

High-resolution CMB Temperature Anisotropies as a Particle Detector: Are We Detecting (Sterile) Neutrinos?

Lloyd Knox (UC Davis)

Zhen Hou & Marius Millea (UCD)

Ryan Keisler (UC), Christian
Reichardt (UCB)

SPT collaboration

Planck collaboration

Review of This Session

- Neal Weiner: Direct Detection

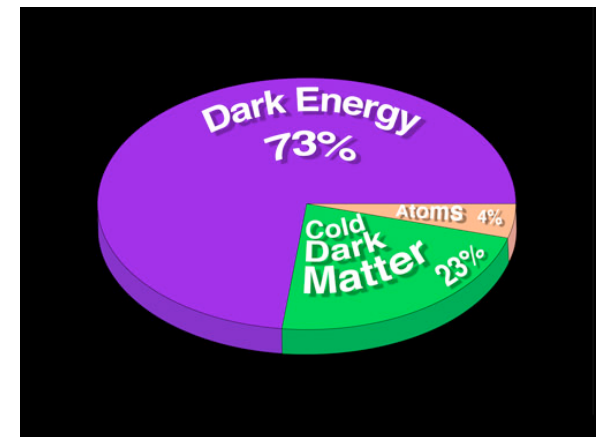
Review of This Session

- Neal Weiner: Direct Detection
- Nicole Bell: Indirect Detection

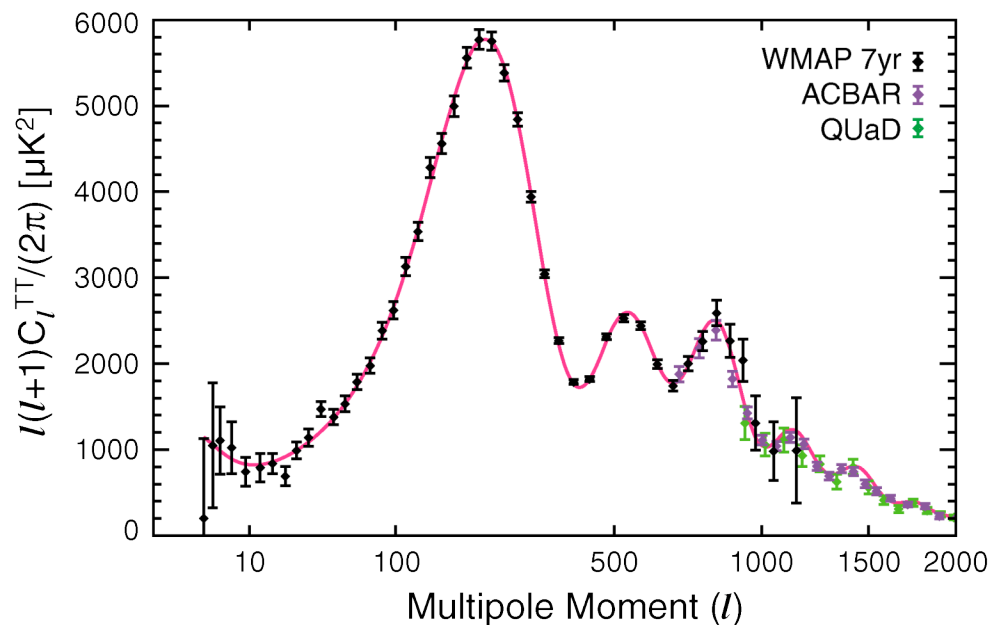
Review of This Session

- Neal Weiner: Direct Detection
- Nicole Bell: Indirect Detection
- Lloyd Knox: Really Indirect Detection
(of dark stuff)

CMB is a well-calibrated detector



Contents of Universe*



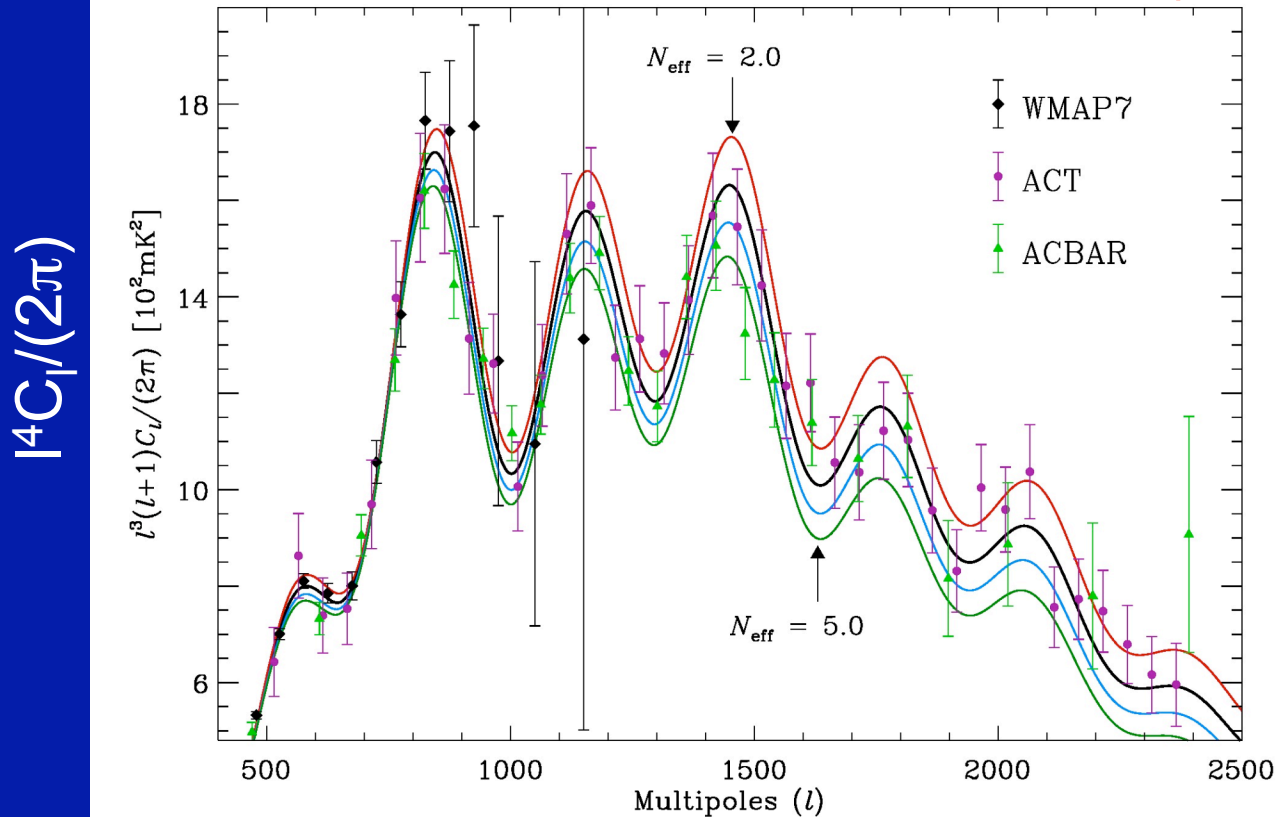
Theory-observation agreement quite striking.

Confirmation and increased precision, but no surprises.

*contains trace amounts of nuts

Surprises may come at higher resolution

Hou, Keisler, LK, Millea & Reichardt (2011)



Extra neutrino species?

98.4%
confidence
that $N_{\text{eff}} >$
standard
model value
(Hou et al.
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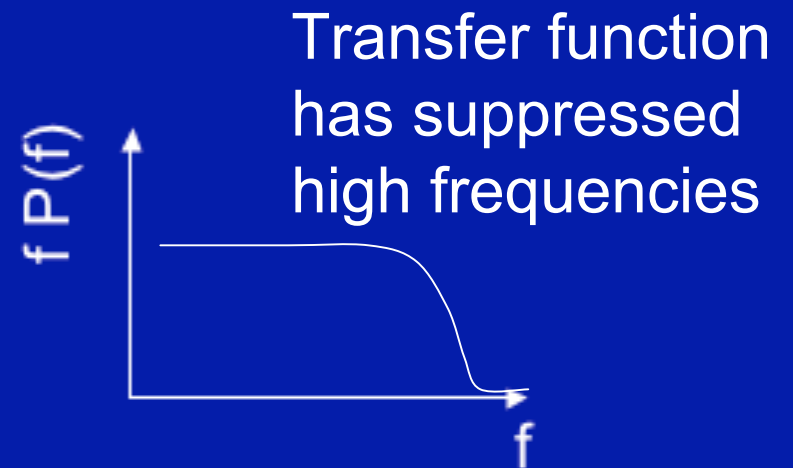
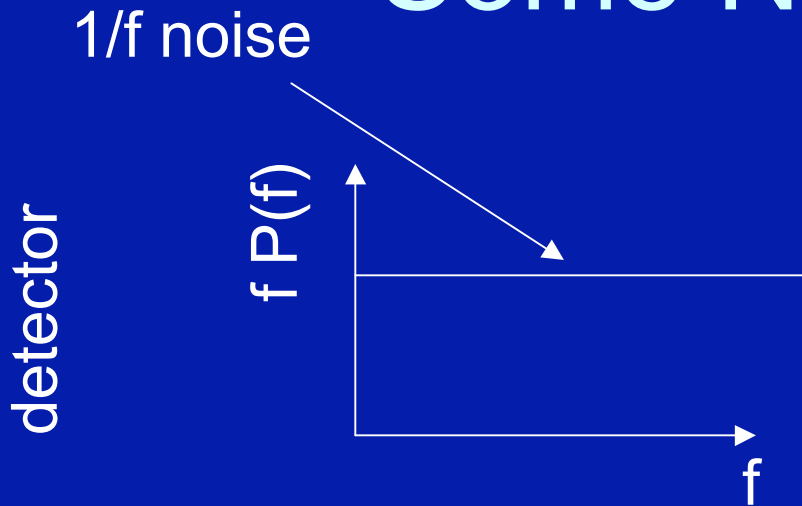
New results soon from SPT and ACT shrinking errors by ~ 3 .

Planck in \sim Jan 2013 will shrink errors by a further factor of ~ 5 .

Outline

- CMB as a detector of a stochastic process
- The stochastic process (inflation)
- Review of BBN and neutrinos
- Impact of extra neutrinos on the detector's (CMB's) transfer function
- Relation to other evidence for extra (sterile) neutrinos
- Discriminating changes to the stochastic process, from changes to the transfer function

The CMB is Like a Detector of Some Noise Source

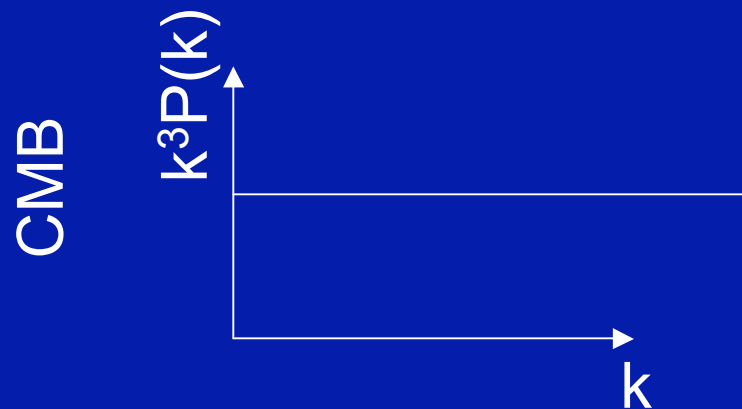
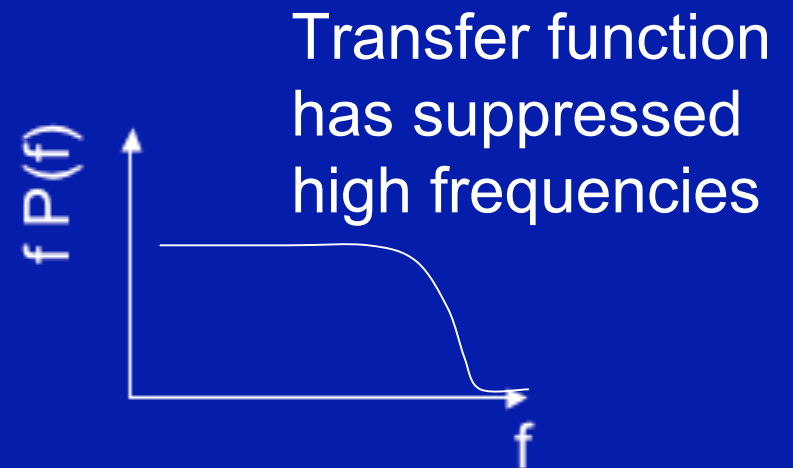
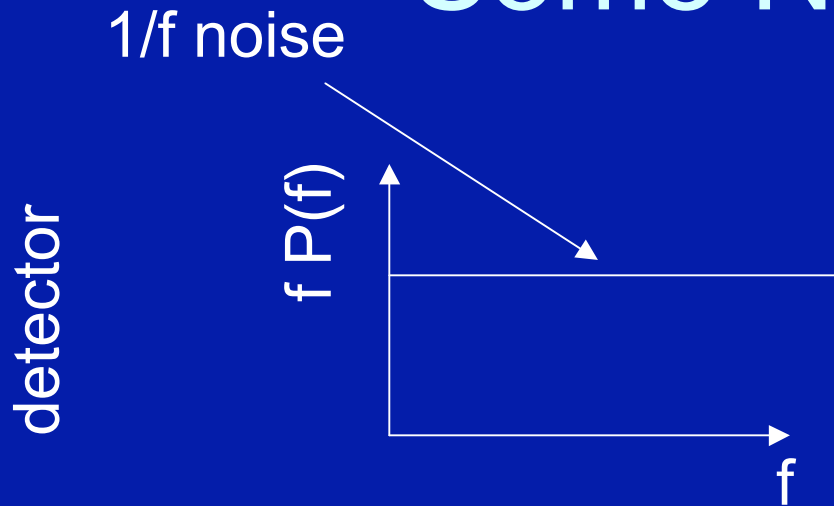


CMB

Noise source (input)

Detector output

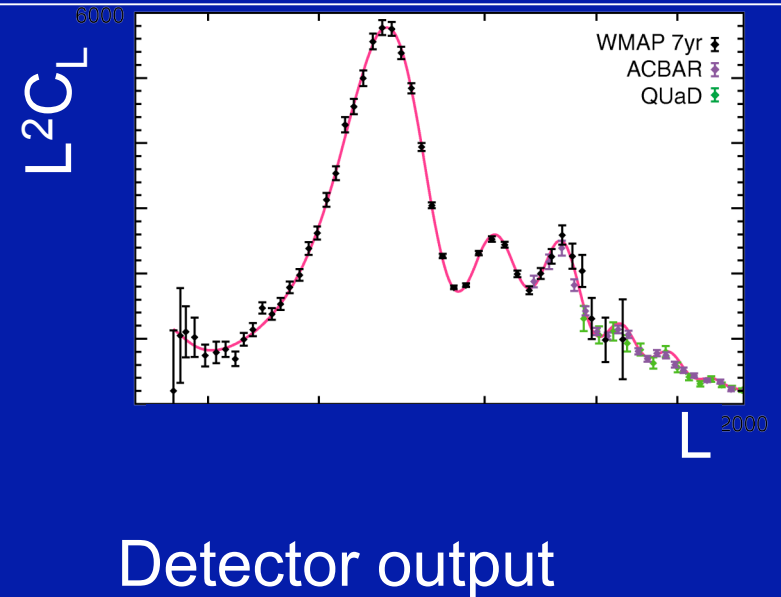
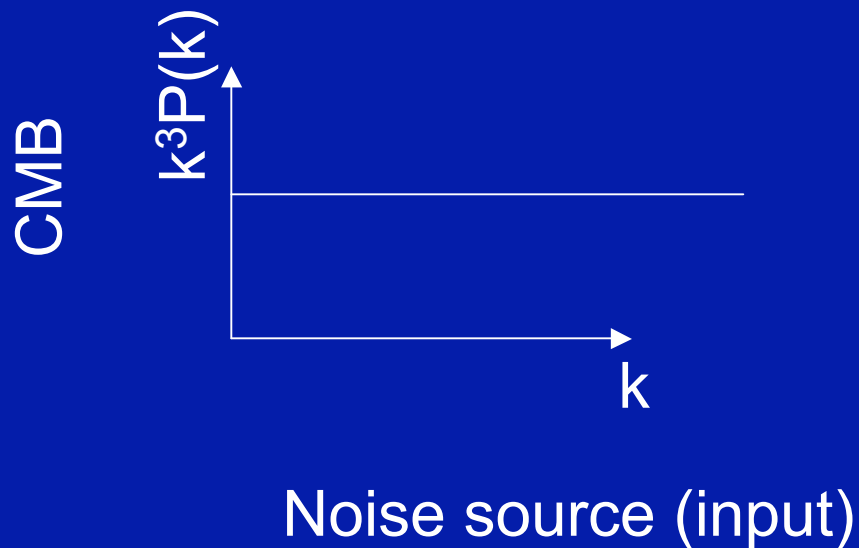
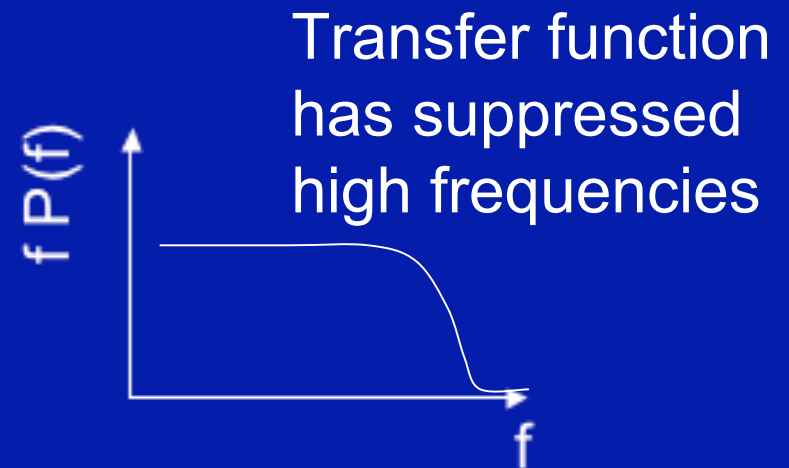
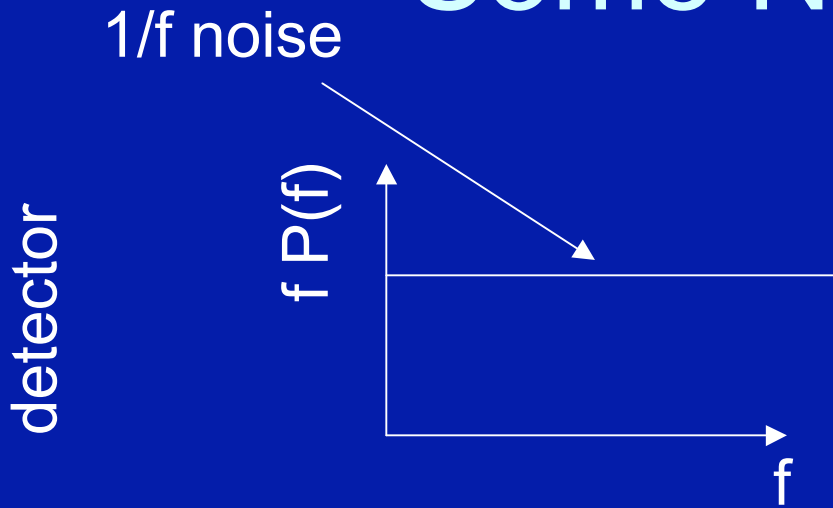
The CMB is Like a Detector of Some Noise Source



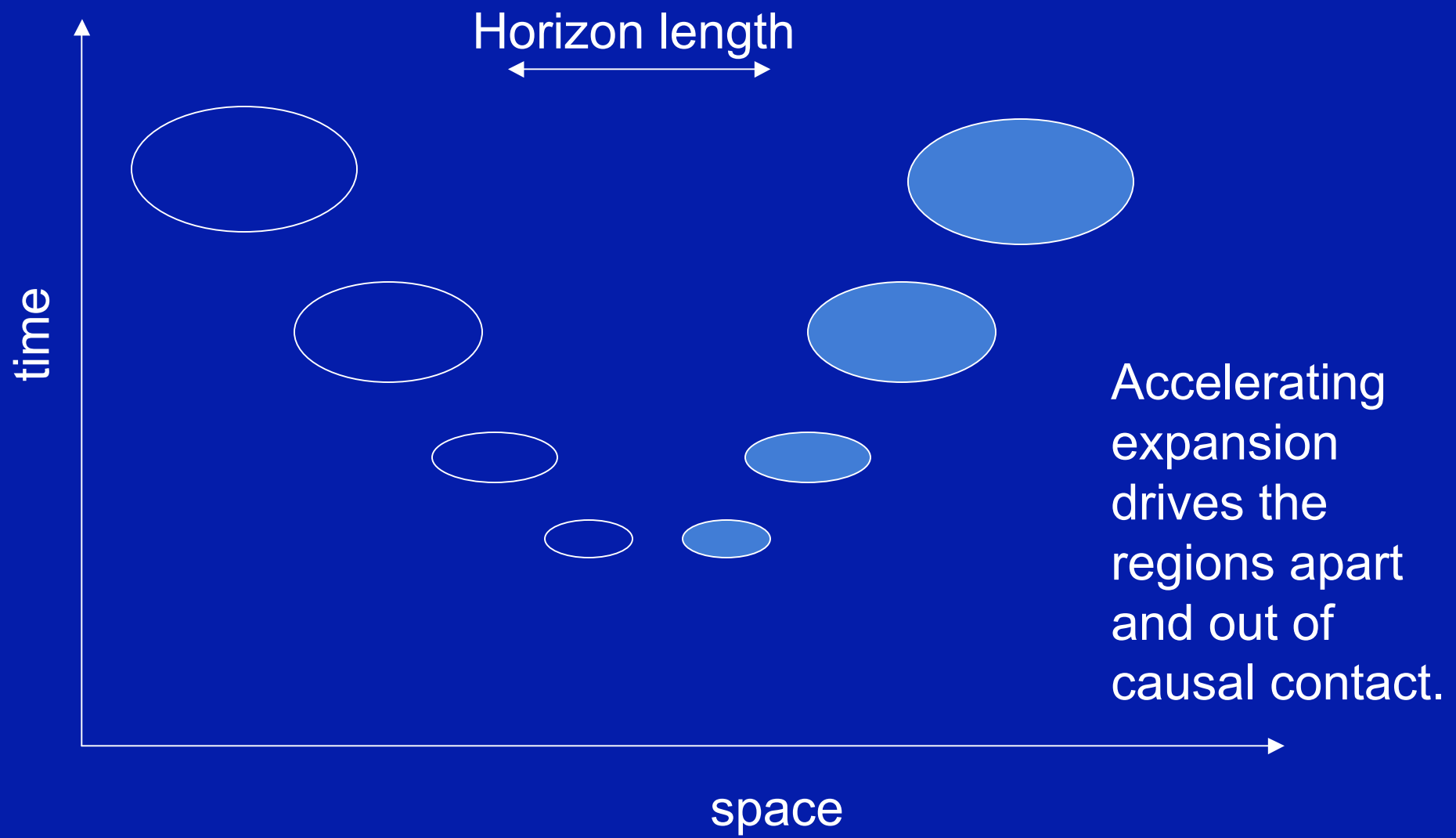
Noise source (input)

Detector output

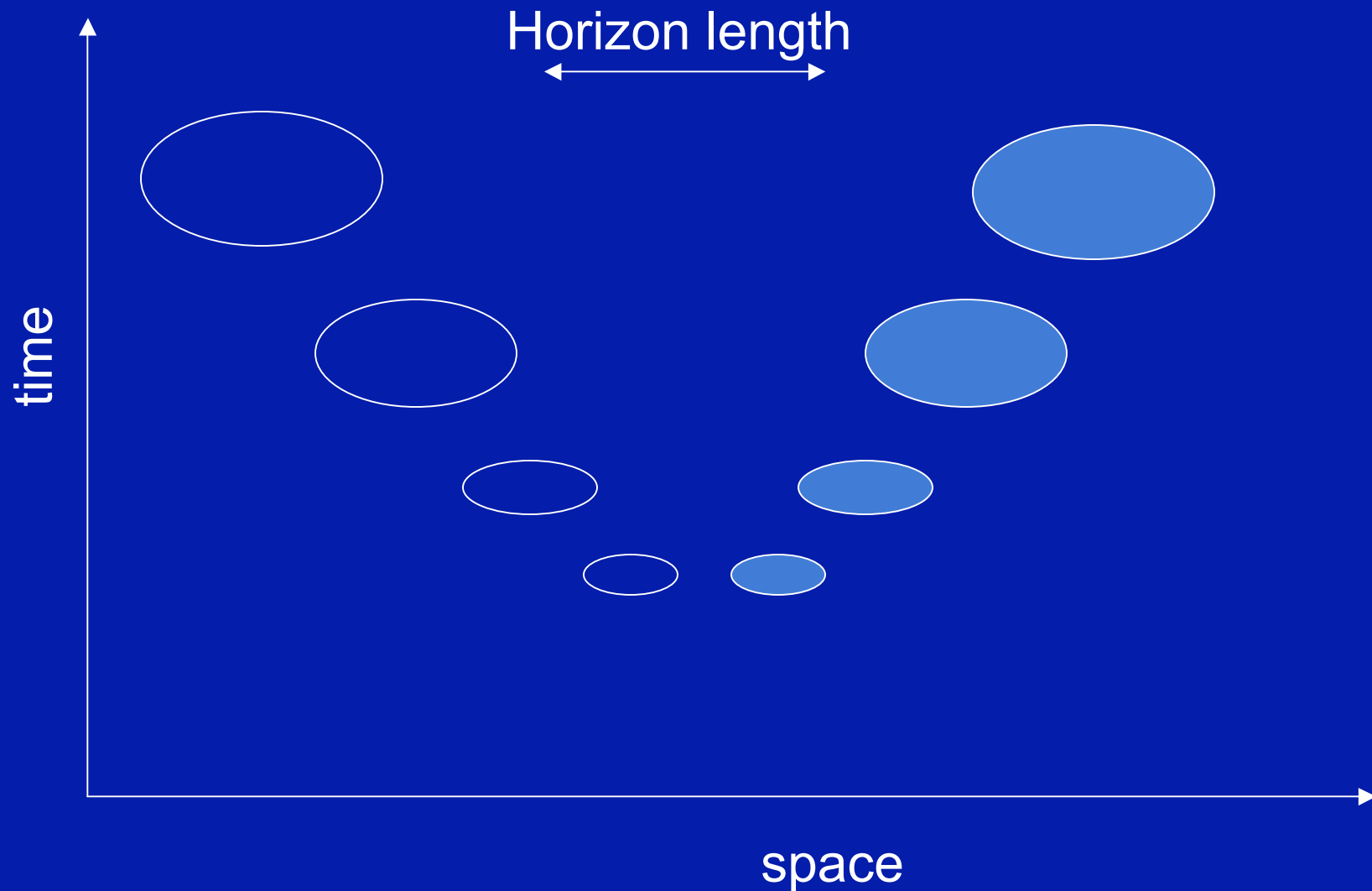
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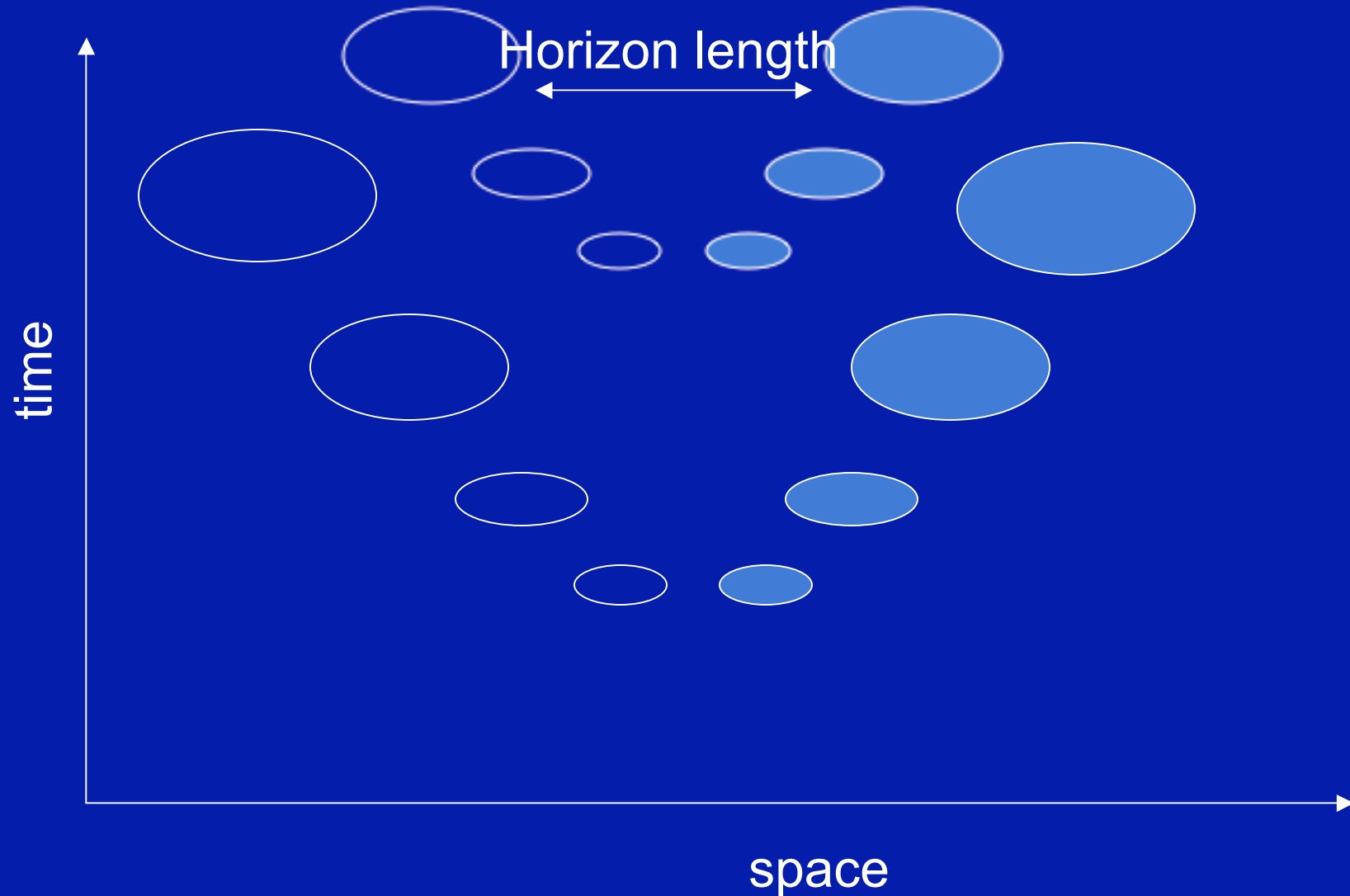
Expansion prevents quantum fluctuation from becoming undone



Smaller-scale perturbations are made later



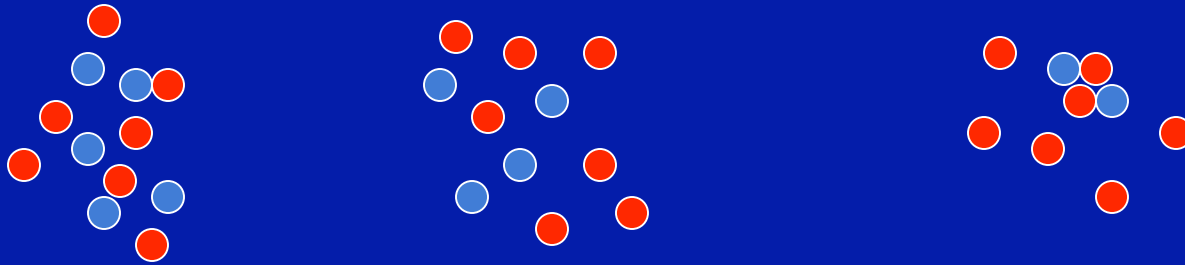
Smaller-scale perturbations are made later



$$H^2 = 8\pi G\rho/3$$

● p
● n

BBN and Y^*_p Review



↑
 $T_{p,n \text{ freeze-out}}$

↑
 T_{nuc}

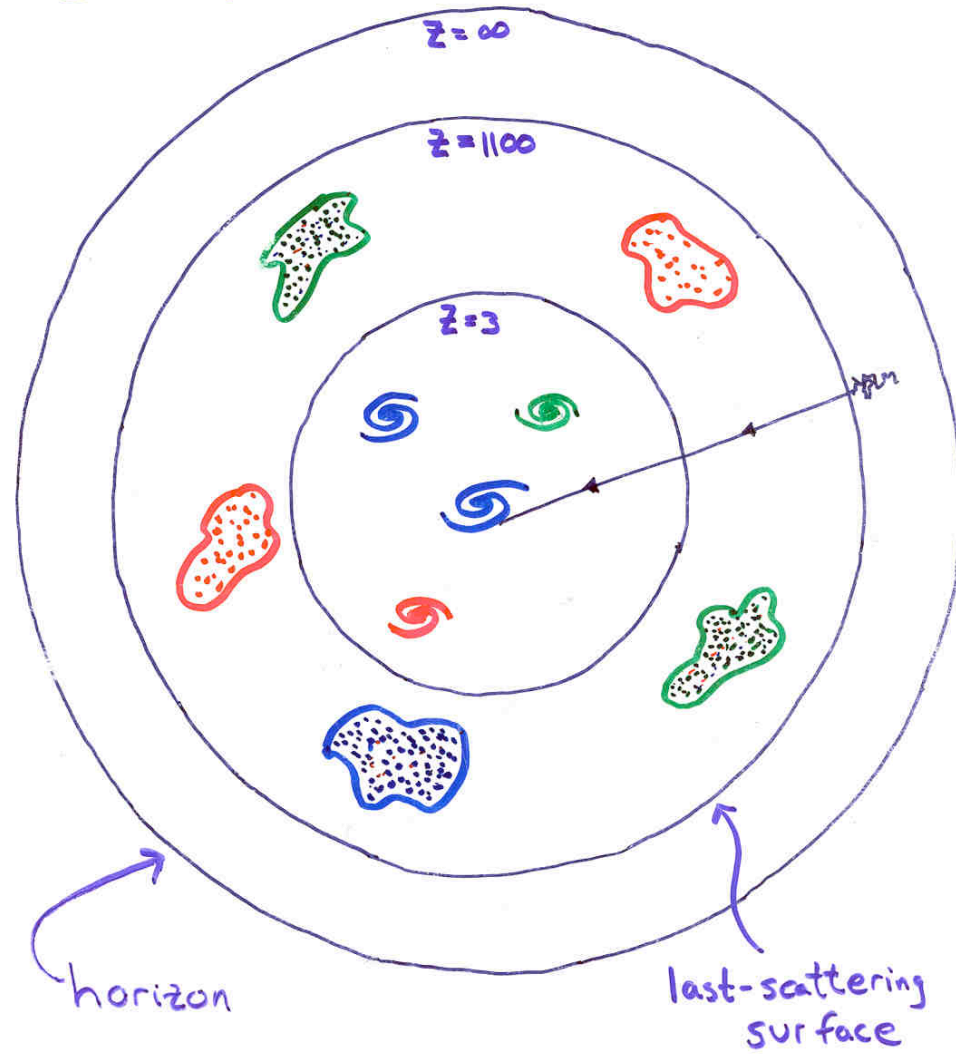
Cooler ==>

More ν species ==> higher H (at given T) ==> less time to cool to T_{nuc} ==> fewer decays ==> more Helium (also freeze-out n/p ratio increases with H)

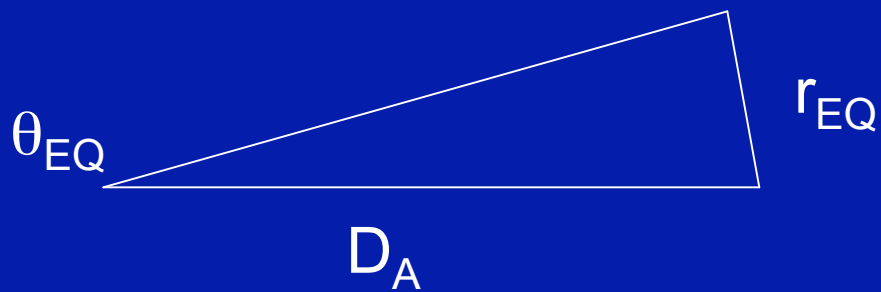
*Fraction of baryonic mass in Helium

THE HISTORY OF A SINGLE PHOTON

(12)



Three Scales in the CMB Transfer Function

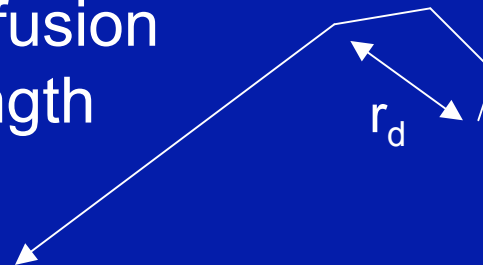


r_{EQ} included for completeness

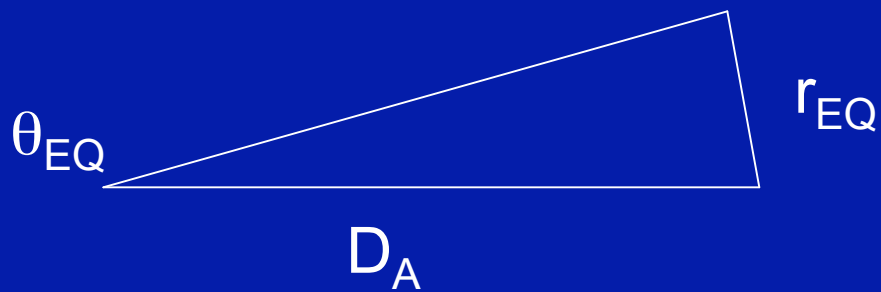
sound horizon: distance sound could travel by the time of last scattering. θ_s controls peak locations.



diffusion length



Three Scales in the CMB Transfer Function

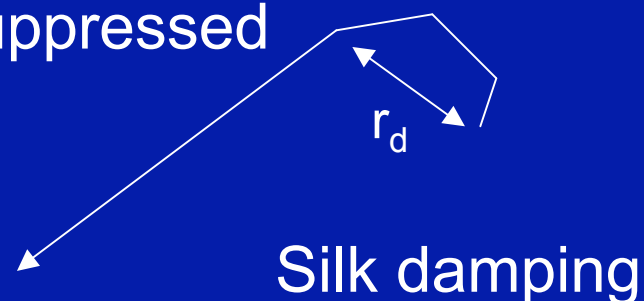


r_{EQ} included for completeness

sound horizon: distance sound could travel by the time of last scattering. θ_s controls peak locations.



Modes with $\lambda < r_d$ are suppressed



Three Scales in the CMB Transfer Function

$$100 \theta_s = 1.04 \pm 0.0016$$

sound horizon: distance sound could travel by the time of last scattering. θ_s controls peak locations.



$$r_s \propto 1/H$$

Extra $\nu \implies$ higher $\rho \implies$ higher $H \implies$ takes less time to cool to $T_{\text{rec}} \implies r_s$ is smaller

$$H^2 = 8\pi G\rho/3$$

Effect of extra ν on r_s

$$100 \theta_s = 1.04 \pm 0.0016$$

sound horizon: distance sound could travel by the time of last scattering. θ_s controls peak locations.



$r_s = \int_0^{a^*} c_s da / (a^2 H)$ Extra $\nu \implies$ higher $\rho \implies$ higher $H \implies$ takes less time to cool to $T_{\text{rec}} \implies r_s$ is smaller

$$H^2 = 8\pi G\rho/3$$

If we knew D_A we could find $r_s = \theta_s D_A$ and determine H

Effect of extra ν on r_d

Random-walk so goes as sq. root of time $\implies r_d \sim 1/H^{0.5}$

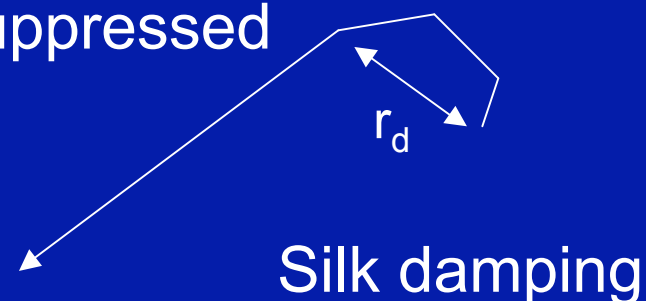
(Remember $r_s \sim 1/H$)

$$\theta_d/\theta_s = r_d/r_s \sim H^{0.5}$$

Dependence on D_A has
dropped out!

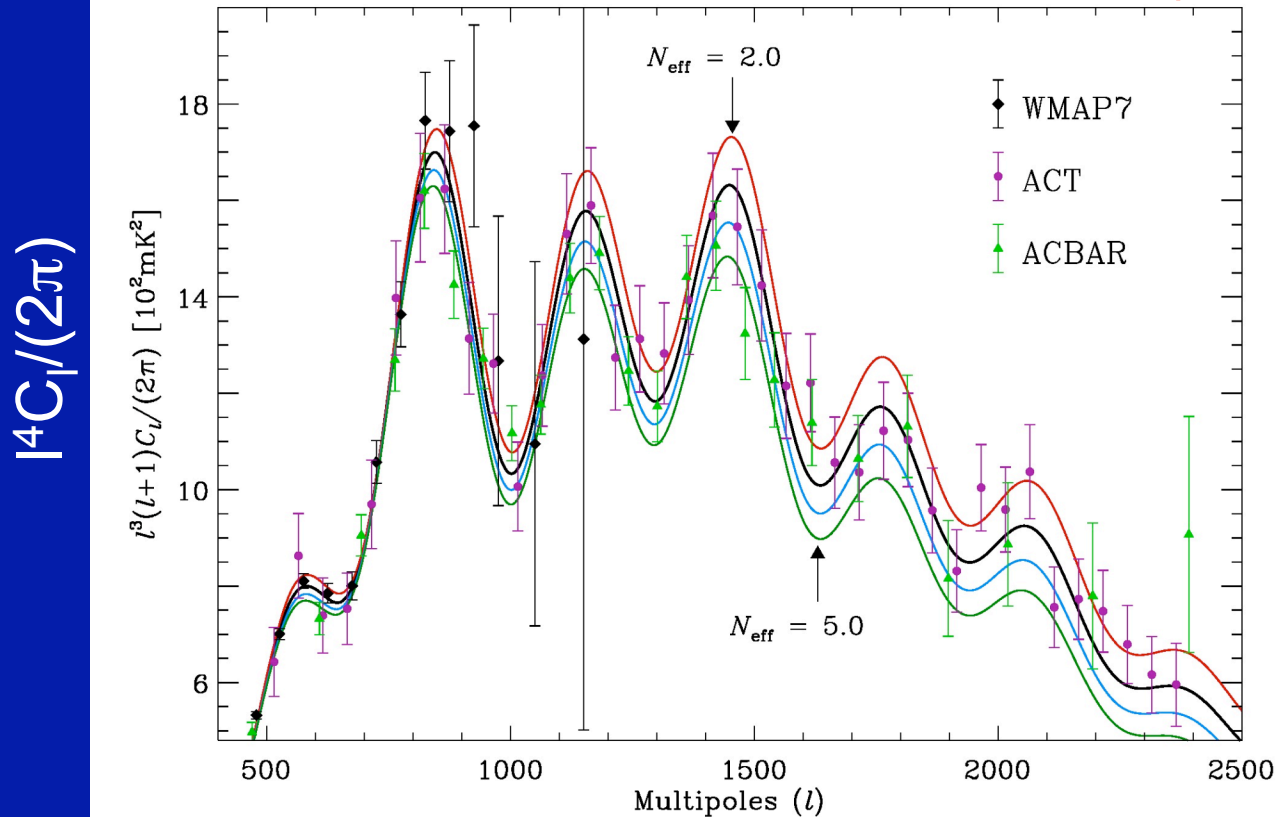


Modes with $\lambda < r_d$ are
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Surprises may come at higher resolution

Hou, Keisler, LK, Millea & Reichardt (2011)

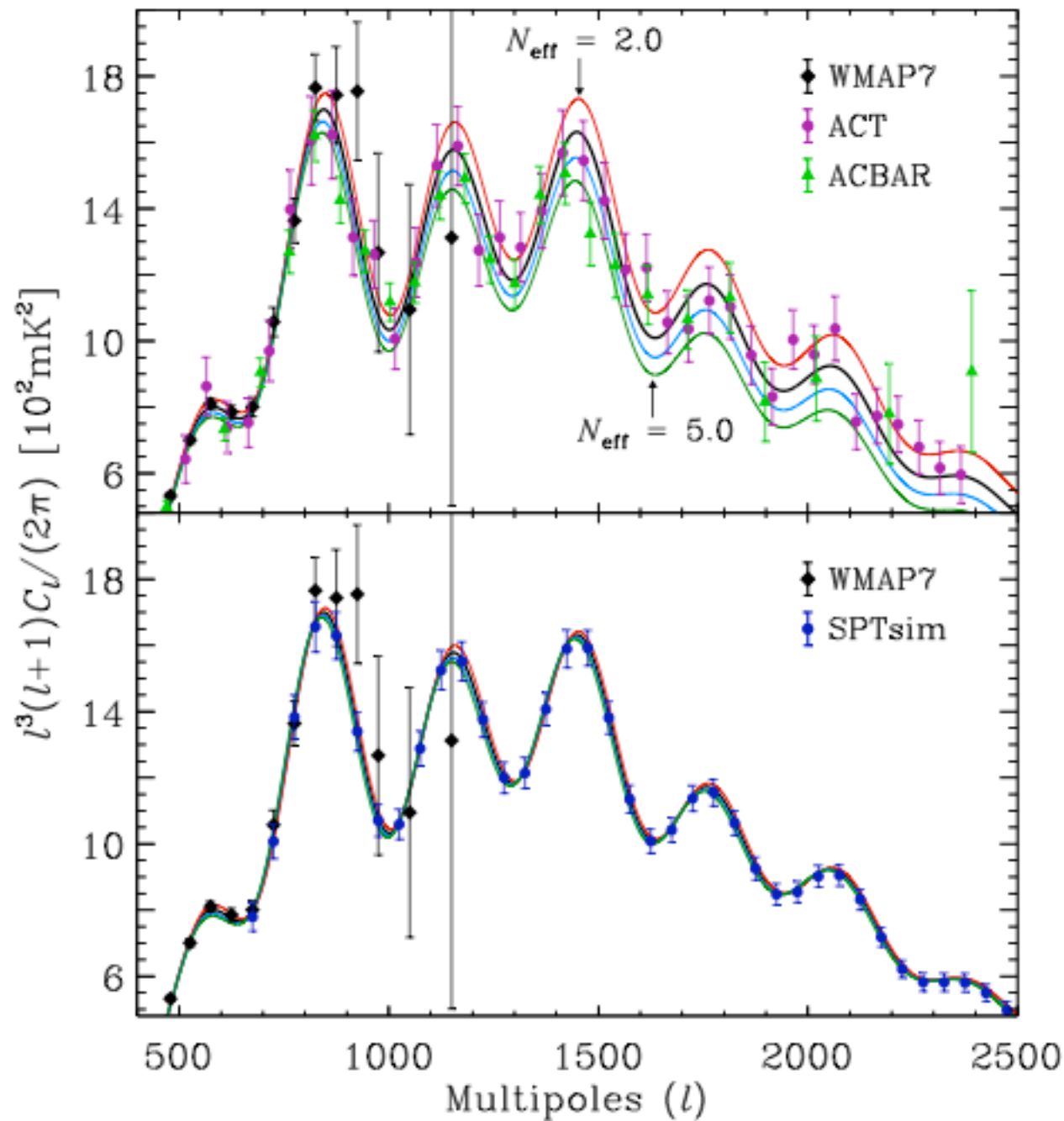


Extra neutrino species?

98.4%
confidence
that $N_{\text{eff}} >$
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model value
(Hou et al.
2011)

N_{eff} is increased here from 2 to 5 with fixed θ_{EQ} and θ_s .

To fix θ_{EQ} we increase ρ_{cdm} . To fix θ_s we adjust ρ_{Λ} to change D_A .



Same models but with θ_d fixed as well.

The effect is indeed due to change to θ_d

Bashinsky & Seljak (2004)

Extra Cosmological Neutrinos?

Arguments For

- 98.4% confidence that $N_{\text{eff}} >$ standard model value of 3.046
- Oscillation evidence for sterile neutrinos from mini-Boone / LSND
- Oscillation to sterile neutrinos can explain reactor anomalies too.
- Measurements of Y have increased in magnitude and uncertainty allowing $N_{\text{eff}} = 4$ to be consistent with BBN and perhaps preferred (Izotov & Thuan 2010, Aver, Olive & Skillman 2010, 2011)

Neutrino Fever

Neutrino Fever

6. [arXiv:1006.5276](#) [[pdf](#), [ps](#), [other](#)]

Cosmology seeking friendship with sterile neutrinos

[Jan Hamann](#), [Steen Hannestad](#), [Georg G. Raffelt](#), [Irene Tamborra](#), [Yvonne Y.Y. Wong](#)

Comments: 4 pages, 1 figure, matches version published in PRL

Journal-ref: Phys.Rev.Lett.105:181301,2010

Subjects: **High Energy Physics – Phenomenology (hep-ph)**; Cosmology and Extragalactic Astrophysics (astro-ph.CO)

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Phys. Rev. Lett. 105, 181301 (2010) [4 pages]

Cosmology Favoring Extra Radiation and Sub-eV Mass Sterile Neutrinos as an Option

Abstract

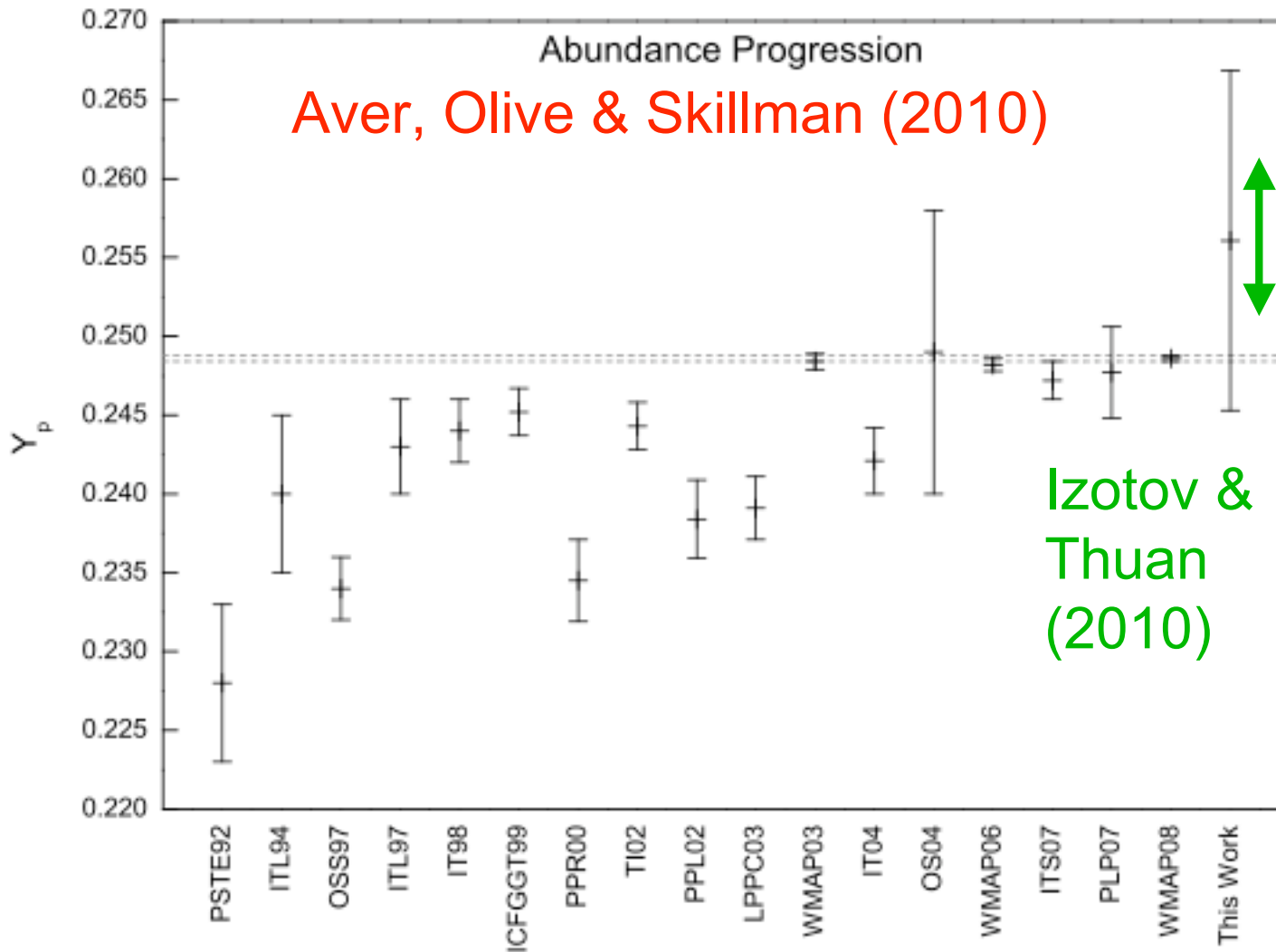
References

Citing Articles (10)

Download: [PDF \(149 kB\)](#) [Buy this article](#) [Export: BibTeX or EndNote \(RIS\)](#)

[Jan Hamann](#)¹, [Steen Hannestad](#)¹, [Georg G. Raffelt](#)², [Irene Tamborra](#)^{2,3,4}, and [Yvonne Y. Y. Wong](#)⁵

Y_p Measurements



From
extragalactic
regions of
ionized low-
metallicity
gas

(except
for WMAP
points)

Extra Neutrinos ?

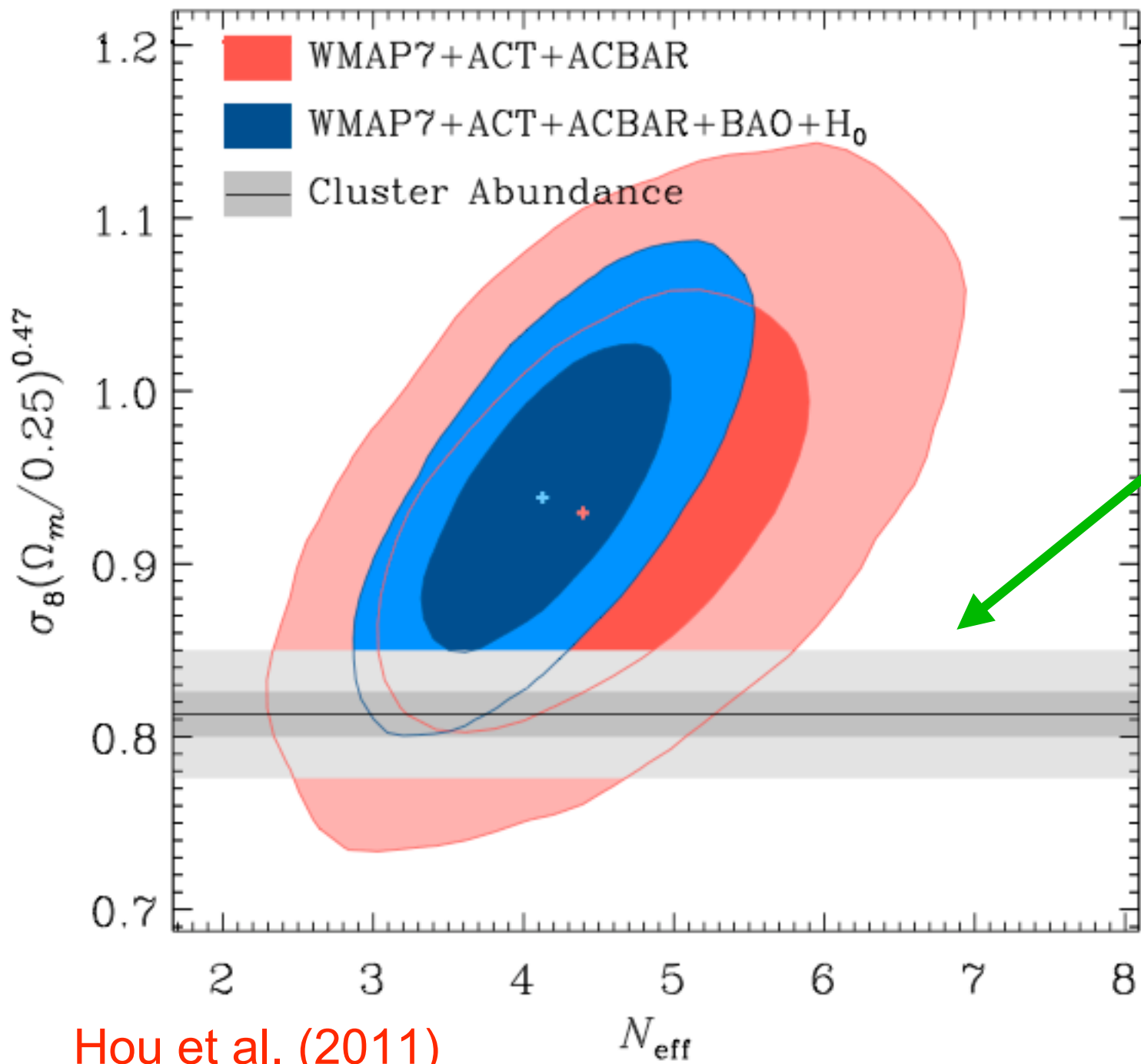
Arguments against

- 1) There are other solutions to the damping tail power deficit
 - Statistical fluke
 - Increasing Y_p (btw, that's how we kept θ_d fixed)
 - Primordial power spectrum ($dn_s/d\ln k$ non-zero)

Extra Neutrinos ?

Arguments against

- 1) There are other solutions to the damping tail power deficit
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- 2) They lead to too many galaxy clusters.



Vikhlinin et al. (2009) constraint from X-ray cluster abundances

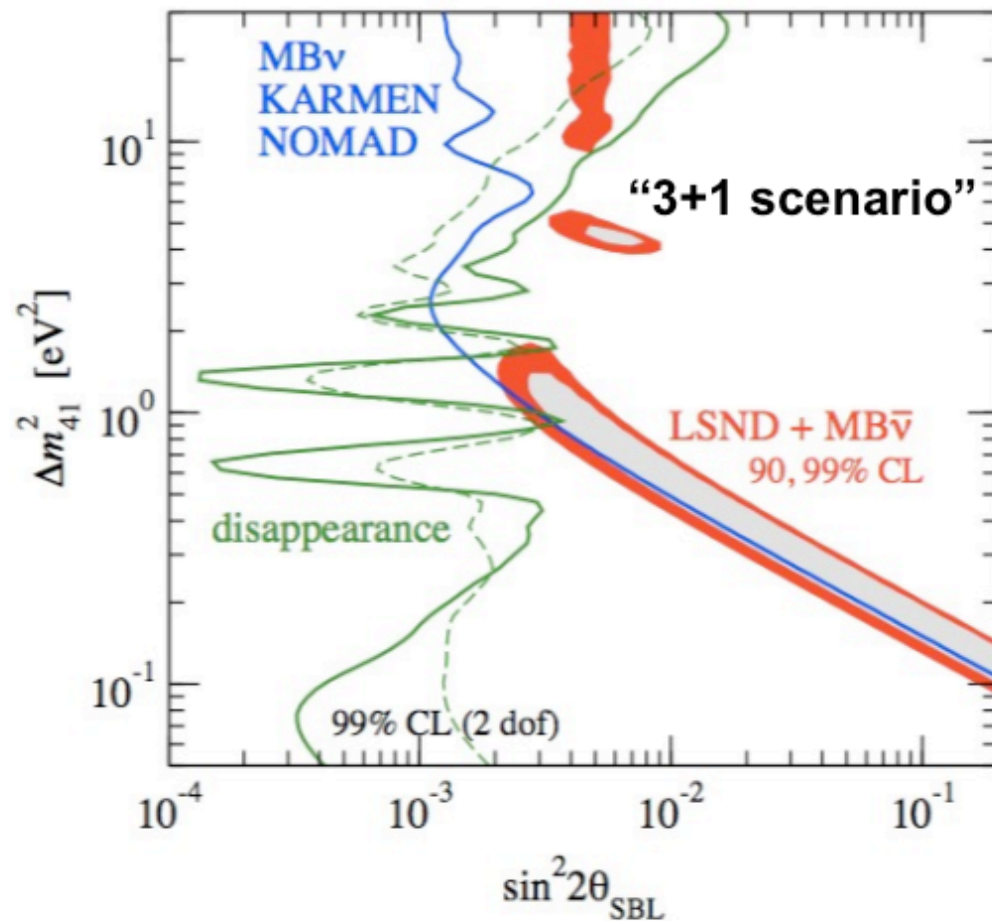
Hou et al. (2011)

Extra Neutrinos ?

Arguments against

- 1) There are other solutions to the damping tail power deficit
 - Statistical fluke
 - Increasing Y_p (btw, that's how we kept θ_d fixed)
 - Primordial power spectrum ($dn_s/d\ln k$ non-zero)
- 2) They lead to too many galaxy clusters.
- 3) Laboratory and Reactor neutrino oscillation solutions require ~ 1 eV masses. Too massive. They would cluster on large scales, altering shape of matter power spectrum.

From Yvette Wong's Avignon presentation



Kopp, Maltoni &
Schwetz 1103.4570

— New reactor fluxes
- - - Old reactor fluxes

- **Tension** between LSND/MiniBooNE and reactor disappearance experiments.

- **New analysis** of reactor fluxes finds 3% higher mean flux.
→ disappearance @98.6% confidence (old: 68% CL).

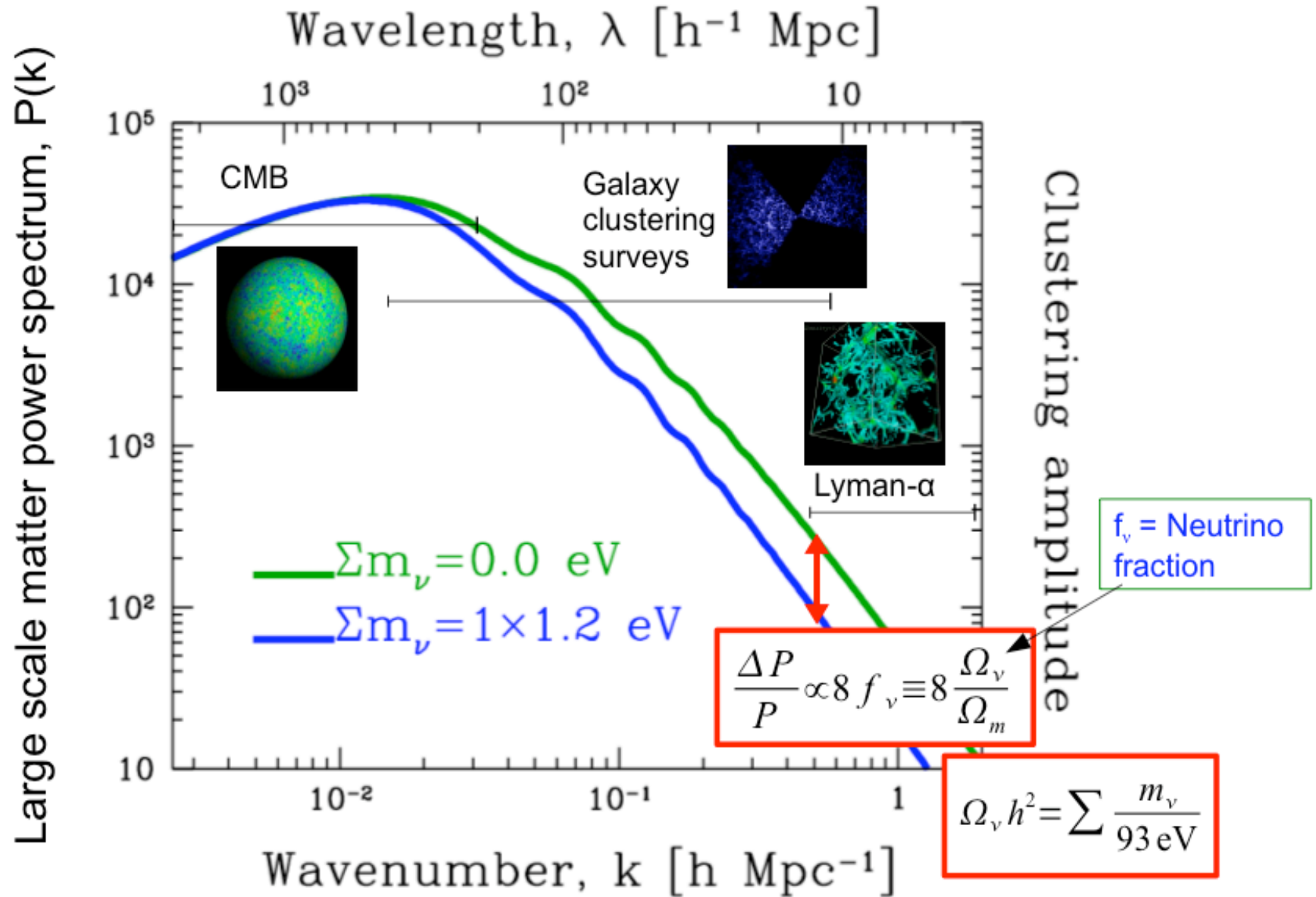
Mention et al. 1101.2755

- “3+1” best-fit: $\Delta m_{41}^2 \sim 1 \text{ eV}^2$

→ $m_s \sim 1 \text{ eV}$

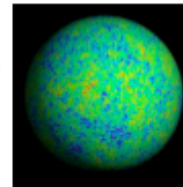
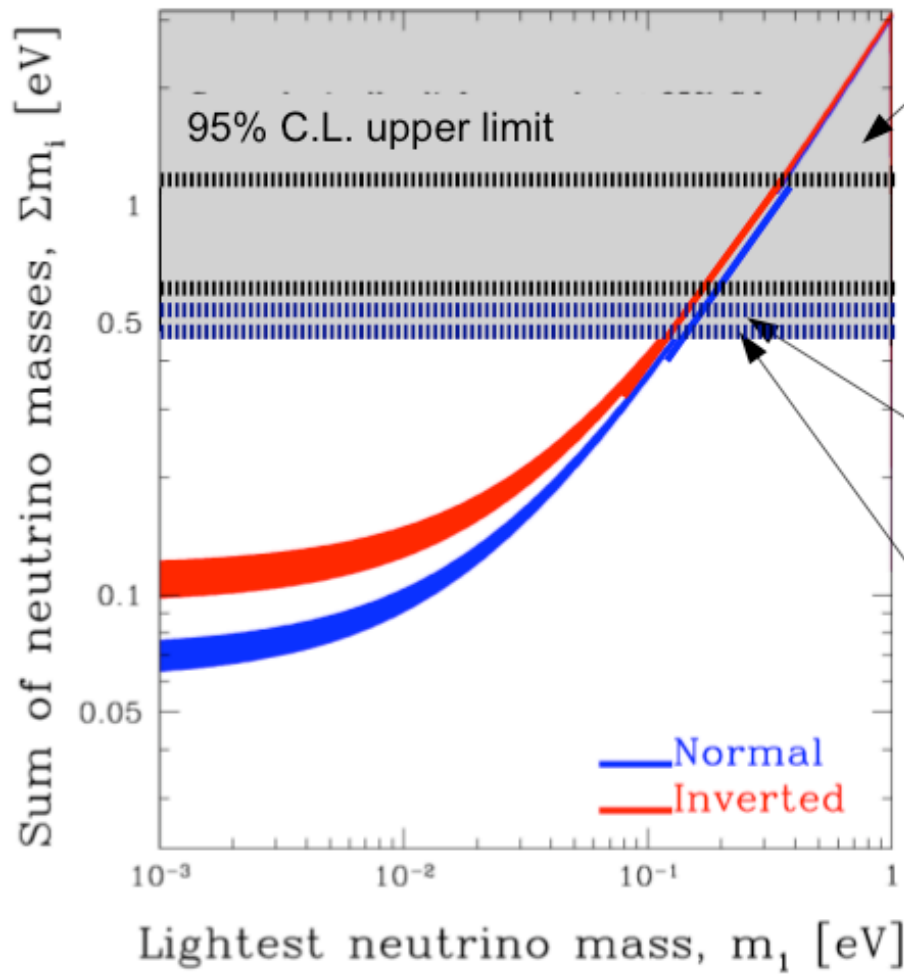
If lightest neutrino mass $\sim 0 \text{ eV}$

From Yvette Wong's Avignon presentation

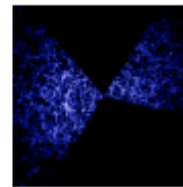


From Yvette Wong's Avignon presentation

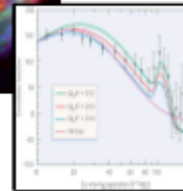
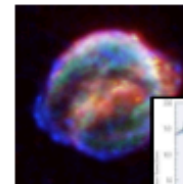
Present status...



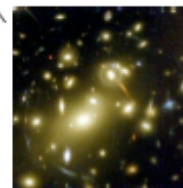
WMAP7 only*
Komatsu et al. 2010



WMAP7+SDSS-HPS*
Hannestad, Mirizzi, Raffelt
& Y³W 2010



WMAP5+SDSS-HPS
+SN+HST
Reid et al. 2009
(extended models)



WMAP5+Weak lensing*
Tereno et al. 2008
Ichiki et al. 2008

* Λ CDM+ m_ν

So what's going on?

- If there are extra sterile neutrinos (as lab and nuclear experiments indicate), they better not get thermally produced in the early Universe because their masses are too large.
- Lab and nuclear experiments indicate two extra species. Could there be a third sterile neutrino, produced in early U that is massive enough to be the dark matter?

So what's going on?

- If there are extra sterile neutrinos (as lab and nuclear experiments indicate), they better not get thermally produced in the early Universe because their masses are too large.
- Lab and nuclear experiments indicate two extra species. Could there be a third sterile neutrino, produced in early U that is massive enough to be the dark matter?

16. [arXiv:0812.3256](#) [[pdf](#), [ps](#), [other](#)]

Realistic sterile neutrino dark matter with keV mass does not contradict cosmological bounds

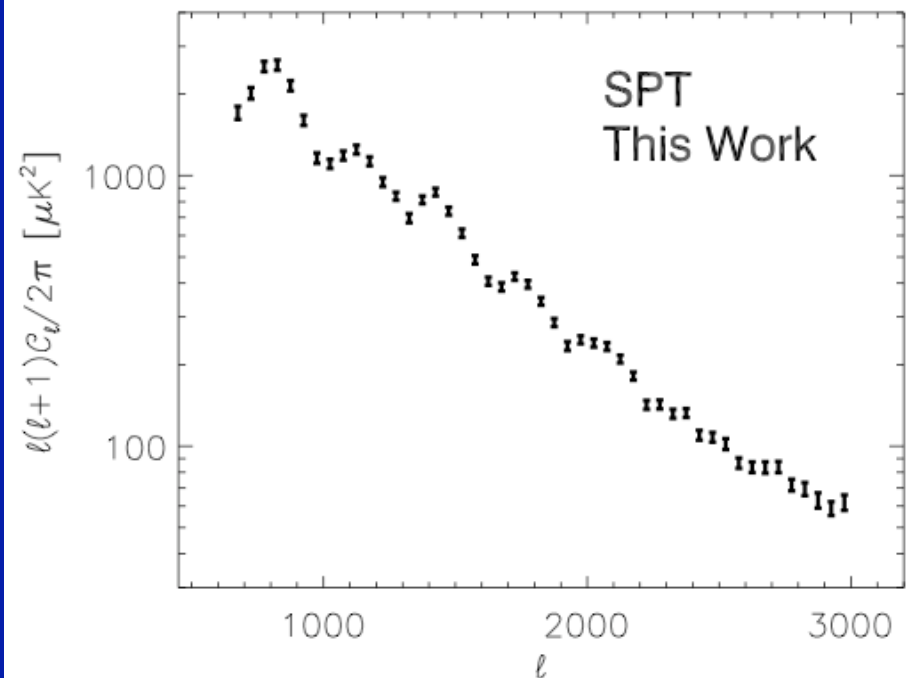
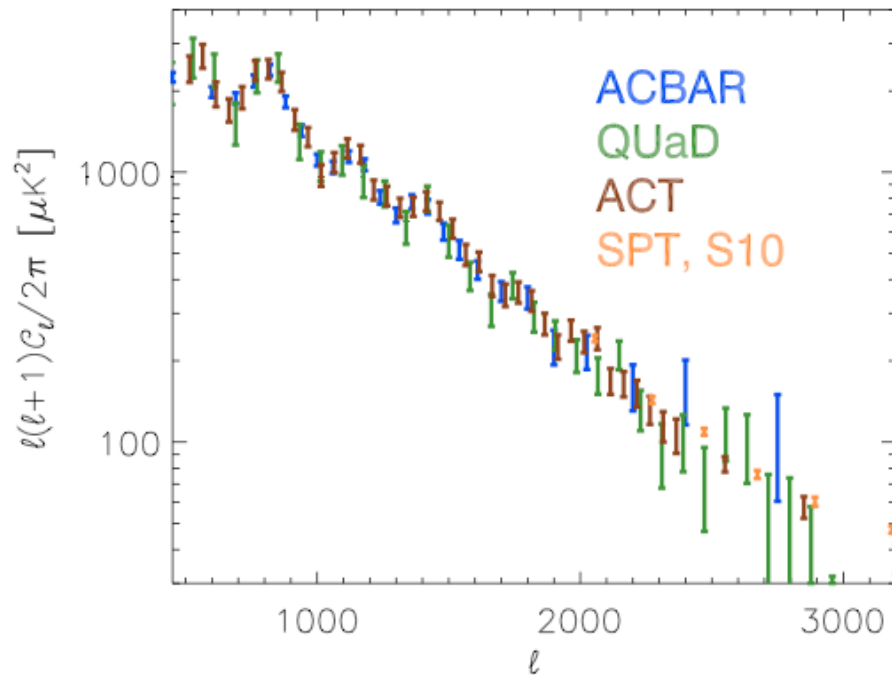
Alexey Boyarsky, Julien Lesgourgues, Oleg Ruchayskiy, Matteo Viel

Comments: 4 pages, 4 figures

Journal-ref: Phys.Rev.Lett.102:201304,2009

Subjects: **High Energy Physics - Phenomenology (hep-ph)**; Astrophysics (astro-ph)

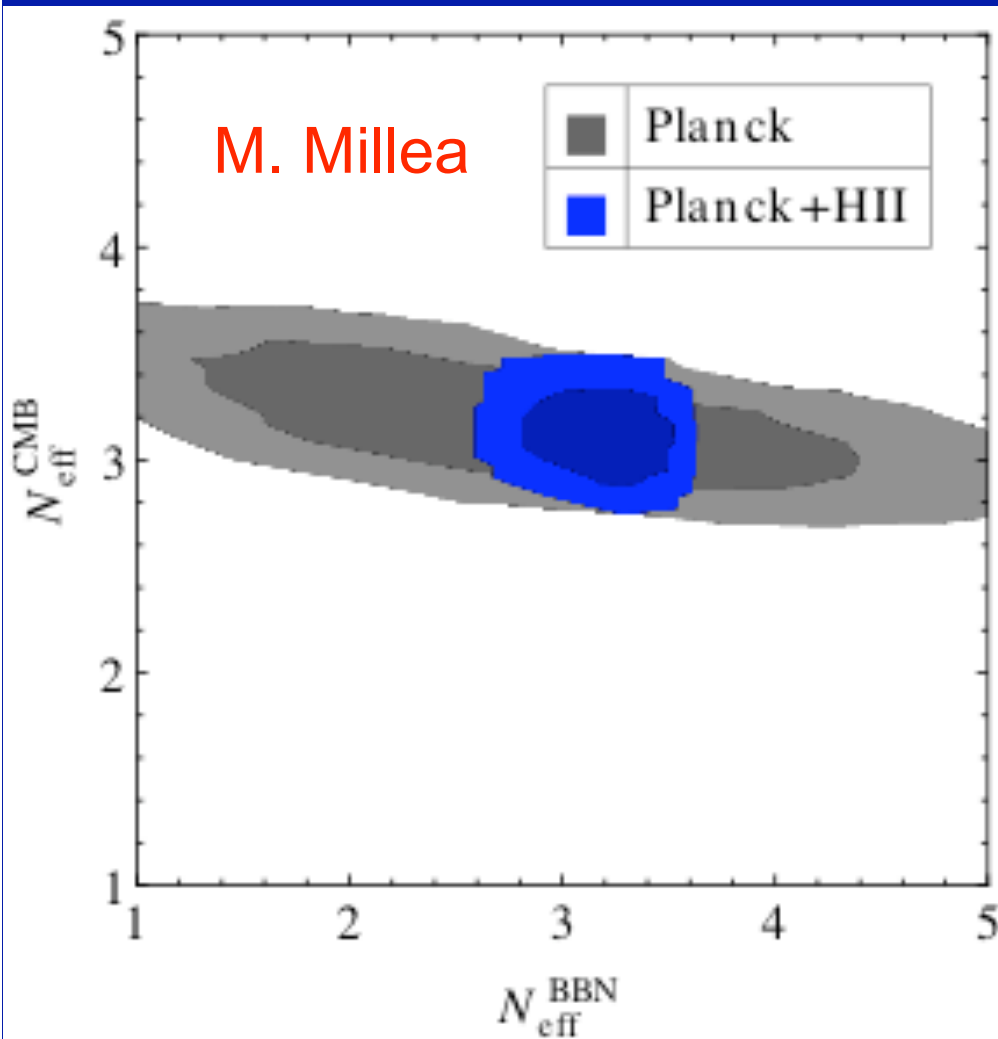
Advertisement for Keisler et al. (SPT, in prep)



Data available now

Data available soon

The Future



With better data we can relax assumption that $N_{\text{eff}}^{\text{BBN}} = N_{\text{eff}}^{\text{CMB}}$ (so far assumed implicitly throughout this talk).

Forecast for Planck

Forecast for Planck + Y_{p} measurement with error same size as reported by Izotov & Thuan (2010).

With luck, these will disagree! (e.g. Fischler & Myers (2010))

High-resolution CMB Temperature Anisotropies as a Particle Physics Detector

- High-resolution CMB data, by being sensitive to the diffusion scale, allow us to determine expansion rate up to recombination -- similar to Helium abundance sensitivity to expansion rate during BBN.
- There are interesting things happening now regarding BBN, high-res CMB, and lab and reactor neutrino data.
- It is not obvious how this all fits together.
- New damping tail measurements will be out very soon, and dramatically improved ones will be out from Planck in ~ Jan 2013.
- Planck measurements will provide tight constraints on (or detections of) hot relics.