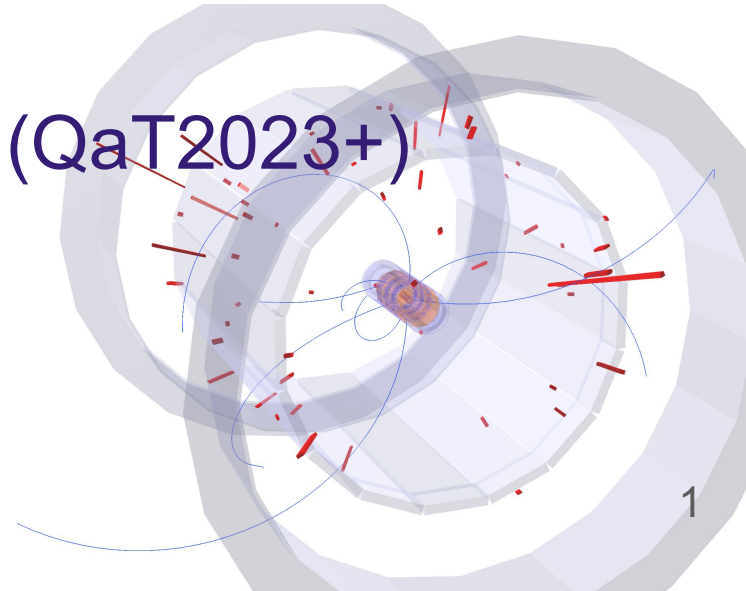


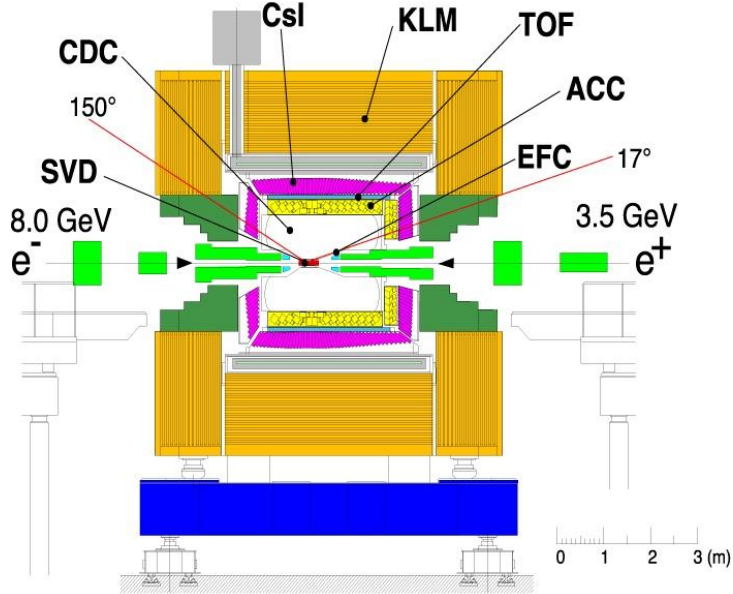
# Quarkonium measurements at Belle-II

Quarkonia as Tools 2023 (QaT2023+)

Pavel Oskin  
(IJCLab)

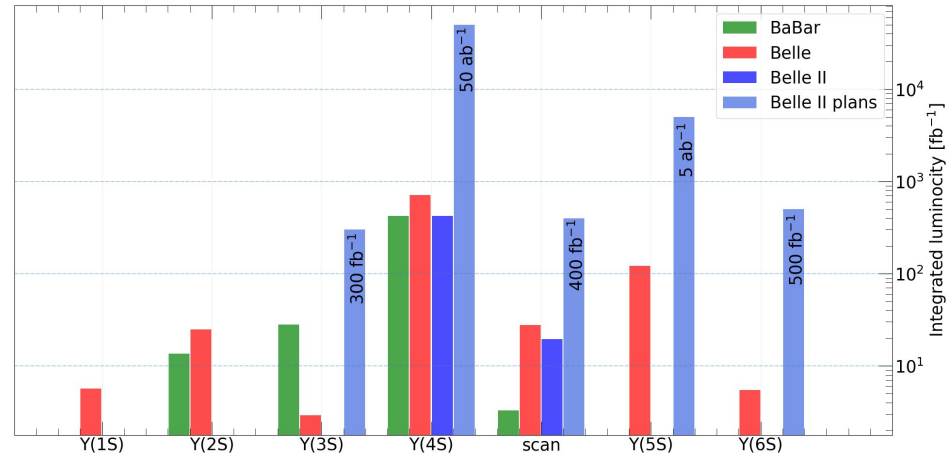
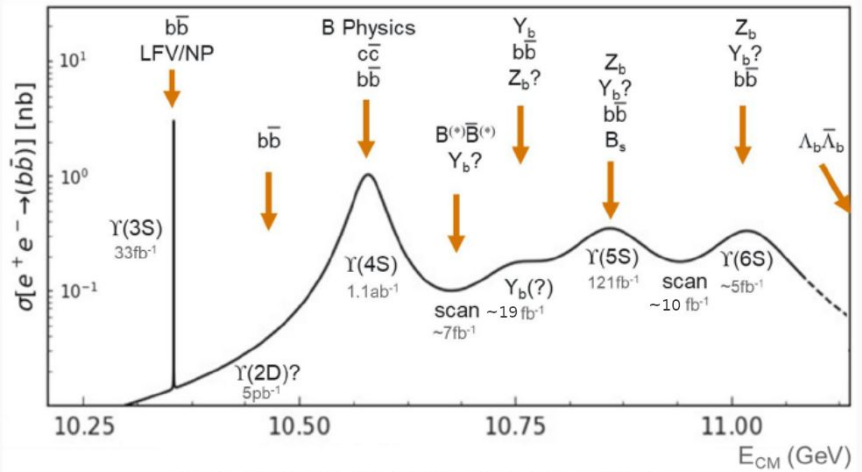


# KEKB and Belle



Collected data from 1999 to 2010:

- $121 \text{ fb}^{-1}$  at  $Y(5S) \sim 7.11 \times 10^6 B_s B_s$
- $711 \text{ fb}^{-1}$  at  $Y(4S) \sim 771 \times 10^6 BB$
- $3 \text{ fb}^{-1}$  at  $Y(3S)$
- $24 \text{ fb}^{-1}$  at  $Y(2S)$
- $6 \text{ fb}^{-1}$  at  $Y(1S)$
- $26 \text{ fb}^{-1}$  scan above  $Y(4S)$



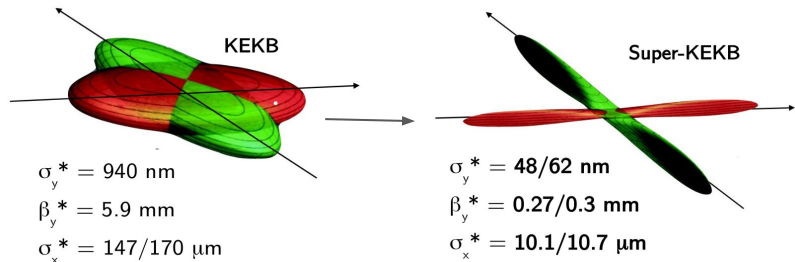
# SuperKEKB and Belle II

Beam current increased by **x2**

$$L = \frac{\gamma_{\pm}}{2er_e} \left( 1 + \frac{\sigma_y^*}{\sigma_x^*} \right) \left( \frac{I_{\pm} \xi_{y\pm}}{\beta_y^*} \right) \left( \frac{R_L}{R_{\xi_y}} \right)$$

Vertical beta function at IP reduced by **1/20**  
"Nano-beam" scheme

**x40** instant luminosity increase

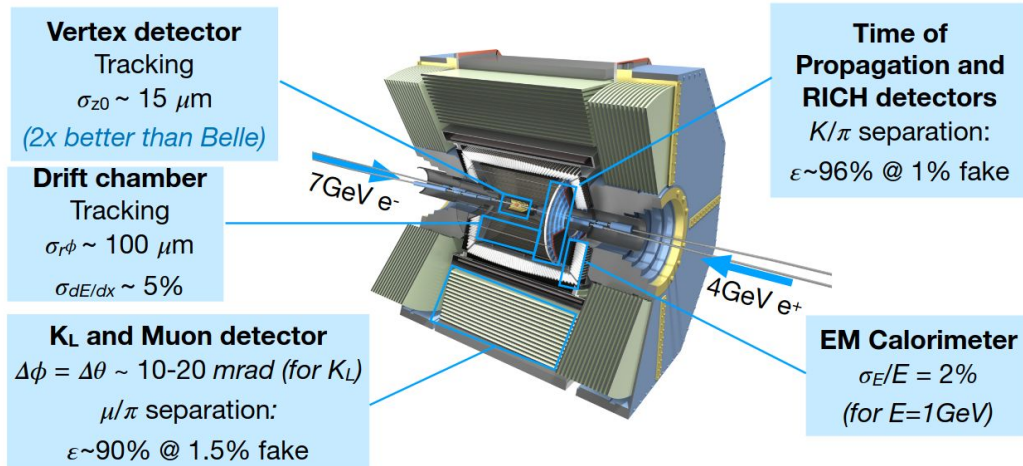
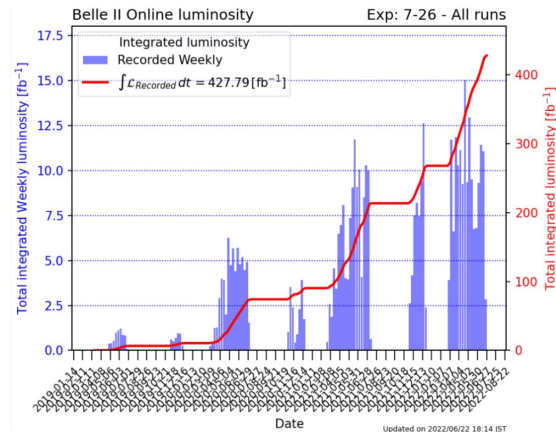


Collected data:

- $424 \text{ fb}^{-1}$  at Y(4S)
- $19 \text{ fb}^{-1}$  scan above Y(4S)

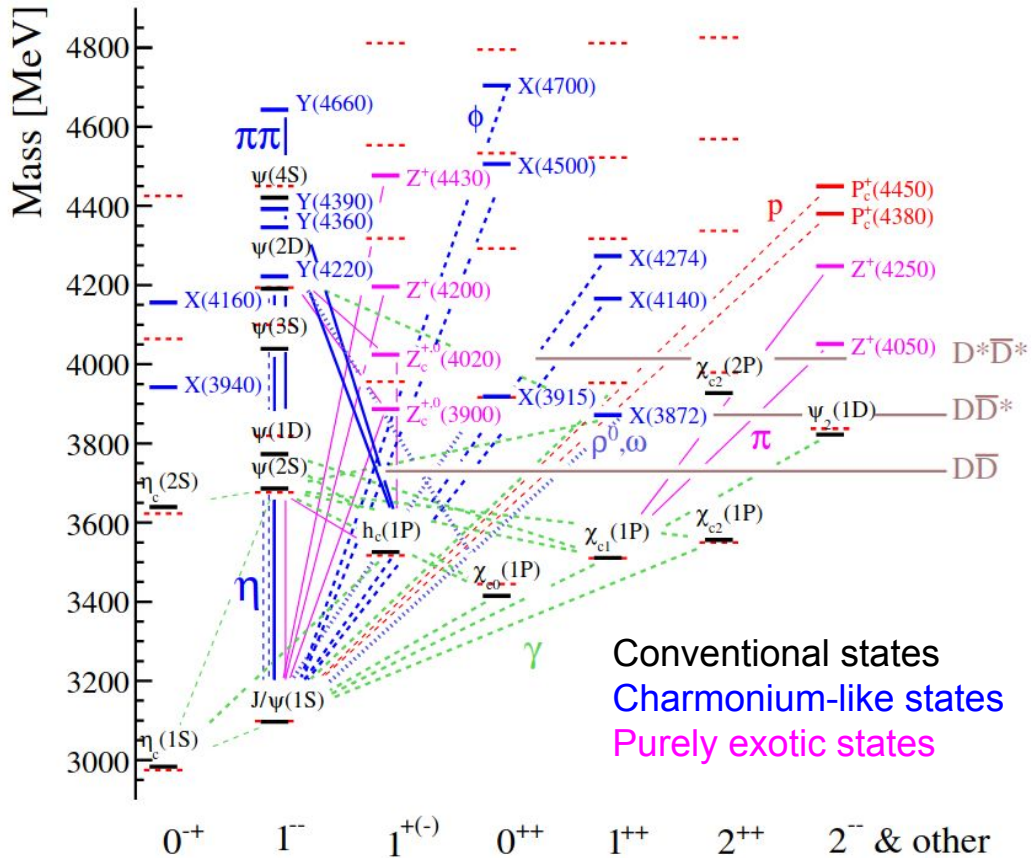
Goal:

- $50 \text{ ab}^{-1}$  at Y(4S)



PTEP 2020 (2020) 2, 029201

# Charmonium nomenclature



Below  $D\bar{D}$  threshold: well described by potential models

Above  $D\bar{D}$  threshold:

**X** states such as X(3872)

- neutral non-vector non-conventional states

**Y** states such as Y(4260) and Y(4660)

- neutral vector non-conventional states  $J^{PC} = 1^{--}$  states.

**Z** states such as  $Z_c^+(3900)$  and  $Z_b^+(10650)$

- charged states, can not be pure quarkonium
- similar for b- and c- onium

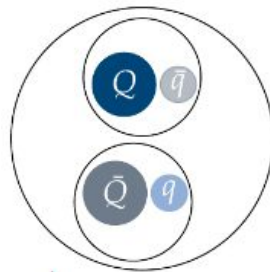
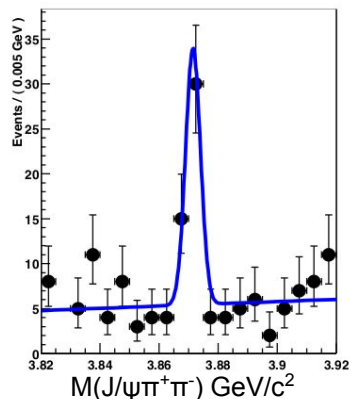
# Theoretical models

## Hadronic molecule

Compound state of two hadrons. The most promising model.

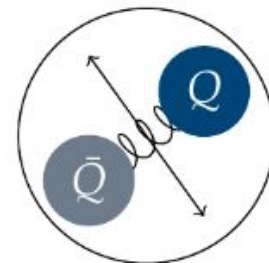
The charmonium-like states can be described as a mixture of pure charmonium and a molecular component:

$$X(3872) = C_1 \cdot |c\bar{c}\rangle + C_2 \cdot |D^0 \bar{D}^{0*}\rangle$$



## Hybrid

Conventional quark-antiquark mesons with excited gluon degrees of freedom.



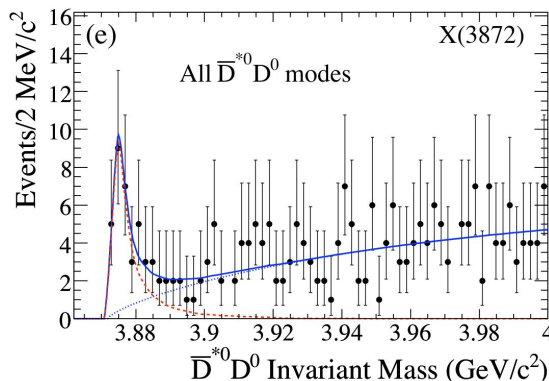
## Compact tetraquark

States containing four constituent quarks irrespective of their clustering.



## Hadroquarkonium

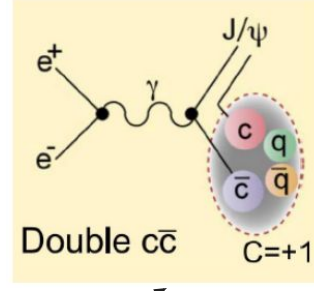
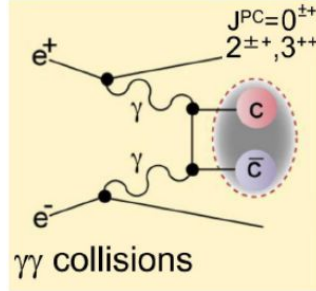
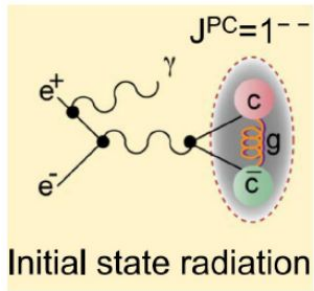
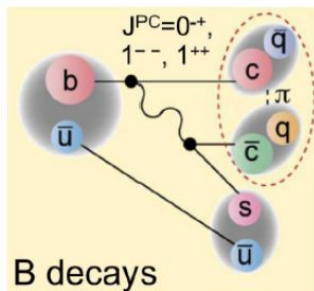
Compact quarkonium core surrounded by an excited light-quark cloud.



# Charmonium production at Belle II

At Belle / Belle II charmonium events comes “for free”:

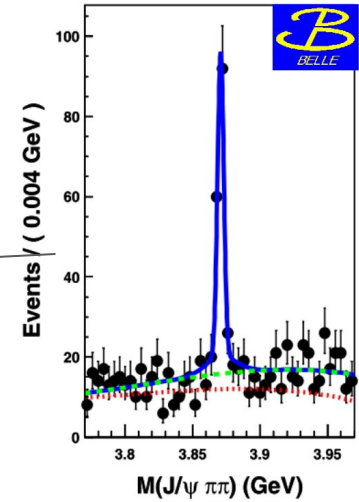
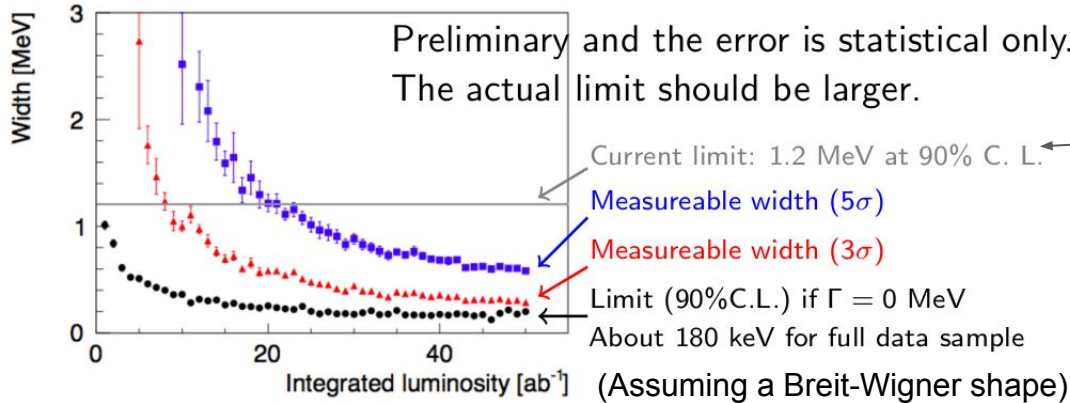
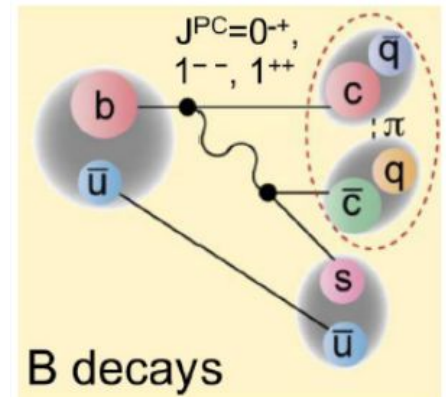
- We don't need to develop special triggers.
- We don't need to tune accelerator.
- Uniquely clean environment compare to pp production.



**Unique for Belle II !**

# B decays

- $B \rightarrow K X_{cc^-}$  is CKM favoured process with large branching fractions  $\sim 10^{-3} - 10^{-4}$ .
- Due to the known number of B pairs, it is possible to measure absolute BF[ $B \rightarrow X_{cc^-} K$ ].
- Access to X(3872) lineshape with  $D^0 \bar{D}^0 \pi^0$  channel.
- Access to unknown / hardly reconstructable modes with the B-tagging technique.



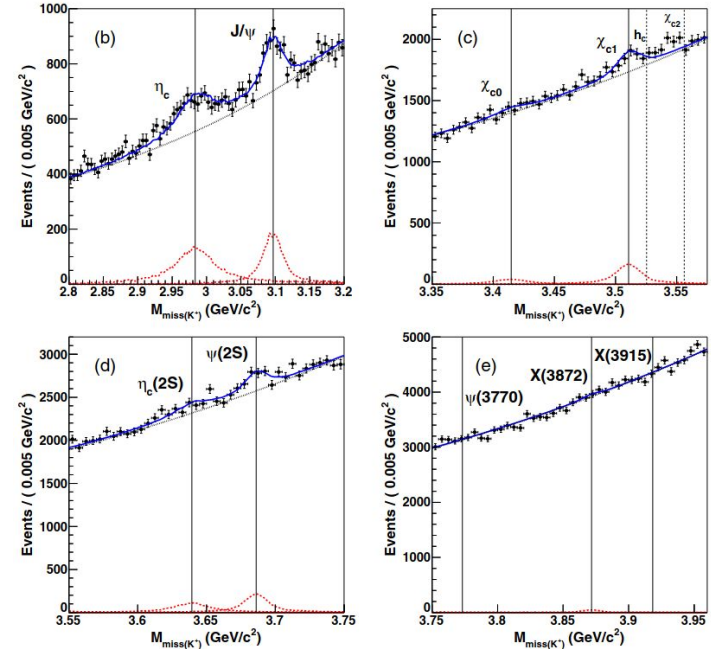
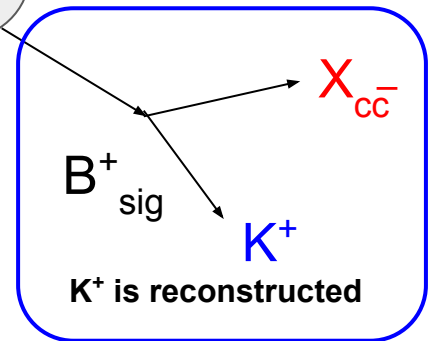
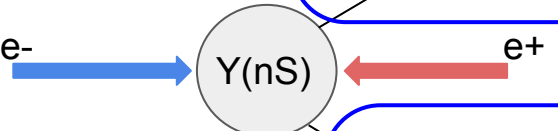
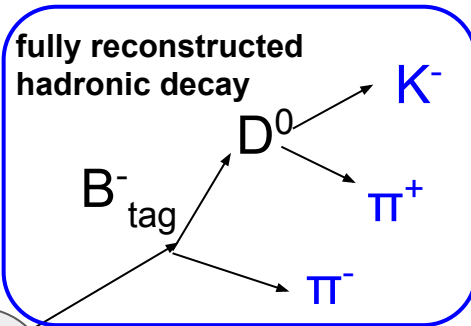
Phys. Rev. D 84, 052004 (2011)

# Charmonium + B-tagging

PRD 97, 012005 (2018)

Precise measurement of the  $e^+e^-$  energy allows to observe mass peaks with partial reconstruction:

$$M_{X_{cc}} = \sqrt{p_{e^+e^-} - p_{B_{tag}} - p_K}$$

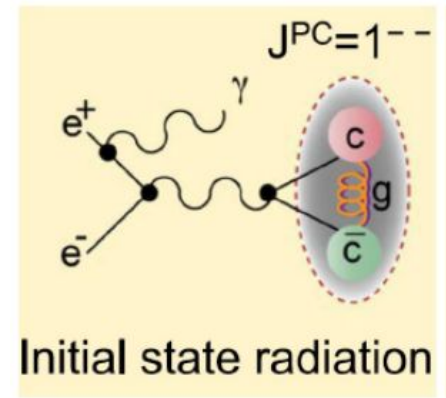


Mode	Yield	Significance ( $\sigma$ )	$\epsilon (10^{-3})$	$\mathcal{B} (10^{-4})$	World average for $\mathcal{B} (10^{-4})$ [10]
$\eta_c$	$2590 \pm 180$	14.2	$2.73 \pm 0.02$	$12.0 \pm 0.8 \pm 0.7$	$9.6 \pm 1.1$
$J/\psi$	$1860 \pm 140$	13.7	$2.65 \pm 0.02$	$8.9 \pm 0.6 \pm 0.5$	$10.26 \pm 0.031$
$\chi_{c0}$	$430 \pm 190$	2.2	$2.67 \pm 0.02$	$2.0 \pm 0.9 \pm 0.1 (<3.3)$	$1.50^{+0.15}_{-0.14}$
$\chi_{c1}$	$1230 \pm 180$	6.8	$2.68 \pm 0.02$	$5.8 \pm 0.9 \pm 0.5$	$4.79 \pm 0.23$
$\eta_c(2S)$	$1050 \pm 240$	4.1	$2.77 \pm 0.02$	$4.8 \pm 1.1 \pm 0.3$	$3.4 \pm 1.8$
$\psi(2S)$	$1410 \pm 210$	6.6	$2.79 \pm 0.02$	$6.4 \pm 1.0 \pm 0.4$	$6.26 \pm 0.24$
$\psi(3770)$	$-40 \pm 310$	-	$2.76 \pm 0.02$	$-0.2 \pm 1.4 \pm 0.0 (<2.3)$	$4.9 \pm 1.3$
$X(3872)$	$260 \pm 230$	1.1	$2.79 \pm 0.01$	$1.2 \pm 1.1 \pm 0.1 (<2.6)$	$(<3.2)$
$X(3915)$	$80 \pm 350$	0.3	$2.79 \pm 0.01$	$0.4 \pm 1.6 \pm 0.0 (<2.8)$	-

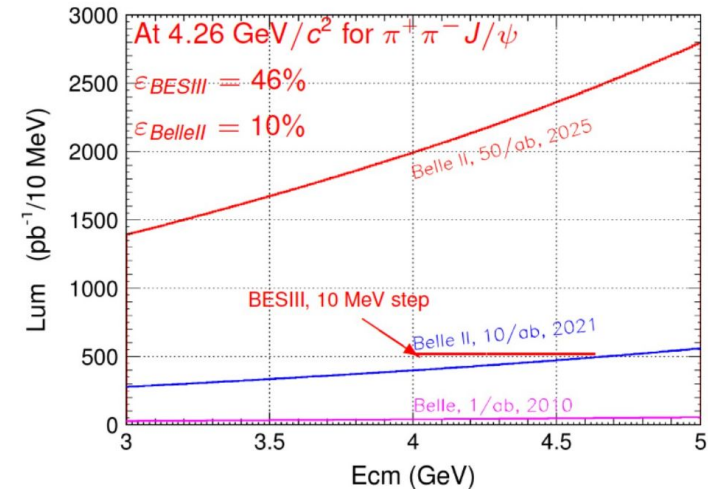


# Initial state radiation

- 50  $\text{ab}^{-1}$  data corresponds to 2000-2800  $\text{pb}^{-1}$  /10 MeV at 4-5 GeV, which is compatible with BES III (500  $\text{pb}^{-1}$  /10 MeV, but higher efficiency).
- Access to  $1^{--}$  states with masses  $> 4.6$  GeV.
- Effective luminosity and detection efficiency are relatively low.
- c-baryons ( $\Lambda_c^+ \Sigma_c^-, \Sigma_c^+ \Sigma_c^-$ ) and cs-mesons pairs production ( $D_s^+ D_{s2}^- (2573), D_s^- D_{s0}^* (2317)$ ).



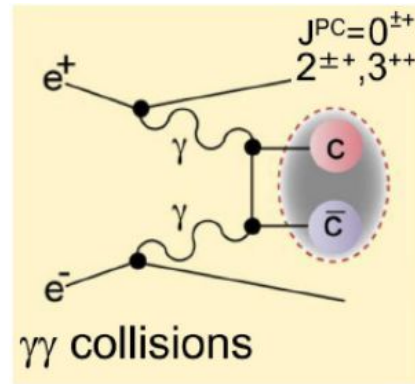
Golden Channels	$E_{c.m.}$ (GeV)	Statistical error (%)	Related XYZ states
$\pi^+ \pi^- J/\psi$	4.23	7.5 (3.0)	$Y(4008), Y(4260), Z_c(3900)$
$\pi^+ \pi^- \psi(2S)$	4.36	12 (5.0)	$Y(4260), Y(4360), Y(4660), Z_c(4050)$
$K^+ K^- J/\psi$	4.53	15 (6.5)	$Z_{cs}$
$\pi^+ \pi^- h_c$	4.23	15 (6.5)	$Y(4220), Y(4390), Z_c(4020), Z_c(4025)$
$\omega \chi_{c0}$	4.23	35 (15)	$Y(4220)$
		10 $\text{ab}^{-1}$ 50 $\text{ab}^{-1}$	



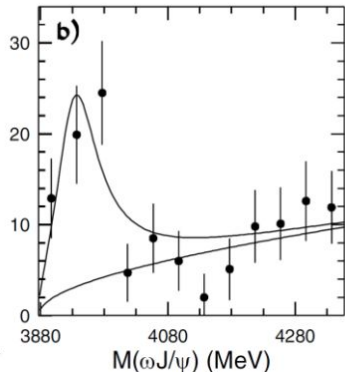
PTEP 2020 (2020) 2, 029201

# Two-photon production

- Access to a wide range of quantum numbers.  $J/\psi\omega$  and  $J/\psi\phi$  in particular.
- Provides a better S/B ratio compare to B-decays.
- Allows to measure the  $Q^2$  dependence of  $c\bar{c}$  production.

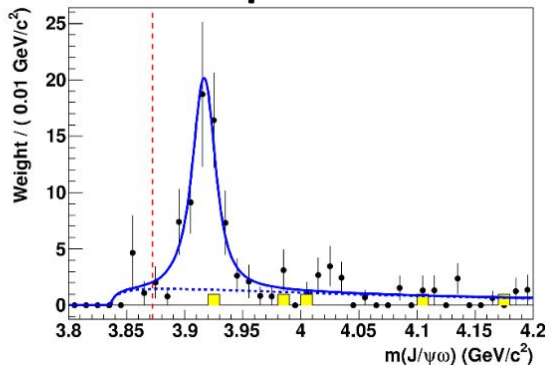


## B-decay



PRL 94, 182002 (2005)

## Two photon

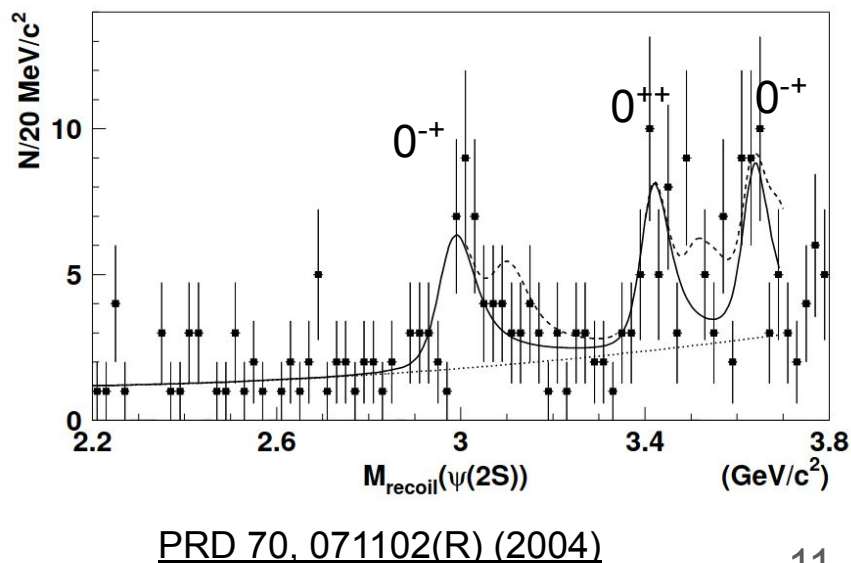
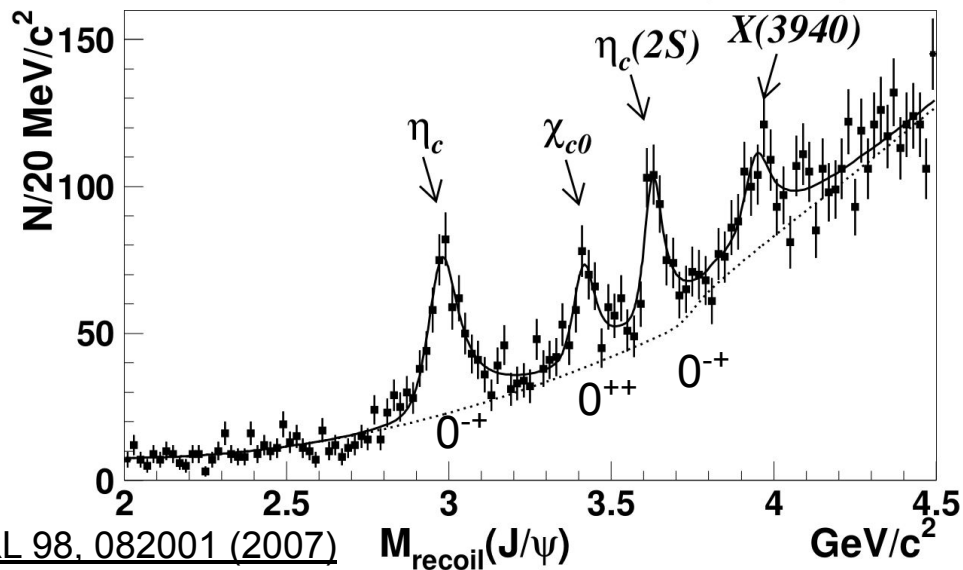
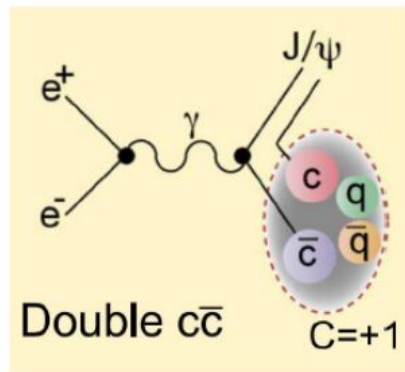


PRD 86, 072002 (2012)

The  $e^+e^- \rightarrow e^+e^- H$  cross section is largest in the kinematic regions where the  $Q^2$  value is very close to zero and  $M_H$  is small when compared with the beam energy. In contrast, in regions where either  $W$  or  $Q^2$  is much larger than the typical QCD energy scale ( $\sim 1$  GeV), the cross section decreases rapidly

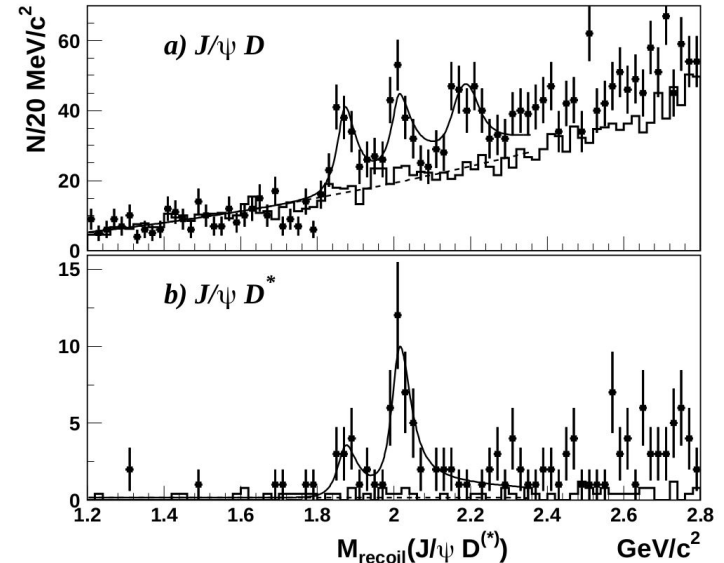
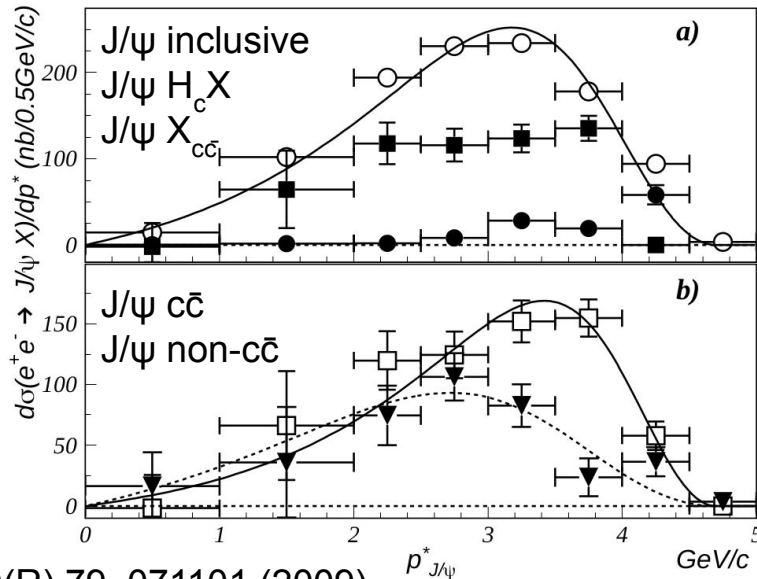
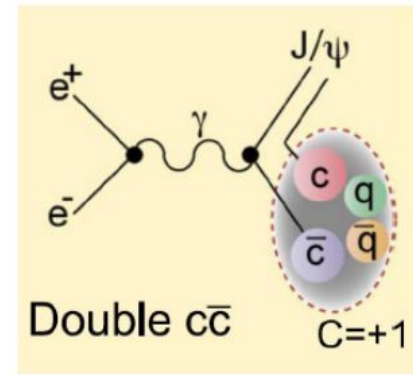
# Double charmonium production

- All observed processes are of the type  $e^+e^- \rightarrow c\bar{c} (J=1) c\bar{c} (J=0)$ . Belle II will allow to study  $e^+e^- \rightarrow \eta_c X_{(c\bar{c})}$  and  $e^+e^- \rightarrow \chi_{c0} X_{(c\bar{c})}$  processes.
- $e^+e^- \rightarrow \psi(2S) X_{(c\bar{c})}$  can be studied at Belle II with larger statistics.
- Large cross section revealed importance of the next order corrections in NRQCD.



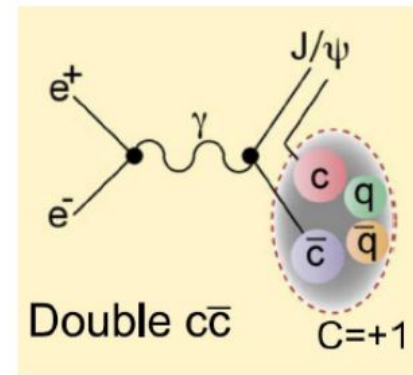
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# Double charmonium production

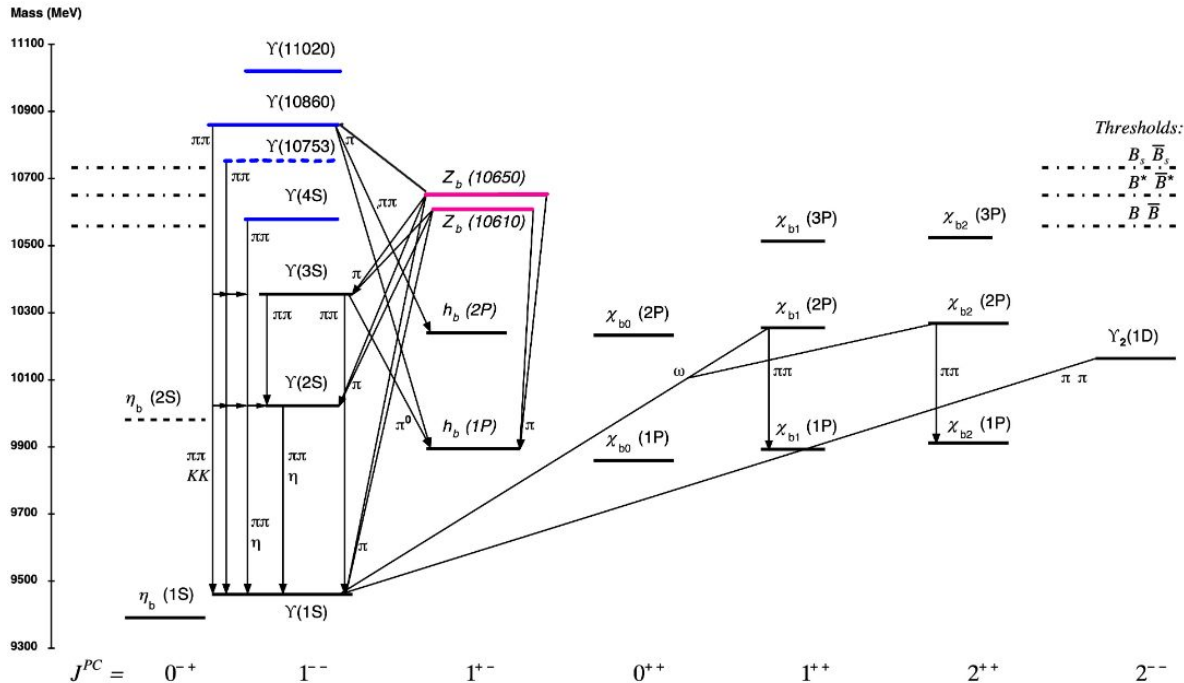
- All observed processes are of the type  $e^+e^- \rightarrow c\bar{c}$  ( $J=1$ )  $c\bar{c}$  ( $J=0$ ). Belle II will allow to study  $e^+e^- \rightarrow \eta_c X_{(c\bar{c})}$  and  $e^+e^- \rightarrow \chi_{c0} X_{(c\bar{c})}$  processes.
- $e^+e^- \rightarrow \psi(2S) X_{(c\bar{c})}$  can be studied at Belle II with larger statistics.
- Large cross section revealed importance of the next order corrections in NRQCD.



**$N_{\text{charged}} > 4$**

**Dominance of the  $e^+e^- \rightarrow J/\psi c\bar{c}$  production mechanism. It should be noted that in this analysis (unlike that of PRL 89, 142001 (2002)) no correction for the charged track multiplicity ( $N_{\text{charged}} > 4$ ) requirement was applied for any of the processes. For  $e^+e^- \rightarrow J/\psi$  non- $cc$ , such corrections are only possible by relying on a model, while for  $e^+e^- \rightarrow J/\psi c\bar{c}$  they are close to unity.**

# Bottomonium



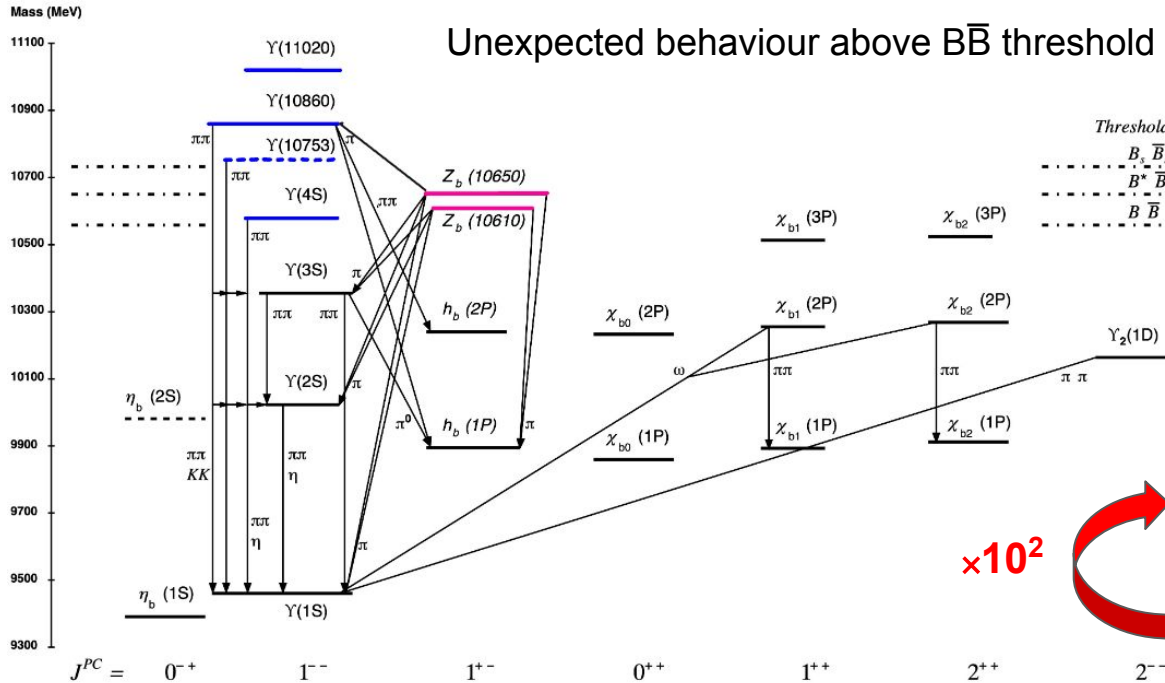
We have two type of states:

- Below  $B\bar{B}$  threshold
- Above  $B\bar{B}$  threshold

The states below  $B\bar{B}$  mesons threshold are well described by the potential models.

Conventional bottomonium (pure  $bb$  states)  
 Bottomonium-like states (mixture of  $bb$  and  $B\bar{B}$ )  
 Purely exotic states

# Bottomonium



Process	$\Gamma$ , MeV
$\Upsilon(5S) \rightarrow \Upsilon(1S)\eta$	$0.039 \pm 0.011$
$\Upsilon(5S) \rightarrow \Upsilon(2S)\eta$	$0.204 \pm 0.44$
$\Upsilon(4S) \rightarrow \Upsilon(1S)\eta$	$0.004 \pm 0.0008$
$\Upsilon(4S) \rightarrow h_b(1P)\eta$	$0.045 \pm 0.007$
$\Upsilon(3S) \rightarrow \Upsilon(1S)\eta$	$< 0.002 \cdot 10^{-3}$
$\Upsilon(2S) \rightarrow \Upsilon(1S)\eta$	$(0.0093 \pm 0.0015) \cdot 10^{-3}$

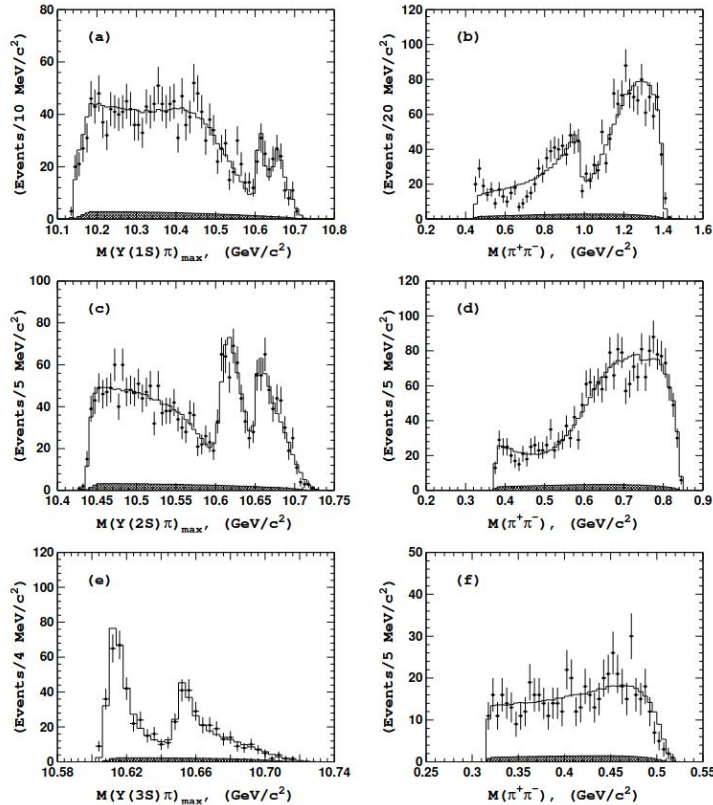
Process	$\Gamma$ , MeV
$\Upsilon(5S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	$0.59 \pm 0.04 \pm 0.09$
$\Upsilon(5S) \rightarrow \Upsilon(2S)\pi^+\pi^-$	$0.85 \pm 0.07 \pm 0.16$
$\Upsilon(5S) \rightarrow \Upsilon(3S)\pi^+\pi^-$	$0.52^{+0.20}_{-0.17} \pm 0.10$
$\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	$0.0060 \pm 0.0005$
$\Upsilon(3S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	$0.0009 \pm 0.00008$
$\Upsilon(4S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	$0.0019 \pm 0.0002$

Conventional bottomonium (pure bb states)  
 Bottomonium-like states (mixture of bb and BB)  
 Purely exotic states

PRL 100, 112001 (2008)

# Bottomonium

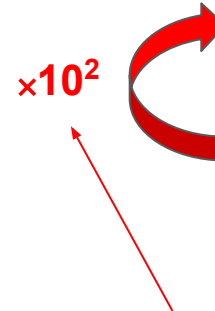
Mod. Phys. Lett. A 32, No04, 1750025



PRL 108 (2012) 122001

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$\Upsilon(4S) \rightarrow \Upsilon(1S)\pi^+\pi^-$	$0.0019 \pm 0.0002$



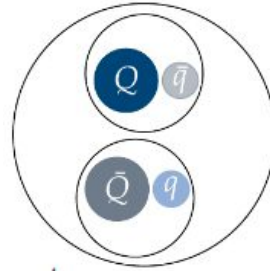
PRL 100, 112001 (2008)

Proceed via intermediate exotic state  
 $Y(5S) \rightarrow [Z_b^+ \rightarrow Y(nS)\pi^+]\pi^-$



## Hadronic molecule

Compound state of two hadrons. The most promising model. The bottomonium-like states can be described as a mixture of pure bottomonium and a molecular component:



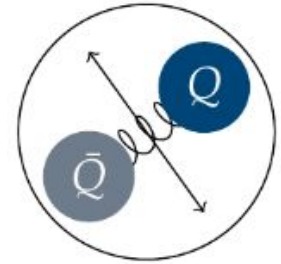
$$Y(10860) = C_1 \cdot |bb\rangle + C_2 \cdot |B_{(s)}^{((*)*)} B_{(s)}^{((*)*)}\rangle$$

$Z_b$ decay mode	Branching fraction
$Z_b^+(10610) \rightarrow \Upsilon(nS)/h_b(mP)\pi^+$	$14.4_{-1.9}^{+2.5}\%$
$Z_b^+(10610) \rightarrow B^+\bar{B}^{*0}/B^+\bar{B}^0$	$85.6_{-2.9}^{+2.1}\%$
$Z_b^+(10650) \rightarrow \Upsilon(nS)/h_b(mP)\pi^+$	$26.6_{-4.7}^{+5.0}\%$
$Z_b^+(10650) \rightarrow B^{*+}\bar{B}^{*0}$	$74_{-6}^{+4}\%$

PRL, 108, 122001 (2012)

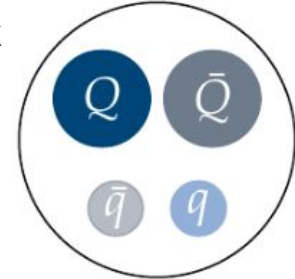
## Hybrid

Conventional quark-antiquark mesons with excited gluon degrees of freedom.



## Compact tetraquark

States containing four constituent quarks irrespective of their clustering.



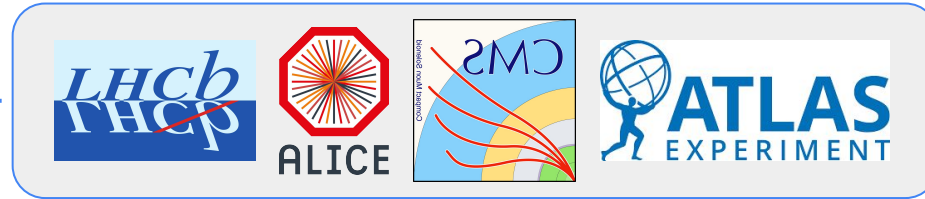
## Hadroquarkonium

Compact quarkonium core surrounded by an excited light-quark cloud.



# Bottomonium production

Prompt production

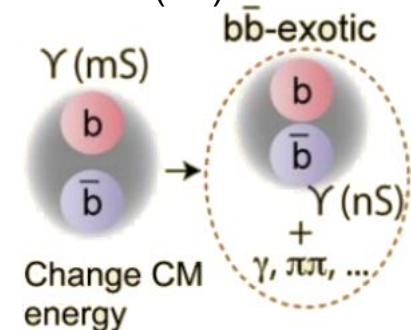


Direct production in  $e^+e^-$  collisions

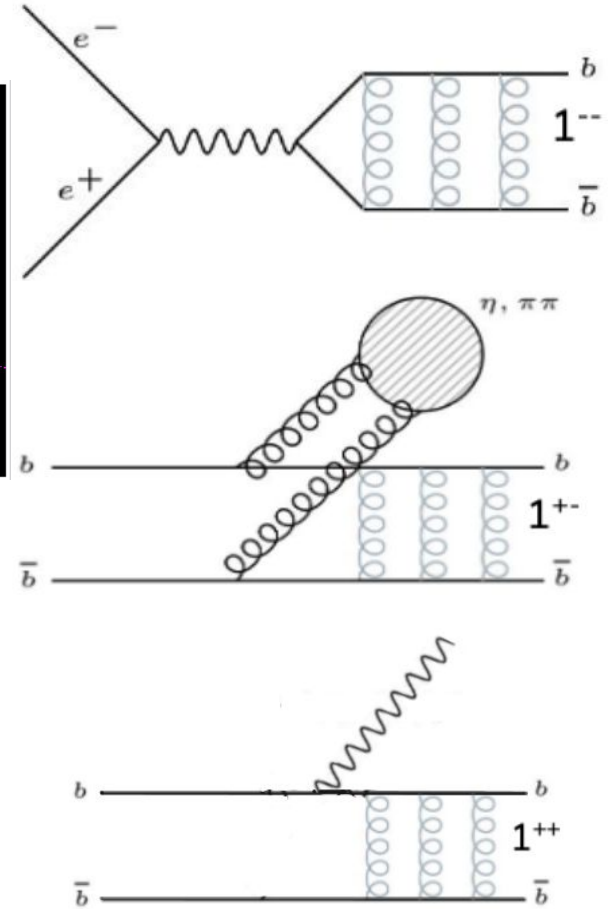
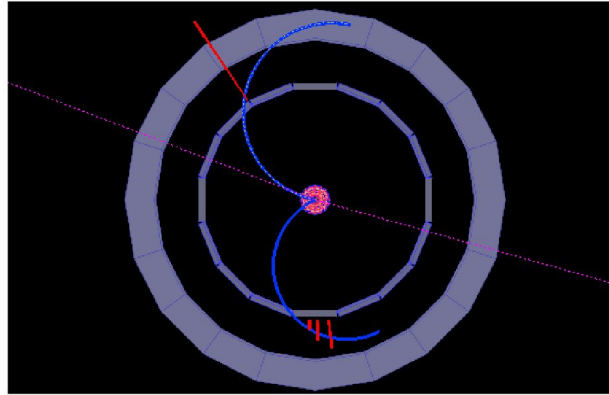
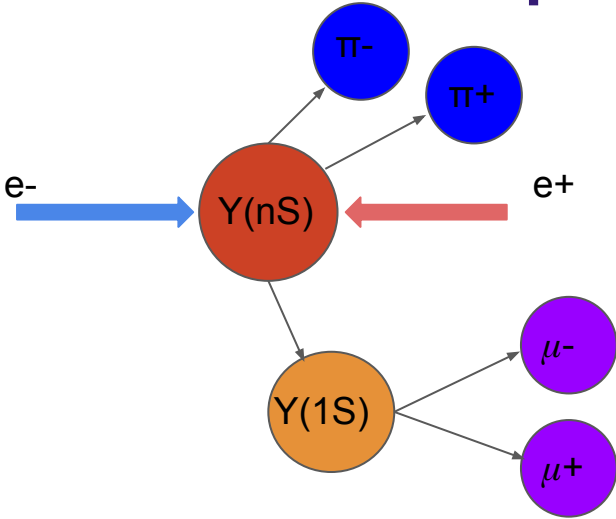


- Provides a unique clean environment (we have only bottomonium in the event);
- Precise measurement of the beam energy gives access to “unreconstructable” particles and decay modes;
- Allows tuning CM energy
- **Only  $Y(nS)$  states can be produced with quantum numbers of the photon  $1^-$**
- Other quantum numbers can be obtained via hadronic or radiative transitions from  $1^-$  states;

gives access to exotic states above  $Y(4S)$



# Bottomonium production



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- Precise measurement of the beam energy gives access to “unreconstructable” particles and decay modes;
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- **Only  $Y(nS)$  states can be produced with quantum numbers of the photon  $1^-$**
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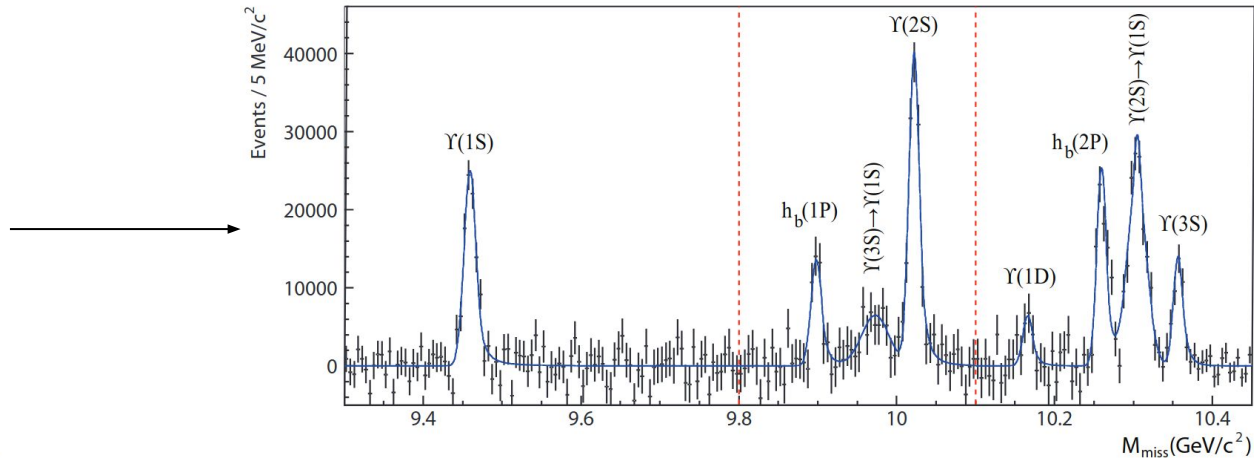
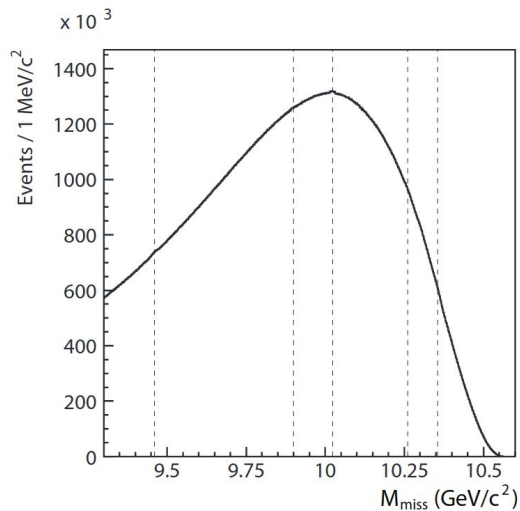
# Missing mass technique for bottomonium

Precise measurement of the  $e^+e^-$  energy allows to observe mass peaks with partial reconstruction:

$\pi^+ \pi^-$  missing mass has better resolution than  $M(\mu^+ \mu^-)$

$$M_{\text{miss}}(\pi\pi) = \sqrt{(E_{\text{c.m.}} - E_{\pi\pi}^*)^2 - p_{\pi\pi}^{*2}}$$

$\Upsilon(5S) \rightarrow X_{b\bar{b}} \pi^+ \pi^-$  transitions:



PRL 108 (2012) 032001

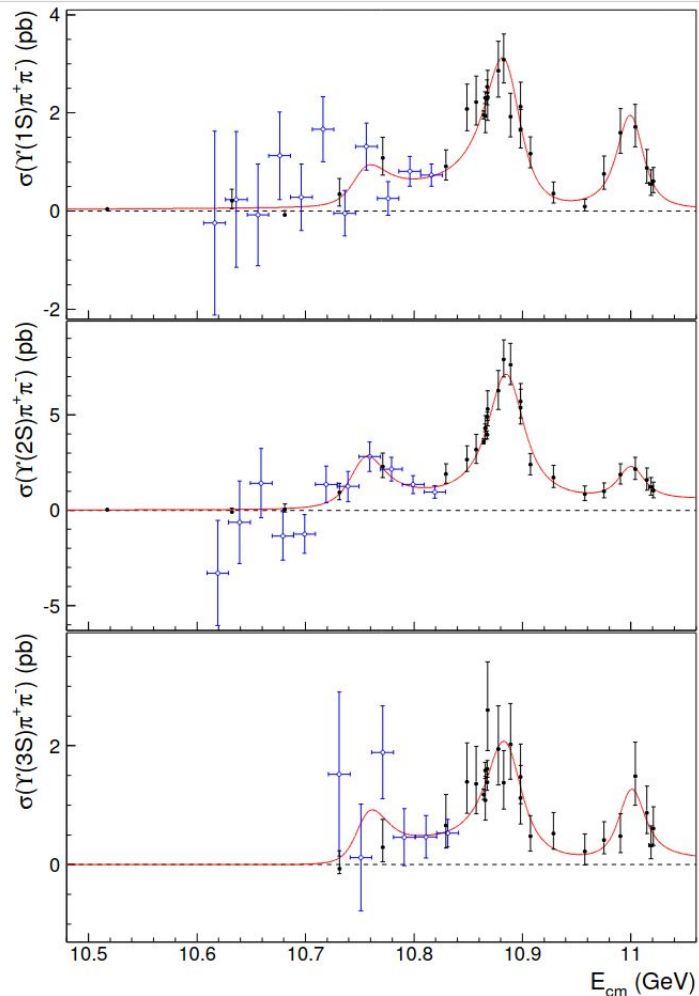
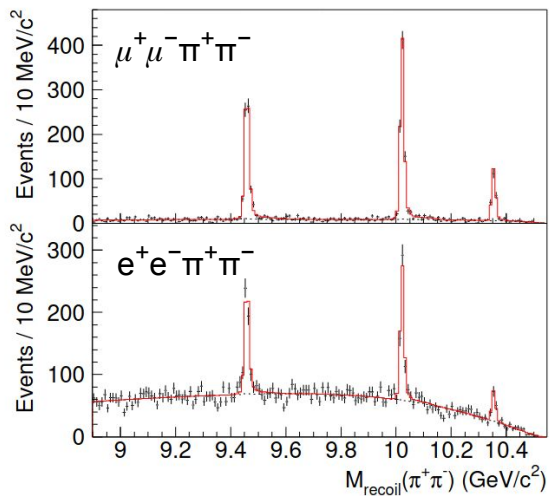
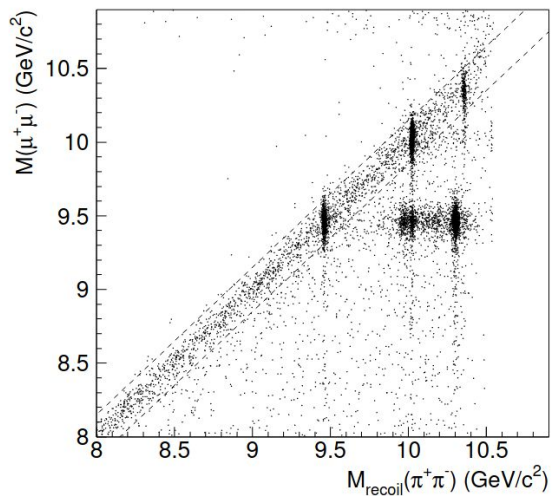
Very important for unreconstructable bottomonium states  $\eta_b(nS)$ ,  $h_b(nP)$

# Discovery of the $Y(10753)$

$Y(10753)$  state was observed in the  $e^+e^- \rightarrow Y(nS) \pi^+ \pi^-$  ( $n = 1, 2, 3$ ) cross section energy dependence.

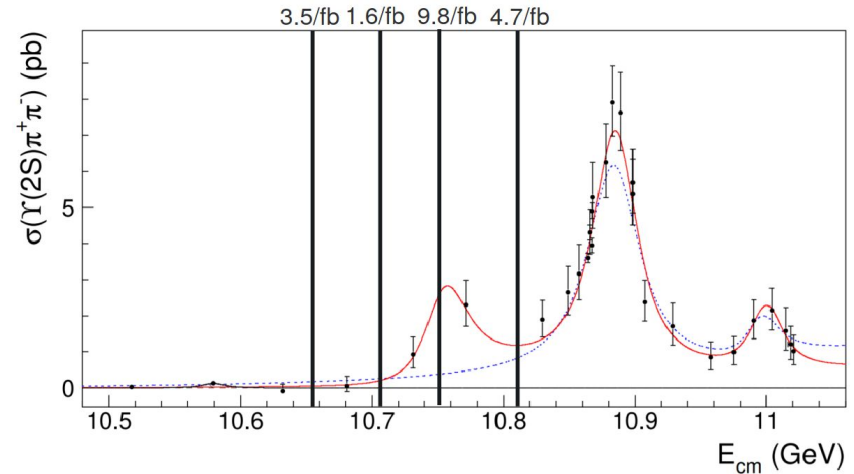
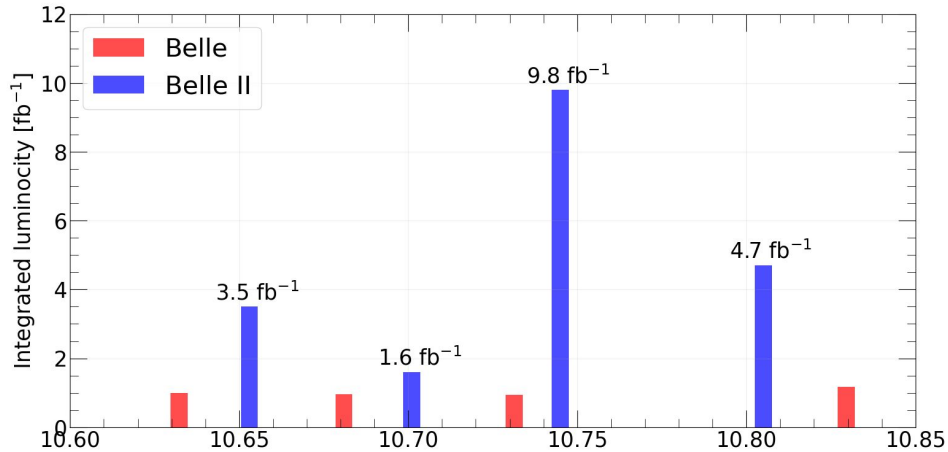
JHEP 10 (2019) 220

measurement at  $\sqrt{s} = 10.864$  GeV with  $L = 47.647$  fb $^{-1}$ :



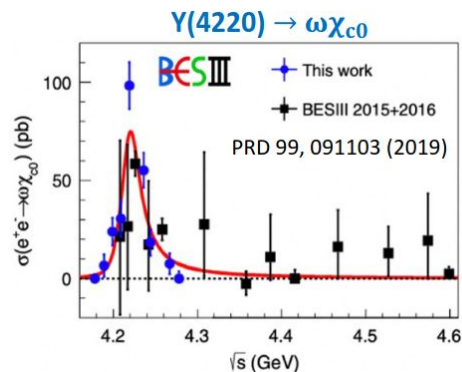
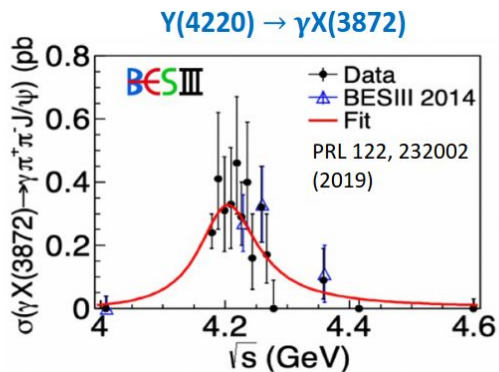
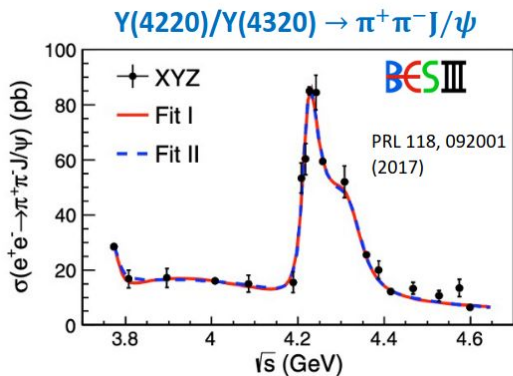
# Belle II energy scan above $Y(4S)$

- Unique data provide an opportunity to study  $Y(10753)$  in different final states and understand its nature.
- Scan above  $Y(4S)$  has a good potential for early physics impact by Belle II even with small statistics.

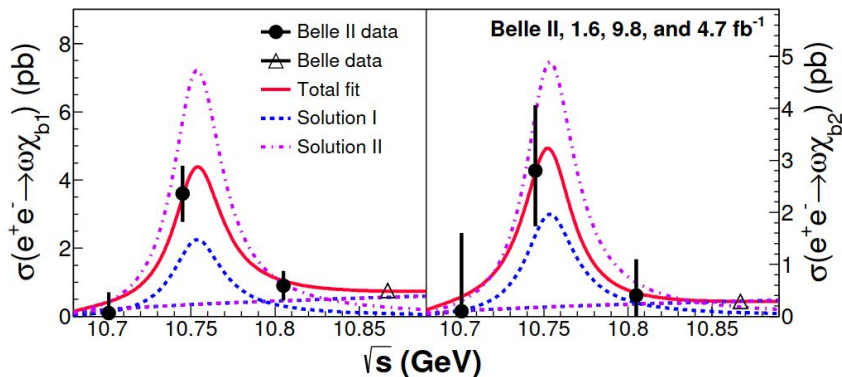


# Observation of $e^+e^- \rightarrow \omega\chi_{bJ}(1P)$ at $\sqrt{s} = 10.745$ GeV

- Similar to Y(10753) structure named Y(4220) was observed in  $e^+e^- \rightarrow J/\psi \pi^+ \pi^-$  cross section dependence by BES III (PRL 118, 092001 (2017)).
- BES III also observed the Y(4220) peak in  $\gamma X(3872)$  and  $\omega\chi_{c0}$  final states (PRL 122, 232002 (2019), PRD 99, 091103(R) (2019)).
- $\omega\chi_{b1,2}$  production was found to be enhanced near Y(5S) (PRL 113, 142001 (2014)).
- We expect Y(10753) to decay into  $\gamma[X_b \rightarrow \omega Y(1S)]$  and  $\omega\chi_{bJ}$  final states.



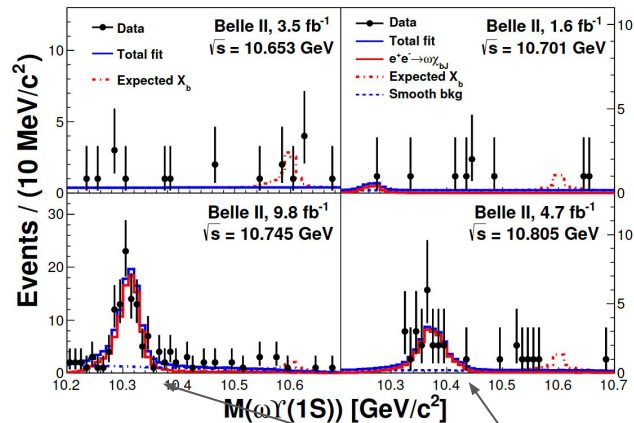
# Observation of $e^+e^- \rightarrow \omega\chi_{bJ}(1P)$ at $\sqrt{s} = 10.745$ GeV



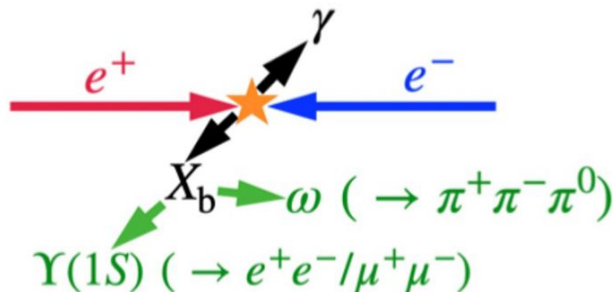
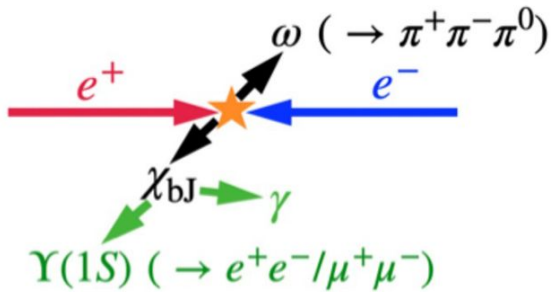
Strongly enhanced at 10.745 GeV.

$$\sigma(e^+e^- \rightarrow \omega\chi_{b1}) = 3.6 \pm 0.7 \pm 0.5 \text{ pb}$$

$$\sigma(e^+e^- \rightarrow \omega\chi_{b2}) = 2.8_{-1.0}^{+1.2} \pm 0.4 \text{ pb}$$



No evidence of  $X_b$  ( $b\bar{b}$  partner of  $X(3872)$ ) state. These are reflections of  $\omega\chi_{bJ}(1P)$ .



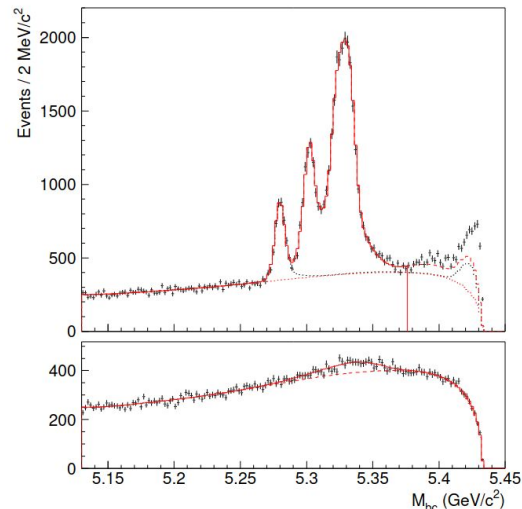
arXiv:2208.13189



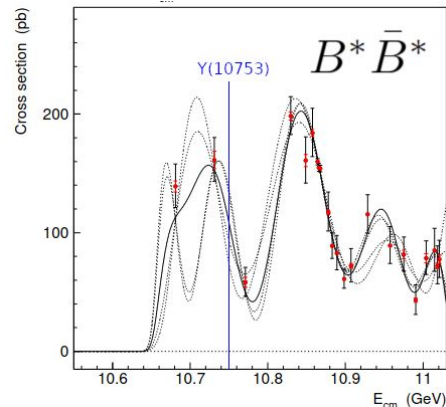
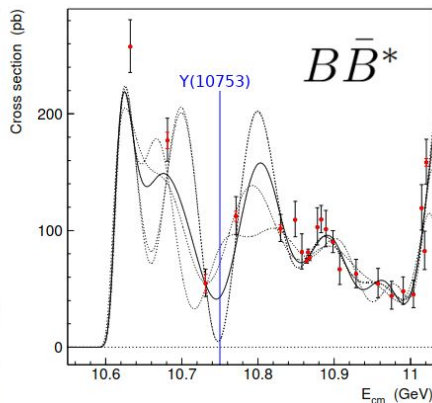
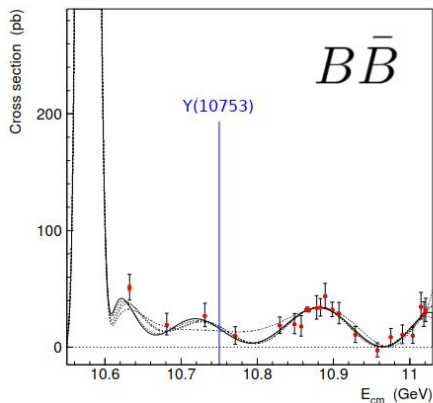
# $B\bar{B}$ decomposition with B-tagging

Study the energy dependence of the  $B\bar{B}$  pairs production.

- B-tagging can be used to measure the  $B_{(s)}^{(*)} \bar{B}_{(s)}^{(*)}$  cross section energy dependence.
- Yet another unique way to study bottomonium at Belle / Belle II.
- A good probe for bottomonium models (especially the molecular model).



$$M_{bc} = \sqrt{(E_{cm}/2)^2 - p_B^2}$$



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# Belle II potential charmonium

The main source of improvements and new results is the increase in statistics.

- Measure double charmonium production cross sections with  $e^+e^- \rightarrow \psi(2S)X_{(c\bar{c})}$  and search for  $e^+e^- \rightarrow \eta_c X_{(c\bar{c})}$  and  $e^+e^- \rightarrow \chi_{c0} X_{(c\bar{c})}$ .
- Measure  $X(3872)$  lineshape with  $D^0\bar{D}^0\pi^0$  channel.
- Study high energy region ( $E > 4.6$  GeV) unapproachable for BES III with ISR production.
- Search for new charmonium-like states or new decay modes. Confirm the states and transitions obtained with low statistics. Measure quantum numbers of observed charmonium-like states.
- Search for new charmonium state with improved B-tagging method.
- Larger statistics will allow to measure  $p^*_{c\bar{c}}$ ,  $Q^2$  production dependence.

# Belle II potential bottomonium

Scan above  $Y(4S)$  gives an opportunity for a lot of unique studies:

- $Y(10753)$  decays to different exclusive and inclusive final states. Study of its properties;
- Energy dependence of the various final state cross sections;
- $B\bar{B}$  decomposition and its cross section dependence on CM energy;

Wide range of long-term non- $Y(4S)$  possibilities:

- Increase the above- $Y(4S)$  scan statistics;
- $Y(6S)$  region study with high statistics after accelerator upgrade;

## Golden Modes

$$e^+e^- \rightarrow \pi^+\pi^-\Upsilon(pS)(\rightarrow \ell^+\ell^-)$$

$B\bar{B}$  decomposition

$$\pi^+\pi^-\text{ Dalitz}$$

$$Y_b \rightarrow \omega\eta_b(1S)$$

$$Y_b \rightarrow \omega\chi_{bJ}(1P)$$

## Silver Modes

$$Y_b \rightarrow \pi^+\pi^-X \text{ (inclusive)}$$

$$Y_b \rightarrow \eta X \text{ (inclusive)}$$

$$Y_b \rightarrow \eta\Upsilon(1S, 2S)(\rightarrow \ell^+\ell^-)$$

$$Y_b \rightarrow \eta'\Upsilon(1S)(\rightarrow \ell^+\ell^-)$$

$$Y_b \rightarrow \Upsilon(1S) \text{ (inclusive)}$$

## Bronze Modes

$$Y_b \rightarrow \gamma X_b$$

$$Y_b \rightarrow \pi^0\pi^0\Upsilon(pS)(\rightarrow \ell^+\ell^-)$$

$$Y_b \rightarrow KK(\phi)\Upsilon(pS)(\rightarrow \ell^+\ell^-)$$

$$Y_b \rightarrow \pi^0\pi^0X \text{ (inclusive)}$$

$$Y_b \rightarrow \pi^0X \text{ (incl. or excl.)}$$

...

Backup

# X(3872) lineshape with $D^0\bar{D}^0\pi^0$ channel

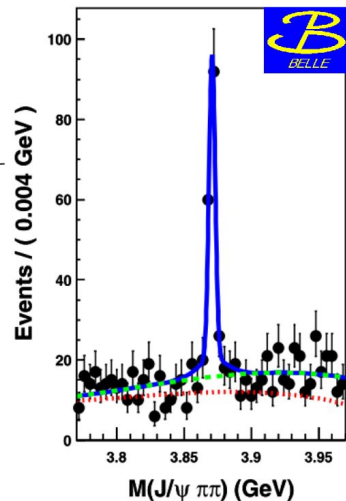
$\Gamma_{\text{tot}}$   
< 1.2 MeV (90% C.L.)

<

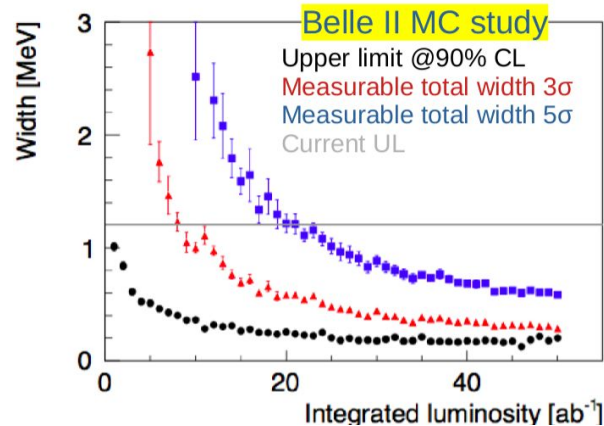
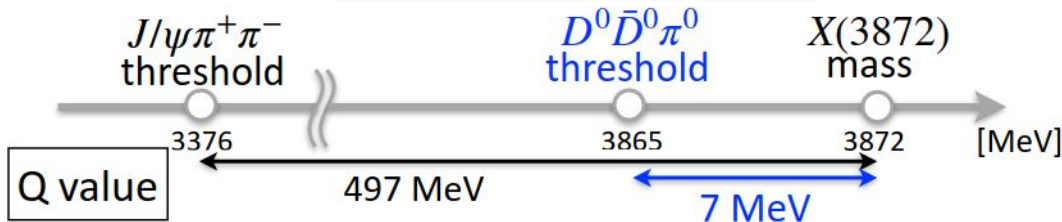
Mass resolution  
 $1.86 \pm 0.01 \text{ MeV}/c^2$

Phys. Rev. D 84, 052004 (2011)

- Improvement of mass resolution is more essential than that of statistics.
- We can use channel with good mass resolution:
  - Mass resolution:  $684 \pm 8 \text{ keV}$
  - Signal yield with  $1 \text{ ab}^{-1}$ :  $64.5 \pm 23.9$



Comes from large error of  
 $\text{Br}(B^\pm \rightarrow K^\pm X(3872))$   
 $\times \text{Br}(X(3872) \rightarrow D^0\bar{D}^0\pi^0)$



(Assuming a Breit-Wigner shape)

# B decays

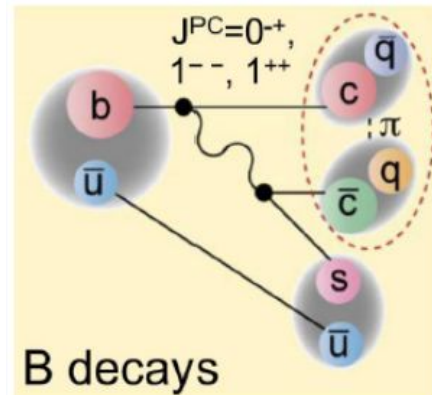
**S** - seen (observed)

**NS** - not seen or excluded

**N** - not performed (in a given mass range)

**MF** - missing fit (fit have not been extended to this mass range)

– no search has been performed in this mode



State	$J^{PC}$	$\psi\pi\pi$	$\psi\omega$	$\psi\gamma$	$\psi\phi$	$\psi\eta$	$\psi'\pi\pi$	$\psi'\omega$	$\psi'\gamma$	$\chi_{c\gamma}$	$p\bar{p}$	$\Lambda\bar{\Lambda}$	$\Lambda_c\bar{\Lambda}_c$	$D\bar{D}$	$D\bar{D}^*$	$2D^*$	$2D_s^{(*)}$	$\gamma\gamma$
X(3872)	$1^{++}$	<b>S</b>	<b>S</b>	<b>S</b>	—	<b>NS</b>	—	—	<b>S</b>	<b>NS</b>	<b>MF</b>	<b>MF</b>	—	—	<b>S</b>	—	—	<b>NS</b>
Y(3940)	$J^{P+}$	<b>MF</b>	<b>S</b>	<b>NS</b>	—	—	—	—	<b>MF</b>	—	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>NS</b>	—	<b>N</b>	<b>N</b>
Z(3930)	$2^{++}$	<b>MF</b>	<b>MF</b>	<b>NS</b>	—	—	—	—	<b>MF</b>	—	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>MF</b>	—	<b>N</b>	<b>N</b>
Y(4140)	$J^{P+}$	<b>MF</b>	<b>MF</b>	<b>N</b>	<b>S</b>	—	<b>N</b>	—	<b>N</b>	—	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>
X(4160)	$0^{P+}$	<b>MF</b>	<b>MF</b>	<b>N</b>	<b>MF</b>	—	<b>N</b>	—	<b>N</b>	—	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>
Y(4260)	$1^{--}$	<b>NS</b>	—	—	—	<b>MF</b>	<b>N</b>	—	—	<b>N</b>	<b>MF</b>	<b>MF</b>	—	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	—
X(4350)	$J^{P+}$	<b>MF</b>	<b>MF</b>	<b>N</b>	<b>MF</b>	—	<b>N</b>	<b>N</b>	<b>N</b>	—	<b>MF</b>	<b>MF</b>	—	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>
Y(4350)	$1^{--}$	<b>MF</b>	—	—	—	<b>MF</b>	<b>N</b>	—	—	<b>N</b>	<b>MF</b>	<b>MF</b>	—	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	—
Y(4660)	$1^{--}$	<b>N</b>	—	—	—	<b>MF</b>	<b>N</b>	—	—	<b>N</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	—

# Initial state radiation

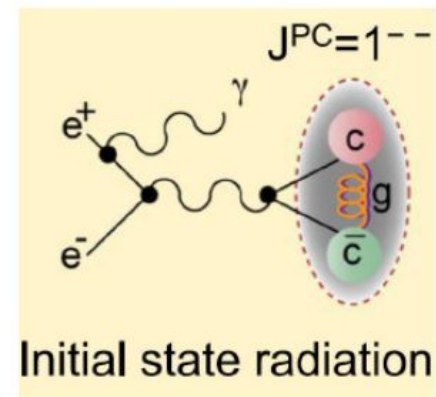
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**N** - not performed (in a given mass range)

**MF** - missing fit (fit have not been extended to this mass range)

– no search has been performed in this mode



State	$J^{PC}$	$\psi\pi\pi$	$\psi'\pi\pi$	$\psi\eta$	$\chi_{c\gamma}$	$p\bar{p}$	$\Lambda\bar{\Lambda}$	$\Lambda_c\bar{\Lambda}_c$	$D\bar{D}$	$D\bar{D}^*$	$2D^*$	$2D_s^{(*)}$
Y(4260)	$1^{--}$	<b>S</b>	<b>NS</b>	<b>NS</b>	<b>NS</b>	<b>NS</b>	<b>MF</b>	—	<b>NS</b>	<b>NS</b>	<b>NS</b>	<b>NS</b>
Y(4350)	$1^{--}$	<b>NS</b>	<b>S</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>
Y(4660)	$1^{--}$	<b>NS</b>	<b>S</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>S</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>	<b>MF</b>

# Two-photon production

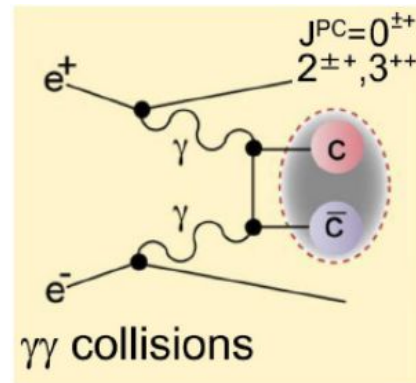
**S** - seen (observed)

**NS** - not seen or excluded

**N** - not performed (in a given mass range)

**MF** - missing fit (fit have not been extended to this mass range)

– no search has been performed in this mode



State	$J^{PC}$	$\psi\pi\pi$	$\psi\omega$	$\psi\gamma$	$\psi\phi$	$\psi'\pi\pi$	$\psi'\omega$	$\psi'\gamma$	$p\bar{p}$	$\Lambda\bar{\Lambda}$	$\Lambda_c\bar{\Lambda}_c$	$D\bar{D}$	$D\bar{D}^*$	$2D^*$	$2D_s^{(*)}$
X(3872)	$1^{++}$	<b>N</b>	hard	hard	—	—	—	hard	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>N</b>	—	—
Y(3915)	$0^{++}$	<b>N</b>	<b>S</b>	hard	—	—	—	hard	<b>MF</b>	<b>MF</b>	—	<b>MF</b>	<b>N</b>	—	<b>N</b>
Z(3930)	$2^{++}$	<b>N</b>	<b>MF</b>	hard	—	—	—	hard	<b>MF</b>	<b>MF</b>	—	<b>S</b>	<b>N</b>	—	<b>N</b>
Y(4140)	$J^{P+}$	<b>N</b>	<b>MF</b>	hard	<b>NS</b>	<b>N</b>	—	hard	<b>N</b>	<b>N</b>	—	<b>MF</b>	<b>N</b>	<b>N</b>	<b>N</b>
X(4160)	$0^{P+}$	<b>N</b>	<b>MF</b>	hard	<b>NS</b>	<b>N</b>	—	hard	<b>N</b>	<b>N</b>	—	<b>MF</b>	<b>N</b>	<b>N</b>	<b>N</b>
X(4350)	$J^{P+}$	<b>N</b>	<b>N</b>	hard	<b>S</b>	<b>N</b>	<b>N</b>	hard	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>	<b>N</b>



# Double charmonium production

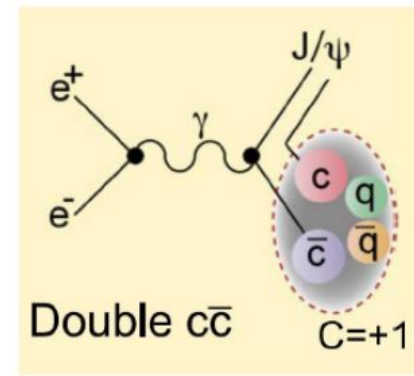
**S** - seen (observed)

**NS** - not seen or excluded

**N** - not performed (in a given mass range)

**MF** - missing fit (fit have not been extended to this mass range)

– no search has been performed in this mode



State	$J^{PC}$	$\psi\pi\pi$	$\psi\omega$	$\psi\gamma$	$\psi\phi$	$\psi'\pi\pi$	$\psi'\omega$	$\psi'\gamma$	$\chi_{c\gamma}$	$p\bar{p}$	$\Lambda\bar{\Lambda}$	$\Lambda_c\bar{\Lambda}_c$	$DD$	$DD^*$	$2D^*$
X(3872)	$1^{++}$	hard	<b>N</b>	hard	—	hard	—	hard	hard	hard	hard	—	<b>MF</b>	<b>MF</b>	—
X(3940)	$0^{-+}$	hard	<b>N</b>	hard	—	hard	—	hard	hard	hard	hard	—	<b>NS</b>	<b>S</b>	—
Z(3930)	$2^{++}$	hard	<b>N</b>	hard	—	hard	—	hard	hard	hard	hard	—	<b>MF</b>	<b>MF</b>	—
Y(4140)	$J^{P+}$	hard	<b>N</b>	hard	<b>N</b>	hard	—	hard	hard	hard	hard	—	<b>MF</b>	<b>MF</b>	<b>MF</b>
X(4160)	$0^{P+}$	hard	<b>N</b>	hard	<b>N</b>	hard	—	hard	hard	hard	hard	—	<b>MF</b>	<b>S</b>	<b>MF</b>
X(4350)	$J^{P+}$	hard	<b>N</b>	hard	<b>N</b>	hard	<b>N</b>	hard	hard	hard	hard	hard	<b>MF</b>	<b>MF</b>	<b>MF</b>