

# Higgs Boson Mass Limit in GMSB with Messenger-Matter Mixing


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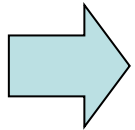
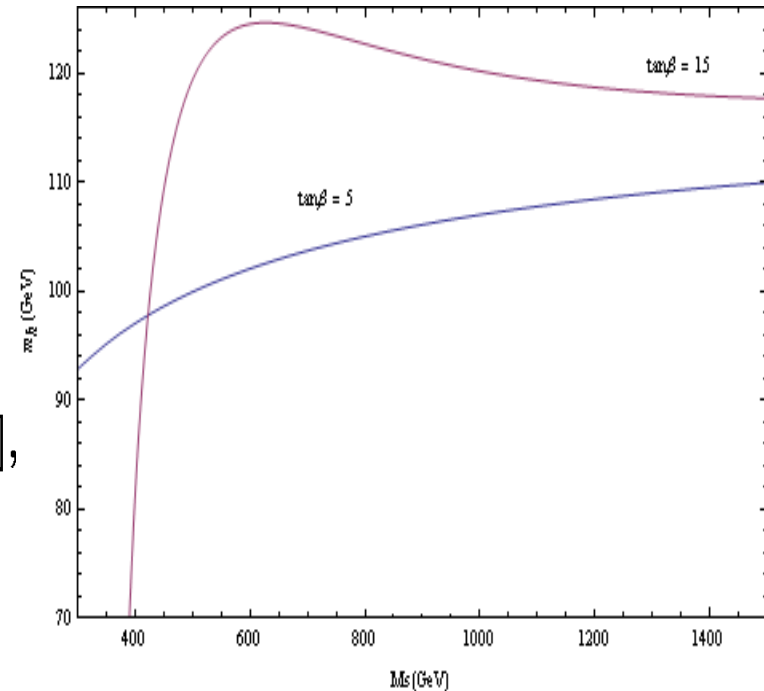
# Introduction: Higgs Mass Bounds in MSSM

Tree Level

$$m_h < M_z \tan(\beta) \quad \text{Excluded by LEP2 (114.4 GeV)}$$

1-and 2- loop

$$m_h^2 = M_z^2 \cos^2 2\beta \left( 1 - \frac{3}{8\pi^2} \frac{m_t^2}{v^2} t \right) + \frac{3}{4\pi^2} \frac{m_t^4}{v^2} \left[ \frac{1}{2} \chi_t + t + \frac{1}{16\pi^2} \left( \frac{3}{2} \frac{m_t^2}{v^2} - 32\pi\alpha_3 \right) (\chi_t t + t^2) \right],$$



# Introduction: GMSB without Messenger-Matter Mixing

$$M_{\lambda_i} = \frac{\alpha_r(M)}{4\pi} \Lambda, \quad \Lambda = \frac{F}{M}$$

$$\tilde{m}^2 = 2N \sum_{r=1}^3 c_r \left( \frac{\alpha_r(M)}{4\pi} \right)^2 \Lambda^2.$$



Universal Scalar Masses

## Features of GMSB

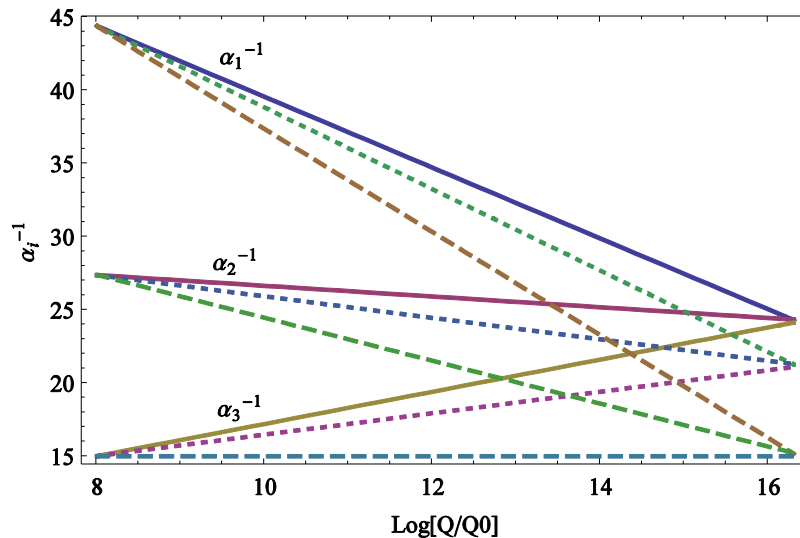


Highly predictive



FCNC are naturally suppressed

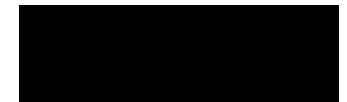
$5 + \bar{5}, 10 + \bar{10}$



Is it possible to obtain maximal mixing in the ordinary GMSB?

No, because  $A_t \approx 0$  at  $M_{mess}$

Messenger-matter mixing with messenger fields belong to  $10 + \bar{10}$  can reproduce



# Introduction: The Objectives

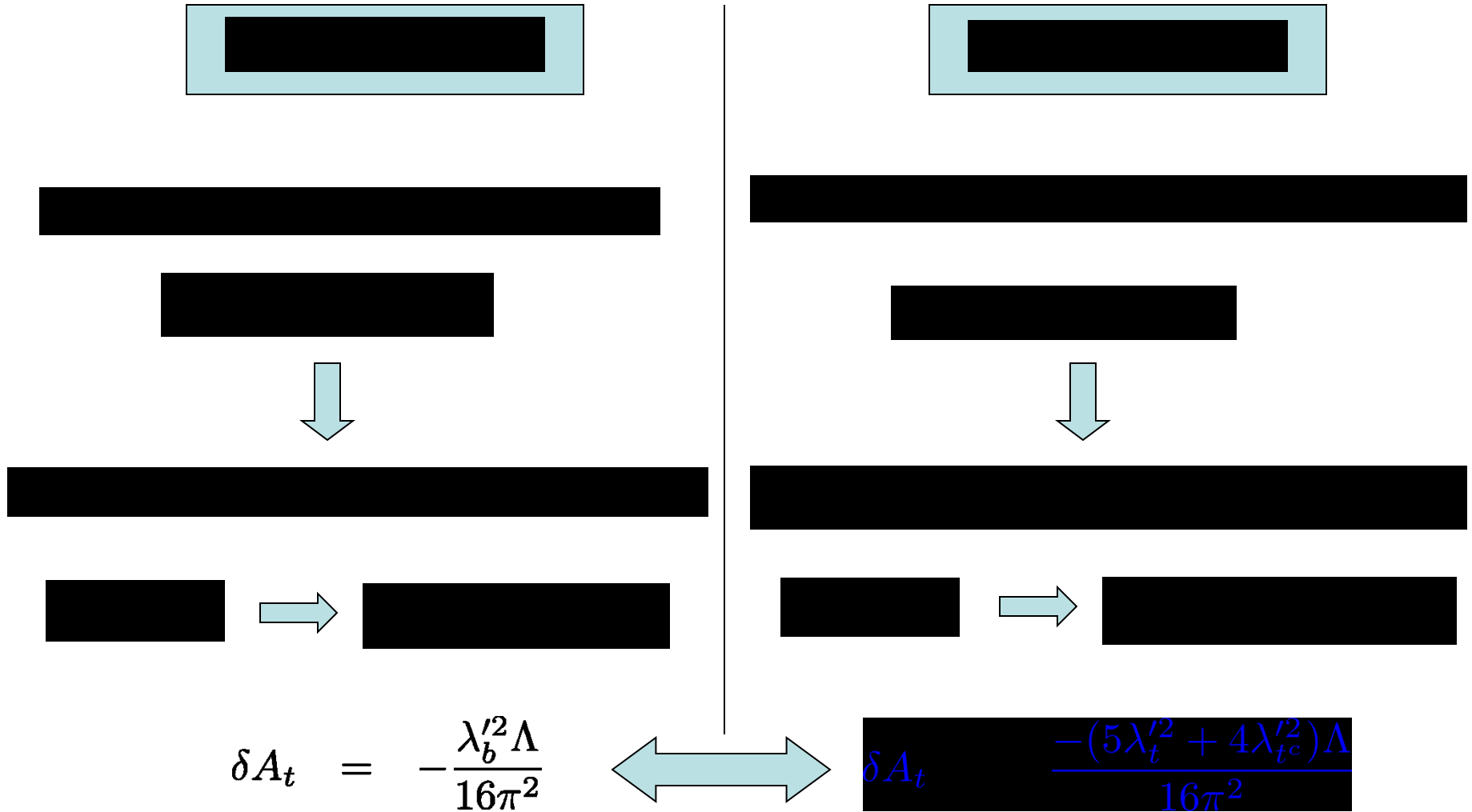
To construct GMSB model with messenger-matter mixing that

- 1- leads to a significant enhancement to the lightest Higgs mass compared to the ordinary GMSB.
- 2- leads to scalar mass spectra below 1 TeV.

The above objectives should be consistent with

- 1- FV processes are suppressed in agreement with experiment .
- 2- Messenger scale below  $3 \times 10^8$  GeV is preferred by Cosmology.

# GMSB with Messenger-Matter Mixing



The expressions for soft terms induced by messenger-matter mixing can be found in [ Giudice and Rattazi, 1998; Chacko and Ponton, 2002 ]

# Higgs Mass Bounds in the XXXXXXXXXX Model

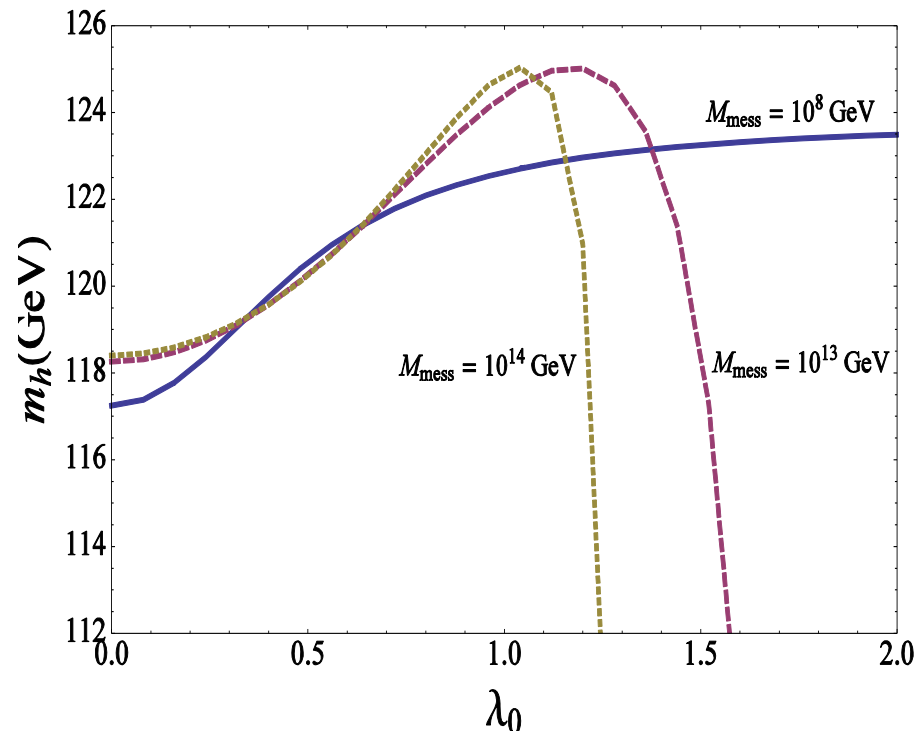
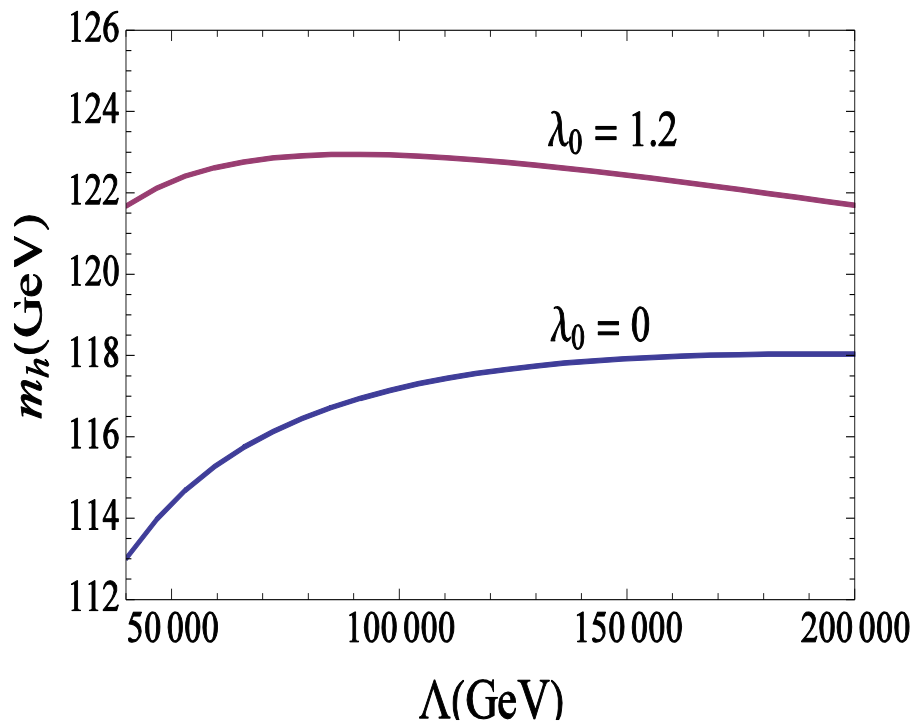
XXXXXXXXXX  $\longrightarrow$  XXXXXX There are three parameters:

XXXXXXXXXX XXXXXXXXXX XXXXXXXXXX

$\lambda_0$	$m_h$ (GeV)	$\Lambda$ ( $10^5$ GeV)	$M_{mess}$ ( $10^8$ GeV)	$\tilde{m}_{t_1}$ (GeV)	$\tilde{m}_{t_2}$ (GeV)	$A_t/M_s$
0	<b>119</b>	1.6	$3.16 \times 10^5$	<b>3590</b>	4145	-0.86
0.4	120	1.36	1	2756	3289	-1.1
0.8	123	0.912	$10^5$	1553	2143	-1.55
1.2	<b>125</b>	0.784	17782	<b>1088</b>	1751	-1.95
1.6	125	0.784	1778	1066	1743	-2
2	125	0.784	177	1138	1762	-1.93

Table 1: We show the values of the GMSB input parameters,  $\Lambda$ ,  $\lambda_0$  and  $M_{mess}$  that lead to the highest  $m_h$  values. These values correspond to  $\lambda_{m0} = 0$  and  $\tan \beta = 10$ .

# Higgs Mass Bounds in the ██████████ Model





# Higgs Mass Bounds in the ██████████ Model

Name		$10 + \overline{10}$	$10 + \overline{10}$	$5 + \overline{5}$
Inputs	$M_{mess}$	$10^8$	$4 \times 10^5$	$10^8$
	$N_{mess}$	3	3	1
	$\Lambda$	$0.3 \times 10^5$	$0.3 \times 10^5$	$0.95 \times 10^5$
	$\tan \beta$	10	5.6	11.6
	$\lambda_0$	1.2	1.2	1.2
Higgs:	$m_h$	121	117.7	114.6
	$m_H^0$	675	675	1107
	$m_A$	675	674	1107
	$m_{H^\pm}$	679	678	1110
Gluino:	$m_{\tilde{g}}$	852	852	899
Neutralinos:	$m_{\chi_1}$	121	127	128
	$m_{\chi_2}$	234	245	248
	$m_{\chi_3}$	667	658	706
	$m_{\chi_4}$	675	668	713
Charginos:	$\chi_1^+$	236	233	250
	$\chi_2^+$	676	667	738
Squarks:	$m_{\tilde{u}_L, \tilde{c}_L}$	810	787	1120
	$m_{\tilde{u}_R, \tilde{c}_R}$	786	765	1071
	$m_{\tilde{d}_L, \tilde{s}_L}$	810	787	1121
	$m_{\tilde{d}_R, \tilde{s}_R}$	782	763	1064
	$m_{\tilde{b}_L}$	692	682	997
	$m_{\tilde{b}_R}$	780	763	1045
	$m_{\tilde{t}_L}$	692	682	997
	$m_{\tilde{t}_R}$	518	531	890
Sleptons:	$m_{\tilde{e}_L, \tilde{\mu}_L}$	224	201	371
	$m_{\tilde{\nu}_{eL}, \tilde{\nu}_{\mu L}}$	224	201	371
	$m_{\tilde{e}_R, \tilde{\mu}_R}$	168	150	182
	$m_{\tilde{\tau}_L}$	224	201	352
	$m_{\tilde{\tau}_R}$	167	150	1014

# Flavor Violation

Mass Insertion Parameters:

The messenger-matter couplings reintroduce the flavor violation

$\delta\tilde{m}^2$

$\delta A$

$\delta\tilde{m}^2$ ,  $\delta A$  are generated by the exotic Yukawa couplings.

Froggatt-Nielsen Mechanism:

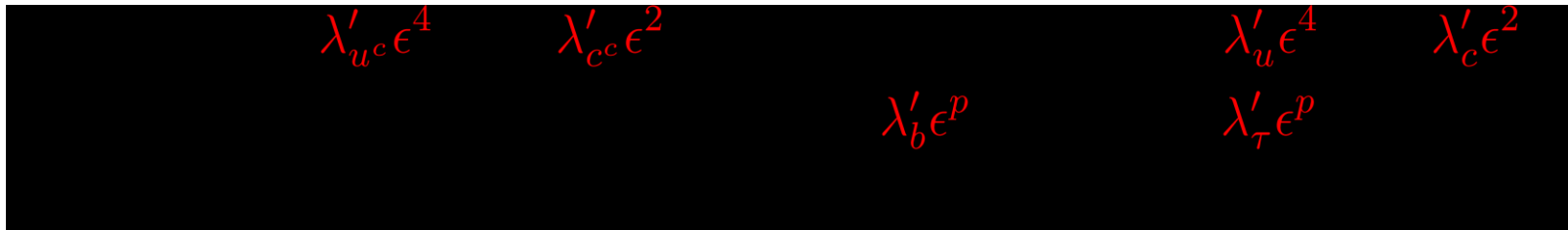
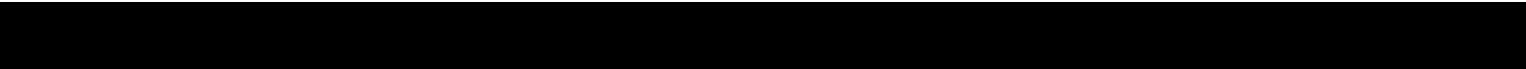
- ✧ U(1) flavor symmetry is assumed.
- ✧ U(1) is broken at high scale  $M^*$  by  $\langle S \rangle$
- ✧ The hierarchy of the masses and mixing can be explained as power expansion of



# Flavor Violation

SU(5)	$10_1$	$10_2$	$10_3$	$\bar{5}_1$	$\bar{5}_2, \bar{5}_3$	$5_u, 5_d$	$S$	$5_m$	$\bar{5}_m$	$10_m$	$\bar{10}_m$	$Z$
$U(1)$	4	2	0	p+1	p	0	-1	$-\alpha$	0	0	$-\alpha$	$\alpha$

Table 1: The  $U(1)$  charge assignments to the messenger, MSSM, Z and S fields.



$$M_d \sim M_e^T \sim \epsilon^p \begin{pmatrix} \epsilon^5 & \epsilon^3 & \epsilon \\ \epsilon^4 & \epsilon^2 & 1 \\ \epsilon^4 & \epsilon^2 & 1 \end{pmatrix}, \quad \longrightarrow \quad \text{Lopsided Structure}$$

$$U_L^e \sim U_R^d \sim \begin{pmatrix} 1 & \epsilon & \epsilon \\ \epsilon & \omega & \omega \\ \epsilon & \omega & \omega \end{pmatrix}, \quad U_R^e \sim U_L^d \sim \begin{pmatrix} 1 & \epsilon^2 & \epsilon^4 \\ \epsilon^2 & 1 & -\epsilon^2 \\ \epsilon^4 & \epsilon^2 & 1 \end{pmatrix}.$$

# Flavor Violation

Process	Mass Insertion ( $\delta$ )	$5 + \bar{5}$	$10 + \bar{10}$	Exp. Bounds*
$\mu \rightarrow e\gamma$	$(\delta_{12}^l)_{LL}$	-	$\epsilon^{4p+1}$	<b>0.02</b>
	$(\delta_{12}^l)_{RR}$	$\epsilon^6$	-	0.0087
	$(\delta_{12}^l)_{RL,LR}$	$\kappa_5^l(\epsilon^{p+4}, \epsilon^{p+3})$	$\kappa_{10}^l \epsilon^{3p+1}$	<b><math>7 \times 10^{-6}</math></b>
$\tau \rightarrow e\gamma$	$(\delta_{13}^l)_{LL}$	-	$\epsilon^{4p+1}$	75
	$(\delta_{13}^l)_{RR}$	$\epsilon^4$	-	22
	$(\delta_{13}^l)_{RL,LR}$	$\kappa_5^l(\epsilon^{p+4}, \epsilon^{p+1})$	$\kappa_{10}^l \epsilon^{3p+1}$	0.44
$\tau \rightarrow \mu\gamma$	$(\delta_{23}^l)_{LL}$	-	$\epsilon^{4p}$	14
	$(\delta_{23}^l)_{RR}$	$\epsilon^2$	-	4
	$(\delta_{23}^l)_{RL,LR}$	$\kappa_5^d(\epsilon^{p+2}, \epsilon^p)$	$\kappa_{10}^l \epsilon^{3p}$	0.08
$K - \bar{K}$	$(\sqrt{\text{Re}(\delta_{12}^d)_{LL}^2}, \sqrt{\text{Im}(\delta_{12}^d)_{LL}^2})$	$\epsilon^6$	$\epsilon^6$	(0.065, 0.0052)
	$(\sqrt{\text{Re}(\delta_{12}^d)_{RR}^2}, \sqrt{\text{Im}(\delta_{12}^d)_{RR}^2})$	-	$\epsilon^{1+4p}$	(0.065, 0.0052)
	$(\sqrt{\text{Re}(\delta_{12}^d)_{LR}^2}, \sqrt{\text{Im}(\delta_{12}^d)_{LR}^2})$	$\kappa_5^d \epsilon^{4+p}$	$\kappa_{10}^d \epsilon^{1+3p}$	(0.007, $5.2 \times 10^{-5}$ )
	$(\sqrt{\text{Re}(\delta_{12}^d)_{LR}^2}, \sqrt{\text{Im}(\delta_{12}^d)_{LR}^2})$	$\kappa_5^d \epsilon^{3+p}$	$\kappa_{10}^d \epsilon^{1+3p}$	(0.007, $5.2 \times 10^{-5}$ )
	$\sqrt{\text{Re}(\delta_{12}^d)_{LL}(\delta_{12}^d)_{RR}}$	-	$\epsilon^{3.5+2p}$	0.00453
	$\sqrt{\text{Im}(\delta_{12}^d)_{LL}(\delta_{12}^d)_{RR}}$	-	$\epsilon^{3.5+2p}$	0.00057
$B_d - \bar{B}_d$	$(\text{Re}\delta_{13}^d, \text{Im}\delta_{13}^d)_{LL}$	$\epsilon^4$	$\epsilon^4$	(0.238, 0.51)
	$(\text{Re}\delta_{13}^d, \text{Im}\delta_{13}^d)_{RR}$	-	$\epsilon^{1+4p}$	(0.238, 0.51)
	$(\text{Re}\delta_{13}^d, \text{Im}\delta_{13}^d)_{LR,RL}$	$\kappa_5^d(\epsilon^{4+p}, \epsilon^{1+p})$	$\kappa_{10}^d \epsilon^{1+3p}$	(0.0557, 0.125)
$B_s - \bar{B}_s$	$(\delta_{23}^d)_{LL}$	$\epsilon^2$	$\epsilon^2$	1.19
	$(\delta_{23}^d)_{RR}$	-	$\epsilon^{1+4p}$	1.19
$b \rightarrow s\gamma$	$(\delta_{23}^d)_{LR,RL}$	$\kappa_5^d(\epsilon^{p+2}, \epsilon^p)$	$\kappa_{10}^d(\epsilon^2, 1)$	0.04

\* F. Gabbiani, E. Gabrielli, A. Masiero and L. Silvestrini, Nucl. Phys. B **477**, 321 (1996) [arXiv:hep-ph/9604387]; D. Becirevic *et al.*, Nucl. Phys. B **634**, 105 (2002) [arXiv:hep-ph/0112303]; K. S. Babu and Y. Meng, Phys. Rev. D **80**, 075003 (2009) [arXiv:0907.4231 [hep-ph]];

# Conclusion

- ✦ The maximal mixing condition that leads to the upper limit of the lightest Higgs mass of MSSM is obtained in the model where the messenger fields belong to  $\mathbf{5}_H$  of SU(5).
- ✦ Consistent with cosmology preference  $m_0 \leq 3 \times 10^8$  GeV. This model can lead to the lightest Higgs mass of around 121 GeV with all superparticle masses below 1 TeV.
- ✦ These results are consistent with the gauge and exotic Yukawa couplings being perturbative and unified at the GUT scale as well as the FCNC being suppressed in agreement with experimental bounds.