

# Introdução à Física de Partículas

Introduction to particle Physics

(3/4)



TÉCNICO  
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3 – 8 Setembro, CERN, Geneva

1931

# Modelo do núcleo durante os anos 20:

## Protões + electrões

Inconsistente com a "nova" MQ.

range, estimated visually, was sometimes as much as 3 mm. at N.T.P.

These results, and others I have obtained in the course of the work, are very difficult to explain on the assumption that the radiation from beryllium is a quantum radiation, if energy and momentum are to be conserved in the collisions. The difficulties disappear, however, if it be assumed that the radiation consists of particles of mass 1 and charge 0, or neutrons. The capture of the  $\alpha$ -particle by the  $\text{Be}^9$  nucleus may be supposed to result in the formation of a  $\text{C}^{12}$  nucleus and the emission of the



O prémio Nobel da Física foi atribuído a James Chadwick em 1935;

"for the discovery of the neutron".



Letters to the Editor

*[The Editor does not hold himself responsible for opinions expressed by his correspondents. Neither can he undertake to return, nor to correspond with the writers of, rejected manuscripts intended for this or any other part of NATURE. No notice is taken of anonymous communications.]*

Possible Existence of a Neutron

It has been shown by Bothe and others that beryllium when bombarded by  $\alpha$ -particles of polonium emits a radiation of great penetrating power, which has an absorption coefficient in lead of about  $0.3 \text{ (cm.)}^{-1}$ . Recently Mme. Curie-Joliot and M. Joliot found, when measuring the ionisation produced by this beryllium radiation in a vessel with a thin window, that the ionisation increased when matter containing hydrogen was placed in front of the window. The effect appeared to be due to the ejection of protons with velocities up to a maximum of nearly  $3 \times 10^9 \text{ cm. per sec.}$  They suggested that the transference of energy to the proton was by a process similar to the Compton effect, and estimated that the beryllium radiation had a quantum energy of  $50 \times 10^6 \text{ electron volts.}$

I have made some experiments using the valve counter to examine the properties of this radiation excited in beryllium. The valve counter consists of a small ionisation chamber connected to an amplifier, and the sudden production of ions by the entry of a particle, such as a proton or  $\alpha$ -particle, is recorded by the deflexion of an oscillograph. These experiments have shown that the radiation emits particles from hydrogen, helium, lithium, beryllium, carbon, air, and argon. The particles ejected from hydrogen behave, as regards range and ionising power, like protons with speeds up to about  $3.2 \times 10^9 \text{ cm. per sec.}$  The particles from the other elements have a large ionising power, and appear to be in each case recoil atoms of the elements.

If we ascribe the ejection of the proton to a Compton recoil from a quantum of  $52 \times 10^6 \text{ electron volts,}$  then the nitrogen recoil atom arising by a similar process should have an energy not greater than about 400,000 volts, should produce not more than about 10,000 ions, and have a range in air at N.T.P. of about 1.3 mm. Actually, some of the recoil atoms in nitrogen produce at least 30,000 ions. In collaboration with Dr. Feather, I have observed the recoil atoms in an expansion chamber, and their range, estimated visually, was sometimes as much as 3 mm. at N.T.P.

These results, and others I have obtained in the course of the work, are very difficult to explain on the assumption that the radiation from beryllium is a quantum radiation, if energy and momentum are to be conserved in the collisions. The difficulties disappear, however, if it be assumed that the radiation consists of particles of mass 1 and charge 0, or neutrons. The capture of the  $\alpha$ -particle by the  $\text{Be}^9$  nucleus may be supposed to result in the formation of a  $\text{C}^{12}$  nucleus and the emission of the neutron. From the energy relations of this process the velocity of the neutron emitted in the forward direction may well be about  $3 \times 10^9 \text{ cm. per sec.}$  The collisions of this neutron with the atoms through which it passes give rise to the recoil atoms, and the observed energies of the recoil atoms are in fair agreement with this view. Moreover, I have observed that the protons ejected from hydrogen by the radiation emitted in the opposite direction to that of the exciting  $\alpha$ -particle appear to have a much smaller range than those ejected by the forward radiation.

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This again receives a simple explanation on the neutron hypothesis.

If it be supposed that the radiation consists of quanta, then the capture of the  $\alpha$ -particle by the  $\text{Be}^9$  nucleus will form a  $\text{C}^{12}$  nucleus. The mass defect of  $\text{C}^{12}$  is known with sufficient accuracy to show that the energy of the quantum emitted in this process cannot be greater than about  $14 \times 10^6 \text{ volts.}$  It is difficult to make such a quantum responsible for the effects observed.

It is to be expected that many of the effects of a neutron in passing through matter should resemble those of a quantum of high energy, and it is not easy to reach the final decision between the two hypotheses. Up to the present, all the evidence is in favour of the neutron, while the quantum hypothesis can only be upheld if the conservation of energy and momentum be relinquished at some point.

J. CHADWICK.

Cavendish Laboratory,  
Cambridge, Feb. 17.

The Oldoway Human Skeleton

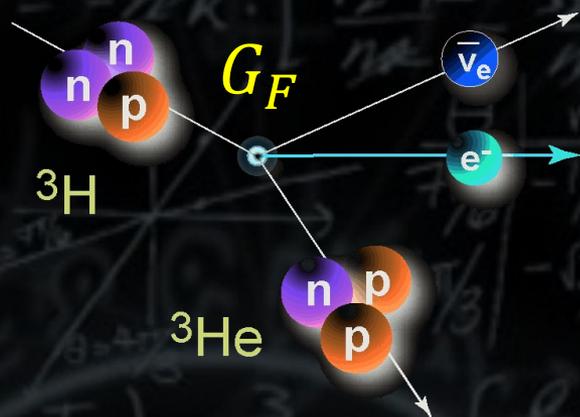
A LETTER appeared in NATURE of Oct. 24, 1931, signed by Messrs. Leakey, Hopwood, and Beck, in which, among other conclusions, it is stated that "there is no possible doubt that the human skeleton came from Bed No. 2 and not from Bed No. 4". This must be taken to mean that the skeleton is to be considered as a natural deposit in Bed No. 2, which is overlaid by the later beds Nos. 3 and 4, and that all consideration of human interment is ruled out.

If this be true, it is a most unusual occurrence. The skeleton, which is of modern type, with fibulae, was found completely articulated down even to the phalanges, and in a position of extraordinary contraction. Complete mammalian skeletons of any age are, as field palaeontologists know, of great rarity. When they occur, their perfection can usually be explained as the result of sudden death and immediate covering by volcanic dust. Many of the more or less perfect skeletons which may be seen in museums have been rearticulated from bones found somewhat scattered as the result of death from floods, or in the neighbourhood of drying water-holes. We know of no case of a perfect articulated skeleton being found in company with such broken and scattered remains as appear to be abundant at Oldoway. Either the skeletons are all complete, as in the *Stenomylus* quarry at Sioux City, Nebraska, or are all scattered and broken in various degrees, as in ordinary bone beds. The probability, therefore, that the Oldoway skeleton represents an artificial burial is thus one that will occur to palaeontologists.

The skeleton was exhumed in 1913, and published photographs show that the excavation made for its disinterment was extensive. It is, therefore, very difficult to believe that in 1931 there can be reliable evidence left at the site as to the conditions under which it was deposited. If naturally deposited in Bed No. 2, the skeleton is of the highest possible importance, because it would be of pre-Monsterian age, and would be in the company of *Pithecanthropus* and the Pittdown, Heidelberg, and Peking men, all of whose remains are fragmentary to the last degree. Of the few other human remains for which such antiquity is claimed, the Galley Hill skeleton and the Ipswich skeleton are, or apparently were, complete. The first of these was never seen *in situ* by any trained observer, and the latter has, we believe, been withdrawn by its discoverer. The other fragments, found long ago, are entirely without satisfactory evidence as to their mode of occurrence.

1934

Em 1934, Fermi propõe a primeira teoria para explicar o decaimento  $\beta$  dos núcleos.



Esta descrição não é válida para toda a gama de energias...

Tempo de vida do muão:

$$\tau_\mu = \frac{192\pi^3}{G_F^2 m_\mu^5}$$

Constante de Fermi:

$$G_F = 1.166364 \times 10^5 \text{ GeV}^{-2}$$

## TENTATIVO DI UNA TEORIA DEI RAGGI $\beta$

Nota <sup>(1)</sup> di ENRICO FERMI

**Sunto.** - Si propone una teoria quantitativa dell'emissione dei raggi  $\beta$  in cui si ammette l'esistenza del « neutrino » e si tratta l'emissione degli elettroni e dei neutrini da un nucleo all'atto della disintegrazione  $\beta$  con un procedimento simile a quello seguito nella teoria dell'irradiazione per descrivere l'emissione di un quanto di luce da un atomo eccitato. Vengono dedotte delle formule per la vita media e per la forma dello spettro continuo dei raggi  $\beta$ , e le si confrontano coi dati sperimentali.

### Ipotesi fondamentali della teoria.

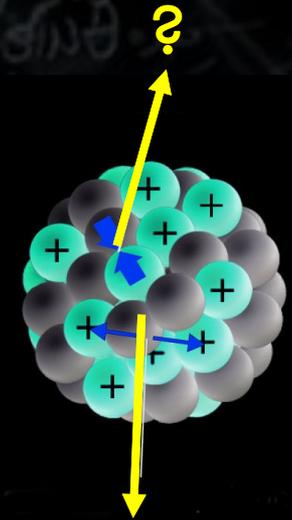
§ 1. Nel tentativo di costruire una teoria degli elettroni nucleari e dell'emissione dei raggi  $\beta$ , si incontrano, come è noto, due difficoltà principali. La prima dipende dal fatto che i raggi  $\beta$  primari vengono emessi dai nuclei con una distribuzione continua di velocità. Se non si vuole abbandonare il principio della conservazione dell'energia, si deve ammettere perciò che una frazione dell'energia che si libera nel processo di disintegrazione  $\beta$  sfugga alle nostre attuali possibilità di osservazione. Secondo la proposta di PAULI si può p. es. ammettere l'esistenza di una nuova particella, il così detto « neutrino », avente carica elettrica nulla e massa dell'ordine di grandezza di quella dell'elettrone o minore. Si ammette poi che in ogni processo  $\beta$  vengano emessi simultaneamente un elettrone, che si osserva come

1934

O núcleo é constituído por prótons e neutrões...

Mas qual a força que os mantém ligado no núcleo?

To remove this defect, it seems natural to modify the theory of Heisenberg and Fermi in the following way. The transition of a heavy particle from neutron state to proton state is not always accompanied by the emission of light particles, i. e., a neutrino and an electron, but the energy liberated by the transition is taken up sometimes by another heavy particle, which in turn will be transformed from proton state into neutron state. If the probability of occurrence of the latter process is much larger than that of the former, the interaction between

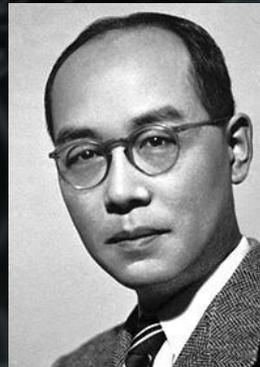


Repulsão electrostática

Potencial de Yukawa:

$$V(r) = -g^2 \frac{e^{-\frac{mcr}{\hbar}}}{r^2}$$

$$m \sim 100 \text{ MeV}$$



O prémio Nobel da Física foi atribuído a Hideki Yukawa em 1949;

"for his prediction of the existence of mesons on the basis of theoretical work on nuclear forces".

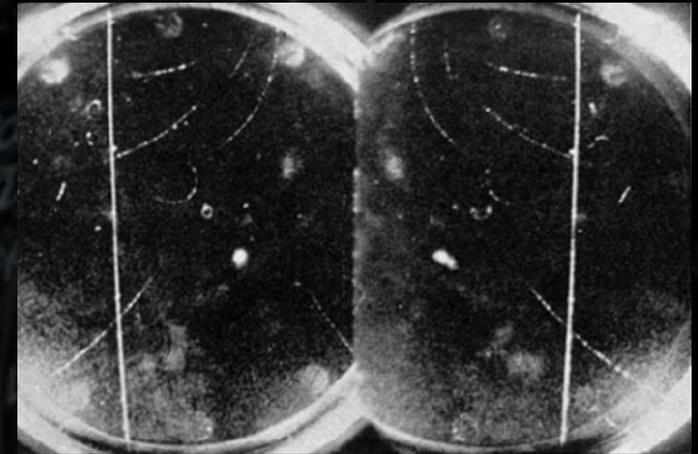


1936

Em 1936 Anderson observam novos rastros numa câmara de Wilson com massa "intermédia".

Seria esta partícula de Yukawa?

Não. Esta nova partícula não tinha afinidade para se ligar ao núcleo.



Anderson chamou a esta partícula o "mesotrão" (hoje conhecida como muão  $\mu$ ). O pião viria a ser descoberto por Cecil Powell em 1947.



O prémio Nobel da Física foi atribuído a Cecil Powell em 1950;

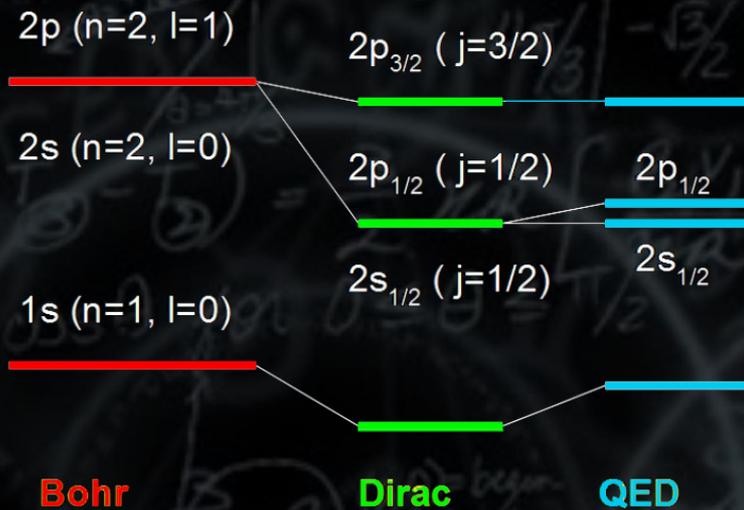
"for his development of the photographic method of studying nuclear processes and his discoveries regarding mesons made with this method".



1947

$$E_{n,j} = mc^2 \left[ 1 + \frac{Z^2}{(n - \varepsilon_j)^2} \right]^{\frac{1}{2}}, \quad \varepsilon_j = j + \frac{1}{2} - \sqrt{\left( j + \frac{1}{2} \right)^2 - Z^2 e^4}$$

Segundo as previsões de Dirac, os estados  $^2S_{1/2}$  e  $^2P_{1/2}$  são degenerados



Em 1947 Willis Lamb e Robert Retherford mediram uma diferença de energia entre estes dois níveis.



O prêmio Nobel da Física foi atribuído a Willis Lamb em 1955;

"for his discoveries concerning the fine structure of the Hydrogen spectrum".



1947

1936



Bethe calculou pela primeira vez o desvio de Lamb obtendo o valor de:

1040 MHz



Conferência em Shelter Island (1947)

... figure and  $K = mc^2$ , the logarithm has the value 7.63, and we find

$$W_{ns'} = 136 \ln[K/(E_n - E_m)] = 1040 \text{ megacycles.} \quad (12)$$

<sup>7</sup> I am indebted to Dr. Stehn and Miss Steward for the numerical calculations.

Bethe foi o primeiro a determinar as correcções relevantes para o desvio de Lamb.

PHYSICAL REVIEW

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The Electromagnetic Shift of Energy Levels

H. A. BETHE  
Cornell University, Ithaca, New York  
(Received June 27, 1947)

BY very beautiful experiments, Lamb and Retherford<sup>1</sup> have shown that the fine structure of the second quantum state of hydrogen does not agree with the prediction of the Dirac theory. The  $2s$  level, which according to Dirac's theory should coincide with the  $2p_1$  level, is actually higher than the latter by an amount of about  $0.033 \text{ cm}^{-1}$  or 1000 megacycles. This discrepancy had long been suspected from spectroscopic measurements.<sup>2,3</sup> However, so far no satisfactory theoretical explanation has been given. Kemble and Present, and Pasternack<sup>4</sup> have shown that the shift of the  $2s$  level cannot be

explained by a nuclear interaction of reasonable magnitude, and Uehling<sup>5</sup> has investigated the effect of the "polarization of the vacuum" in the Dirac hole theory, and has found that this effect also is much too small and has, in addition, the wrong sign.

Schwinger and Weisskopf, and Oppenheimer have suggested that a possible explanation might be the shift of energy levels by the interaction of the electron with the radiation field. This shift comes out infinite in all existing theories, and has therefore always been ignored. However, it is possible to identify the most strongly (linearly) divergent term in the level shift with an electromagnetic *mass* effect which must exist for a bound as well as for a free electron. This effect should

<sup>1</sup> Phys. Rev. 72, 241 (1947).

<sup>2</sup> W. V. Houston, Phys. Rev. 51, 446 (1937).

<sup>3</sup> R. C. Williams, Phys. Rev. 54, 558 (1938).

<sup>4</sup> E. C. Kemble and R. D. Present, Phys. Rev. 44, 1031 (1932); S. Pasternack, Phys. Rev. 54, 1113 (1938).

<sup>5</sup> E. A. Uehling, Phys. Rev. 48, 55 (1935).

# Mais problemas...

Apesar de resolver algumas questões em aberto, a MQR ainda não era uma teoria satisfatória.

$$E_c \sim mc^2 \rightarrow p \sim mc \rightarrow \lambda \sim \frac{h}{mc} = \lambda_c$$

$\lambda_c$  - Comprimento de onda de Compton

**PRINCIPIO DA INCERTEZA:**  $\Delta p \geq \frac{h}{\Delta x} = mc$

**Conclusão:** Se tentarmos localizar uma partícula de massa  $m$  numa região do espaço de dimensões menores que  $\lambda_c$ , então as flutuações na energia são suficientes para criar um par partícula-antipartícula.

UMA TEORIA COMPLETAMENTE RELATIVISTA NÃO SE PODE BASEAR NA IDEIA DE QUE UM SISTEMA PODE SER DESCRITO PELA FUNÇÃO DE ONDA A 1 PARTÍCULA.

# TEORIA QUÂNTICA DE CAMPOS (TQC)

MECÂNICA QUÂNTICA: Posição e momento são tratados como operadores:

$$\hat{x} \psi = x \psi \quad , \quad \hat{p} \psi = -i\hbar \frac{\partial \psi}{\partial x}$$

## E O TEMPO † ?

Numa teoria relativista espaço deveriam estar em pé de igualdade. Mas em MQ o tempo é um parâmetro.

~~$$\hat{T} \psi = t \psi$$~~

EM TQC, A POSIÇÃO E O MOMENTO SÃO DESPROMOVIDOS A PARÂMETROS QUE SÃO ARGUMENTOS DE UM CAMPO

$\hat{\phi}(x, t)$  SEGUNDA QUANTIZAÇÃO

# TEORIA QUÂNTICA DE CAMPOS (TQC)

$$\hat{\phi}(x, t) = f(\hat{a}, \hat{a}^+)$$

$\hat{a}, \hat{a}^+$  - operadores de criação e aniquilação

Estados em TQC:  $|\vec{p}\rangle = \hat{a}^+(\vec{p}) | 0 \rangle$

Estados com várias partículas:

Estados com várias partículas:  $|\vec{p}_1, \vec{p}_2\rangle = \hat{a}^+(\vec{p}_1) \hat{a}^+(\vec{p}_2) | 0 \rangle$

Aniquilação:  $\hat{a}(\vec{p}_1) |\vec{p}_1, \vec{p}_2\rangle = |\vec{p}_2\rangle$

O operador campo actua em estados de tal modo que se podem criar ou aniquilar partículas!

# Electrodinâmica quântica

TQC DOS FOTÕES, ELECTRÕES, POSITRÕES E SUAS INTERAÇÕES

## ElectroDinâmica Quântica - EDQ (QED)



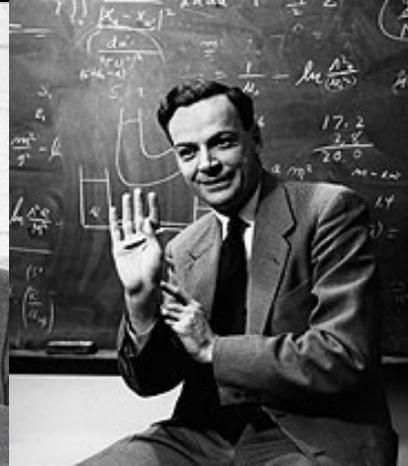
Schwinger



Tomonaga



Dyson



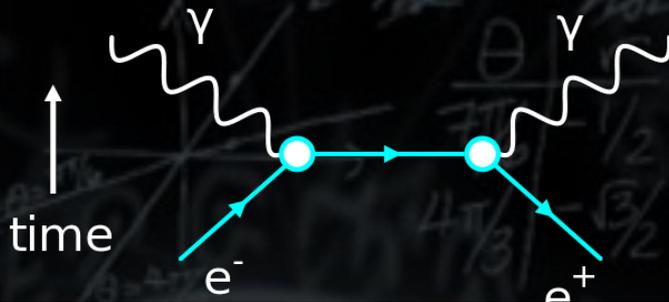
Feynman



The Nobel Prize in Physics 1965 was awarded jointly to Sin-Itiro Tomonaga, Julian Schwinger and Richard P. Feynman "for their fundamental work in quantum electrodynamics, with deep-ploughing consequences for the physics of elementary particles".

1948

Em 1948 Feynman desenvolveu um método “gráfico” que permite calcular processos que envolvem interações entre partículas elementares.



A furgoneta Feynman



Para o diagrama acima:  $\mathcal{M} = \varepsilon_{\mu}^{*}(p_1, \lambda_1)\varepsilon_{\nu}^{*}(p_2, \lambda_2)\bar{v}(e^{+})(ie\gamma^{\nu})\frac{i}{\not{q} - m}(ie\gamma^{\mu})u(e^{-})$

A TÉCNICA DOS DIAGRAMAS DE FEYNMAN É USADA PARA CALCULAR AS “PROBABILIDADES” DE OCORRÊNCIA DE PROCESSOS ENVOLVENDO PARTÍCULAS ELEMENTARES.

# O VÁCUO QUÂNTICO

## vácuo

(latim *vacuus*, -a, -um)

*adj.*

1. Que não contém nada; que não se acha ocupado por coisa alguma. = OCO, VAZIO
2. [Jurídico, Jurisprudência] Que se possui mas que não se desfruta ou não se goza.

*s. m.*

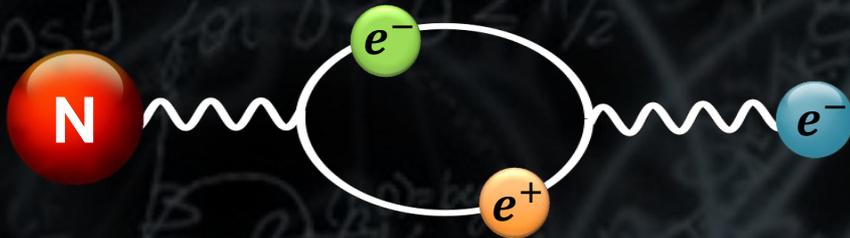
3. Espaço circunscrito que não contém ar ou que se supõe vazio.
4. O espaço entre os corpos celestes que se supõe vazio.

EM FÍSICA QUÂNTICA O VÁCUO ESTÁ LONGE DE SER ESPAÇO VAZIO...

PIH Energia-tempo

$$\Delta E \Delta t \geq \frac{\hbar}{2}$$

**2ª Quantização:** O campo E.M. pode ser interpretado como um conjunto infinito de osciladores harmônicos cujo estado fundamental tem energia não nula.



**Polarização do vácuo**

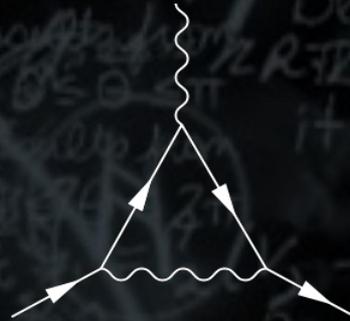
À semelhança da polarização de um dielétrico.

# O DESVIO DE LAMB

Polarização do vácuo



Correcção ao vértice



Correcção à energia própria



2p (n=2, l=1)

2p<sub>3/2</sub> (j=3/2)

-27 MHz

1017 MHz

2s (n=2, l=0)

2p<sub>1/2</sub> (j=1/2)

2p<sub>1/2</sub>

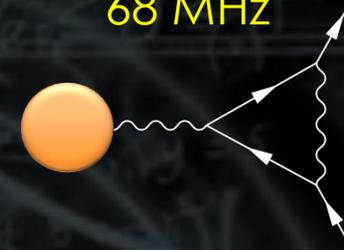


68 MHz

1s (n=1, l=0)

2s<sub>1/2</sub> (j=1/2)

2s<sub>1/2</sub>



Teoria: 1058 MHz , Exp: 1057.9

Bohr

Dirac

QED

# O EFEITO CASIMIR

**Segunda Quantização:** O campo electromagnético consiste num conjunto de osciladores harmónicos.

$$\text{CASIMIR (1948): } \frac{F(a)}{A} = \frac{\pi^2 \hbar c}{240 a^4}$$



Entre as placas condutoras só alguns modos são permitidos.  
**EFEITO:** Força entre as placas condutoras.

VOLUME 78, NUMBER 1

PHYSICAL REVIEW LETTERS

6 JANUARY 1997

## Demonstration of the Casimir Force in the 0.6 to 6 $\mu\text{m}$ Range

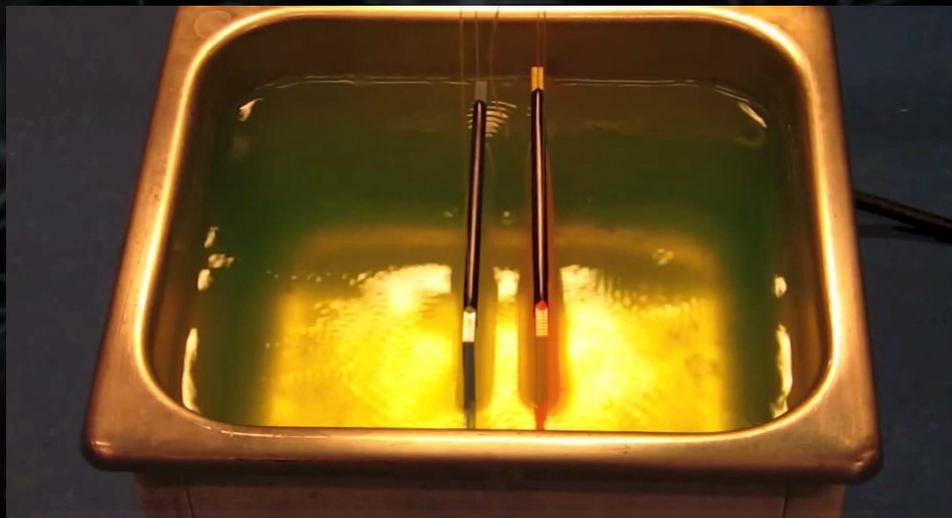
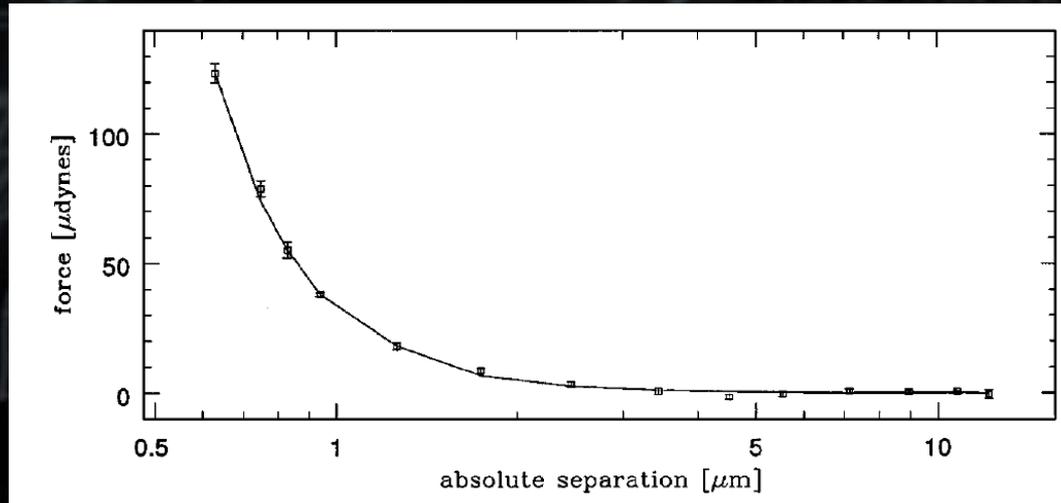
S. K. Lamoreaux\*

*Physics Department, University of Washington, Box 35160, Seattle, Washington 98195-1560*  
(Received 28 August 1996)

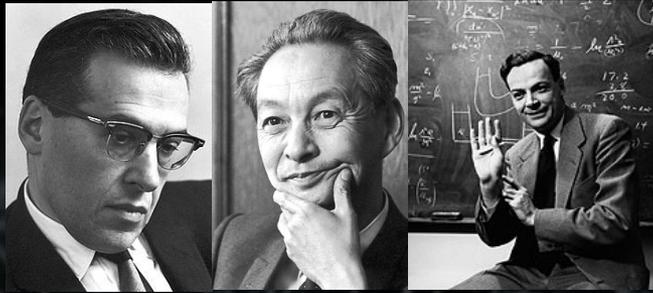
The vacuum stress between closely spaced conducting surfaces, due to the modification of the zero-point fluctuations of the electromagnetic field, has been conclusively demonstrated. The measurement employed an electromechanical system based on a torsion pendulum. Agreement with theory at the level of 5% is obtained. [S0031-9007(96)02025-X]

PACS numbers: 12.20.Fv, 07.07.Mp

# O EFEITO CASIMIR

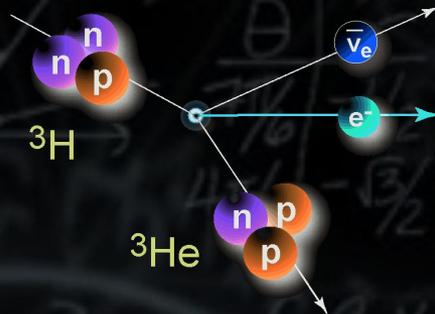


# SITUAÇÃO NO FINAL DOS ANOS 40



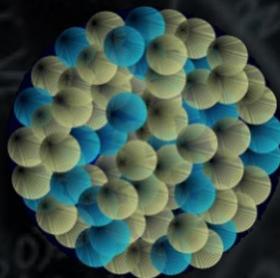
## ELECTRODINÂMICA QUÂNTICA

Teoria quântica dos electrões, positrões, fótons e da interacção electromagnética.  
( $e^-$ ,  $e^+$ ,  $\gamma$ )



## FORÇA FRACA

Teoria do decaimento radioactivo descrita pela interacção de Fermi.



## FORÇA FORTE

Força responsável pela coesão do núcleo descrita pelo potencial de Yukawa.

( $n$ ,  $p$ ,  $\pi$ )