## Dark matter in Warped Universal Extra Dimensions

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## **Outline**

- Motivation
- 2 Backgrounds with Z<sub>2</sub> Symmetry
- KK-Radion Dark Matter
  - Freeze-out KK-Radion DM ( $\Lambda_r \approx 2 \text{ TeV}$ )
  - Coannihilations with t' ( $\Lambda_r \approx 10 100 \text{ TeV}$ )
  - KK-radion as a SuperWIMP ( $\Lambda_r > 100 \text{ TeV}$ )
- Direct and Indirect Detection
- Conclusions

#### Motivation

- The possible link between dark matter (DM) and particle physics phenomena at the TeV scale has prompted many models with a suitable DM candidate in the TeV mass range such as Supersymmetry with R-parity, little Higgs with T-parity among others.
- In extra dimensional theories the best known example of a theory that naturally provides a DM candidate is Universal Extra Dimensions (UED).
- We seek to construct a minimal scenario that combines the "nice" features of UED and RS.
- We study a Z<sub>2</sub> warped geometry, addressing the stabilization of the extra dimension.
- Through the stabilization mechanism, we dynamically generate a mass for the lightest radion even state and the lightest odd excitation of the radion field becomes our DM candidate, with a mass parametrically smaller than the KK scale.

# Backgrounds with Z<sub>2</sub> Symmetry

- Consider a 5D real scalar Φ minimally coupled to gravity,
- Interest in the stabilization of the extra dimension leading to a symmetric background about y = 0,  $y \in [-L, L]$ .
- Solve the coupled gravity/scalar system taking into account the backreaction of the scalar VEV on the geometry (generate a mass for the zero mode radion).
- Restrict to backgrounds with 4D Lorentz symmetry,

$$ds^2 = e^{-2A(y)} \eta_{\mu\nu} dx^{\mu} dx^{\nu} - dy^2$$
, (1)

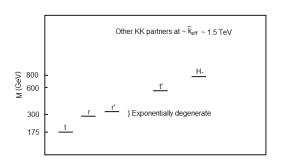
An approximate solution is given by (IR-UV-IR)

$$\phi(y)/\phi_0 \approx s_\phi y/L$$
,  $A(y) \approx s_A k y^2/(2L) = k_{\text{eff}} y^2/(2L)$ , (2)

• For the radion we find that its profile is given by  $F_0 \approx e^{2(A(y)-A(L))}$ , and that its mass scales as,

$$m_0 pprox rac{2}{\sqrt{k_{eff}L}} k_{eff} e^{-A(L)} pprox 0.25 \, \tilde{k}_{eff} \,,$$
 (3)

 Profiles that are localized near the IR branes have exponentially degenerate mass for odd and even modes. • Typical spectrum in this kind of scenario.



## The KK-Radion as a WIMP

Radion interactions are non-renormalizable and controlled by the decay constant

$$\Lambda_r = \sqrt{\frac{3M_5^3}{k_{\rm eff}^3}} imes \tilde{k}_{
m eff}.$$

- : Allows to interpolate between RS-like "strong warping" scenarios (when  $k_{\rm eff} \sim M_5$ ) and "UED-like" scenarios (when  $k_{\rm eff} \ll M_5$ ).
- In the strong warping scenario, KK-radions can annihilate into SM fermion (mainly top) pairs and Higgses (annihilation into massless gauge bosons is suppressed),

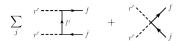
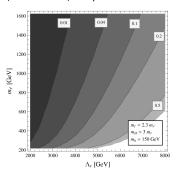


Figure: Annihilation into fermions



Figure: Annihilation into Higgs and longitudinal gauge bosons

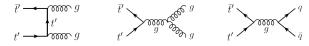
• Contours of constant  $\Omega_{r'}h^2$  in the  $m_{r'}-\Lambda_r$  plane



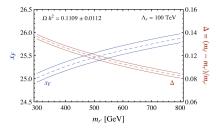
• Correct relic density of  $\Omega_{r'}h^2\approx 0.1$  can be obtained for natural values of the parameters ( $\Lambda_r\approx 3.5$  TeV,  $m_{r'}\approx 300$  GeV).

## Coannihilations with t'

• If  $\Lambda_r \sim 100$  TeV and  $m_{t'} \gtrsim m_{r'}$ , coannihilations with the strongly interacting t' can become dominant and the  $\Omega_r h^2$  can be completely controlled by QCD.



• WMAP constraint on the degree of degeneracy,  $\Delta = (m_{t'} - m_{r'})/m_{r'}$  as a function of  $m_{r'}$ 



#### KK-Radion as a Non-Thermal Relic

- Situation arises when  $\Lambda_r \gg k_{eff}$  (superweak radion coupling).
- Assume that  $T_R \lesssim T_c \approx 1$  TeV, with  $T_c$  the critical temperature for the deconfinement/confinement phase transition. Otherwise, the universe trapped eternally in the false vacuum.
- Production channels from scattering and decays:  $gq^j \to r'q^k$ ,  $f^j \to r'f^k$ ,  $V^{j,a}_{j} \to r'V^{k,b}_{j}$

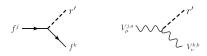


Figure: Production through decays

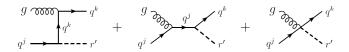
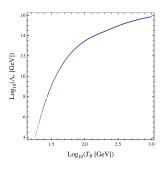


Figure: Production through scattering

• Decays are the dominant production mechanism. We study formally up to  $T_R \sim 10 k_{\rm eff}$  (4D cut-off of KK-theory).



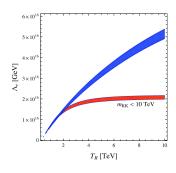


Figure: WMAP constraint in the  $\Lambda_R$ - $T_R$  plane in Log scale. Figure: WMAP constraint in the  $\Lambda_R$ - $T_R$  up to  $T_R = 10 \text{ TeV}$ 

BBN constraints are satisfied as long as T<sub>R</sub> ≤ few hundred GeV.

## **Direct Detection**

- Consider brane Higgs-Radion mixing term,  $\frac{1}{2} \left[ \delta(y-L) + \delta(y+L) \right] \sqrt{g_{\text{ind}}} \, \xi \mathcal{R}_4 H^{\dagger} H$ .
- Interactions relevant for scattering of DM against nuclei are  $r_- h_+ h_-$  and  $r_+ h_- h_-$ .

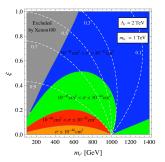


Figure: Contours of constant  $\sigma_{\phi_{\rm DM}N\to\phi_{\rm DM}N}$  in the  $\xi$ - $m_{r_-}$  plane for  $\Lambda_r=2~{\rm TeV}$ 

• There is a region where the DM is mostly  $r_-$ , with a sizable nucleon cross section (below the dashed line of  $|U_{r_-,L}|^2 = 0.9$ ).

## Indirect detection

- Focus on the photon flux from KK-radion annihilations (other signals are small).
- Consider the brane localized operator  $[\delta(y+L)+\delta(y-L)]\sqrt{g_{ind}}(\eta/8\Lambda_5^3)K^2F_{\mu\nu}F^{\mu\nu}$ , where  $\Lambda_5$  is the 5D cutoff and K is the trace of the extrinsic curvature.
- Reduction to 4D leads to  $-\frac{e^2\kappa}{8\Lambda^2} r'^2 F_{\mu\nu} F^{\mu\nu}$  with  $e^2\kappa \approx 64\eta/(k_{\rm eff}L) \times (k_{\rm eff}/\Lambda_5)^3$ , with  $\eta$  as large as  $3/2\pi$ .

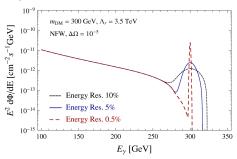


Figure: Continuum photon flux and  $2\gamma$ -line signal from KK-radion annihilation at the center of the galaxy. For the line signal, we show three detector energy resolutions: 10% (black line), 5% (blue line) and 0.5% (red line)



## Conclusions

- We discussed a novel scenario where the warped geometry provides the symmetry for a suitable dark matter candidate which we identify after stabilization of the extra dimension with lightest odd radion mode (the LKP).
- Depending on the value of the decay constant  $\Lambda_r$ , standard freeze out DM ( $\Lambda_r \approx 2$  TeV), coannihilation ( $\Lambda_r \approx 10-100$  TeV or non-thermal ( $\Lambda_r \gg 100$  TeV) scenarios arise.
- In the presence of mixing with the Higgs sector → possible direct detection signals.
- Indirect detection signals that may be probed by FERMI-LAT.