

# $E_{\text{CM}}$ determination at future circular $e^+e^-$ colliders

Guy Wilkinson  
University of Oxford  
ECFA  $e^+e^-$  WG1-PREC meeting, 2/11/22

with many thanks to colleagues in FCC EPOL WG

# $E_{CM}$ determination at ~~future circular $e^+e^-$ colliders~~

FCC-ee

Almost all statements  
equally applicable to CEPC

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# Outline

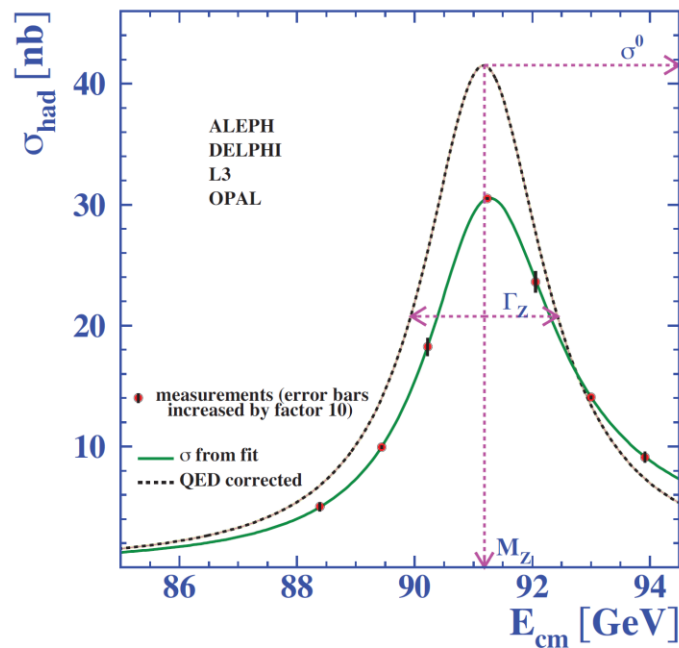
- Lessons from LEP
  - what was achieved, and what were the limitations
- Looking towards FCC-ee
  - a baseline strategy for coping with  $\sim 5 \times 10^{12}$  Z (and physics at other energies), including *selected* highlights and ongoing challenges. Not comprehensive !

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  - what was achieved, and what were the limitations
- Looking towards FCC-ee
  - a baseline strategy for coping with  $\sim 5 \times 10^{12}$  Z (and physics at other energies), including *selected* highlights and ongoing challenges. Not comprehensive !

# Challenges of Z metrology

Measuring  $M_Z$  and  $\Gamma_Z$  at LEP had many challenges (e.g. luminosity measurement, theoretical understanding of lineshape etc.) but knowledge of horizontal scale, i.e.  $E_{CM}$  calibration, was identified as limiting factor right from outset.



Horizontal-scale uncertainty set by knowledge of collision energy, also common between experiments.

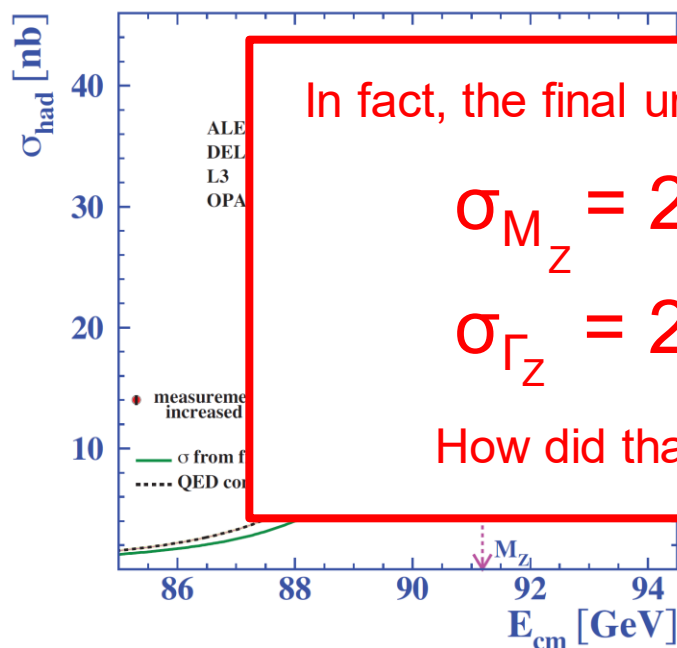
It was guessed that  $\sim 10$  MeV uncertainty *might* be possible.

Outlook shortly before LEP turn on: *“The overall conclusion is that at LEP the  $Z^0$  mass and width can be measured with relative ease down to ... +/- 50 MeV. A factor of 2-3 improvement can be reached with a determined effort...”*

CERN 86-02 ‘Physics at LEP’, ed. Ellis and Peccei.

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# Collision-energy calibration

Knowledge of collision energy leading systematic in mass and width measurement:

$m_Z$  total uncertainty = 2.1 MeV, of which  $E_{CM}$  contribution = 1.7 MeV

$\Gamma_Z$  total uncertainty = 2.3 MeV, of which  $E_{CM}$  contribution = 1.2 MeV

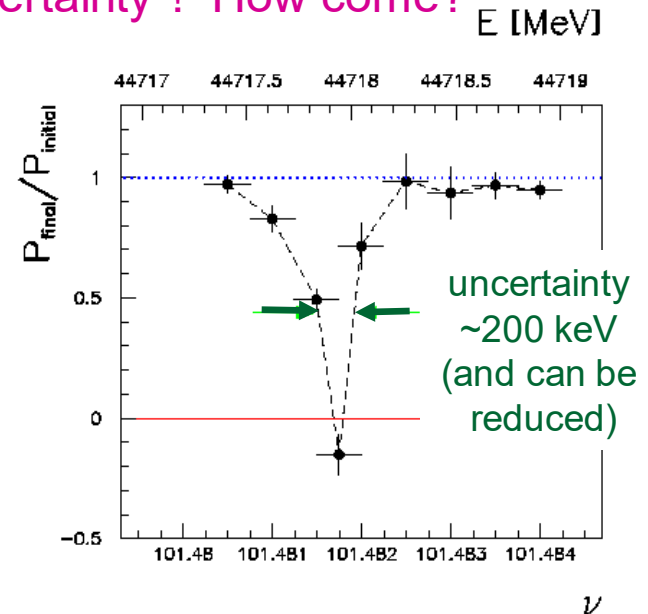
But *much* better than anticipated, and < stat. uncertainty ! How come?

High level of precision achieved through miracle of resonant de-polarisation (RDP), which is unique to circular  $e^+e^-$  machines.

- Wait for transverse polarisation to build up;
- Precession frequency,  $\nu_s$ , directly proportional to  $E_b$ :

$$E_b = 2 \nu_s m_e c^2 / (g_e - 2)$$

- Monitor polarisation with Compton scattering from laser whilst exciting beam with transverse oscillating B field. Find frequency at which depol<sup>n</sup> occurs.



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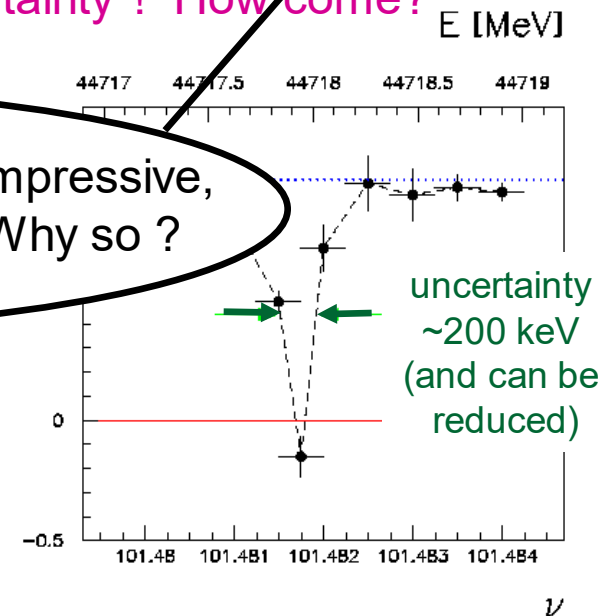
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Hang on, these uncertainties, though impressive, are >> intrinsic uncertainty of RDP. Why so ?

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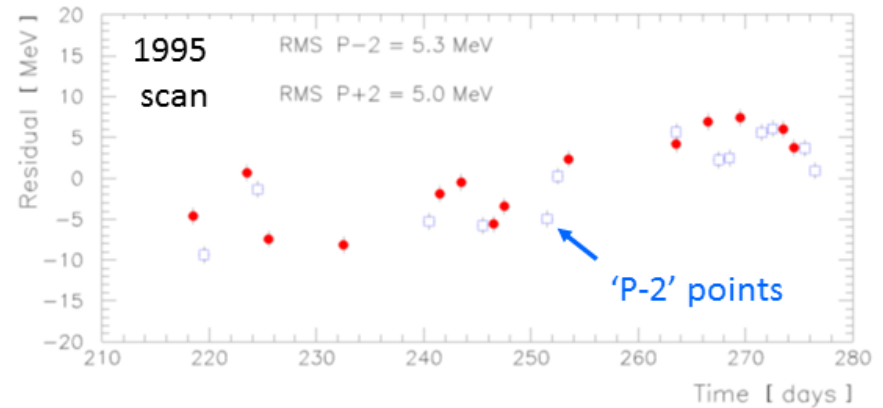
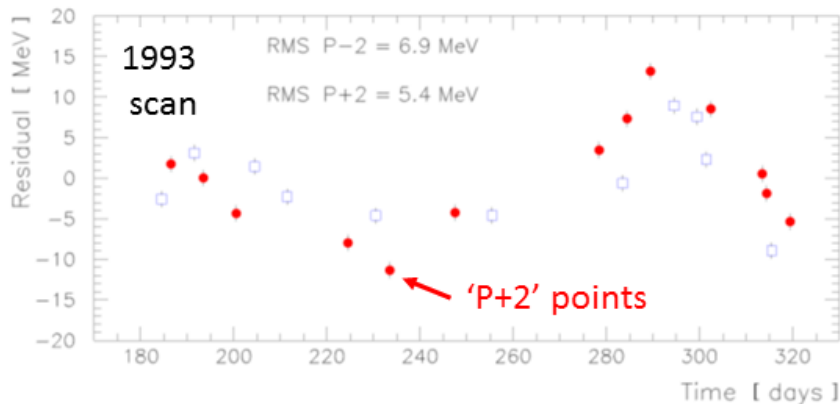
- Monitor polarisation with Compton scattering from laser whilst exciting beam with transverse oscillating B field. Find frequency at which depol<sup>n</sup> occurs.





# Challenge of $E_{CM}$ calibration at LEP

At LEP RDP could not be performed during physics operation. Time-consuming procedure carried out at the end of certain fills, involving dedicated optics. these measurements showed scatter indicating considerable evolution in  $E_b$ .



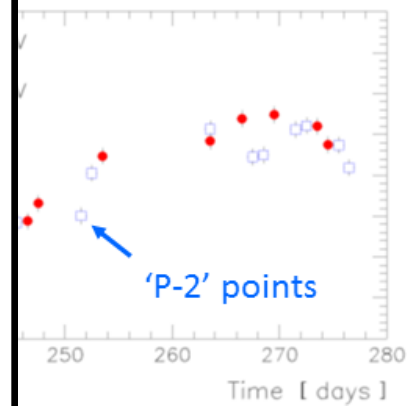
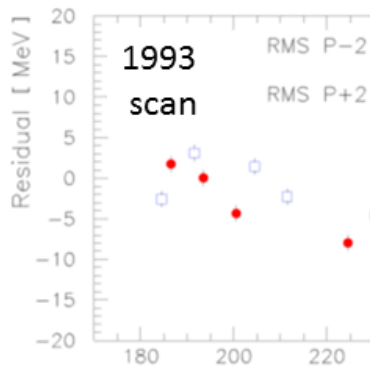
To calibrate the physics data-taking period, necessary to understand and model this evolution – a long and painful process that took many years. Ingredients:

- Bright ideas and machine theory;
- Dedicated instrumentation *e.g.* NMRs in magnets, BPMs *etc.*;
- Lots of machine time for studies (~50 full days in period 1993-2009);
- Mechanisms parameterised in models, used to calibrate physics data periods.

# Challenge of $E_{CM}$ calibration at LEP

At LEP RDP calibration of the centre-of-mass energy is a time-consuming procedure carried out during the 1993-1995 energy scan. These measurements are used to calibrate the centre-of-mass energy for the physics data periods.

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## Calibration of centre-of-mass energies at LEP1 for precise measurements of Z properties

The LEP Energy Working Group

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**Abstract.** The determination of the centre-of-mass energies from the LEP1 data for 1993, 1994 and 1995 is presented. Accurate knowledge of these energies is crucial in the measurement of the Z resonance parameters. The improved understanding of the LEP energy behaviour accumulated during the 1995 energy scan is detailed, while the 1993 and 1994 measurements are revised. For 1993 these supersede the previously published values. Additional instrumentation has allowed the detection of an unexpectedly large energy rise during physics fills. This new effect is accommodated in the modelling of the beam-energy in 1995 and propagated to the 1993 and 1994 energies. New results are reported on the magnet temperature behaviour which constitutes one of the major corrections to the average LEP energy.

The 1995 energy scan took place in conditions very different from the previous years. In particular the interaction-point specific corrections to the centre-of-mass energy in 1995 are more complicated than previously: these arise from the modified radiofrequency-system configuration and from opposite-sign vertical dispersion induced by the bunch-train mode of LEP operation.

Finally an improved evaluation of the LEP centre-of-mass energy spread is presented. This significantly improves the precision on the Z width.

[EPJC 6 (1999) 187]

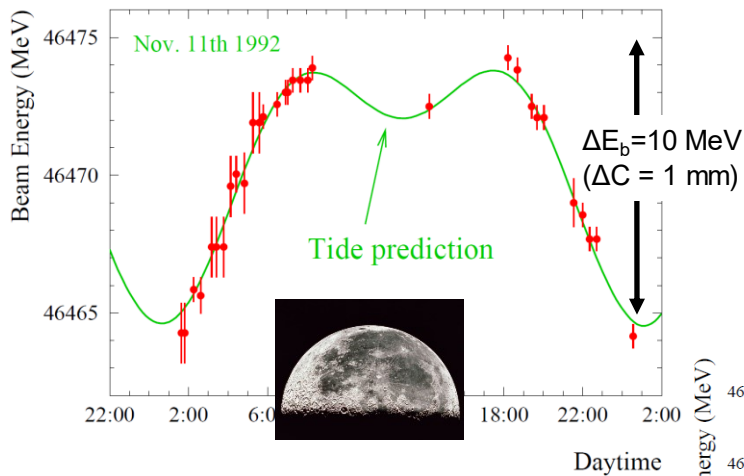
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- Bright idea
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stand and model this evolution –  
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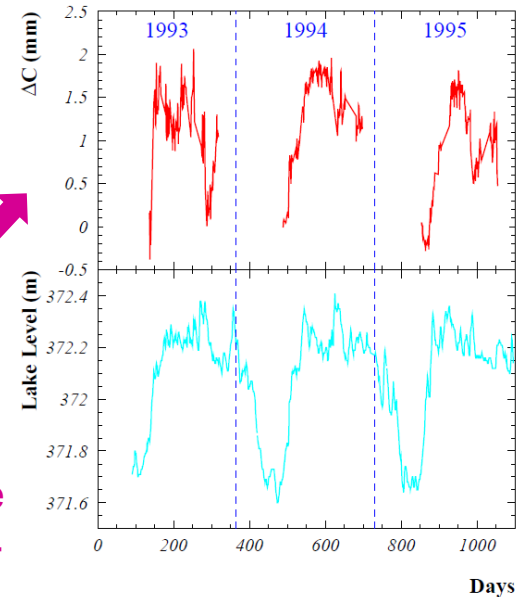
etc.;  
(1993-2009);  
physics data periods.

# Some mechanisms of $E_b$ variation

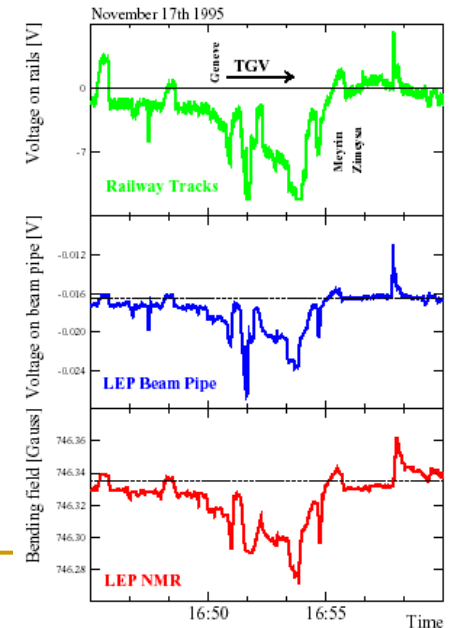
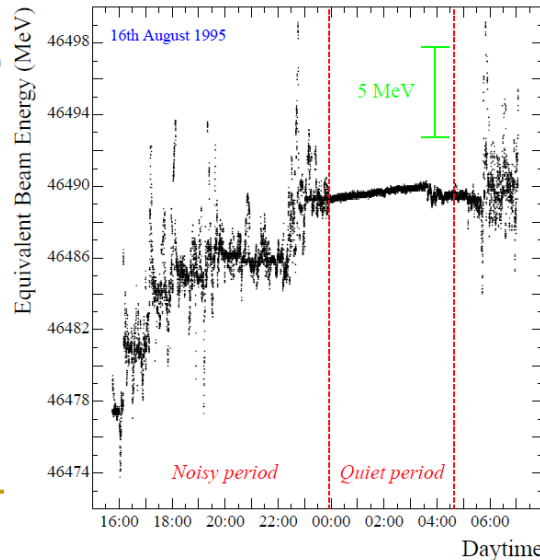


Short- (tide) and long- (lake) term ring distortions.

NB at FCC-ee effects will be  $\sim 10x$  larger due to smaller momentum-compaction factor !



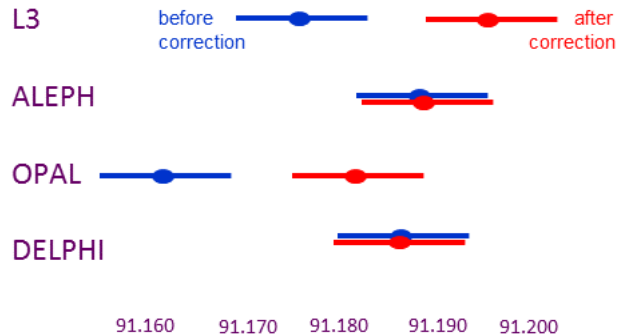
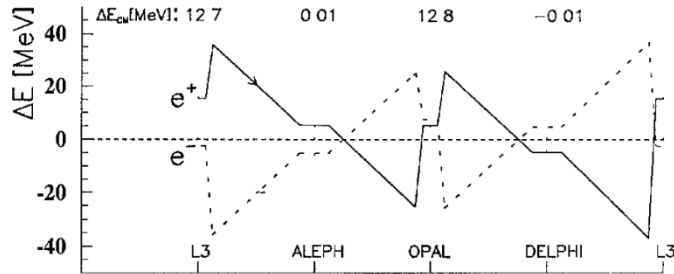
Rise of dipole fields due to stimulation from returning current from TGV.



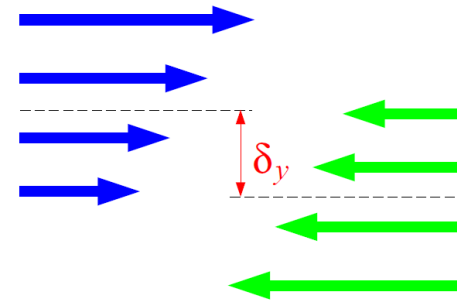
# Interaction-point specific corrections

RDP gives  $E_b$ , but  $E_{CM} \neq 2E_b$ . Interaction-point specific corrections must be applied.

'RF sawtooth' – synchrotron radiation loss and RF boosts.



Opposite-sign vertical dispersion (present in 1995) combined with collision offsets  $\rightarrow E_{CM}$  bias.



$$\Delta E_{cm} = \frac{-1}{2} \frac{\delta y}{\sigma_y^2} \frac{\sigma_{E_b}^2}{E_b} \Delta D_y^*$$

$\sigma_{E_b}$  = Energy spread

$\Delta E_b^*$  = Difference in dispersion between  $e^+$  and  $e^-$

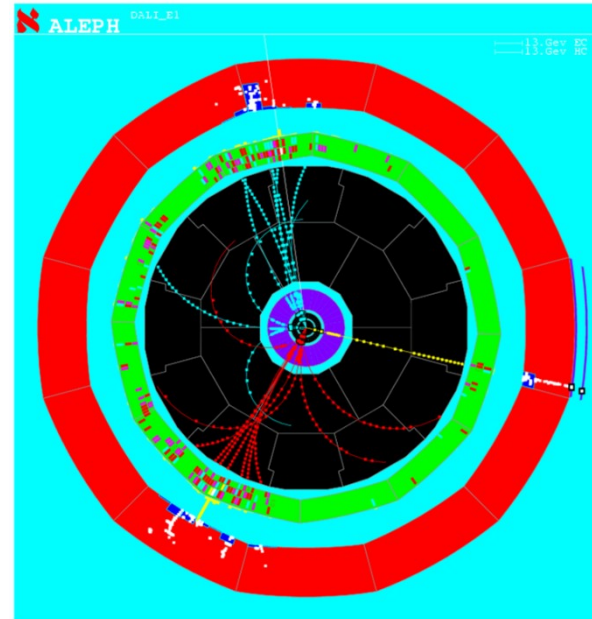
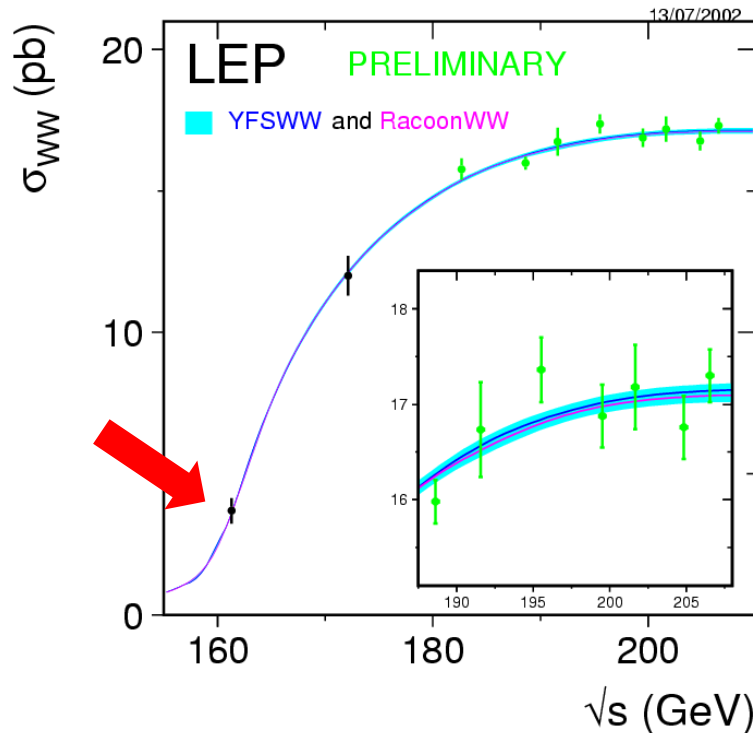
e.g. if  $\Delta D_y^* \sim 2$  mm and  $\delta y = 1$   $\mu$ m  $\rightarrow \Delta E_{CM} = 2$  MeV.

Discovered in 1991, from comparison of individual  $m_Z$  results [[PLB 307 \(1993\) 187](#)].

Control by continual optimisation of luminosity vs  $\delta y$  ('Vernier scans').

# Measuring $m_W$ in $e^+e^- \rightarrow W^+W^-$

Two methods available: measure  $WW$  cross-section at threshold, or fully reconstruct event. Former has fewer systematics, and will probably be the method of choice at FCC-ee, but lower statistical uncertainty gave latter higher weight at LEP.



In both cases a leading systematic uncertainty comes from collision energy (yes, that again).

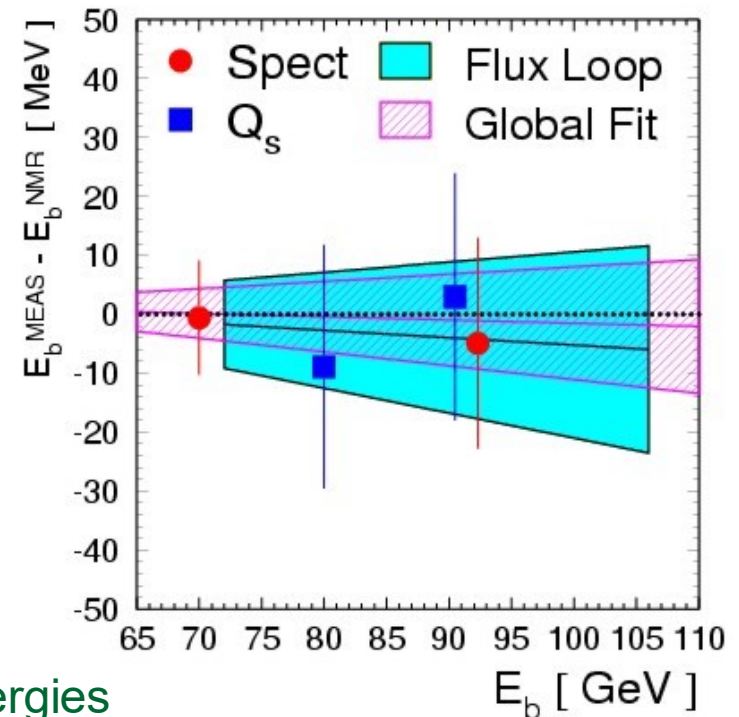
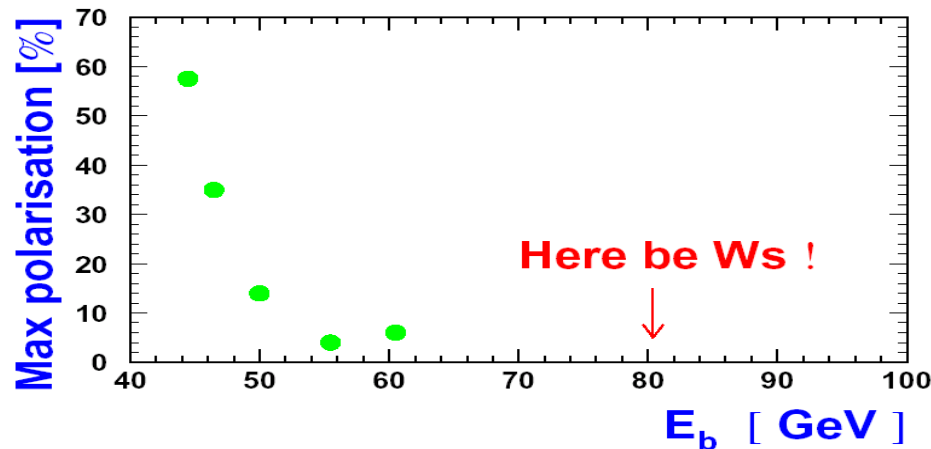
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# Measuring $m_W$ in $e^+e^- \rightarrow W^+W^-$

$$\frac{\Delta m_W}{m_W} = \frac{\Delta E_{CM}}{E_{CM}}$$

Surely not a problem? Many fewer W's than Z's – statistical precision at LEP a few  $10^{-4}$ , and  $E_{CM}$  measured to  $2 \times 10^{-5}$  at  $Z^0$ . What's the worry?

Growth of beam spread with energy means depolarising resonances destroy polarisation. Ergo, no RDP at  $E_b > 60$  GeV at LEP2...



...instead must use a variety of methods (e.g. spectrometer) to extrapolate from RDP energies to  $W^+W^-$  regime. Very difficult, but it was done [EPJC 39 (2005) 253].

# Outline

- Lessons from LEP
  - what was achieved, and what were the limitations
- Looking towards FCC-ee
  - a baseline strategy for coping with  $\sim 5 \times 10^{12}$  Z (and physics at other energies), including *selected* highlights and ongoing challenges. Not comprehensive !

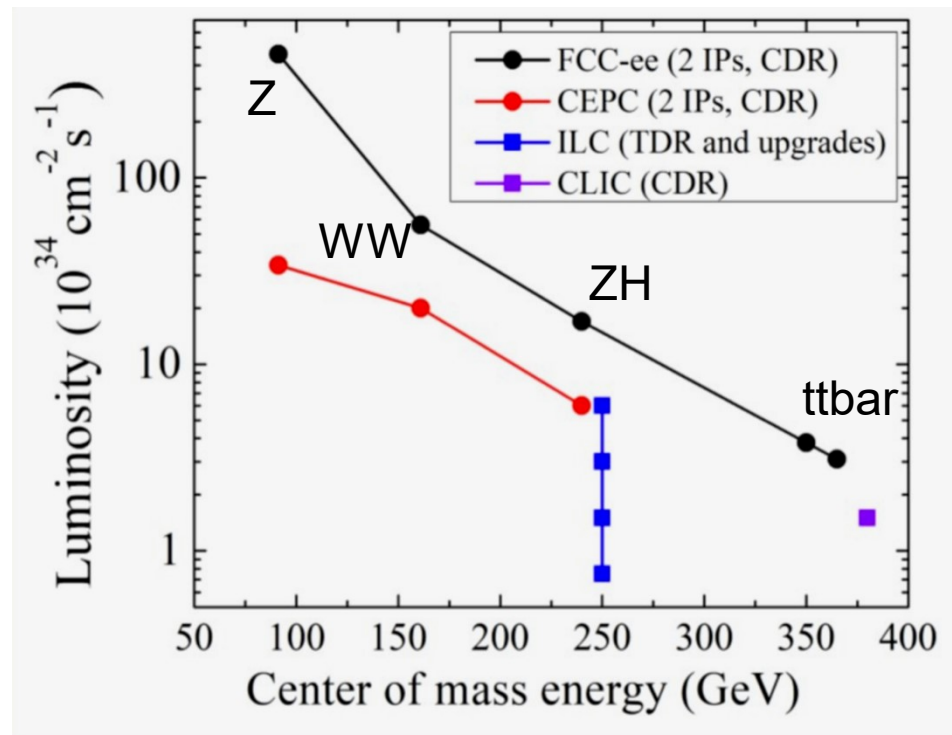
# $E_{CM}$ calibration requirements at FCC-ee

$E_{CM}$  systematics at LEP for  $M_Z$ ,  $\Gamma_Z$  &  $M_W$  were similar to statistical uncertainties. Enormous increase in sample sizes at FCC-ee means corresponding improvements are required for  $E_{CM}$ . Will also be necessary for other observables (e.g.  $A_{FB}^\mu$ ). Big improvements also required in knowledge of related quantities, e.g.  $E_b$  spread.

$5 \times 10^{12}$   $Z^0$  produced  
=  $3 \times 10^5$  x LEP

$10^8$   $W^+W^-$  produced  
= 3000 x LEP

(numbers based on  
CDR running strategy  
& two interaction points)



Will present key points of developing strategy for meeting this challenge. Follow up with some remarks about other  $E_{CM}$  points, including possible  $m_H$  run.



# Strategy for $E_{\text{CM}}$ calibration at FCC-ee – key points

In contrast to LEP, build  $E_{\text{CM}}$  calibration requirements into machine design and planning from start. We already have a baseline strategy, which addresses the main problems that afflicted the LEP energy calibration.

- Perform RDP ‘continuously’ (~5 times per hour). This is done on ~250 out of 16600 non-colliding pilot bunches.



Removes to first order all time-dependent effects !!!

- Measure separately for  $e^+$  &  $e^-$ .
- Adjust RF frequency at short intervals to suppress tide-like effects.

PREPARED FOR SUBMISSION TO JHEP

## Polarization and Centre-of-mass Energy Calibration at FCC-ee

The FCC-ee Energy and Polarization Working Group:

Alain Blondel,<sup>1,2,3</sup> Patrick Janot,<sup>2</sup> Jörg Wenninger<sup>2</sup> (Editors)

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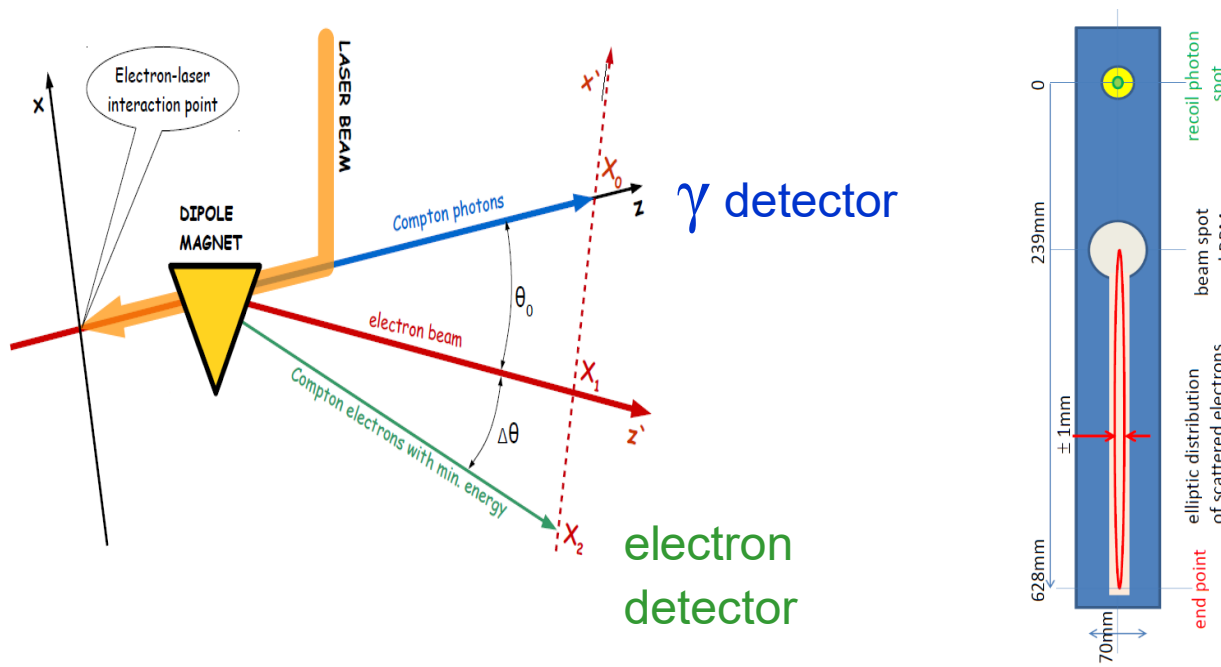
Jorg.Wenninger@cern.ch

arXiv:1909.12245v1 [physics.acc-ph] 26 Sep 2019

[arXiv:1909.12245]

# Polarization measurements

Polarimetry measurement will be based on inverse Compton scattering, where both electrons and photons will be detected, in contrast to LEP where only backscattered photon was measured. Backscattered photon sensitive to only transverse polarization, but electrons can access complete polarization vector.



Need to measure polarization level of electrons and positrons, so (at least) two polarimeters required. Full specifications currently under discussion.

# Resonant depolarization strategy

- Require 5-10% level of polarization for reliable RDP measurement.
- Natural polarization build-up time at Z pole is around 29 hours for 10%. Therefore use wigglers for 1-2 hours at start of fill to polarize pilot bunches, before injecting physics bunches (unavoidable dead time – to be minimised).
- Depolarize one bunch at a time, and rely on natural polarization build-up to replenish polarization level of used bunches.
- Aim to perform measurement every ~10 minutes.
- Under consideration: also performing free spin precession measurements, where spin vector is rotated into horizontal plane and its precession frequency is directly measured. May have complementary systematics to RDP.

# Longitudinal polarization considerations

Any residual longitudinal polarization will bias cross sections & forward-backward asymmetries (indeed, high longitudinal polarization is actually useful, but we assume we are not in that regime – rather longitudinal polarization is a nuisance).

Consider forward-backward asymmetry of  $b\bar{b}$  at Z pole:  $A_{\text{FB}}^b = \frac{3}{4} \mathcal{A}_e \mathcal{A}_b$

where in the SM  $\mathcal{A}_e \approx 0.15$ ,  $\mathcal{A}_b \approx 0.95 \Rightarrow A_{\text{FB}}^b \approx 0.11$

Now, if there is longitudinal polarization, asymmetry becomes:  $(A_{\text{FB}}^b)' = \frac{3}{4} \mathcal{A}'_e \mathcal{A}_b$

where  $\mathcal{A}'_e = -\left(\frac{\mathcal{A}_e - P}{1 - \mathcal{A}_e P}\right)$  with  $P = \frac{(P_z)_{e^-} - (P_z)_{e^+}}{1 - (P_z)_{e^-} (P_z)_{e^+}}$

and  $(P_z)_{e^\pm}$  the longitudinal polarization of the  $e^\pm$ .

# Longitudinal polarization considerations

Any residual longitudinal polarization will bias cross sections & forward-backward asymmetries (indeed, high longitudinal polarization is actually useful, but we assume we are not in that regime – rather longitudinal polarization is a nuisance).

So, if  $(P_Z)_{e^-} = (P_Z)_{e^+}$  (no reason to be so) =  $10^{-5}$  (ballpark guess)

$$P = 2 \times 10^{-5} \implies \frac{(A_{FB}^b)' - A_{FB}^b}{A_{FB}^b} = 1.3 \times 10^{-4}$$

Statistical uncertainty on  $A_{FB}^b$  around  $2 \times 10^{-5}$  (relative), and QCD uncertainty which will probably be larger. Still, to be safe we would want to control  $P_Z$  to  $< 10^{-5}$ .

- Likely strategy:
- 1) Measure polarization levels of physics bunches, as well as pilot bunches. (As transverse polarization  $\gg$  longitudinal, a  $10^{-3}$  measurement of former is probably sufficient.)
  - 2) If necessary, continually depolarize physics bunches.

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Removes to first order all time-dependent effects !!!

- Measure separately for  $e^+$  &  $e^-$ .
- Adjust RF frequency at short intervals to suppress tide-like effects.
- Frequent scans to suppress dispersion biases at IP.

PREPARED FOR SUBMISSION TO JHEP

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[arXiv:1909.12245]

# Interaction-point specific corrections

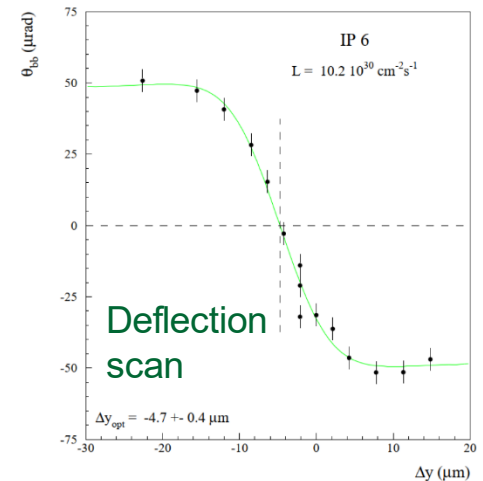
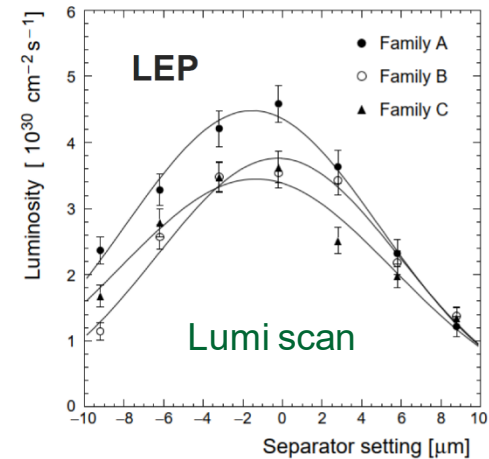
As at LEP, we will need to worry about possible dispersion and RF sawtooth effects.

$$\Delta\sqrt{s} = -u_0 \frac{\sigma_E^2 \Delta D^*}{E_0 \sigma_u^2}$$

So, for  $\Delta D^* = 10 \mu\text{m}$ ,  $\Delta E_{\text{CM}} \sim 1 \text{ MeV} / \text{nm}$ .

Therefore must keep offset to  $\ll 1 \text{ nm}$  (at least, on average) & also measure dispersion of beams (NB it is *difference* in dispersion that matters).

Both offset and dispersion can be measured through luminosity scan or through beam-beam deflection scan, using angles found from BPMs.

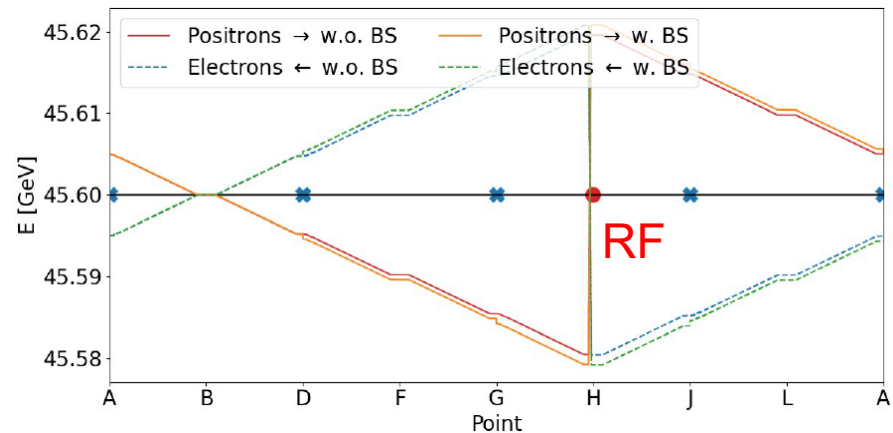


# Interaction-point specific corrections

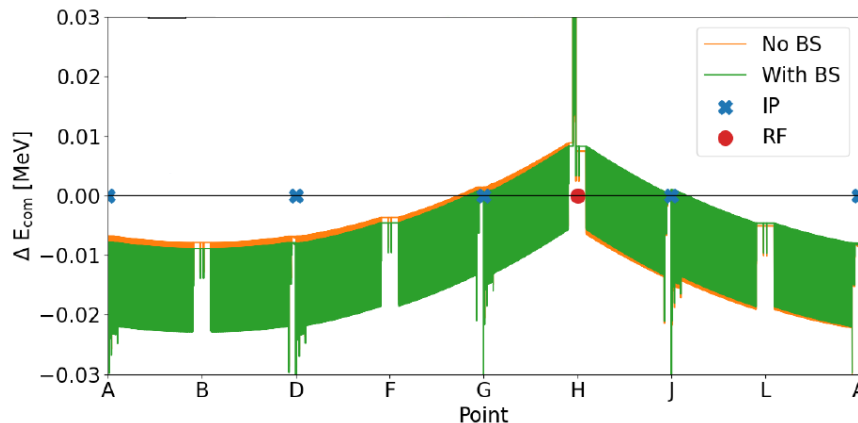
As at LEP, we will need to worry about possible dispersion and RF sawtooth effects.

Optimal configuration for minimizing shift in  $E_{CM}$  around ring is to have single location for RF station (but for  $t\bar{t}$  running, two are necessary).

Can monitor through  $e^+ e^-$  separation in BPMs, and through measurement of boost in experiments (see later).



J. Keintzel



IP	$\Delta E_{CM}$ [keV]	Boost [MeV]
PA	- 7.851	10.665
PD	- 7.931	- 10.108
PG	0.570	- 30.883
PJ	0.844	31.439



# Strategy for $E_{\text{CM}}$ calibration at FCC-ee – key points

In contrast to LEP, build  $E_{\text{CM}}$  calibration requirements into machine design and planning from start. We already have a baseline strategy, which addresses the main problems that afflicted the LEP energy calibration.

- Perform RDP ‘continuously’ (~5 times per hour). This is done on ~250 out of 16600 non-colliding pilot bunches.



Removes to first order all time-dependent effects !!!

- Measure separately for  $e^+$  &  $e^-$ .
- Adjust RF frequency at short intervals to suppress tide-like effects.
- Frequent scans to suppress dispersion biases at IP.
- Invest in extensive instrumentation and logging of all machine parameters.
- Exploit measurements from experiments that provide complementary information.

PREPARED FOR SUBMISSION TO JHEP

## Polarization and Centre-of-mass Energy Calibration at FCC-ee

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# Measurement of crossing angle and $E_{CM}$ spread

Even with energy loss / RF and dispersion effects under control, another ingredient is needed to go from the local beam energies to  $E_{CM}$ . The crossing angle !

$$\sqrt{s} = 2\sqrt{E_{e^+}E_{e^-}} \cos \alpha/2$$

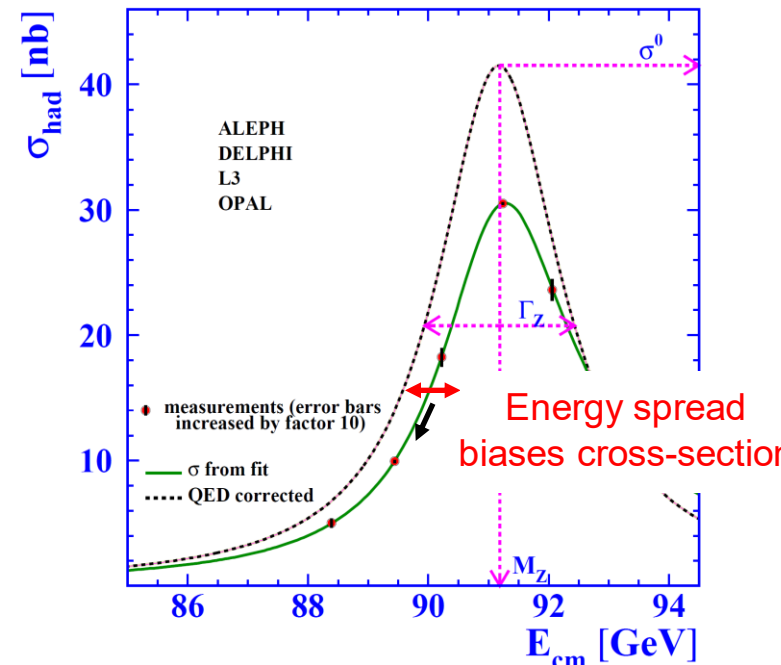
Crossing angle  $\alpha$  is around 30 mrad at FCC-ee (in fact there are beam-beam effects that modify both local beam energy and  $\alpha$  - won't discuss through lack of time).

Aside from this, need to know  $E_{CM}$  spread well. Not a great concern at LEP but a dangerous systematic on several measurements at FCC-ee (e.g. Z, W, top widths, offpeak  $A_{FB}^\mu$  for  $\alpha_{QED}$ ).

Beam energy is not monochromatic, but has a spread of  $\sim 50$  MeV at Z.

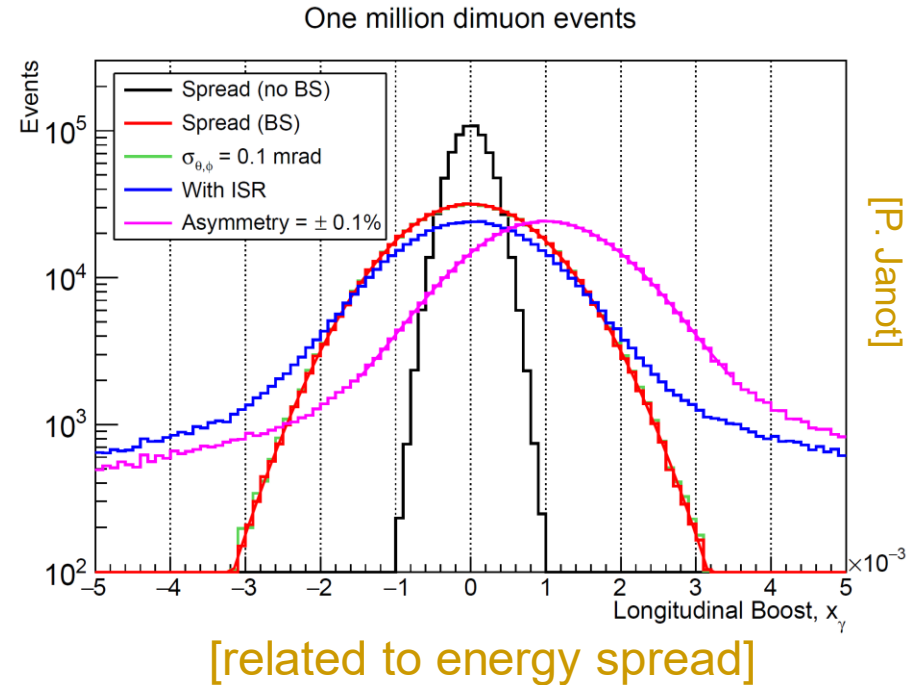
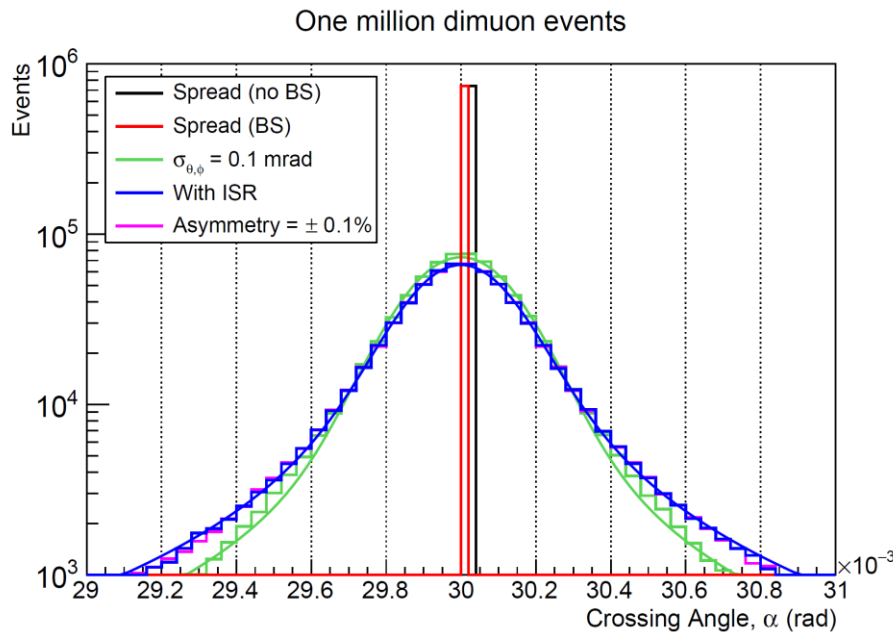
Spread in collision energy,  $\sigma_{E_{CM}}$  will shift cross-section measurements by  $\delta_\sigma$  as line shape is (clearly!) not linear.

$$\delta\sigma = -0.5 \frac{d^2\sigma}{dE^2} \sigma_{E_{CM}}^2$$



# Measurement of crossing angle and $E_{CM}$ spread

These effects can be controlled to necessary precision through monitoring topology of  $Z^{(*)} \rightarrow \mu\mu(\gamma)$  events, of which million will be collected every  $\sim 5$  minutes at Z pole.



Results documented in [arXiv:1909.12245](https://arxiv.org/abs/1909.12245), but should be revisited with more sophisticated simulations, consideration of ISR knowledge, detector resolution *etc.*

# Initial estimates of what is achievable at Z

Table compiled for [arXiv:1909.12245](https://arxiv.org/abs/1909.12245) for key Z observables

Observable	statistics	$\Delta\sqrt{s}_{\text{abs}}$ 100 keV	$\Delta\sqrt{s}_{\text{syst-ptp}}$ <b>40 keV</b>	calib. stats. 200 keV/ $\sqrt{N^i}$	$\sigma_{\sqrt{s}}$ 85 $\pm$ <b>0.05</b> MeV
$m_Z$ (keV)	4	100	<b>28</b>	1	–
$\Gamma_Z$ (keV)	4	2.5	<b>22</b>	1	<b>10</b>
$\sin^2 \theta_W^{\text{eff}} \times 10^6$ from $A_{\text{FB}}^{\mu\mu}$	2	–	<b>2.4</b>	0.1	–
$\frac{\Delta\alpha_{\text{QED}}(m_Z^2)}{\alpha_{\text{QED}}(m_Z^2)} \times 10^5$	3	0.1	<b>0.9</b>	–	<b>0.1</b>

Much better than LEP, but  $E_{\text{CM}}$  systematics do not match stat precision for  $M_Z$  &  $\Gamma_Z$ .

We aim to improve on this ! Many advances reported at [2<sup>nd</sup> FCC EPOL Workshop](#) at CERN in September, and further updates will be presented at FCC Midterm Review in summer next year. Help is very much welcome.

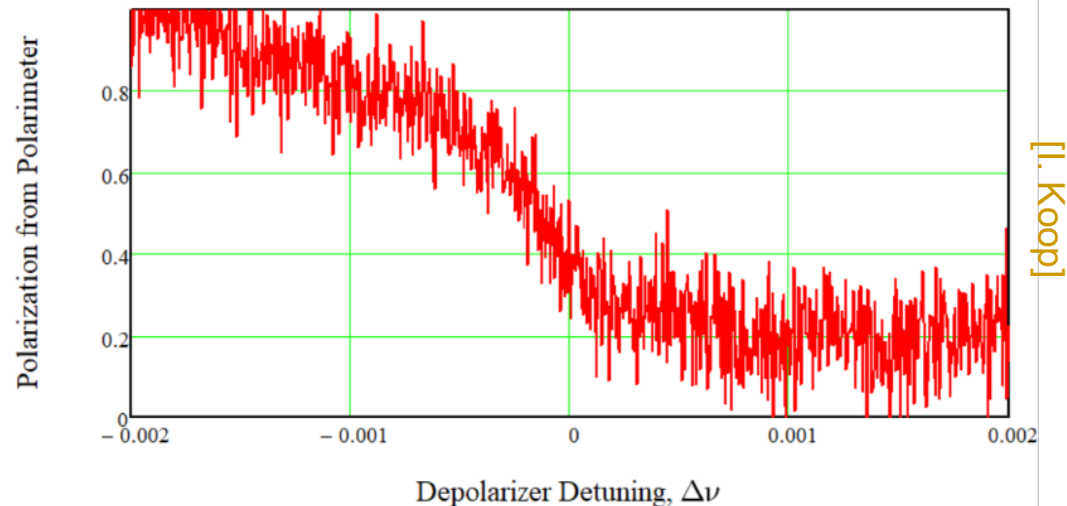
# $E_{\text{CM}}$ calibration for $m_W$

In contrast to LEP, at FCC-ee resonant depolarisation should be possible at the  $W^+W^-$  threshold ( $E_b \sim 80$  GeV). This is because the beam-energy spread

$$\sigma_{E_b} \sim E_b^2 / \sqrt{\rho}$$

and the magnetic bending radius  $\rho$  at FCC-ee is larger than at LEP. (Furthermore, the improvements in instrumentation will be helpful in increasing polarization levels.)

Recent simulations confirm that polarization, and clear depolarization signals, will be achievable. Although precision is unlikely to match that at Z, it should be perfectly adequate for  $m_W$ .



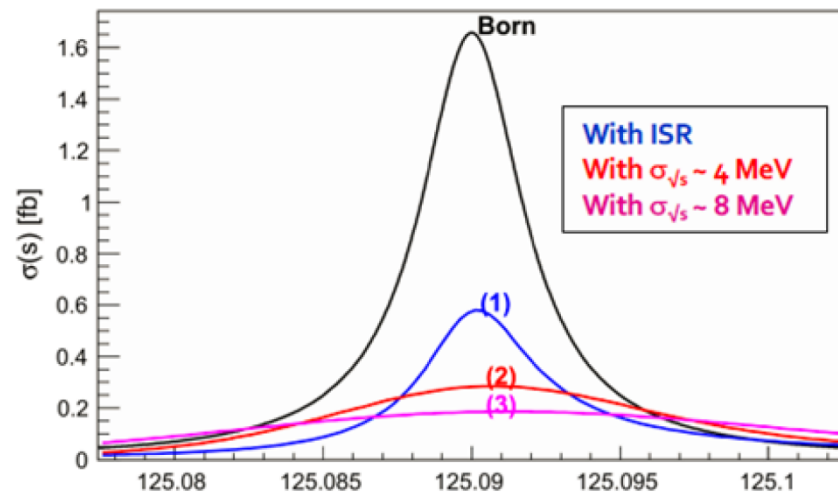
# Other energy points

It is assumed that RDP will not be possible when collecting ZH and  $t\bar{t}$  data.

Here can rely on  $Z^* \rightarrow f\bar{f}\gamma$  events and  $e^+e^- \rightarrow W^+W^-$  events, which will have sufficient statistical power to meet the physics goals at both energy points.

And what about if a run is made at  $E_{\text{CM}} = m_H$  to measure electron Yukawa ?

- Monochromatization is required to reduce  $E_{\text{CM}}$  spread to same order as Higgs width ( $\sim 4$  MeV)
- Need to ensure high stability for  $E_{\text{CM}}$  and very good knowledge of its value event-to-event.



→ RDP certainly needed, with additional requirements different to those at Z !

# Conclusions

- Precise knowledge of  $E_{\text{CM}}$  is a key asset of circular  $e^+e^-$  colliders.
- The Z-scan campaigns at LEP demonstrated both the power of RDP, and also the additional challenges that are required to meet physics goals.
- The huge sample sizes at FCC-ee mandates a corresponding improvement in the control of  $E_{\text{CM}}$ . A baseline strategy has been developed which already shows great promise....
- ...but much work remains to understand requirements better, firm up specifications, and pursue further developments to improve outlook further.
- RDP will also be available in  $W^+W^-$  regime – opportunity then available to improve precision of CDF measurement by factor 20-30.

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# Backups

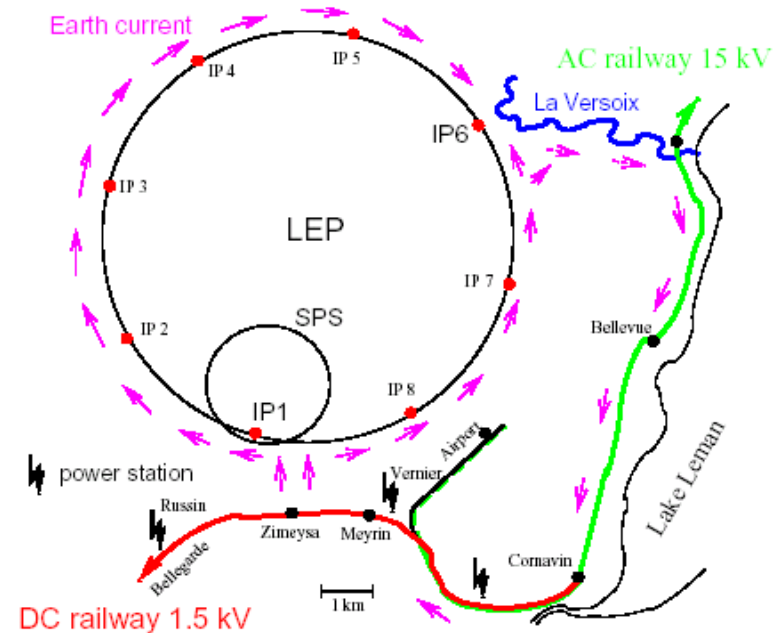
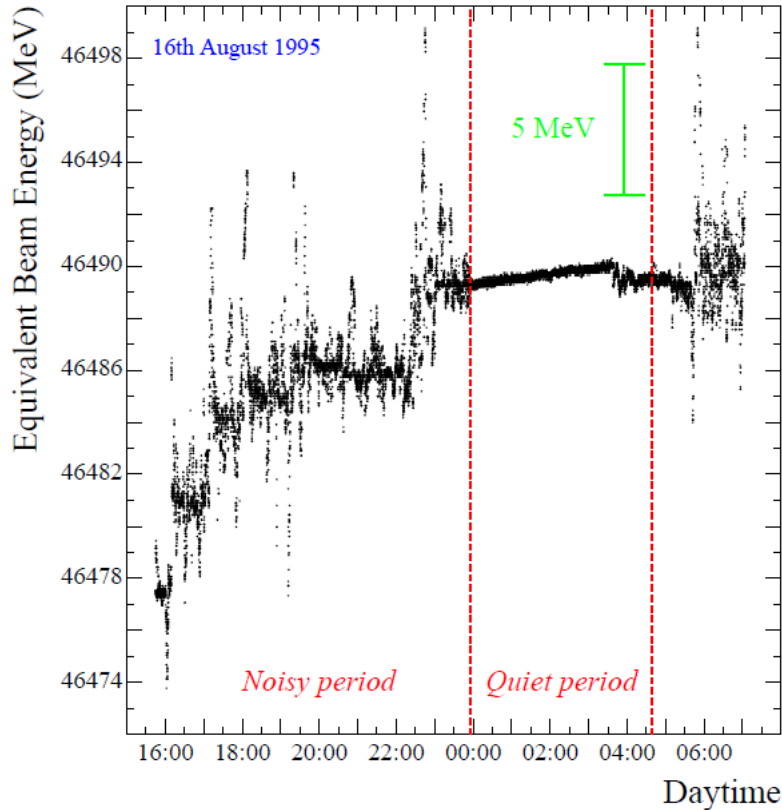
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# (Selected) mechanisms of $E_b$ variation

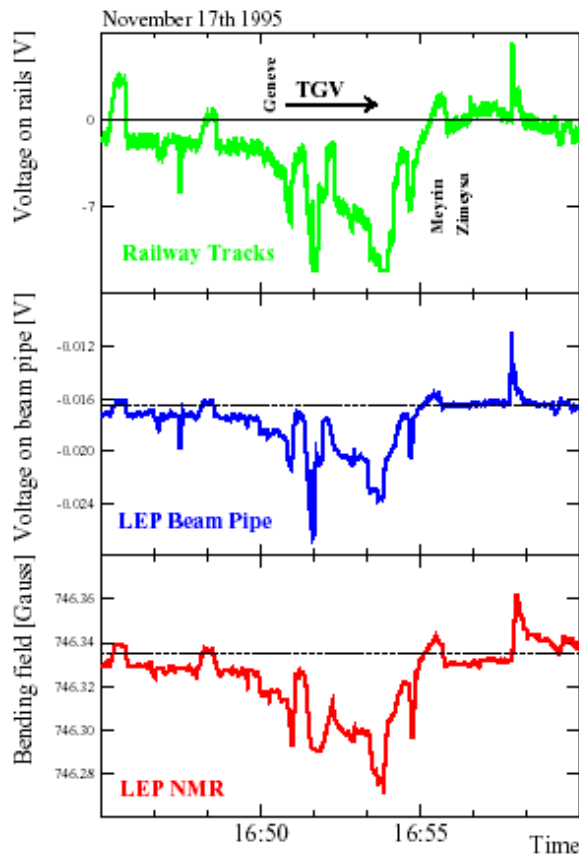
Strange noise and field rises in magnets correlated to time of day and time in fill.

Found to be due to magnets being 'tickled' by current on beam pipe from passing trains.

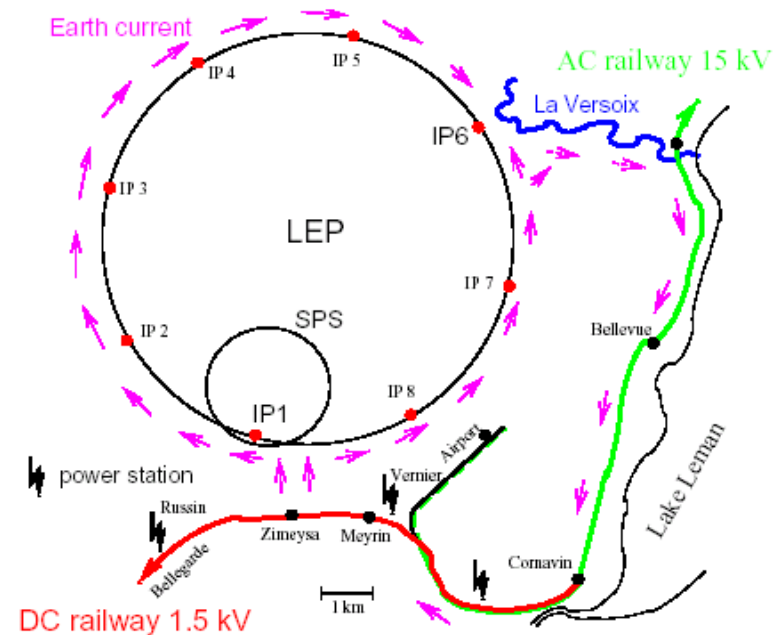


# (Selected) mechanisms of $E_b$ variation

Strange noise and field rises in magnets correlated to time of day and time in fill.



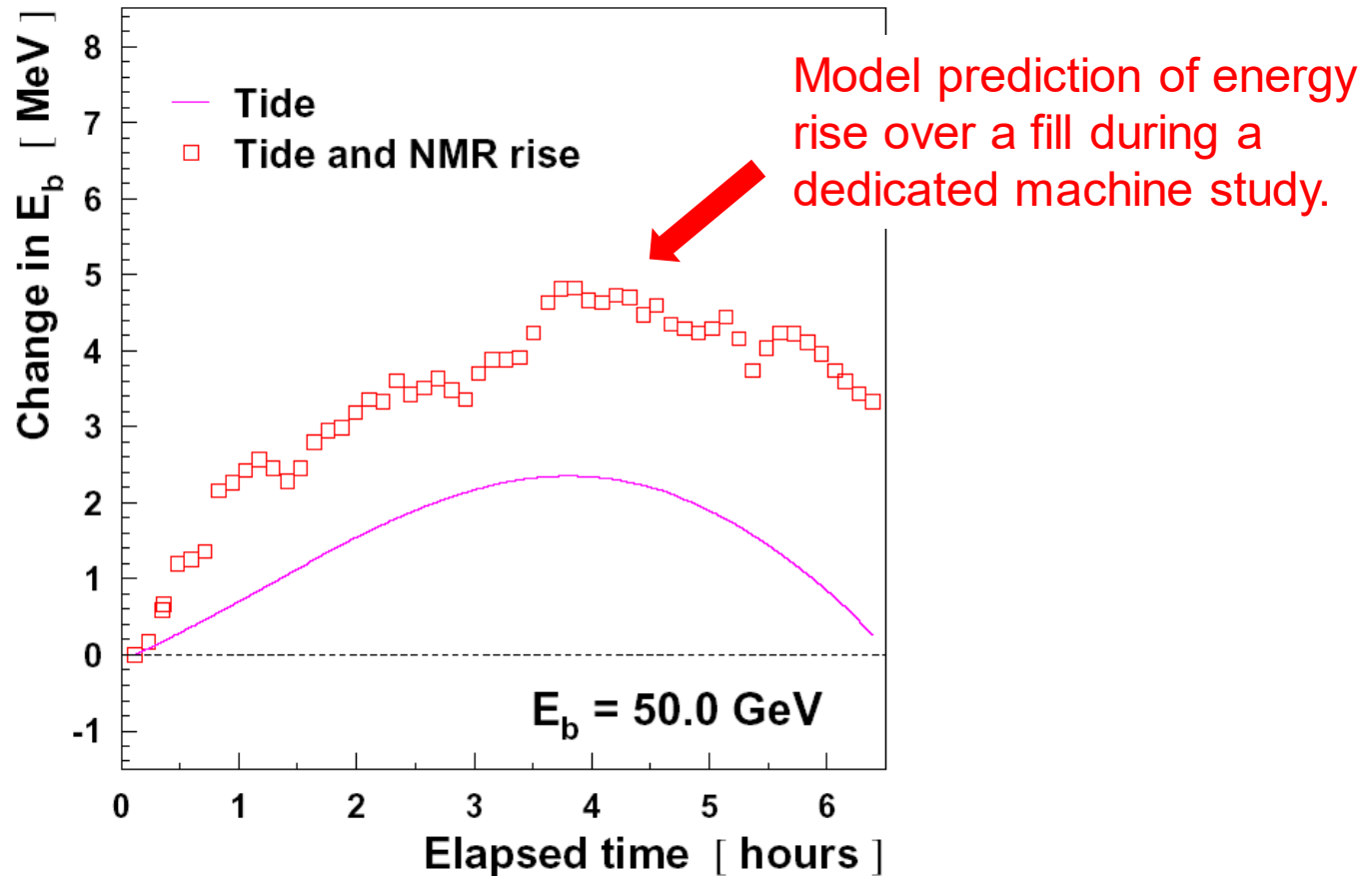
Found to be due to magnets being 'tickled' by current on beam pipe from passing trains.



Compelling correlation between current on track, on beam pipe & noise in magnets.

# (Selected) mechanisms of $E_b$ variation

Energy rise modelled with great precision.



# (Selected) mechanisms of $E_b$ variation

Energy rise modelled with great precision, in excellent agreement with RDP.

