

# Magnetic Resonance Imaging

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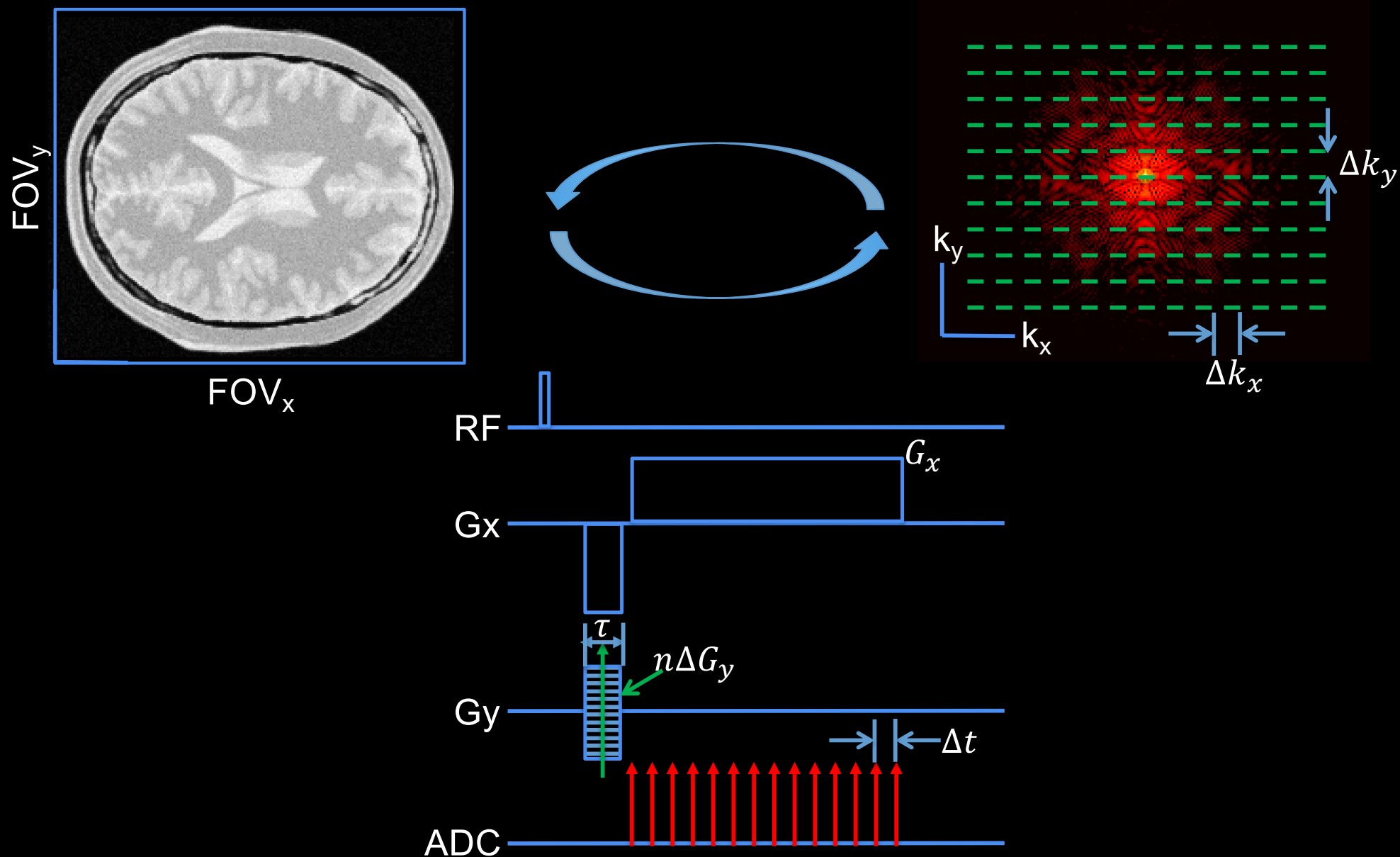
Gordon Center for Medical Imaging

Massachusetts General Hospital

Harvard Medical School



# Outline



# Outline

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- Introduction
- MR physics
  - Spin
  - Bloch equation
  - Signal excitation and reception
  - Relaxation
- MR imaging
  - Projection-based MR imaging
  - Fourier transform-based MR imaging
  - k-space sampling
- Summary

# Magnetic Resonance Imaging

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## Introduction



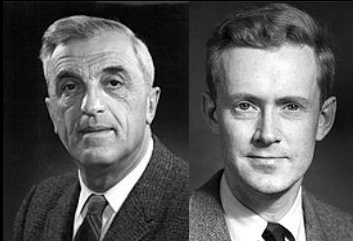


# History



**1944: Rabi**

Physics (Measured magnetic moment of nucleus)



**1952: Felix Bloch and Edward Mills Purcell**

Physics (Basic science of NMR phenomena)



**1991: Richard Ernst**

Chemistry (High-resolution Pulsed FT-NMR)



**2002: Kurt Wüthrich**

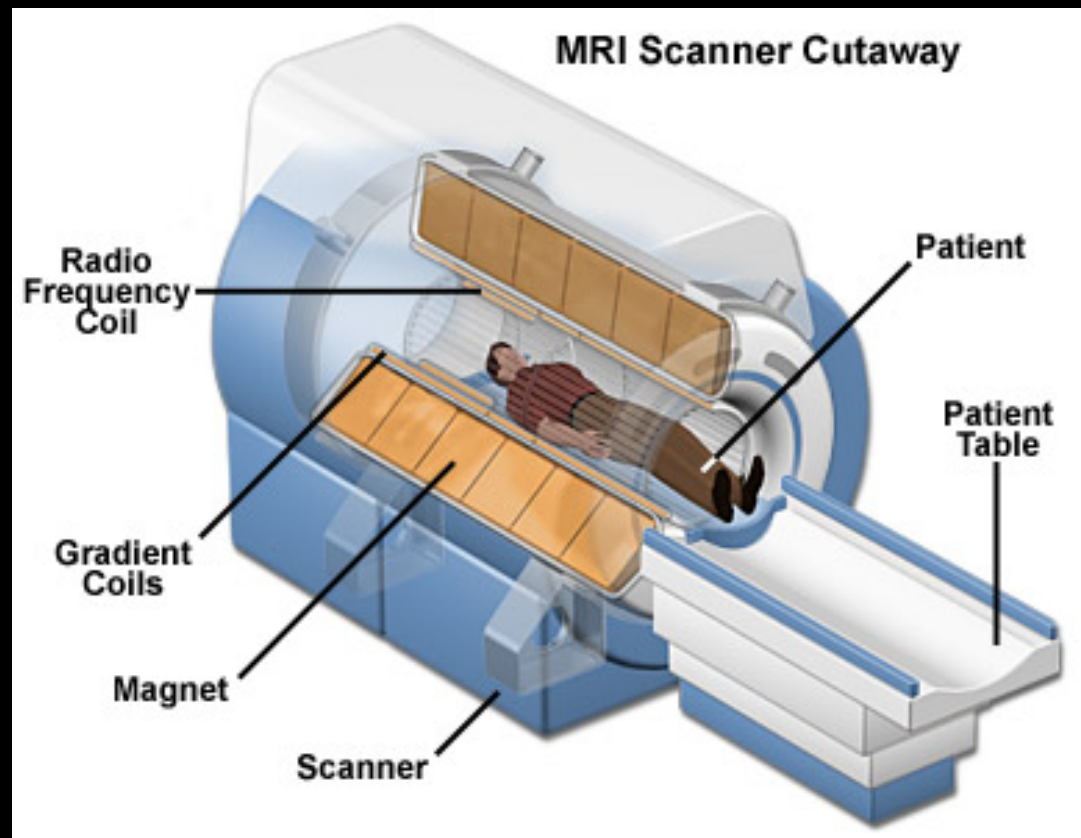
Chemistry (3D molecular structure in solution by NMR)



**2003: Paul Lauterbur & Peter Mansfield**

Physiology or Medicine (MRI technology)

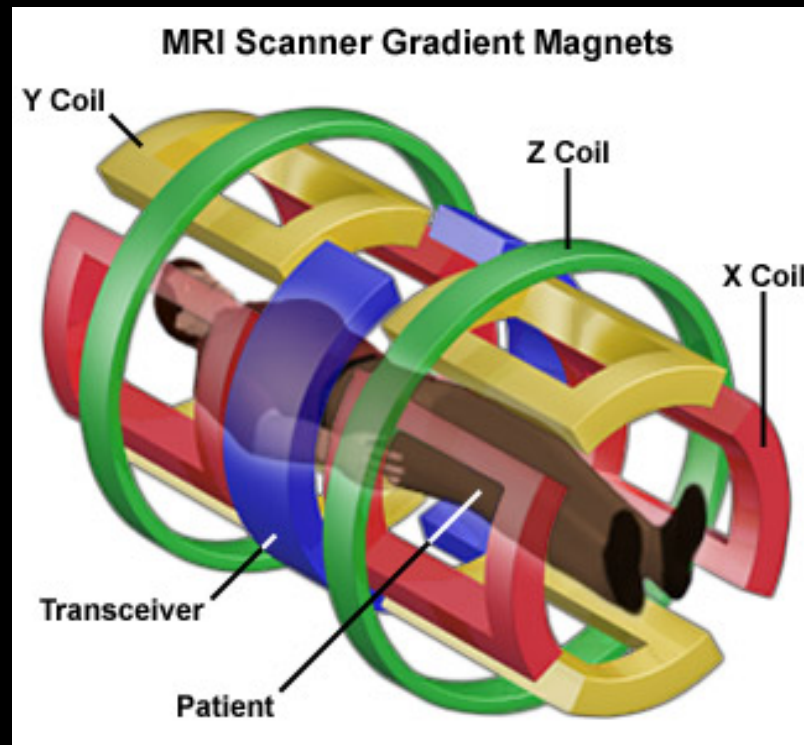
# MRI scanner



## Superconducting magnet

- Strong, static, and homogeneous magnetic field (1.5 ~ 10 T)
- A 3 T magnetic field is 60,000 times stronger than the Earth's magnetic field.
- < 1 ppm in a 25-cm diameter sphere
- **Polarizer**

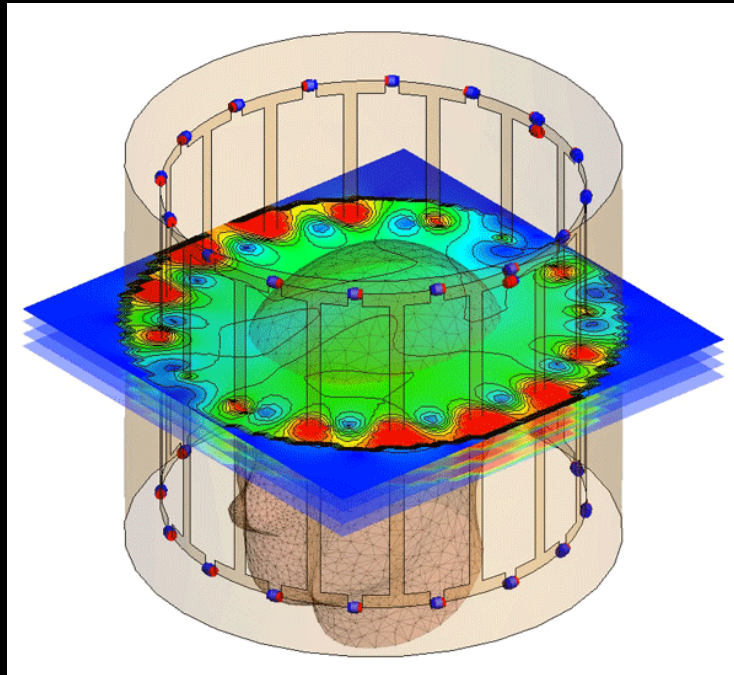
# MRI scanner



## Gradient coils

- An inhomogeneous magnetic field whose z-component varies linearly along a specific direction called the gradient direction.
- Typical specification: 40 mT/m, 200 T/m/s slew rate
- $B_0 \approx 3 \text{ T}$ ,  $G_x \approx 10 \text{ mT/m}$  or  $\max(xG_x) \approx \frac{1}{1000} B_0$
- **Spatial encoding**

# MRI scanner



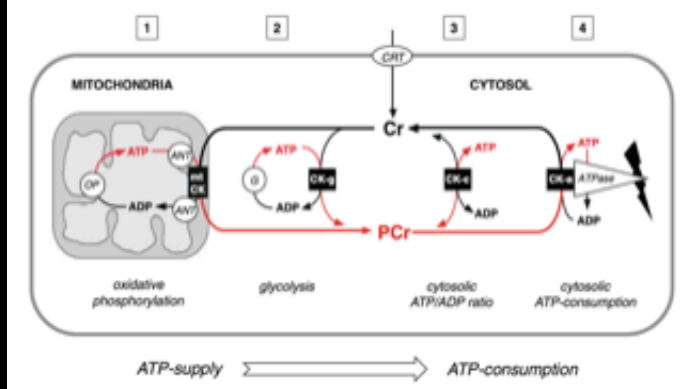
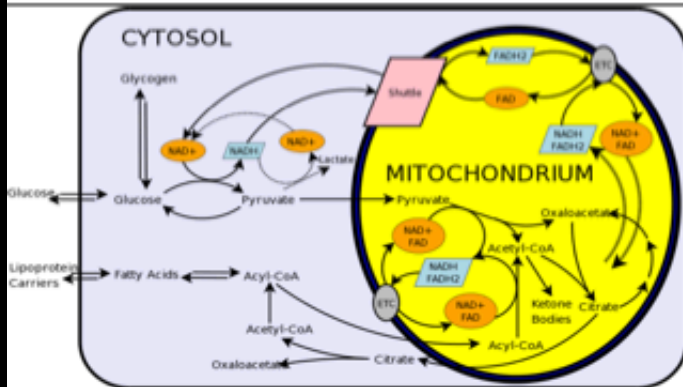
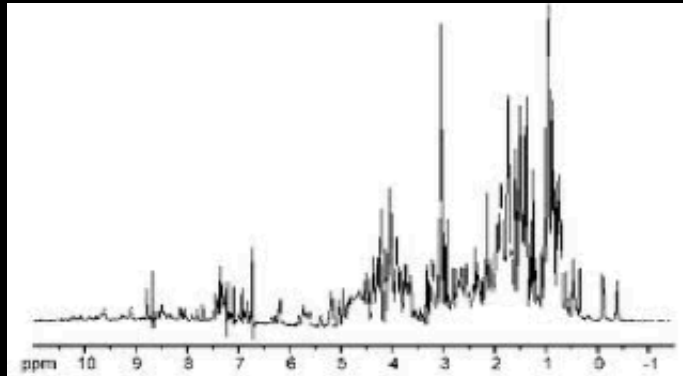
## Radiofrequency (RF) coils

- Time-varying magnetic field at NMR resonance frequency (e.g., 60 ~300 MHz)
- Typical strength 0.1 Gauss
- **Signal excitation and reception**

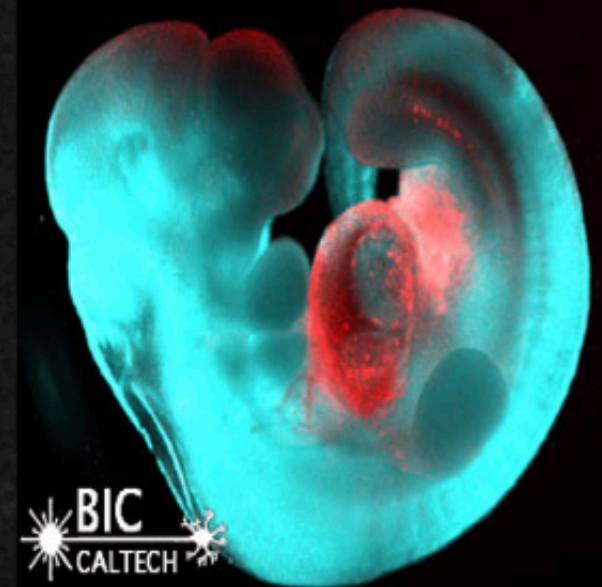
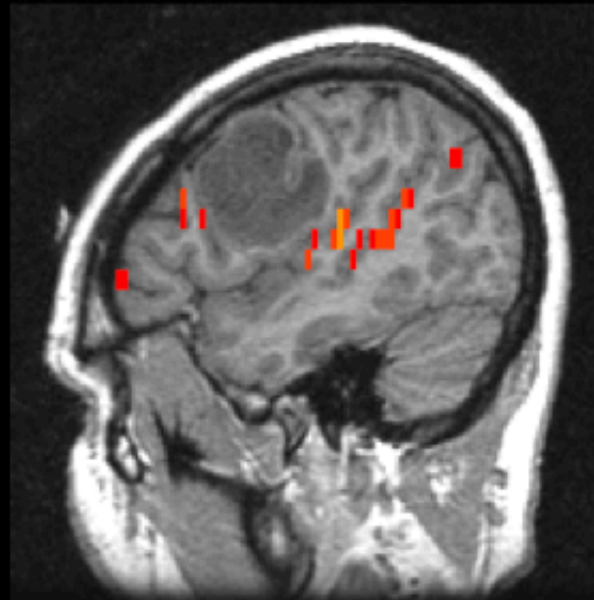
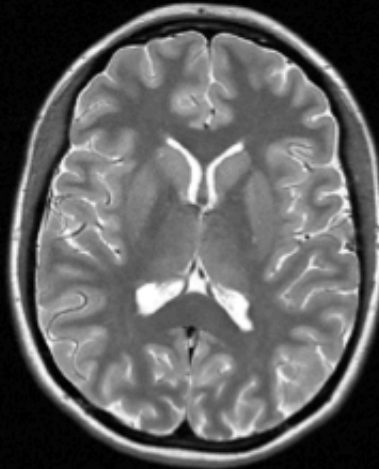
Courtesy of P. Futter, computed with FEKO

# Magic of Spins

## Spectroscopy



## Imaging





# Magnetic Resonance Imaging

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## MR Physics



# Spin

- Spin is an *intrinsic property* of particles just like mass, charge, etc.

- Nuclear magnetic moment



–  $\vec{\mu} = \gamma \vec{J}$ , where  $\vec{J}$  is the spin angular momentum

–  $\vec{\mu} = \gamma \hbar \sqrt{I(I + 1)}$ , where  $I$  is the nuclear spin quantum number.

–  $I = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \frac{5}{2}, \dots$

- Nuclei with an odd mass number have half-integral spin.
- Nuclei with an even mass number and an even charge number have zero spin.
- Nuclei with an even mass number but an odd charge number have integral spin.

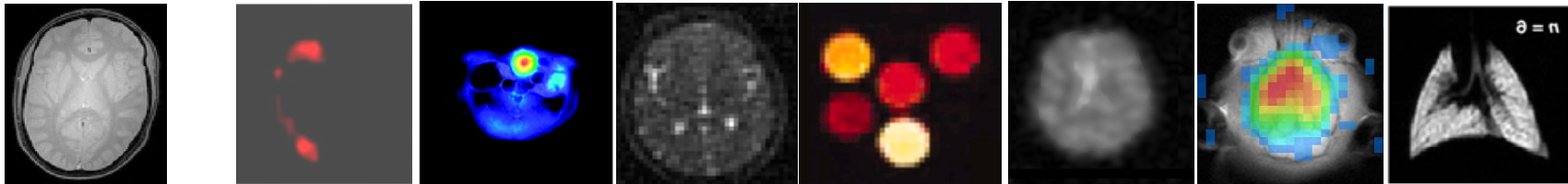
# Nuclei used in MRI

Nucleus	$\gamma / 2\pi$ (MHz/T)
$^1\text{H}$	42.576
$^2\text{H}$	6.53566
$^3\text{He}$	-32.434
$^7\text{Li}$	16.546
$^{13}\text{C}$	10.705
$^{14}\text{N}$	3.0766
$^{15}\text{N}$	-4.3156
$^{17}\text{O}$	-5.7716
$^{23}\text{Na}$	11.262
$^{31}\text{P}$	17.235
$^{129}\text{Xe}$	-11.777

1 H																	2 He				
3 Li	4 Be															5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg															13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr				
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe				
55 Cs	56 Ba	71 Lu	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn				
87 Fr	88 Ra	103 Lr	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Uub	113 Uut	114 Uuq	115 Uup	116 Uuh	117 Uus	118 Uuo				
		*																			
		57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb						
		**																			
		89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No						

40      30      20      10      0      -10      -20      -30      MHz/T

$^1\text{H}$        $^{31}\text{P}$        $^{13}\text{C}$        $^{23}\text{Na}$        $^{15}\text{N}$        $^{17}\text{O}$        $^{129}\text{Xe}$        $^3\text{He}$





# Magnetized nuclear spin system

- In the presence of an external magnetic field

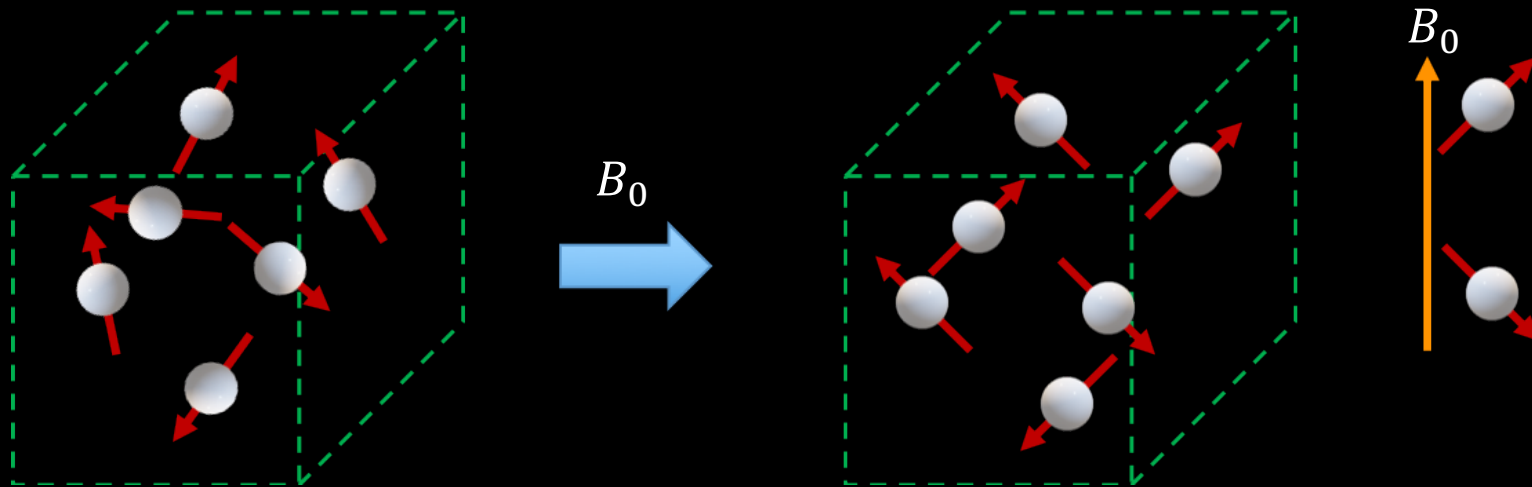
- $\mu_z = \gamma m_I \hbar$

- $m_I = -I, -I + 1, \dots, I$ , is the magnetic quantum number

- For spin-1/2 systems

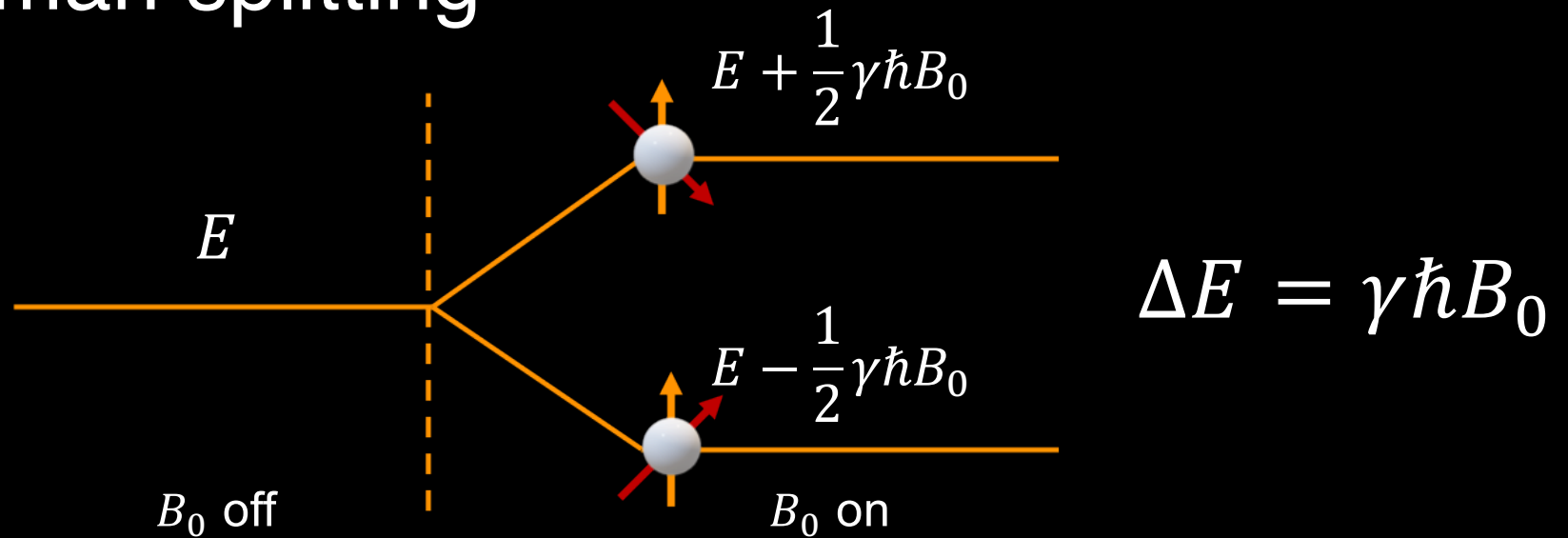
- $\mu_z = \pm \frac{\gamma \hbar}{2}$

- $\mu_{xy} = \frac{\gamma \hbar}{\sqrt{2}} e^{i\phi}$ , where  $\phi$  is randomly distributed.

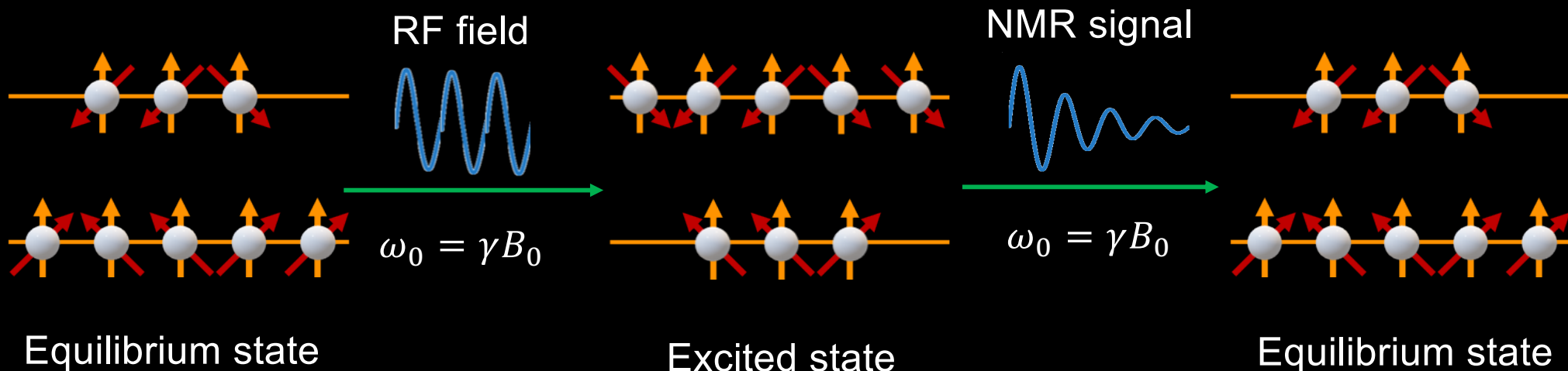


# Magnetized nuclear spin system

- Zeeman splitting



- Nuclear magnetic resonance



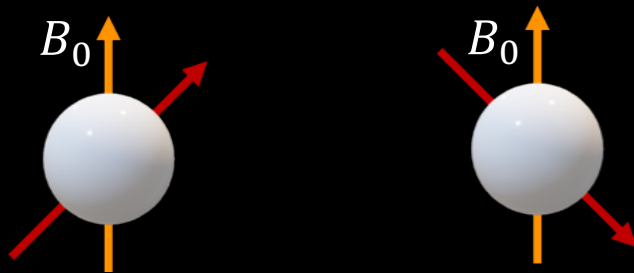
# Bulk magnetization

- To describe the ensemble behavior of a spin system

$$\vec{M} = \sum_{n=1}^{N_s} \vec{\mu}_n$$

– Spin-1/2 system

$$\vec{M} = \left( \sum_{n=1}^{N_{\uparrow}} \frac{1}{2} \gamma \hbar \hat{z} - \sum_{n=1}^{N_{\downarrow}} \frac{1}{2} \gamma \hbar \hat{z} \right) = (N_{\uparrow} - N_{\downarrow}) \frac{1}{2} \gamma \hbar \hat{z}$$



# Bulk magnetization

- Bulk magnetization (spin-1/2)

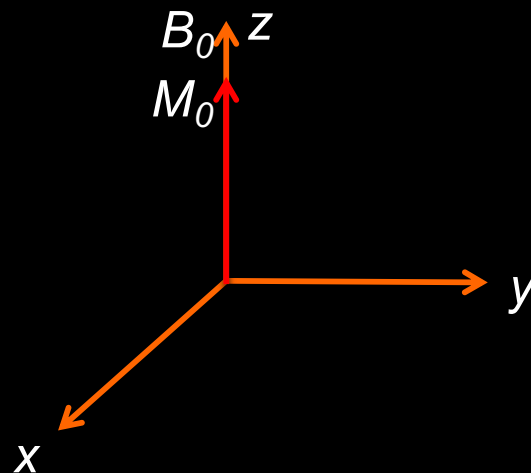
- $\frac{N_{\uparrow} - N_{\downarrow}}{N_S} \approx \frac{\gamma \hbar B_0}{2KT_S}$ ,  $N_S$ : total number of spins

- At 1 Tesla,  $\frac{N_{\uparrow} - N_{\downarrow}}{N_S} \approx 3 \times 10^{-6}$  (three in a million!)

- $\vec{M}$  is parallel to the external magnetic field at equilibrium

- $M_0$  is proportional to  $B_0$

$$\begin{aligned}\vec{M} &= \frac{\gamma^2 \hbar^2 B_0 N_S}{4KT_S} \hat{z} \\ &= M_0 \hat{z}\end{aligned}$$

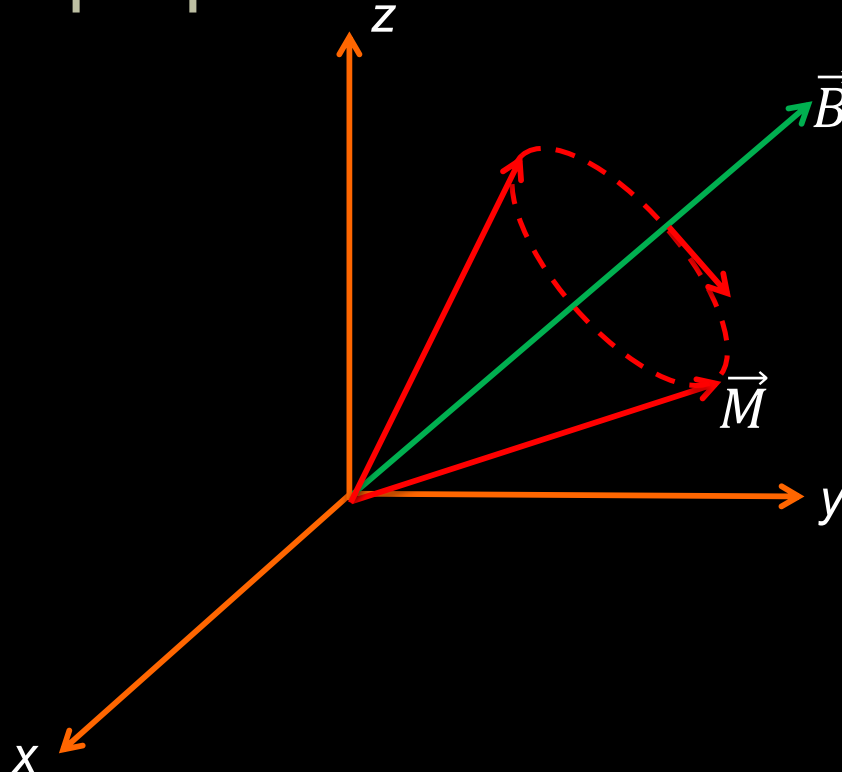


# Bloch equation

- The governing equation of the dynamics of bulk magnetization
  - Ignoring relaxation,  $\vec{M}$  rotates clockwise about the external magnetic field  $\vec{B}$
  - Larmor frequency  $\omega$  is proportional to field strength

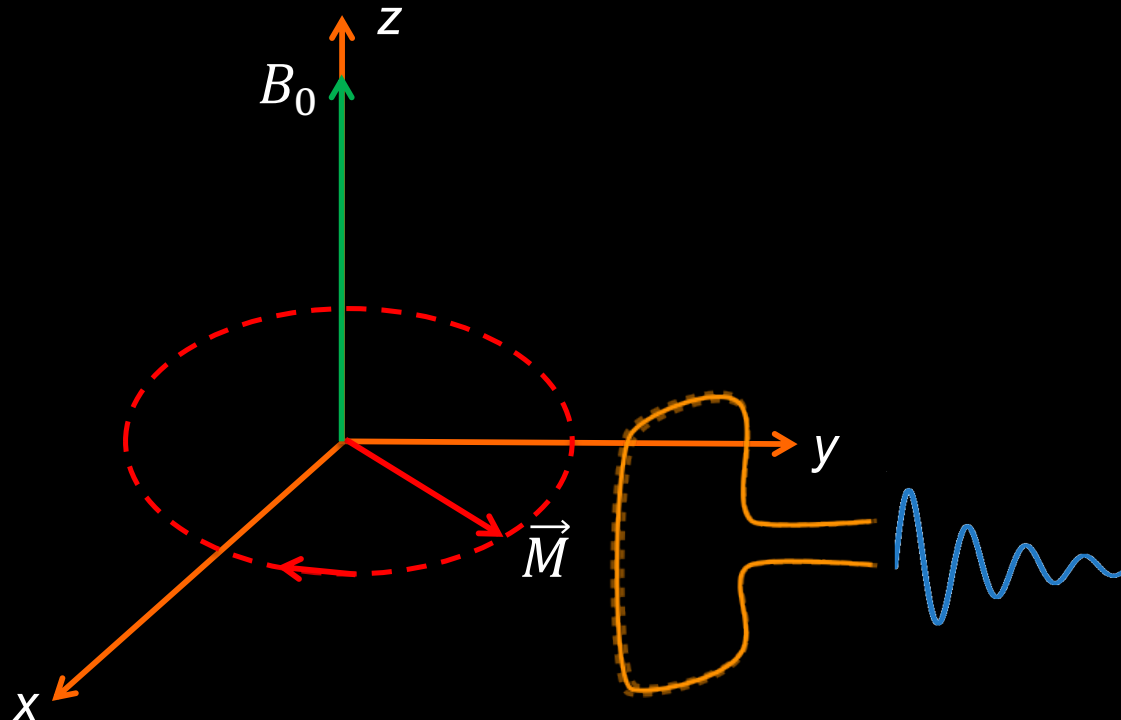
$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}$$

$$\omega = \gamma |\vec{B}|$$



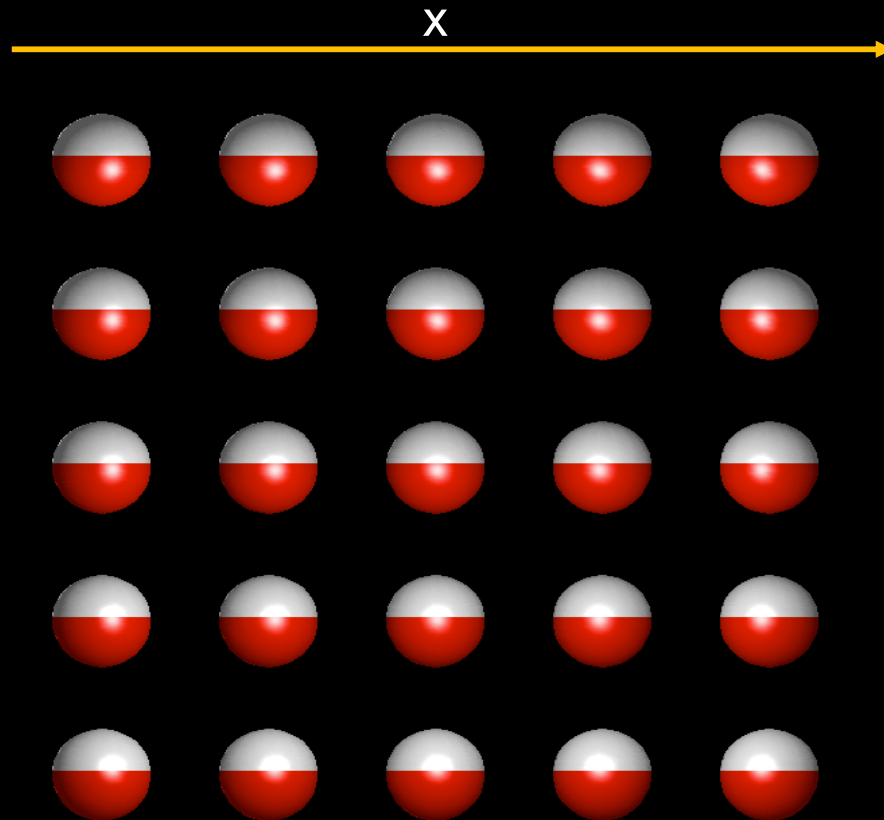
# Example 1: Effects of main field $\vec{B}_0$

- Precession frequency is  $\omega_0 = \gamma B_0$ 
  - $\gamma$  is called the gyromagnetic ratio
  - $\gamma/2\pi = 42.58 \text{ MHz/T}$  for  $^1\text{H}$
  - At  $B_0 = 1.5 \text{ T}$ ,  $f_0 = \omega_0/2\pi \approx 64 \text{ MHz}$  (TV channel 3 in US or channel 2 in Australia)



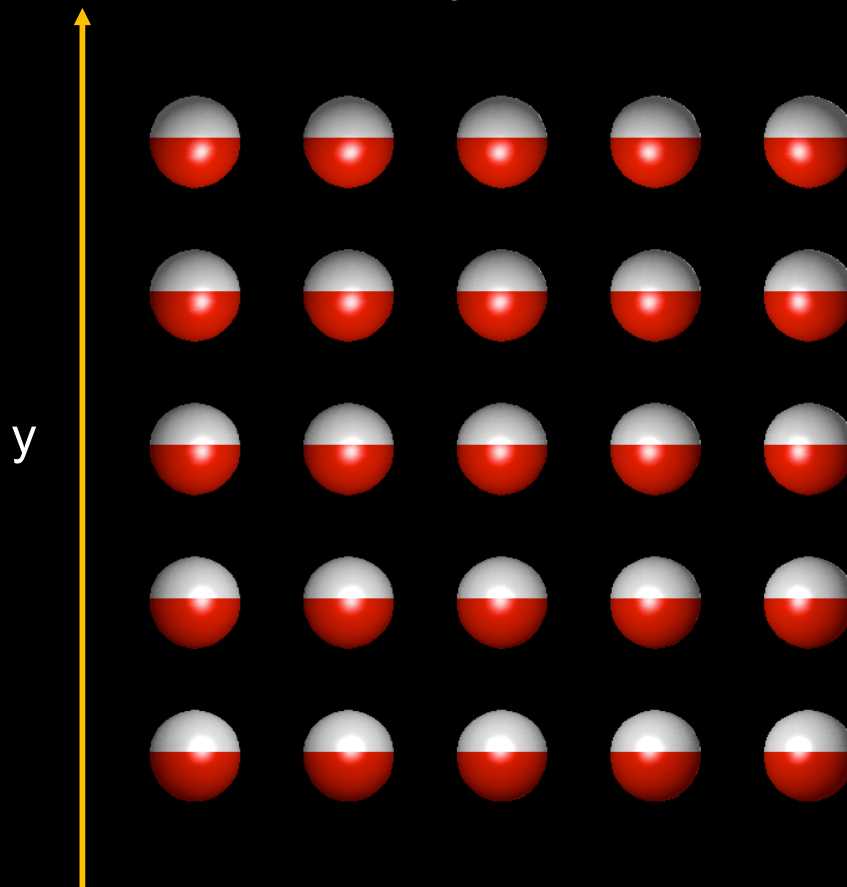
# Example 2: Effects of Gradient field $\vec{G}$

- Gradient field: An *inhomogeneous* magnetic field whose *z-component* varies linearly along a specific direction called the gradient direction.
  - x-gradient:  $\vec{B} = (B_0 + xG_x)\hat{z}$  or  $\omega = (\omega_0 + \gamma xG_x)\hat{z}$



# Example 2: Effects of Gradient field $\vec{G}$

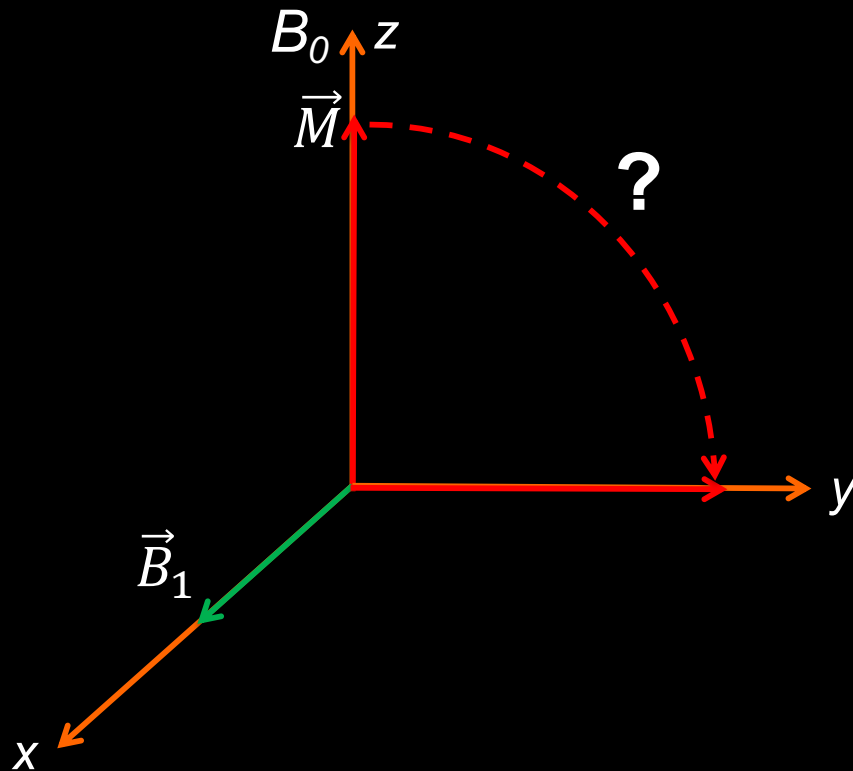
- Gradient field: An *inhomogeneous* magnetic field whose *z-component* varies linearly along a specific direction called the gradient direction.
  - y-gradient:  $\vec{B} = (B_0 + yG_y)\hat{z}$  or  $\omega = (\omega_0 + \gamma yG_y)\hat{z}$





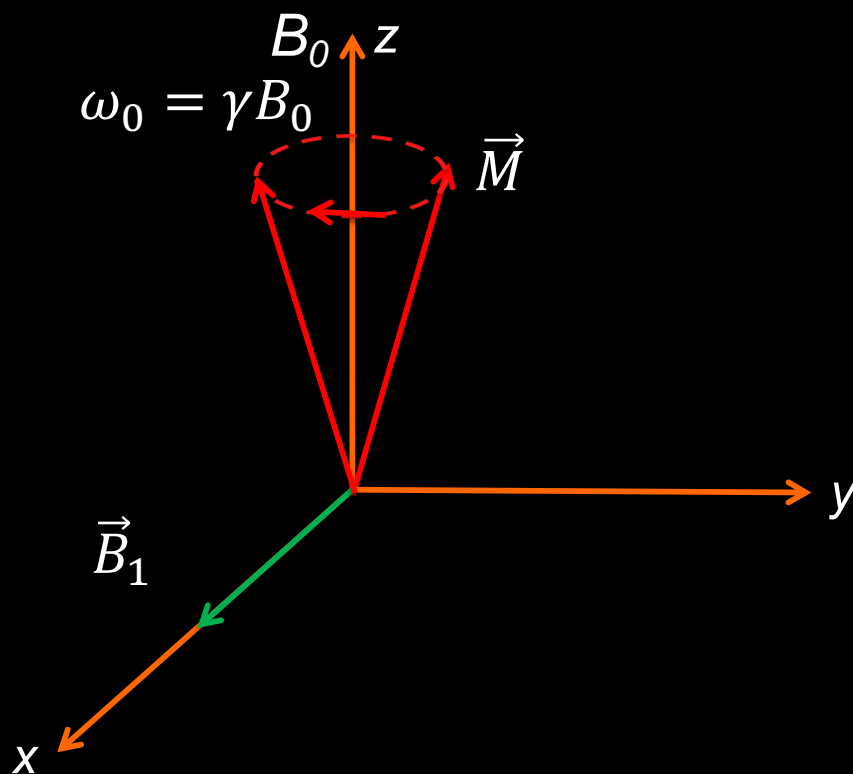
# Non-selective signal excitation

- Excitation: Rotate  $\vec{M}$  from the equilibrium state to the transversal plane.



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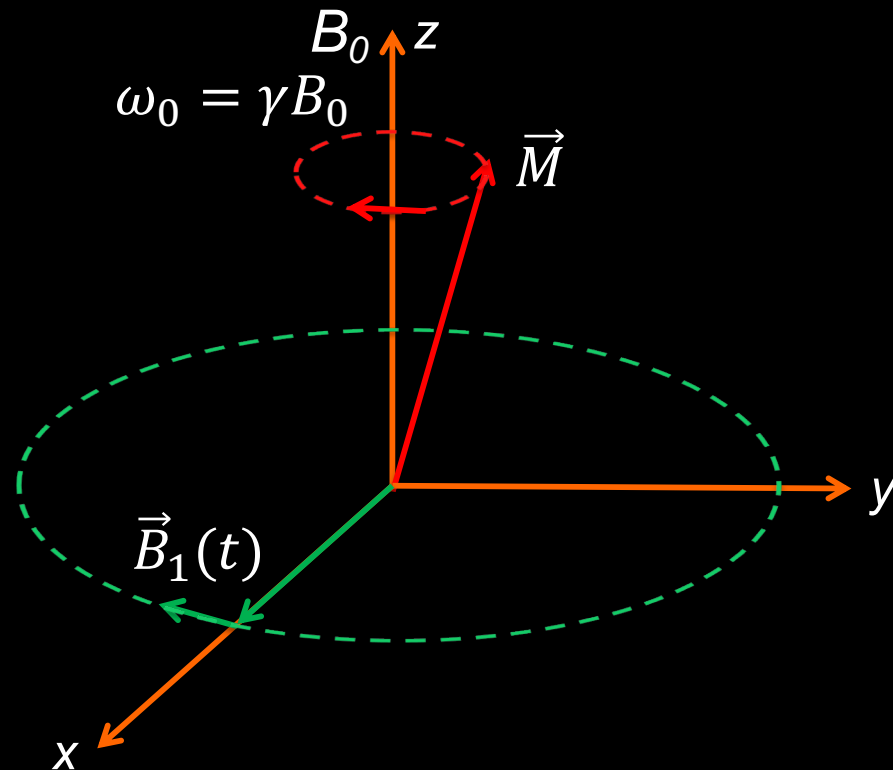
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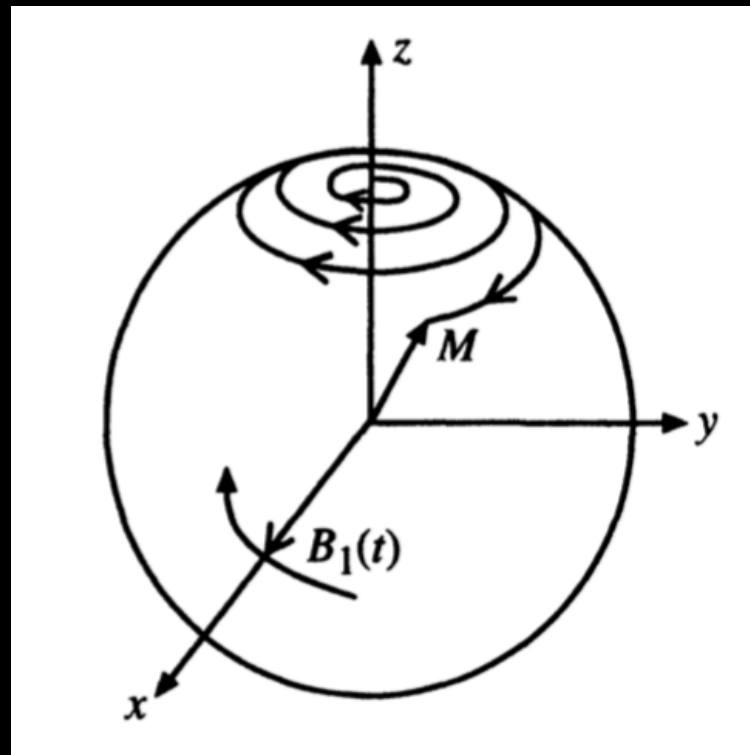
# Non-selective signal excitation

- Excitation: Rotate  $\vec{M}$  from the equilibrium state to the transversal plane.
  - Circular polarized radiofrequency (RF) field
$$\vec{B}_1(t) = B_1(t)(\hat{x}\cos\omega_0t - \hat{y}\sin\omega_0t)$$
  - $B_1 \approx 0.1$  Gauss,  $B_1(t)$  lasts several milliseconds



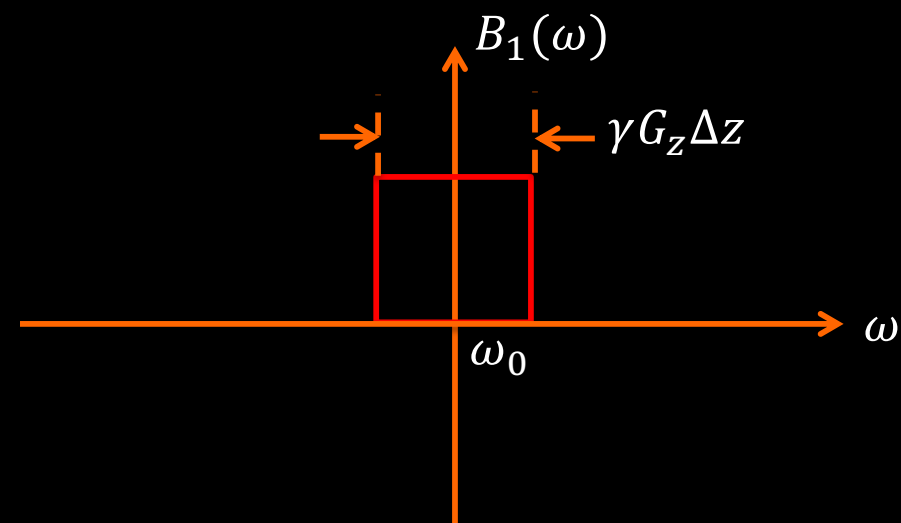
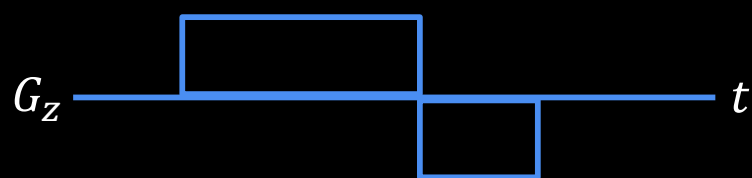
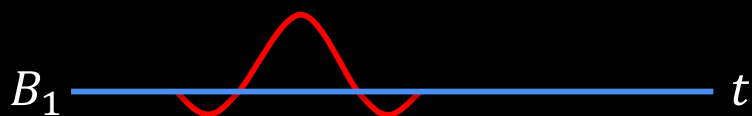
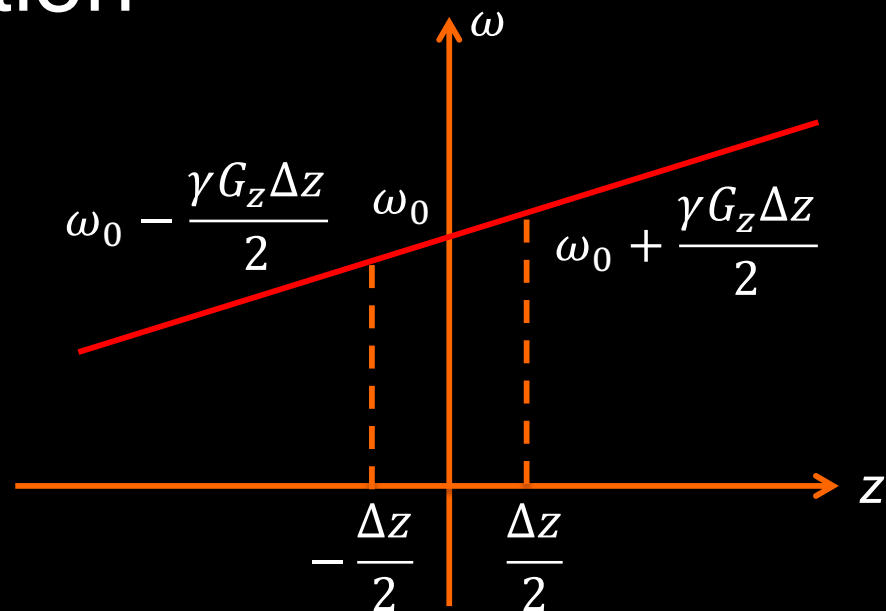
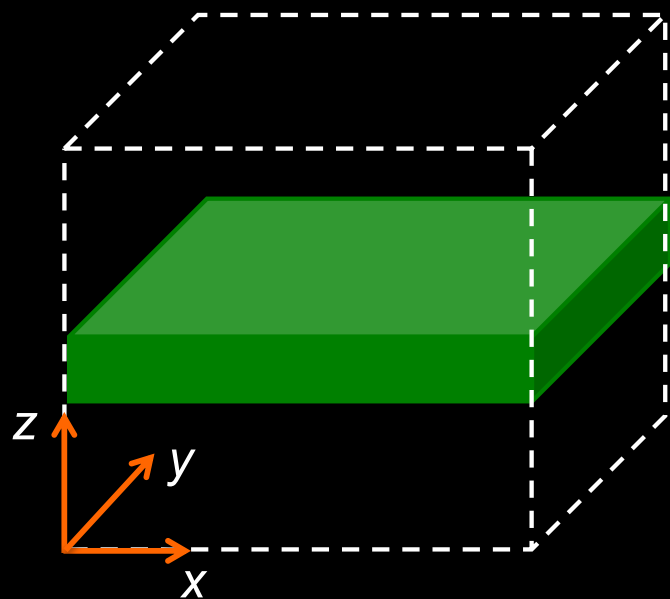
# Non-selective signal excitation

- Excitation: Rotate  $\vec{M}$  from the equilibrium state to the transversal plane.
  - Circular polarized radiofrequency (RF) field
$$\vec{B}_1(t) = B_1(t)(\hat{x}\cos\omega_0t - \hat{y}\sin\omega_0t)$$
  - $B_1(t)$  typically lasts several milliseconds



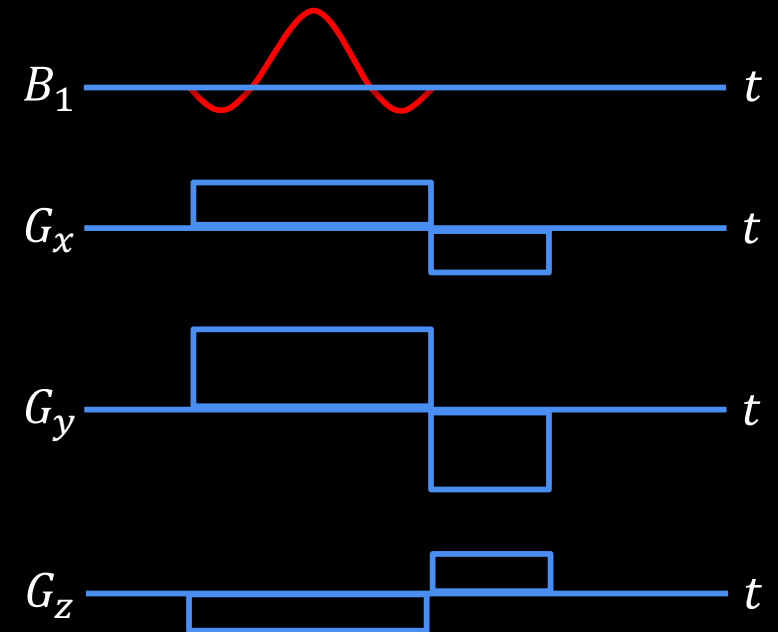
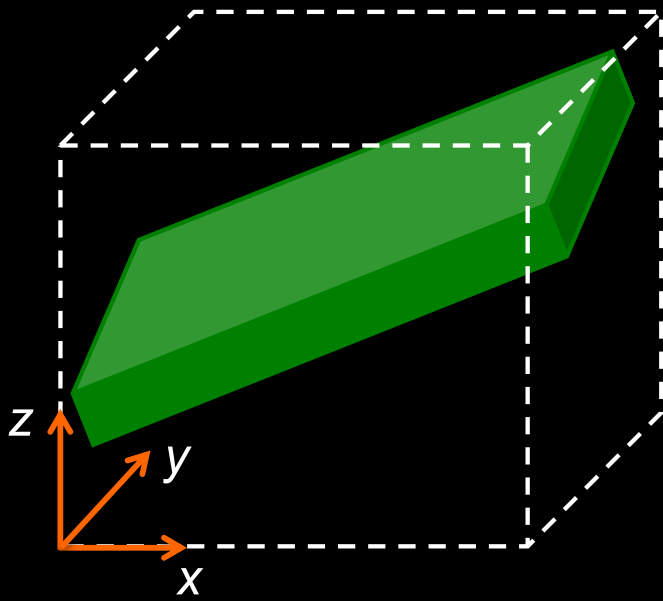
# Selective signal excitation

- Slice-selective excitation



# Selective signal excitation

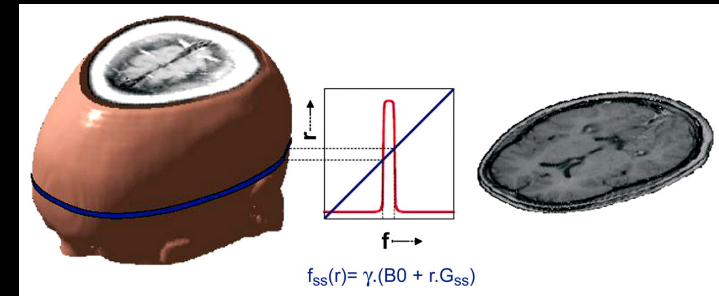
- Slice-selective excitation
  - 2D/3D imaging in an oblique plane/volume



# Selective signal excitation

## One-dimensional RF pulse

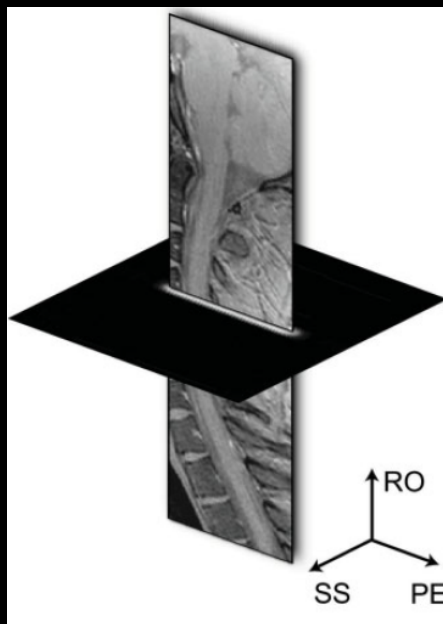
- Slice selection (Garroway 1974; Lauterbur, 1975; Mansfield, 1976; Hoult, 1977)



## Multidimensional RF pulse: more flexible excitation pattern

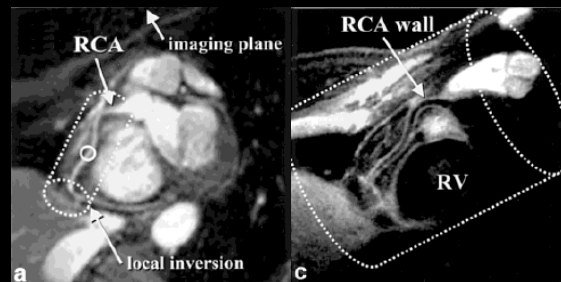
### Reduced FOV imaging

Saritas et al., 2008.



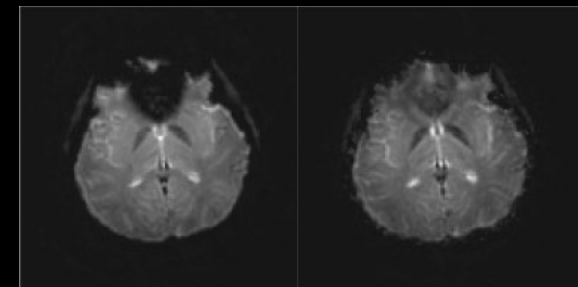
### Local excitation

Botnar et al., 2001.



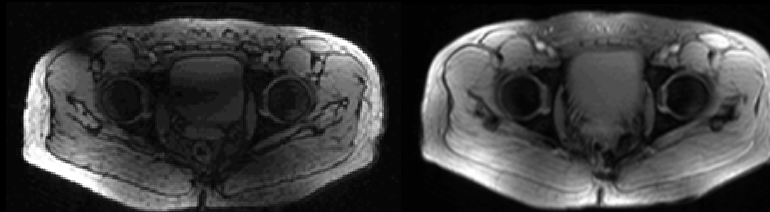
### Signal recovery

Yip et al., 2006.

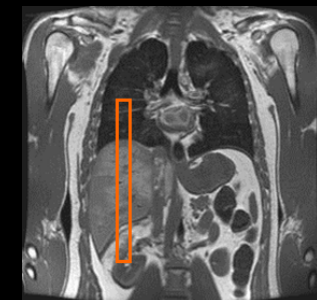


### B<sub>1</sub> inhomogeneity correction

Kerr et al., 2008.



### Motion tracking





# Signal reception

- According to the principle of reciprocity:

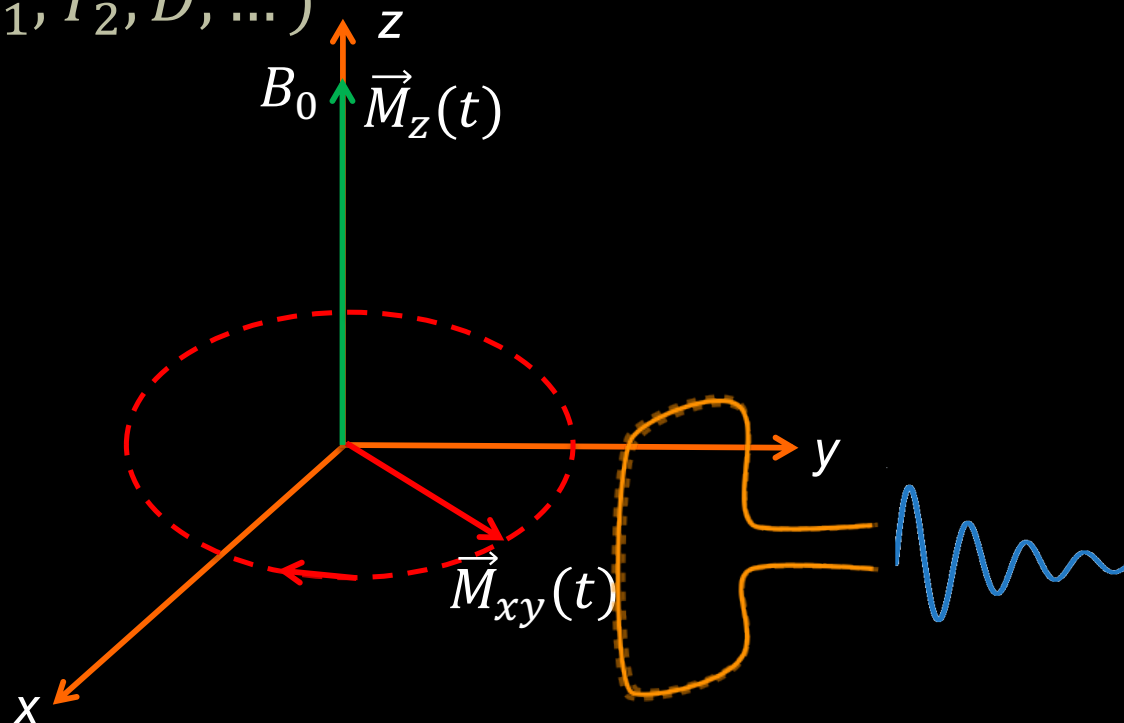
$$V(t) = \omega_0 \int_{\text{Object}} |\vec{M}_{xy}(\vec{r})| \cdot \vec{B}_1^-(\vec{r}) \cos(-\omega_0 t + \phi(\vec{r}, t)) d\vec{r}$$

–  $\vec{B}_1^-(\vec{r})$ :  $B_1$  field generated by a hypothetical unit current flowing in a receive coil.

–  $|\vec{M}_{xy}(\vec{r})| = M_0 \sin \theta f(T_1, T_2, D, \dots)$

–  $\phi(\vec{r}, t)$

- Gradient
- $B_0$  inhomogeneity
- Susceptibility
- Chemical shift
- Velocity
- ...

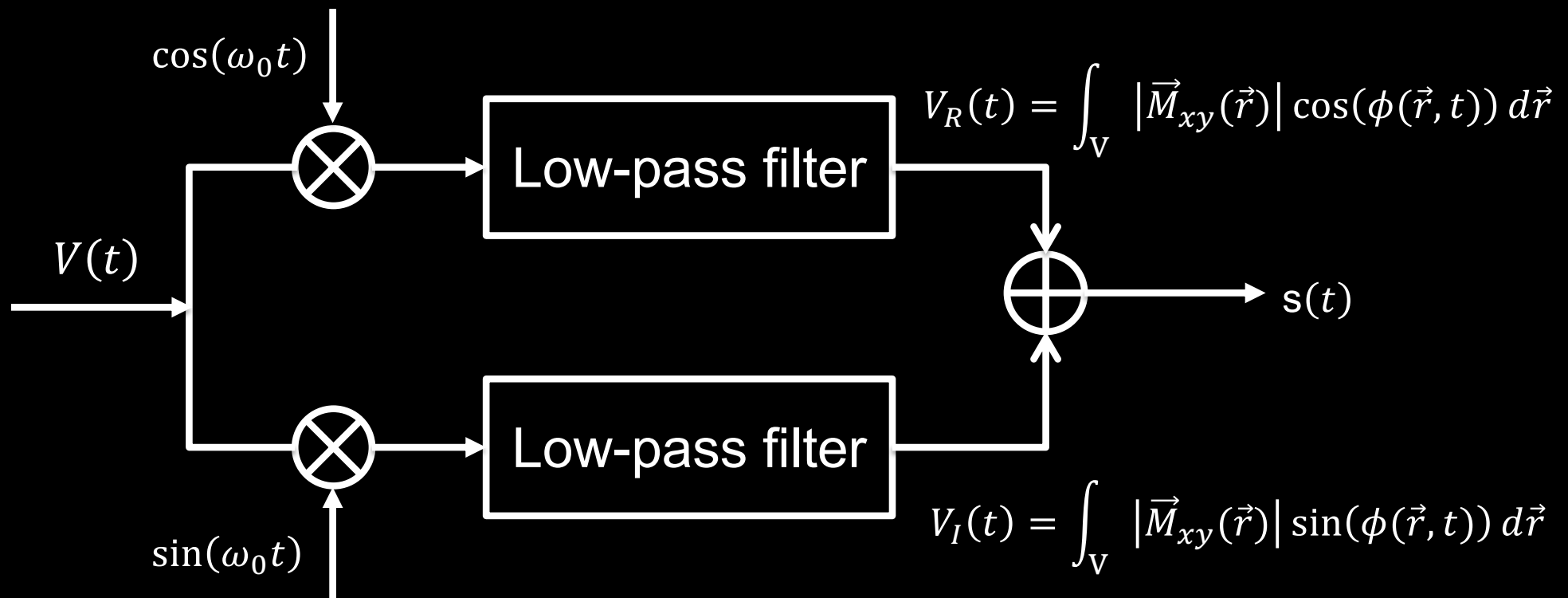


# Signal reception

- Phase-sensitive detection

$$V(t) = \int_{\mathbf{V}} |\vec{M}_{xy}(\vec{r})| \cos(-\omega_0 t + \phi(\vec{r}, t)) d\vec{r}$$

$$s(t) = V_R(t) + iV_I(t) = \int_{\mathbf{V}} |\vec{M}_{xy}(\vec{r})| e^{i\phi(\vec{r}, t)t} d\vec{r}$$



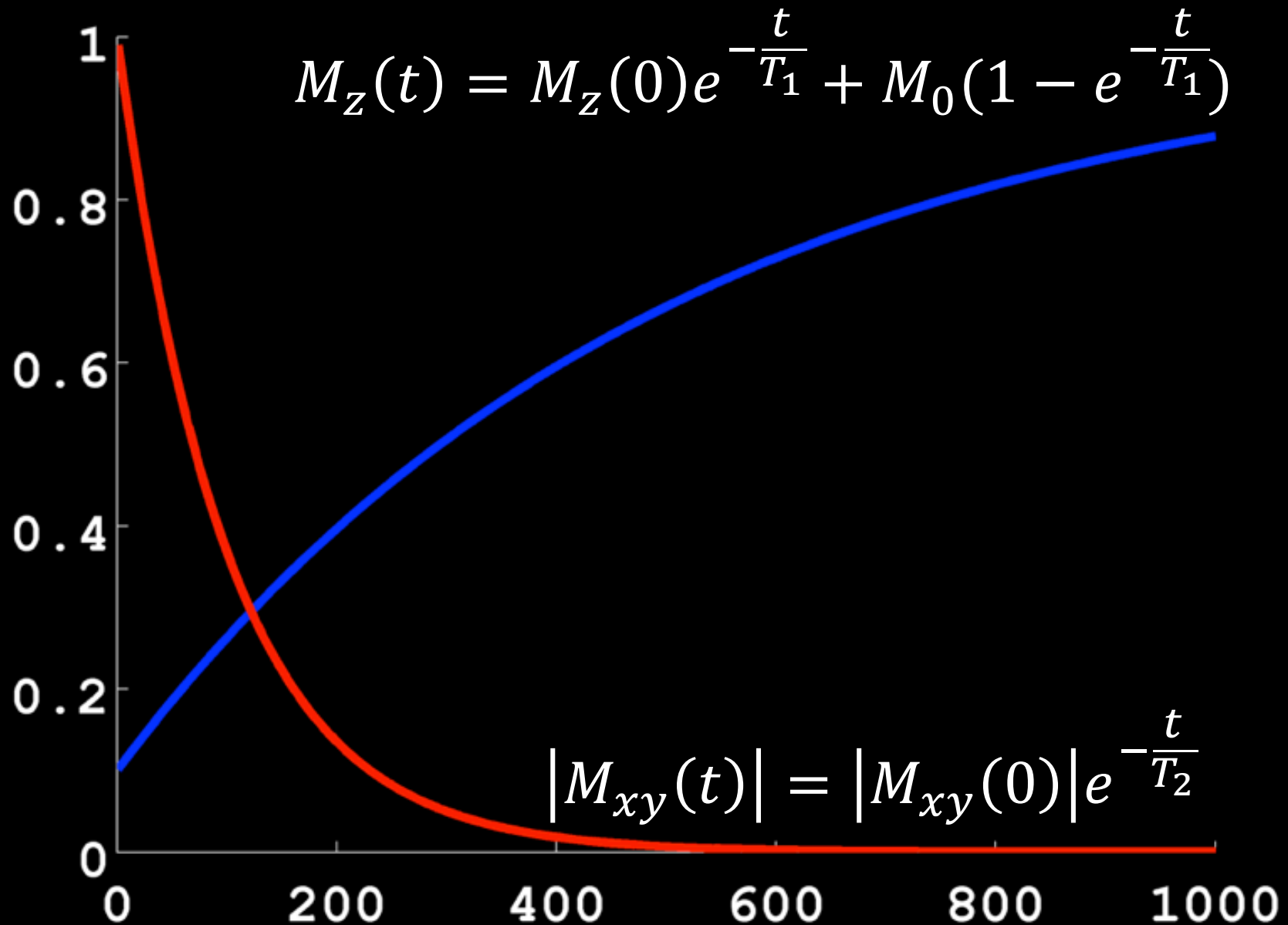
# Relaxation

- Bloch equation

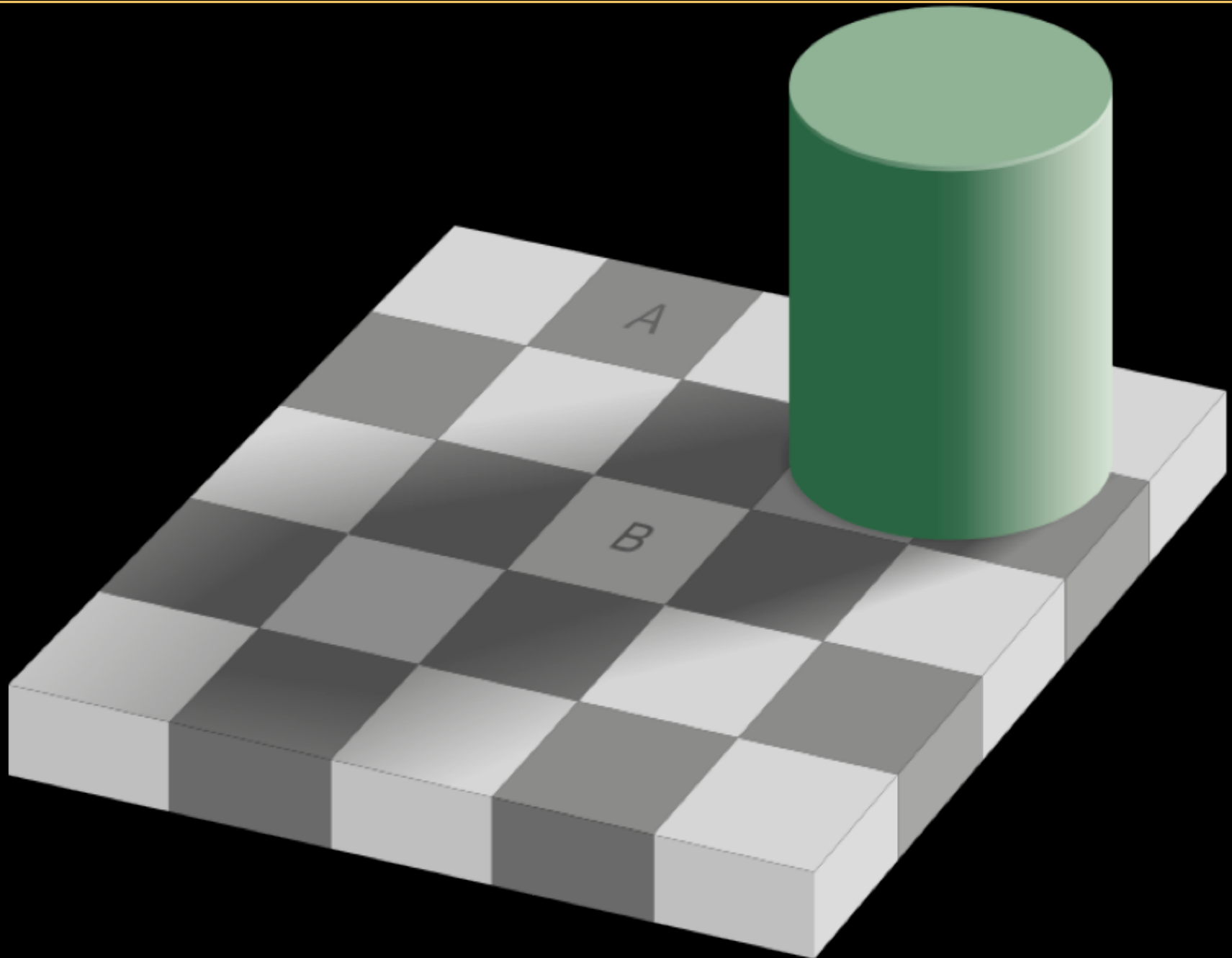
$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B} - \frac{M_x \hat{x} + M_y \hat{y}}{T_2} - \frac{(M_z - M_0) \hat{z}}{T_1}$$

- $T_1$  : Longitudinal relaxation time
- $T_2$  : Transverse relaxation time
- $M_0 \hat{z}$  : Equilibrium state of the bulk magnetization

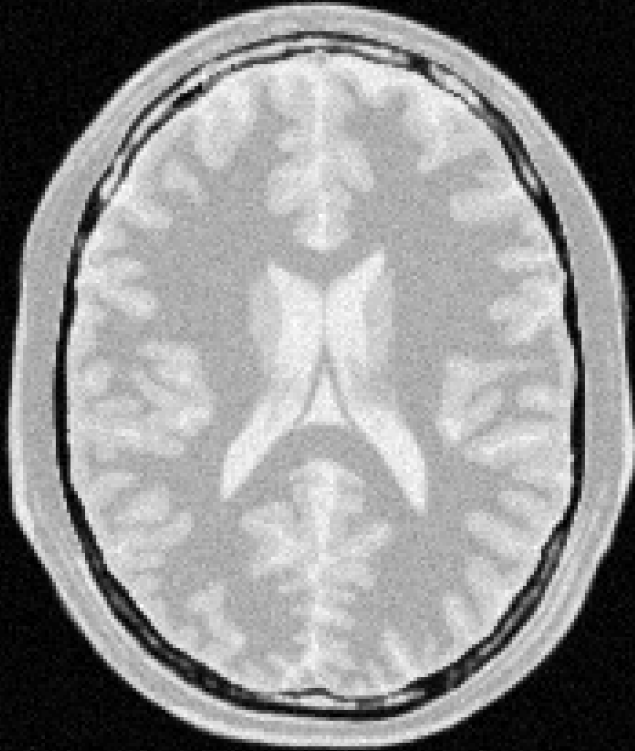
# Relaxation



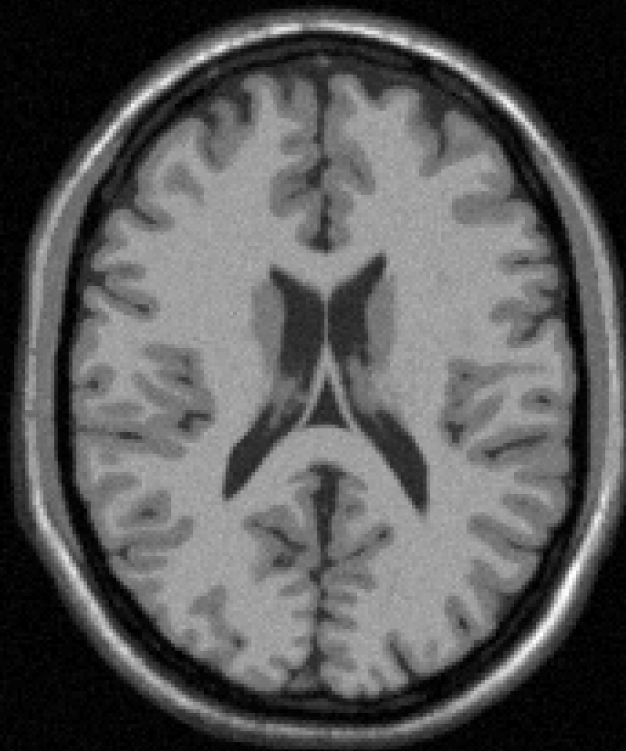
# Contrast



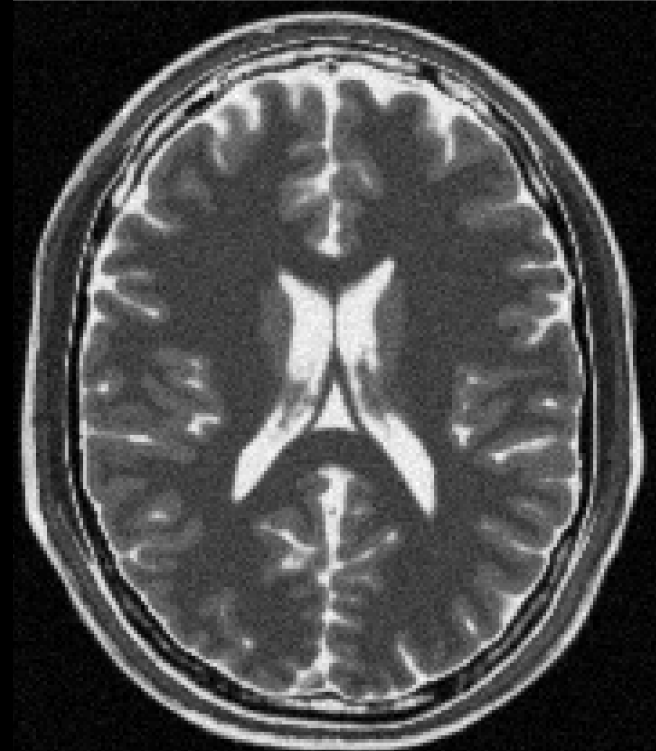
# Contrast



Proton density



T1-weighted



T2-weighted

# Review: MRI physics

- Bulk magnetization

- $M_0 = \sum_n^{N_s} \mu_{z,n}$

- $M_0 = \frac{\gamma^2 \hbar^2 B_0 N_s}{16\pi^2 K T_s}$  (1/2-spin system)

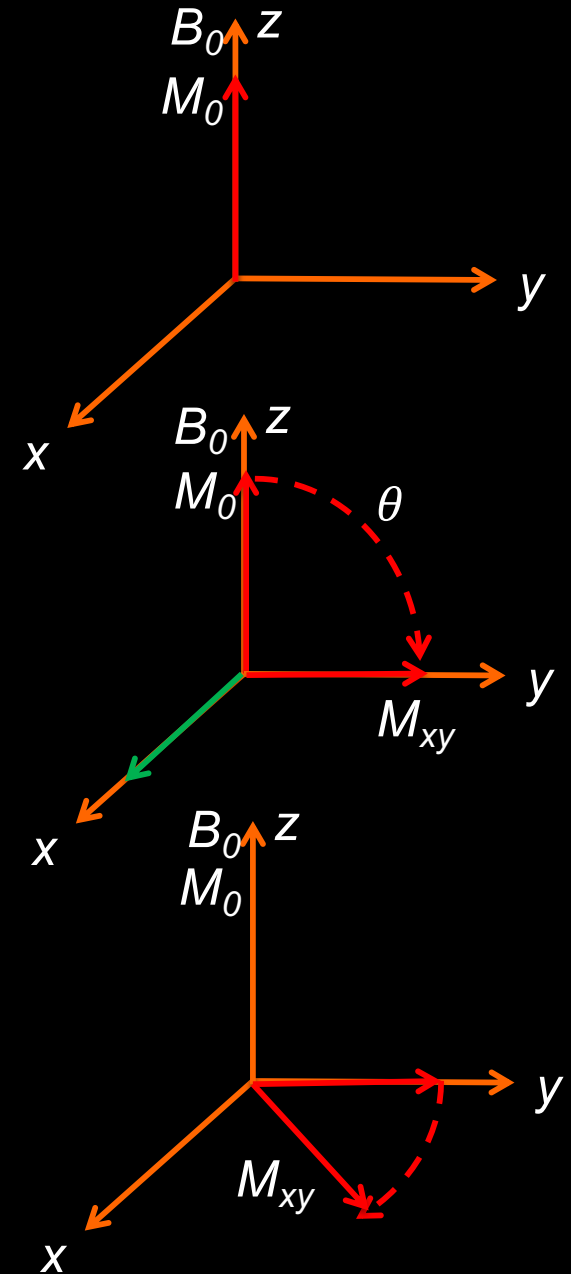
- Signal excitation

- $|M_{xy}| = M_0 \sin \theta$

- Signal detection

- $s(t) = M_0 \sin \theta e^{-i\omega_0 t}$

- $\omega_0 = \gamma B_0$



# Magnetic Resonance Imaging

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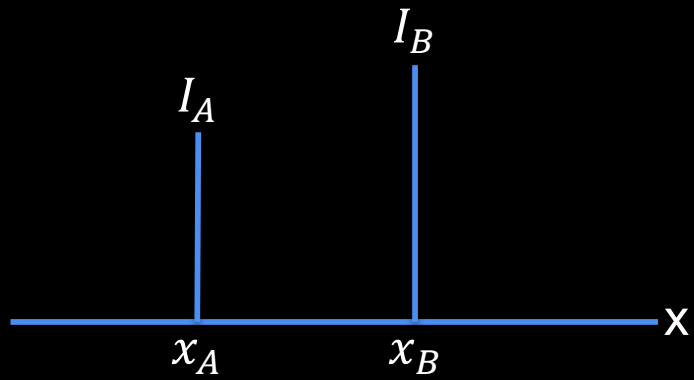
## MR Imaging



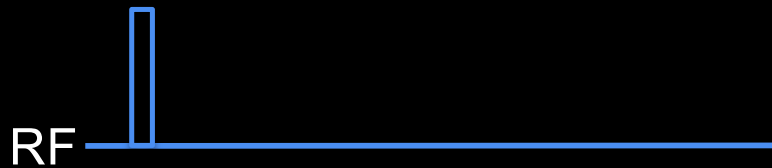


# 1D Imaging

$I(x)$

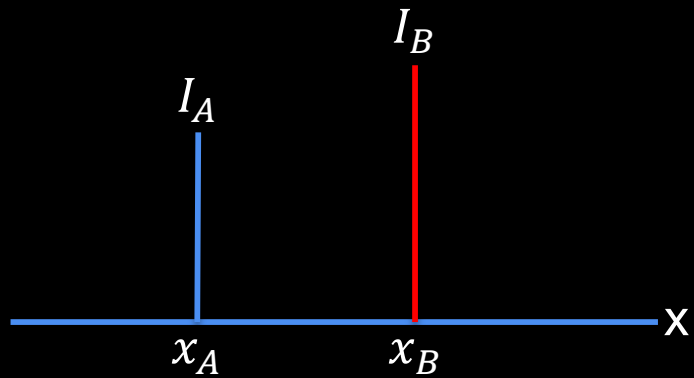


$$s(t) = I_A e^{-i\omega_0 t} + I_B e^{-i\omega_0 t}$$



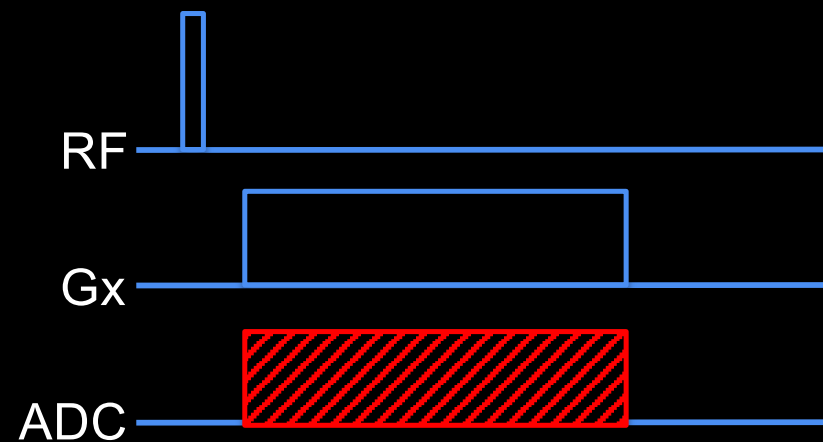
# 1D Imaging

$I(x)$



$$s(t) = (I_A e^{-i\gamma G_x x_A t} + I_B e^{-i\gamma G_x x_B t}) e^{-i\omega_0 t}$$

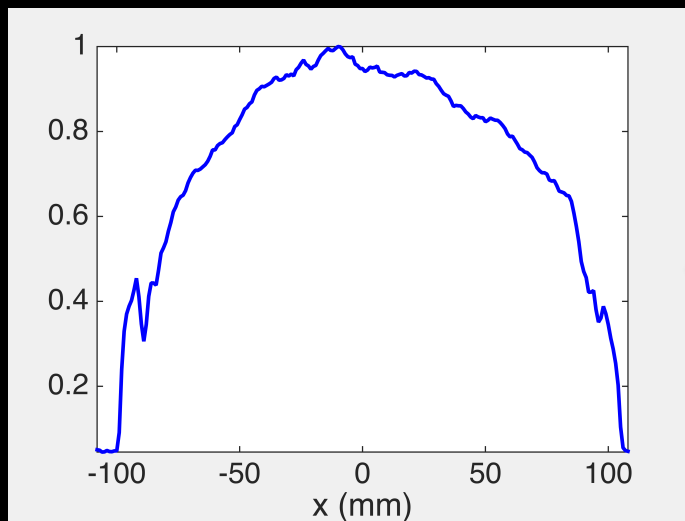
$$s(t) = I_A e^{-i\gamma G_x x_A t} + I_B e^{-i\gamma G_x x_B t}$$



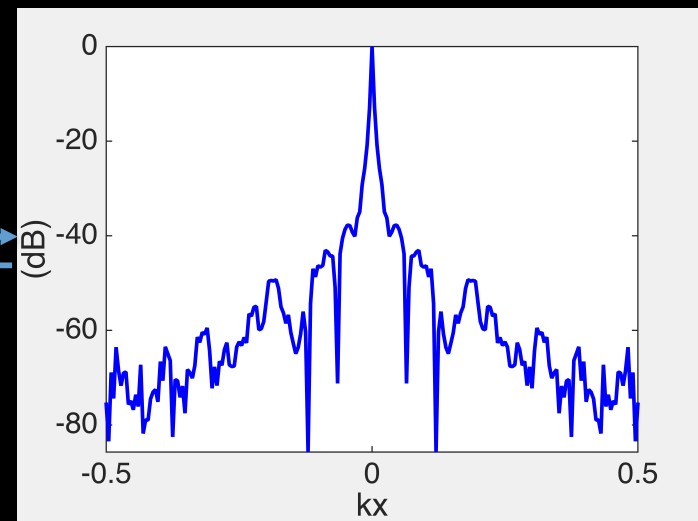
# 1D Imaging

$$s(t) = \int_x I(x) e^{-i\gamma G_x x t} dx$$

$$s(k_x) = \int_x I(x) e^{-i2\pi x \cdot k_x} dx, k_x = \frac{\gamma G_x t}{2\pi}$$



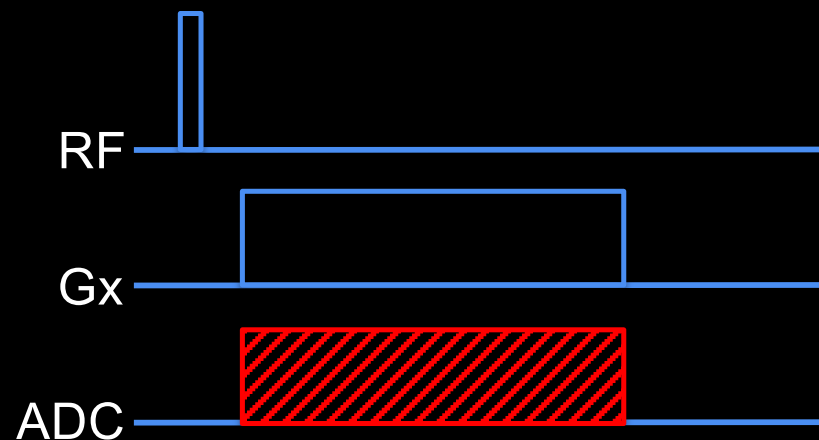
$I(x)$



$s(k_x)$

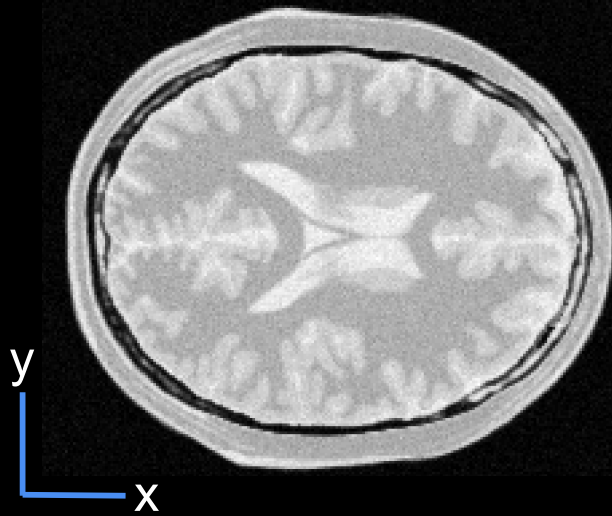
FT

Inverse FT



# 2D imaging

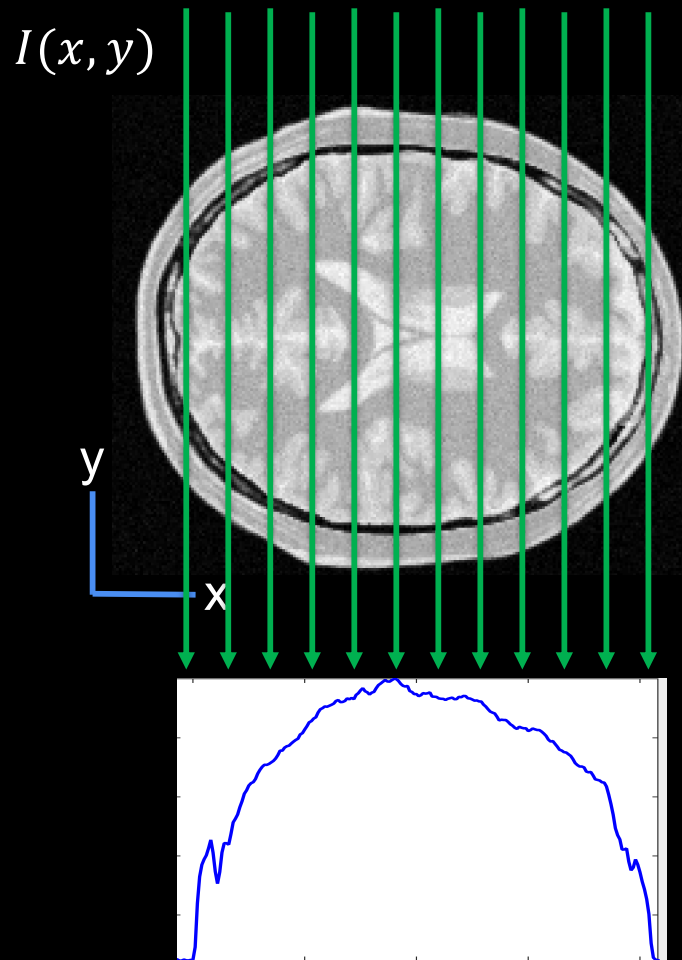
$I(x, y)$



$$s(t) = \iint_{x,y} I(x, y) e^{-i2\pi x \cdot k_x} dx dy, \quad k_x = \frac{\gamma G_x t}{2\pi}$$



# 2D imaging

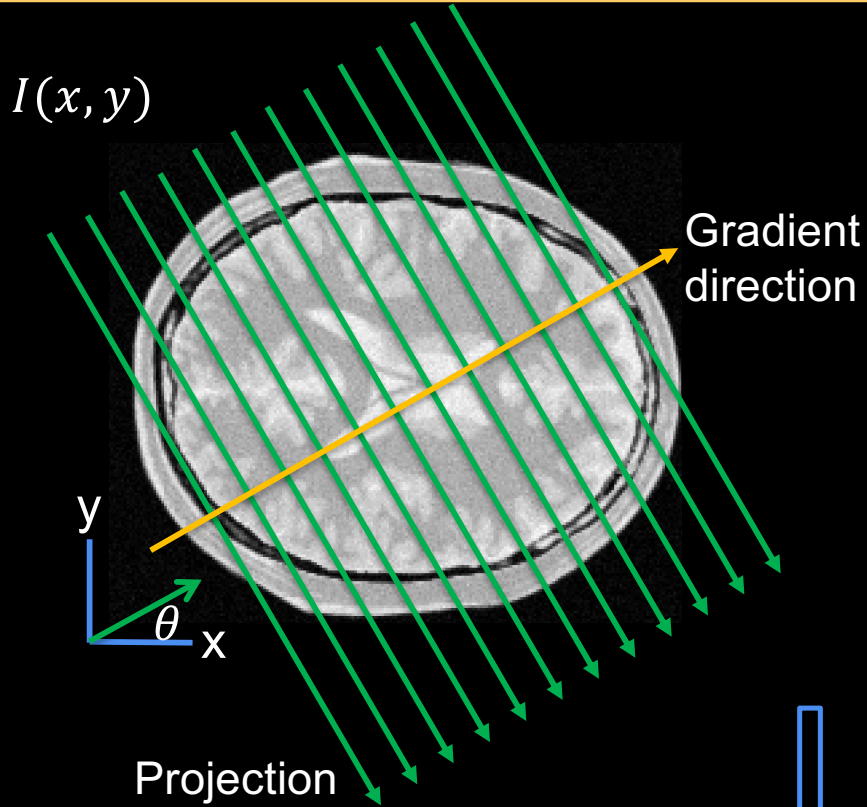


$$s(t) = \iint_{x,y} I(x, y) e^{-i2\pi x \cdot k_x} dx dy, \quad k_x = \frac{\gamma G_x t}{2\pi}$$

$$s(k_x) = \int_x \underbrace{\left( \int_y I(x, y) dy \right)}_{\text{Projection along } y!} e^{-i2\pi x \cdot k_x} dx$$



# 2D imaging



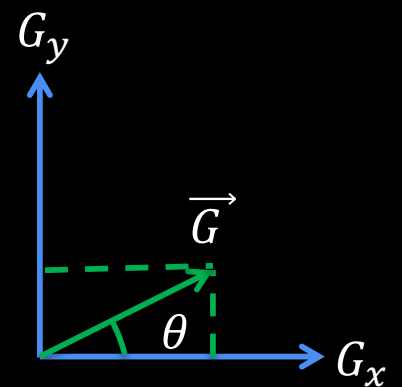
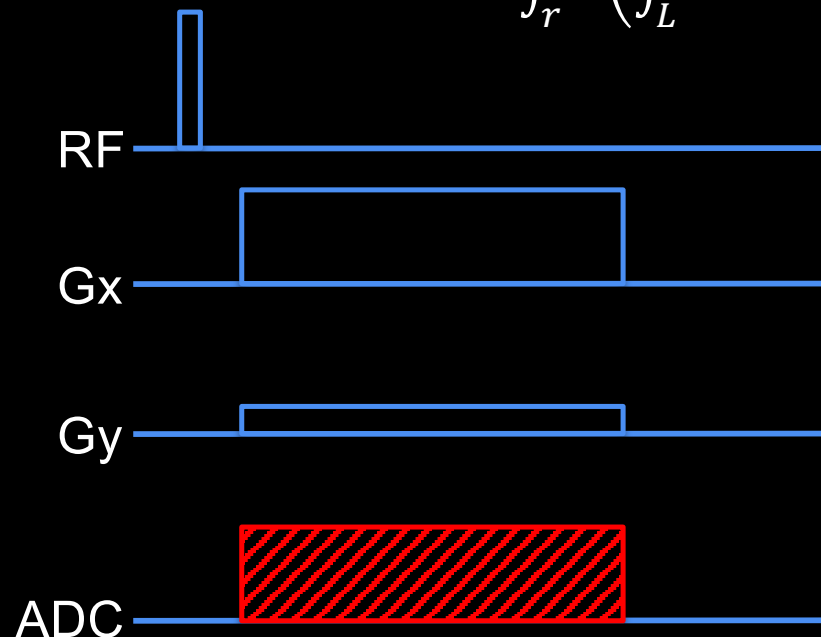
$$s(t, \theta) = \iint_{x,y} I(x, y) e^{-i\gamma|\vec{G}|(x \cos \theta + y \sin \theta)t} dx dy$$

$$= \int_r \left( \int_L I(x, y) dl \right) e^{-i\gamma|\vec{G}|rt} dr$$

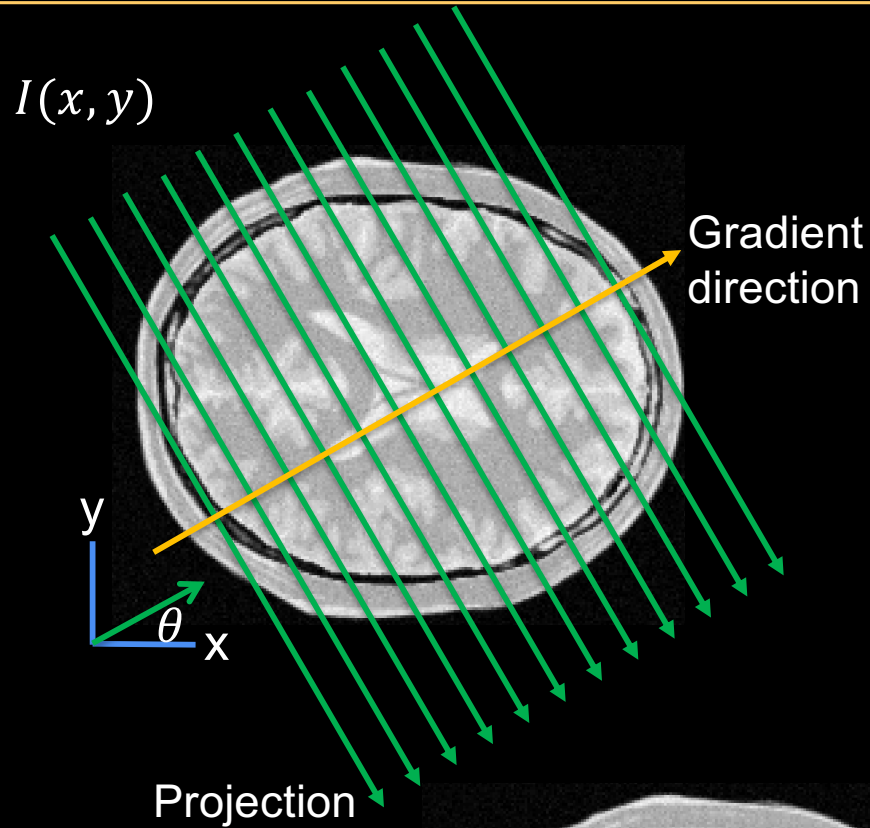
Projection along:

$$x \cos \theta + y \sin \theta = r$$

$$s(k, \theta) = \int_r \left( \int_L I(x, y) dl \right) e^{-i2\pi kr} dr, k = \frac{\gamma|\vec{G}|t}{2\pi}$$



# 2D imaging

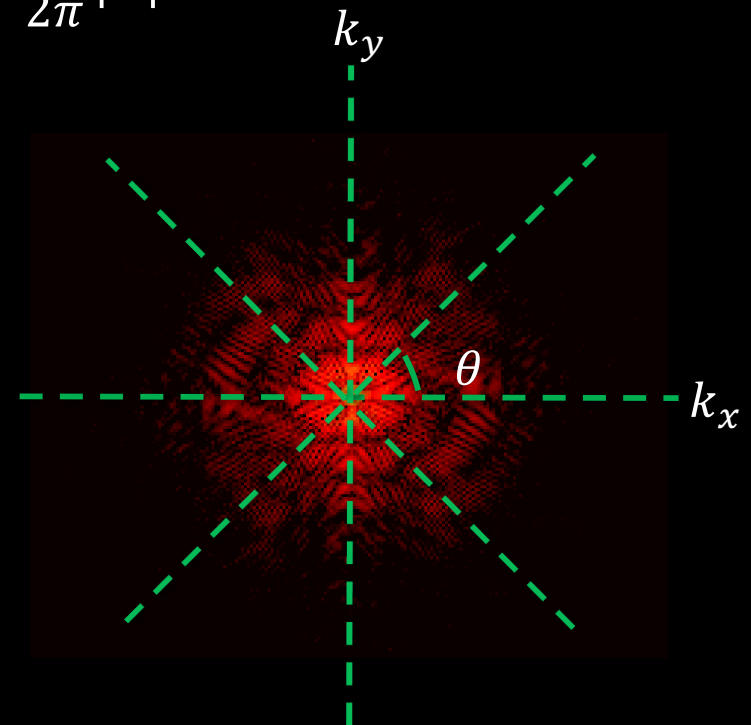
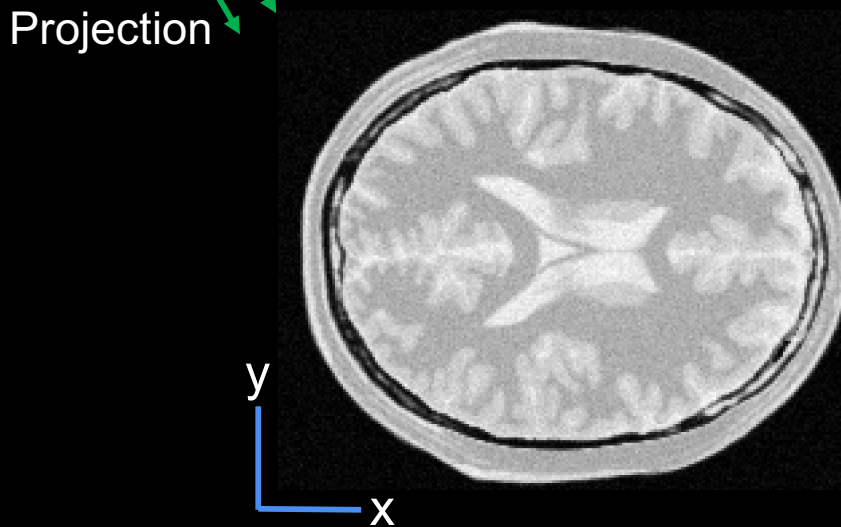


$$s(t, \theta) = \iint_{x,y} I(x, y) e^{-i\gamma |\vec{G}| (x \cos \theta + y \sin \theta) t} dx dy$$

$$s(k_x, k_y) = \iint_{x,y} I(x, y) e^{-i2\pi(xk_x + yk_y)} dx dy$$

$$k_y = \frac{\gamma}{2\pi} |\vec{G}| t \sin \theta$$

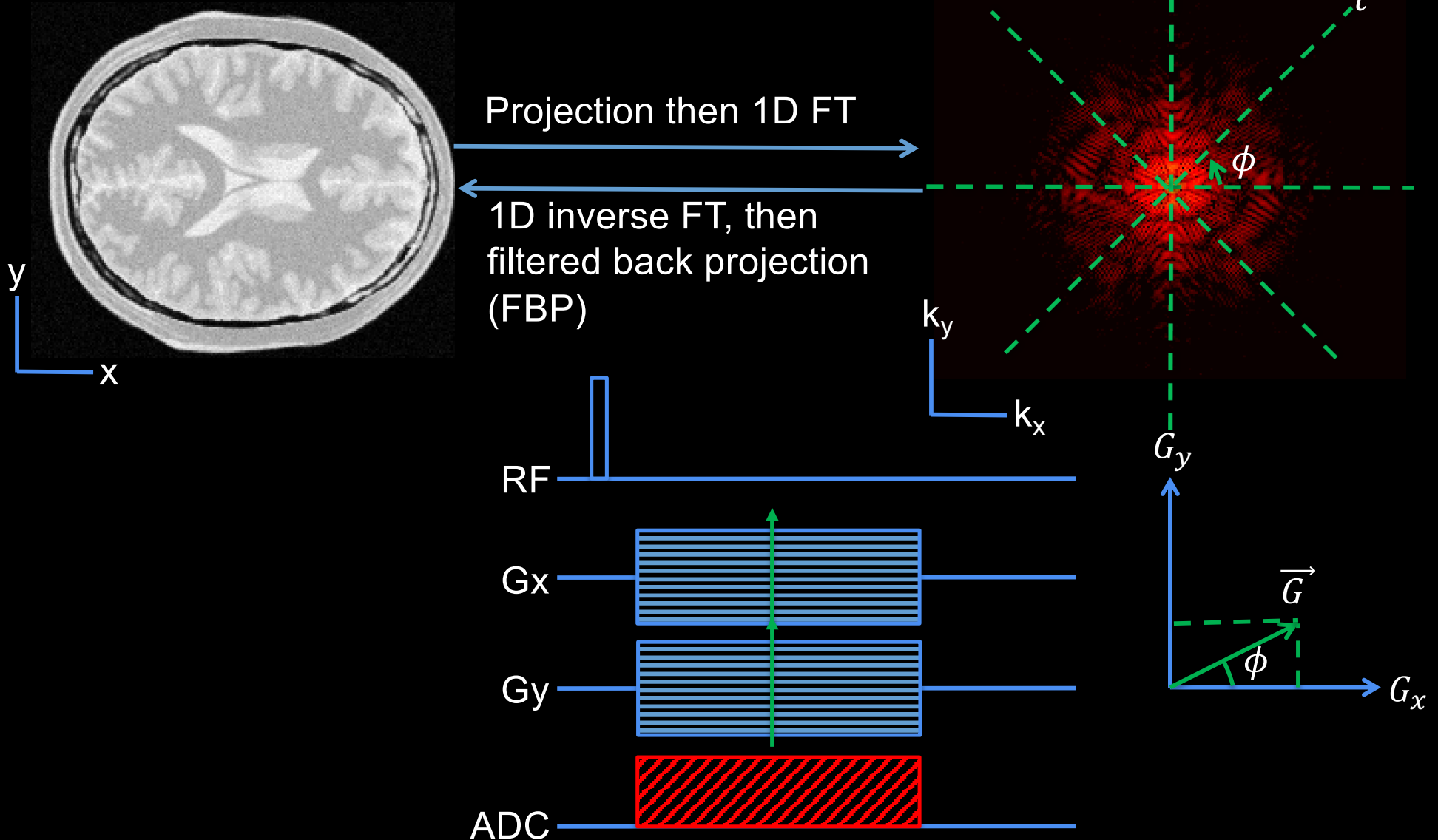
$$k_x = \frac{\gamma}{2\pi} |\vec{G}| t \cos \theta$$



# 2D imaging: Radial sampling

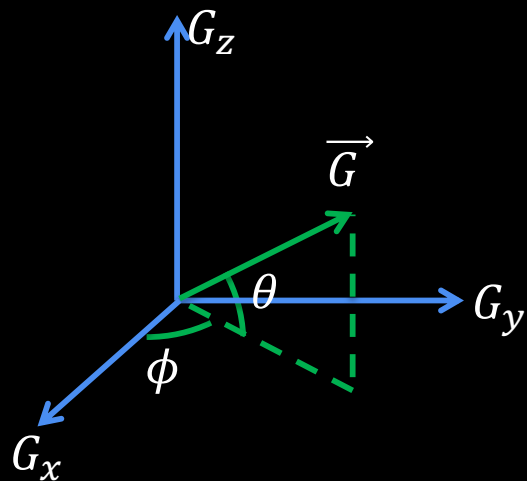
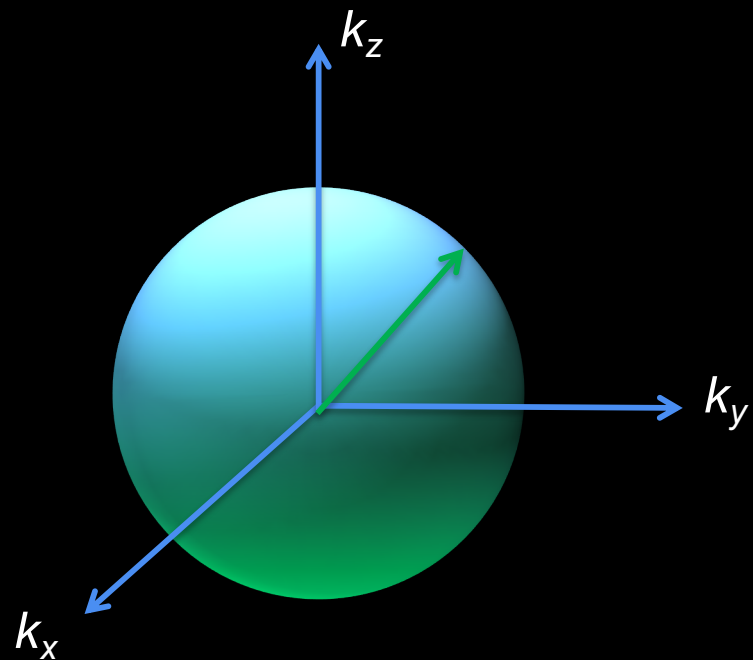
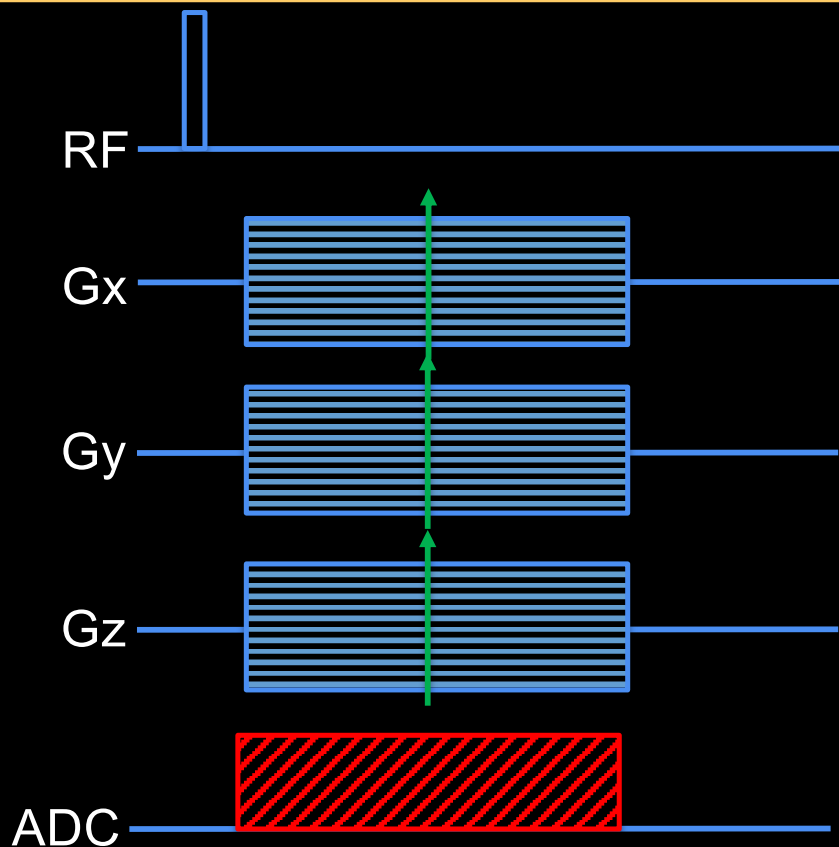
$I(x, y)$

$$s(t, \phi) = \iint_{x,y} I(x, y) e^{-i2\pi(x \cdot k_x + y \cdot k_y)} dx dy$$



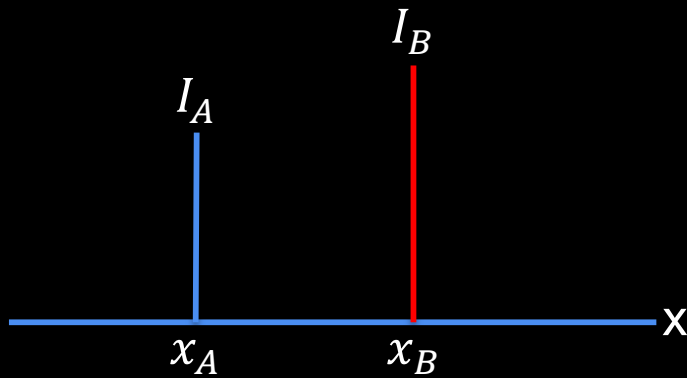


# 3D Imaging: Radial sampling



# 1D Imaging: Revisit

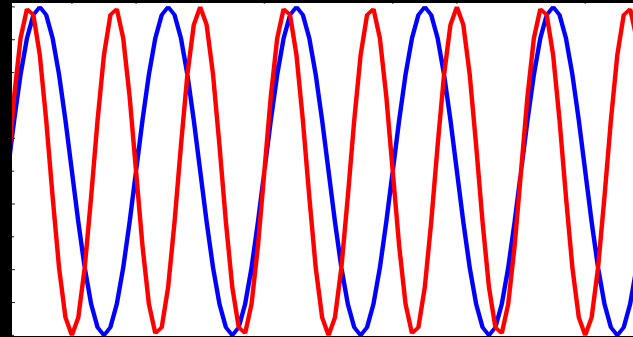
$I(x)$



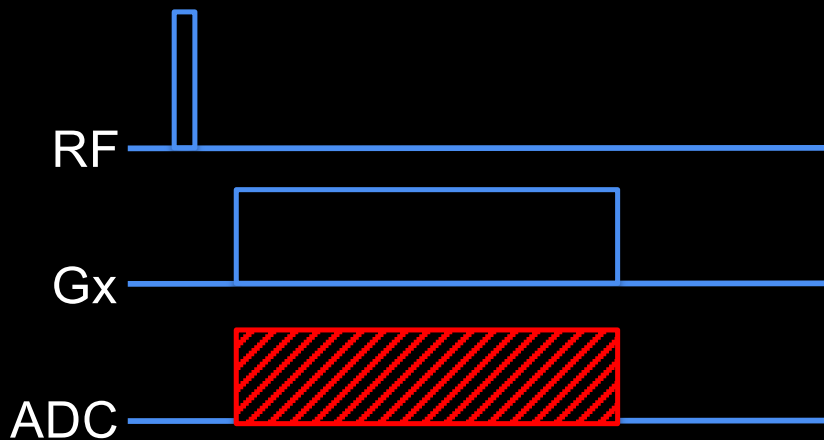
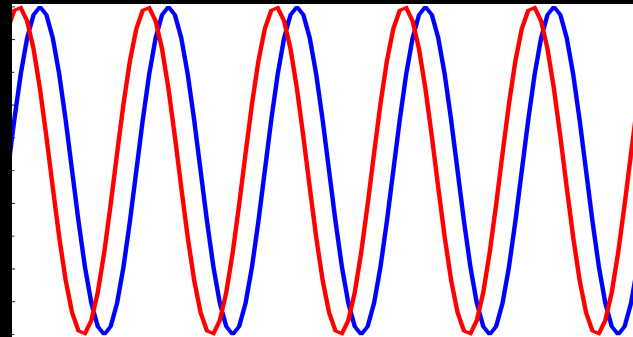
$$s(t) = I_A e^{-i\gamma G_x x_A t} + I_B e^{-i\gamma G_x x_B t}$$

Is it the only way to do it?

Frequency

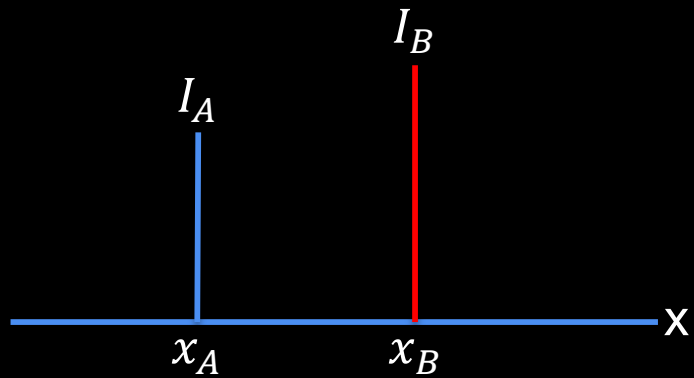


Phase?



# 1D Imaging: Revisit

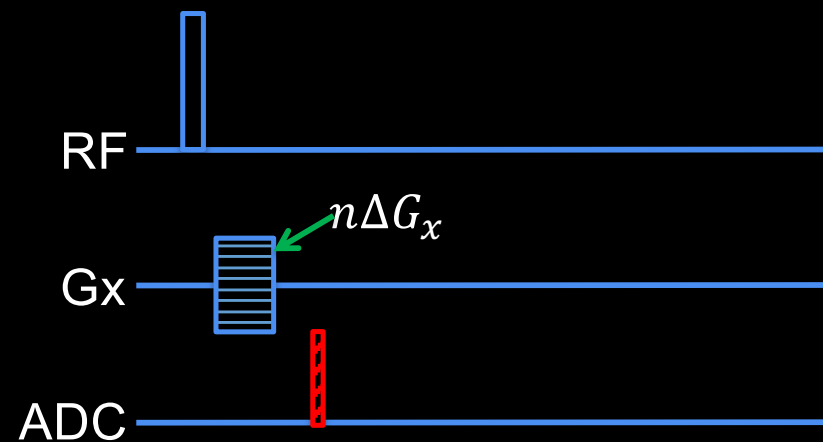
$I(x)$



$$s(n) = I_A e^{-i\varphi_A(n)} + I_B e^{-i\varphi_B(n)}$$

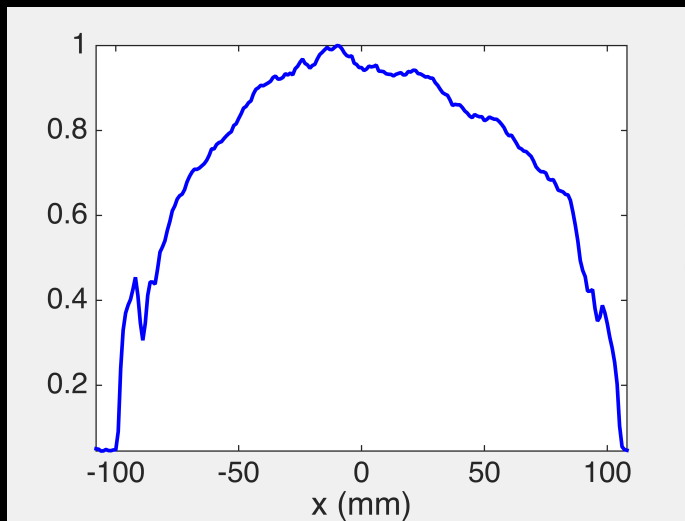
$$\varphi_A(n) = \gamma n \Delta G_x \tau x_A$$

$$\varphi_B(n) = \gamma n \Delta G_x \tau x_B$$



# 1D Imaging: Revisit

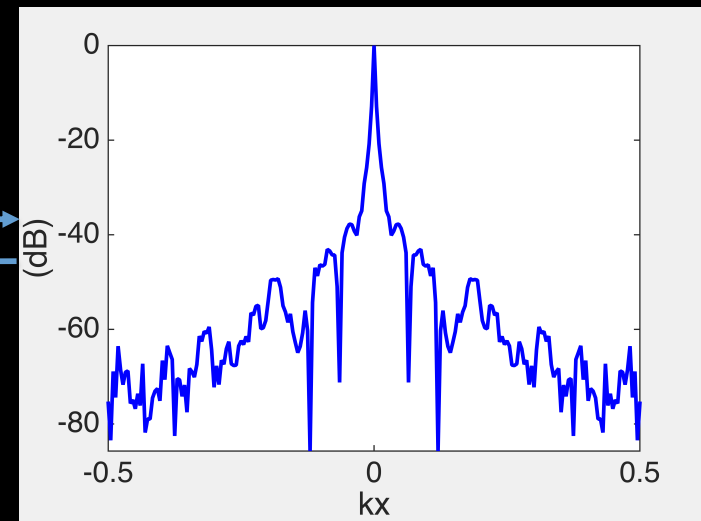
$$s(k_x) = \int_x I(x) e^{-i2\pi x \cdot k_x} dx,$$
$$k_x = \frac{\gamma n \Delta G_x \tau}{2\pi}$$



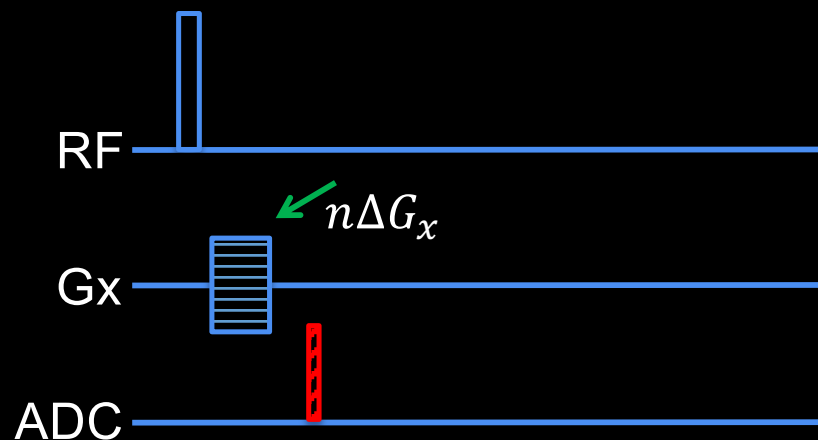
$I(x)$

FT

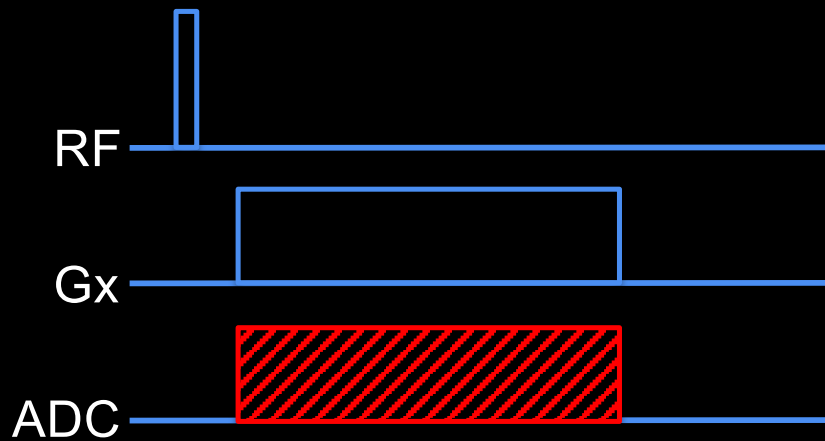
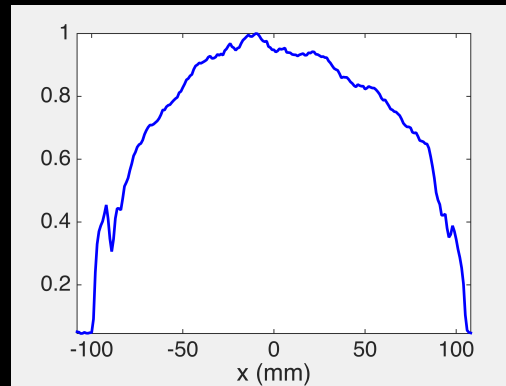
Inverse FT



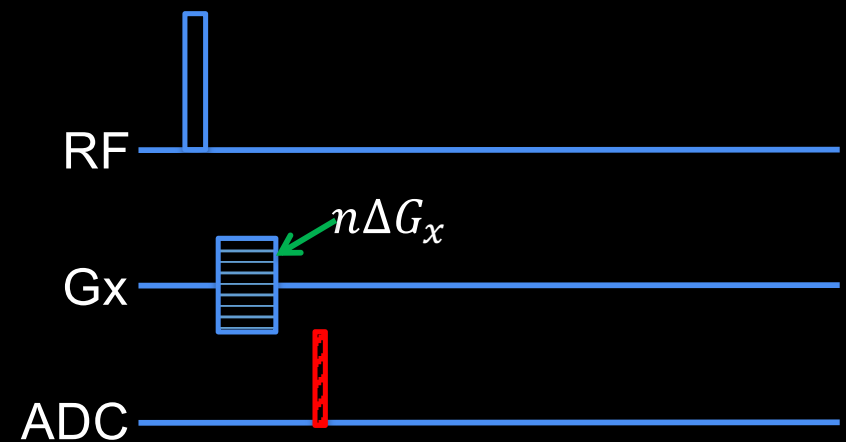
$s(k_x)$



# 1D imaging: Frequency vs. Phase encoding



$$s(n) = \int_x I(x) e^{-i2\pi x \cdot k_x} dx, k_x = \frac{\gamma n G_x \Delta t}{2\pi}$$

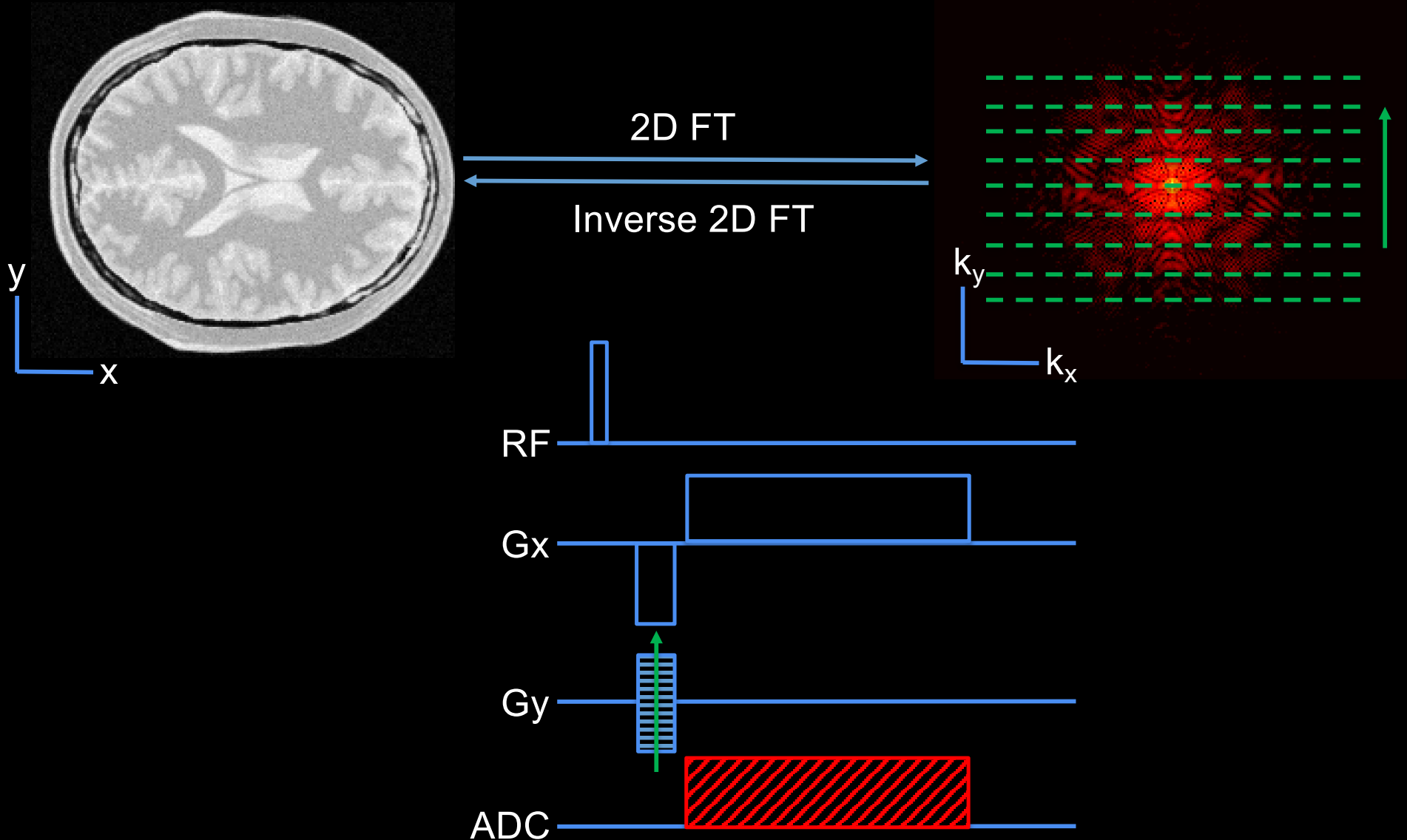


$$s(n) = \int_x I(x) e^{-i2\pi x \cdot k_x} dx, k_x = \frac{\gamma n \Delta G_x \tau}{2\pi}$$

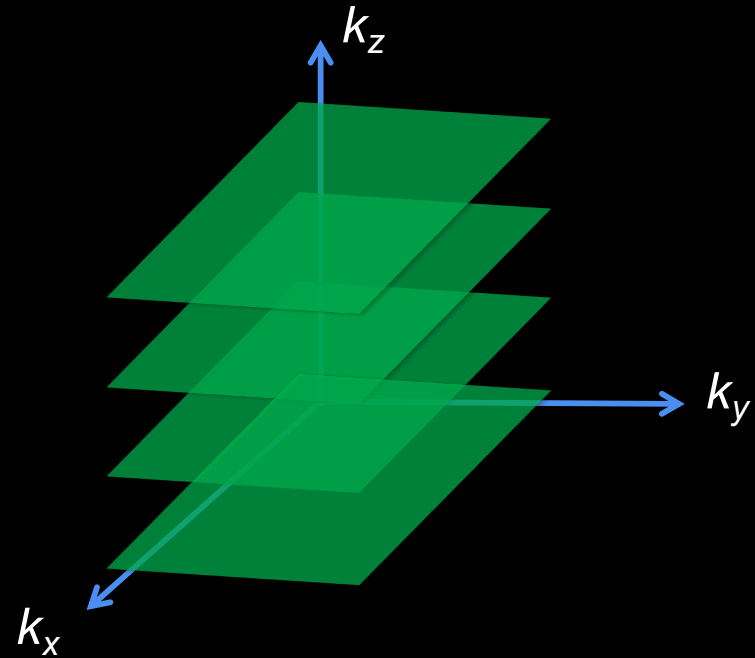
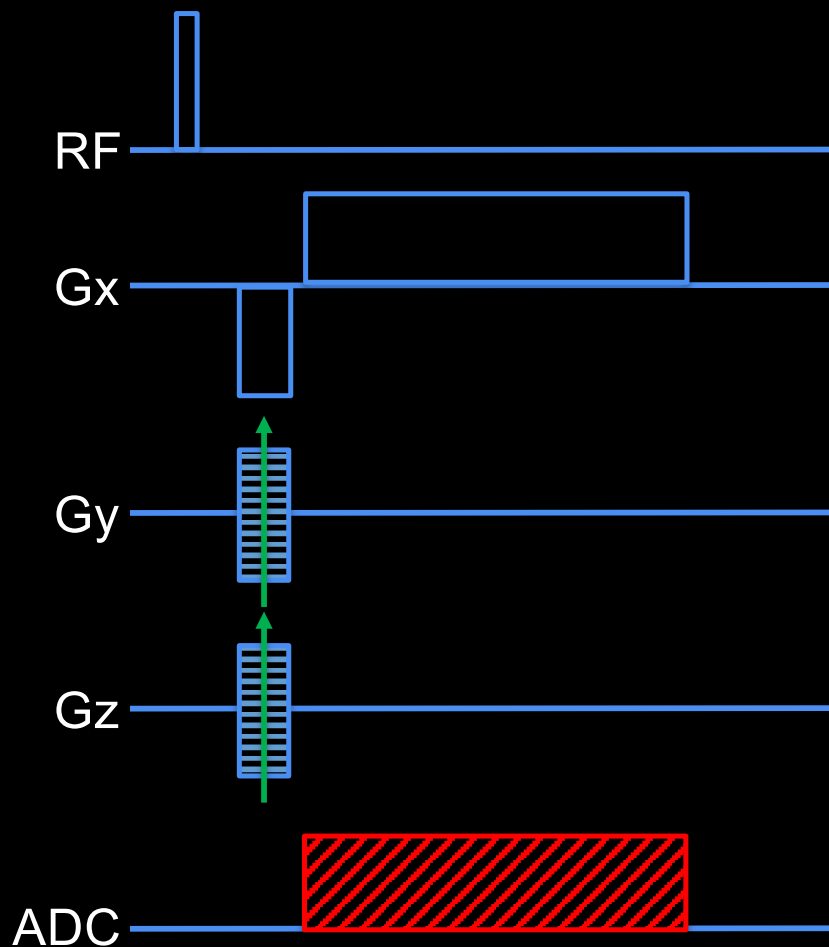
# 2D imaging: Cartesian sampling

$I(x, y)$

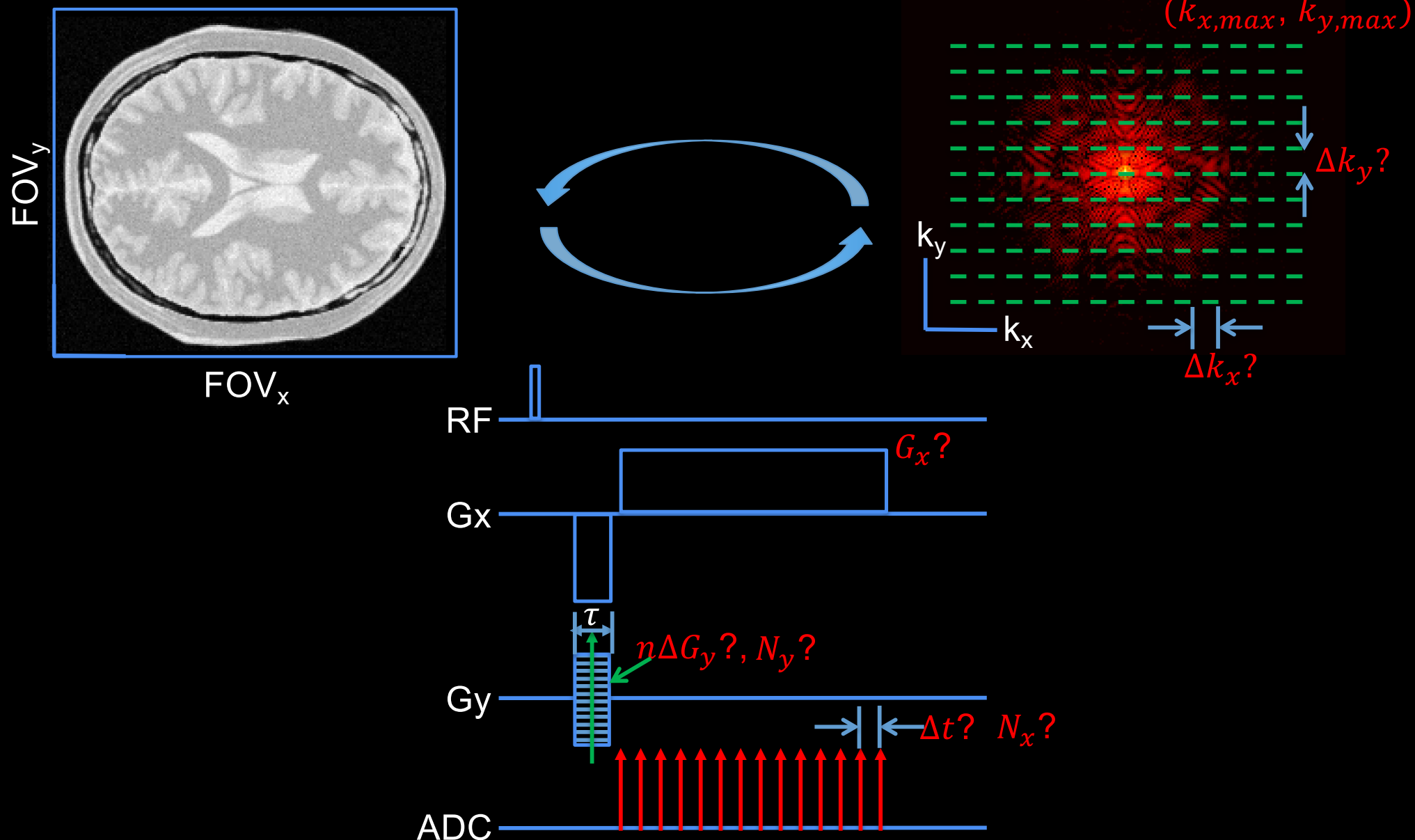
$$s(k_x, k_y) = \iint_{x,y} I(x, y) e^{-i2\pi(x \cdot k_x + y \cdot k_y)} dx dy$$



# 3D Imaging: Cartesian sampling

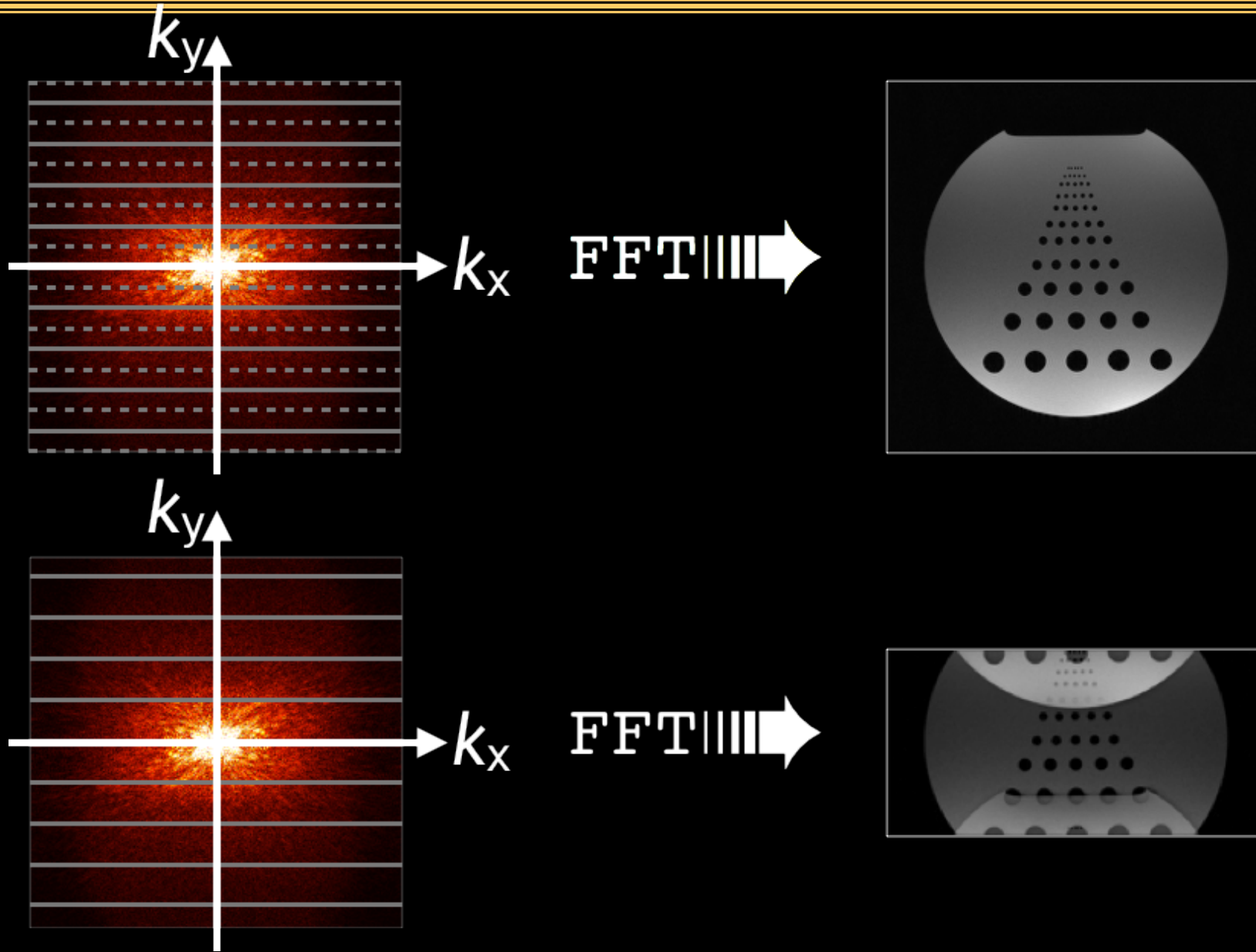


# k-space sampling

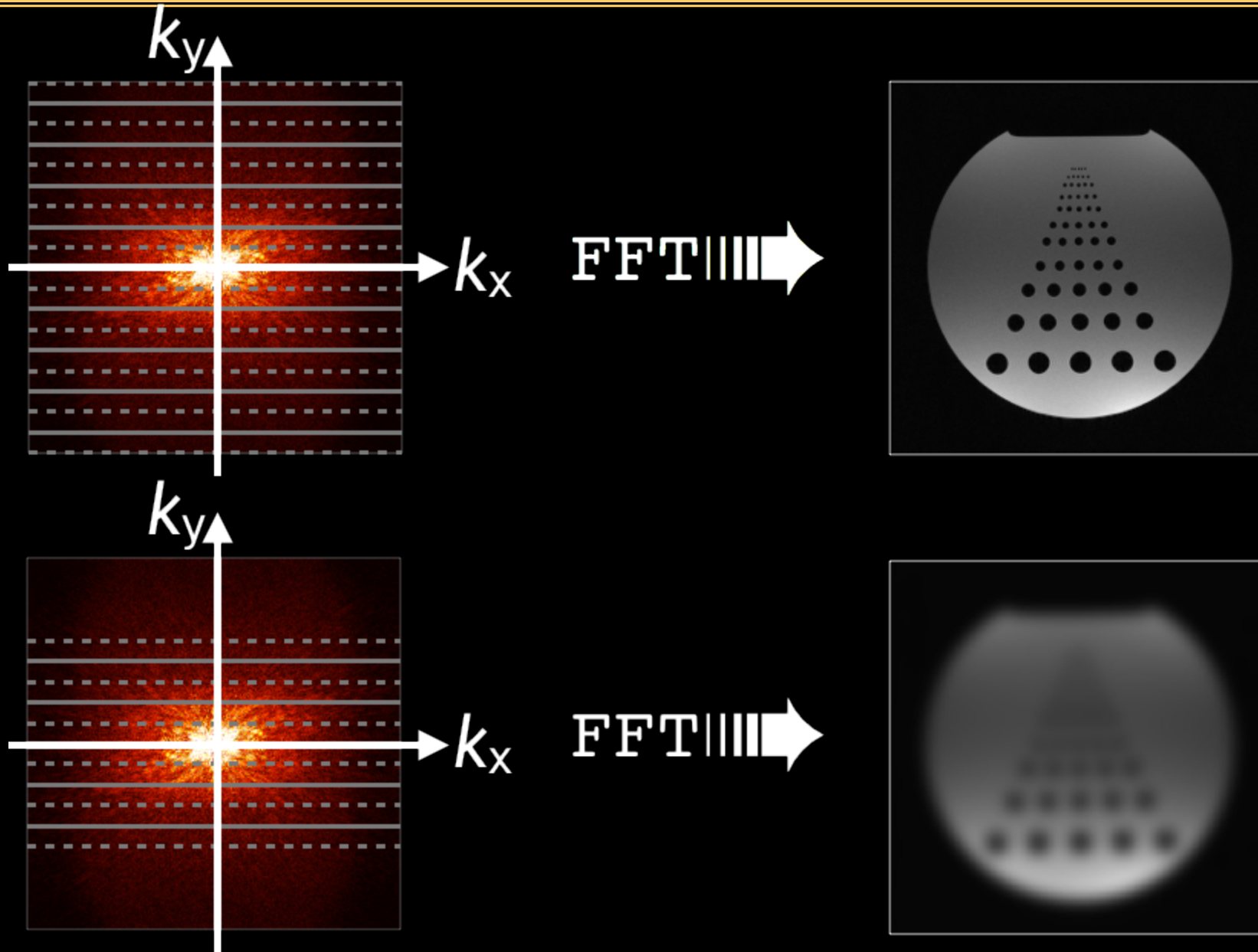




# k-space: Aliasing



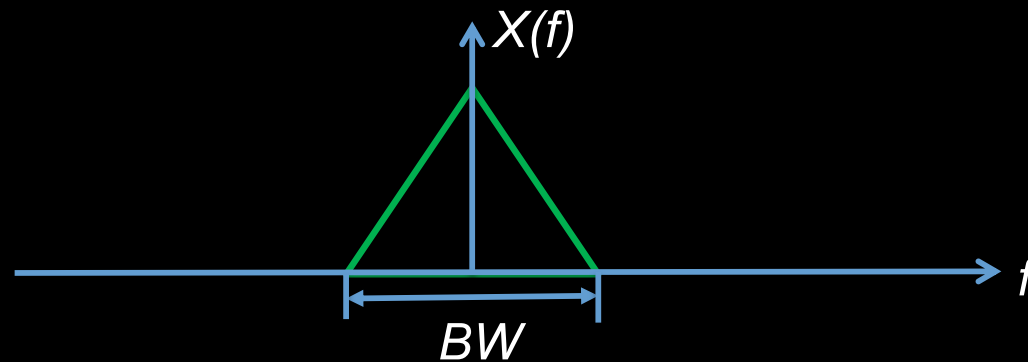
# k-space: Resolution



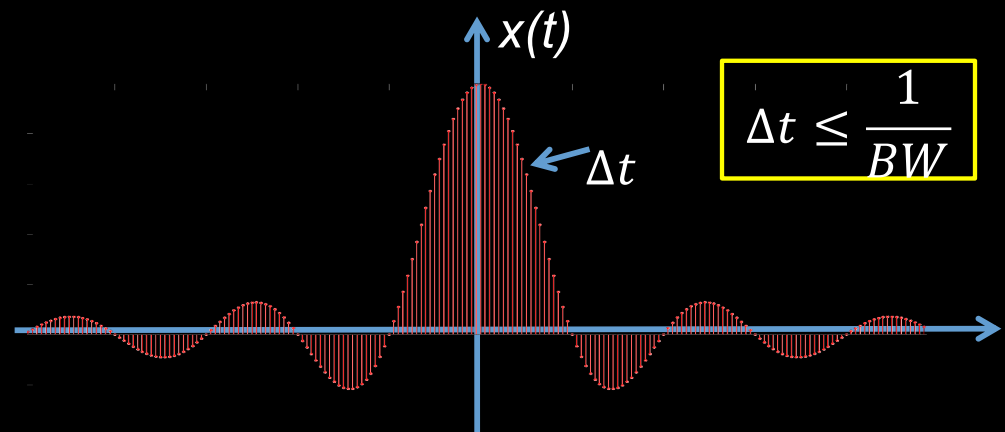
# k-space sampling: Field of View

- Nyquist sampling theorem

Frequency domain:

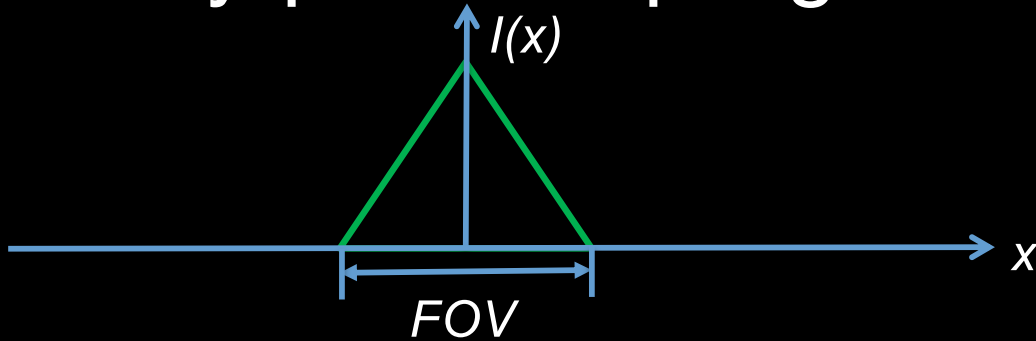


Time domain:

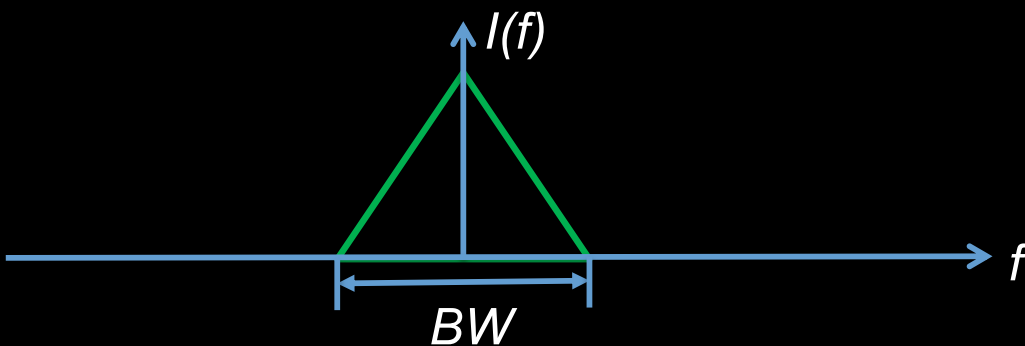


# k-space sampling: Field of View

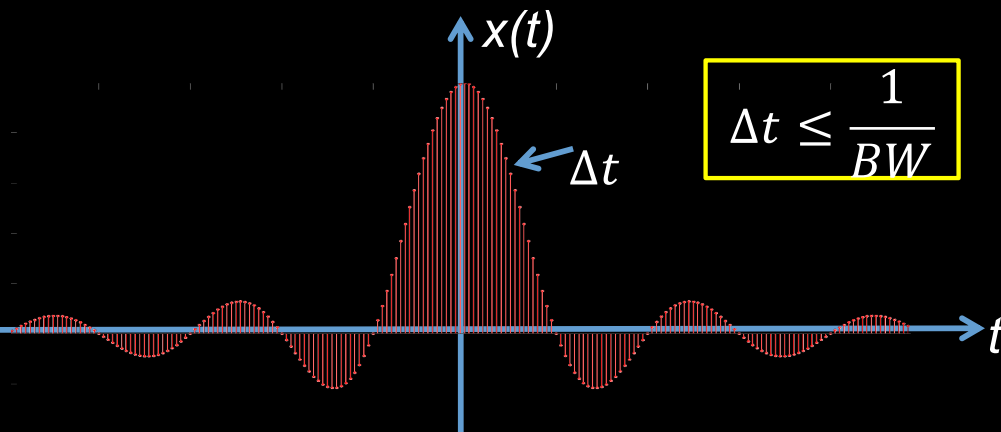
- Nyquist sampling theorem



$$I(x) = 0, \text{ if } |x| \leq \frac{FOV}{2}$$



$$I(f) = 0, \text{ if } |f| \leq \frac{BW}{2} = \frac{1}{2} \frac{\gamma G_x FOV}{2\pi}$$

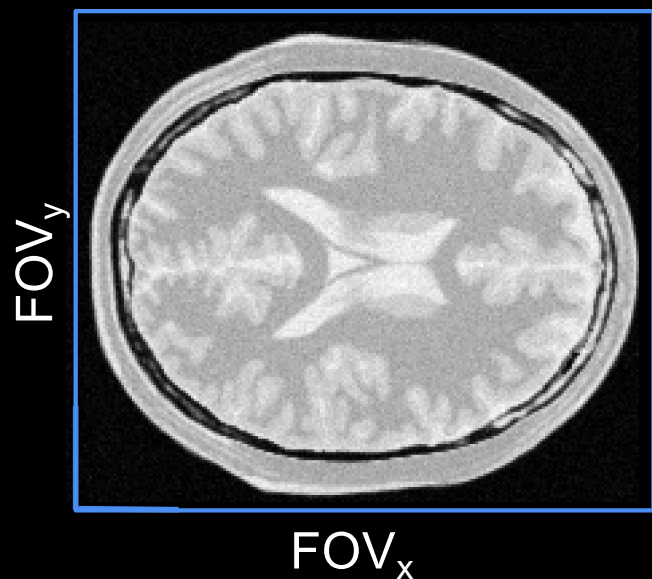


$$\Delta t \leq \frac{1}{BW} = \frac{1}{\gamma G_x FOV / 2\pi}$$

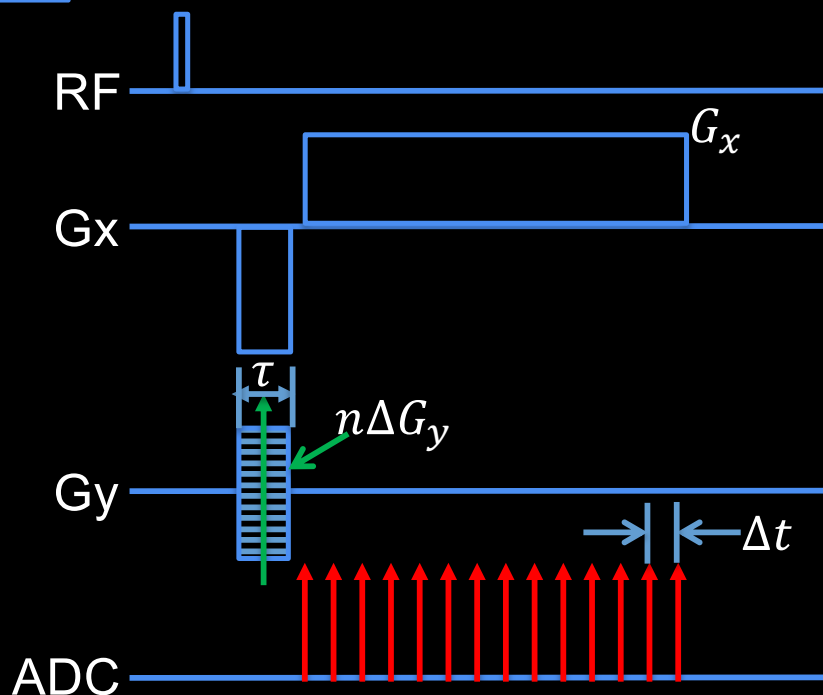
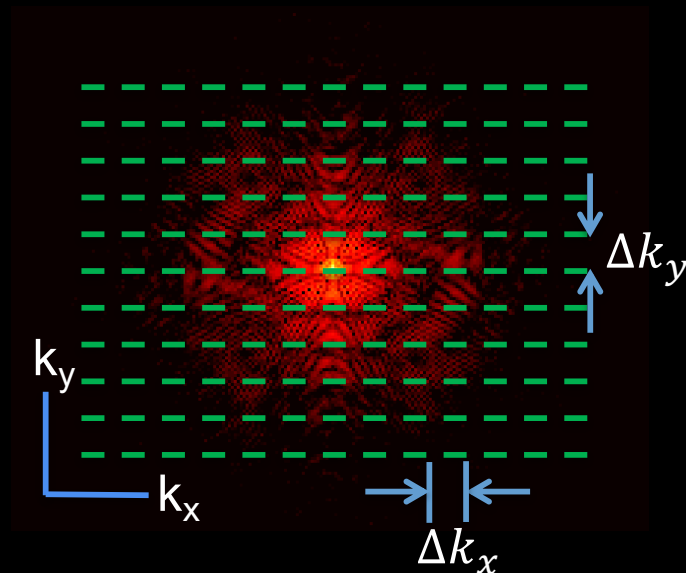
$$\Delta k_x = \frac{\gamma G_x \Delta t}{2\pi} \leq \frac{1}{FOV}$$

$$s(n\Delta k_x) = \int_x I(x) e^{-i2\pi x \cdot k_x} dx, k_x = \frac{\gamma n G_x \Delta t}{2\pi}$$

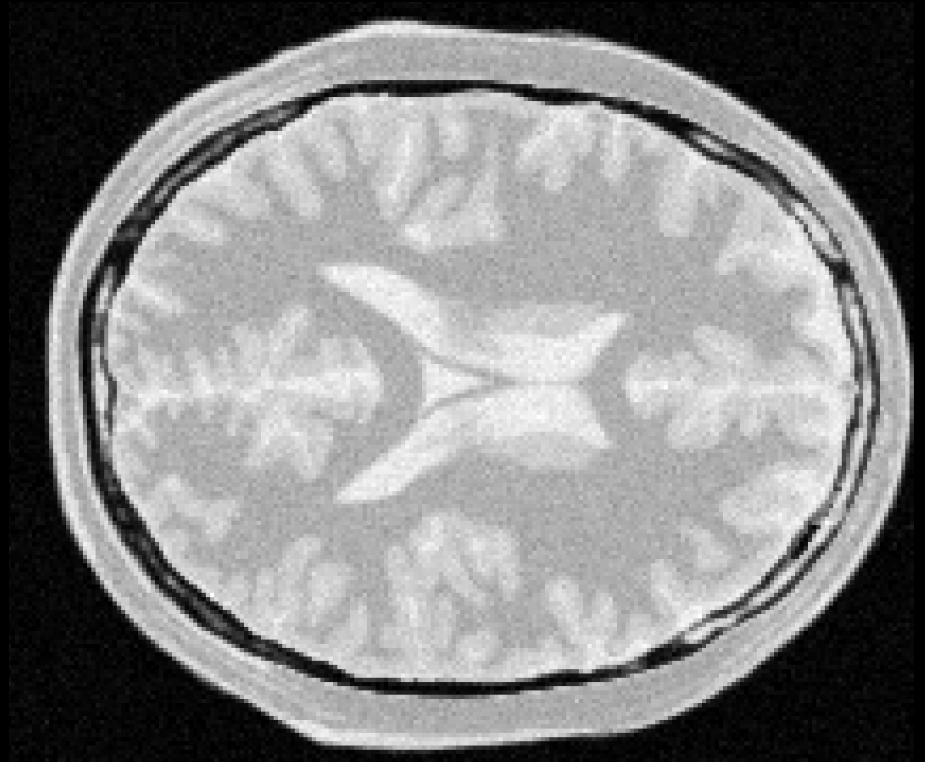
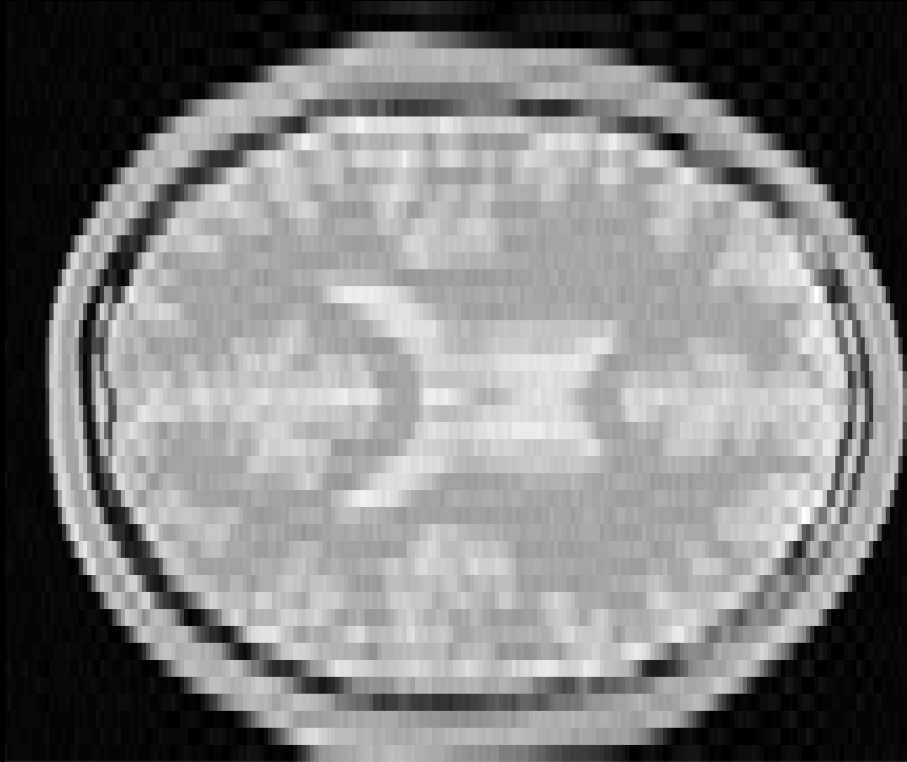
# k-space sampling: Field of View



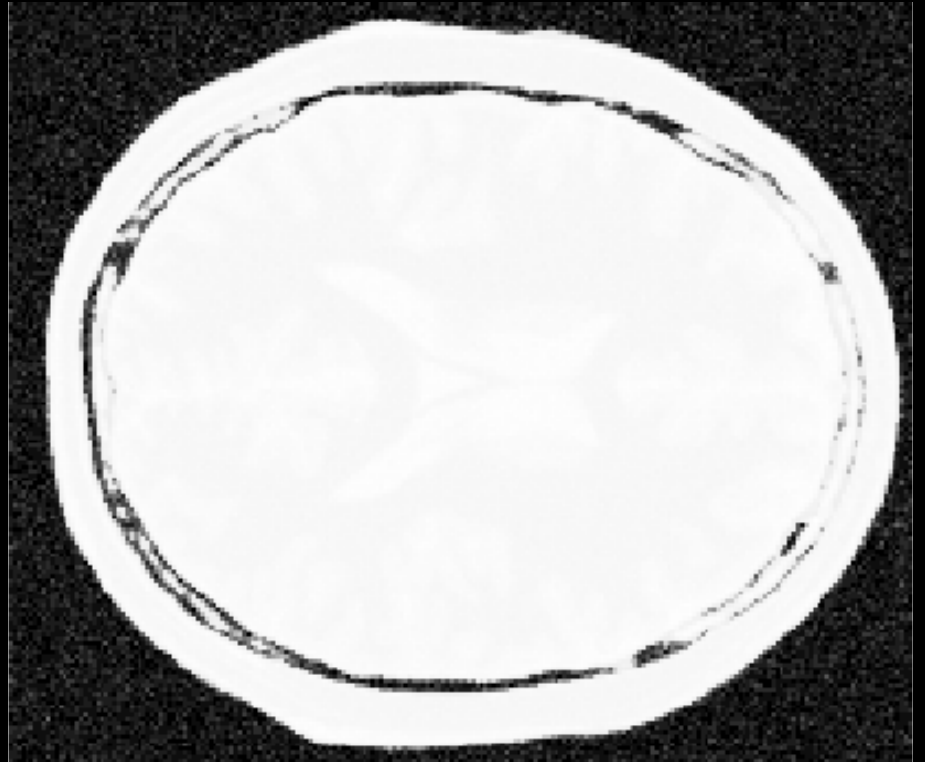
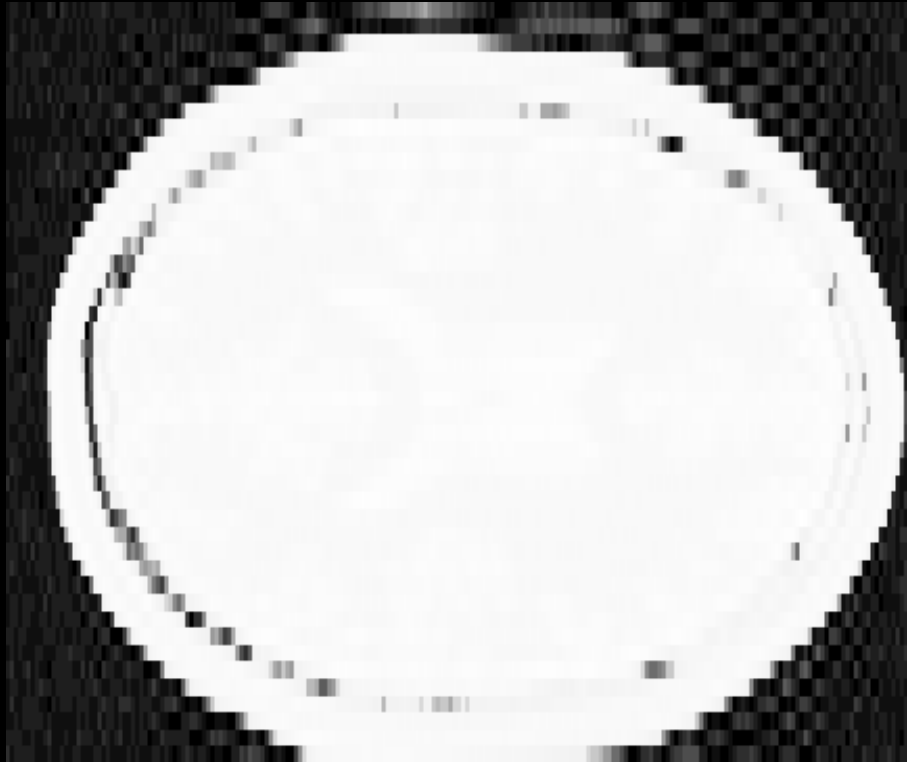
$$\Delta k_x = \frac{\gamma G_x \Delta t}{2\pi} \leq \frac{1}{FOV_x}$$
$$\Delta k_y = \frac{\gamma \Delta G_y \tau}{2\pi} \leq \frac{1}{FOV_y}$$



# k-space sampling: Resolution



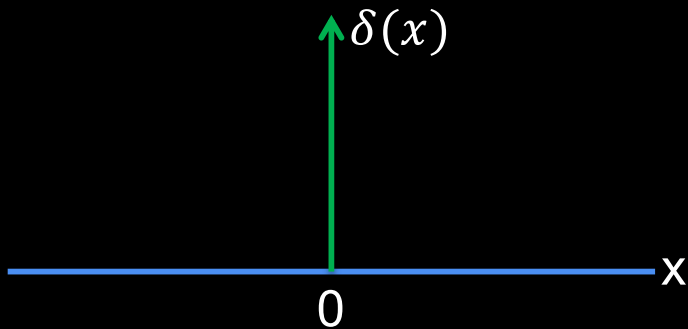
# k-space sampling: Resolution



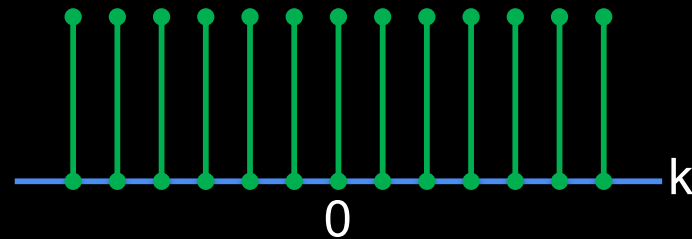
# k-space sampling: Resolution

- Point-spread-function

A point source in the image domain  $k$ -space



$$s(n\Delta k) = \int_{-\infty}^{\infty} \delta(x) e^{-i2\pi n\Delta kx} dx \equiv 1$$



Properties of Dirac delta function

$$\delta(x) = 0, \text{ for } x \neq 0$$

$$\int_{-\infty}^{\infty} \delta(x) dx = 1$$

$$\int_{-\infty}^{\infty} \delta(x) \varphi(x) dx = \varphi(0)$$

Reconstruction

$$h(x) = \Delta k \sum_{n=-N/2}^{N/2-1} e^{i2\pi n\Delta kx} = e^{-i\pi\Delta kx} \Delta k \frac{\sin(\pi N\Delta kx)}{\sin(\pi\Delta kx)}$$



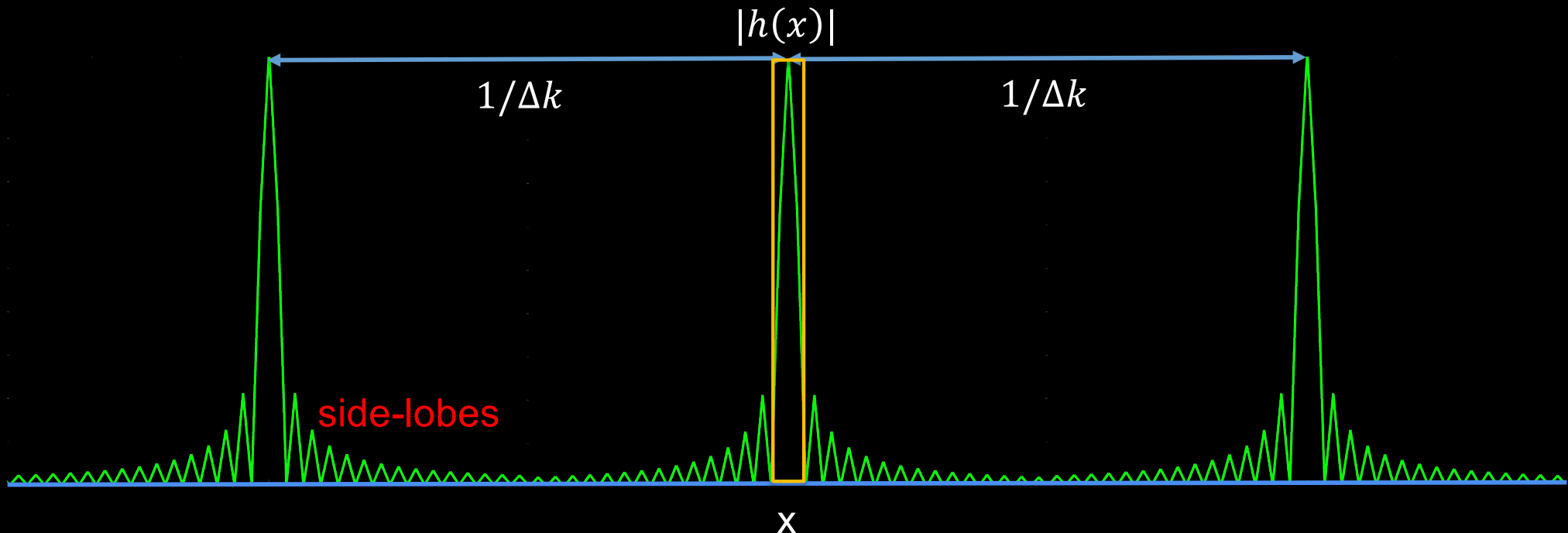
# k-space sampling: Resolution

- Point-spread-function

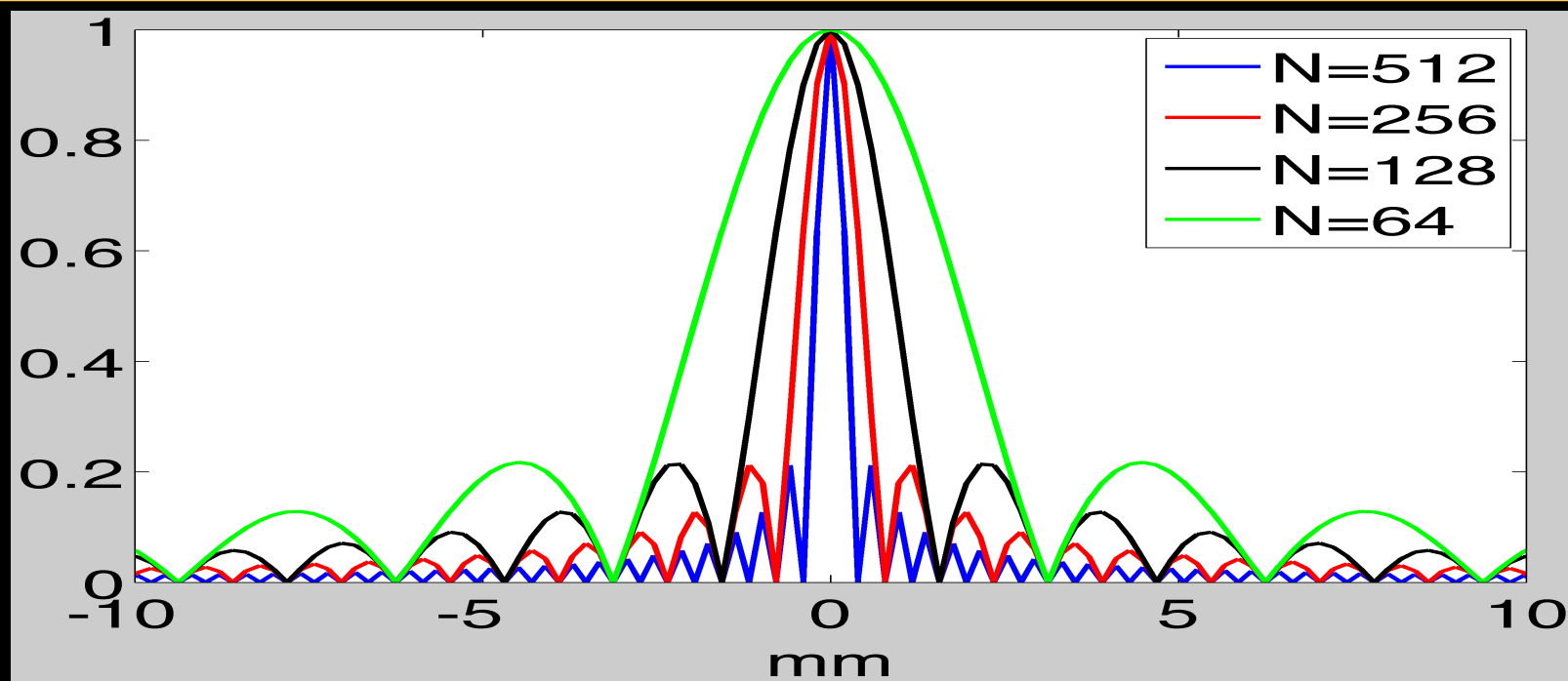
$$h(x) = e^{-i\pi\Delta kx} \Delta k \frac{\sin(\pi N\Delta kx)}{\sin(\pi\Delta kx)}$$

$$\hat{I}(x) = I(x) * h(x)$$

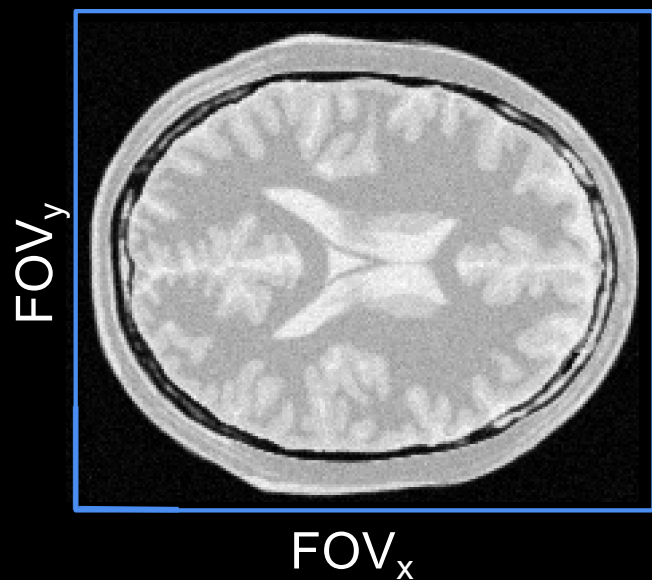
$$W_h = \frac{1}{h(0)} \int_{-FOV/2}^{FOV/2} h(x) dx = \frac{FOV}{N}$$



# k-space sampling: Resolution



# k-space sampling

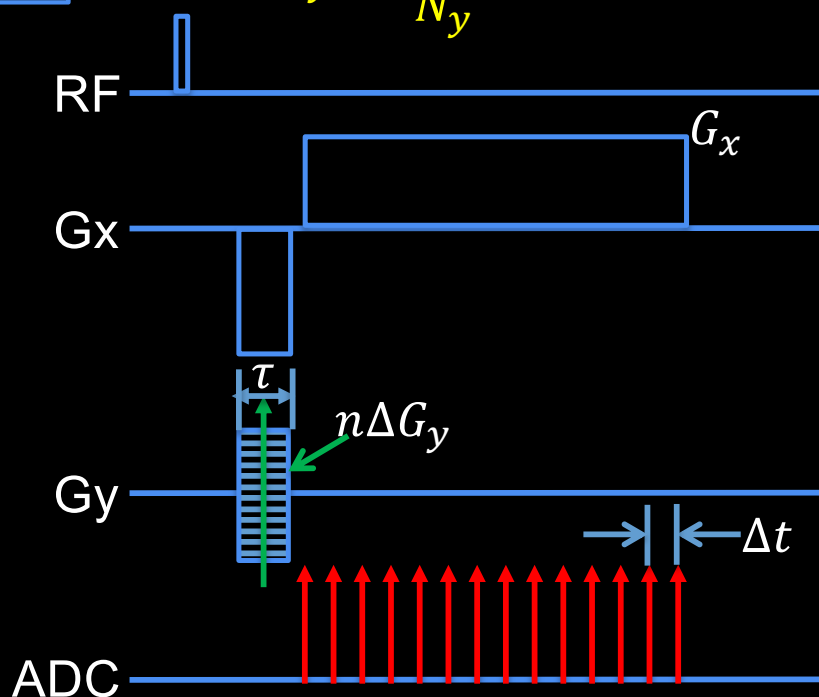
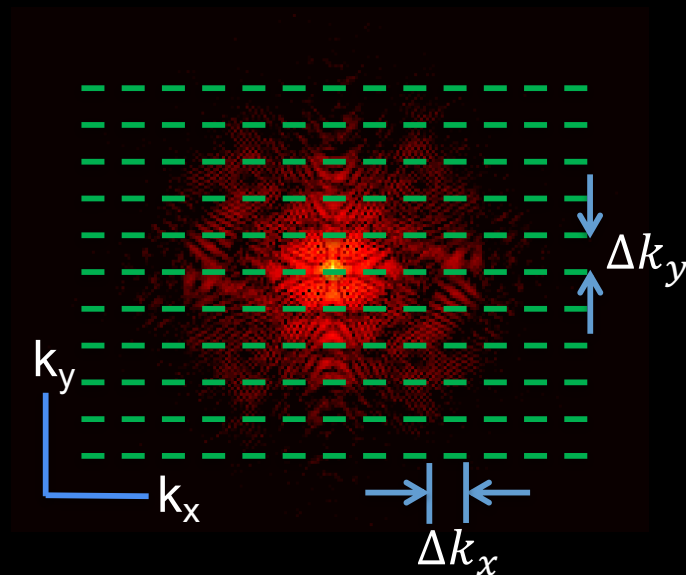


$$\Delta k_x = \frac{\gamma G_x \Delta t}{2\pi} \leq \frac{1}{FOV_x}$$

$$\Delta k_y = \frac{\gamma \Delta G_y \tau}{2\pi} \leq \frac{1}{FOV_y}$$

$$\Delta x = \frac{FOV_x}{N_x}$$

$$\Delta y = \frac{FOV_y}{N_y}$$



# Summary

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- NMR physics
  - Spins and Bloch equation
- MR signal excitation and reception
- MR signal encoding & decoding
  - Frequency encoding
  - Phase encoding
  - Projection-based MR imaging
  - Fourier transform-based MR imaging
- K-space sampling
  - Field of View
  - Resolution

# Omitted topics

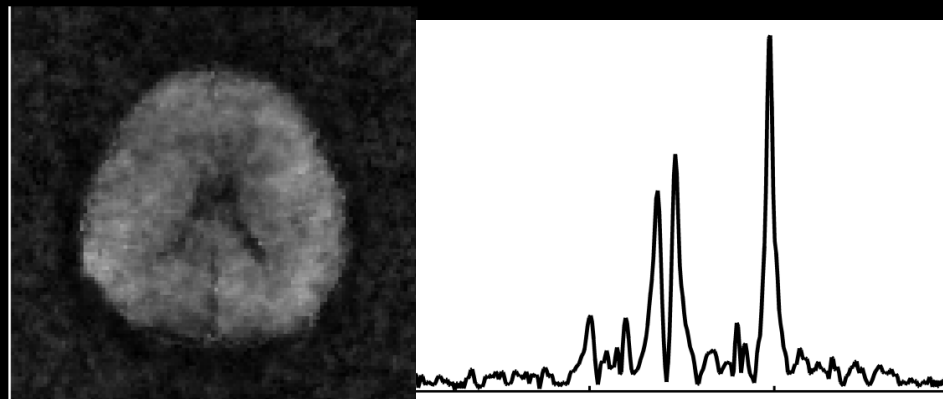
- MR physics
  - Noise and signal-to-noise ratio
  - RF pulse design
  - Fast imaging sequence
    - RARE, FLASH, bSSFP, BURST, ...
  - Contrast
    - BOLD signal, diffusion, perfusion, susceptibility, spectroscopic imaging, flow imaging, ...
- MR imaging
  - k-space sampling trajectories and non-Cartesian reconstruction
  - Constrained image reconstruction
    - Parallel imaging, compressed sensing, low-rank matrix/tensor based reconstruction, model based reconstruction, ...
  - Parameter mapping
    - T1, T2, T2\*, B0, B1, water/fat separation, MR fingerprinting, ...
- MR hardware
  - Gradient coil design
  - RF coil design
  - Hardware imperfection and compensation
  - MRI using Earth's field
- Applications of MR

# Future readings

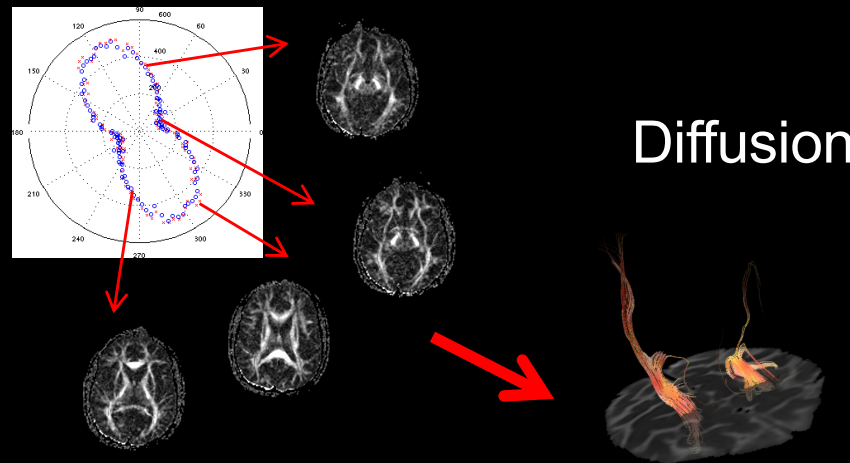
- Quantum mechanics
  - M. H. Levitt, *Spin Dynamics: Basics of Nuclear Magnetic Resonance* (2<sup>nd</sup> Edition), John Wiley & Sons, Ltd., 2018.
- EM field in MR
  - J. M. Jin, *Analysis and Design in Magnetic Resonance Imaging*, CRC Press, 1998.
- MRI
  - Z.-P. Liang & P. C. Lauterbur, *Principles of Magnetic Resonance Imaging: A Signal Processing Perspective*, IEEE Press/John Wiley, 1999.
  - R. W. Brown, et al., *Magnetic Resonance Imaging: Physical Principles and Sequence Design* (2<sup>nd</sup> Edition), Wiley-Blackwell, 2014
- MRI sequence programming
  - M. A. Bernstein, et al., *Handbook of MRI Pulse Sequences*, Academic Press, 2004.

# There are much more!

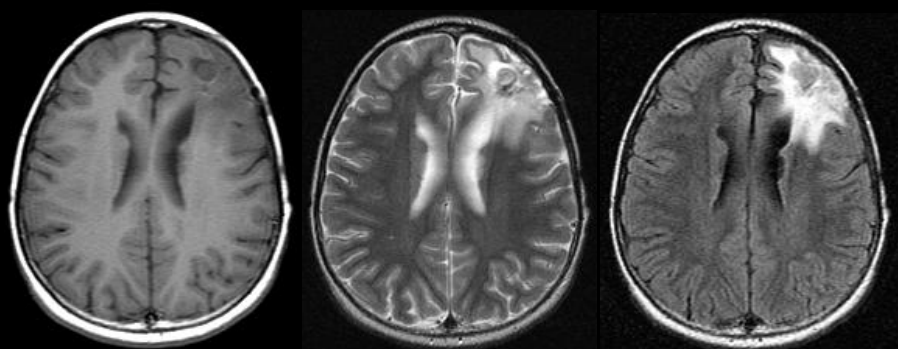
*There is nothing that nuclear spins will not do for you, as long as you treat them as human beings.* -- Erwin Hahn



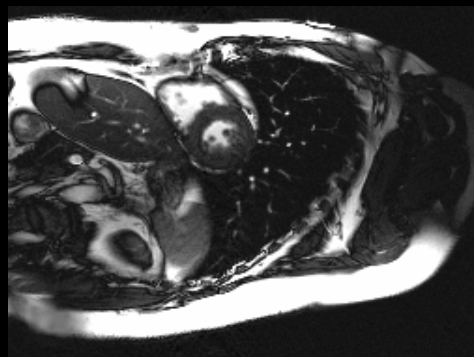
MRSI



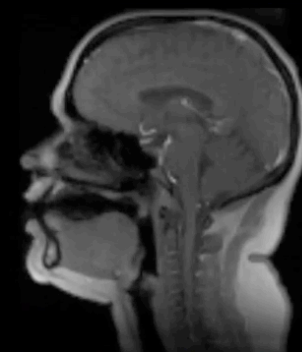
Diffusion



T1, T2 changes



Cardiac imaging



Speech imaging

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*Thank you for your attention.*