MATCHING AND MERGING

COMPARING FIXED ORDER AND PARTON SHOWER

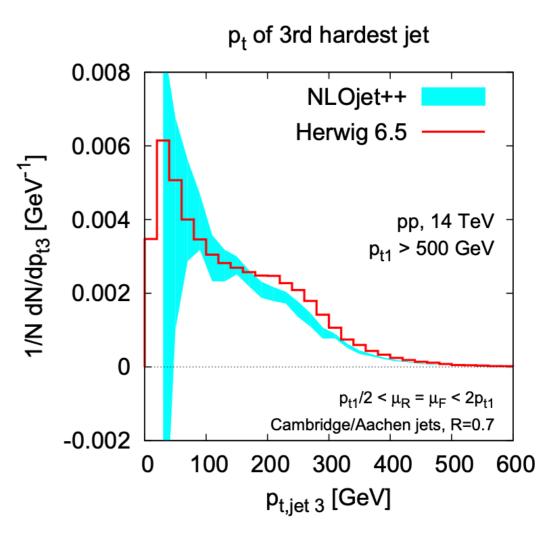
Parton shower	Fixed order
Correct only for soft/ collinear radiation	Hard radiation correctly described
High multiplicity final states possible	At most ~10 particles in final state
Realistic, hadronic final states	Only partonic final states
Hard to improve accuracy	Known how to systematically improve accuracy

BRIEF ASIDE ON JETS

- Fixed order calculations describe single partons, but they are not what is observed in a detector
- Soft-collinear radiation causes a 'spray' of particles from each initiating parton
- The spray forms a cone-like shape of radiation, called a jet
- Many different ways to define jets, which must be IR-safe definitions use a jet radius *R* which quantifies jet size
- Jet algorithms take partons and cluster them into jets

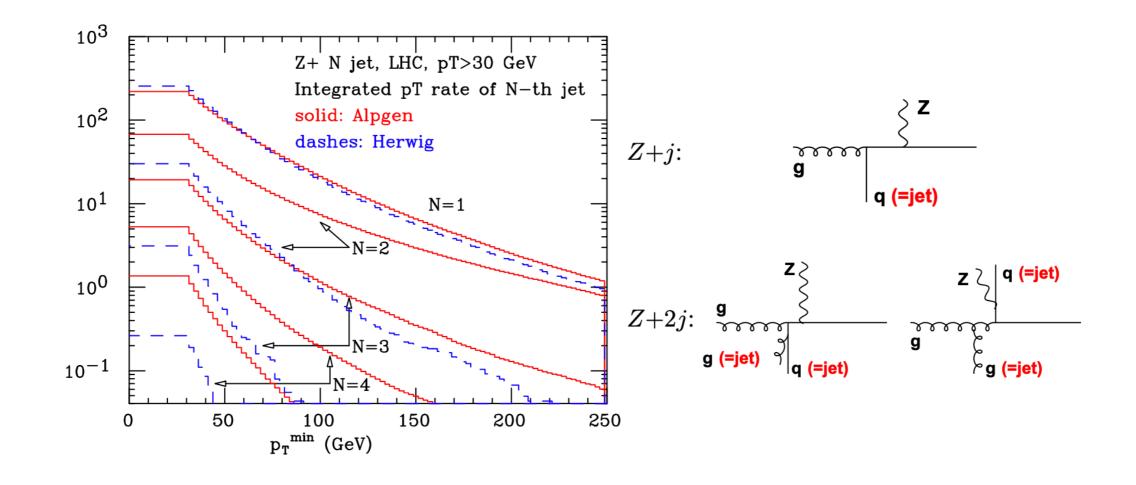
JET PRODUCTION

- Tree level processes are qq → qq,
 qg → qg etc.
- Parton shower starts with these and adds extra radiation
- Only correct in soft/collinear limit, but sometimes adds hard extra emissions
- Pretty good from PS large uncertainties on FO at small values as large logs blow up



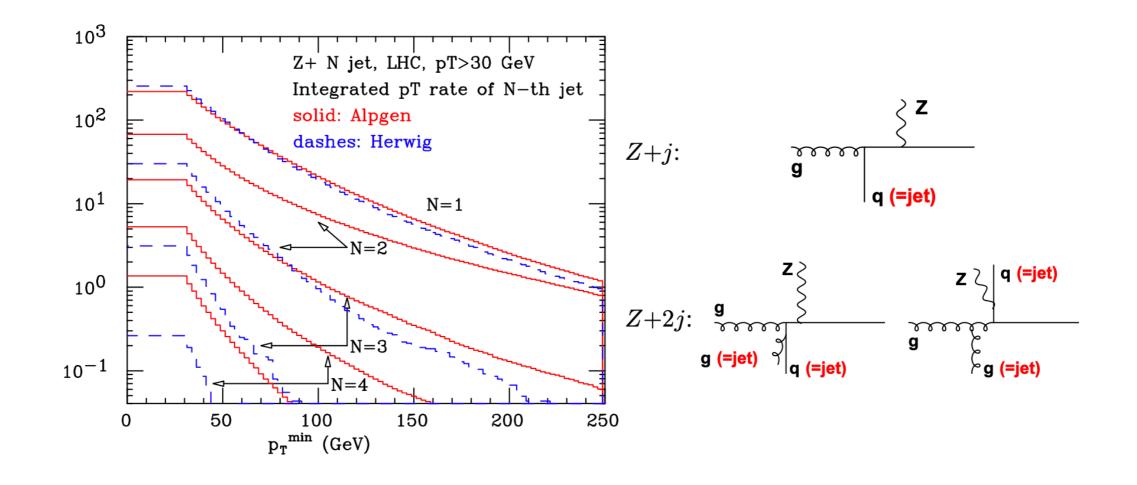
VECTOR BOSON + JET PRODUCTION

- Fig. shows cross section for Nth jet to have transverse energy above E_T
- PS and FO in agreement for 1st jet, but terrible for >2



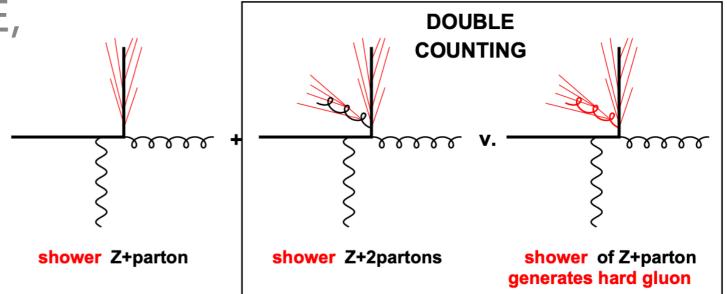
VECTOR BOSON + JET PRODUCTION

- Explanation: HERWIG generates hard Z + j configs
- But: also soft/coll. enhanced events where Z is radiated off a dijet config, not captured by QCD shower alone



MATRIX ELEMENTS WITH PARTON SHOWERS (MEPS)

- How do we combine PS and FO?
- Consider Z + j production as the underlying hard process. PS does a good job for Z + 1 parton, but terribly for more
- Naïve solution: generate
 Z + 2 with correct LO ME,
 then shower
- Problem: double counting!



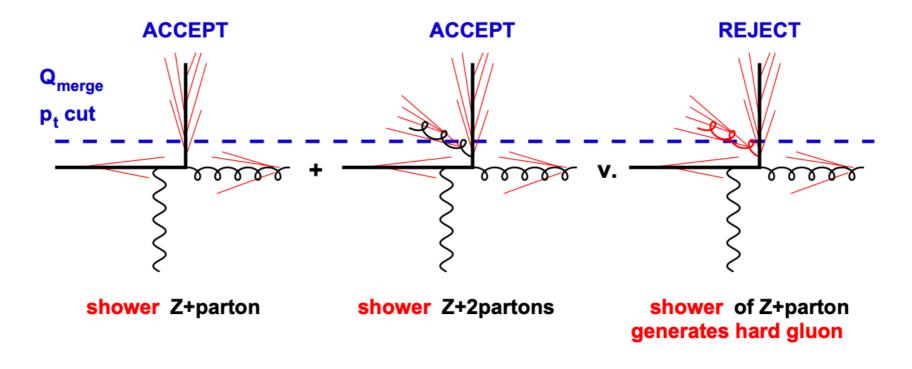
Credit: G. Salam

MLM MERGING

- > Introduce transverse momentum cutoff $Q_{\rm ME}$ and angular cutoff $R_{\rm ME}$
- Generate tree-level events for up to Z + N partons, where all partons have $p_T > Q_{\rm ME}$ and separation $\theta > R_{\rm ME}$
- Shower all events
- Apply a jet algorithm to the showered event with $R > R_{\rm ME}$, and identify all jets with $p_T > Q_{\rm merge} \gtrsim Q_{\rm ME}$
- If each jet corresponds to one of the partons and there are no extra jets above Q_{merge} then accept, otherwise reject

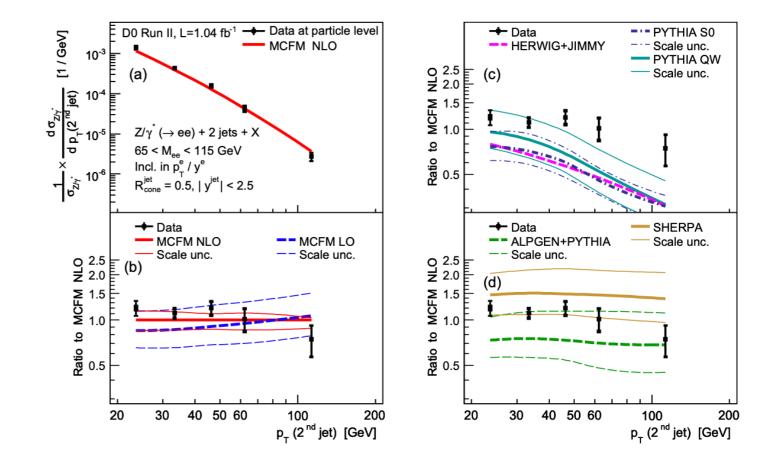
MLM MERGING

- Double-counting removed by rejection of hard radiation
- Hard jets come only from the matrix element



MLM VS FIXED ORDER AND PARTON SHOWER

- MLM (green) gets shape right
- Large scale uncertainty and normalisation wrong, much worse than NLO (red)



 Ideally, need a way to combine NLO calculations with the parton shower (or even NNLO)

MATCHING NLO TO PARTON SHOWER

- Criteria for a successful combination of NLO+PS:
 - Total cross section inherited from NLO
 - Radiation pattern (first order) follows NLO real emission
 - Logarithmic accuracy of PS is maintained
- Recall NLO structure:

$$\sigma_{N}^{\text{NLO}} = \int d\Phi_{\mathscr{B}} \left[\mathscr{B}_{N}(\Phi_{\mathscr{B}}) + \mathscr{V}_{N}(\Phi_{\mathscr{B}}) + \mathscr{I}_{N}^{\mathscr{S}}(\Phi_{\mathscr{B}}) \right] + \int d\Phi_{\mathscr{R}} \left[\mathscr{R}_{N}(\Phi_{\mathscr{R}}) - \mathscr{S}_{N}(\Phi_{\mathscr{R}}) \right]$$

IMPROVING THE PARTON SHOWER – MATRIX ELEMENT CORRECTIONS

- Parton shower good for soft/collinear, bad for hard emissions
- Can we correct it to get the hardest emission right?
- In many processes, parton shower is an overestimate of exact ME:

$$\mathscr{R}_{N}(\Phi_{\mathscr{B}} \otimes \Phi_{1}) \leq \mathscr{B}_{N}(\Phi_{\mathscr{B}}) \otimes \mathscr{K}_{N}(\Phi_{1})$$

> \mathcal{K} is PS soft and collinear splitting kernel (we discussed $P_{ij'}$ the collinear case only)

MATRIX ELEMENT CORRECTIONS

First emission pattern looks like:

$$d\sigma_{N} = d\Phi_{\mathscr{B}} \mathscr{B}_{N}(\Phi_{\mathscr{B}}) \left\{ \Delta_{N}(\mu_{Q}^{2}, t_{c}) + \int_{t_{c}}^{\mu_{Q}^{2}} d\Phi_{1} \left[\mathscr{K}_{N}(\Phi_{1}) \Delta_{N}(\mu_{Q}^{2}, t(\Phi_{1})) \right] \right\}$$

No emission probability Single emission probability at a given time t

- Terms in curly brackets integrate to 1 (shower is unitary)
- Let's modify the splitting kernel to make it look more like the real matrix element, at least for the first emission:

$$\tilde{\mathcal{K}}_{N}(\Phi_{1}) = \mathcal{R}_{N}(\Phi_{\mathcal{B}} \otimes \Phi_{1}) / \mathcal{B}_{N}(\Phi_{\mathcal{B}})$$

MATRIX ELEMENT CORRECTIONS

First emission pattern modified to:

$$\mathrm{d}\sigma_{N} = \mathrm{d}\Phi_{\mathscr{B}}\mathcal{B}_{N}(\Phi_{\mathscr{B}}) \left\{ \tilde{\Delta}_{N}(\mu_{Q}^{2}, t_{c}) + \int_{t_{c}}^{\mu_{Q}^{2}} \mathrm{d}\Phi_{1} \left[\frac{\mathscr{R}_{N}(\Phi_{\mathscr{B}} \otimes \Phi_{1})}{\mathscr{B}_{N}(\Phi_{\mathscr{B}})} \tilde{\Delta}_{N}(\mu_{Q}^{2}, t(\Phi_{1})) \right] \right\}$$

- Now first emission follows real matrix element!
- Practically, use normal shower kernels and simply accept/ reject points with a probability

$$\mathcal{P}_{\text{MEC}} = \frac{\mathcal{R}_{N}(\Phi_{\mathcal{B}} \otimes \Phi_{1})}{\mathcal{B}(\Phi_{\mathcal{B}}) \otimes \mathcal{K}_{N}(\Phi_{1})}$$

NLO MATCHING – THE POWHEG METHOD

Define Born-like configurations which give NLO-accurate cross section:

$$\overline{\mathscr{B}}_{N}(\Phi_{\mathscr{B}}) = \mathscr{B}_{N}(\Phi_{\mathscr{B}}) + \overline{\mathscr{V}}_{N}(\Phi_{\mathscr{B}}) + \int d\Phi_{1} \left[\mathscr{R}_{N}(\Phi_{B} \otimes \Phi_{1}) - \mathscr{S}_{N}(\Phi_{B} \otimes \Phi_{1}) \right]$$

IR-subtracted, UV-renormalised virtual piece is

$$\overline{\mathcal{V}}_{N}(\Phi_{\mathscr{B}}) = \mathcal{V}_{N}(\Phi_{\mathscr{B}}) + \mathcal{I}_{N}^{\mathscr{S}}(\Phi_{\mathscr{B}})$$

• Works if $\Phi_{\mathscr{R}} = \Phi_{\mathscr{B}} \otimes \Phi_1$. $\overline{\mathscr{B}}$ terms are fully differential cross sections of Born configurations with NLO weight.

NLO MATCHING – THE POWHEG METHOD

- Unitary PS cannot spoil NLO cross section
- Still need pattern of first emission to be correct up to $\mathcal{O}(\alpha_s)$
- Get this by applying matrix element corrections!
- POWHEG formula given by

$$\mathrm{d}\sigma_{N} = \mathrm{d}\Phi_{\mathscr{B}}\overline{\mathscr{B}}_{N}(\Phi_{\mathscr{B}}) \left\{ \tilde{\Delta}_{N}(\mu_{Q}^{2}, t_{c}) + \int_{t_{c}}^{\mu_{Q}^{2}} \mathrm{d}\Phi_{1} \left[\frac{\mathscr{R}_{N}(\Phi_{\mathscr{B}} \otimes \Phi_{1})}{\mathscr{B}_{N}(\Phi_{\mathscr{B}})} \tilde{\Delta}_{N}(\mu_{Q}^{2}, t(\Phi_{1})) \right] \right\}$$

NLO MATCHING – THE POWHEG METHOD

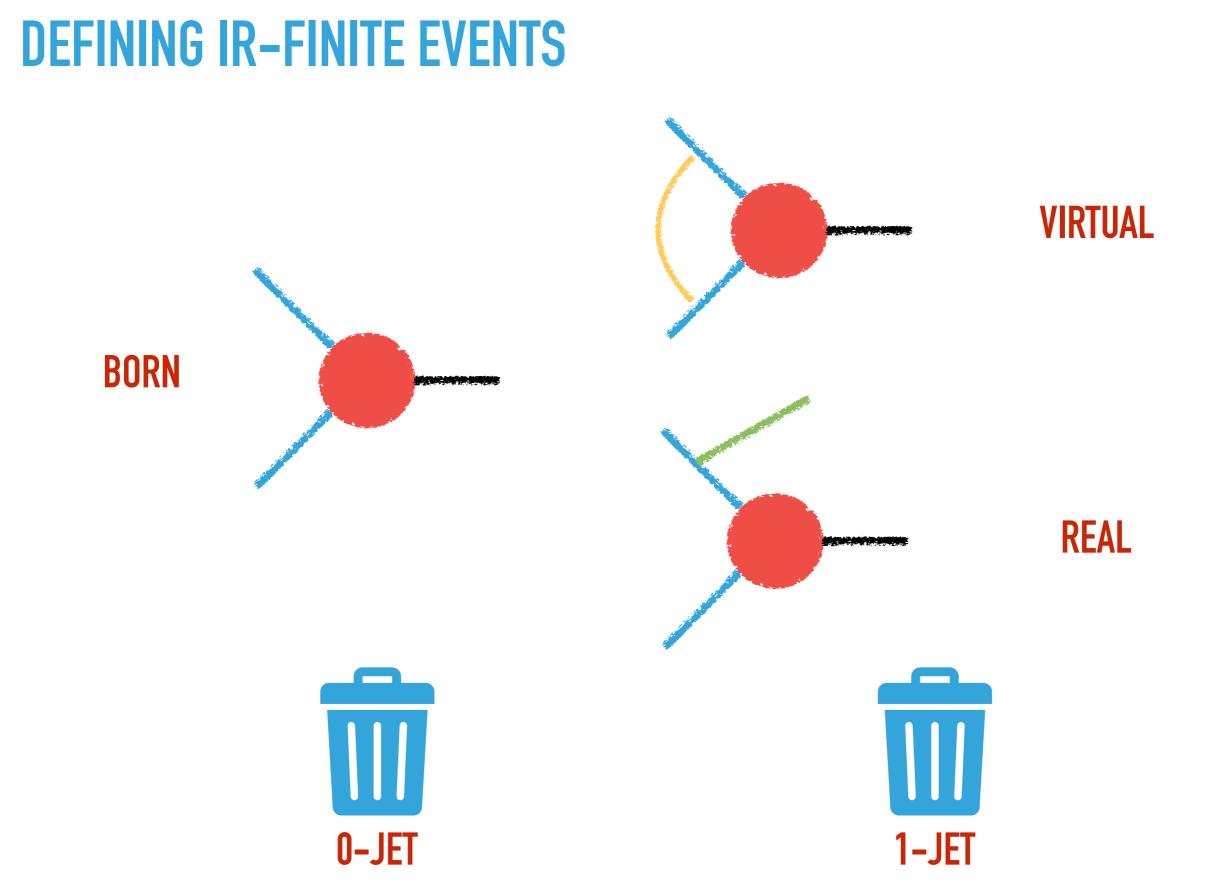
POWHEG formula given by

$$\mathrm{d}\sigma_{N} = \mathrm{d}\Phi_{\mathscr{B}}\overline{\mathscr{B}}_{N}(\Phi_{\mathscr{B}}) \left\{ \tilde{\Delta}_{N}(\mu_{Q}^{2}, t_{c}) + \int_{t_{c}}^{\mu_{Q}^{2}} \mathrm{d}\Phi_{1} \left[\frac{\mathscr{R}_{N}(\Phi_{\mathscr{B}} \otimes \Phi_{1})}{\mathscr{B}_{N}(\Phi_{\mathscr{B}})} \tilde{\Delta}_{N}(\mu_{Q}^{2}, t(\Phi_{1})) \right] \right\}$$

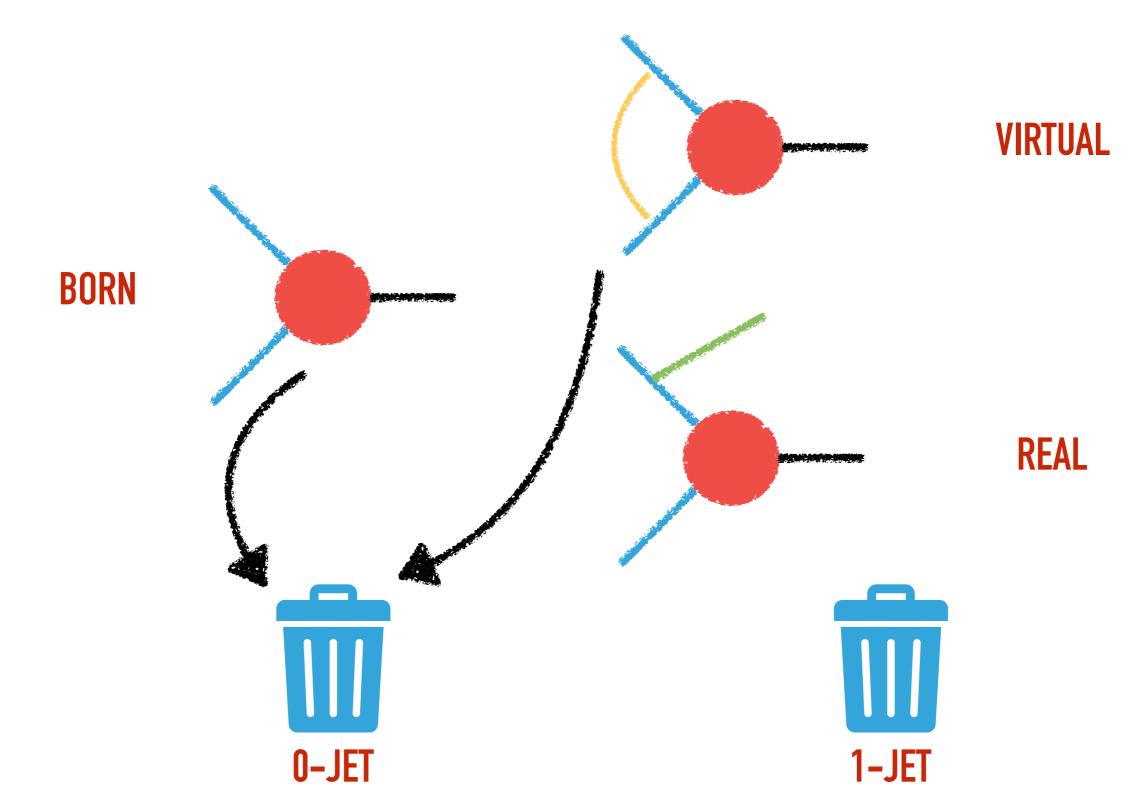
- Gets NLO cross section right (term in curly braces integrates to unity)
- Gets real radiation right at $\mathcal{O}(\alpha_s)$ NLO terms in $\overline{\mathscr{B}}$ hitting $\mathscr{R}_N/\mathscr{B}_N$ are $\mathcal{O}(\alpha_s^2)$
- Subtleties in scale choices, starting scale of PS

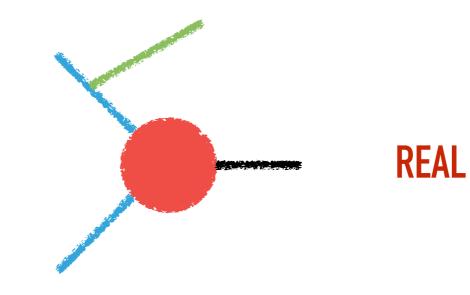
MATCHING TO NNLO

- LO gives order of magnitude estimate, NLO is reliable, but need NNLO for precision.
- NNLO calculations are much harder than NLO.
- Many overlapping divergences up to 2 extra emissions, can be soft and/or collinear in different combinations
- Cancellation of divergences between real and virtual diagrams is still guaranteed by the KLN theorem
- Let's take a step back how do we define an 'event' which is IRfinite?



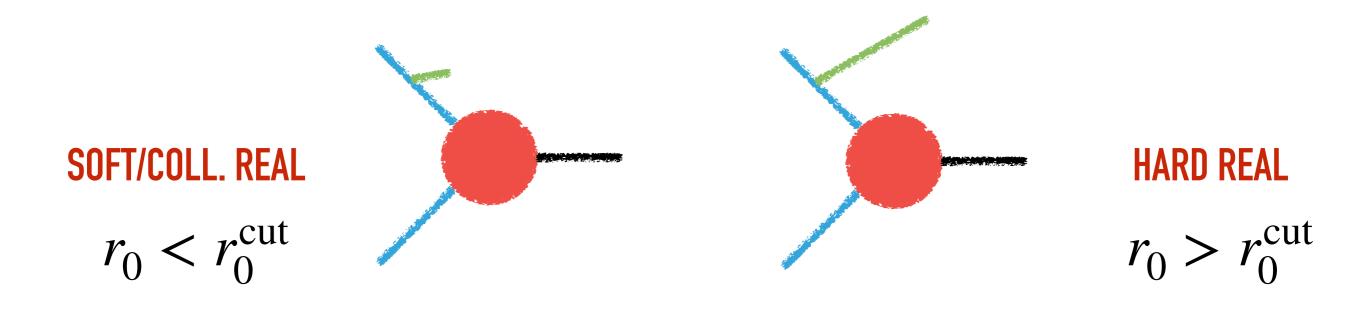






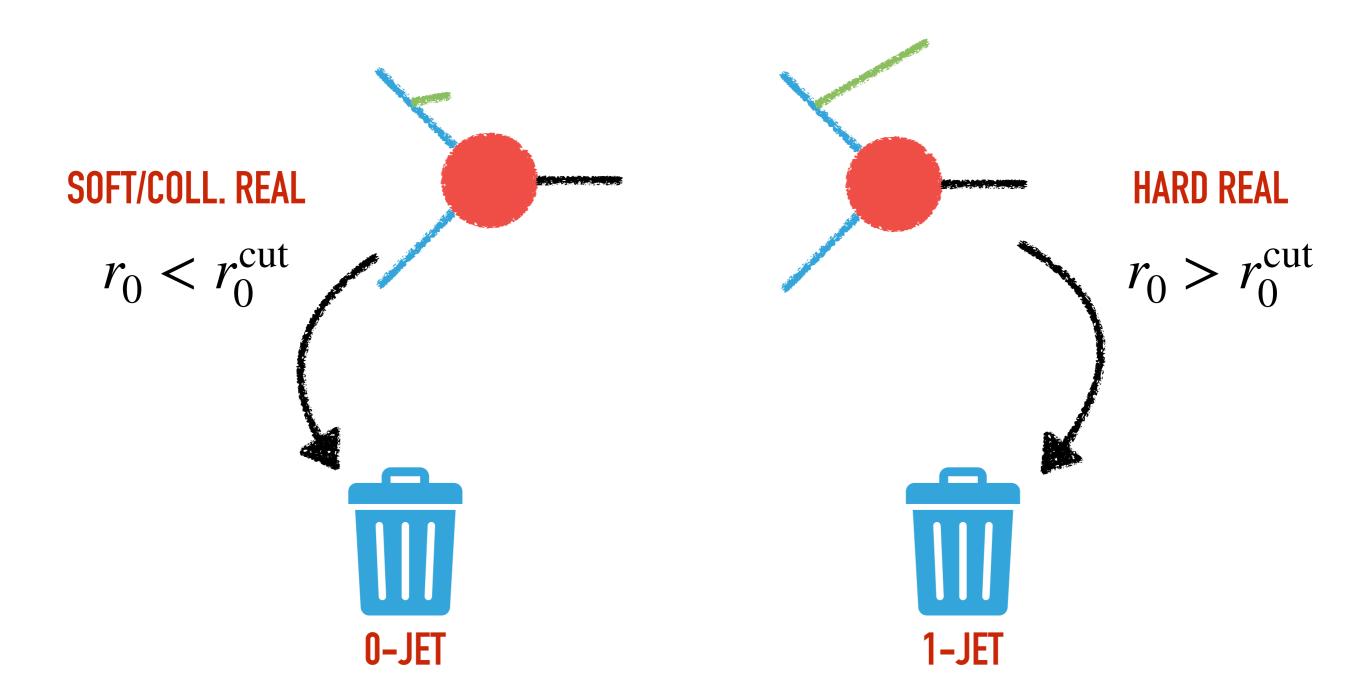


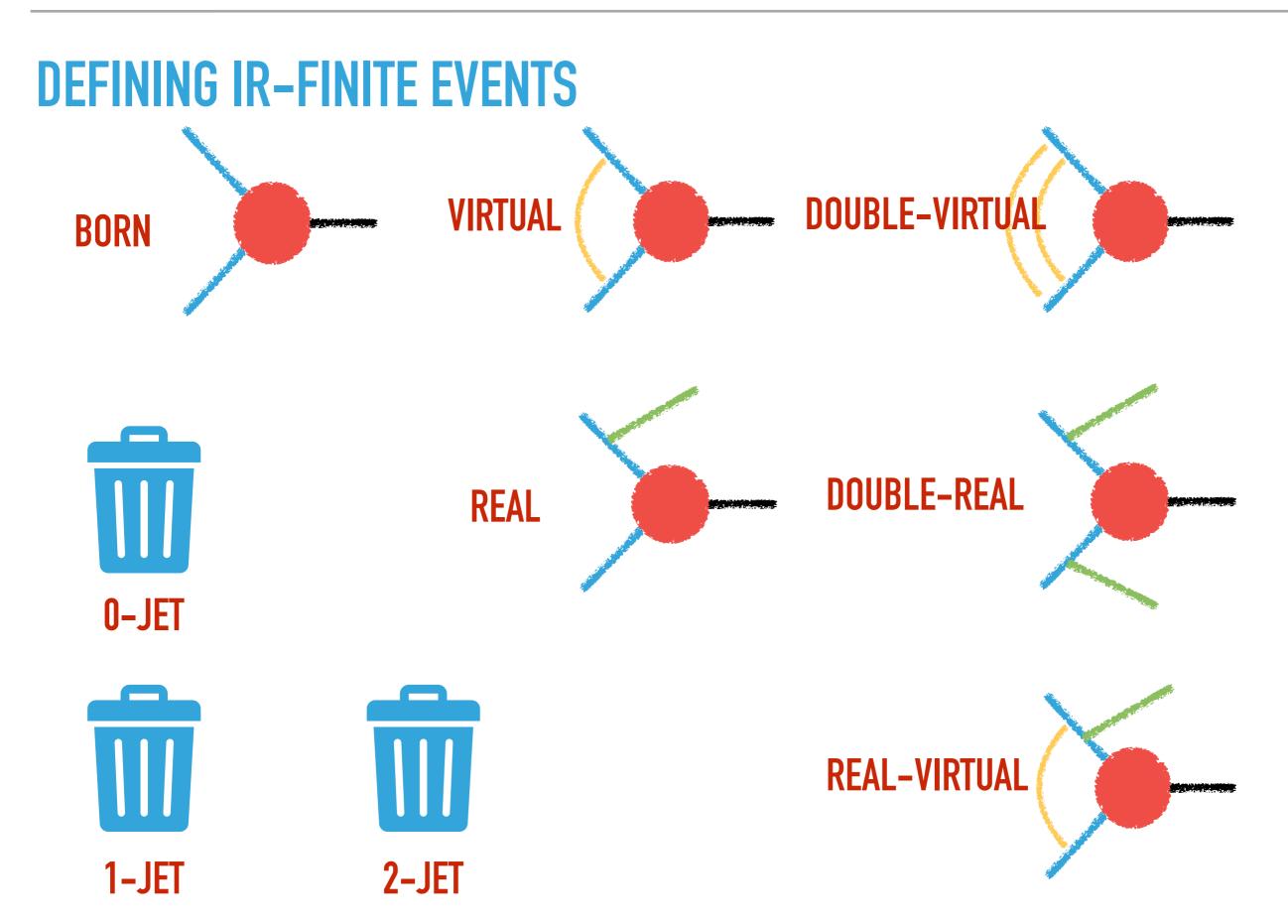


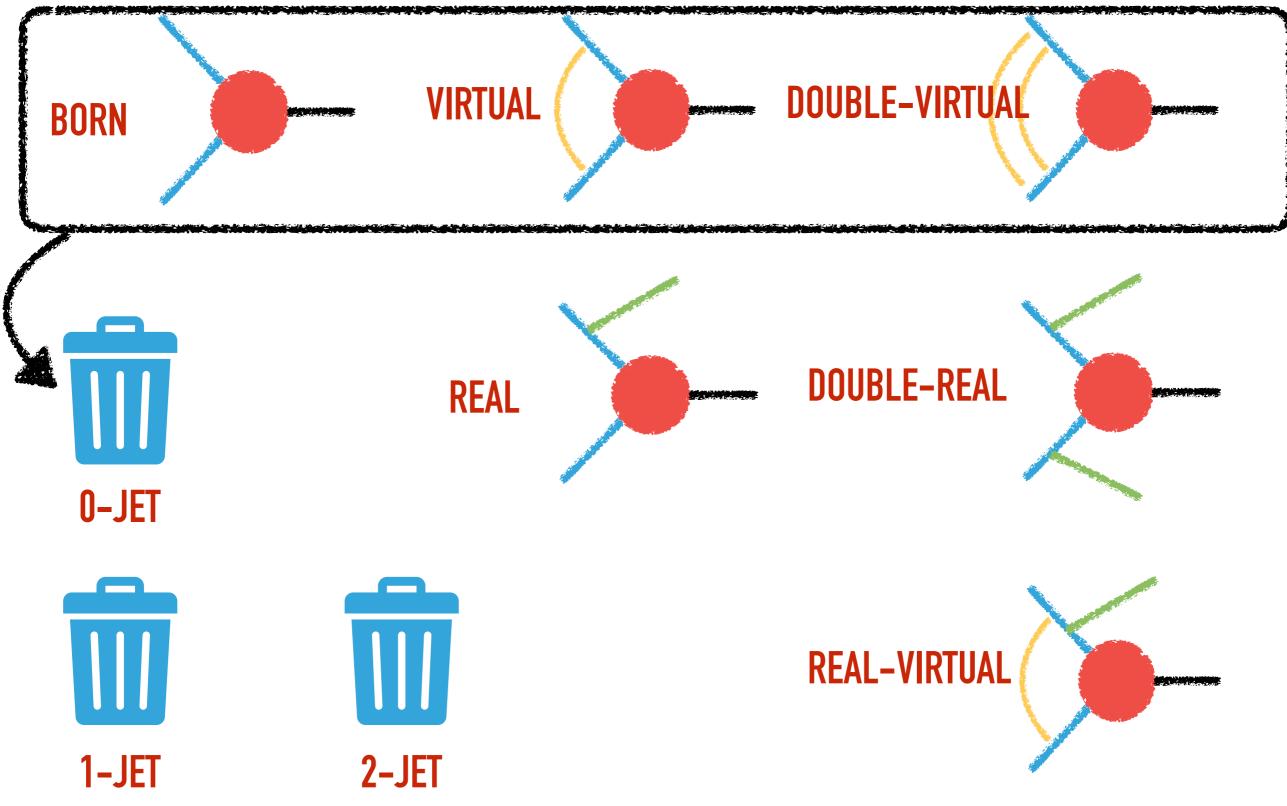


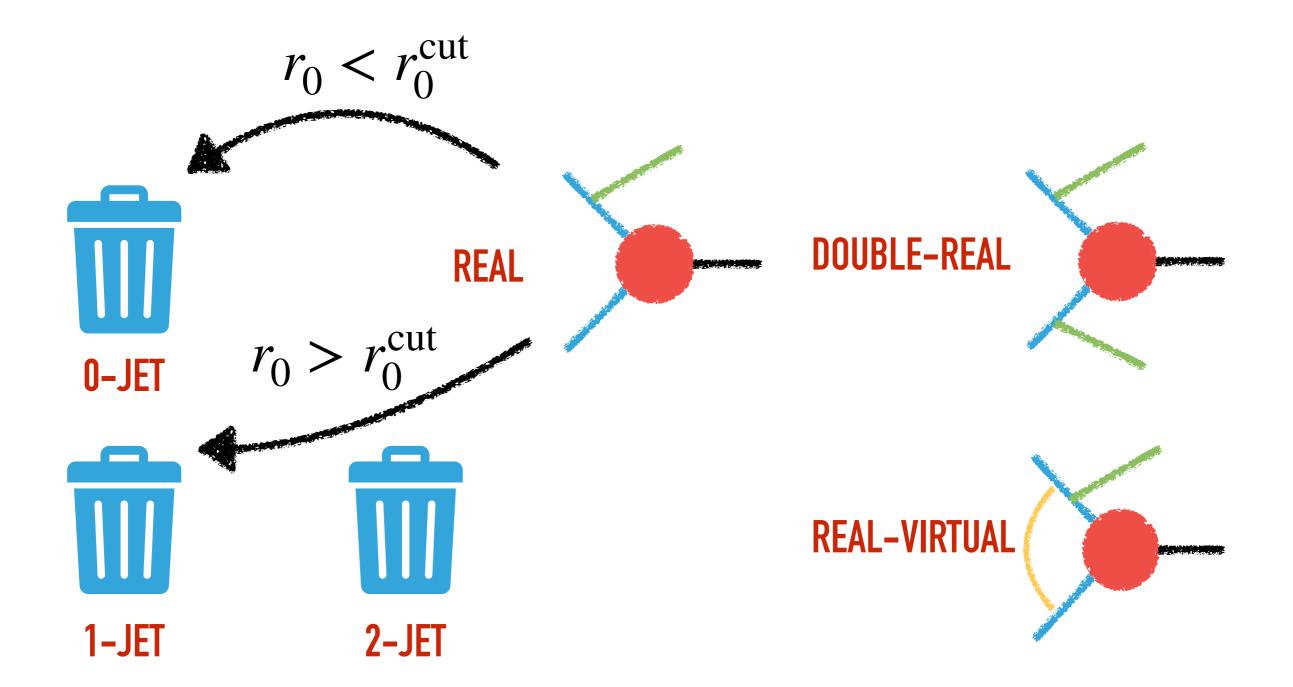


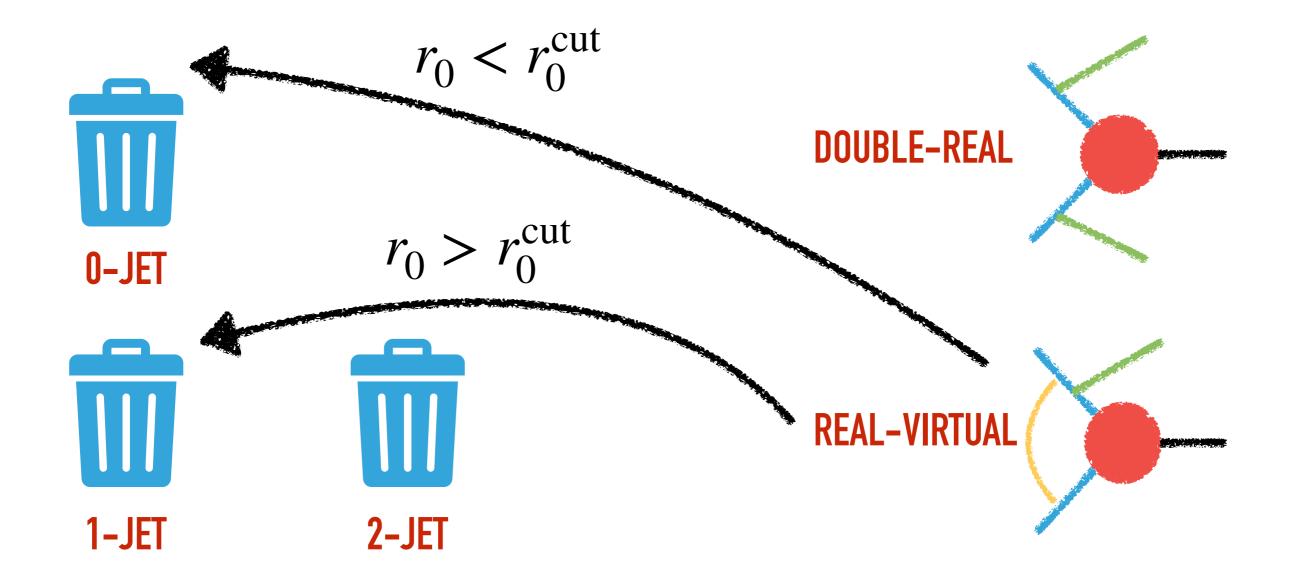


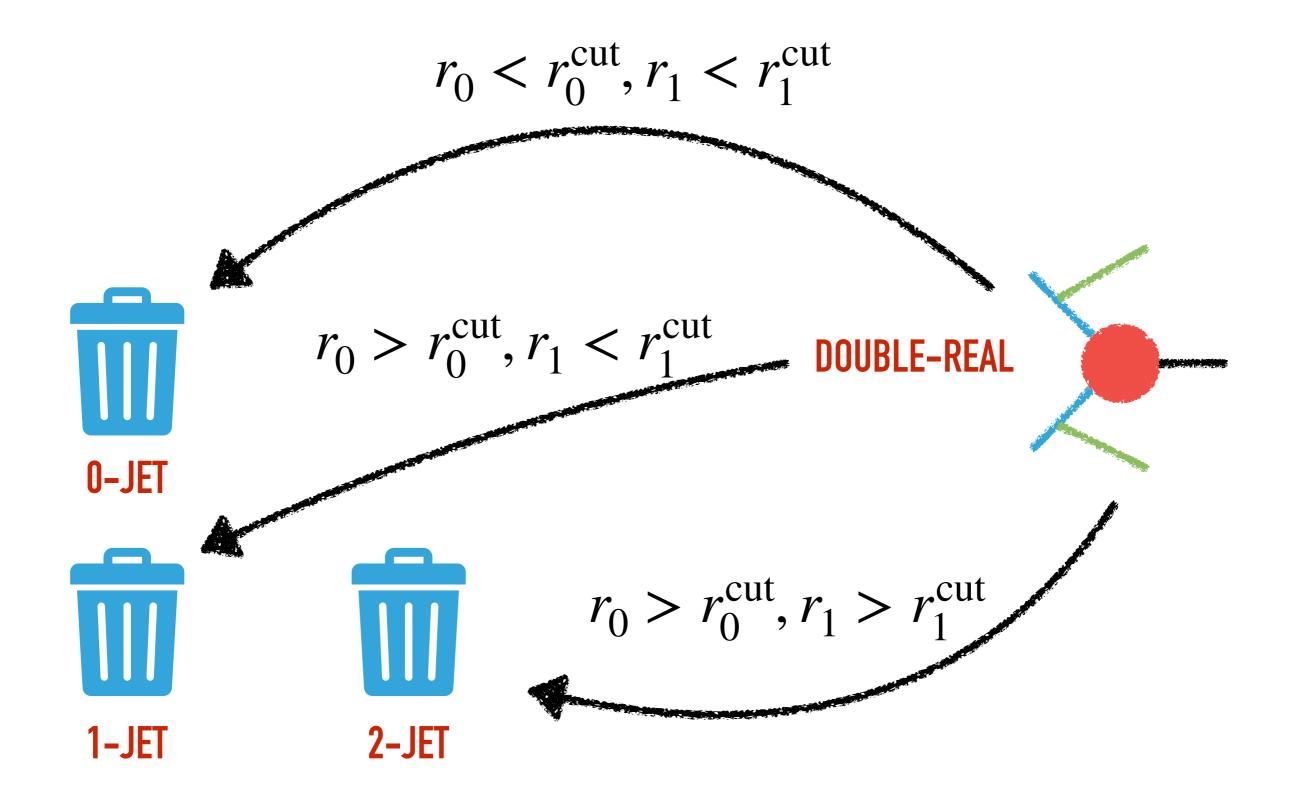










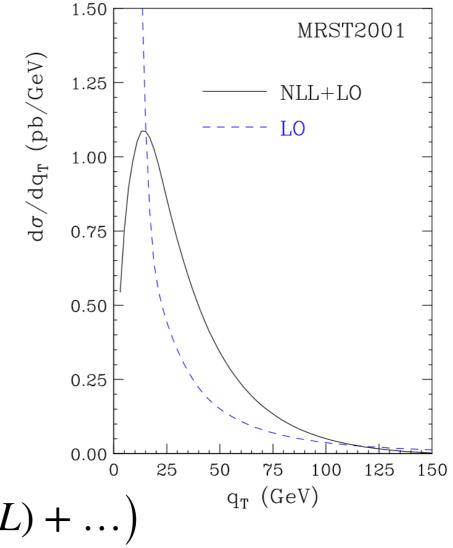


- Defining events this way introduced a projection from a higher multiplicity to a lower multiplicity phase space
- Results are only (N)NLO accurate up to power corrections in r_0^{cut} - as $r_0^{\text{cut}} \rightarrow 0$, exact fixed order result is recovered
- Causes large logarithms to appear which spoil perturbative convergence!

 $L = \log(Q/r_0^{\text{cut}})$ becomes large...

RESUMMATION – THE CURE FOR LARGE LOGS

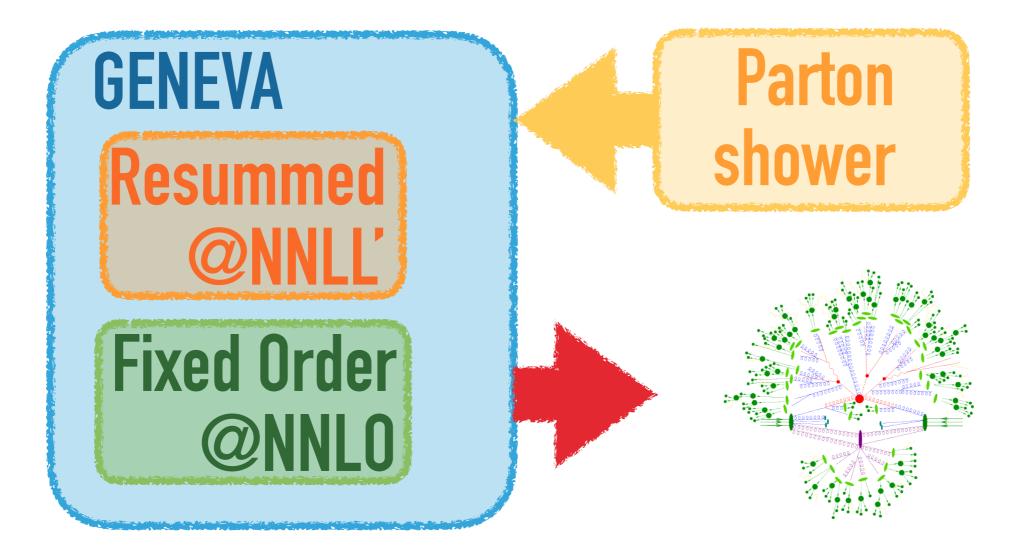
- Large logs signal the breakdown of the perturbative series in the coupling, leading term $\alpha L^2 \sim 1 \Rightarrow \alpha L \ll 1$
- Reordering the series to expand in a genuinely small parameter cures behaviour



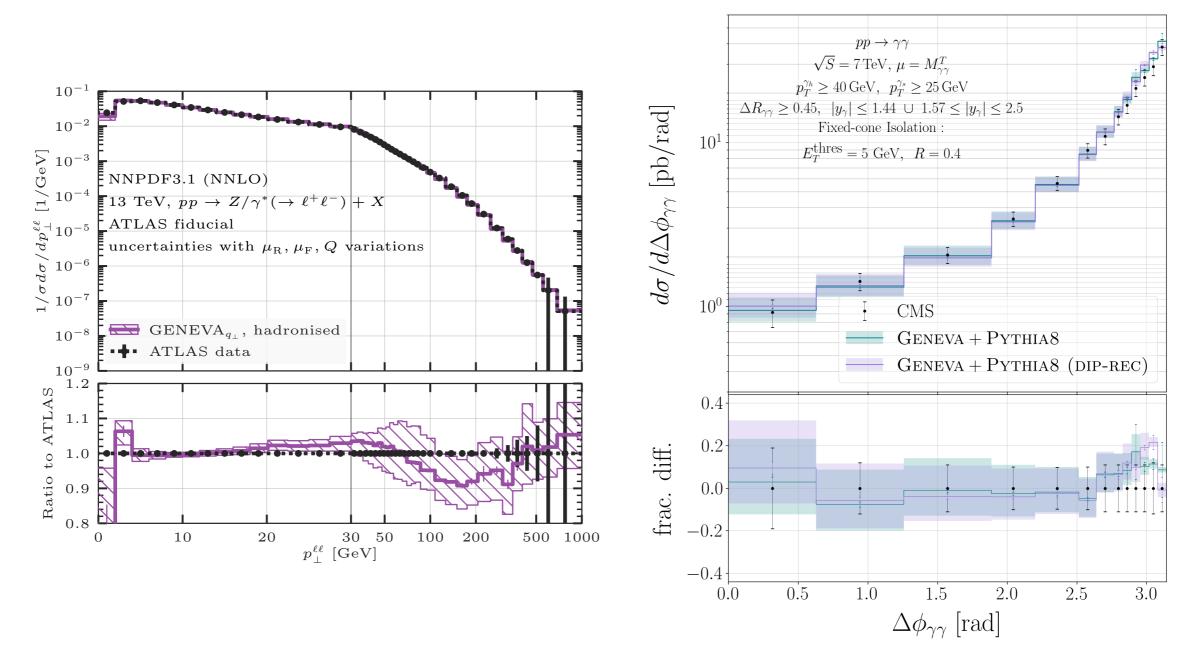
 $d\sigma = C(\alpha_s) \exp \left(Lg_1(\alpha_s L) + g_2(\alpha_s L) + \alpha_s g_3(\alpha_s L) + \dots \right)$

Different formalisms available to achieve this

COMBINING RESUMMED AND FIXED ORDER CALCULATIONS IN GENEVA

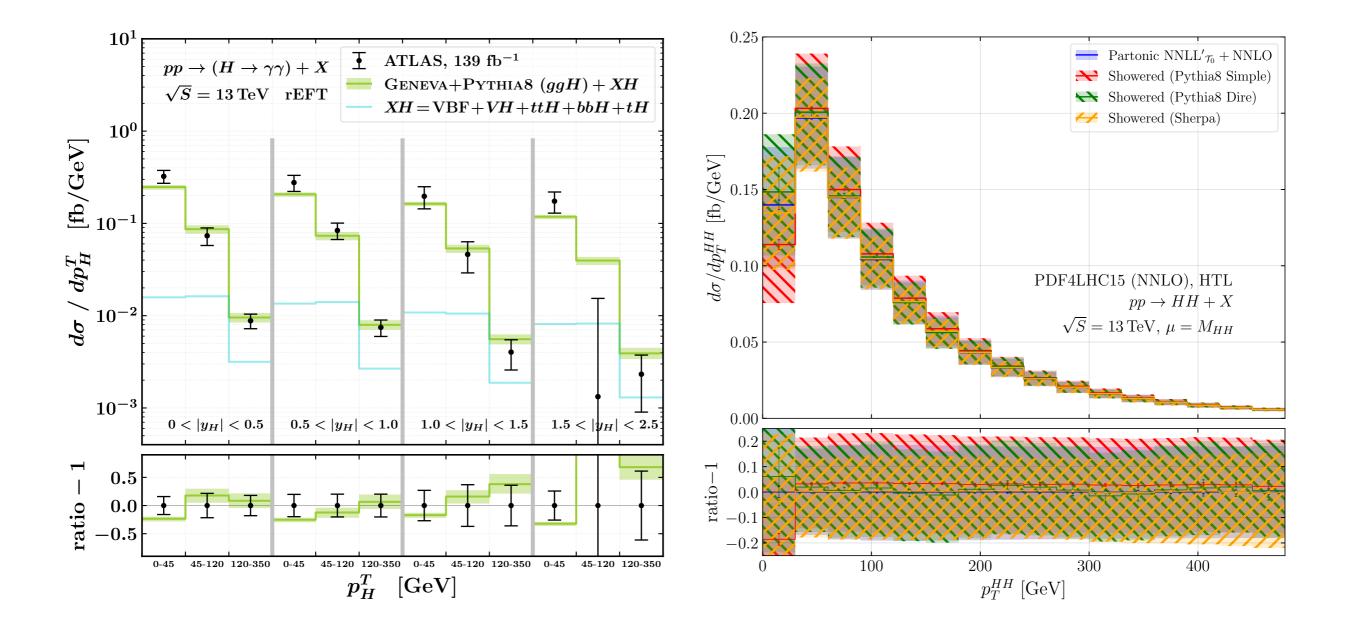


GENEVA AS AN NNLO+PS EVENT GENERATOR



2010.10498, 2102.08390, S. Alioli, C. Bauer, A. Broggio, A. Gavardi, S. Kallweit, MAL, R. Nagar, D. Napoletano, L. Rottoli

GENEVA AS AN NNLO+PS EVENT GENERATOR



2212.10489, 2301.11875, S. Alioli, G. Billis, A. Broggio, A. Gavardi, S. Kallweit, MAL, G. Marinelli, R. Nagar, D. Napoletano

SUMMARY

- Fixed order and parton shower calculations have different advantages - important to be able to combine them to achieve best theoretical description
- Merging combines samples with different multiplicities at FO and showers them without double counting
- Matching corrects first emissions of parton shower to be (N)NLO accurate and gives events with (N)NLO weight

OVERVIEW

- I have not been exhaustive by a long shot many topics uncovered and details omitted, see initial references for more info
- Aim has not been to bring you up-to-speed with cutting edge developments or list all available tools, but to peek inside the black-box
- Hopefully now you appreciate the power and limitations of event generators, and can debug more successfully!