

# Testing the Electroweak Theory in Multiboson Measurements in ATLAS

(ATLAS public results from 2023)

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On behalf of ATLAS Collaboration

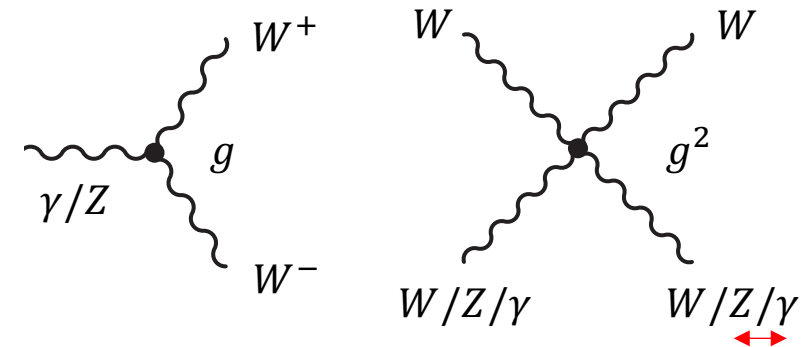
Lake Louise Winter Institute

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# Electroweak Multiboson Interactions

□ The Standard Model (SM) of elementary particles predicts triple and quartic gauge couplings between the electroweak bosons due to the non-Abelian structure of the electroweak interaction

- ❖ Experimental studies of multiboson interactions are therefore important tests of the SM electroweak theory
- ❖ Effective Field Theory (EFT) frameworks are typically used in model independent searches for new physics effects



□ According to the SM, the massive electroweak bosons obtain their masses and hence, can have longitudinal polarisations via the Higgs mechanism of the spontaneously broken electroweak symmetry

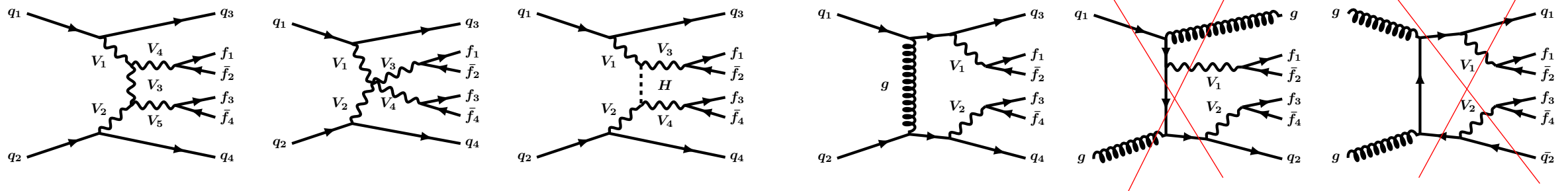
- ❖ Measurements of polarisation observables in the multiboson interactions are the direct probes of this mechanism

□ The latest measurements and first observations of the production of different multiboson final states in the ATLAS detector are presented

# Same Sign $W^\pm W^\pm jj$ Measurement

$W^\pm W^\pm jj$  has the largest ratio of the electroweak to QCD production cross sections among all vector boson scattering (VBS) sensitive  $VVjj$  final states

As the QCD leading order diagrams with initial gluons are forbidden

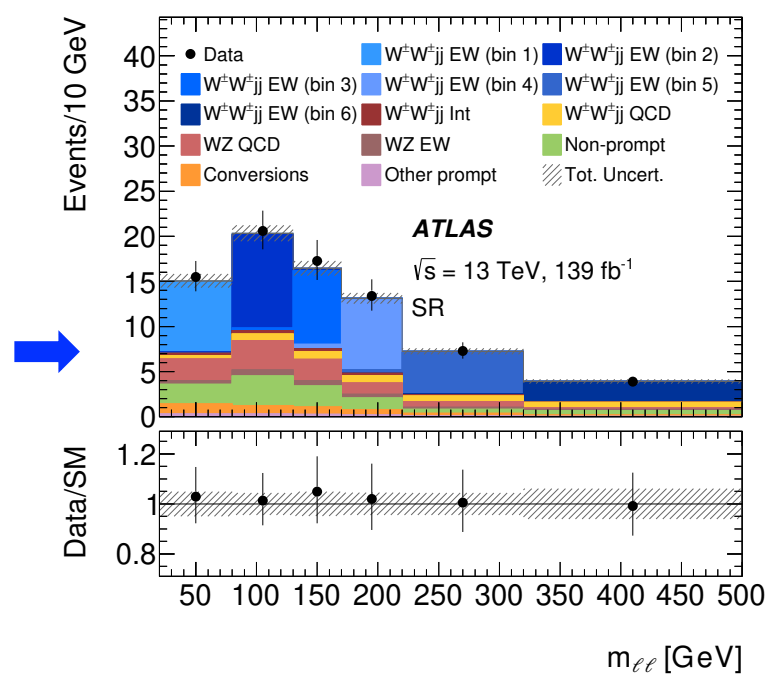


Fiducial and differential cross sections as functions of different observables are measured for both electroweak and inclusive (EW+QCD)  $W^\pm W^\pm jj$  production

Good agreement is found for the fiducial cross sections with the SM predictions

Post-fit SR distribution of event yields from differential cross section extraction as a function of  $m_{ll}$

Description	$\sigma_{\text{fid}}^{\text{EW}}$ [fb]	$\sigma_{\text{fid}}^{\text{EW+Int+QCD}}$ [fb]
Measured cross section	$2.92 \pm 0.22$ (stat.) $\pm 0.19$ (syst.)	$3.38 \pm 0.22$ (stat.) $\pm 0.19$ (syst.)
MG5_AMC+HERWIG7	$2.53 \pm 0.04$ (PDF) $^{+0.22}_{-0.19}$ (scale)	$2.92 \pm 0.05$ (PDF) $^{+0.34}_{-0.27}$ (scale)
MG5_AMC+PYTHIA8	$2.53 \pm 0.04$ (PDF) $^{+0.22}_{-0.19}$ (scale)	$2.90 \pm 0.05$ (PDF) $^{+0.33}_{-0.26}$ (scale)
SHERPA	$2.48 \pm 0.04$ (PDF) $^{+0.40}_{-0.27}$ (scale)	$2.92 \pm 0.03$ (PDF) $^{+0.60}_{-0.40}$ (scale)
SHERPA $\otimes$ NLO EW	$2.10 \pm 0.03$ (PDF) $^{+0.34}_{-0.23}$ (scale)	$2.54 \pm 0.03$ (PDF) $^{+0.50}_{-0.33}$ (scale)
POWHEG BOX+PYTHIA	2.64	-



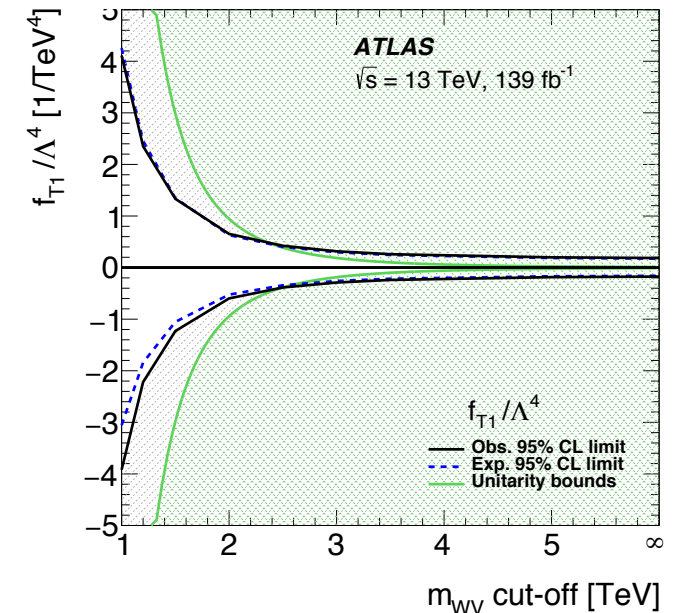
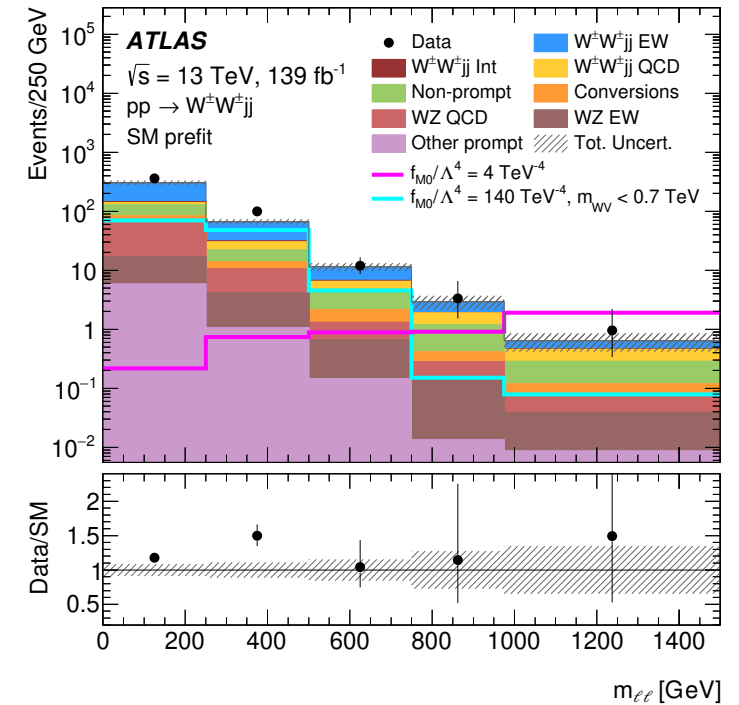
# Same Sign $W^\pm W^\pm jj$ Measurement

Competitive limits (@ 95% C.L.) are set on the Wilson coefficients of the relevant EFT dimension-8 operators that have large effects on the  $WWWW$  coupling

The reconstructed  $m_{ll}$  distribution is used in the EFT measurements

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{f_i^{(6)}}{\Lambda^2} O_i^{(6)} + \sum_j \frac{f_j^{(8)}}{\Lambda^4} O_j^{(8)}$$

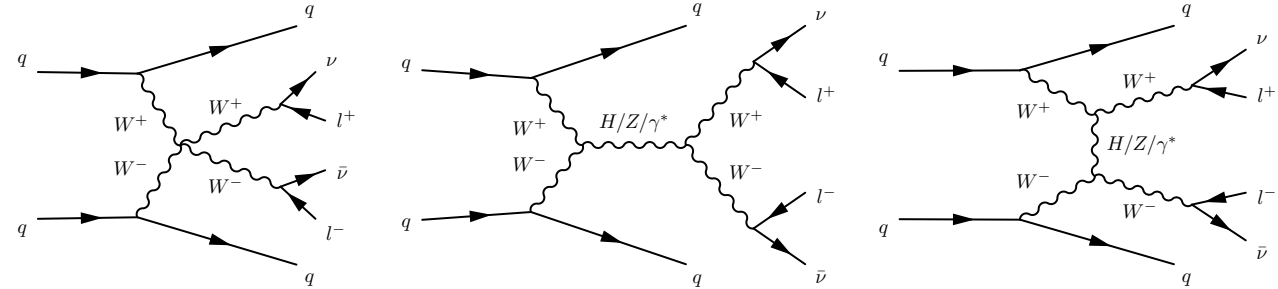
Coefficient	Type	No unitarisation cut-off [TeV <sup>-4</sup> ]	Lower, upper limit at the respective unitarity bound [TeV <sup>-4</sup> ]
$f_{M0}/\Lambda^4$	Exp.	[-3.9, 3.8]	-64 at 0.9 TeV, 40 at 1.0 TeV
	Obs.	[-4.1, 4.1]	-140 at 0.7 TeV, 117 at 0.8 TeV
$f_{M1}/\Lambda^4$	Exp.	[-6.3, 6.6]	-25.5 at 1.6 TeV, 31 at 1.5 TeV
	Obs.	[-6.8, 7.0]	-45 at 1.4 TeV, 54 at 1.3 TeV
$f_{M7}/\Lambda^4$	Exp.	[-9.3, 8.8]	-33 at 1.8 TeV, 29.1 at 1.8 TeV
	Obs.	[-9.8, 9.5]	-39 at 1.7 TeV, 42 at 1.7 TeV
$f_{S02}/\Lambda^4$	Exp.	[-5.5, 5.7]	-94 at 0.8 TeV, 122 at 0.7 TeV
	Obs.	[-5.9, 5.9]	-
$f_{S1}/\Lambda^4$	Exp.	[-22.0, 22.5]	-
	Obs.	[-23.5, 23.6]	-
$f_{T0}/\Lambda^4$	Exp.	[-0.34, 0.34]	-3.2 at 1.2 TeV, 4.9 at 1.1 TeV
	Obs.	[-0.36, 0.36]	-7.4 at 1.0 TeV, 12.4 at 0.9 TeV
$f_{T1}/\Lambda^4$	Exp.	[-0.158, 0.174]	-0.32 at 2.6 TeV, 0.44 at 2.4 TeV
	Obs.	[-0.174, 0.186]	-0.38 at 2.5 TeV, 0.49 at 2.4 TeV
$f_{T2}/\Lambda^4$	Exp.	[-0.56, 0.70]	-2.60 at 1.7 TeV, 10.3 at 1.2 TeV
	Obs.	[-0.63, 0.74]	-



# Observation of Opposite Sign $W^+W^-jj$

- ATLAS observed the electroweak VBS  $W^+W^-jj$  production in fully leptonic final states

Leptons are required to have different flavours



- Top quark (mainly the  $t\bar{t}$ ) along with QCD  $W^+W^-jj$  production make huge background to the signal

66% and 24% contributions to the total (post-fit) event prediction in the inclusive signal region, respectively



Process	Event yields	
	$n_{\text{jets}} = 2$	$n_{\text{jets}} = 3$
EWK $W^+W^-jj$	$158 \pm 27$	$54 \pm 13$
Top quark	$2885 \pm 214$	$1851 \pm 131$
Strong $W^+W^-jj$	$1214 \pm 256$	$514 \pm 121$
$W$ +jets	$37 \pm 97$	$19 \pm 48$
$Z$ +jets	$216 \pm 62$	$65 \pm 25$
Multiboson	$101 \pm 5$	$42 \pm 3$
SM prediction	$4610 \pm 77$	$2546 \pm 48$
Data	4610	2533

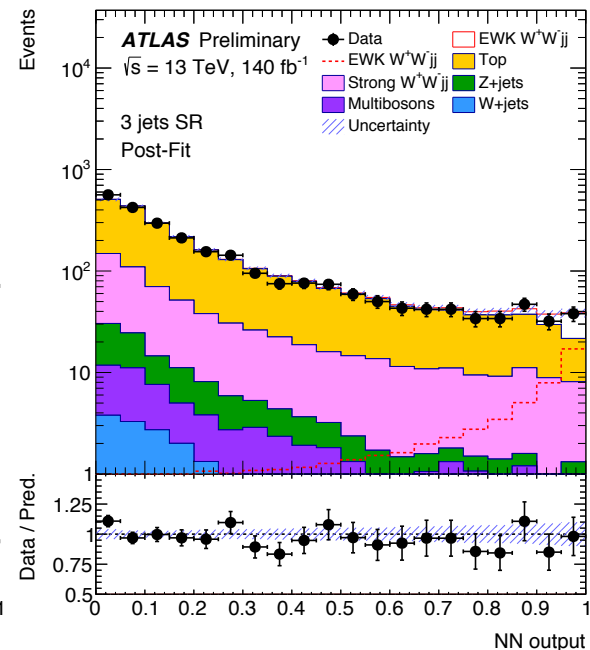
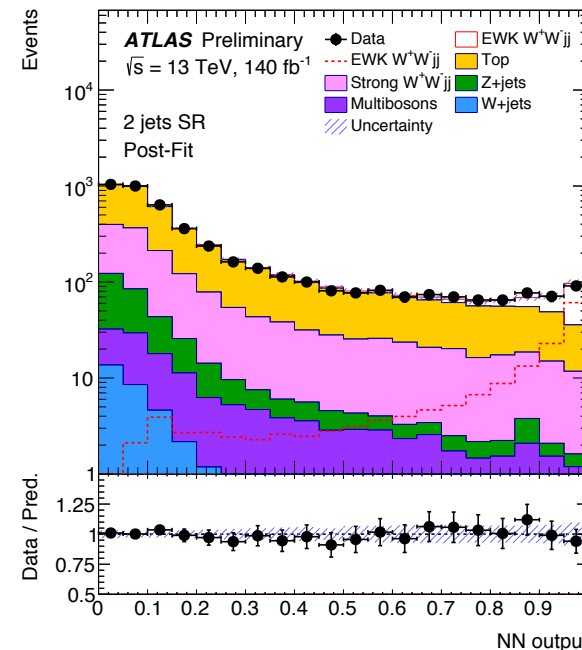
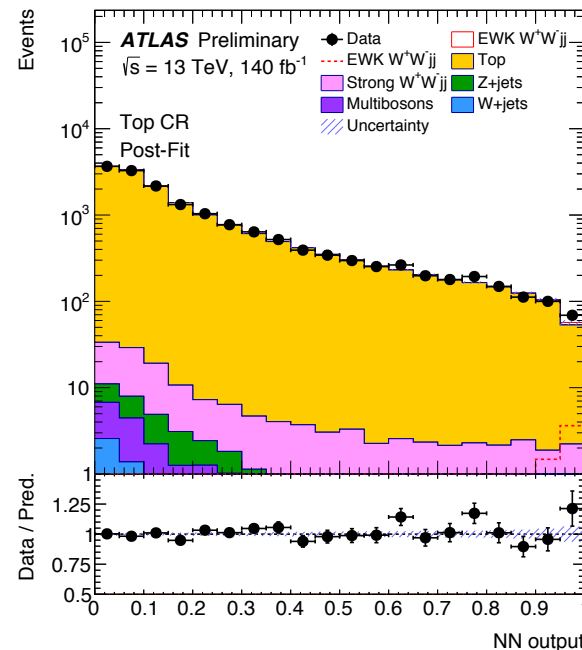
- Signal region is split into the exclusive 2- and 3-jet event categories to enhance the sensitivity

- Control region for the top quark background is defined by requiring one of the two leading jets to be b-tagged

# Observation of Opposite Sign $W^+W^-jj$

- Neural Network based discriminant is used to distinguish the signal from background
  - ❖ Signal, top quark and QCD background events are used in the NN training
- Profile-likelihood fit method is used to fit simultaneously the signal, top and QCD background normalisations in the NN output in the 1 control and 2 signal regions

- Observed (expected) signal significance is  $7.1\sigma$  ( $6.2\sigma$ )
  - ❖ Statistical uncertainty of the measured signal normalisation is  $12.3\%$  with  $18.5\%$  total uncertainty



- Signal fiducial cross section is measured to  $2.65_{-0.48}^{+0.52} fb$  vs. predicted  $2.20_{-0.13}^{+0.14} fb$ 
  - ❖ Fiducial volume defined closely to detector level selection but requiring  $m_{jj} > 500 GeV$

# Opposite Sign $W^+W^-$ Cross-Sections

$\sqrt{s} = 13 \text{ TeV}, 140 \text{ fb}^{-1}$

□ Fiducial and differential cross sections are measured in  $W^+W^- \rightarrow e^\pm \nu \mu^\mp \bar{\nu}$  final states

❖ The fiducial cross section is extrapolated to the full phase-space of  $W^+W^-$  production

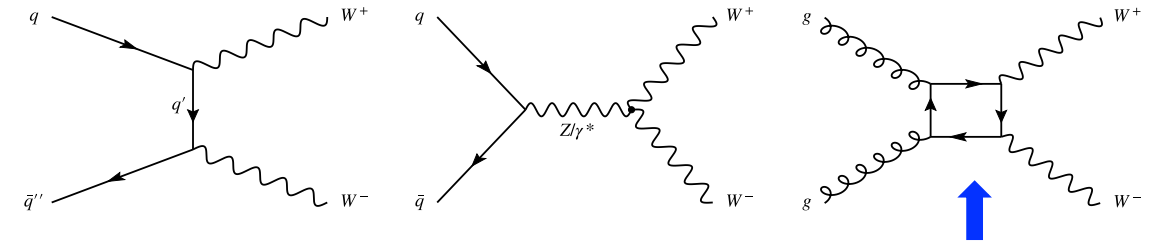
□ Top-quark background is precisely estimated in bins of the signal region with the help of two dedicated control regions

❖ Defined by requiring exactly 1 and exactly 2 b-jets, respectively

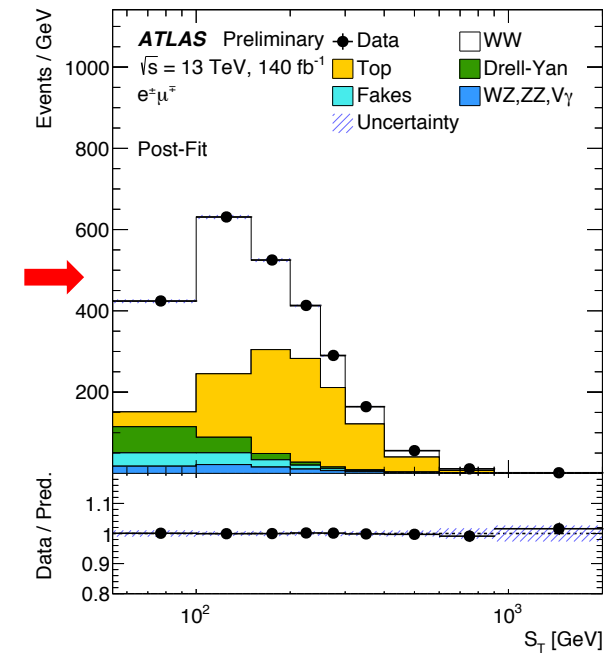
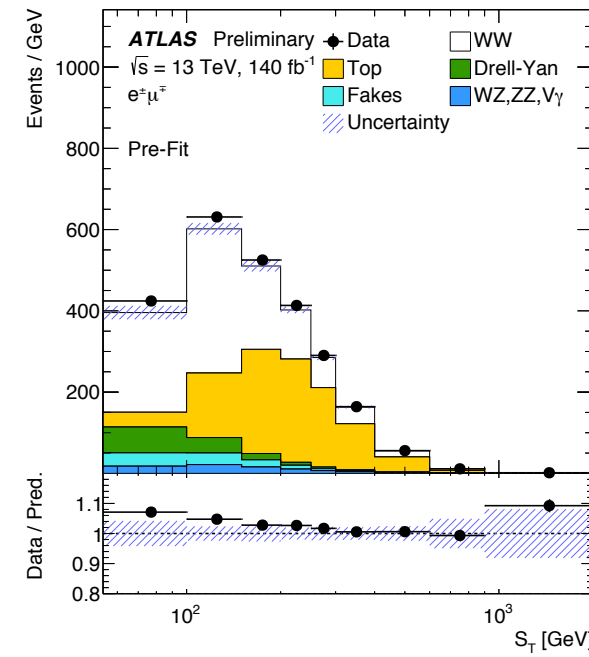
❖ No constraints on the jet multiplicity

□ Fiducial cross section is measured using profile-likelihood fit with the free signal normalisation to the detector-level  $S_T$  distribution

❖  $S_T$  - scalar sum of all jet and lepton transverse momenta



Provides ~5% contribution in the  $WW$  production rate

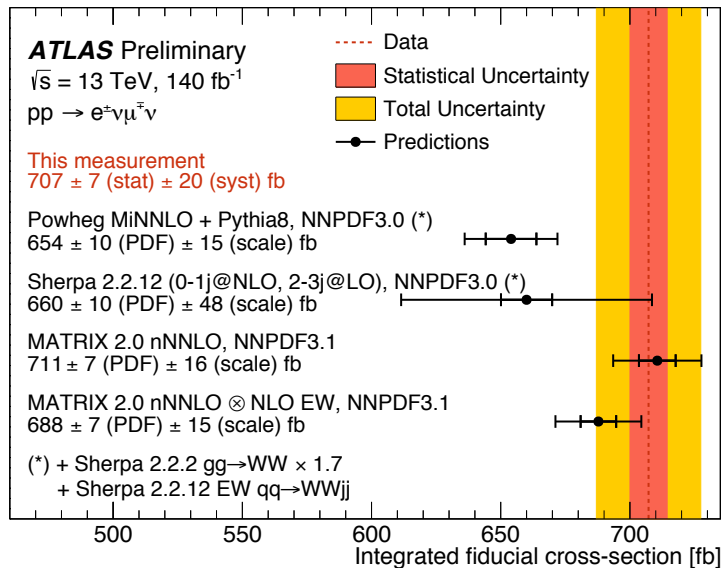


# Opposite Sign $W^+W^-$ Cross-Sections

□ Fiducial cross section is measured with the 3.1% (!) total uncertainty

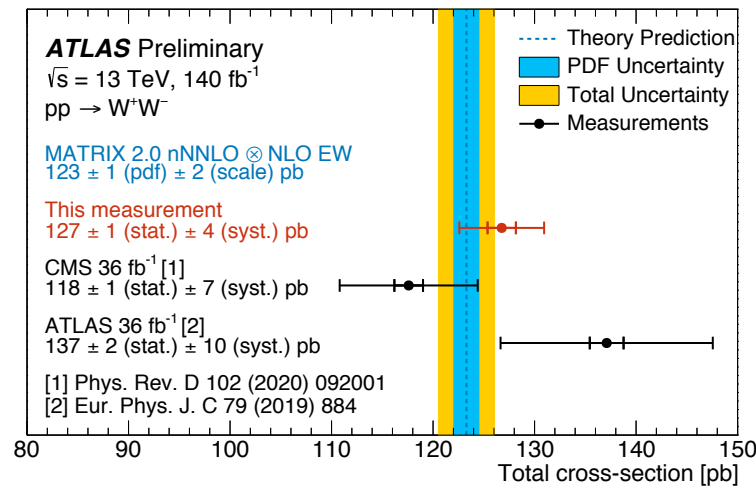
❖ Excellent agreement with the MATRIX 2.0.1 prediction at nNNLO in QCD NNLO qq → WW + NLO gg → WW

$$\sigma_{\text{fid}} = 707 \pm 7 \text{ (stat.)} \pm 20 \text{ (syst.) fb}$$



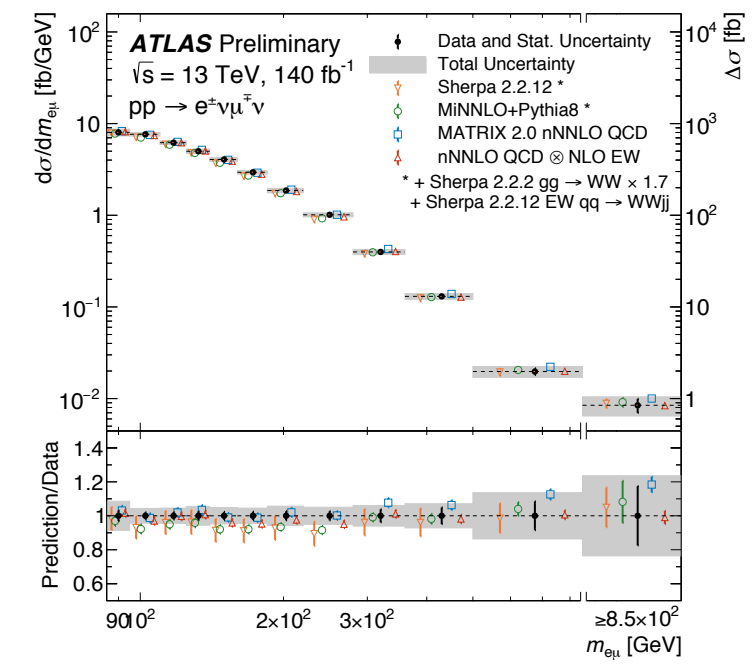
□ Total cross section of  $W^+W^-$  production is calculated using the acceptance of the  $W^+W^- \rightarrow e^\pm\nu\mu^\mp\nu$  events: 23.7% ± 0.3%, and the leptonic W branching ratio: 10.86%

$$\sigma_{\text{total}} = 127 \pm 1 \text{ (stat.)} \pm 4 \text{ (syst.) pb}$$



□ Differential cross sections are measured for **twelve** observables

❖ Iterative Bayesian unfolding used  
 ❖ Good agreement with predictions observed





# ZZ( $\rightarrow 4l$ )jj Differential Cross Sections

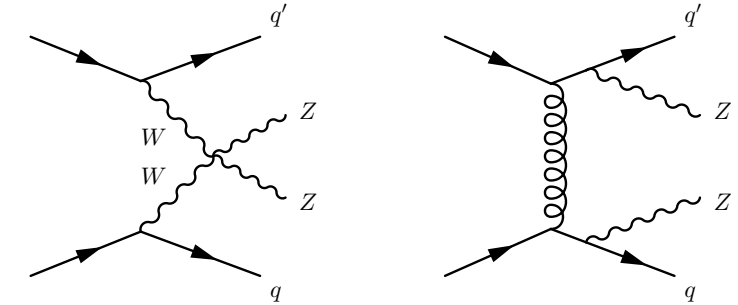
□ Differential cross sections are measured in VBS-enhanced ( $\zeta < 0.4$ ) and VBS-suppressed ( $\zeta > 0.4$ ) regions

❖ Three types of observables are measured

- VBS observables
- Polarisation, charge conjugation and parity observables
- QCD-sensitive observables

❖ Both EW and QCD production mechanisms are probed

❖ Centralities,  $m_{jj}$ ,  $|\Delta y_{jj}|$  and multiplicity of jets ( $n_{jets}^{gap}$ ) in between the two leading jets are most important observables in EW VBS measurements

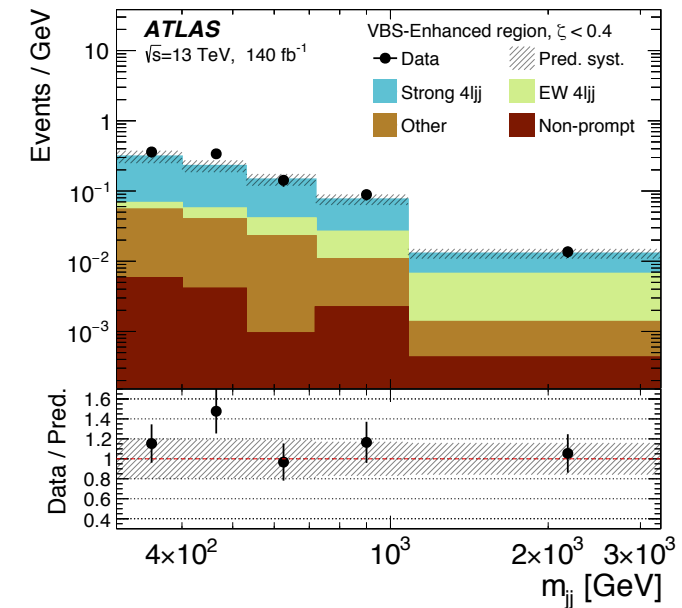


$$\zeta = \left| \frac{[y_{4\ell} - 0.5(y_{j_1} + y_{j_2})]}{\Delta y_{jj}} \right|$$

□ Two Z bosons are selected from the same-flavour opposite-charge lepton pairs

- ❖ Have smallest  $|m_{ll} - m_Z|$
- ❖ Are formed from different leptons

Process	Event yield $\pm$ stat. $\pm$ syst.	
	VBS-enhanced	VBS-suppressed
strong $4\ell jj$ (SHERPA)	$98.9 \pm 0.5 \pm 25.2$	$45.5 \pm 0.3 \pm 12.9$
EW $4\ell jj$ (MG5+PY8)	$24.1 \pm 0.1 \pm 1.8$	$2.12 \pm 0.02 \pm 0.14$
Prompt background	$18.8 \pm 0.2 \pm 2.2$	$5.5 \pm 0.1 \pm 0.4$
Non-prompt background	$3.0 \pm 0.6 \pm 3.2$	$1.1 \pm 0.5 \pm 1.2$
Total prediction	$144 \pm 1 \pm 26$	$54 \pm 1 \pm 13$
Data	169	53

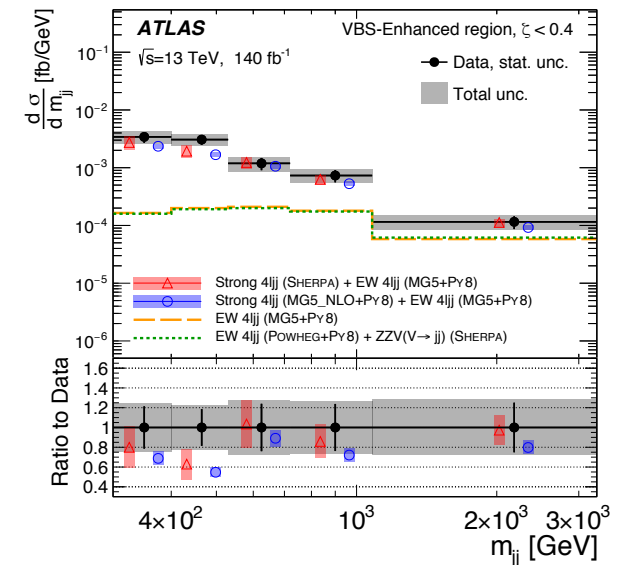
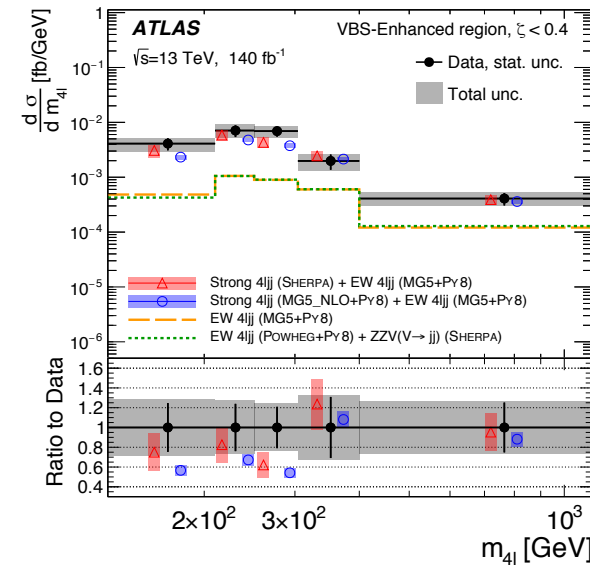


# ZZ( $\rightarrow 4l$ )jj Differential Cross Sections

Iterative Bayesian unfolding is used to measure differential cross sections

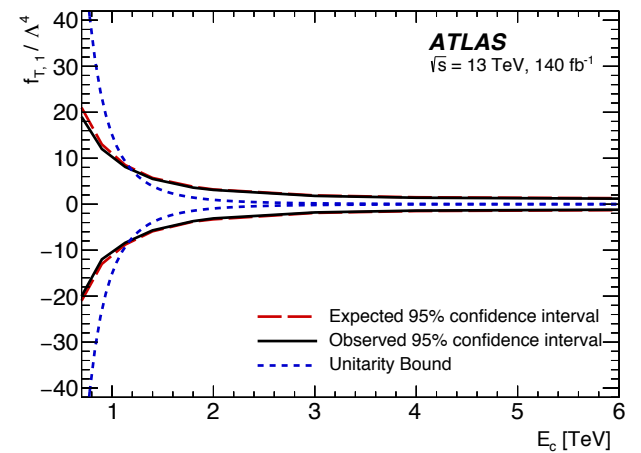
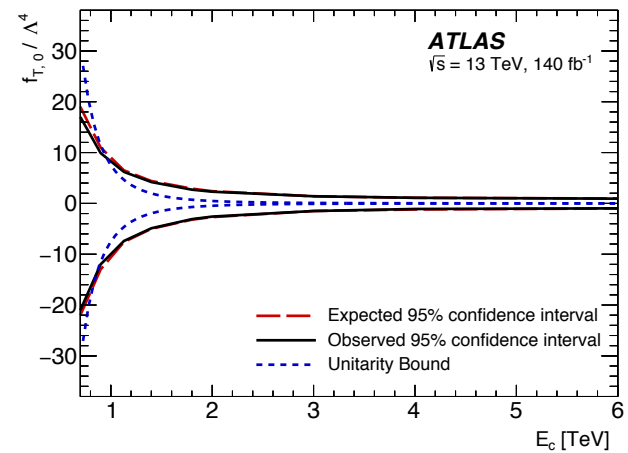
EFT samples combined with the SM signal are fitted to the unfolded  $m_{4l}$  and  $m_{jj}$  distributions and limits (@ 95% C.L.) on anomalous couplings of dimension-8 operators are obtained

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_i \frac{f_{T,i}}{\Lambda^4} \mathcal{O}_{T,i}$$



Much stronger limits are set on these three couplings by the same sign  $WWjj$  measurement (presented on the slide #4)

Wilson coefficient	$ \mathcal{M}_{d8} ^2$ Included	95% confidence interval [TeV $^{-4}$ ]	
		Expected	Observed
$f_{T,0}/\Lambda^4$	yes	[-0.98, 0.93]	[-1.00, 0.97]
	no	[-23, 17]	[-19, 19]
$f_{T,1}/\Lambda^4$	yes	[-1.2, 1.2]	[-1.3, 1.3]
	no	[-160, 120]	[-140, 140]
$f_{T,2}/\Lambda^4$	yes	[-2.5, 2.4]	[-2.6, 2.5]
	no	[-74, 56]	[-63, 62]
$f_{T,5}/\Lambda^4$	yes	[-2.5, 2.4]	[-2.6, 2.5]
	no	[-79, 60]	[-68, 67]
$f_{T,6}/\Lambda^4$	yes	[-3.9, 3.9]	[-4.1, 4.1]
	no	[-64, 48]	[-55, 54]
$f_{T,7}/\Lambda^4$	yes	[-8.5, 8.1]	[-8.8, 8.4]
	no	[-260, 200]	[-220, 220]
$f_{T,8}/\Lambda^4$	yes	[-2.1, 2.1]	[-2.2, 2.2]
	no	$[-4.6, 3.1] \times 10^4$	$[-3.9, 3.8] \times 10^4$
$f_{T,9}/\Lambda^4$	yes	[-4.5, 4.5]	[-4.7, 4.7]
	no	$[-7.5, 5.5] \times 10^4$	$[-6.4, 6.3] \times 10^4$



# $ZZ \rightarrow 4l$ Polarisation and CP Properties

Production of longitudinally polarised  $Z_L Z_L$  bosons is measured in  $ZZ \rightarrow l^+ l^- l'^+ l'^-$  final states with  $l, l' = e, \mu$

Profile-likelihood fit to the Boosted Decision Tree classifier is employed to measure the  $Z_L Z_L$  and the joined  $Z_L Z_T + Z_T Z_T$  polarisation components of the signal

The polarisation templates are obtained at the NLO accuracy in both QCD and EW

Advanced re-weighting scheme of the LO templates

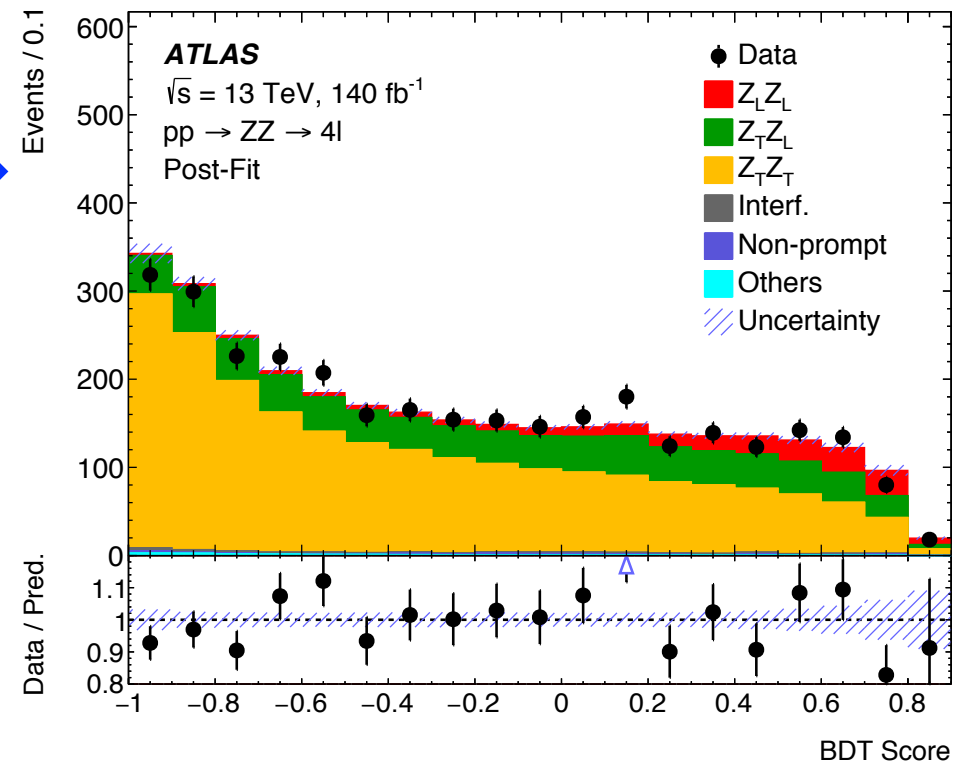
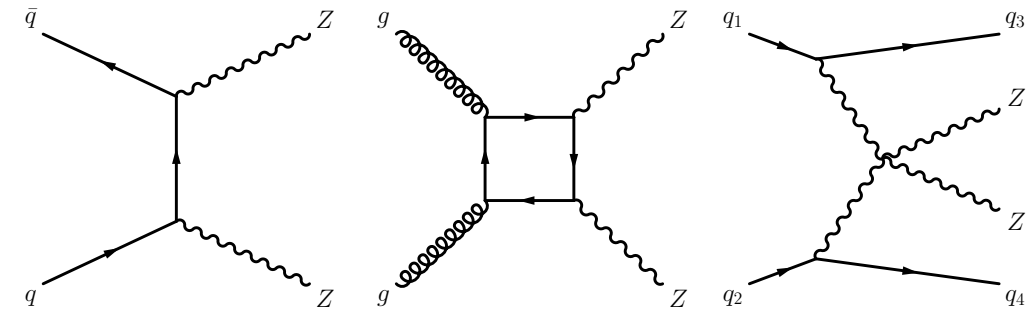
Observed significance of  $Z_L Z_L$  is  $4.3\sigma$

$3.8\sigma$  expected

Fiducial cross section of  $Z_L Z_L$  is  $2.45 \pm 0.60 \text{ fb}$

Consistent with the SM value of  $2.10 \pm 0.09 \text{ fb}$

Includes QCD and electroweak NLO corrections

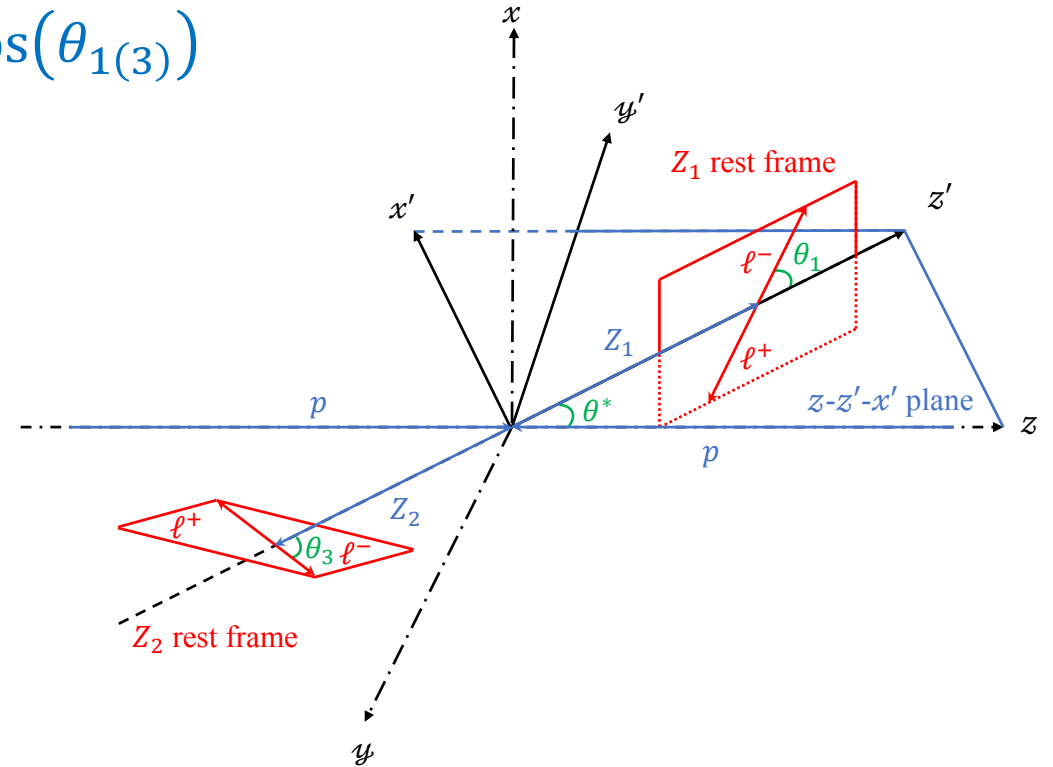
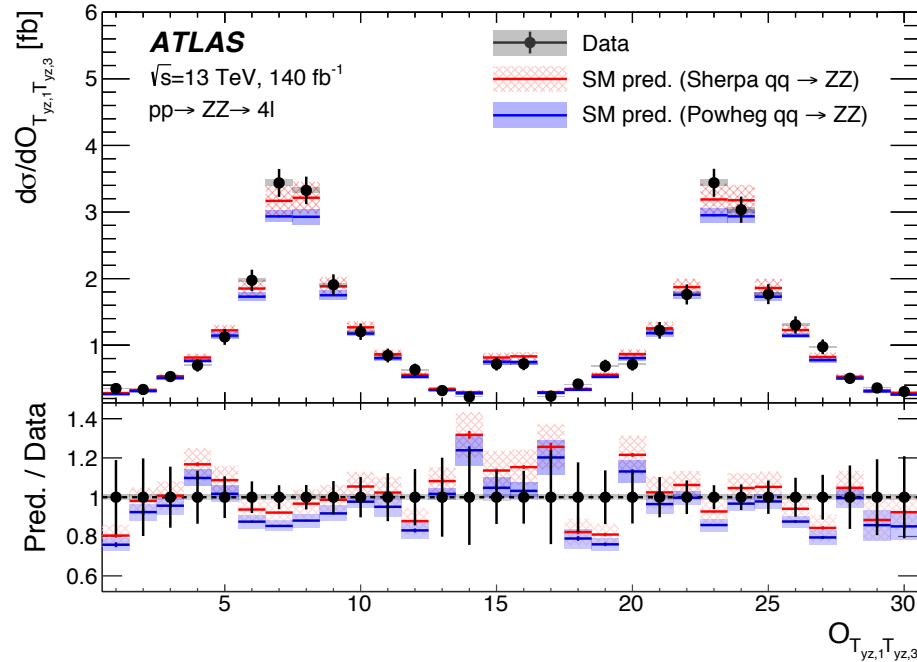


# ZZ → 4l Polarisation and CP Properties

□ Differential cross section as a function of **CP-sensitive Optimal Observable** is measured

❖ Angular observables:  $T_{yz,1(3)} = \sin(\phi_{1(3)}) \times \cos(\theta_{1(3)})$

❖ 2D → 1D mapping:  $T_{yz,1}:T_{yz,3} \rightarrow O_{T_{yz,1}T_{yz,3}}$



$$\sigma^i = \sigma_{\text{SM}}^i + c \cdot \sigma_{\text{interference}}^i + c^2 \cdot \sigma_{\text{quadratic}}^i$$

□ First limits (@ 95% C.L.) on anomalous CP-odd neutral triple gauge couplings,  $f_Z^4$  and  $f_\gamma^4$ , using only the linear interference terms are obtained

aNTGC parameter	Interference only		Full	
	Expected	Observed	Expected	Observed
$f_Z^4$	[-0.16, 0.16]	[-0.12, 0.20]	[-0.013, 0.012]	[-0.012, 0.012]
$f_\gamma^4$	[-0.30, 0.30]	[-0.34, 0.28]	[-0.015, 0.015]	[-0.015, 0.015]

# ZZ → 4l Cross Sections at $\sqrt{s} = 13.6$ TeV

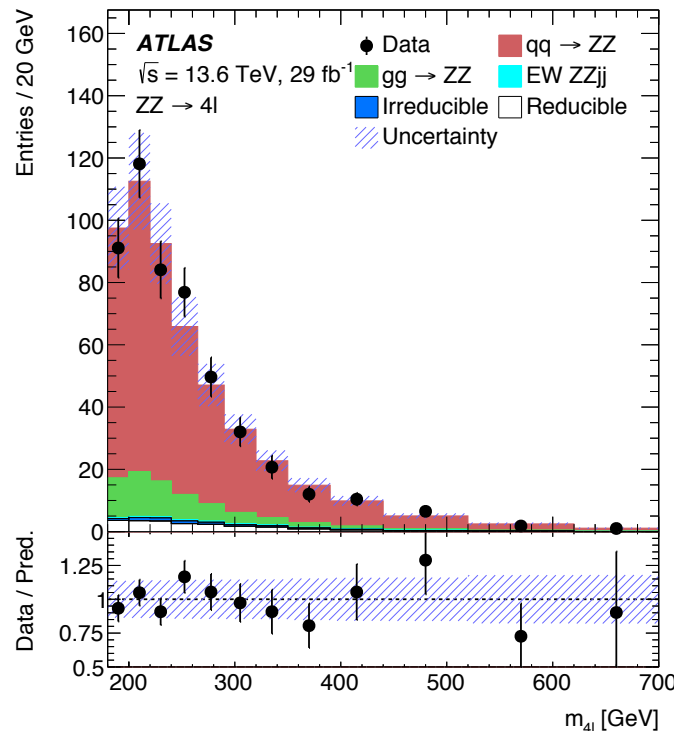
arXiv:2311.09715

$\sqrt{s} = 13.6$  TeV, 29 fb<sup>-1</sup>

□ ZZ production is measured first time by the ATLAS in pp-collisions at  $\sqrt{s} = 13.6$  TeV

❖ Fiducial and differential cross sections are measured in 4l final states

○ Fiducial cross-section is extrapolated to the total cross section satisfying  $66 < m_Z < 116$  GeV for both Z



□ Cut-and-count approach is used

$$\sigma_{fid} = \frac{N_{obs} - N_{bkg}}{\mathcal{L} \times C_{ZZ}}$$

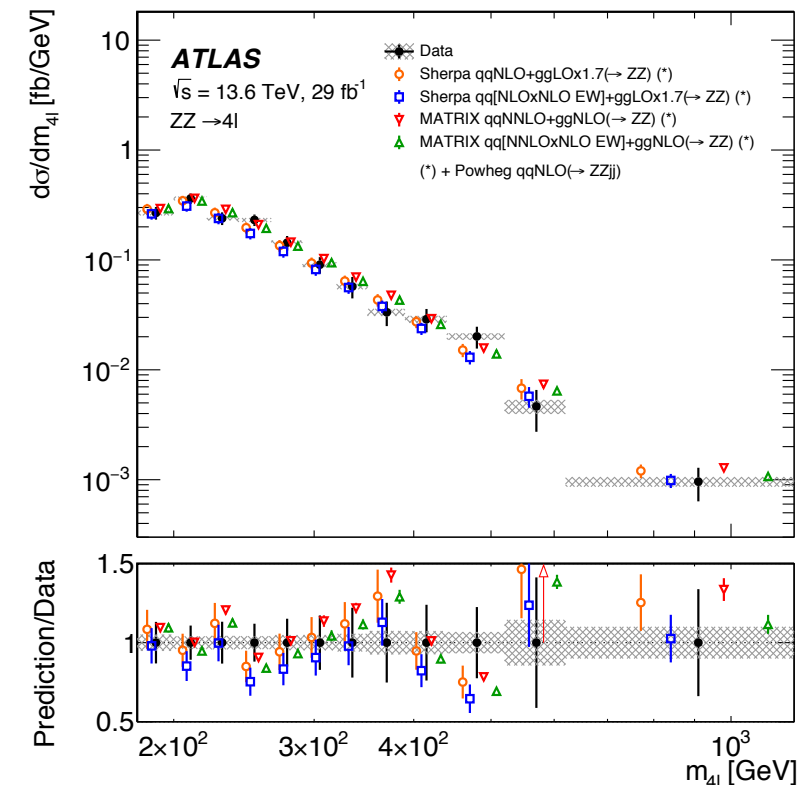
$$C_{ZZ} = 0.555 \pm 0.022$$

$$\sigma_{total} = \frac{\sigma_{fid}}{A_{ZZ} \times BR(ZZ \rightarrow 4l)}$$

$$A_{ZZ} = 0.482 \pm 0.003$$

MATRIX calculations @  
nNNLO QCD ⊗ NLO EW

□ Iterative Bayesian unfolding is used



	Measurement	MC prediction	MATRIX prediction
Fiducial	$36.7 \pm 1.6(\text{stat}) \pm 1.5(\text{syst}) \pm 0.8(\text{lumi})$ fb	$36.8^{+4.3}_{-3.5}$ fb	$36.5 \pm 0.7$ fb
Total	$16.8 \pm 0.7(\text{stat}) \pm 0.7(\text{syst}) \pm 0.4(\text{lumi})$ pb	$17.0^{+1.9}_{-1.4}$ pb	$16.7 \pm 0.5$ pb

# Zγjj Measurement

EW (sensitive to VBS) and extended EW (EW+QCD) production fiducial and differential cross sections are measured in leptonic final states,  $ee$  and  $\mu\mu$

The EW measurements employ a control region to constrain the QCD background

SR:  $\zeta(Z\gamma) < 0.4$ , CR:  $\zeta(Z\gamma) > 0.4$

$$\zeta(Z\gamma) = \left| \frac{y_{Z\gamma} - (y_{j_1} + y_{j_2})/2}{y_{j_1} - y_{j_2}} \right|$$

Profile-likelihood fit to  $m_{jj}$  distributions in the SR and CR (in case of the EW measurement) is used to extract signal normalisation → evaluate fiducial cross sections

Both EW and extended EW fiducial cross sections are in a good agreement with the SM predictions

EW	$\sigma_{EW} = 3.6 \pm 0.5 \text{ fb}$
$(m_{jj} > 500 \text{ GeV})$	$\sigma_{EW}^{pred} = 3.5 \pm 0.2 \text{ fb}$

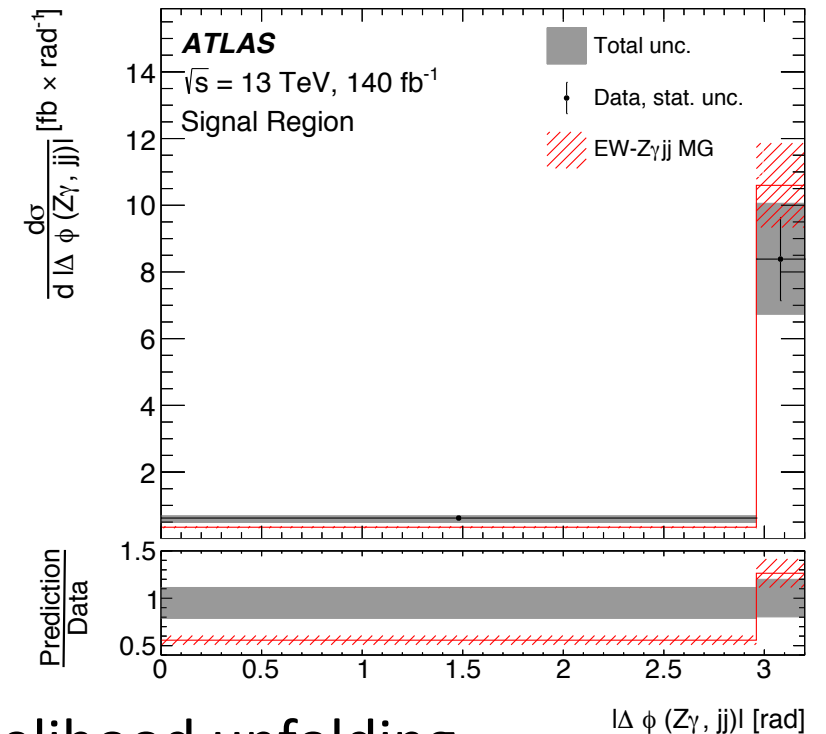
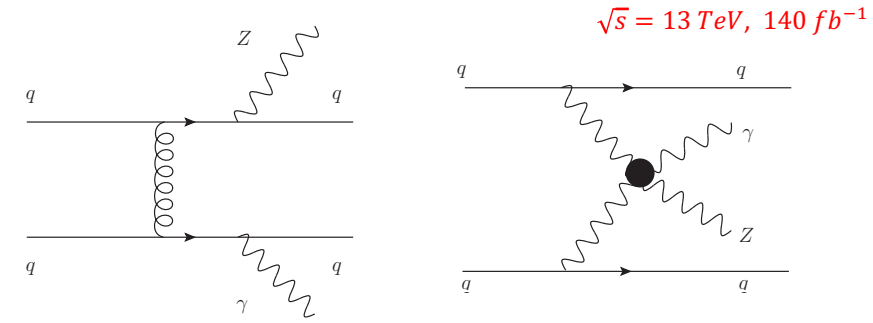
Extended EW	$\sigma_{Z\gamma} = 16.8^{+2.0}_{-1.8} \text{ fb}$
$(m_{jj} > 150 \text{ GeV})$	$\sigma_{Z\gamma}^{pred} = 15.7^{+5.0}_{-2.6} \text{ fb}$

Differential cross sections are measured using profile-likelihood unfolding

Unfolded observables are in a good agreement with SM distributions except of  $|\Delta\phi(Z\gamma, jj)|$

About two standard deviation is observed in the lowest bin of the EW measurement

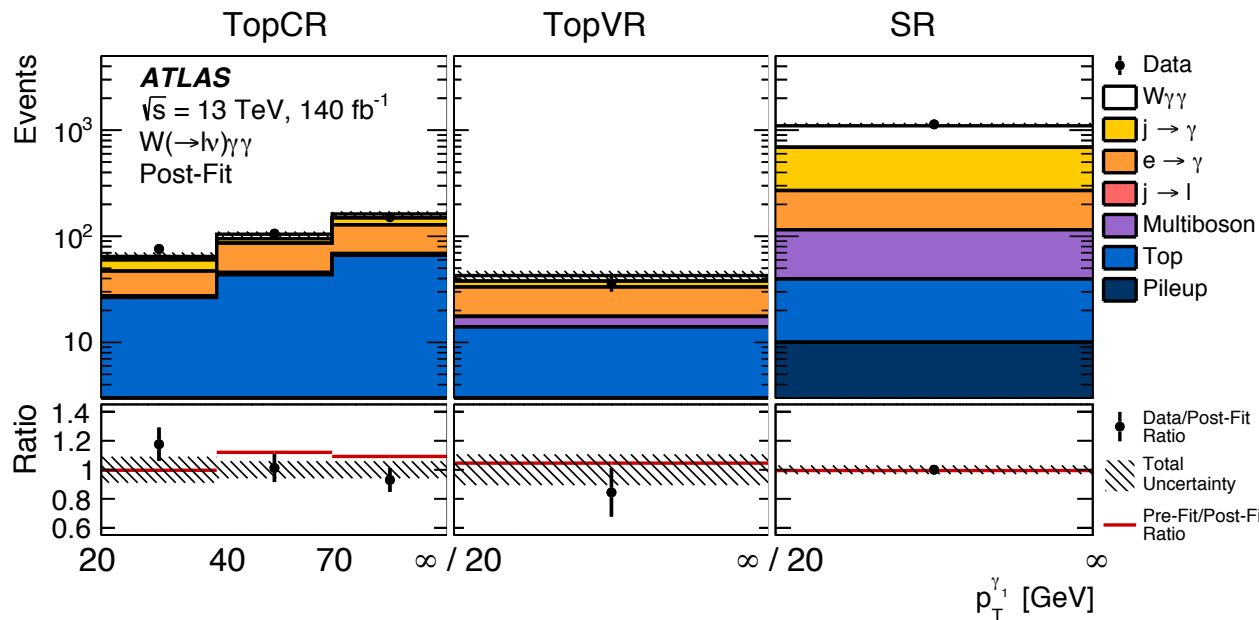
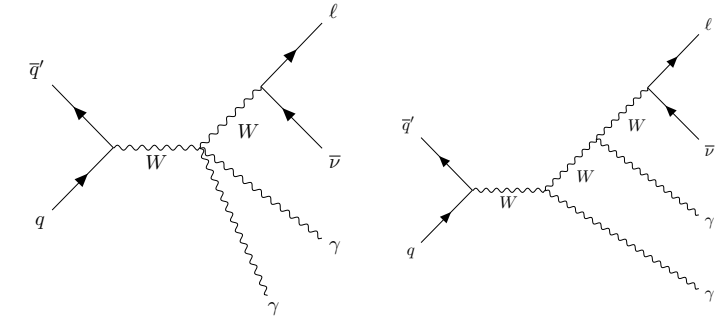
Phys. Lett. B 846 (2023) 138222



# Observation of $W\gamma\gamma$

$W\gamma\gamma$  production in  $e/\mu$  final states is observed first time with significance of  $5.6\sigma$  (observed and expected)

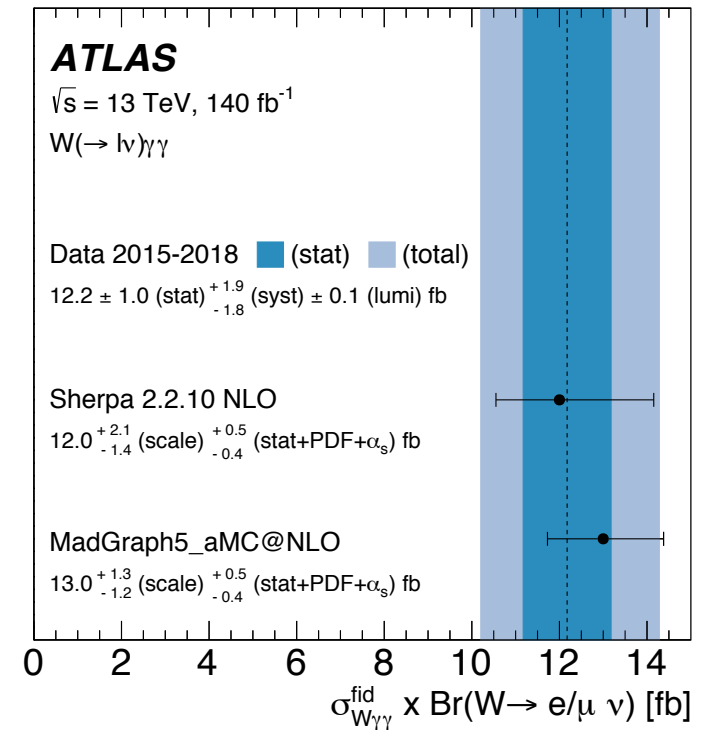
Main background is due to the jet-to-photon and electron-to-photon misidentification



$$\sigma_{fid}^{e/\mu} = \frac{N_{signal}}{\mathcal{L} \times C}$$

Fiducial cross section is measured for combined  $W \rightarrow e\nu/\mu\nu$  events with total 17% uncertainty

Leptonically decaying  $\tau$ -leptons are not considered



# Observation of $WZ\gamma$

□  $WZ\gamma \rightarrow l'^{\pm}\nu l^+l^-\nu\gamma$  production, where  $l^{(\prime)} = e, \mu$ , is observed first time with significance of  $6.3\sigma$  (expected  $5.0\sigma$ )

❖ Dominant background stems from the non-prompt leptons and photons from hadronic decays and from misidentified jets

□ Profile-likelihood fit is done to the signal and two control regions for  $ZZ\gamma$  and  $ZZ(e \rightarrow \gamma)$  backgrounds

❖ Three normalization parameters for the signal and the two backgrounds are fitted simultaneously

❖ All leptonic final state events ( $eee, \mu\mu\mu, eee, e\mu\mu$ ) are combined in the fit

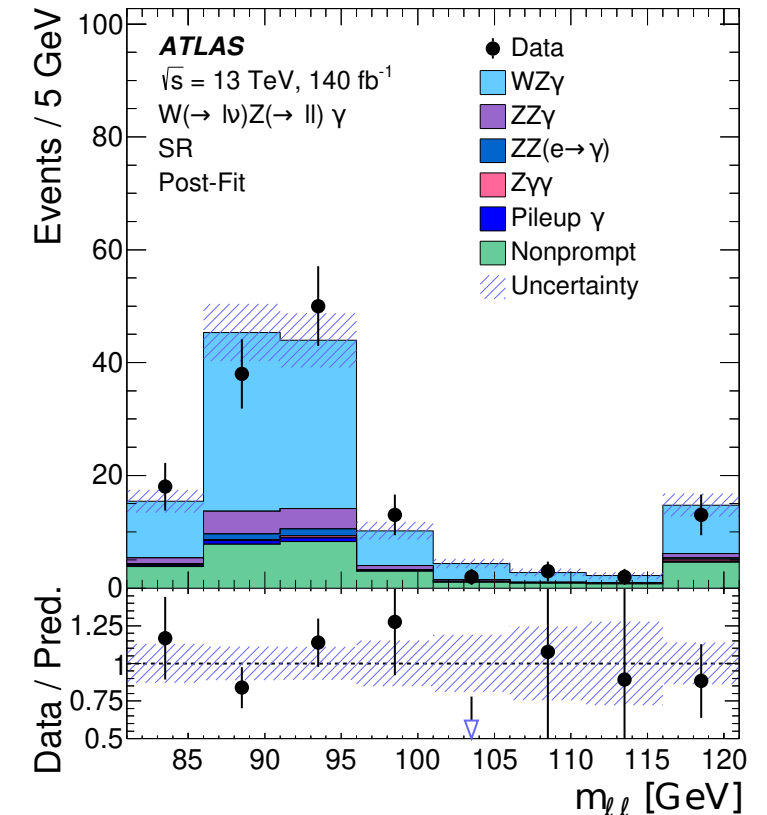
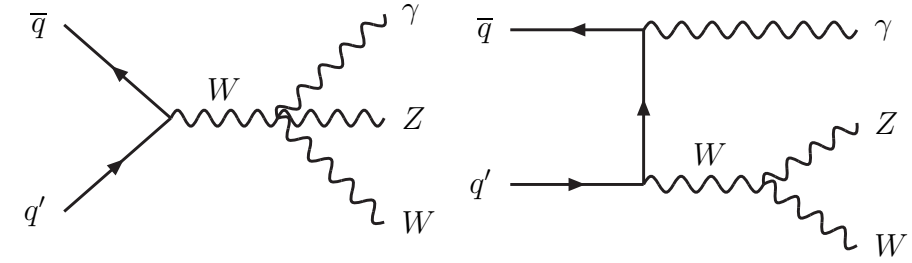
□ Fitted signal normalization parameter  $\mu_{WZ\gamma}$  is used to measure the fiducial cross section

❖  $\sigma_{fid}^{obs} = 2.0 \pm 0.30 (stat.) \pm 0.16 (syst.) fb$

○ Consistent with the  $\sigma_{fid}^{SM} = 1.50 \pm 0.06 (tot.) fb$  within  $1.5\sigma$

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$\sqrt{s} = 13 TeV, 140 fb^{-1}$





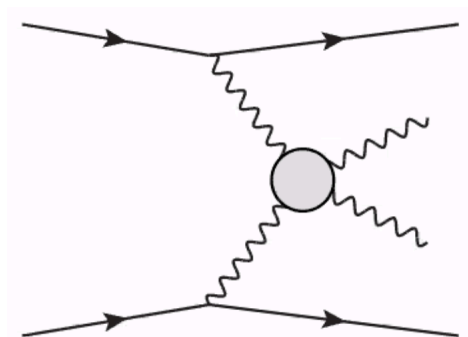
# Summary

- ❑ Measurements of multiboson production processes allow to test the gauge interactions of the SM electroweak theory and its symmetry breaking mechanism
- ❑ Multiboson production final states have relatively small cross sections even at the LHC energy and can be very sensitive to new physics effects leading to anomalous triple and quartic gauge couplings
- ❑ First time observations and precise measurements of various multiboson production processes became possible in proton-proton collision data collected by the LHC experiments, CMS and ATLAS
- ❑ Latest results of the experimental studies of multiboson interactions in the ATLAS detector were discussed
  - ❖ All presented results are consistent with the Standard Model predictions within the measurement uncertainties

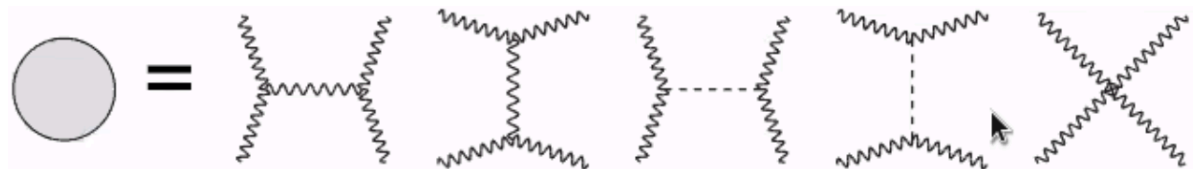
*Thank you!*

# Electroweak Vector Boson Scattering

- Electroweak VBS processes  $V_1 V_2 \rightarrow V_3 V_4$  have not been studied experimentally before the LHC experiments
  - ❖ Low production cross sections even at the LHC energies
    - Sensitivity to possible new physics effects leading to anomalous quartic gauge couplings

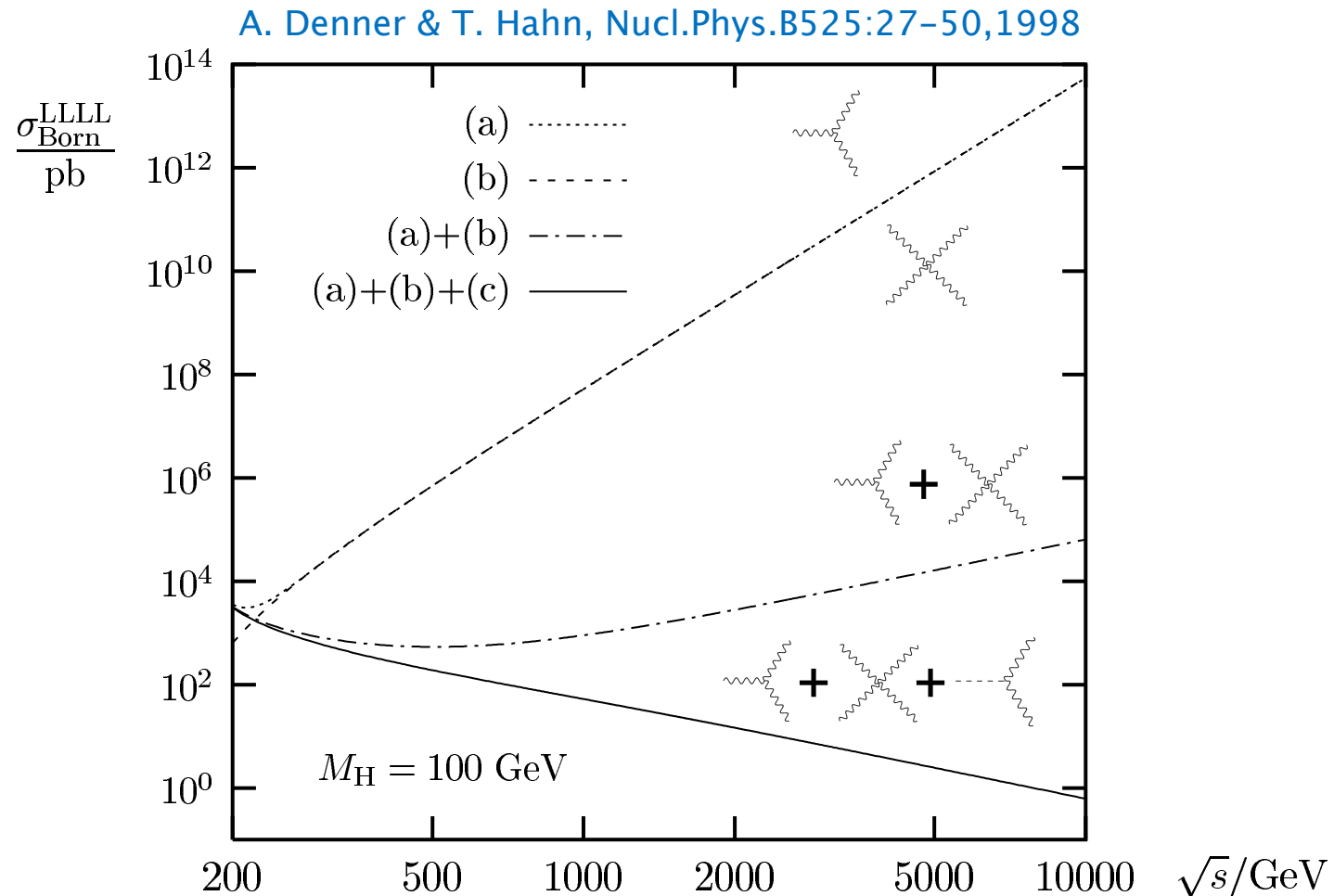


- VBS can involve all or some of the leading order diagrams shown below
  - ❖ Depending on the involved vector bosons



# Electroweak Vector Boson Scattering

- Diagrams with Higgs boson are crucial to avoid a unitarity violation due to the rising scattering cross section of the longitudinally polarized  $W$  bosons:  $W_L W_L \rightarrow W_L W_L$



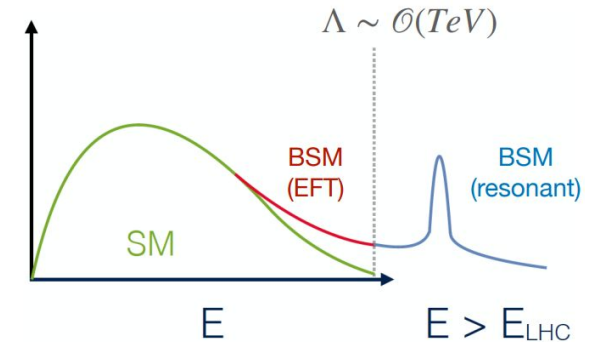
# New Physics Searches – Anomalous Gauge Couplings

□ Low energy effective field theory (EFT) to parameterize indirect new physics effects with the help of high dimension ( $n > 4$ ) operators  $\mathcal{O}_i^{(n)} \rightarrow$  aTGC and aQGC

❖ Linear realization of the SM  $SU(2)_L \otimes U(1)_Y$  gauge symmetry

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{n=5}^{\infty} \sum_i \frac{f_i^{(n)}}{\Lambda^{n-4}} \mathcal{O}_i^{(n)}$$

[Eboli et al., 2020](#)



□ Lowest order operators that generate aQGC (but not aTGC) are dimension-8 operators

	WWWW	WWZZ	ZZZZ	WWAZ	WWAA	ZZZA	ZZAA	ZAAA	AAAA
$\mathcal{O}_{S,0}, \mathcal{O}_{S,1}$	X	X	X						
$\mathcal{O}_{M,0}, \mathcal{O}_{M,1}, \mathcal{O}_{M,6}, \mathcal{O}_{M,7}$	X	X	X	X	X	X	X		
$\mathcal{O}_{M,2}, \mathcal{O}_{M,3}, \mathcal{O}_{M,4}, \mathcal{O}_{M,5}$		X	X	X	X	X	X		
$\mathcal{O}_{T,0}, \mathcal{O}_{T,1}, \mathcal{O}_{T,2}$	X	X	X	X	X	X	X	X	X
$\mathcal{O}_{T,5}, \mathcal{O}_{T,6}, \mathcal{O}_{T,7}$		X	X	X	X	X	X	X	X
$\mathcal{O}_{T,8}, \mathcal{O}_{T,9}$			X			X	X	X	X

TABLE II: Quartic vertices modified by each dimension-8 operator are marked with X.

[Degrande et al., 2013](#)

□ Precise measurements of VBS processes allow to measure the corresponding Wilson coefficients  $f_i^{(8)}$  (look for significant deviations from zero, or set limits)

# EFT Samples in EW Measurements

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{f_i^{(6)}}{\Lambda^2} O_i + \sum_j \frac{f_j^{(8)}}{\Lambda^4} O_j + \dots$$

- EFT “model” for the new physics: only dimension-8 operators have non-zero coefficients
  - ◆ New physics affects only the quartic gauge couplings

- Amplitude of a VBS process with EFT contributions:  $|A_{\text{SM}} + \sum_i c_i A_i|$   $c_i = \frac{f_i^{(8)}}{\Lambda^4}$

- Standard Model, interference, quadratic and cross terms of the total squared amplitude:

$$|A_{\text{SM}} + \sum_i c_i A_i|^2 = |A_{\text{SM}}|^2 + \sum_i c_i 2\text{Re}(A_{\text{SM}}^* A_i) + \sum_i c_i^2 |A_i|^2 + \sum_{ij, i \neq j} c_i c_j 2\text{Re}(A_i A_j^*)$$

- MC samples are generated using only individual terms at a time
- Only one  $c_i$  or one pair of  $c_i$  and  $c_j$  (for generation of cross term samples) are set to nonzero values at a time
  - ◆ Respective sample can be scaled by appropriate  $c_i$ ,  $c_i^2$ , or  $c_i c_j$

# Same Sign $W^\pm W^\pm jj$ Measurement

## □ Monte-Carlo signal and background samples

Process, short description	ME Generator + parton shower	Order	Tune	PDF set in ME
EW, Int, QCD $W^\pm W^\pm jj$ , nominal signal	MADGRAPH5_AMC@NLO2.6.7 + HERWIG7.2	LO	HERWIG	NNPDF3.0NLO
EW, Int, QCD $W^\pm W^\pm jj$ , alternative shower	MADGRAPH5_AMC@NLO2.6.7 + PYTHIA8.244	LO	A14	NNPDF3.0NLO
EW $W^\pm W^\pm jj$ , NLO pQCD approx.	SHERPA2.2.11 & SHERPA2.2.2( $WWW$ ) & POWHEG Box2+PYTHIA8.235 ( $WH$ )	+0,1j@LO NLO	SHERPA A14	NNPDF3.0NNLO
EW $W^\pm W^\pm jj$ , NLO pQCD approx.	POWHEG Boxv2 + PYTHIA8.230	NLO (VBS approx.)	AZNLO	NNPDF3.0NLO
QCD $W^\pm W^\pm jj$ , NLO pQCD approx.	SHERPA2.2.2	+0,1j@LO	SHERPA	NNPDF3.0NNLO
QCD $VVjj$	SHERPA2.2.2	+0,1j@NLO; +2,3j@LO	SHERPA	NNPDF3.0NNLO
EW $W^\pm Z/\gamma^* jj$	MADGRAPH5_AMC@NLO2.6.2+PYTHIA8.235	LO	A14	NNPDF3.0NLO
EW $Z/\gamma^* Z/\gamma^* jj$	SHERPA2.2.2	LO	SHERPA	NNPDF3.0NNLO
QCD $V\gamma jj$	SHERPA2.2.11	+0,1j@NLO; +2,3j@LO	A14	NNPDF3.0NNLO
EW $V\gamma jj$	MADGRAPH5_AMC@NLO2.6.5+PYTHIA8.240	LO	A14	NNPDF3.0NLO
$VVV$	SHERPA2.2.1 (leptonic) & SHERPA2.2.2 (one $V \rightarrow jj$ )	+0,1j@LO	SHERPA	NNPDF3.0NNLO
$t\bar{t}V$	MADGRAPH5_AMC@NLO2.3.3.p0 + PYTHIA8.210	NLO	A14	NNPDF3.0NLO
$tZq$	MADGRAPH5_AMC@NLO2.3.3.p1 + PYTHIA8.212	LO	A14	NNPDF2.3LO
$W^\pm W^\pm jj$ EFT	MADGRAPH5_AMC@NLO 2.6.5 + PYTHIA8.235	LO	A14	NNPDF3.0NLO
$H_5^{\pm\pm}$	MADGRAPH5_AMC@NLO 2.9.5 + PYTHIA8.245	LO	A14	NNPDF3.0NLO

# Same Sign $W^\pm W^\pm jj$ Measurement

## Event selection signal and control regions

- CRs for constraining  $WZjj$ , non-prompt and charge miss-identified lepton backgrounds

Requirement	SR	Low- $m_{jj}$ CR	WZ CR
Leading and subleading lepton $p_T$		$> 27 \text{ GeV}$	
Electron $ \eta $	$< 2.47$ (1.37 in $ee$ ), excluding $1.37 \leq  \eta  \leq 1.52$		
Muon $ \eta $		$< 2.5$	
Leading (subleading) jet $p_T$		$> 65$ (35) $\text{ GeV}$	
Additional jet $p_T$		$> 25 \text{ GeV}$	
Jet $ \eta $		$< 4.5$	
$m_{\ell\ell}$		$> 20 \text{ GeV}$	
$E_T^{\text{miss}}$		$> 30 \text{ GeV}$	
Charge misid. $Z \rightarrow ee$ veto		$ m_{ee} - m_Z  > 15 \text{ GeV}$	–
$b$ -jet veto		$N_{b\text{-jet}} = 0, p_T^{b\text{-jet}} > 20 \text{ GeV},  \eta^{b\text{-jet}}  < 2.5$	–
$N_{\text{veto leptons}}$	$= 0$	$= 0$	$= 1, p_T > 15 \text{ GeV}$
$m_{\ell\ell\ell}$	–	–	$> 106 \text{ GeV}$
$m_{jj}$	$> 500 \text{ GeV}$	$200 < m_{jj} < 500 \text{ GeV}$	$> 200 \text{ GeV}$
$ \Delta y_{jj} $		$> 2$	

## Data, signal and background pre-fit event yields in the SR

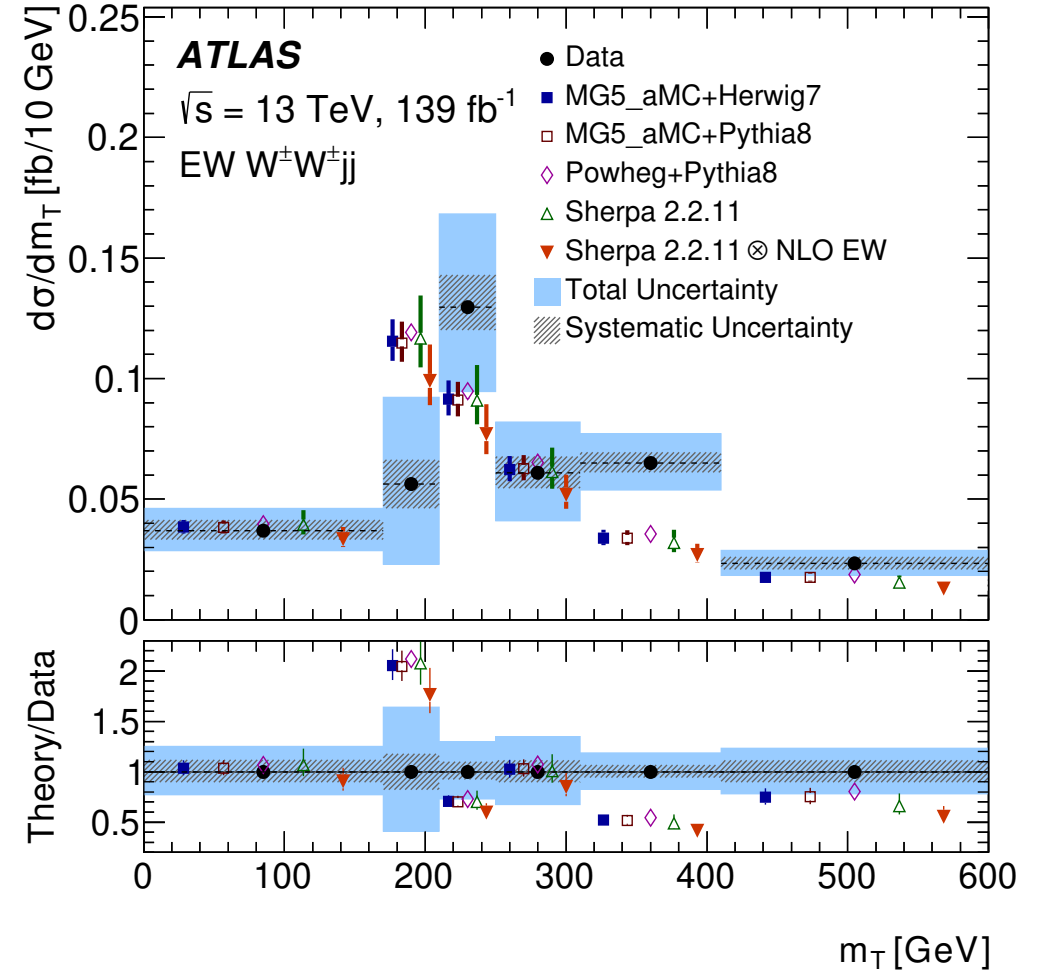
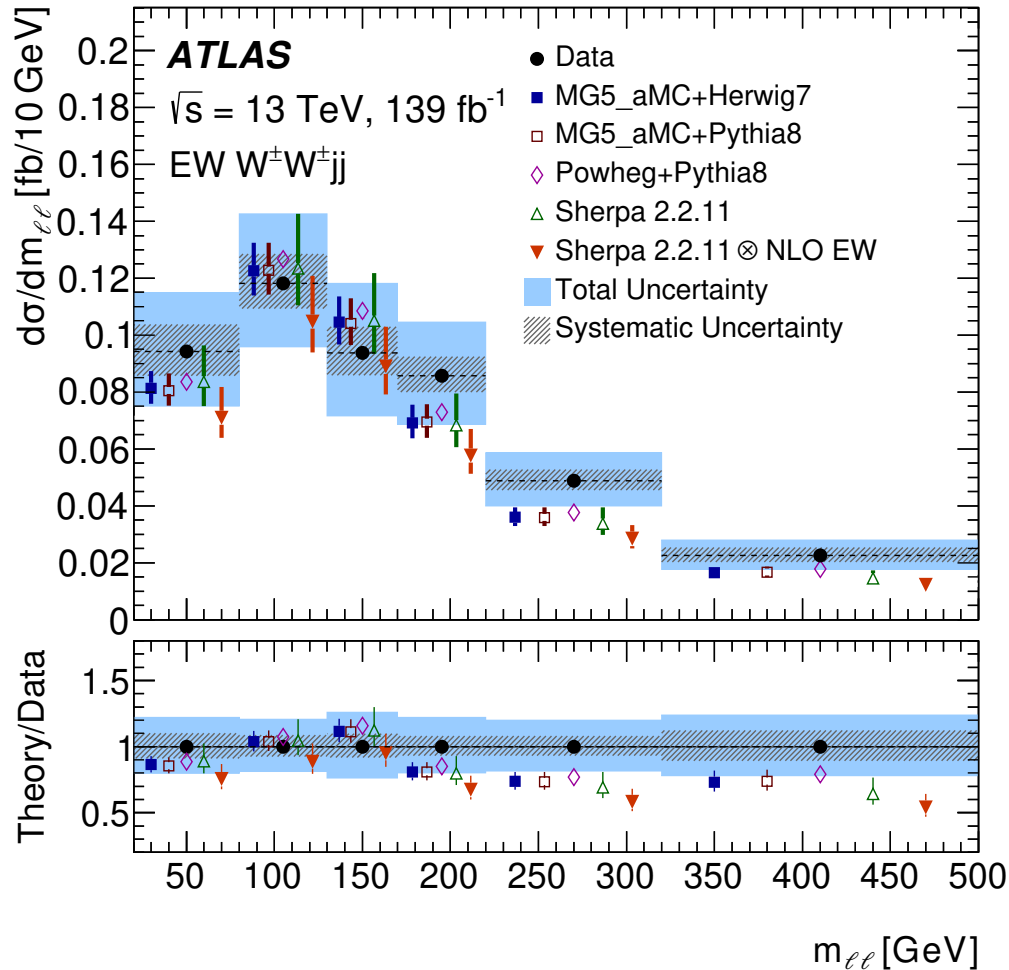
- Sub-regions defined by the same sign lepton flavours
  - 2 sub-regions for electron-muon pairs by distinguished by the leading  $p_T$  lepton flavour

Process	$ee$	$e\mu$	$\mu e$	$\mu\mu$	Combined
$W^\pm W^\pm jj$ EW	27.6 ± 0.9	68.2 ± 1.6	61.3 ± 1.5	77.8 ± 1.7	235 ± 5
$W^\pm W^\pm jj$ QCD	1.6 ± 0.5	7.3 ± 2.2	6.4 ± 1.9	8.8 ± 2.5	24 ± 7
$W^\pm W^\pm jj$ Int	0.93 ± 0.20	2.2 ± 0.5	2.0 ± 0.4	2.5 ± 0.5	7.6 ± 1.6
$W^\pm Zjj$ QCD	8.4 ± 1.0	26.8 ± 3.0	26.7 ± 3.0	20.9 ± 2.2	83 ± 9
$W^\pm Zjj$ EW	1.71 ± 0.14	4.9 ± 0.4	4.1 ± 0.4	4.2 ± 0.4	14.9 ± 1.2
Non-prompt	8.9 ± 2.6	15 ± 4	10.2 ± 3.2	21 ± 7	56 ± 12
$V\gamma$	1.3 ± 0.8	5.1 ± 2.2	4.6 ± 2.6	–	11 ± 5
Charge misid.	3.8 ± 2.0	5.0 ± 1.3	1.2 ± 0.4	–	10 ± 4
Other prompt	1.02 ± 0.29	2.5 ± 0.6	1.8 ± 0.5	1.7 ± 2.2	7.1 ± 2.8
Total expected	55 ± 4	137 ± 7	118 ± 6	137 ± 8	448 ± 20
Data	52	149	127	147	475

# Same Sign $W^\pm W^\pm jj$ Measurement

□ Unfolding results obtained with maximum likelihood fit method

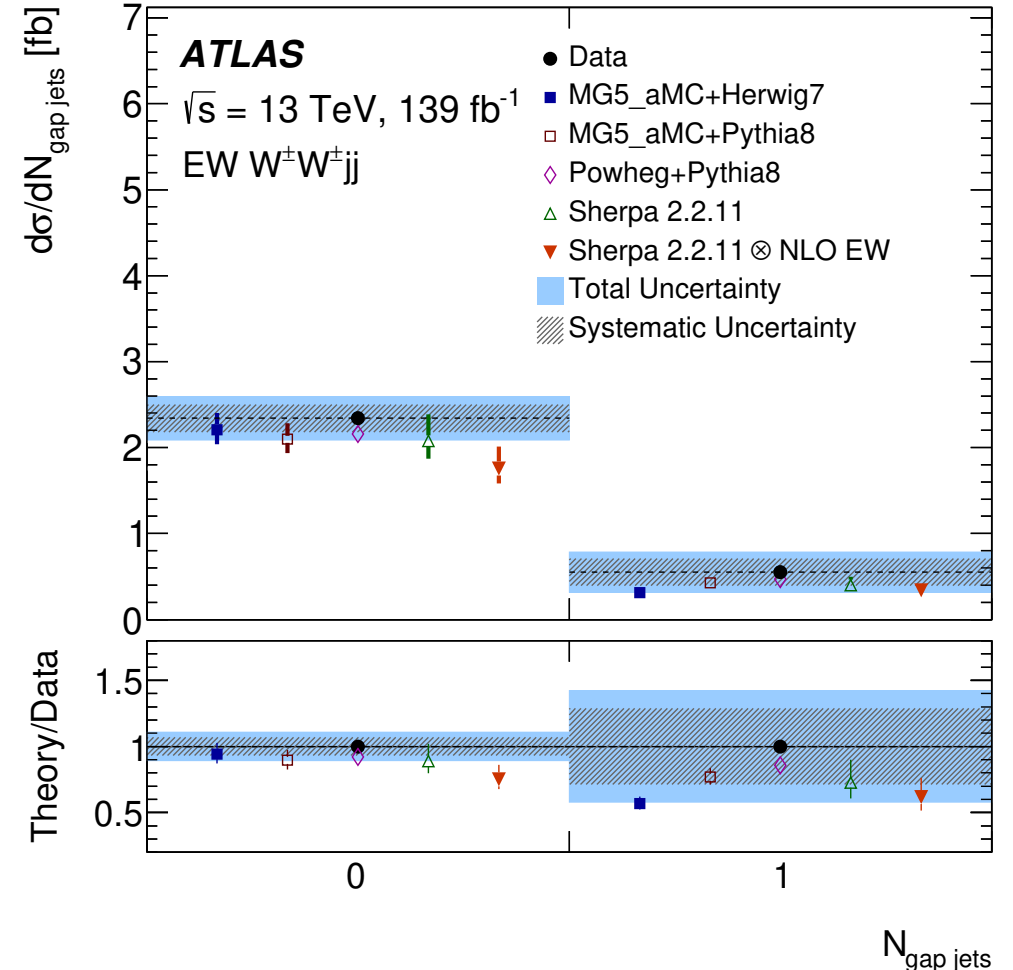
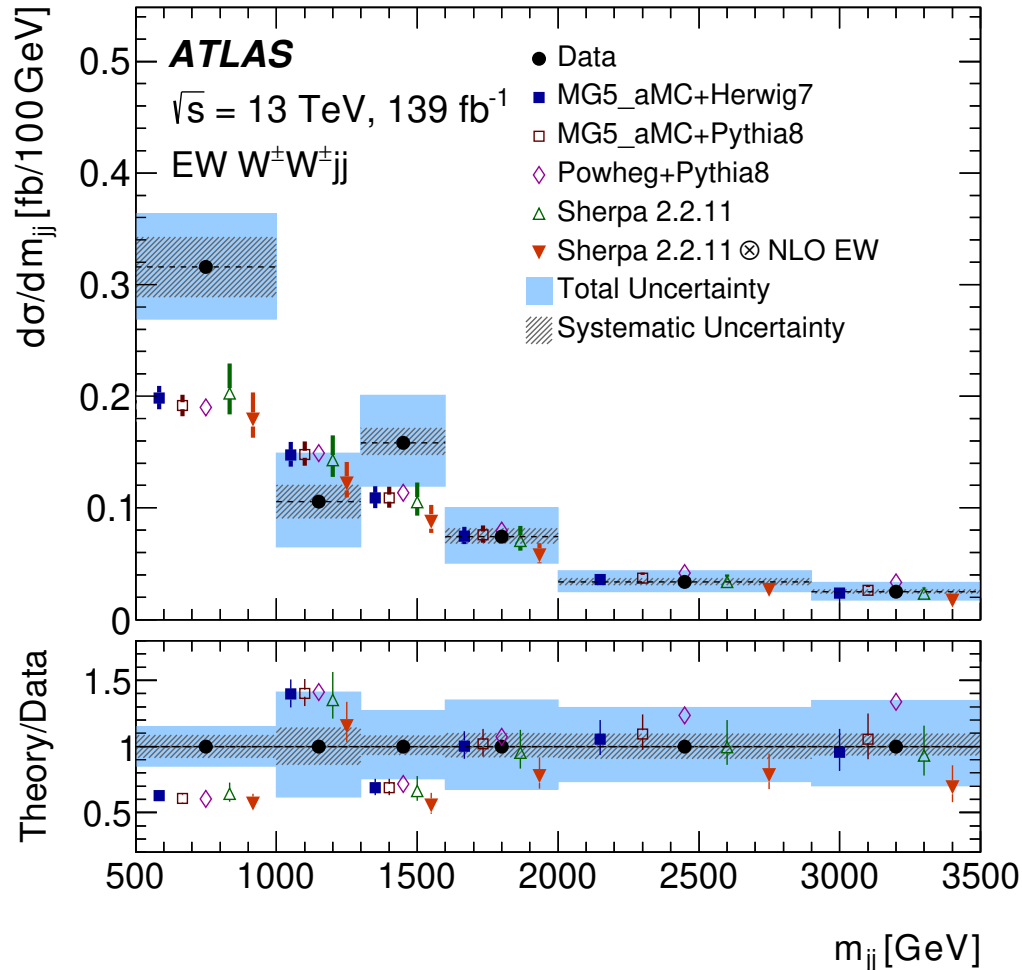
$$m_T = \sqrt{(E_T^{\ell\ell} + E_T^{\text{miss}})^2 - |\vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}}|^2}$$





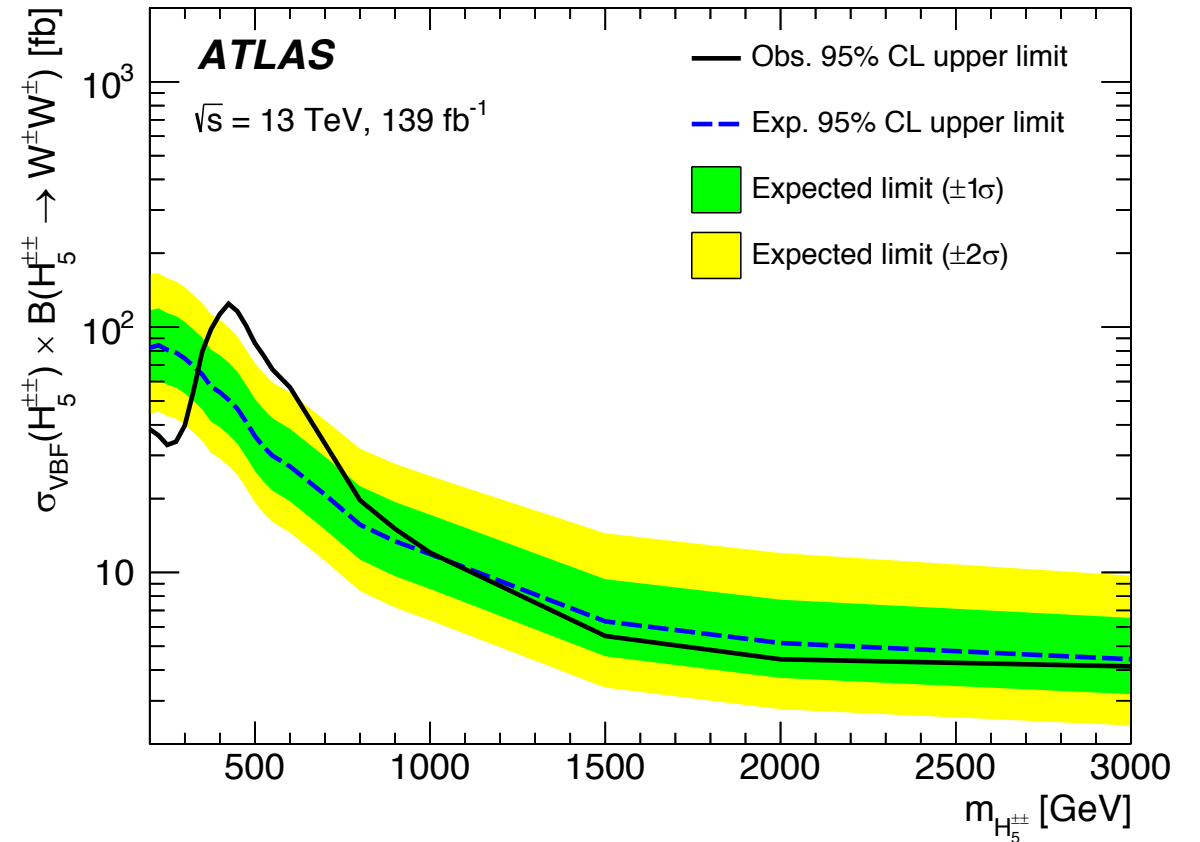
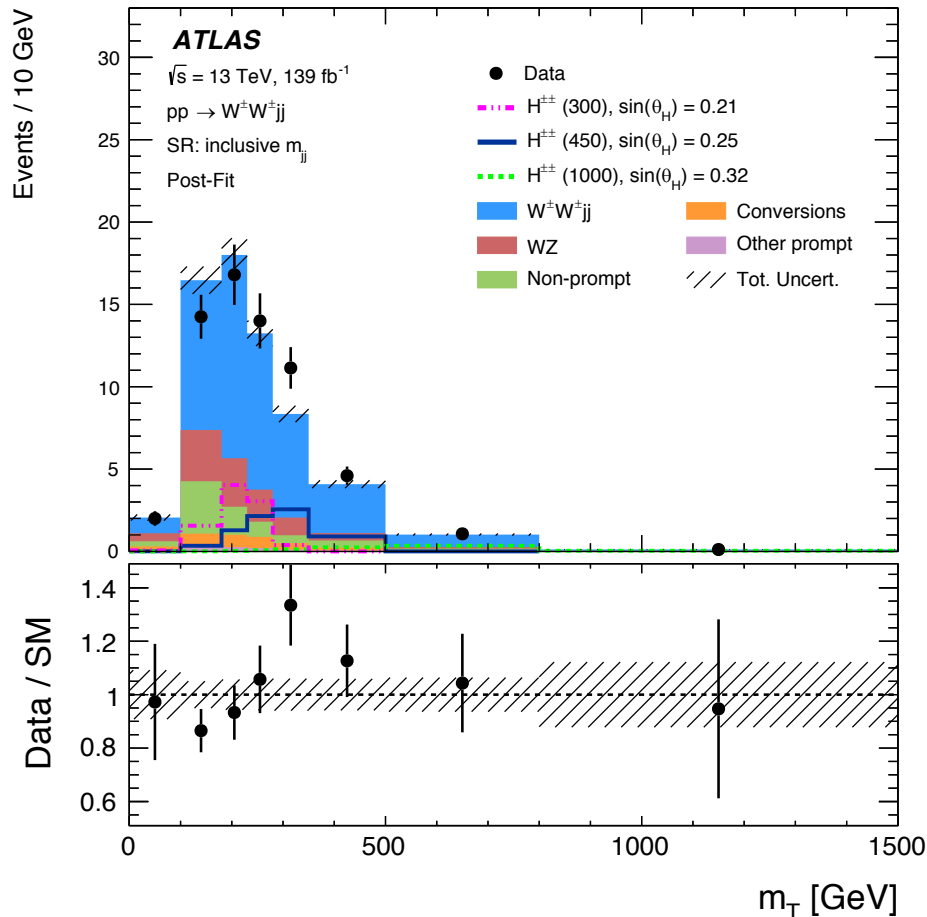
# Same Sign $W^\pm W^\pm jj$ Measurement

□ Unfolding results obtained with maximum likelihood fit method



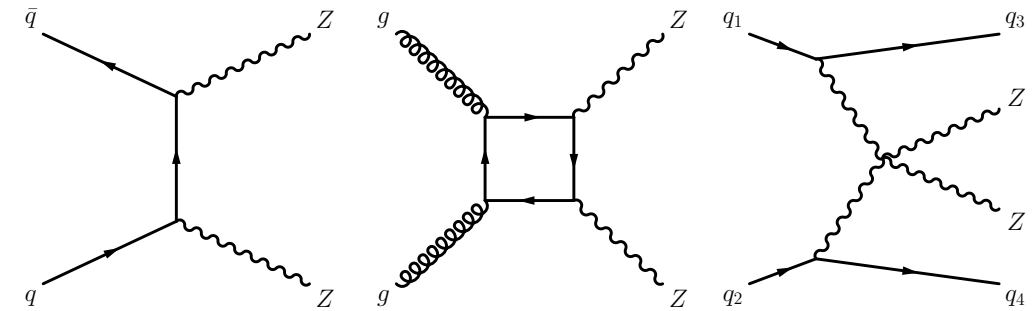
# Same Sign $W^\pm W^\pm jj$ Measurement

- Search for the doubly charged Higgs boson of the GM (Georgi and Machacek) model
- Excess observed at the  $m_T = 450$  GeV with a local (global) significance of  $3.3\sigma$  ( $2.5\sigma$ )

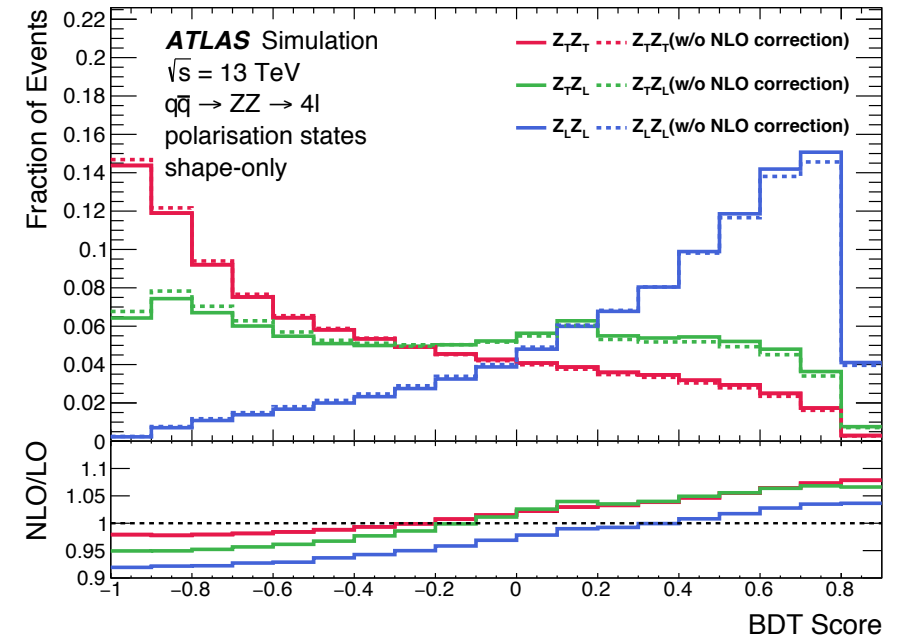


# $ZZ \rightarrow 4l$ Polarisation and CP Properties

- $qq \rightarrow ZZ$  polarisation samples for three helicity states  $Z_L Z_L$ ,  $Z_L Z_T$  and  $Z_T Z_T$  are generated with MadGraph5 at the LO accuracy in QCD
- $gg \rightarrow ZZ$  polarisation samples are obtained with dedicated MC event reweighting method from the inclusive Sherpa sample
- Polarisation and their interference samples are corrected to the NLO accuracy in both QCD and EW



- Boosted Decision Tree classifier is used to discriminate signal from background



# ZZ → 4l Polarisation and CP Properties

## Supporting material

	Pre-fit	Post-fit	
ZZ	$Z_L Z_L$	$189.3 \pm 8.7$	$220 \pm 54$
	$Z_T Z_L$	$710 \pm 29$	$711 \pm 29$
	$Z_T Z_T$	$2170 \pm 120$	$2147 \pm 60$
	Interference	$33.7 \pm 2.8$	$33.4 \pm 2.7$
Non-prompt	$18.7 \pm 7.1$	$18.5 \pm 7.0$	
Others	$20.0 \pm 3.7$	$19.9 \pm 3.7$	
Total	$3140 \pm 150$	$3149 \pm 57$	
Data	3149	3149	

