

Gauge Coupling Unification in Heterotic String Models with Gauge Mediated SUSY Breaking

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Work with Prof. Stuart Raby

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Motivation

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In string theories with compactified extra-dimensions, there generically exist extra non-standard model particles, usually called “exotics”.

To mediate SUSY breaking with vector-like “exotic” particles arising from heterotic string theory, and produce a “consistent” low-energy spectrum.

Mini-Landscape Search¹

- Search for MSSM spectrum at low energies starting with $E_8 \times E_8$ heterotic string models compactified on the orbifold, T^6/\mathbb{Z}_6

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- Look for GUTS with the Standard Model Gauge group embedded

$$E_8 \supset E_6 \supset SO(10) \supset SU(5) \supset G_{SM}$$

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- Spectrum: Three families + Vector-like “exotics”
- 15 models with promising phenomenology.

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Gauge Coupling Unification ²

- Gauge Coupling Unification was studied in 2 of these 15 models.

Model 1 and Model 2A

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- This required some of the vector-like exotics to be massive around $10^9 - 10^{13}$ GeV.
- Solutions were constrained by the value of proton lifetime in these models.

$$\tau(p \rightarrow \pi^0 e^+) \gtrsim 10^{34} \text{ yr}^*$$

- * Current bound from Super Kamikande.

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Matter Content and Energy scales

- $\vec{n} = (n_3, n_2, (n_1, n'_1))$ defines the 'light' exotic matter content of the theory.

$$n_3 \times [(3, 1)_{1/3,*} + (\bar{3}, 1)_{-1/3,*}] + n_2 \times [(1, 2)_{0,*} + (1, 2)_{0,*}] + n_1 \times [(1, 1)_{1,*} + (1, 1)_{-1,*}]$$

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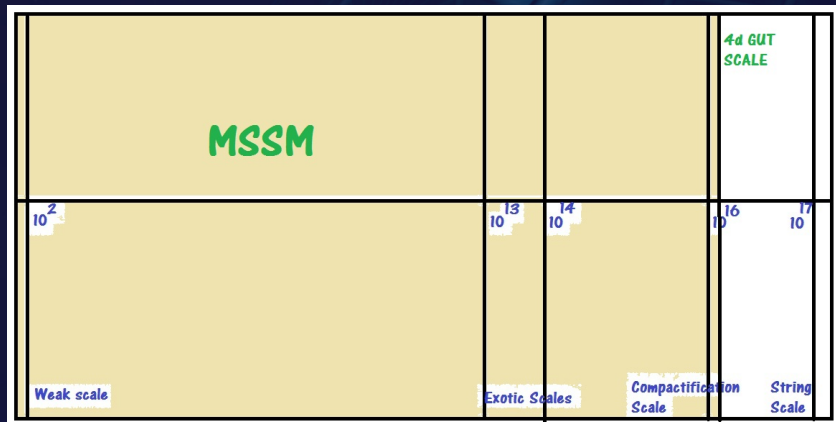
- M_{EX1} - Mass scale of the triplet exotics.
- M_{EX2} - Mass scale of the doublet exotics.
- M_C - The compactification scale of the extra-dimensions.

Heterotic Theory on Orbifold

10^2	10^{13}	10^{14}	10^{16} 10^{17}
Weak scale	Exotic Scales	Compactification Scale	String Scale

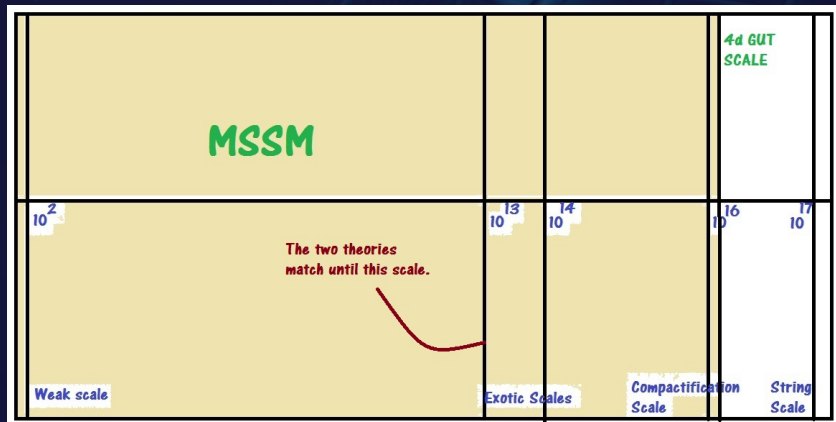
*Figure not drawn to scale.

4D MSSM



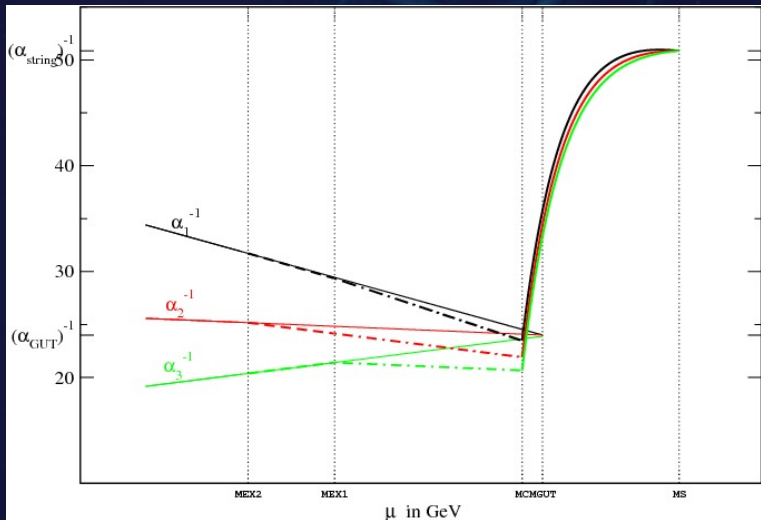
*Figure not drawn to scale.

4D MSSM



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Gauge Coupling Unification



Gaugino Masses

- The gauginos obtain mass at one loop from the exotics:

$$M_i = b_i^{EX3} \frac{\alpha_i}{4\pi} \frac{F\phi}{M_{EX1}} + b_i^{EX2} \frac{\alpha_i}{4\pi} \frac{F\phi}{M_{EX2}}$$

ignoring the sub-dominant contribution from the gravitino.

$$\frac{F\phi}{M_{EX}} \gg m_{3/2}$$

$$b^{EX3} = (n_3, 0, \frac{n_3 + 3n_1}{10}) \quad b^{EX2} = (0, n_2, \frac{3n'_1}{10})$$

Scalar Masses

- The scalars obtain mass at two-loops:

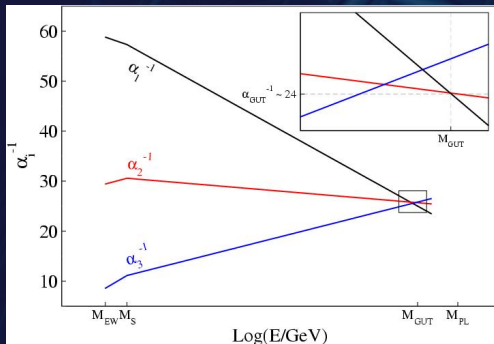
$$m_{\phi_i}^2 = m_{3/2}^2 + 2 \left(b_3^{EX3} \frac{\alpha_3}{4\pi} \frac{F\phi}{M_{EX1}} \right)^2 C_3(i) + 2 \left(b_2^{EX2} \frac{\alpha_2}{4\pi} \frac{F\phi}{M_{EX2}} \right)^2 C_2(i) + 2 \left(\frac{\alpha_1}{4\pi} \left(b_1^{EX3} \frac{F\phi}{M_{EX1}} + b_1^{EX2} \frac{F\phi}{M_{EX2}} \right) \right)^2 C_1(i) + dQ_a^X M_2^2$$

- $dQ_a^X M_2^2$ is a D - term contribution ³ from an anomalous $U(1)_X$ that gives a contribution proportional to GMSB.

³S. Raby and K. Tobe, Nucl. Phys. B539, 3 (1999)

Precision Unification?

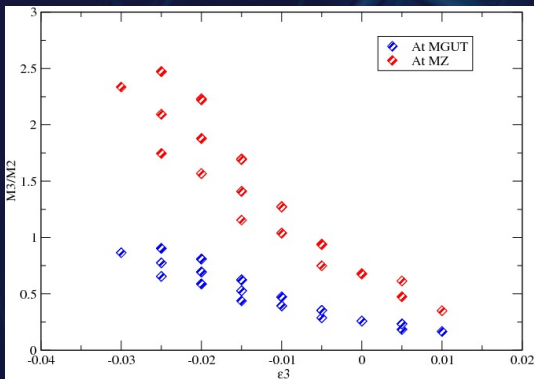
Standard SUSY breaking scenarios require a -4% threshold corrections at the GUT scale.



$$\alpha_i^{-1}(\mu) = \alpha_{GUT}^{-1} + \frac{b_i}{2\pi} \log \frac{M_{GUT}}{\mu} - \alpha_{GUT}^{-1} \frac{\epsilon_3}{(1 + \epsilon_3)} \delta_{i3}$$

Effect of ϵ_3

- We study the effect of threshold corrections on the spectrum of exotics as well as the low energy spectrum.



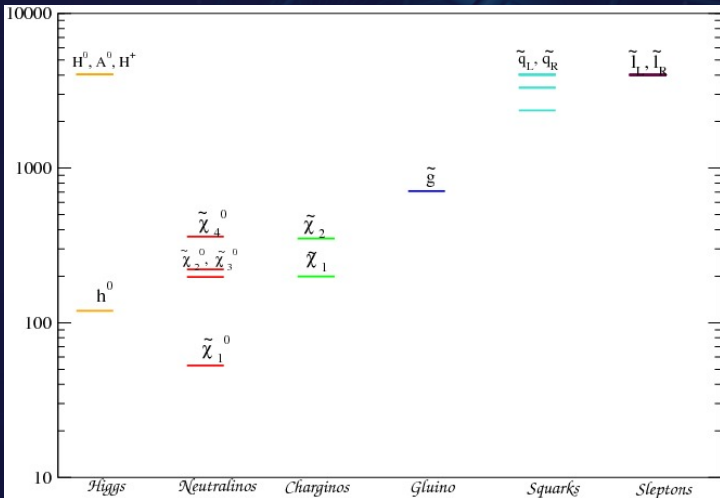
The figure represents the correlation for one particular model with $\vec{n} = (4, 2, (2, 1))$

Two Cases

Observable	Case 1	Case 2
$m_{3/2}$	4 TeV	10 TeV
d	0	5
M_S	6.04×10^{17}	6.05×10^{17}
M_C	1.2×10^{16}	1.2×10^{16}
M_{EX1}	5.03×10^{13}	1.10×10^{14}
M_{EX2}	1.69×10^{13}	8.54×10^{13}
M_{GUT}	2.5×10^{16}	2.0×10^{16}
ϵ_3	-2.5 %	0 %
$\tan \beta$	7	4
mu	-206.217	-1932.930

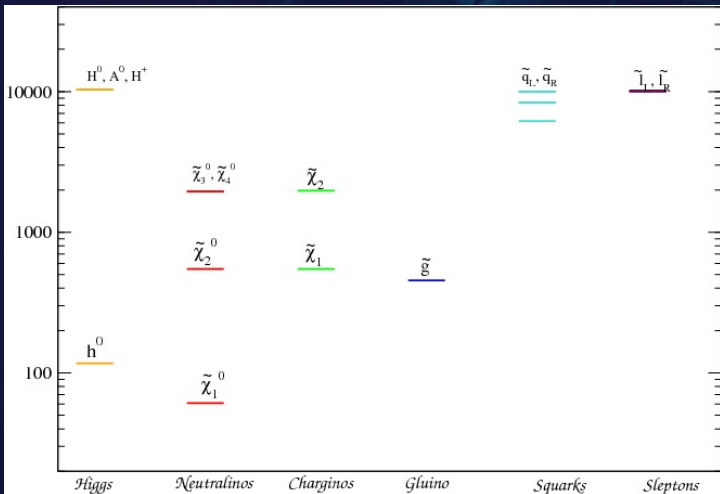
MSSM Spectrum - Case 1

$m_{3/2} = 4 \text{ TeV}$, $d = 0$, $\epsilon_3 = -2.5 \%$



MSSM Spectrum - Case 2

$m_{3/2} = 10$ TeV, $d = 5$, $\epsilon_3 = 0$ %



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- The gaugino masses in the MSSM spectrum depend on the threshold corrections at the GUT scale.
- Interesting Phenomenology - Unique signatures at colliders!!