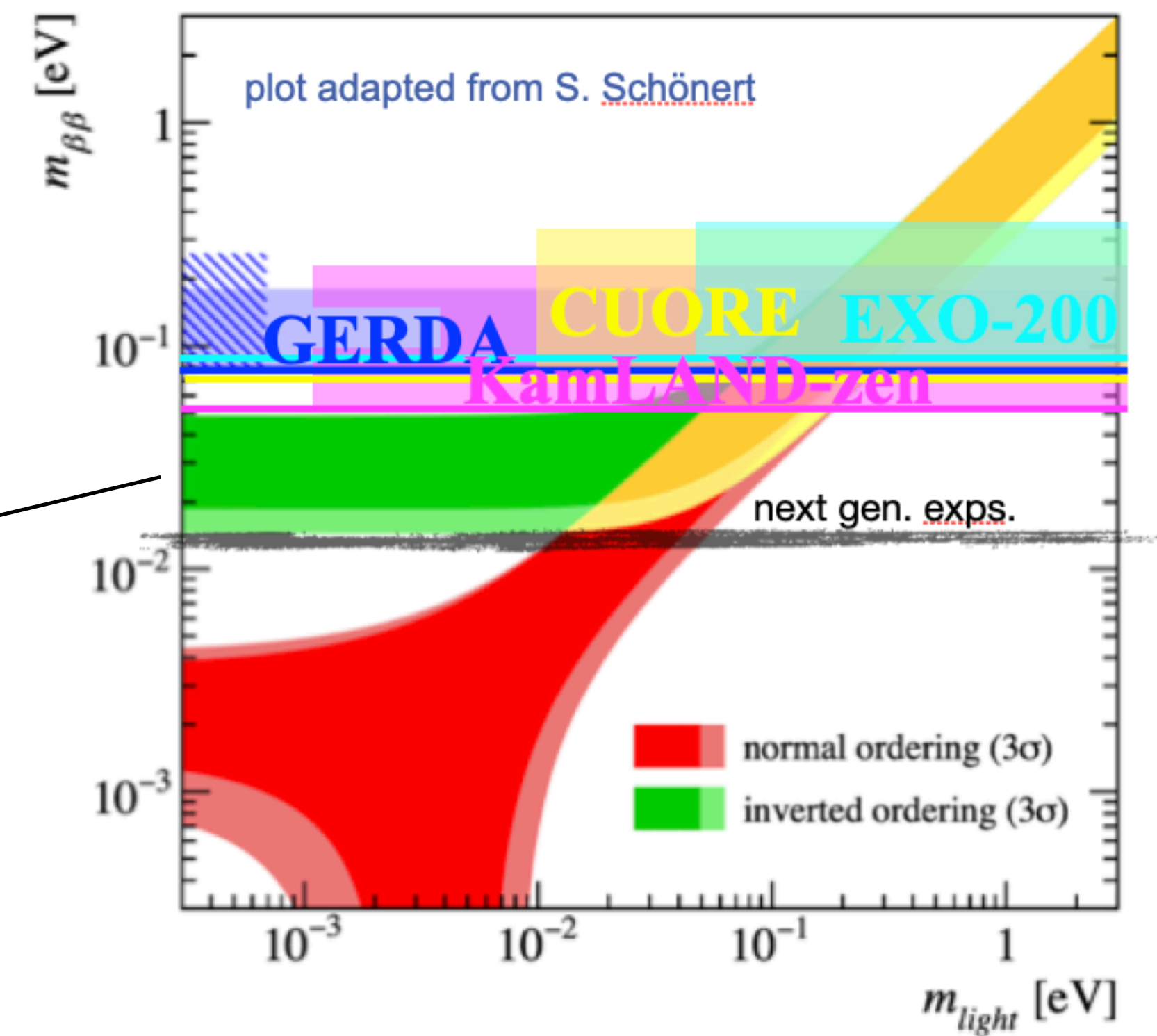
# Neutrinoless double-beta decay $\Rightarrow$ lepton number violation

Combined  $3\sigma$  upper bound from CUORE, EXO, GERDA, KamLLAND-Zen, MAJORANA

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_i \right|$$



neutrino mass interpretation affected by nuclear matrix elements and Majorana phases



new short-range contribution to NME [Cirigliano et al., 1802.10097](#)

see also talk by Simone Quitadamo



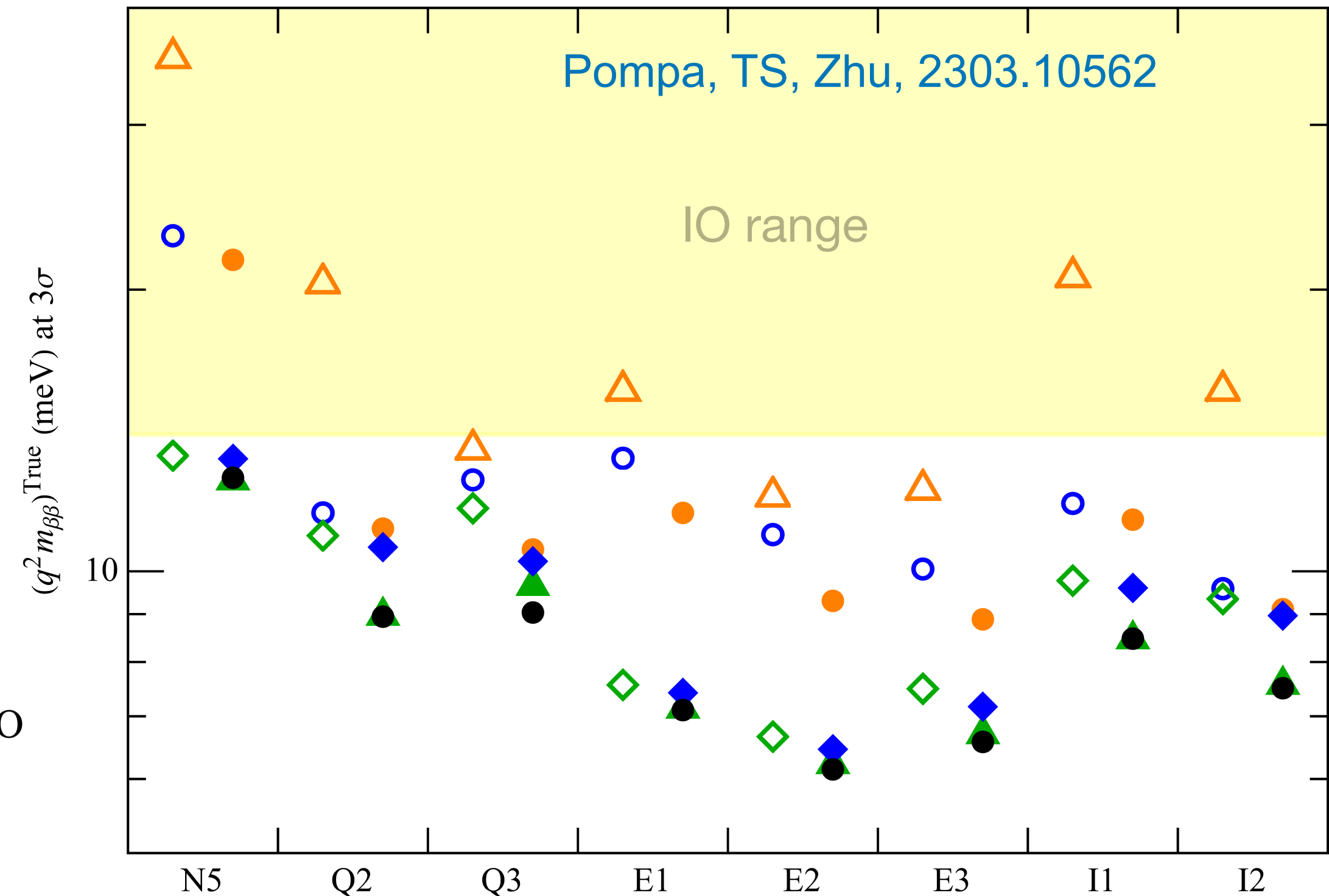
# Sensitivity of selected future $0\nu\beta\beta$ projects

	Isotope	exposure mol yr	backgr #/mol yr
<b>LEGEND</b>	$^{76}\text{Ge}$	8736	$4.9\text{e-}6$
<b>nEXO</b>	$^{136}\text{Xe}$	13700	$4.0\text{e-}5$
<b>CUPID</b>	$^{100}\text{Mo}$	1717	$2.3\text{e-}4$

Agostini et al., 2202.01787

- LEGEND-1000
- ▲ CUPID
- ◇ nEXO
- LEGEND-1000, CUPID
- ▲ LEGEND-1000, nEXO
- ◆ CUPID, nEXO
- LEGEND-1000, CUPID, nEXO

for  $T = 10$  yr



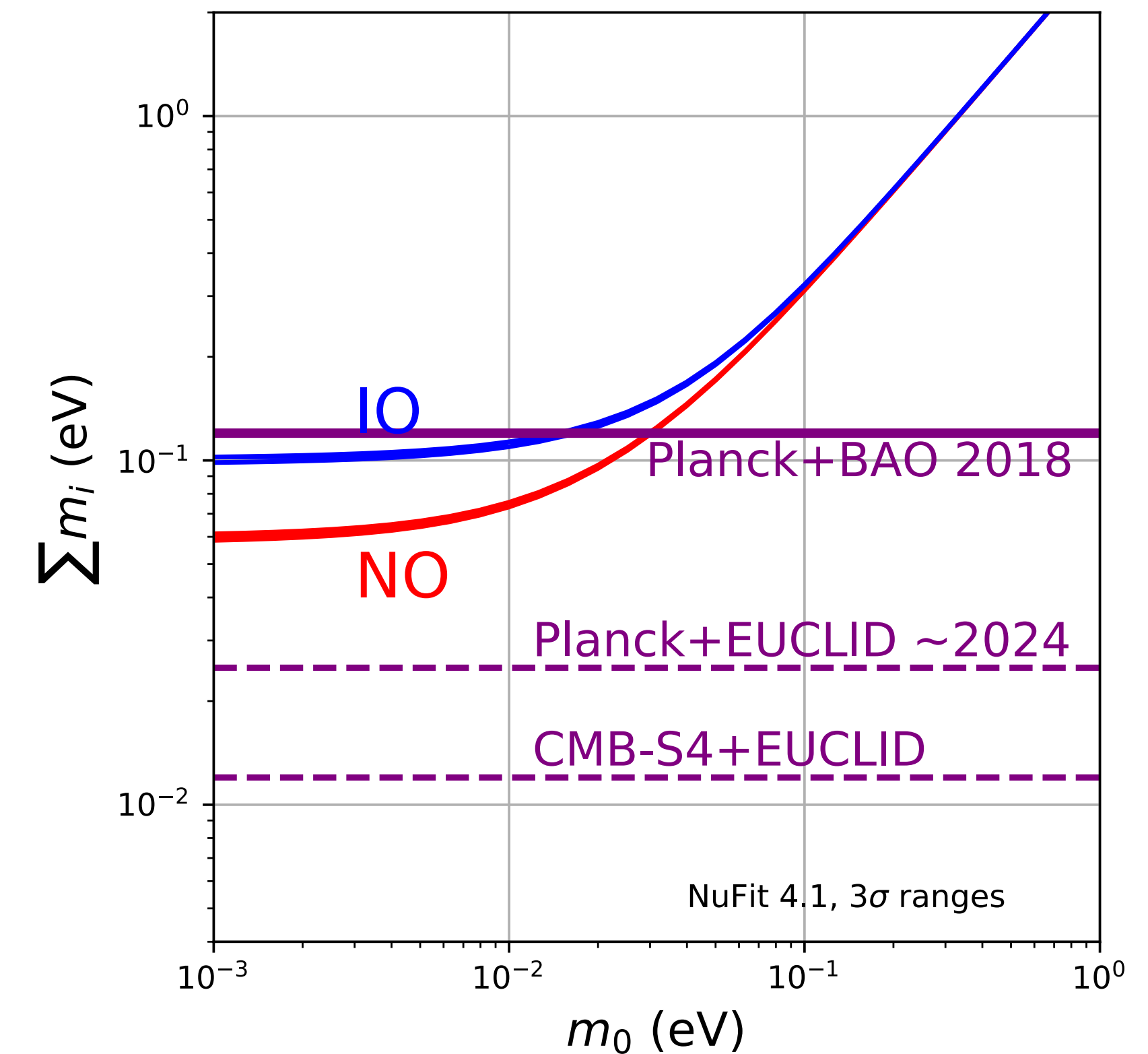
selection of nucl. matrix element calculations



# Neutrino mass from cosmology

see also talk by Vivian Poulin

$$\Sigma \equiv \sum_{i=1}^3 m_i = \begin{cases} m_0 + \sqrt{\Delta m_{21}^2 + m_0^2} + \sqrt{\Delta m_{31}^2 + m_0^2} & \text{(NO)} \\ m_0 + \sqrt{|\Delta m_{32}^2| + m_0^2} + \sqrt{|\Delta m_{32}^2| - \Delta m_{21}^2 + m_0^2} & \text{(IO)} \end{cases}$$



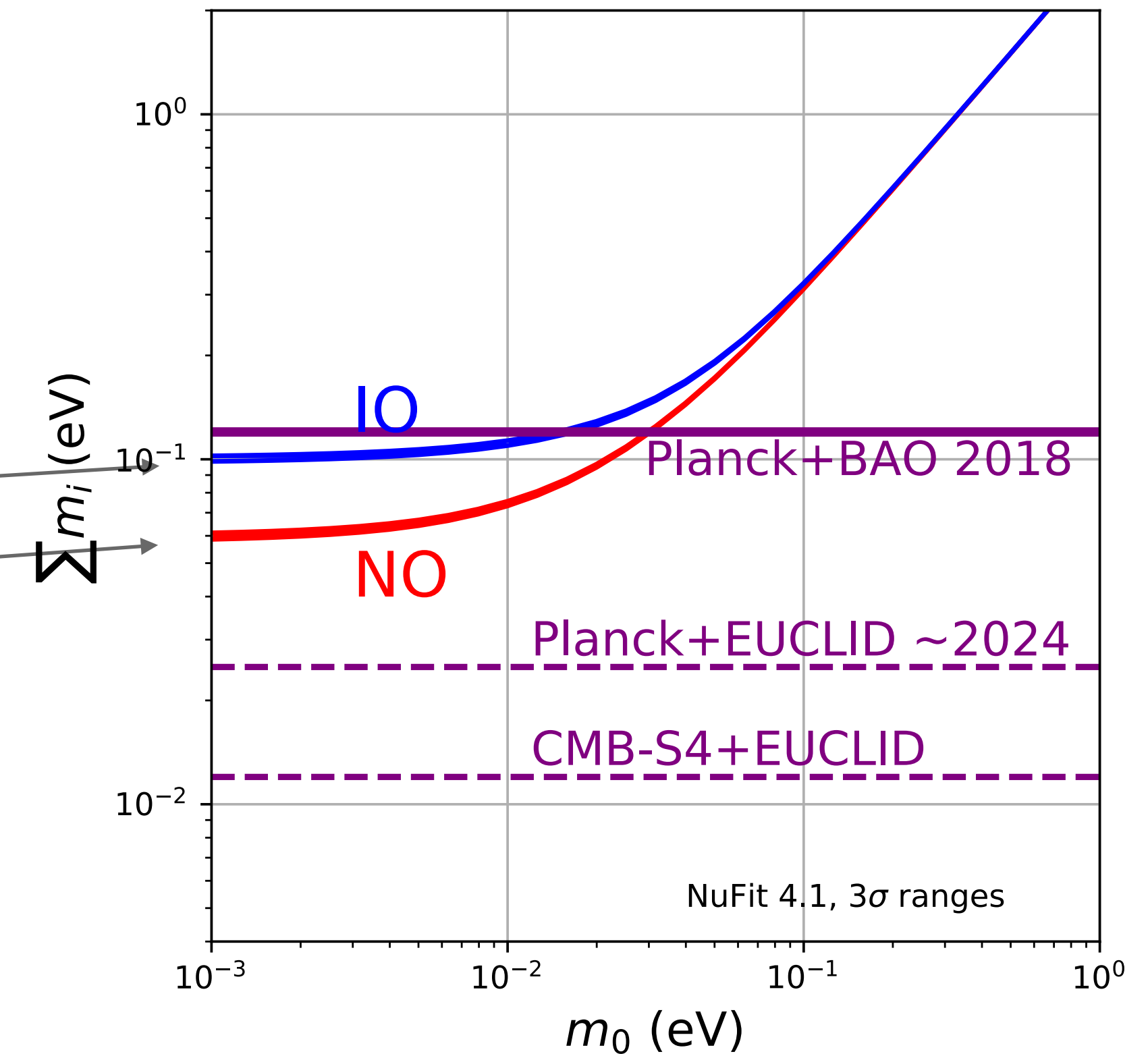


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- minimal values predicted from oscillation data for  $m_0 = 0$ :

$$\Sigma_{\min} = \begin{cases} 98.6 \pm 0.85 \text{ meV} & \text{(IO)} \\ 58.5 \pm 0.48 \text{ meV} & \text{(NO)} \end{cases}$$





# Neutrino mass from cosmology

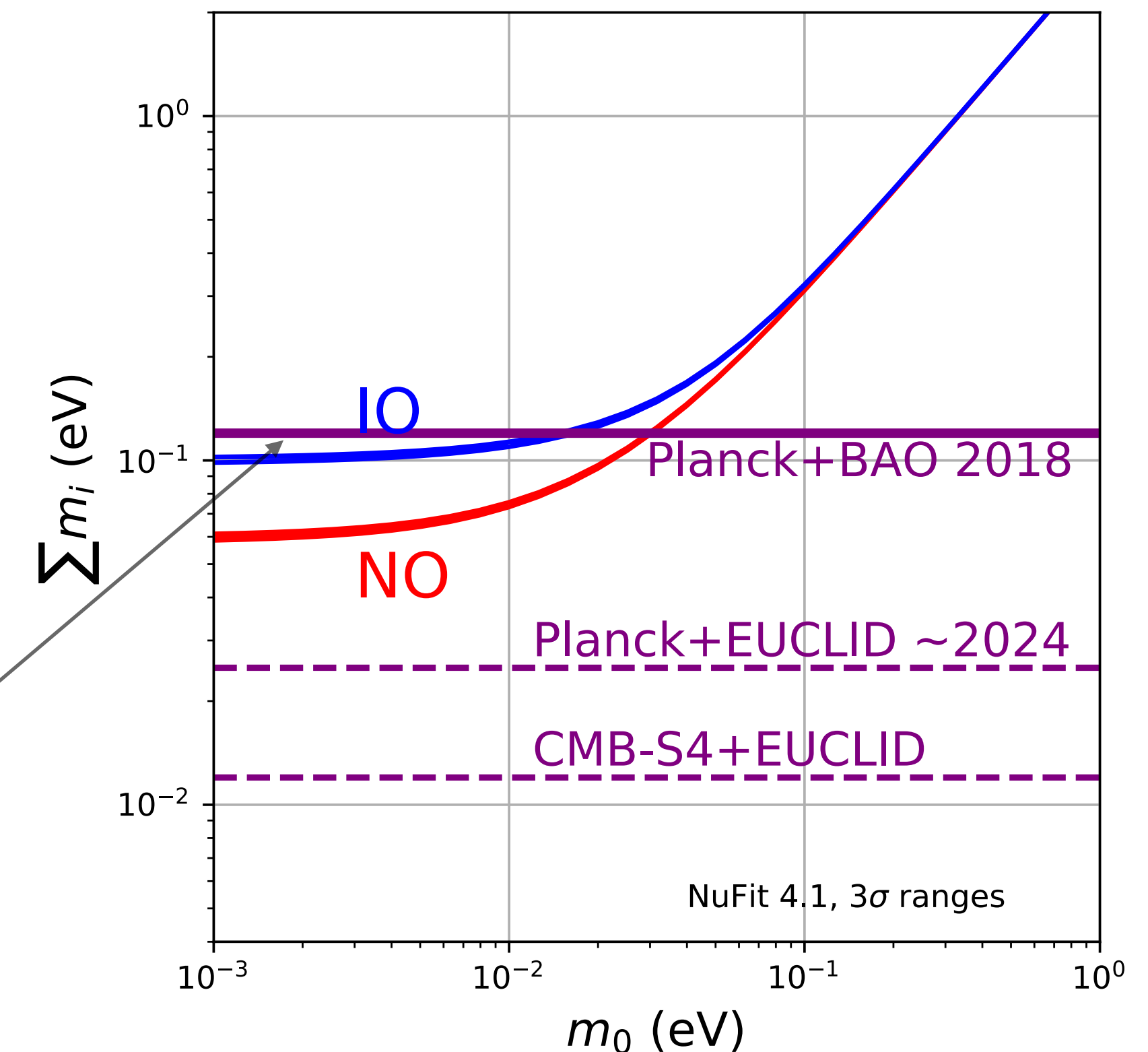
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- Upper bounds from current data:**

- $\Sigma m_\nu < 0.12 \text{ eV}$  (95 % CL) **Planck CMB+BAO**
- $\Sigma m_\nu < 0.09 \text{ eV}$  (95 % CL)  
**DiValentino, Gariazzo, Mena, 21; many papers**



# Excluding inverted ordering with cosmology?

- **Strong Bayesian Evidence for the Normal Neutrino Hierarchy**  
Simpson et al., 1703.03425;  
Jimenez et al., 2203.14247
- **No conclusive evidence for normal ordering:** TS et al. 1703.04585;  
Vagnozzi et al., 1701.08172; Gariazzo et al., 1801.04946; Heavens, Sellentin,  
1802.09450; deSalas et al., 1806.11051; Mahony et al., 1907.04331;  
Hannestad, Roy Choudhury, 1907.12598; Gariazzo et al., 2205.02195

## Bayesian model comparison:

$$B_{\text{NO,IO}} \equiv \frac{\mathcal{Z}_{\text{NO}}}{\mathcal{Z}_{\text{IO}}}$$

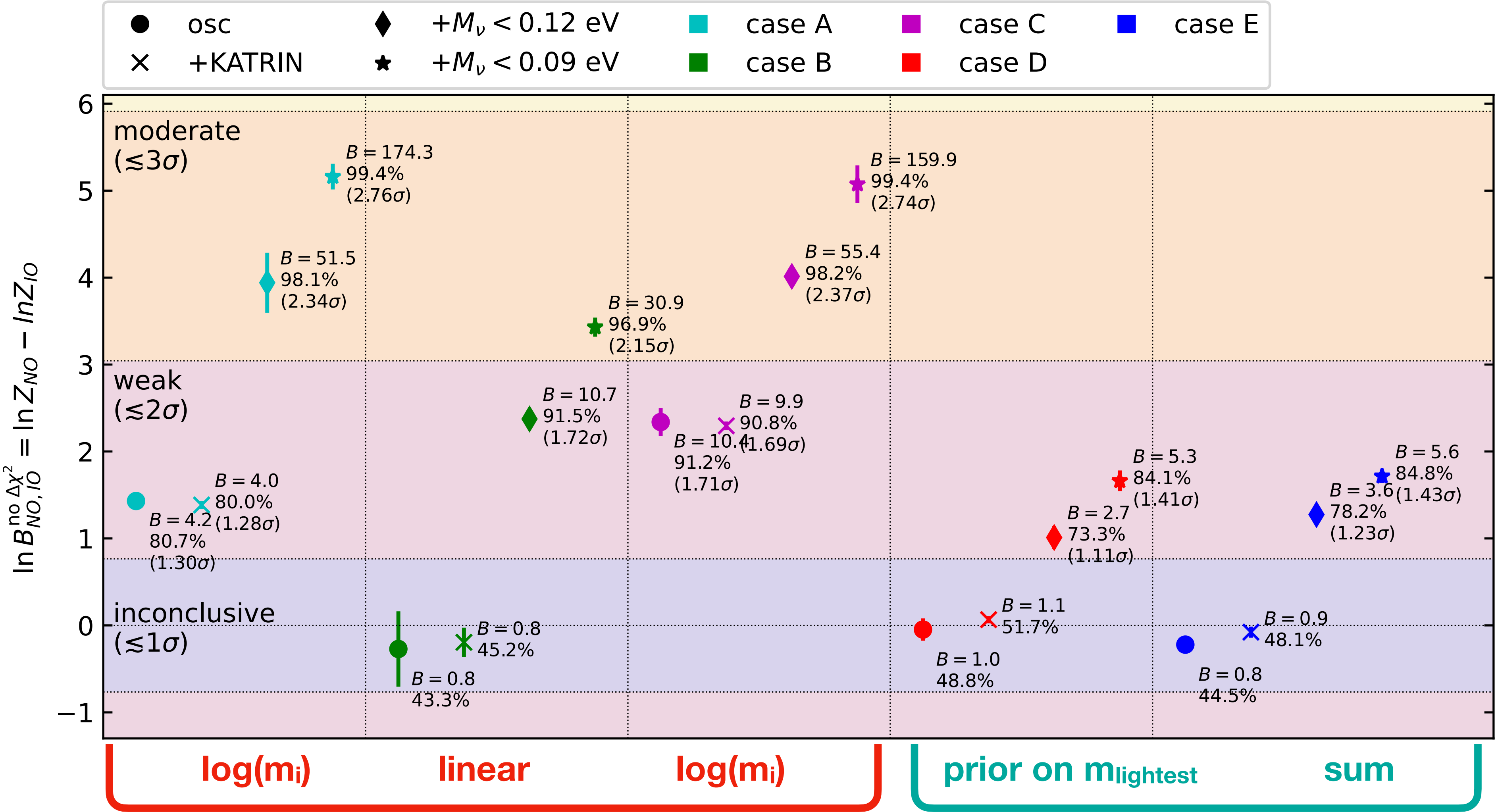
$$\mathcal{Z}_D = P(D|M) = \int d\theta \mathcal{L}_D(\theta) \Pi(\theta)$$

- current preference for NO
  - from cosmology is **prior driven** (not data driven)
  - in combined analysis dominated by **oscillation data**



# Preference for normal ordering (w/o $\Delta\chi^2_{IO/NO}$ from oscillation)

log Bayesian evidence ratio



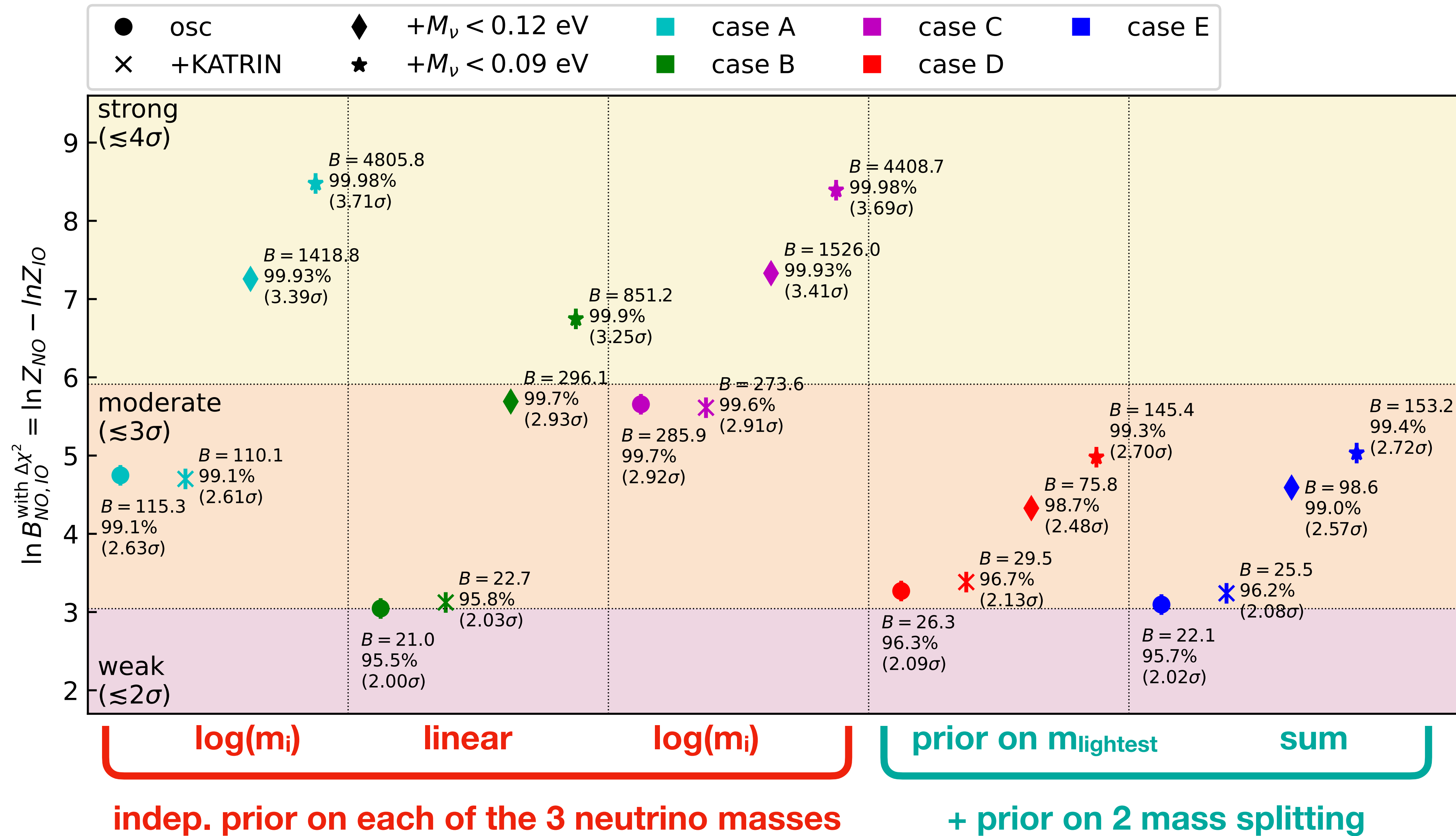
indep. prior on each of the 3 neutrino masses

+ prior on 2 mass splitting

Gariazzo et al., 2205.02195

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log Bayesian evidence ratio

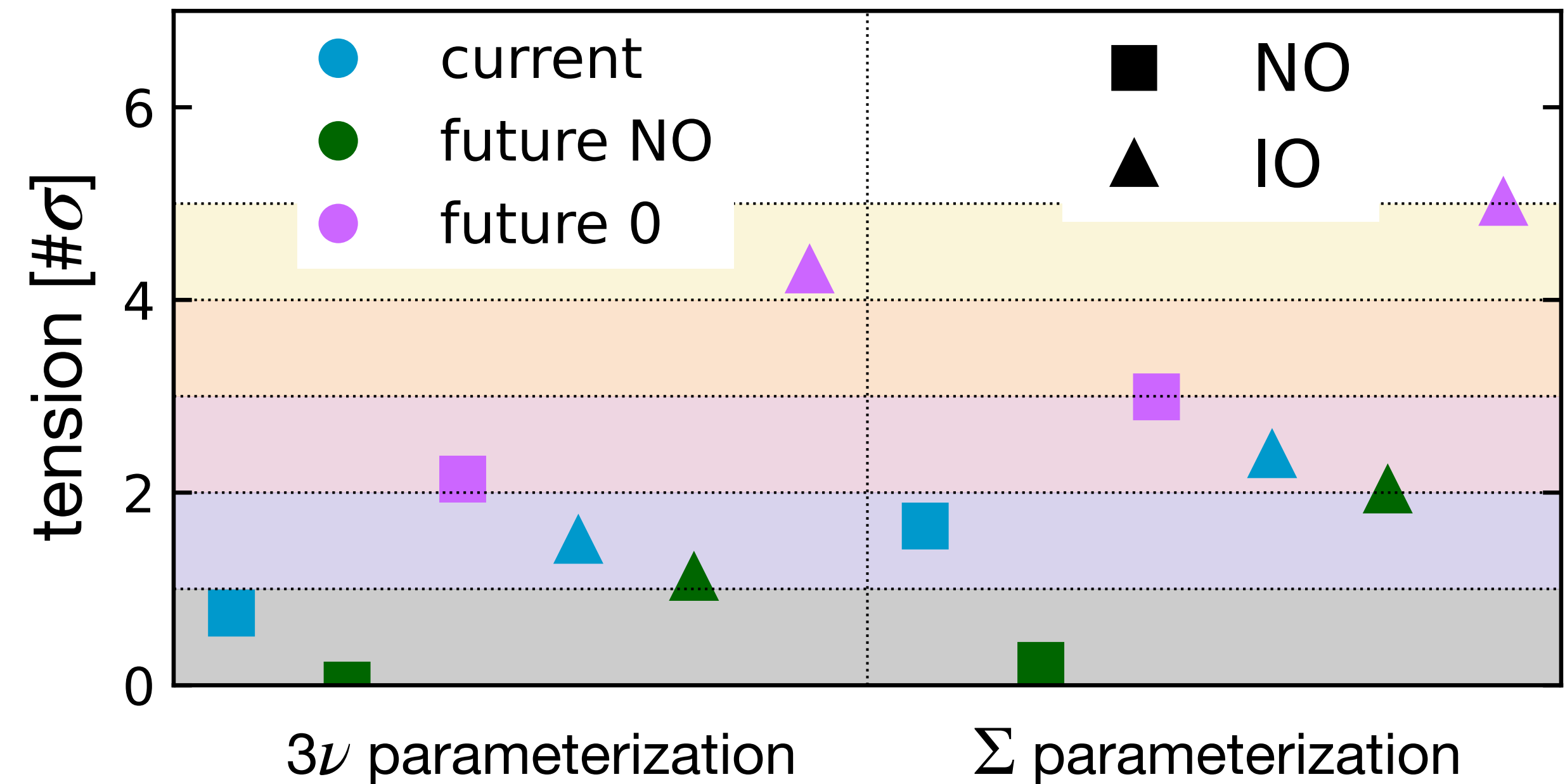
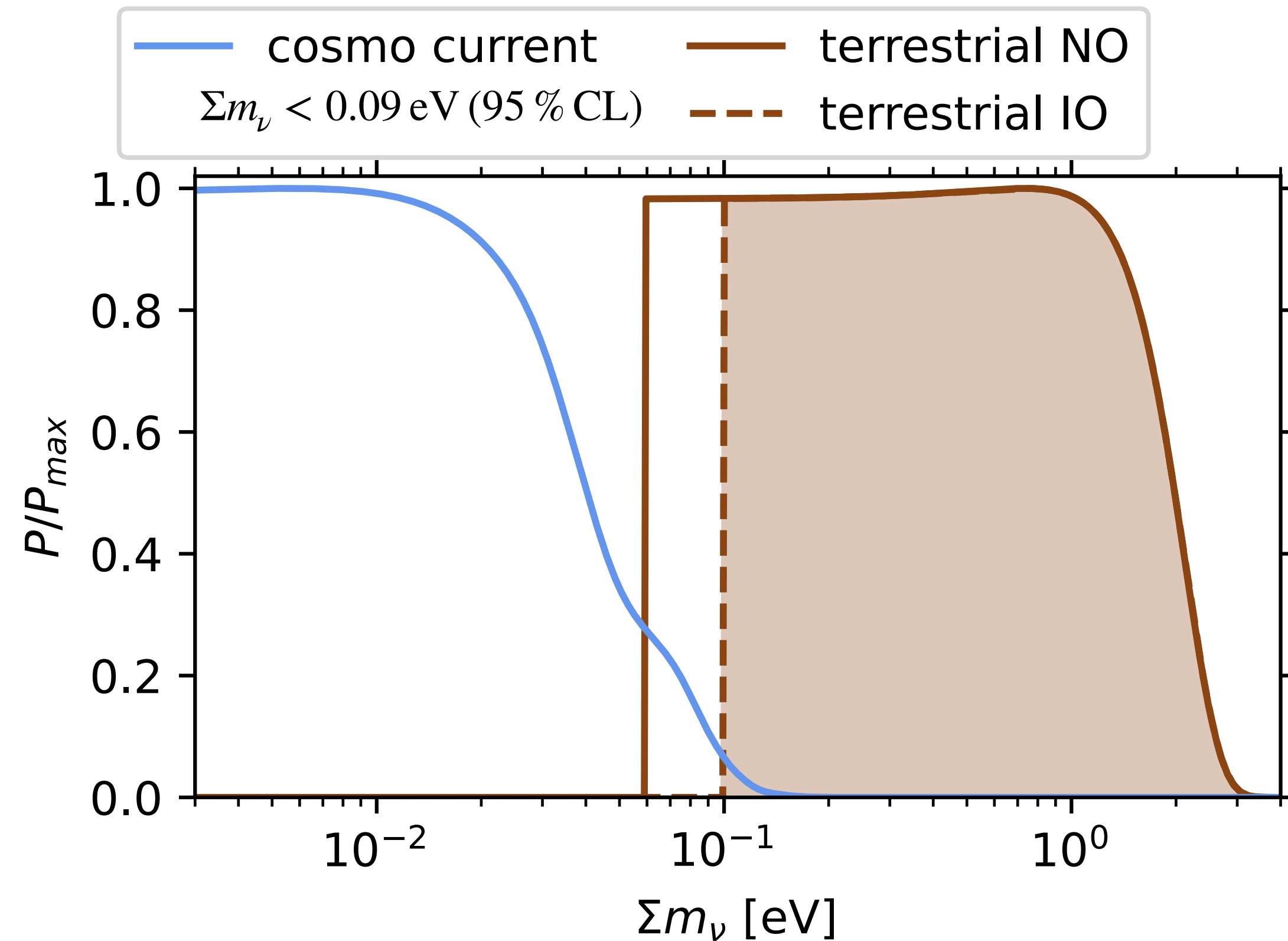


Gariazzo et al., 2205.02195



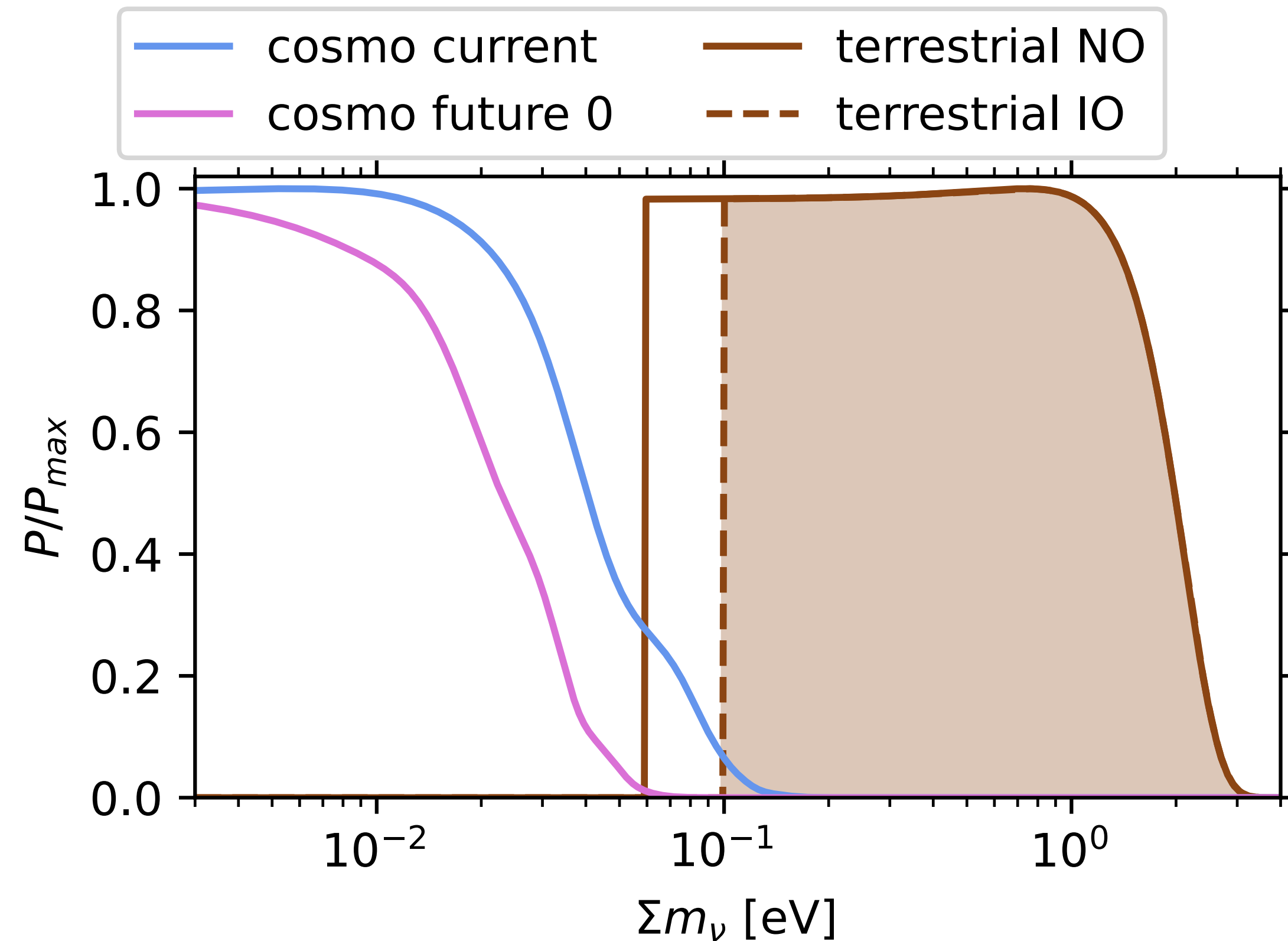
# Tension between cosmology and terrestrial data

Gariazzo, Mena, Schwetz, 2302.14159

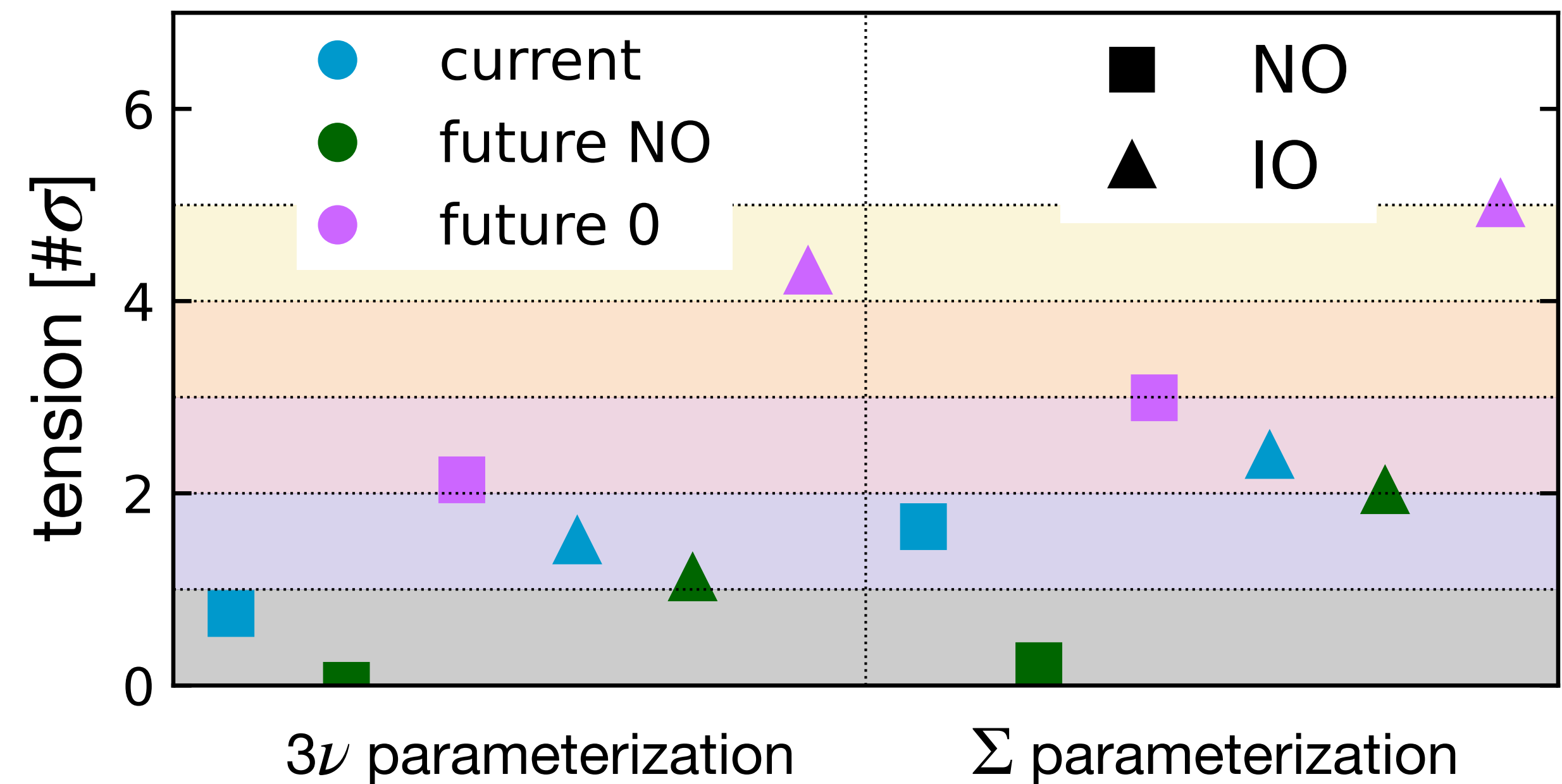


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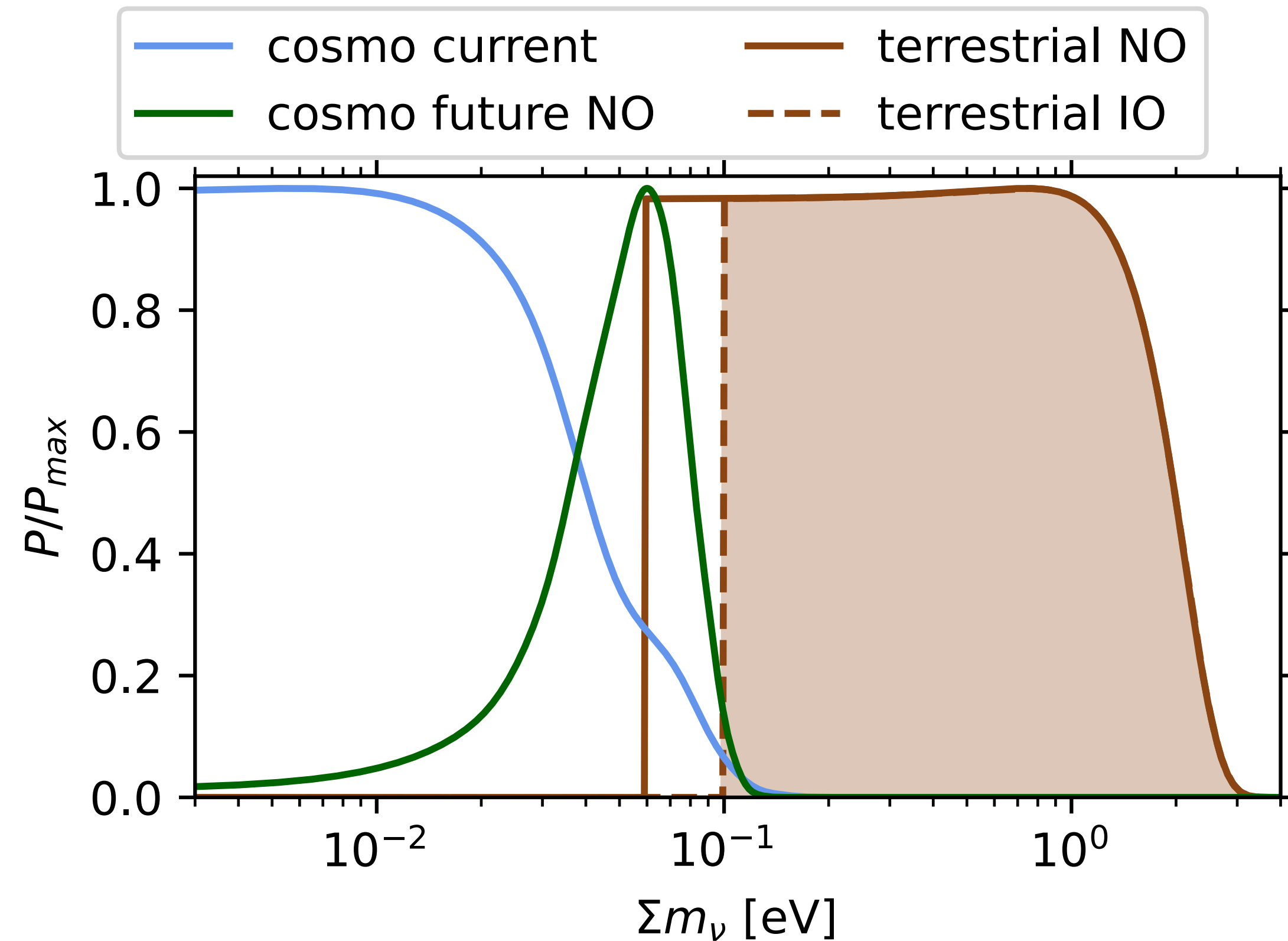
future 0:  $\Sigma m_\nu < 0.02 \text{ eV} (1\sigma)$



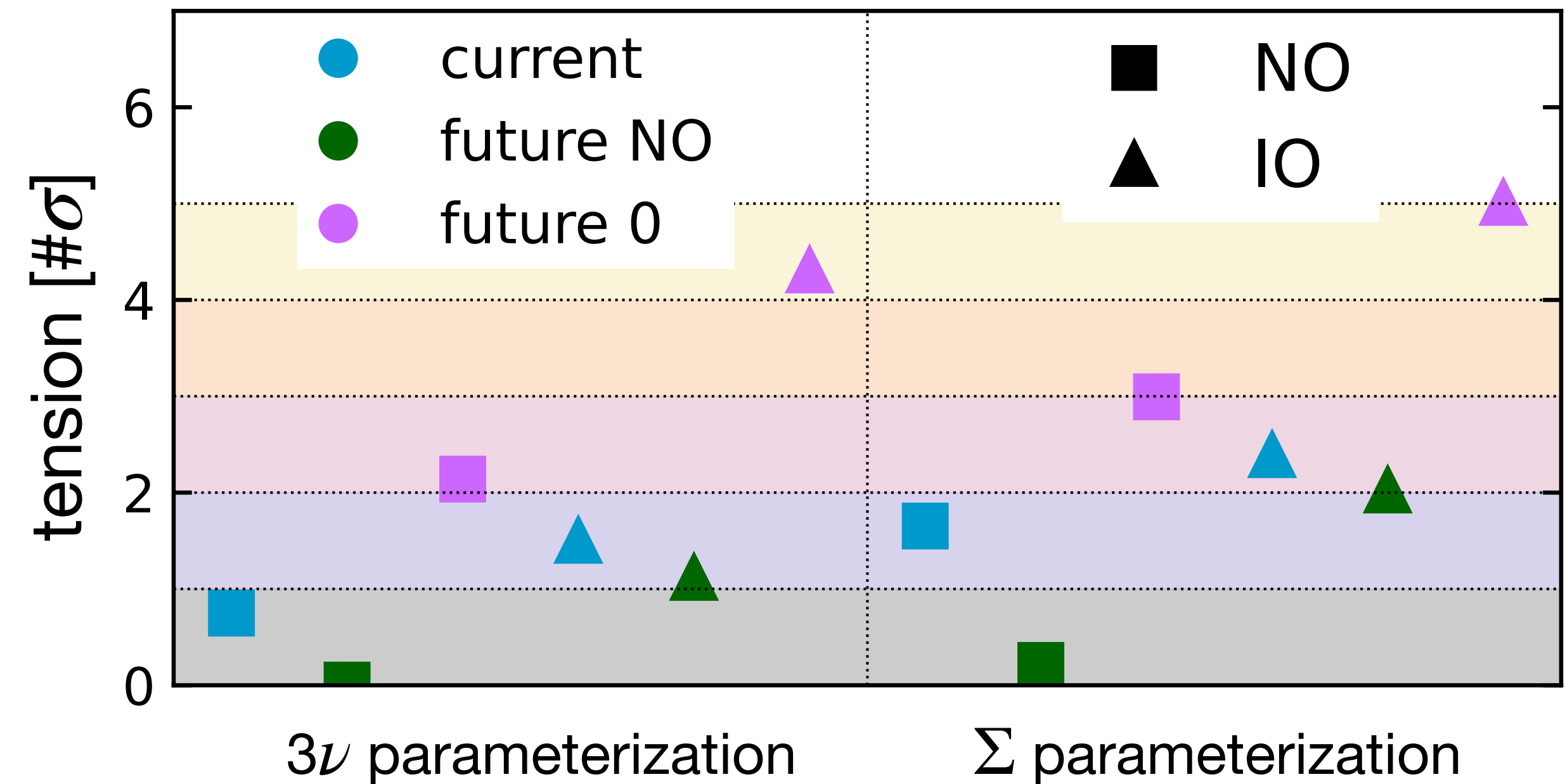


# Tension between cosmology and terrestrial data

Gariazzo, Mena, Schwetz, 2302.14159



future NO:  $\Sigma m_\nu = (0.06 \pm 0.02) \text{ eV} (1\sigma)$



- What if cosmology does not see finite neutrino mass and upper bounds become tighter than the minimal value predicted by neutrino oscillation?
- Can we relax cosmological bounds such that neutrino mass can be in reach for terrestrial experiments?



# Cosmology bounds can be relaxed in non-standard scenarios

- neutrino decay into dark radiation

Chacko et al. 1909.05275; 2002.08401; Escudero et al., 2007.04994;

Barenboim et al., 2011.01502; Chacko et al. 2112.13862:  $\sum m_\nu < 0.42 \text{ eV}$

- time dependent neutrino mass

Lorenz et al. 1811.01991; 2102.13618; Esteban, Salvado, 2101.05804

- modified momentum distribution

Cuoco et al., astro-ph/0502465; Barenboim et al., 1901.04352;

Alvey, Sabti, Escudero, 2111.14870

- reduced neutrino density + dark radiation

Beacom, Bell, Dodelson, 04; Farzan, Hannestad, 1510.02201;

Renk, Stöcker et al., 2009.03286; Escudero, TS, Terol-Calvo, 2211.01729

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# Relaxing the neutrino mass bound from cosmology

## Cosmology is sensitive to:

- energy density in non-relativistic neutrinos (late times)

$$\rho_\nu^{\text{non.rel.}} \approx n_\nu \sum m_\nu < 14 \text{ eV cm}^{-3}$$

- energy density in relativistic neutrinos (early times, BBN, CMB)

$$N_{\text{eff}}^{\text{relat.}} = 2.99 \pm 0.17$$



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- energy density in relativistic neutrinos (early times, BBN, CMB)

$$N_{\text{eff}}^{\text{relat.}} = 2.99 \pm 0.17$$

relax bound on  $m_\nu$  by reducing neutrino number density

$$\sum m_\nu < 0.12 \text{ eV} \left( \frac{n_\nu^{\text{SM}}}{n_\nu} \right)$$

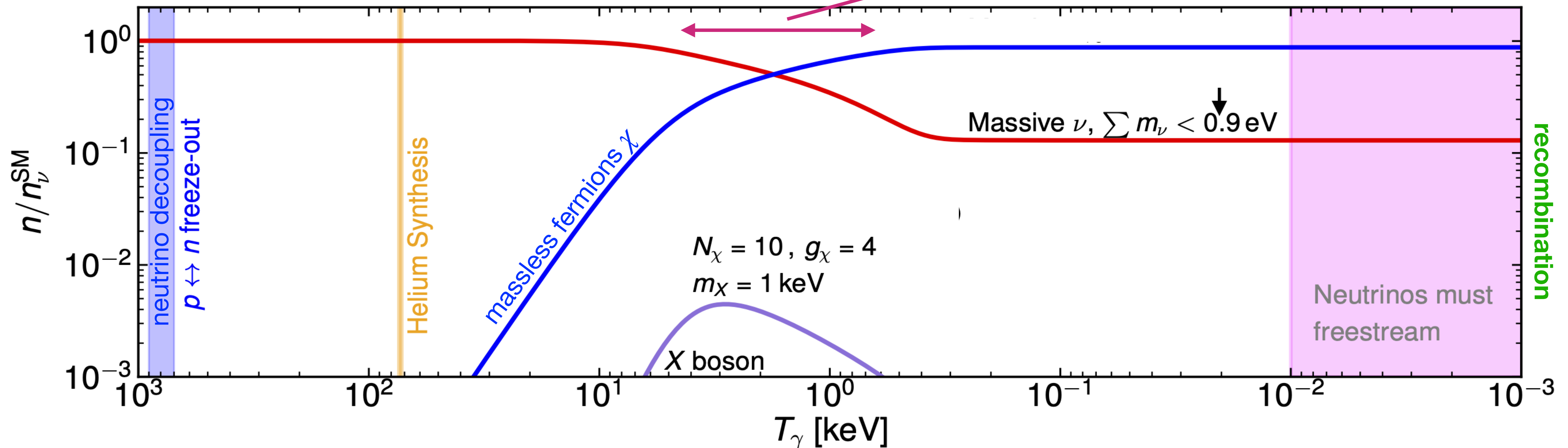
introduce „dark radiation“ to keep  $N_{\text{eff}}^{\text{relat.}} \approx 3$

$$N_{\text{eff}}^{\text{relat.}} = N_{\text{eff}}^\nu + N_{\text{eff}}^{\text{DR}} \approx 3$$

- introduce a set of  $N_\chi$  massless fermions
- a mediator  $X$  coupled to neutrinos (scalar or vector)
- convert active neutrinos into massless fermions after BBN but before CMB decoupling

$$\Gamma(\nu\nu \leftrightarrow X) \gtrsim H$$

$$\Gamma(X \leftrightarrow \chi\chi) \gtrsim H$$



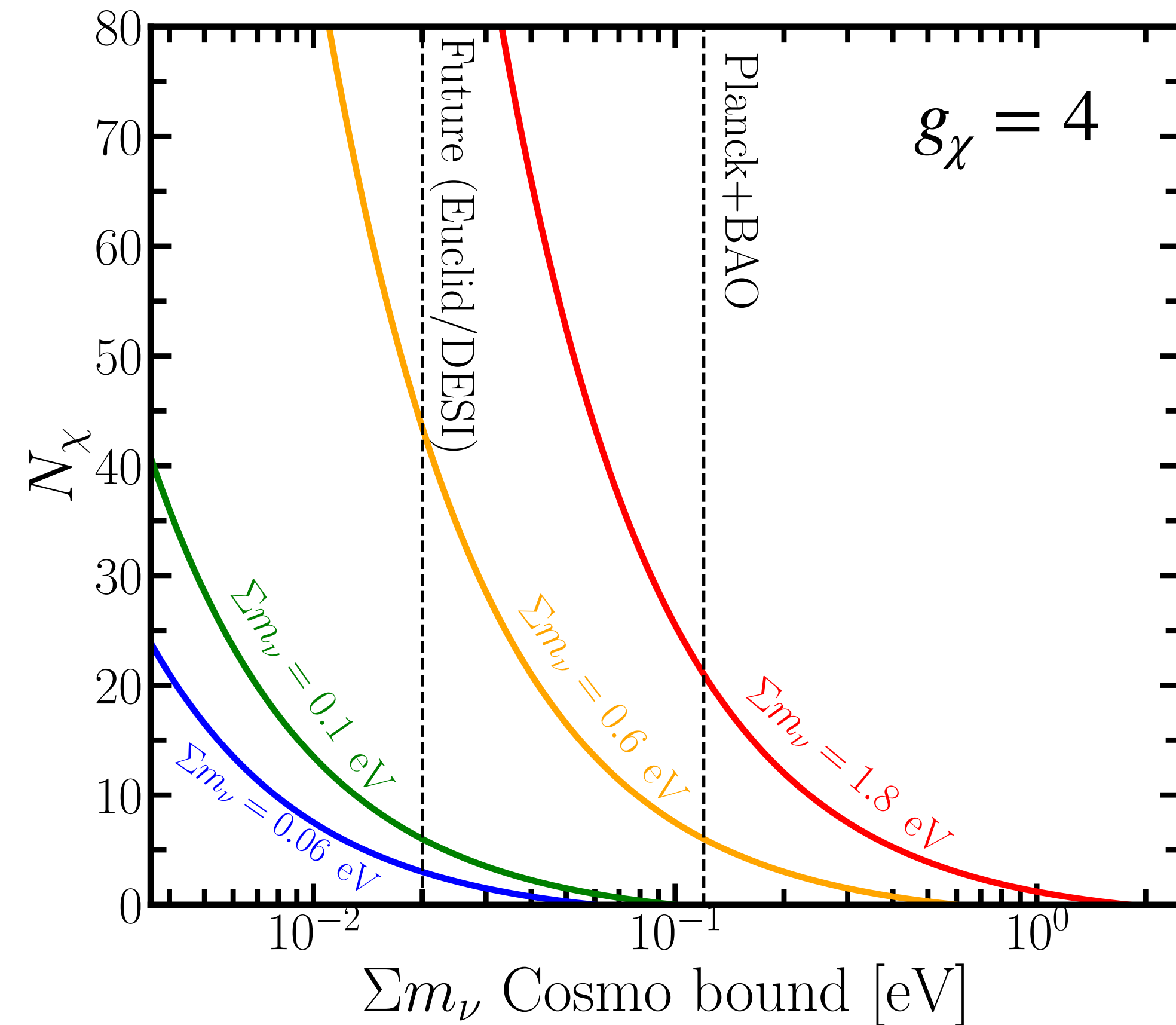
# Relaxed bound from cosmology

Farzan, Hannestad, 1510.02201  
Escudero, TS, Terol-Calvo, 2211.01729

relaxing the present bound by  
converting neutrinos into  $N_\chi$  generations  
of massless fermions with  $g_\chi$  internal  
degrees of freedom:

$$\sum m_\nu < 0.12 \text{ eV} (1 + g_\chi N_\chi / 6)$$

need  $\gtrsim 10$  massless species for  $m_\nu \sim 1 \text{ eV}$





# A seesaw model for large neutrino mass and dark radiation

Escudero, TS, Terol-Calvo, 2211.01729

- 3 heavy right-handed neutrinos (seesaw)
- new abelian symmetry  $U(1)_X$  local or global
- a scalar  $\Phi$  charged under  $U(1)_X$
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**Yukawa sector**

$$-\mathcal{L} = \overline{N}_R Y_\nu \ell_L \tilde{H}^\dagger + \frac{1}{2} \overline{N}_R M_R N_R^c + \overline{N}_R Y_\Phi \chi_L \Phi + \text{h.c.}$$

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**Scalar potential**

$$V = \mu_H^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + \mu_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4 + \lambda_{H\Phi} |\Phi|^2 H^\dagger H$$



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## Gauge interaction

$$\mathcal{L}_{\text{int}} = g_X Z'_\mu \bar{\chi} \gamma^\mu \chi$$

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$$\mathcal{M}_n = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & M_R & \Lambda \\ 0 & \Lambda^T & 0 \end{pmatrix}$$

$$m_D = \frac{v_{EW}}{\sqrt{2}} Y_\nu, \quad \Lambda = \frac{v_\Phi}{\sqrt{2}} Y_\Phi$$

$$\Lambda \ll m_D \ll M_R$$

$$m_{\text{heavy}} \approx M_R$$

$$m_{\text{active}} \approx m_D^2 / M_R$$

$$m_\chi = 0, \quad \theta_{\nu\chi} \approx \Lambda / m_D$$

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Escudero, TS, Terol-Calvo, 2211.01729

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$$-\mathcal{L} = \overline{N_R} Y_\nu \ell_L \tilde{H}^\dagger + \frac{1}{2} \overline{N_R} M_R N_R^c + \overline{N_R} Y_\Phi \chi_L \Phi + \text{h.c.}$$

$$\mathcal{L}_{\text{int}} = g_X Z'_\mu \bar{\chi} \gamma^\mu \chi \quad g_X = \frac{m_{Z'}}{v_\Phi}$$

couplings to neutrinos induced by mixing:  $Z' \leftrightarrow \nu\nu/\nu\chi/\chi\chi$

$$\lambda_{Z'}^{\chi\chi} = g_X$$

$$\lambda_{Z'}^{\chi\nu} = g_X \theta_{\nu\chi}$$

$$\lambda_{Z'}^{\nu\nu} = g_X \theta_{\nu\chi}^2$$

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**indep. params for pheno:**

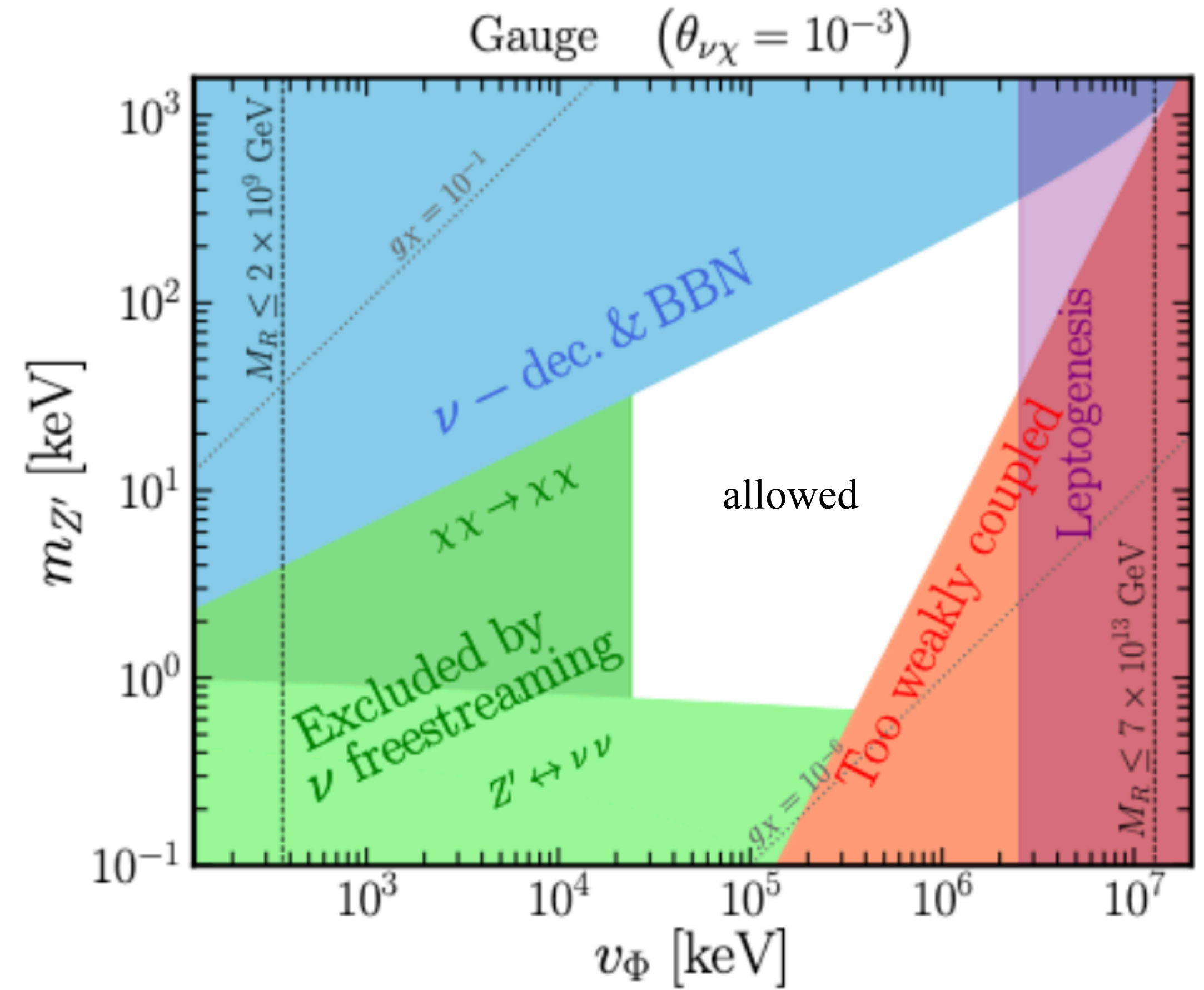
$$m_\nu, M_R, \theta_{\nu\chi}$$

$$v_\Phi, m_{Z'}$$

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$$\mathcal{L}_{\text{int}} = g_X Z'_\mu \bar{\chi} \gamma^\mu \chi \quad g_X = \frac{m_{Z'}}{v_\Phi}$$



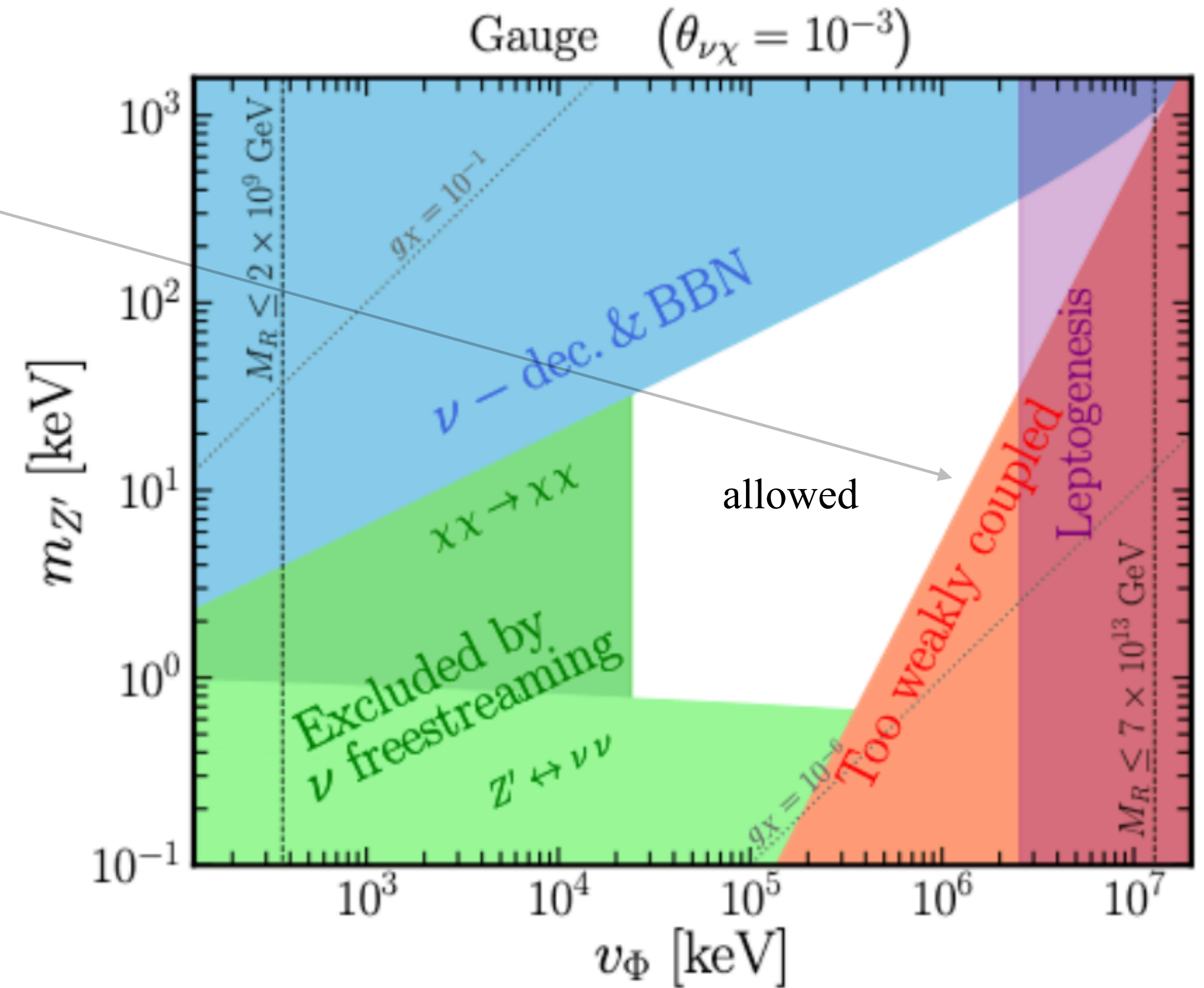
# Available parameter space



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- thermalization of the dark sector:

$$\Rightarrow \langle \Gamma(\nu\nu \rightarrow Z') \rangle \gtrsim H(T = m_{Z'}/3)$$



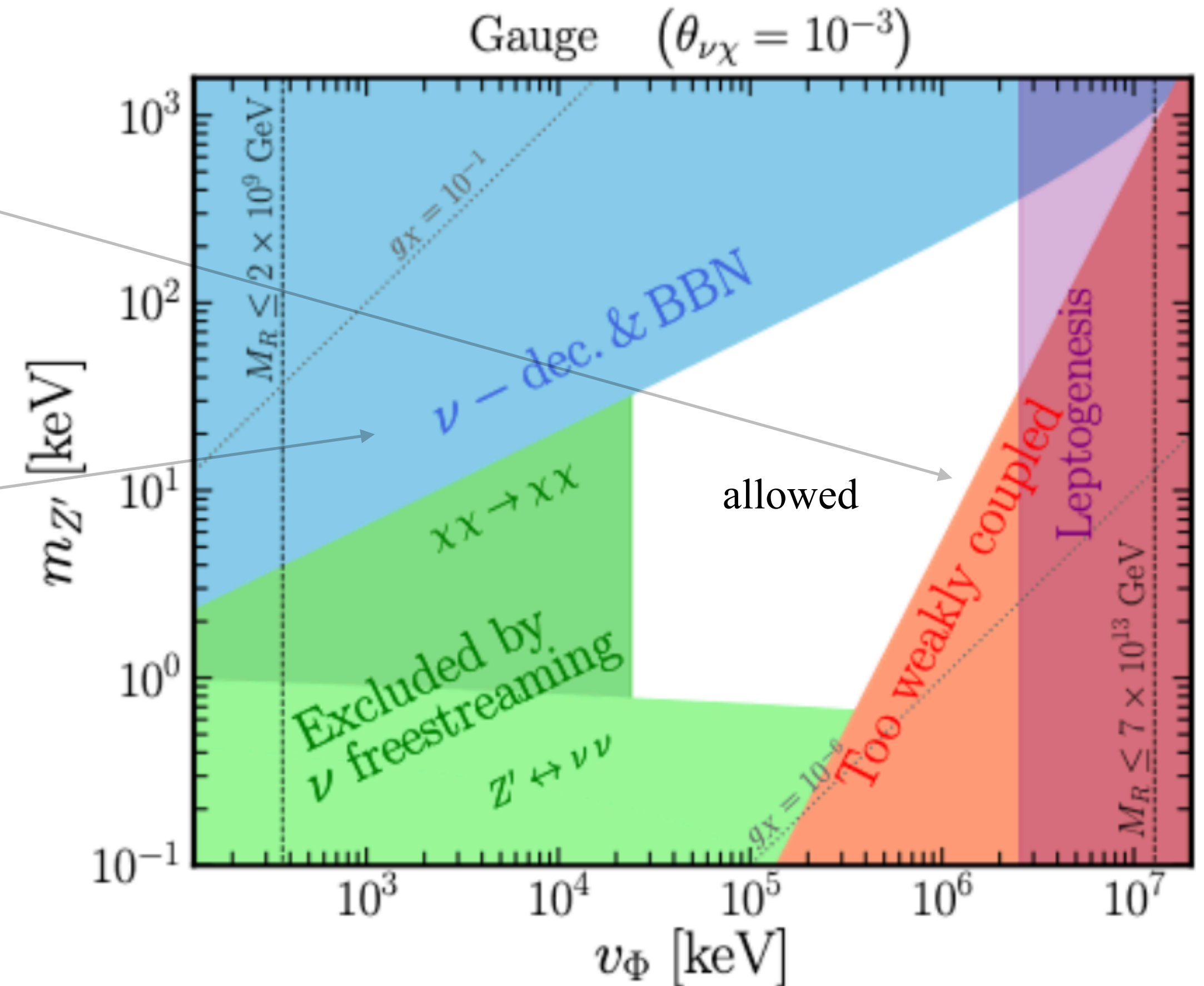
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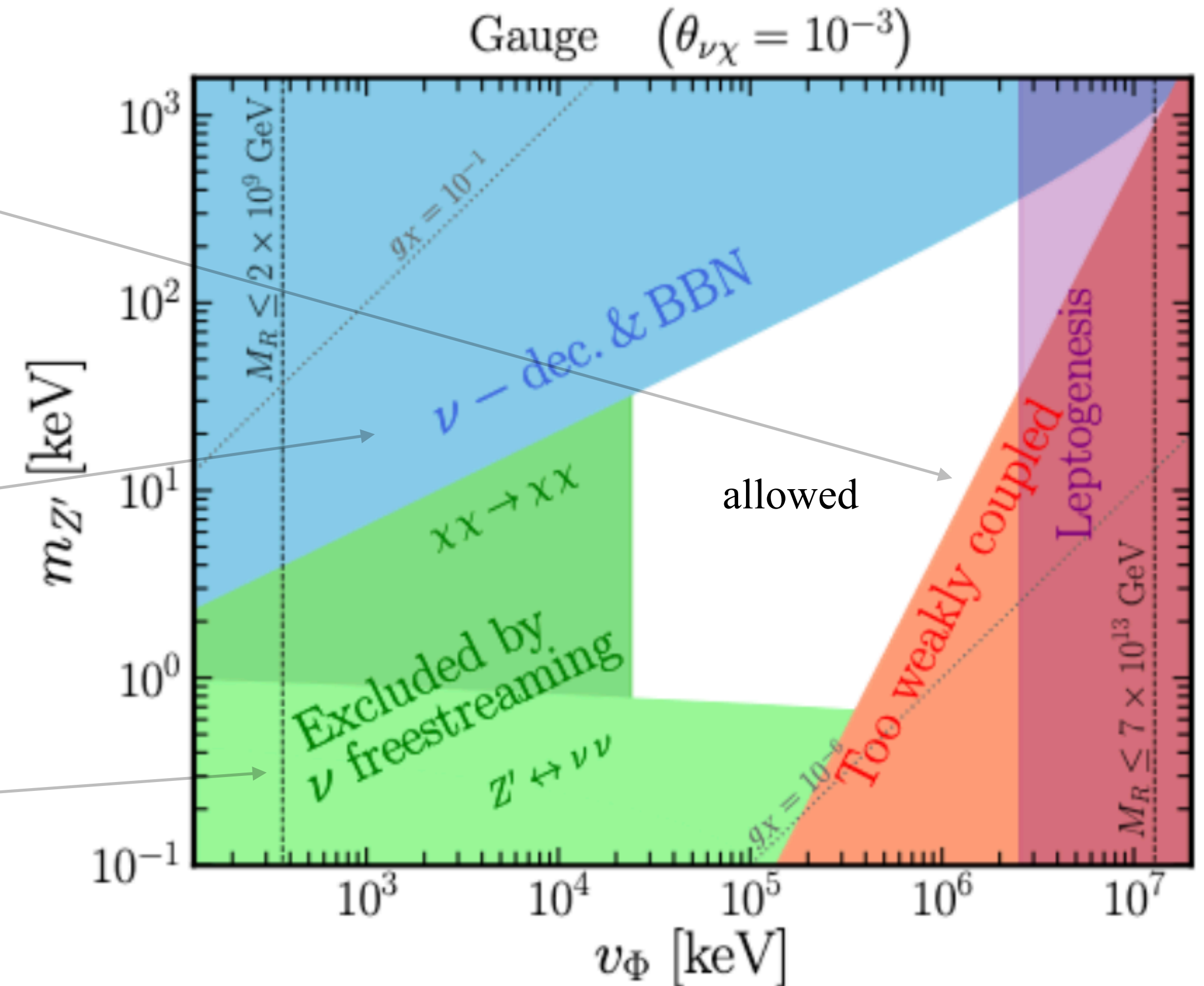
- avoid thermalization of the dark sector before BBN:

$$\langle \Gamma(\nu\nu \rightarrow Z') \rangle < H(T = 0.7 \text{ MeV})$$

- free-streaming of neutrinos & dark radiation before/around recombination

$$\langle \Gamma \rangle < H \text{ for } z < 10^5$$

Taule, Escudero, Garny, 2207.04062



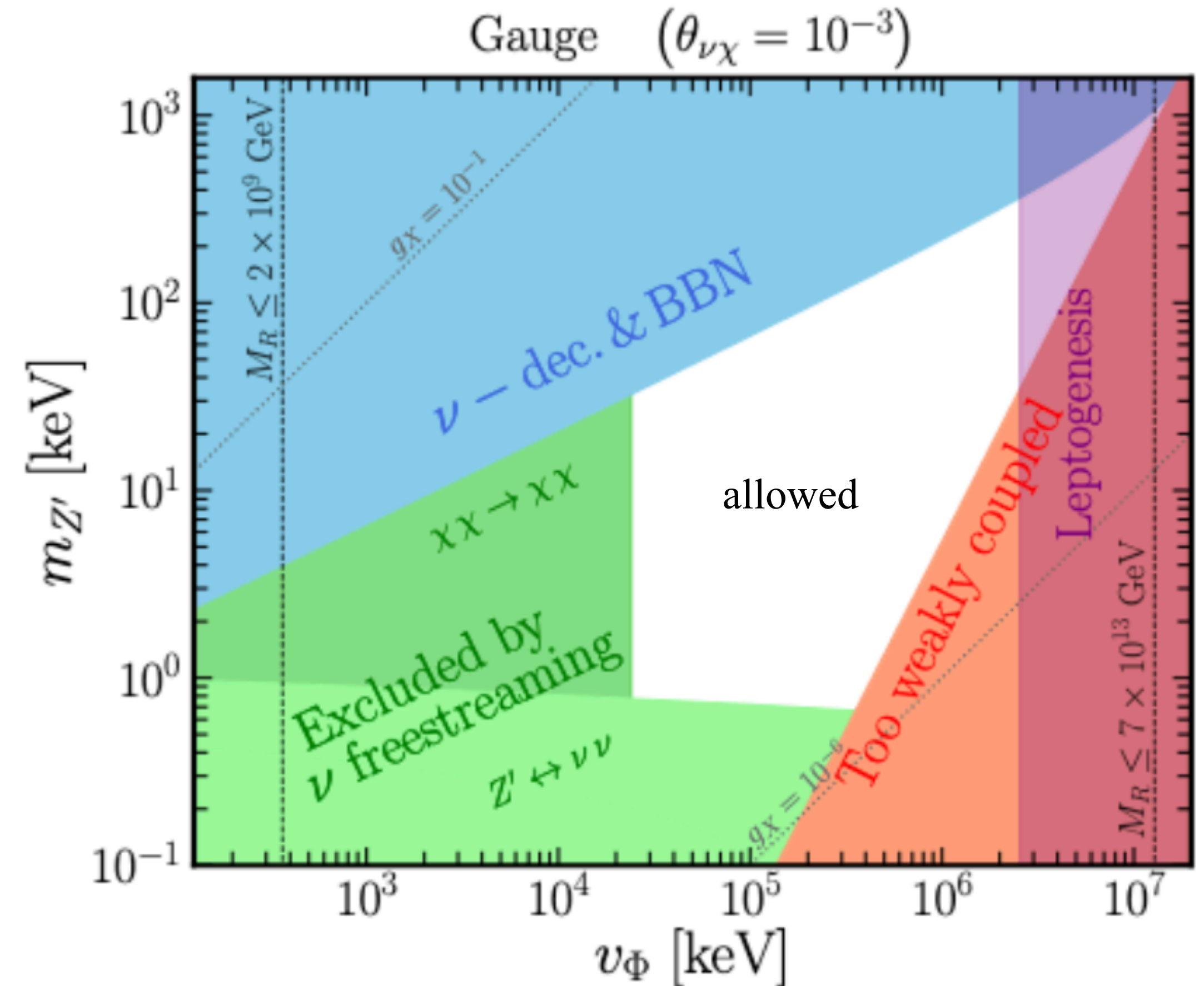


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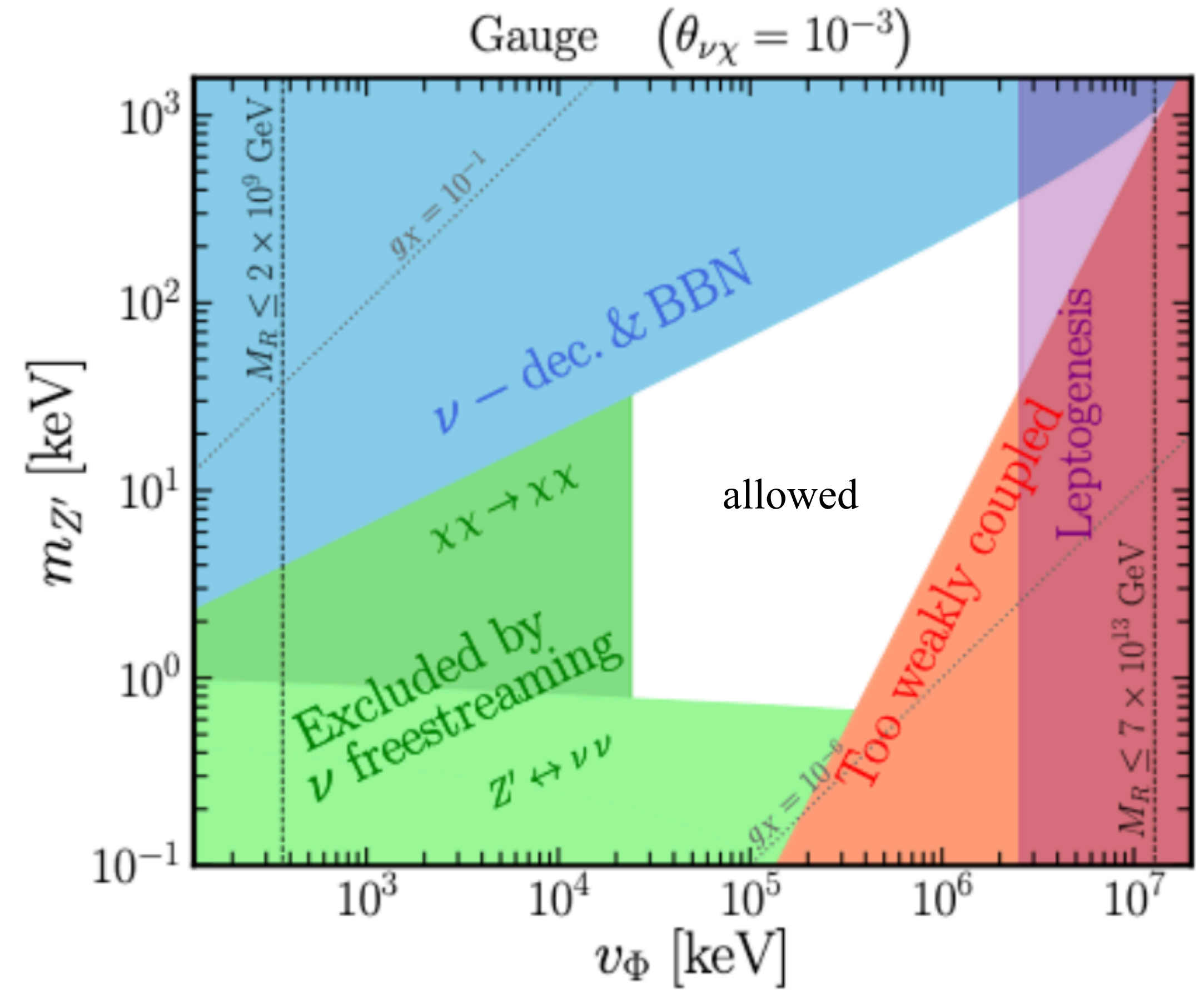
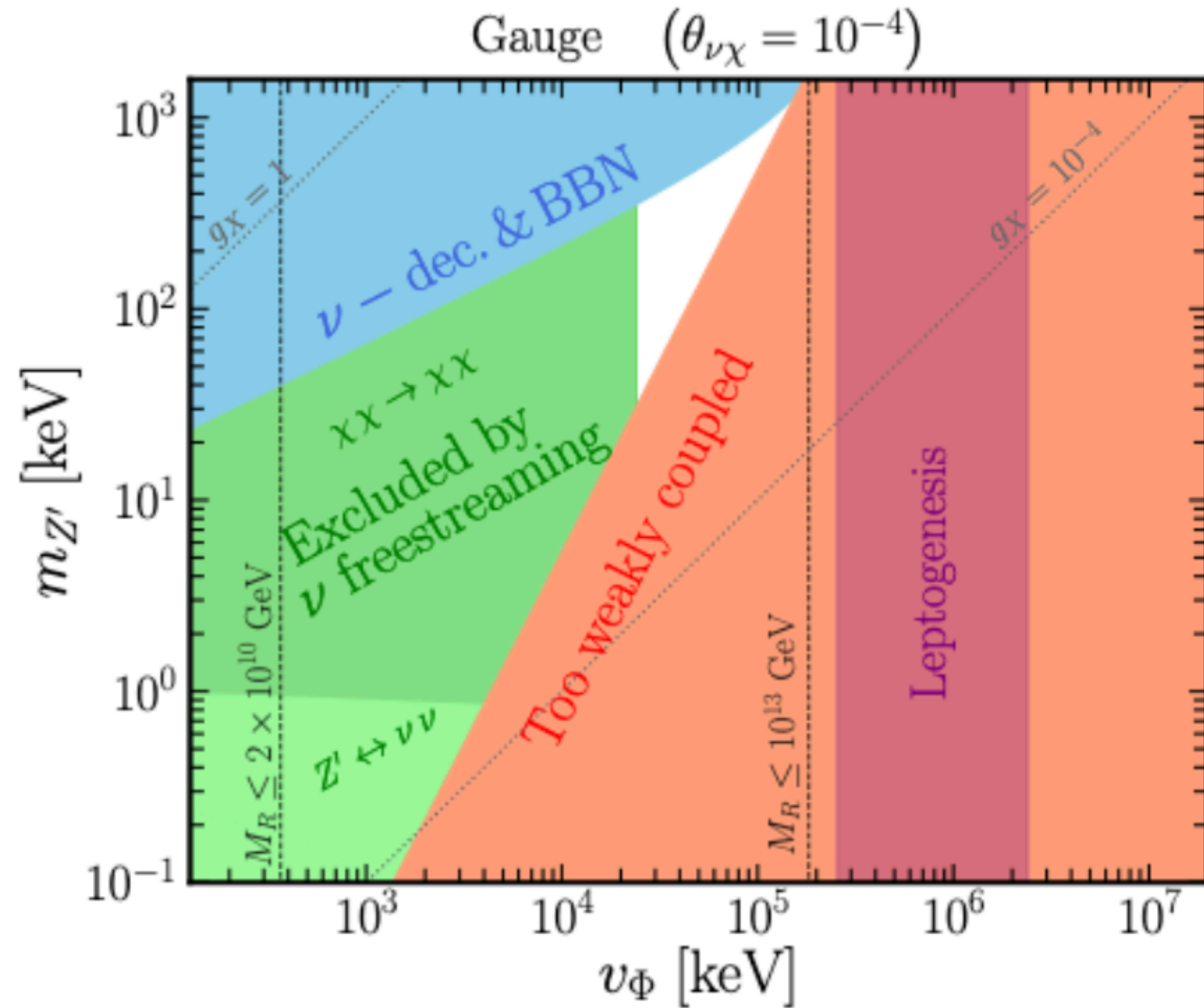
- avoid thermalization of  $\chi$  prior neutrino decoupling due to oscillations

$$|\theta_{\nu\chi}| \lesssim 10^{-3} \sqrt{\frac{10}{N_\chi}} \sqrt{\frac{0.2 \text{ eV}}{m_\nu}}$$

too small to be tested in SBL oscillation experiments



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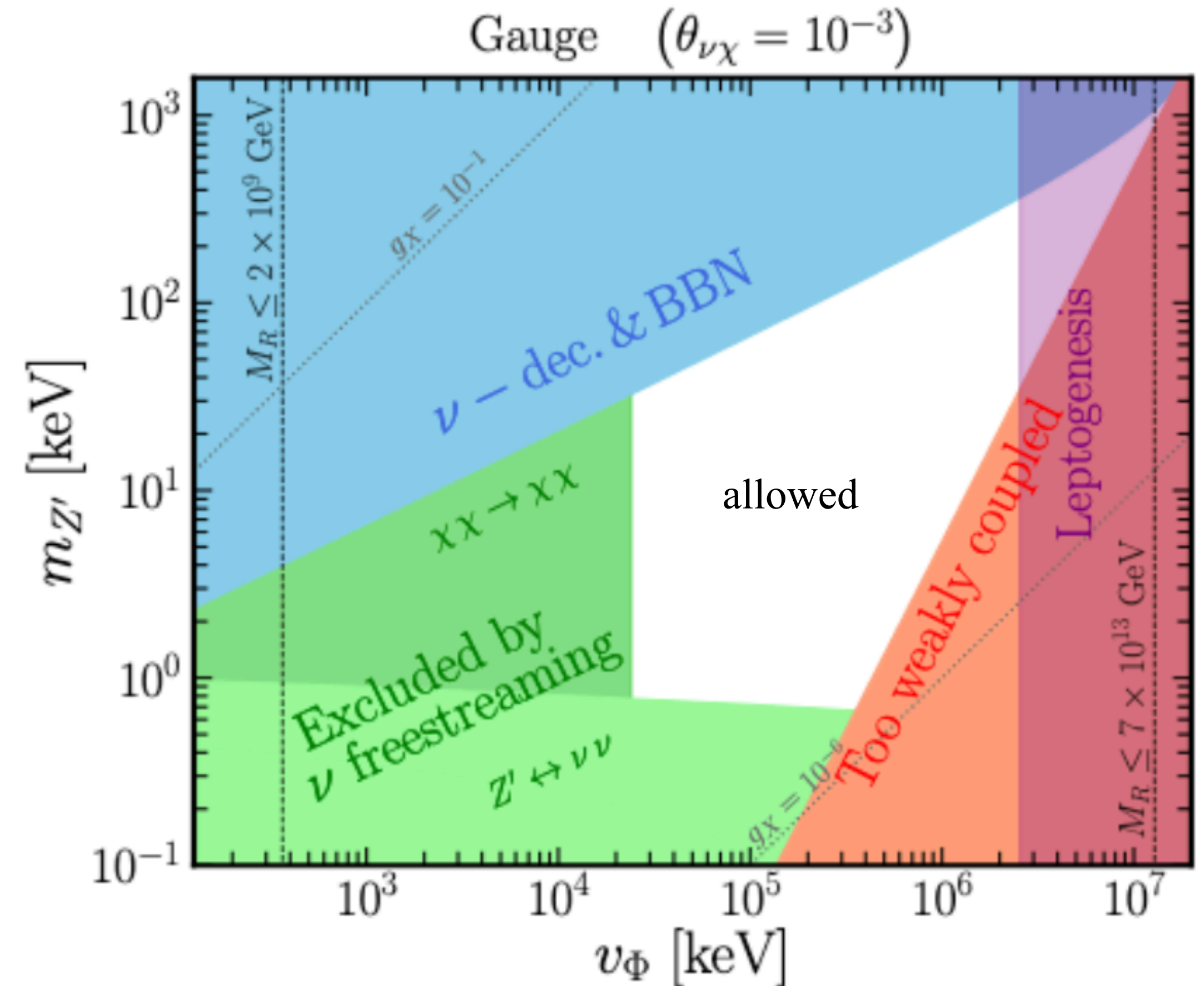


# Available parameter space

- constraints on heavy RH neutrinos:

$$M_R \lesssim 10^{10} - 10^{14} \text{ GeV}$$

- perturbativity of Yukawa  $Y_\Phi \bar{N}_R \chi_L \Phi$
- loop-induced Higgs portal  $\lambda_{\Phi H} |\Phi|^2 H^\dagger H$  remains small enough to avoid thermalization of  $\Phi$  prior BBN



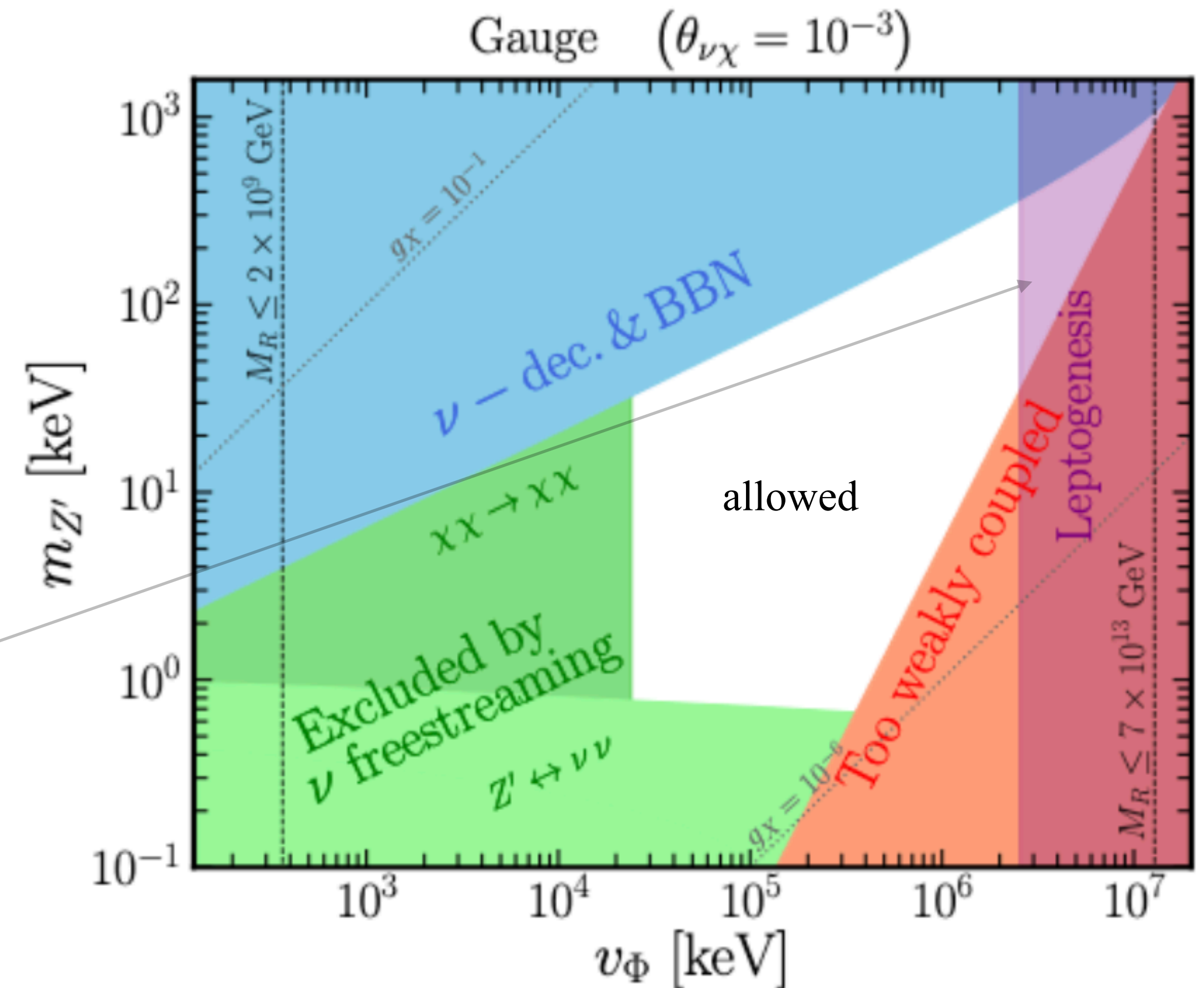


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- standard thermal leptogenesis works if  $N \rightarrow HL$  dominates over  $N \rightarrow \chi\Phi$
- otherwise  $\chi$  would thermalize and conflict with  $N_{\text{eff}}$  during BBN  $\Rightarrow$  require  $T_{RH} < M_R$  (allows still for  $T_{RH} \gg T_{EW}$ )





# Further signatures of the model

- SN cooling arguments for SN1987A exclude

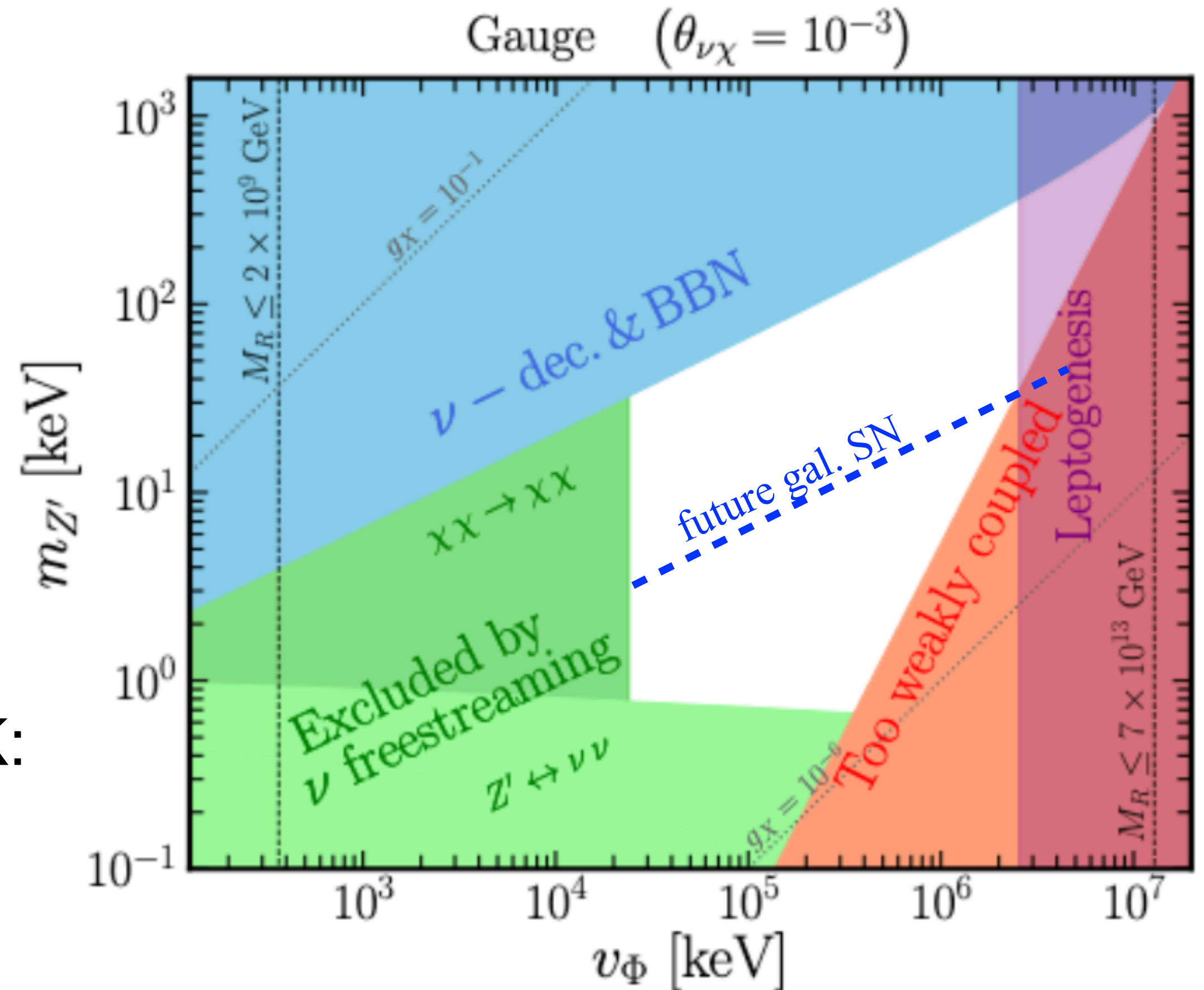
$$3 \times 10^{-7} \frac{\text{keV}}{m_{Z'}} \lesssim \lambda_{Z'}^{\nu\nu} \lesssim 10^{-4} \frac{\text{keV}}{m_{Z'}} \quad \text{Fiorillo, Raffelt, Vitagliano, 2209.11773}$$

weaker than BBN constraint

$$\lambda_{Z'}^{\nu\nu} \lesssim 10^{-7} (\text{keV}/m_{Z'})$$

- Future galactic SN at 10 kpc detected by HyperK: sensitivity down to

$$\lambda_{Z'}^{\nu\nu} \sim 10^{-9} (\text{keV}/m_{Z'}) \quad \text{Akita, Im, Masud, 2206.06852}$$



# Summary

- Exciting interplay of cosmology and terrestrial neutrino mass determinations
- Cosmological bounds reaching minimal values required by oscillations
- Relaxing cosmo bound requires new physics
- Presented simple seesaw model:
  - large number of massless sterile neutrinos ( $N_\chi \gtrsim 10 - 30$ )
  - dark U(1) symmetry with breaking scale between 10 MeV and 10 GeV
  - weakly coupled  $Z'$  with mass 1 — 100 keV with  $\lambda_{Z'}^{\nu\nu} \sim 10^{-9}$

# Complementarity between mass determinations from heaven and earth

link between neutrino mass observables *in the standard scenario*:

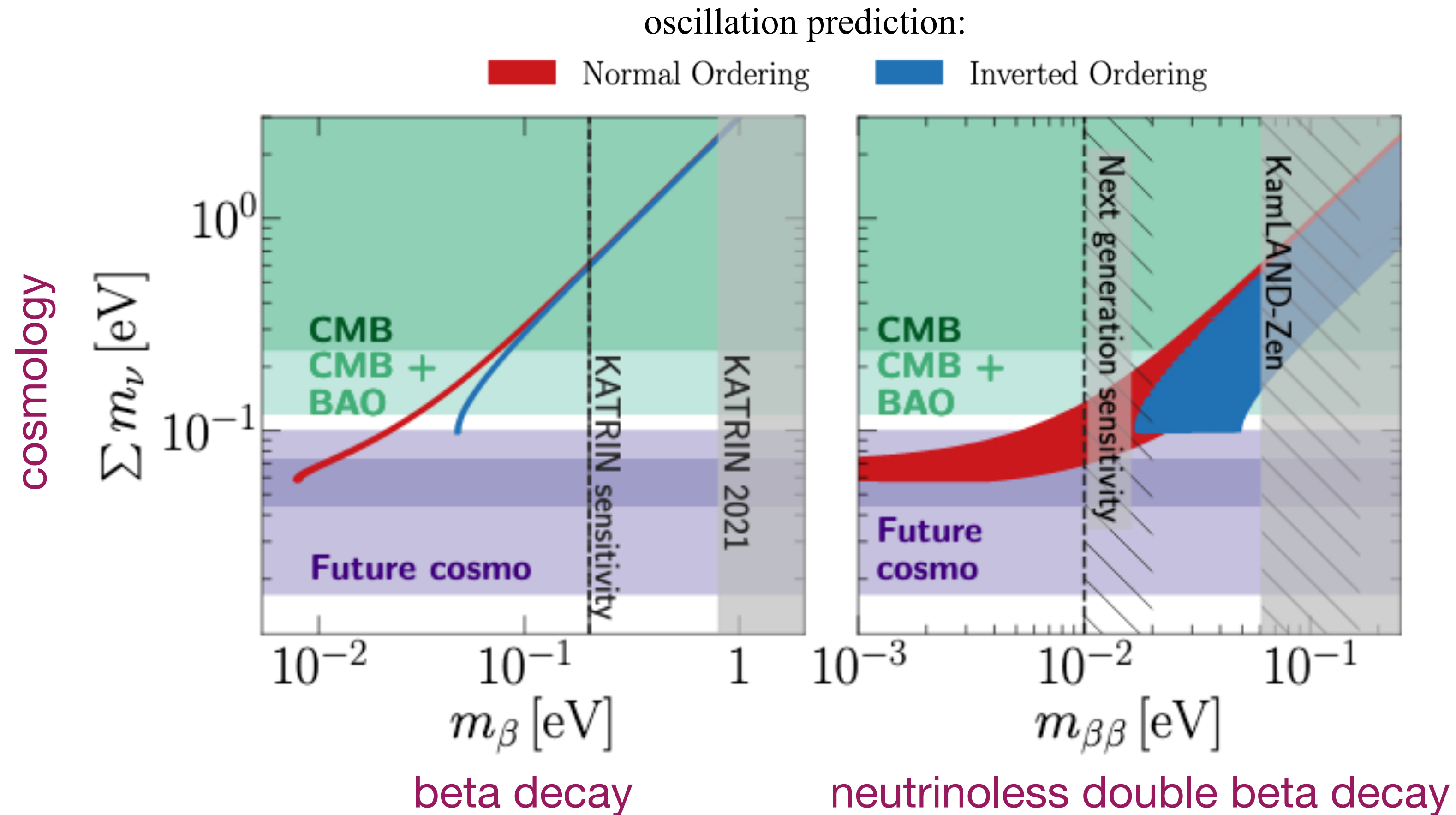


fig. by I. Esteban  
based on NuFit 5.0

# Counting the number of neutrino flavours

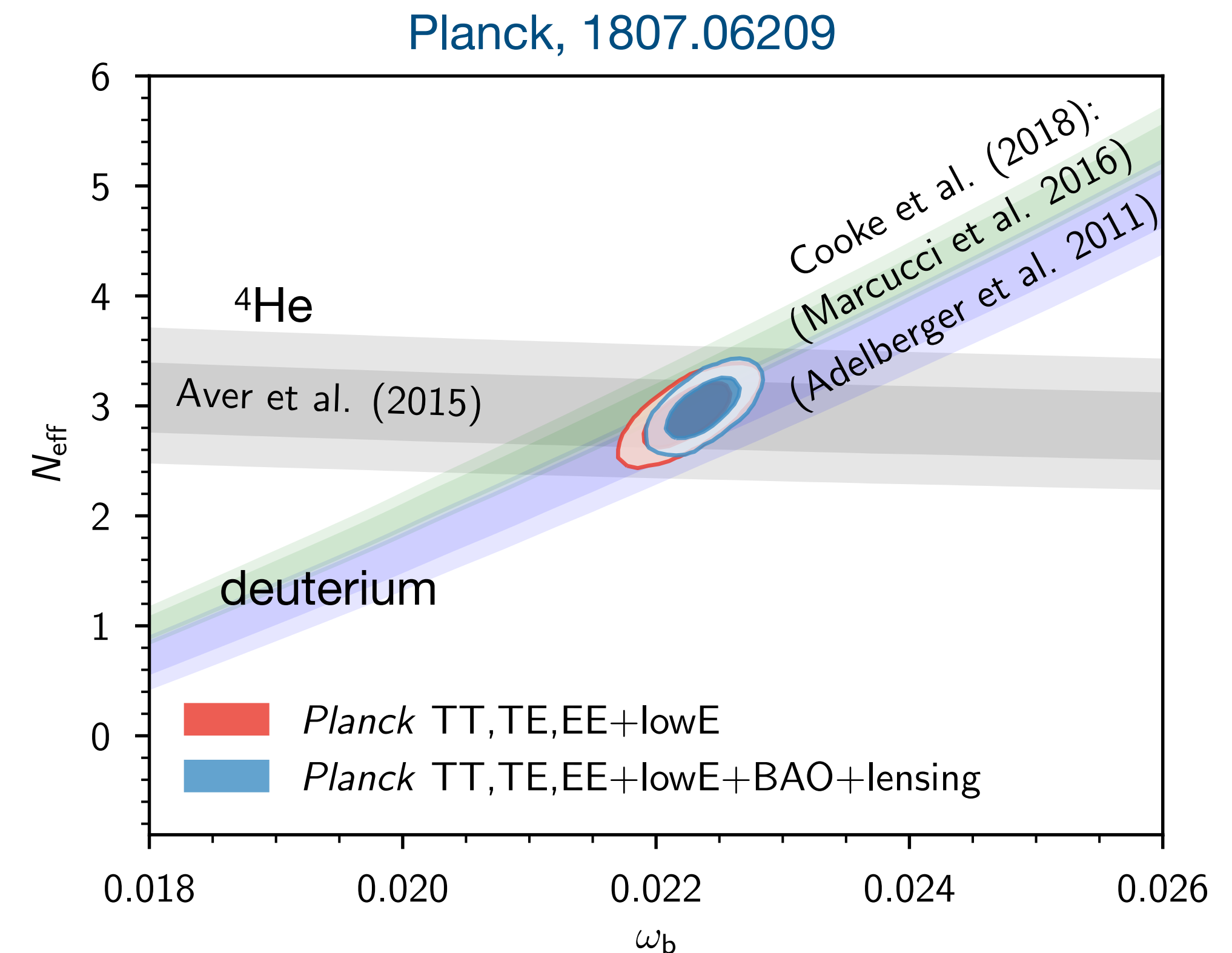
$N_{\text{eff}}$  affects

- formation of light elements (BBN),  
 $T \sim \text{MeV}$ ,  $t \sim 1 \text{ min}$

$$N_{\text{eff}} = 2.78 \pm 0.28 \text{ (68\% CL)}$$

- CMB decoupling,  $T \sim \text{eV}$ ,  $t \sim 400\,000 \text{ yr}$

$$N_{\text{eff}} = 2.99 \pm 0.17 \text{ (68\% CL)}$$





# Neutrino mass from cosmology

$$\sum m_\nu < 0.24 \text{ eV (CMB)}$$

$$\sum m_\nu < 0.12 \text{ eV (CMB+BAO)}$$

limits at 95% CL

Planck 1807.06209

