

# Searching for New Physics at Neutrino Experiments

Asli M. Abdullahi

Dark Interactions

*New Perspectives from Theory and Experiment*

16th November 2022



U.S. DEPARTMENT OF  
**ENERGY**

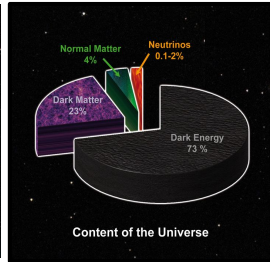
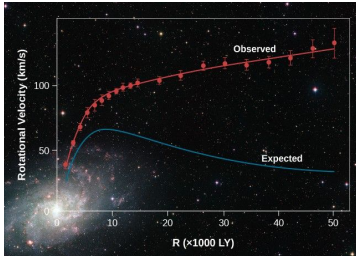
Office of  
Science

# Evidence for New Physics

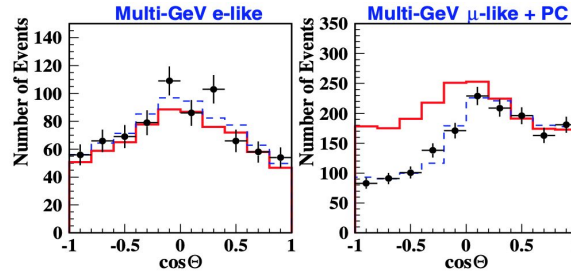
We know that there is **new physics** (see all previous talks)

evidence includes  
e.g.

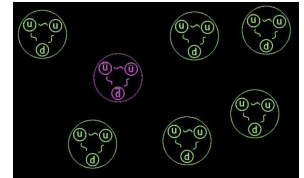
## Dark Matter



## Neutrino Masses



## Baryon asymmetry?



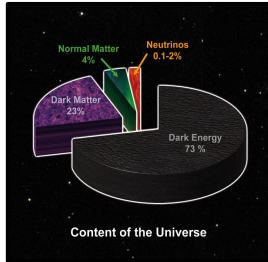
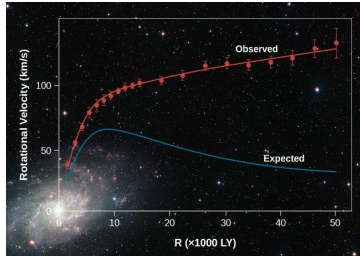
$$\eta_B = (n_B - n_{\bar{B}})/n_\gamma$$
$$\simeq (6.09 \pm 0.06) \times 10^{-10}$$

# Evidence for New Physics

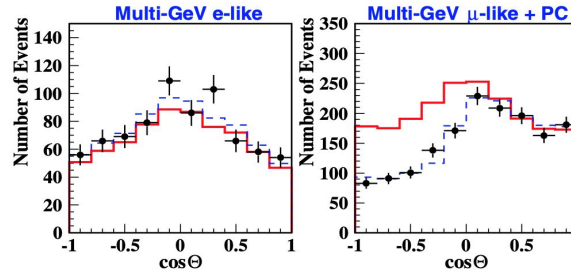
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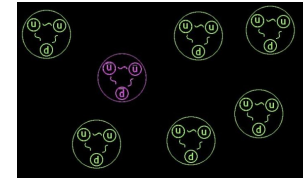
**Dark Matter**



**Neutrino Masses**



**Baryon  
asymmetry?**

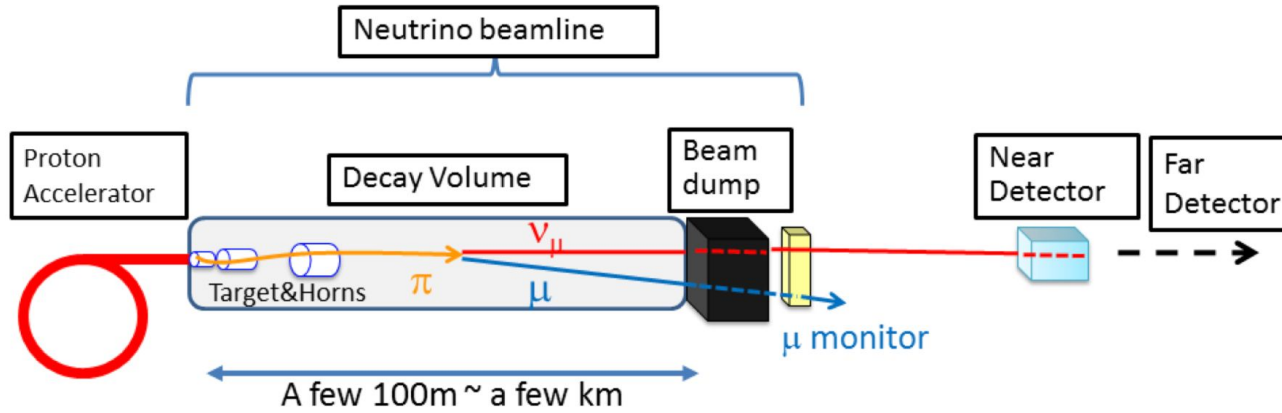


$$\eta_B = (n_B - n_{\bar{B}})/n_\gamma$$
$$\simeq (6.09 \pm 0.06) \times 10^{-10}$$

**How can neutrino experiments help us in finding the new physics?**

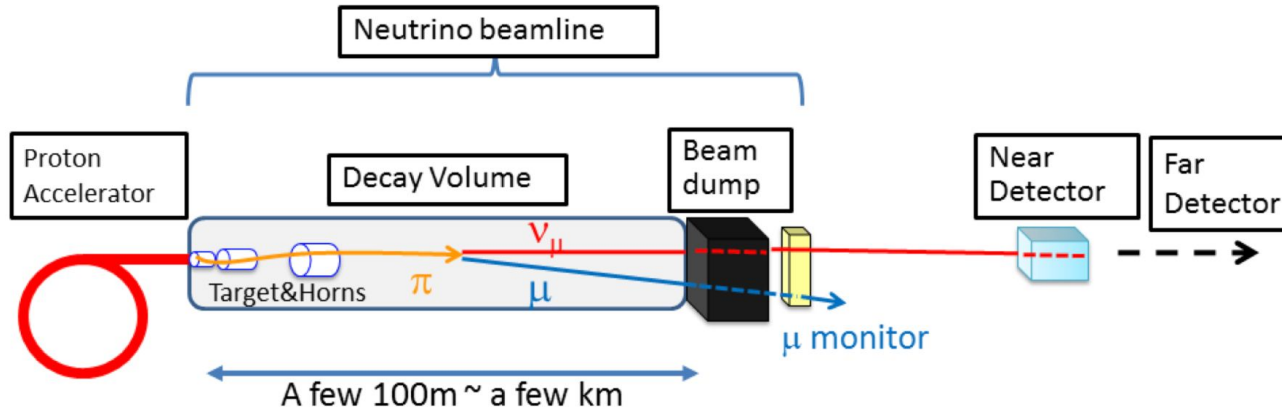
# Searching for New Physics

Typical setup for a pion decay-in-flight accelerator neutrino experiment



# Searching for New Physics

Typical setup for a pion decay-in-flight accelerator neutrino experiment

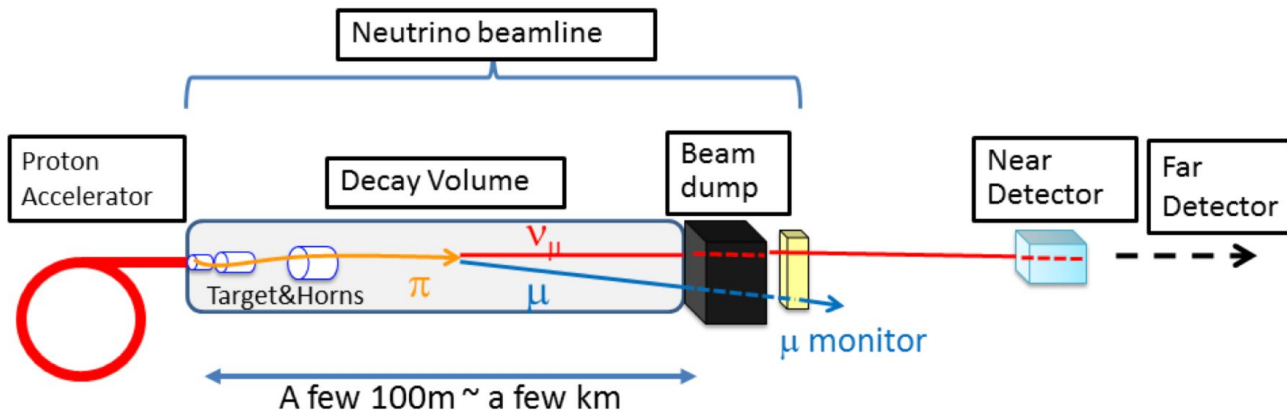


Possible search strategies

1. **Direct production of new particles and mediators**

# Searching for New Physics

Typical setup for a pion decay-in-flight accelerator neutrino experiment

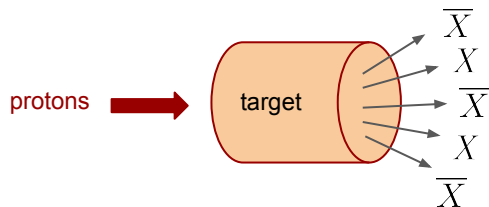


Possible search strategies

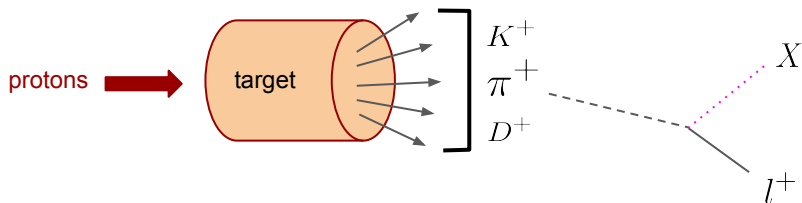
1. Direct production of new particles and mediators
2. Indirect effects on e.g. neutrino oscillation

# Searching for NP: Direct production

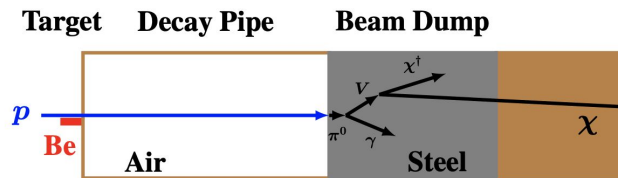
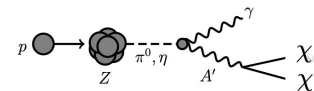
## 1. Drell-Yan/Bremmstrahlung in the proton target



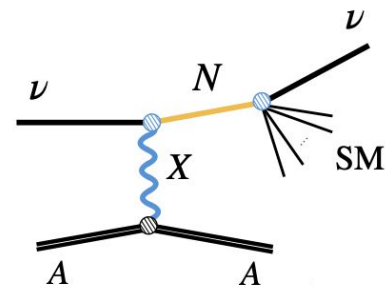
## 2. Meson decay-in-flight



## 3. Beam dump mode



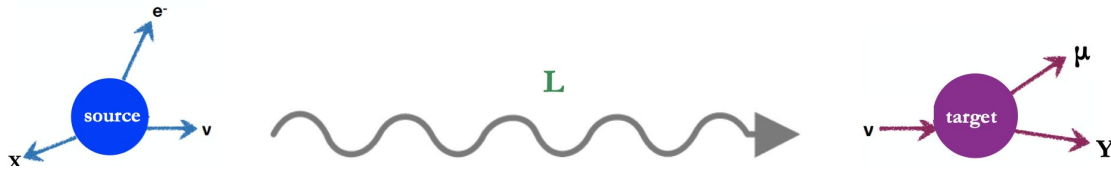
## 4. In-detector neutrino scattering



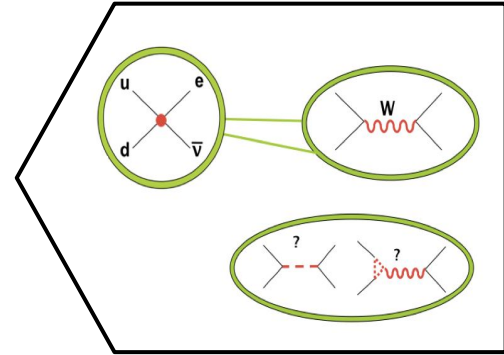
# Searching for NP: Indirect effects

Examples of indirect NP include e.g. NSIs, EFTs (see arXiv: 1901.04553, 1910.02971)

Heavy new physics can modify **neutrino production**, **propagation** and **detection** leading to modified oscillation probabilities



by Martín González-Alonso



Modifications can be computed in the presence of **EFT** and measurements of low-energy observables translated into constraints on higher-dimension SMEFT operators



# Why Neutrino Experiments?

1. High POT ( $\sim 1e21$ ) → **large flux of charged and neutral mesons**
2. Large detector masses  $\sim O(1e2)$  tonnes
3. Good PID ( $p, \mu, e, \gamma$  reconstructed)
4. Good calorimetry
5. High cosmic rejection
6. Parasitic: cost shared with neutrino projects

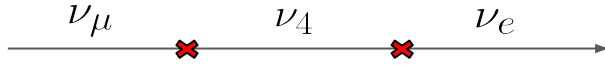
See talk by Richard Van de Water, U.S. Cosmic Visions  
2017

“Future possibilities at Proton fixed target experiments”



**New Physics Searches at  
*Current* and *Future*  
Neutrino Experiments**

# Neutrino Oscillations: Steriles



Among the simplest extensions of the SM that predict neutrino masses

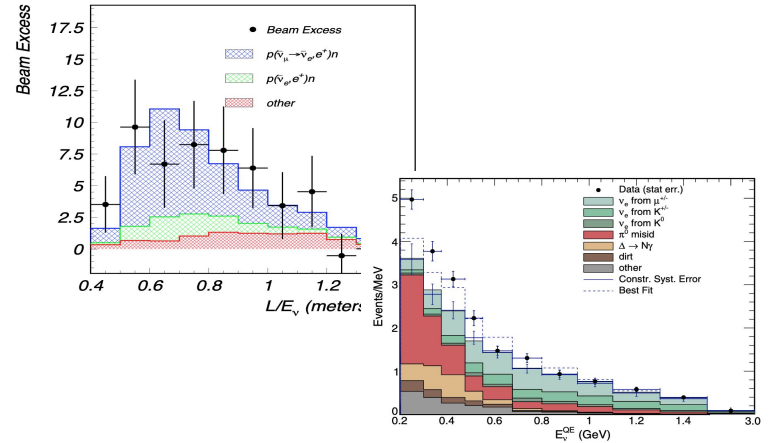
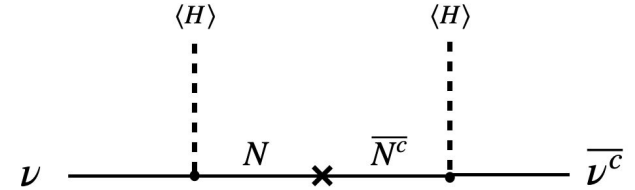
**Type-I seesaw:**

$$\mathcal{L}_{\nu\text{-mass}} \supset -Y_{\alpha i} \overline{L}_{\alpha} \widetilde{H} N_j - \frac{M_{ij}}{2} \overline{N}_i^c N_j + \text{h.c.}$$

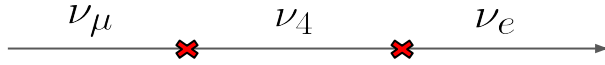
$$M_{\nu} \simeq -\frac{v^2}{2} Y^T M^{-1} Y$$

neutrino portal to a neutral hidden sector

Observation of what appears to be SBL oscillation has motivated the existence of a  $\sim$  eV scale sterile neutrino



# Neutrino Oscillations: Steriles

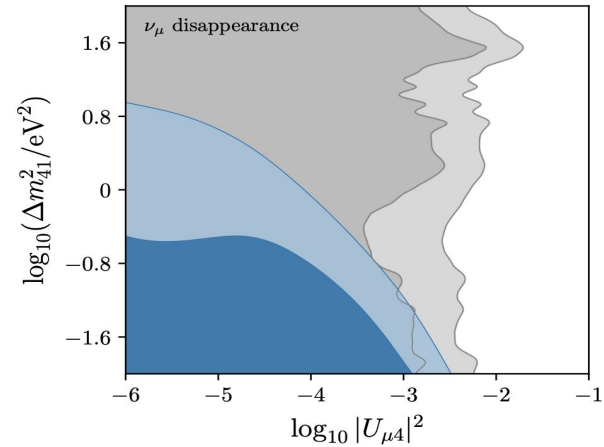


The 3+1 sterile solution to SBL anomalies is now disfavoured

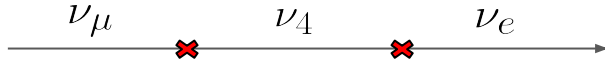
Light sterile neutrino is relativistic at early times and contributes to  $N_{\text{eff}}$

strong constraints from CMB+LSS

S. Hagstotz et al, arXiv:2003.02289



# Neutrino Oscillations: Steriles



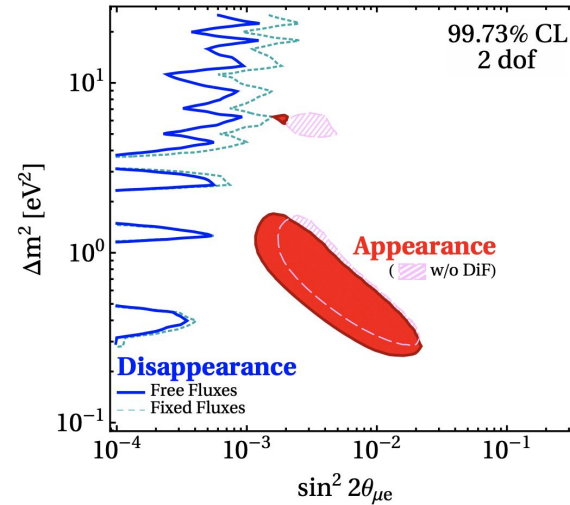
The 3+1 sterile solution to SBL anomalies is now disfavoured

Furthermore, strong tensions between appearance and disappearance data

$$P_{\mu\mu}^{\text{SBL}} = 1 - 4|U_{\mu 4}|^2(1 - |U_{\mu 4}|^2) \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

$$P_{\mu e}^{\text{SBL}} = 4|U_{\mu 4}|^2|U_{e 4}|^2 \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

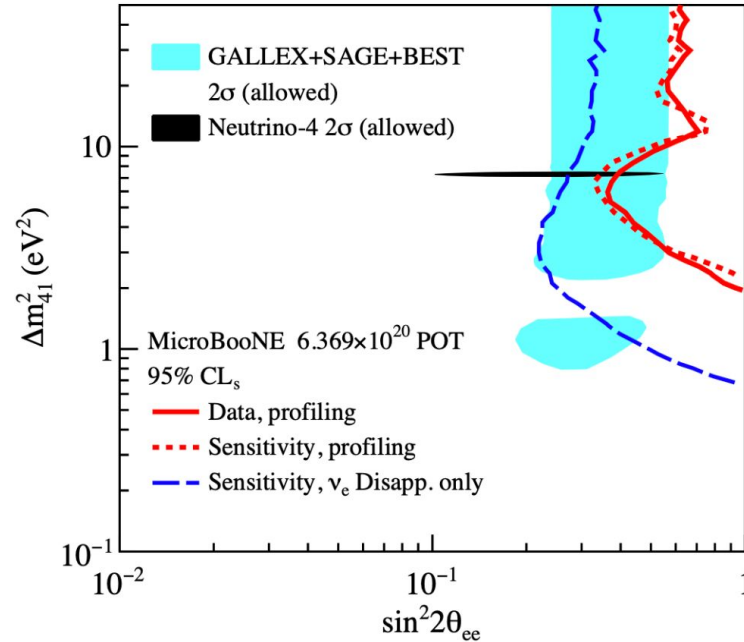
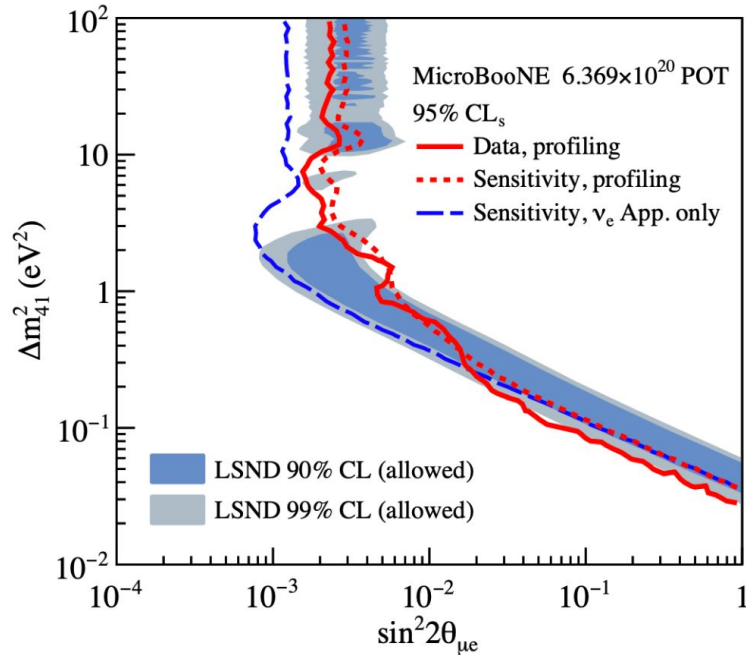
M. Dentler et al, JHEP 08 (2018) 010



# Neutrino Oscillations: Steriles

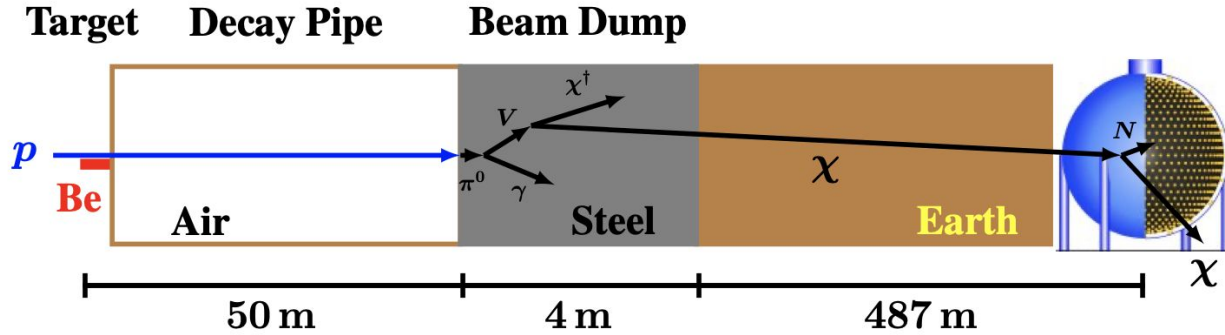
Most recent sterile search comes from MicroBooNE which does not exclude the full LSND (and MiniBooNE) preferred regions

P. Abratenko et al, arXiv:2210.10216

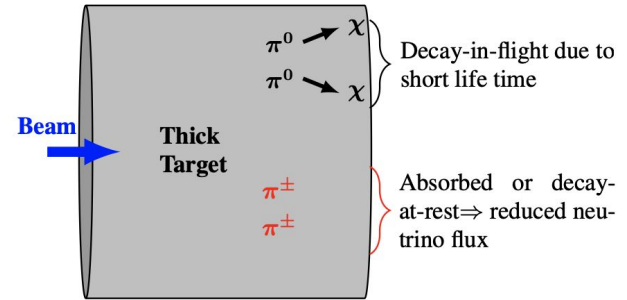


# Beam Dump Mode

BSM searches at neutrino experiments suffer large neutrino bkg and the NP can be hard to detangle



If protons impinge directly on the beam dump target, charged pions are absorbed or decay at rest  $\rightarrow$  neutrino flux suppressed!



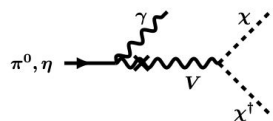
MiniBooNE-DM collab. arXiv: 1807.06137

# Beam Dump: Light Dark Matter

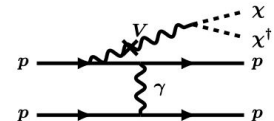
MiniBooNE was the first accelerator neutrino experiment to perform a dark matter search in beam dump mode

## Vector portal LDM

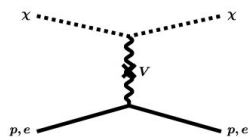
$$\mathcal{L}_V = \mathcal{L}_\chi - \frac{1}{4} V_{\mu\nu} V^{\mu\nu} + \frac{1}{2} m_V^2 V_\mu V^\mu - \frac{\epsilon}{2} F_{\mu\nu} V^{\mu\nu}$$



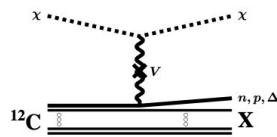
(a) Meson Decay



(b) Proton Bremsstrahlung + Vector-Mixing

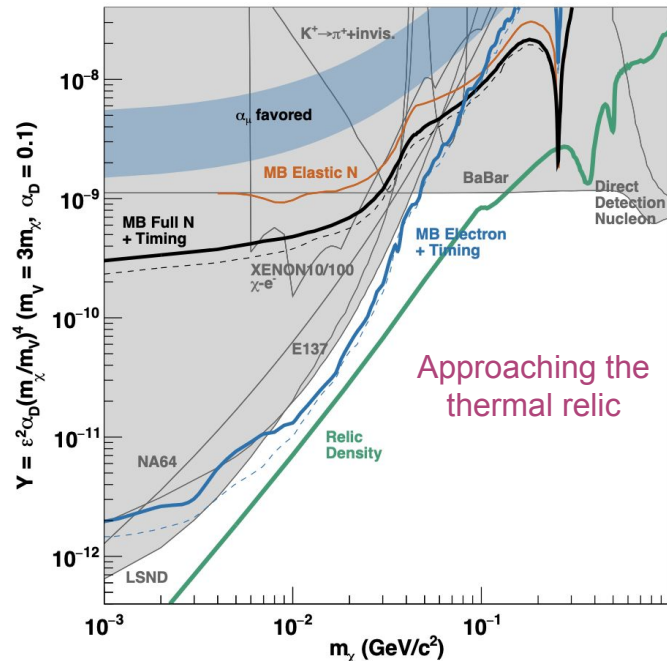


(a) Free Protons or Electrons



(b) Bound Nucleons

MiniBooNE-DM collab. arXiv: 1807.06137





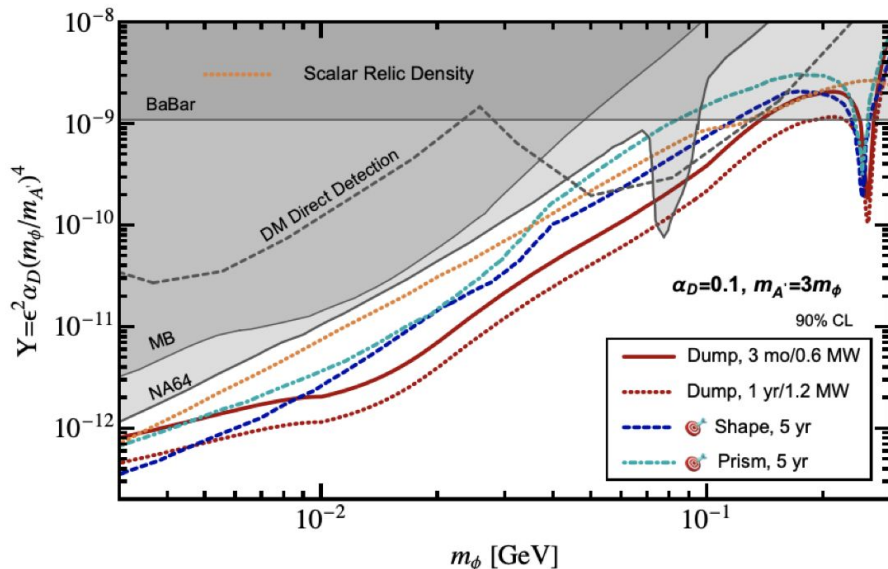
# LDM at a “target-less” DUNE

Recently been proposed that running in beam dump mode at DUNE for < 3 months at half the expected initial beam intensity could give competitive limits on LDM!

Light scalar dark matter coupled to a dark photon

$$\mathcal{L}_{\text{DM}} \supset -\frac{1}{2}\epsilon F_{\mu\nu}F'^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu + |D_\mu\phi|^2 - m_\phi^2|\phi|^2$$

see Brdar et al, arXiv: 2206.06380  
BSM Targets at a “Target-less DUNE”

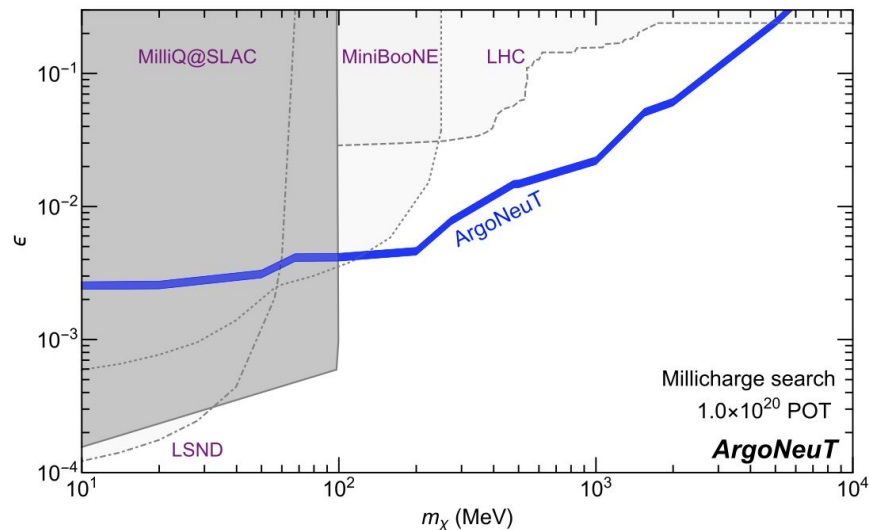
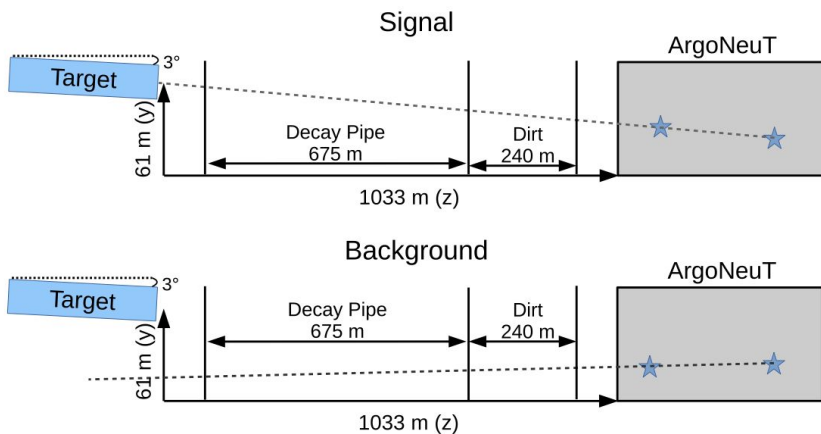


# BSM searches at ArgoNeuT

## Search for Millicharged particles (MCP)

Flux of MCPs from neutral meson decays or Drell-Yan at target scatter off electrons

**Signal:** two soft energy depositions aligned with target direction



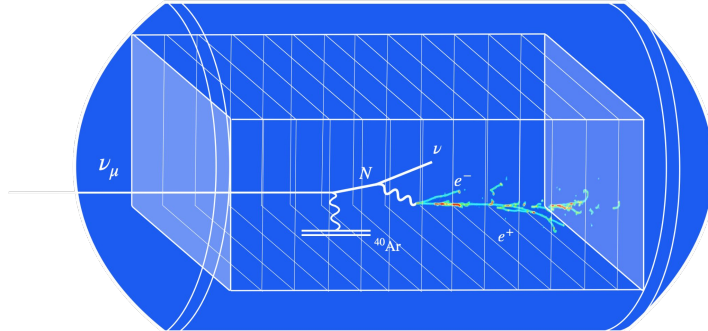
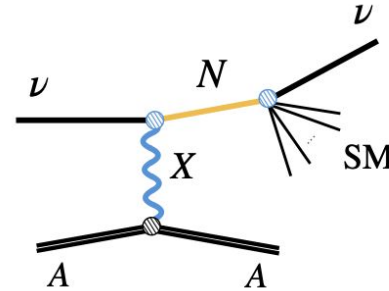
PHYSICAL REVIEW LETTERS 124, 131801 (2020)

# New Particles in Neutrino Scattering

New particles and mediators may be produced in the scattering of neutrinos on the detector material

Distinguishing these processes from SM interactions requires excellent PID and reconstruction → **LArTPC detectors**

e.g. ArgoNeuT, MicroBooNE (running)  
+ ICARUS, SBND, DUNE (upcoming)



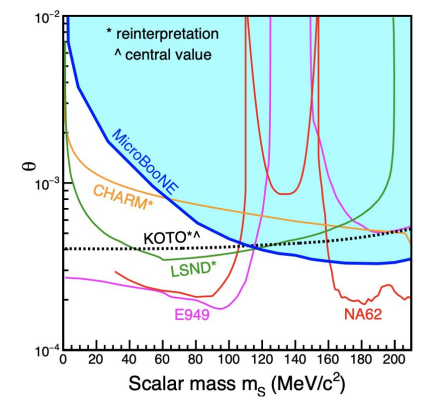
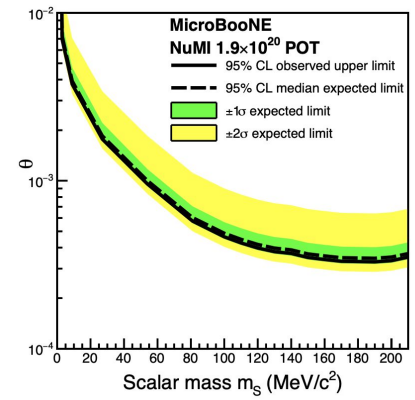
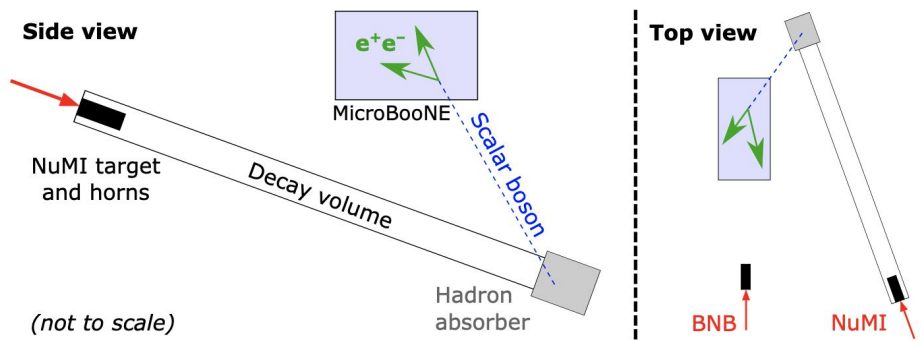
Schematic of MicroBooNE TPC and cryotank

1. Improved hadronic vertex ID
2. Better spatial resolution
3. Improved calorimetry

# BSM searches at MicroBooNE

Higgs portal scalar decaying to  $e^+e^-$  pair (first search for an  $e^+e^-$  final state)

Scalars produced in KDAR at the NuMI\* hadron absorber travel to MicroBooNE detector and produce  $e^+e^-$  pairs



Leading constraint between  
130 - 160 MeV

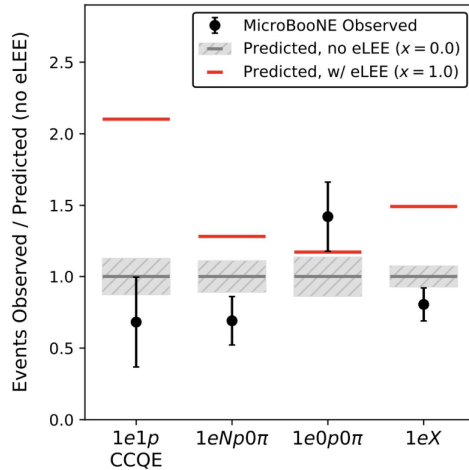
\*MicroBooNE is located 8 deg. off-axis with respect to the NuMI target

# MicroBooNE LEE search

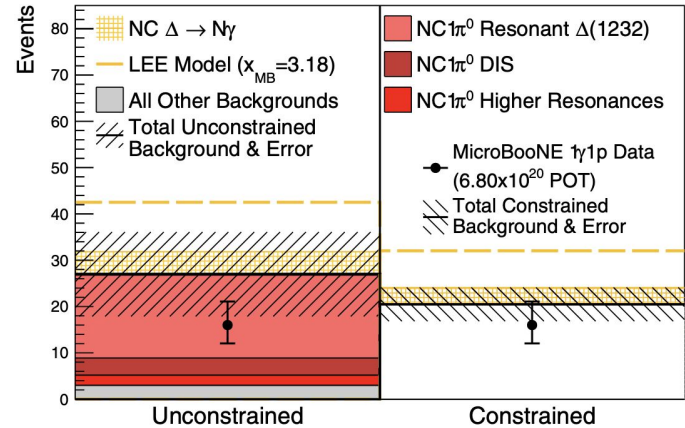
Among MicroBooNE's main physics goals is to determine the origin of the MiniBooNE LEE

Searches for an eLEE and gLEE have already been performed

## Search for excess $e^-$ (PRL 128, 241801 2022)



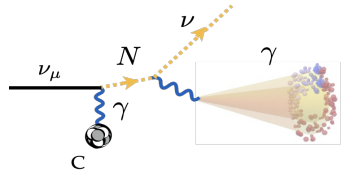
## Search for SM NC $\Delta \rightarrow N\gamma$ (PRL 128, 111801 2022)



**No electron or single photon excess observed compatible with MiniBooNE!**

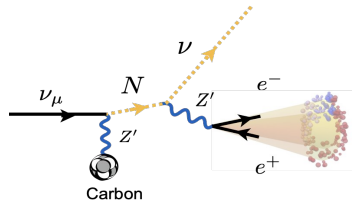
# MicroBooNE LEE search

A number of new physics models which predict EM final states ( $e^-$ ,  $\gamma$ ,  $e^+e^-$ ) have been proposed as alternatives to a SM LEE, e.g.



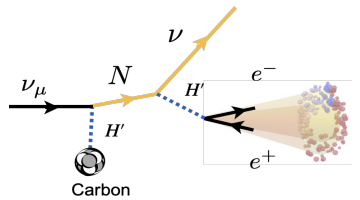
transition magnetic moment

$$F^{\mu\nu} \left( \frac{\mu_\nu^{aj}}{2} \bar{\nu}_\alpha \sigma_{\mu\nu} N_j + \frac{\mu_\nu^{ij}}{2} \bar{N}_i \sigma_{\mu\nu} N_j \right)$$



dark photon

$$Z'_\mu \left( V^{\alpha j} \bar{\nu}_\alpha \gamma^\mu N_j + V^{ij} \bar{N}_i \gamma^\mu N_j + d_V^\ell \bar{\ell} \gamma^\mu \ell \right)$$



dark scalar

$$h' \left( S^{\alpha j} \bar{\nu}_\alpha N_j + S^{ij} \bar{N}_i N_j + d_S^\ell \bar{\ell} \ell \right)$$

While MicroBooNE is yet to perform dedicated searches for these models, we may already be able to constrain them using MicroBooNE public data

**\*\* work-in-progress \*\***



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Columbia Uni.



Daniele Massaro  
Uni. of Bologna



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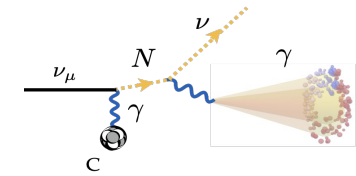
Mark Ross-Lonergan  
LANL



Guanqun Ge  
Columbia Uni.

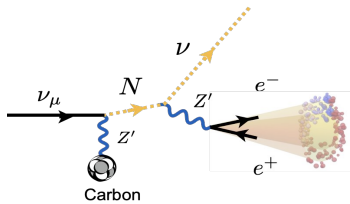
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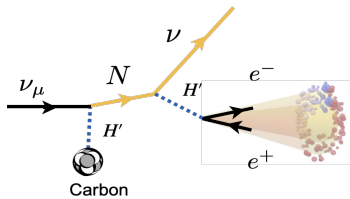
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dark scalar

$$h' \left( S^{\alpha j} \bar{\nu}_\alpha N_j + S^{ij} \bar{N}_i N_j + d_S^\ell \bar{\ell} \ell \right)$$

Simulation of neutrino-nuclear scattering at MicroBooNE using



`pip install DarkNews`

**DarkNews** is a light-weight Python generator for neutrino-nucleus upscattering to heavy neutrinos.

**Public and documented**

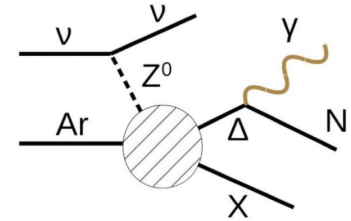
Paper: [arxiv.org/abs/2207.04137](https://arxiv.org/abs/2207.04137)

GitHub: [github.com/mhostert/DarkNews-generator](https://github.com/mhostert/DarkNews-generator)

PyPI: [pypi.org/project/DarkNews/](https://pypi.org/project/DarkNews/)

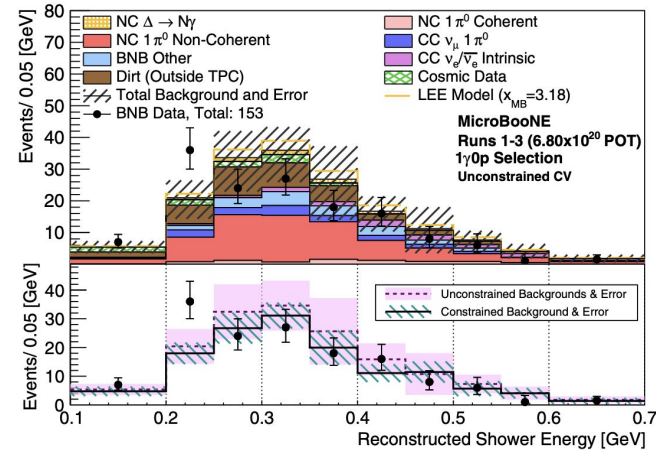
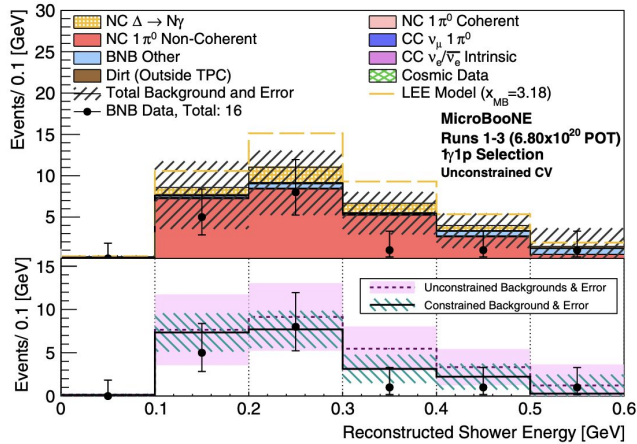
# gLEE: NC $\Delta \rightarrow N\gamma$ search

MicroBooNE constrain two topologies in the gLEE ( $\gamma$ LEE) search:



1.  $\Delta^+ \rightarrow p \gamma \rightarrow 1\gamma 1p$

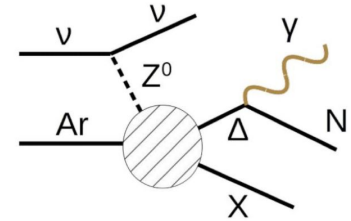
2.  $\Delta^0 \rightarrow n \gamma \rightarrow 1\gamma 0p$



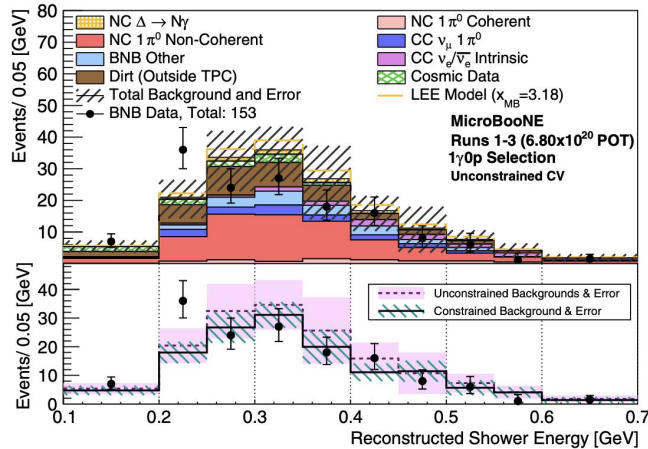


# gLEE: $NC \Delta \rightarrow N\gamma$ search

MicroBooNE constrain two topologies in the gLEE ( $\gamma$ LEE) search:



## 2. $\Delta^0 \rightarrow n \gamma \rightarrow 1\gamma 0p$

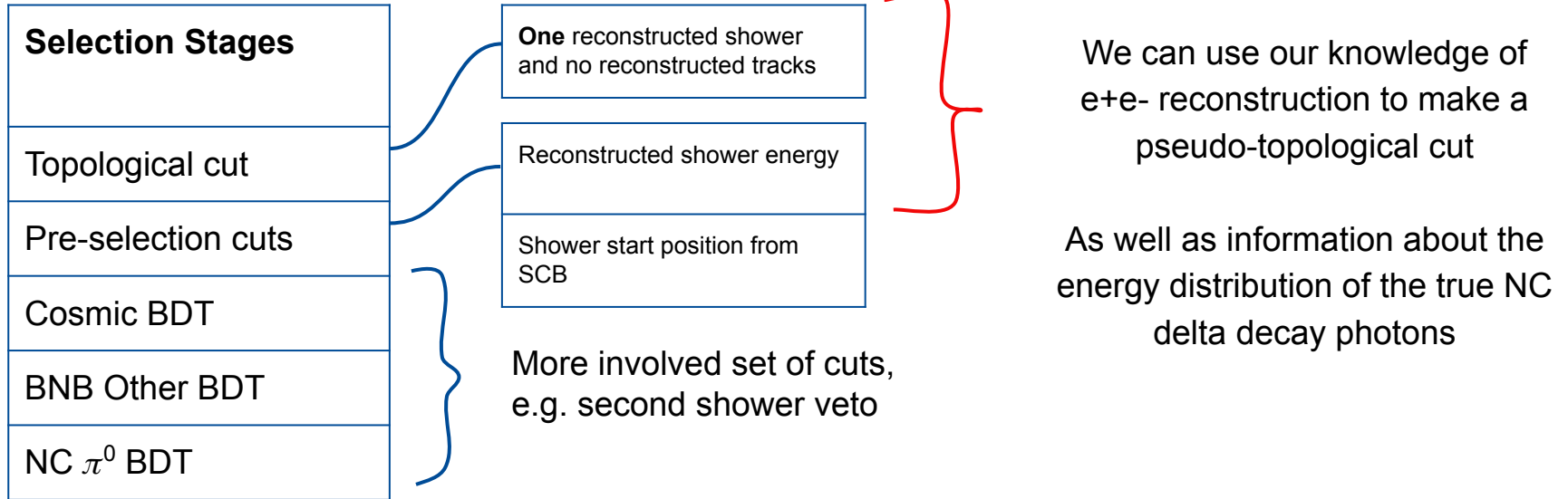


While our models do contribute to both topologies, we **consider only the  $1\gamma 0p$  topology** to avoid issues with nuclear modelling

↖ caveat:  $1\gamma 0p$  selection has larger bkg, as proton kinematics cannot be leveraged for bkg rejection

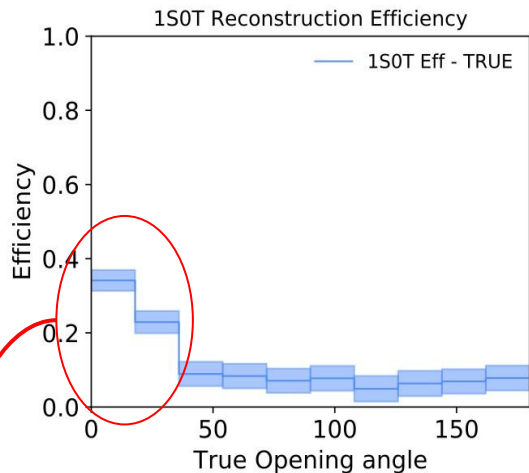
# gLEE 1y 0p selection

Impossible to apply official MicroBooNE selection and so we make some educated guesses

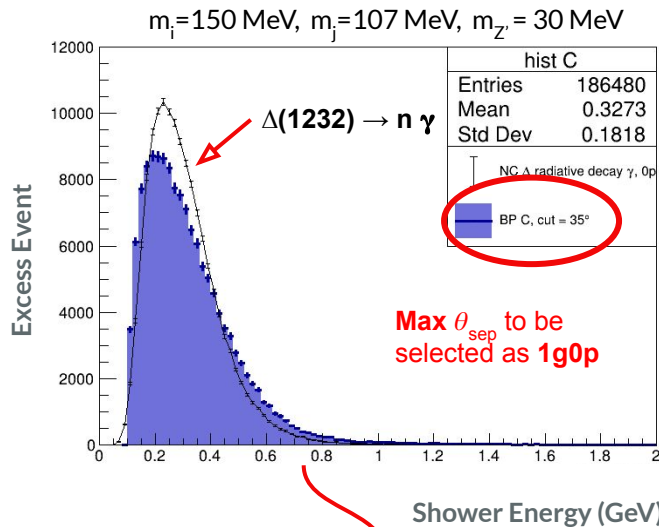


# pseudo-1 $\gamma$ 0p selection

Pairs of uniformly distributed  $e^+e^-$  are passed through MicroBooNE reconstruction and topological reconstruction efficiencies derived

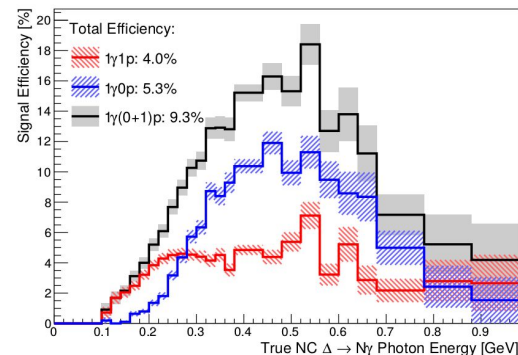


Efficiency to reconstruct a single shower highest for  $e^+e^-$  separation angle **< 35 deg.**



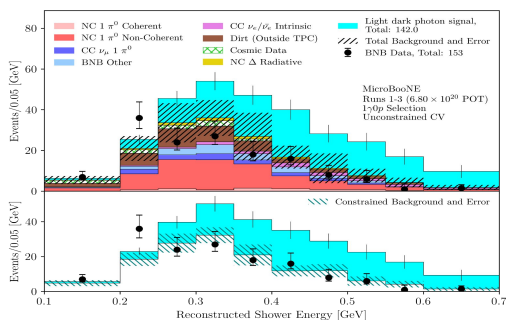
Comparison of truth-level energy distributions for benchmark model suggest it is valid to apply single photon selection efficiency

We can then apply 1g0p efficiency to our remaining events

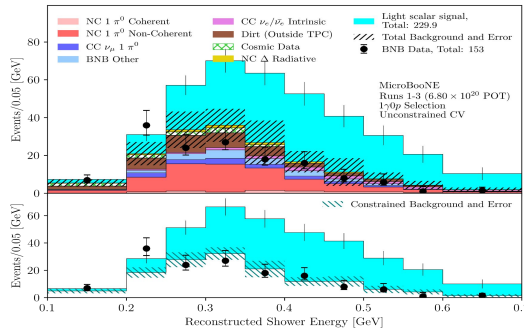


# Preliminary selection

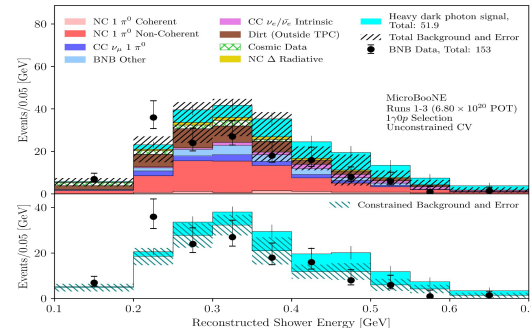
## Light dark photon BP



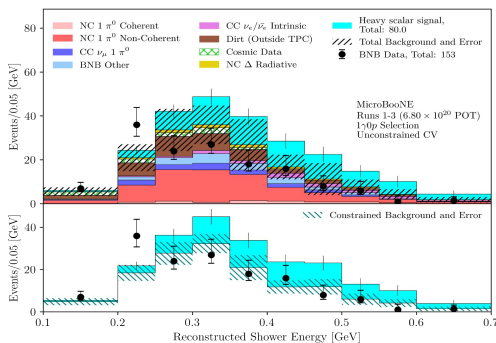
## Light scalar BP



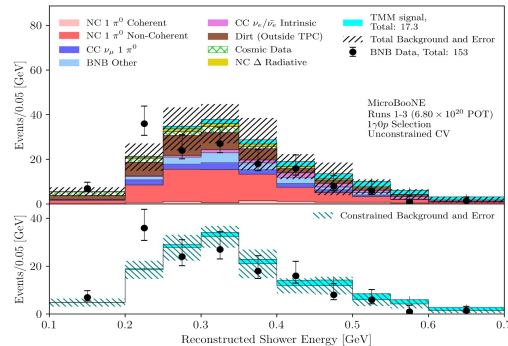
## Heavy dark photon BP



## Heavy scalar BP




## TMM BP



# Takeaways: Why Neutrino Experiments?

1. High POT  $O(1e21)$  → **large flux of charged and neutral mesons** that can decay to new particles which leave signal in the detector
2. Can take advantage of **potential to run in beam dump mode** to search for broader dark sector models
3. Strong **PID** ( $p$ ,  $\mu$ ,  $e$ ,  $\gamma$ ), **spatial** and **energy resolution** (esp. LArTPC-based detectors) make distinguishing NP from SM more possible
4. **Parasitic**: Cost burden is shared with on-going neutrino projects and greater goal of elucidating oscillation parameters



**Thank you for your  
attention!**