Review on Fragmentation Functions

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IWHSS 2011 Paris, April 4



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Fragmentation functions

Describe the collinear transition of a parton *i* into a massless hadron *h* carrying momentum fraction *z*



Represent the probability that a parton hadronizes

- Time-like version of PDFs : information on structure of the nucleon
- Relevant any time a hadron is produced in high energy collisions



Why care about hadrons?

- Hadrons : dominant final state in pp and ep collisions
- Heavy Ion Collisions : QGP by suppression in hadron-production
- The flavor-separation of polarized parton distributions depends crucially on SIDIS results (Hermes-SMC-Compass)
- Extraction of polarized gluon density depends much on PHENIX and STAR pion data

Different FFs very different results for spin dependent distributions DdeF, Navarro, Sassot (2005) wrong D_f wrong Δf

- Look into "unconventional targets" : Lambda polarized fragmentation functions
- Key elements for Target Fragmentation: evolution of fracture functions
- Fundamental role in understanding single spin asymmetries, transversity...
- TH challenge: NLO pQCD fails to describe fixed target experiments!

Some properties of $D^h_i(z,\mu^2)$

• Non-perturbative objects but universal : can be obtained from global analysis!

• **Depend** on energy fraction $z = \frac{E_h}{E_c}$

• **Depend** on factorization scale μ^2 : described by AP equations

 $\frac{d}{d\ln\mu^2}D^h_q(z,\mu^2) = \begin{bmatrix} P_{qq}\otimes D^h_q + P_{gq}D^h_g \end{bmatrix}(z,\mu^2) \qquad \begin{array}{l} \text{NLO}\\ \text{partial NNLO: Moch, Vogt} \end{array}$

Kernels very singular at small z :negative FFs ! $P_{gg} \rightarrow \frac{2C_A}{z} - \frac{\alpha_s}{2\pi} \frac{4C_A^2}{z} \ln^2 z$

• Anyway, mass and higher twist effects not trivial : no systematic treatment

Limits use to z > 0.05 / 0.1

• Energy- Momentum conservation sum rule

$$\sum_{h} \int_{0}^{1} z D_i^h(z,\mu^2) dz = 1$$

without much practical use

Fragmentation Functions

How to obtain them?

Global Analysis : look at different observables



e+e- (SIA) single-inclusive annihilation

$$e^+e^- \to (\gamma, Z) \to H$$

• Cross section depends on two structure functions

$$\frac{1}{\sigma_{tot}}\frac{d\sigma^H}{dz} = \frac{\sigma_0}{\sum_q \hat{e}_q^2} \left[2F_1^H(z,Q^2) + F_L^H(z,Q^2)\right]$$

$$\sigma_{tot} = \sum_{q} \hat{e}_{q}^{2} \sigma_{0} \left[1 + \frac{\alpha_{s}(Q^{2})}{\pi} \right]$$

• At NLO they can be written as

$$2F_{1}^{H}(z,Q^{2}) = \sum_{q} \hat{e}_{q}^{2} \left\{ \begin{bmatrix} D_{q}^{H}(z,Q^{2}) + D_{\bar{q}}^{H}(z,Q^{2}) \end{bmatrix} + \frac{\alpha_{s}(Q^{2})}{2\pi} \begin{bmatrix} C_{q}^{1} \otimes (D_{q}^{H} + D_{\bar{q}}^{H}) & +C_{g}^{1} \otimes D_{g}^{H} \end{bmatrix} (z,Q^{2}) \right\}$$
Known up to NNLO
Coeff. functions
Calculable in pQCD

• Notice that SIA can only give information on the sum $D_q^H(z,Q^2) + D_{ar q}^H(z,Q^2)$

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$$e^{+}$$

$$e^{-}$$

$$Y$$

$$z = \frac{2E_{h}}{Q}$$

e+e- (SIA) single-inclusive annihilation

Advantages • Very precise data from LEP/SLD

Heavy quark tagged data

Only fragmentation functions enter (clean process)

Disadvantages

SIA data dominated by precise LEP/SLD measurements at M_Z

• weak scale dependence (bad resolution for g fragmentation)

mostly determine "singlet" distribution (high precision)

 $\Sigma = D_u + D_{\bar{u}} + D_d + D_{\bar{d}} + D_s + D_{\bar{s}} + D_c + D_{\bar{c}} + D_b + D_{\bar{b}}$

OPAL offers "flavor separated data": no trivial interpretation at NLO

onot precise at large z (relevant for pp collisions)

• Can not separate $D_q^h(z,Q^2)$ from $D_{\overline{q}}^h(z,Q^2)$



current fragmentation

• Distributions in x and z

at LO
$$2F_1^H(x, z_H, Q^2) = \sum_{q, \overline{q}} e_q^2 \quad q(x, Q^2) D_q^H(z_H, Q^2)$$
 expression known up to NLO
"effective charge": different for quarks and antiquarks!

Most valuable observable for flavor/charge separation! Different targets (and x!) can be used



Other possible observables in DIS:

- photoproduction (involves photon pdfs)
- transverse momentum distributions (HERA, EXP+TH uncertainties)

check



Advantages :

allows flavor/charge separation

🖗 larger z

smaller Q², improves scale coverage in evolution : D_g

- relevant for Hermes and Compass kinematics (spin physics)
- almost no HQ fragmentation

Disadvantages :

introduce "dependence" on pdfs (but unpolarized pdfs well constrained from DIS in the same kinematical range, s?)

non-perturbative corrections at small Q²?

almost no HQ fragmentation

Hadron-Hadron collisions

Transverse momentum distribution hard scale $d\sigma(pp \rightarrow h) = \sum_{i,i,k} \int_{0}^{1} dx_{1} f_{i}^{P}(x_{1}, \mu_{FI}^{2}) \int_{0}^{1} dx_{2} f_{j}^{P}(x_{2}, \mu_{FI}^{2}) \int_{0}^{1} dz D_{k}^{h}(z, \mu_{FF}^{2}) d\hat{\sigma}(ij \rightarrow k)$ Hadron-Hadron collisions



Hadron-Hadron collisions

Transverse momentum distribution
hard scale $p \rightarrow f_i^{P(x_1)}$ $q \bar{q} \rightarrow g$ $p \rightarrow g$ $g g \rightarrow g$ $g = f_j^{P(x_1)}$ $g g \rightarrow g$ $g = f_j^{P(x_2)}$ $g g \rightarrow g$ $g = g g \rightarrow g$ $g g \rightarrow g$ $g = g g \rightarrow g$ $g g \rightarrow g$

$$d\sigma(pp \to h) = \sum_{i,j,k} \int_0^1 dx_1 f_i^P(x_1, \mu_{FI}^2) \int_0^1 dx_2 f_j^P(x_2, \mu_{FI}^2) \int_0^1 dz D_k^h(z, \mu_{FF}^2) d\hat{\sigma}(ij \to k)$$

Advantages : allows (very little) flavor/charge separation (gluon dominated) several subprocesses much larger z large (direct!) contribution from D_g relevant for RHIC kinematics (polarized/heavy ion) almost no HQ fragmentation

Disadvantages: problems for fixed target experiments, use only colliders otherwise Threshold resummation needed (not in first approach) larger TH uncertainty (scale dependence) DdeF, W.Vogelsang larger dependence on pdfs almost no HQ fragmentation

NLO analyses for light hadrons

CGGRW(1994), BFGW(2000)

Bourhis et al (2001)

BKK(1995), KKP(2000), AKK(2005,2008)*

Albino, Kniehl, Kramer (1995-2008)

KRE(2000)**

Kretzer (2000)

HKNS(2007)

Hirai, Kumano, Nagai, Sudoh (2007)

use only SIA: no charge separation or ad-hoc assumption

$$D_{\overline{q}}^{h^+}(z,Q^2) = (1-z) D_q^{h^+}(z,Q^2)$$

*AKK2008 includes pp data (with some charge separation)

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First global analysis of fragmentation functions : DSS (2007) DdeF, R.Sassot, M.Stratmann

Global FIT

Advantage : Constrain FF with almost all available data Check of pQCD framework Each experiment spans a different kinematical range Precise determination of distributions and first estimation of uncertainties No need for naive relations between observables and FF's

Disadvantage : more work required !

Feasible if Mellin technique used Tension between different observables requires careful analysis

- Complicated observables:
- without simple NLO interpretation
- involving MC
- "averaged" bins
- large scale dependence
- Weights

DSS global analysis

fragmentation functions for $\pi^+, \pi^-, \pi^0, K^+, K^-, K^0, \overline{K}^0, p, \overline{p}, n, \overline{n}, h^+, h^$ residual $h^+ = \pi^+ + K^+ + p + res^+$

LO and NLO global fits available

SIA data includes: TPC, TASSO, SLD, ALEPH, DELPHI, OPAL + "flavor" tag

SIDIS data from : HERMES, EMC(h) $Q > 1 \,\mathrm{GeV}$

pp data from : PHENIX, STAR, BRAHMS, CDF, UAI, UA2 $p_T > 1 \text{GeV}$ (η) colliders

Estimation of uncertainties using Lagrange multipliers

Technical details

• Flexible parametrization

$$D_i^H(z, Q_0^2) = N_i \, z^{\alpha_i} (1-z)^{\beta_i} \left[1 + \gamma_i (1-z)^{\delta_i} \right]$$

at initial scale

$$Q_0^2 = 1 \operatorname{GeV}^2 \quad u, d, s, g$$
$$Q_0^2 = m_Q^2 \quad c, b$$



with

 α_s and Λ_{QCD} from MRST 2002

- Normalizations for different experiments (if not included in syst.)
- Try to avoid Isospin symmetry assumptions

 Allowing for possible breaking of SU(3) of sea and SU(2) in favored distributions

•unless data can not discriminate for unfavored fragmentations

$$D_{d+\bar{d}}^{\pi^{+}} = N D_{u+\bar{u}}^{\pi^{+}}$$
$$D_{s}^{\pi^{+}} = D_{\bar{s}}^{\pi^{+}} = N' D_{\bar{u}}^{\pi^{+}}$$

$$D_{\bar{u}}^{\pi^+} = D_d^{\pi^+}$$
$$D_{\bar{u}}^{K^+} = D_s^{K^+} = D_d^{K^+} = D_{\bar{d}}^{K^+}$$

D's obtained by global fit : χ^2 minimization



Main problem: how to perform so many evaluations in short time

Solution: work in Mellin space $D_i^n(\mu^2) = \int_0^1 dz \, z^{n-1} \, D_i(z, \mu^2)$

Convolutions become products : simple solution for evolution equations
 Back to "z" space can be done fast
 Rather natural for SIA and SIDIS (analytical)

$$d\sigma = \sum_{abc} \int \int \int f_a f_b \, d\hat{\sigma}_{ab \to cX} D_c \, dx_a dx_b dz_c$$

$$= \frac{1}{2\pi i} \sum_{abc} \int_{\mathcal{C}_n} dn \quad D_c^n \quad \int \int \int \int f_a f_b \, d\hat{\sigma}_{ab \to cX} \, dx_a dx_b dz_c$$

$$= \frac{1}{2\pi i} \sum_{abc} \int_{\mathcal{C}_n} dn \quad D_c^n \quad \int \int \int \int f_a f_b \, d\hat{\sigma}_{ab \to cX} \, dx_a dx_b dz_c$$

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$$= \frac{1}{2\pi i} \sum_{abc} \int_{\mathcal{C}_n} dn \quad \int \int \int f_c f_a f_b \, d\hat{\sigma}_{ab \to cX} \, dx_a dx_b dz_c$$

New technique developed to compute the grids: any observable

 $d\sigma$

deF, Stratmann, Sassot, Vogelsang

allows for >5000 ev./sec.!



Good description of SIDIS multiplicities

Hermes data (not final)







- Large scale uncertainty
- Mainly samples z>0.5



 $\mu_F = \mu_R = p_T$



$$\sigma(pp \to \pi^+) - \sigma(pp \to \pi^-)$$

"Charge-sign asymmetries" from ~ I or 2 orders of magnitude sm and "huge errors"



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Typically $v^2/dof \sim 2$	Experiment	Data type	Relative normalization in fit	Data points fitted	χ^2
Typically $\chi / u 0 J \sim 2$	TPC [15]	Inclusive	0.94	17	18.5
		"uds tag"	0.94	9	1.9
		"c tag"	0.94	9	5.7
		<i>"b</i> tag"	0.94	9	7.4
9	TASSO [38]	Inclusive (34 GeV)	0.94	11	30.1
Large γ^2 from a few isolated data points		Inclusive (44 GeV)	0.94	7	20.5
Large χ morn a few isolated data points	SLD [16]	Inclusive	1.008	28	14.0
(and bre SIDIS and some SIDIS and p)		"uds tag"	1.008	17	11.6
(smail z m siz, and some sizes and pp)		" $c \tan^{\prime\prime}$	1.008	17	11.1
very precise data		<i>b</i> tag	1.008	17	33.2 29.2
very precise data	ALEPH [11] DEI DUI [12]	Inclusive	0.97	22 17	38.3 42.3
	DELFIII $\begin{bmatrix} 12 \end{bmatrix}$	"uds taq"	1.0	17	42.5
		"h tag"	1.0	17	20.4 42.8
	OPAL [13 14]	Inclusive	1.0	21	9.2
		" $u \tan^2$	1.10	5	11.8
		" $d \tan g$ "	1.10	5	9.0
Also tonsion botwoon experiments		"s tag"	1.10	5	49.8
Also tension between experiments		" c tag"	1.10	5	38.3
(like Delphi et large z)		"b tag"	1.10	5	73.0
(like Delphi at large Z)	HERMES [17]	π^+	1.03	32	67.4
		π^-	1.03	32	120.8
	PHENIX [18]	π^0	1.09	23	76.4
	STAR [22]	$\pi^0, \langle \eta \rangle = 3.3$	1.05	4	3.4
		$\pi^0, \langle \eta angle = 3.7$	1.05	5	9.8
	BRAHMS [21]	$\pi^+, \langle \eta angle = 2.95$	1.0	18	28.2
		$\pi^{-},\langle\eta angle=2.95$	1.0	18	43.0
	Total			392	843.7

 χ^2 grows (~25%) for LO fits : mostly from pp data where NLO corrections are very large

Pions + kaons + protons almost saturate charged hadrons: Mainly a prediction

Meet the Distributions : pions



Large differences visible in the gluon at large z : explains pp

Large differences in unfavored distributions : explains SIDIS and pp Similar singlet

For pions u fragmentation smaller than AKK : required by SIDIS

Meet the Distributions : Kaons







Differences with HKNS also sizeable : gluons, unfavored and large z (pp)

No SIDIS : favored/unfavored separation not that reliable (assumptions)

Meet the Distributions : protons



Differences with HKNS also sizeable : gluons, unfavored and large z (pp)

No SIDIS : favored/unfavored separation not that reliable (assumptions)

In general differences look much larger than for pdf fits : comparison between GLOBAL fits

Charged hadrons: Pions + kaons + protons almost saturate charged hadrons: residual only sizable for HQ Mainly a prediction/check : OK



Uncertainties on fragmentation functions

✓ HKNS uncertainties using Hessian Method

M.Hirai et al

only e⁺e⁻ data so uncertainties are not well defined for gluons and flavor/charge separation

Hessian method might fail when far from quadratic dependence

✓ DSS uncertainties using Lagrangian multipliers D.deF, M. Stratmann, R. Sassot

study uncertainties for truncated moments but does not translate easily into uncertainties for z-dependence fragmentation functions

More work (and data!) needed for a precise estimate of uncertainties

Improvements on description of fragmentation functions over the last few years

	Future:
EXP "next summer"	<pre>SIDIS: COMPASS SIDIS: Hermes Final? Belle/Babar: precision at lower energies more pp data (RHIC, LHC?)</pre>
TH	SIDIS at NNLO Splitting at NNLO Threshold resummation for several processes More global fits to compare!

• More precise determination of fragmentation functions

• Realistic estimate of uncertainties and several sets to compute uncertainties for other observables (like pdf's)

Improvements on description of fragmentation functions over the last few years



- More precise determination of fragmentation functions
- Realistic estimate of uncertainties and several sets to compute uncertainties for other observables (like pdf's)

TH effort usually triggered by precise data !



M. Arneodo et al. / Quark distribution function

Older data not very well documented (no multiplicities, only FFs)

MEASUREMENTS OF THE u VALENCE QUARK DISTRIBUTION FUNCTION IN THE PROTON AND u QUARK FRAGMENTATION FUNCTIONS

The European Muon Collaboration

Nuclear Physics B321 (1989) 541-560 North-Holland, Amsterdam



Uncertainties on fragmentation functions

HKNS uncertainties using Hessian Method

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$$\chi^{2}(a) = \chi^{2} - \chi_{0}^{2} = \sum_{i,j} H_{ij} \delta a_{i} \delta a_{j} + \cdots$$

$$\chi^{2} = 15.94$$

$$[\delta D_{i}^{h}(z)]^{2} = \Delta \chi^{2} \sum_{j,k} \left(\frac{\partial D_{i}^{h}(z,a)}{\partial a_{j}} \right)_{\hat{a}} H_{jk}^{-1} \left(\frac{\partial D_{i}^{h}(z,a)}{\partial a_{k}} \right)_{\hat{a}}$$

Main issue : only SIA data

Gluon: uncertainties overestimated due to lack of other data?

Charge separation: uncertainties could be underestimated due to assumptions

For some parameters, χ^2 profile VERY different from quadratic even in global analysis

DSS uncertainties using Lagrangian multipliers

$$\Phi(\lambda_i, \{a_j\}) = \chi^2(\{a_j\}) + \sum_i \lambda_i \mathcal{O}_i(\{a_j\}) \qquad \qquad \int_{0.2}^1 z \, D_i^H(z, Q = 5 \,\text{GeV}) \, dz$$

See how fit deteriorates when FFs forced to give different prediction for \mathcal{O}_i

$$\Delta \chi^2 = 15 (\sim 2\%)$$



-5

-5

Precision



Juon density depends much



Extract

on PHENIX .



DSSV (deF, Stratmann, Sassot, Vogelsang)



Example @ Phenix

Lot of "initial gluon" (gg and qg) "flavor blind"



 low transverse momentum probe gluon fragmentation

central rapidity PHENIX



17. FRAGMENTATION FUNCTIONS IN e^+e^- , ep AND pp COLLISIONS

Revised October 2009 by O. Biebel (Ludwig-Maximilians-Universität, Munich, Germany), D. de Florian (Dep. de Física, FCEyN-UBA, Buenos Aires, Argentina), D. Milstead (Fysikum, Stockholms Universitet, Sweden), and A. Vogt (Dep. of Mathematical Sciences, University of Liverpool, UK).



Figure 17.6: Comparison of up, strange, charm and gluon NLO fragmentation functions for $\pi^+ + \pi^-$ at the mass of the Z. The different lines correspond to the result of the most recent analyses performed in Refs. [93,94,98].

More K multiplicities AKK(08) and Kretzer





Spin-off

Can we say something about the strange pdf? May be, first make sure fragmentation is right!



Better: check directly multiplicities with different strange distributions (even possible to fit them at NLO)

 Present TH challenge: Serious disagreement between TH and EXP for fixed target kinematics in pp collisions



Ad-hoc intrinsic transverse momentum sometimes introduced. But pQCD can solve the problem : resummation

D.deF, W.Vogelsang

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