



## Non-perturbative decoupling of massive fermions

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FOR FUNDAMENTAL PHYSICS

CERN Lattice Coffee Talk, December 13, 2022



HELSINGIN YLIOPISTO  
HELSINGFORS UNIVERSITET  
UNIVERSITY OF HELSINKI

## 1. Motivation

- > flavor-number-dependency of running coupling (massless fermions)
- > fermion mass effects

## 2. Lattice methods

- > used data
- > gradient flow running coupling
- > finite volume effects

## 3. Results I

- > lattice gradient flow coupling data fitted

## 4. Using different renormalization schemes

- > massive GF scheme 1-loop beta function
- > relating the massive GF and massive BF-MOM schemes

## 5. Results II

- > Comparing lattice GF data with 2-loop BF-MOM
- > Potential effects from 3-loop non-universality

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## Beta function and running coupling in $SU(N)$ gauge theories

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$$\mu \frac{dg^2}{d\mu} = \beta(g^2)$$

> running gauge coupling  $g^2$  at energy scale  $\mu$ .

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[F. Herzog et al. JHEP02(2017)090]

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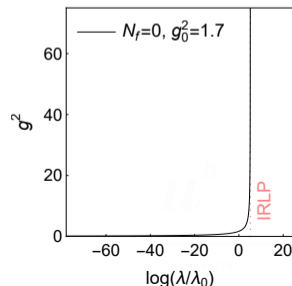
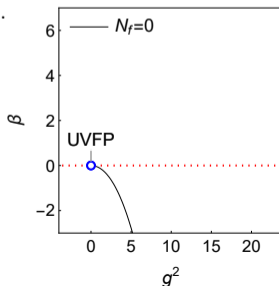
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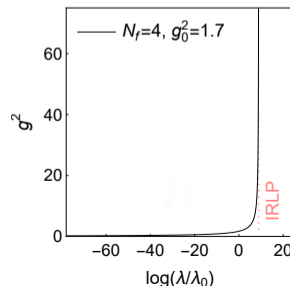
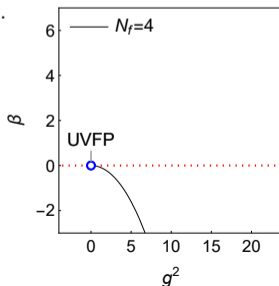
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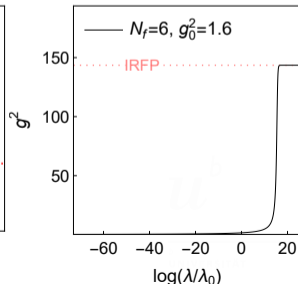
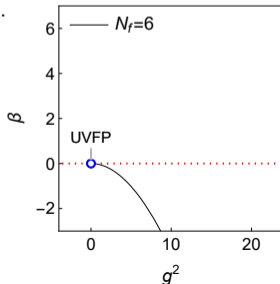
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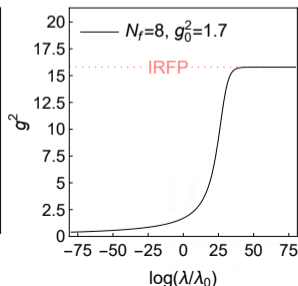
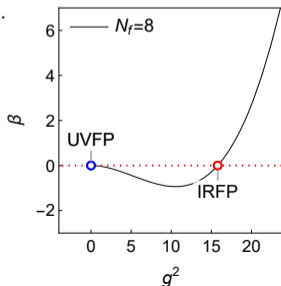
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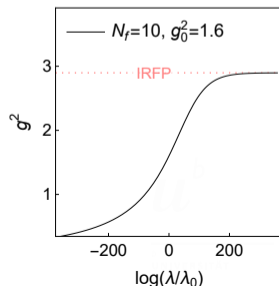
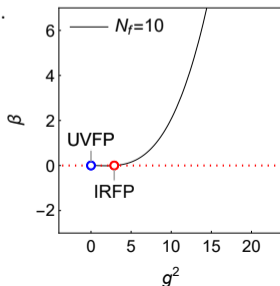
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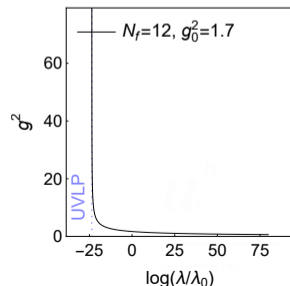
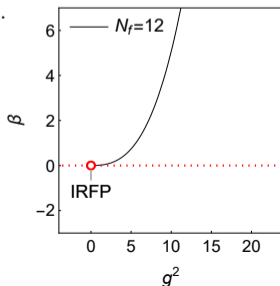
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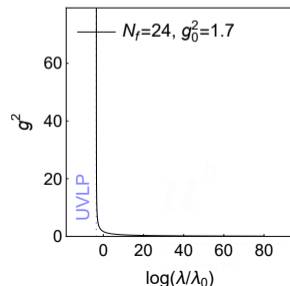
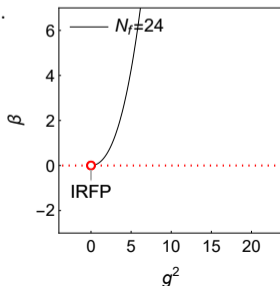
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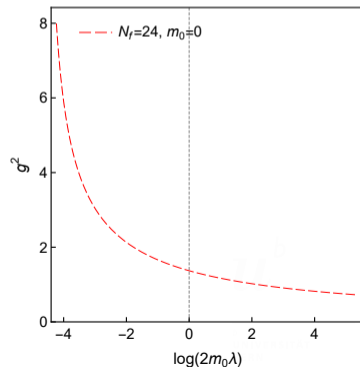
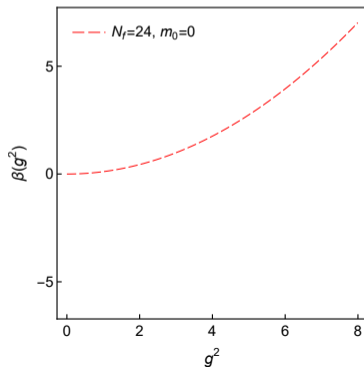
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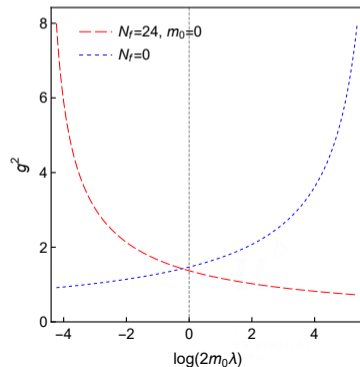
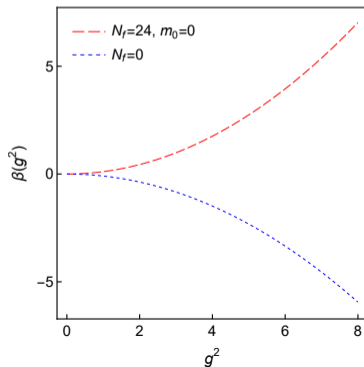
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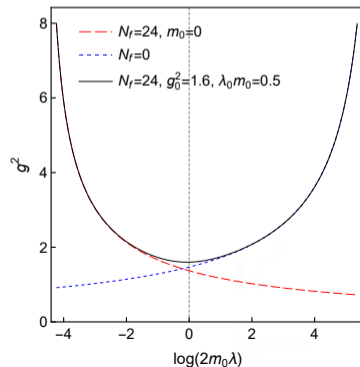
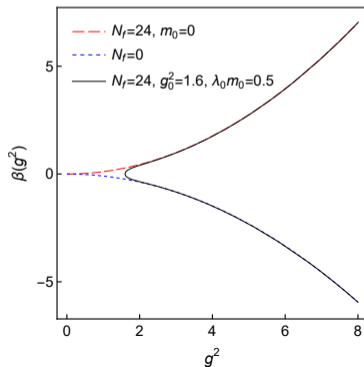
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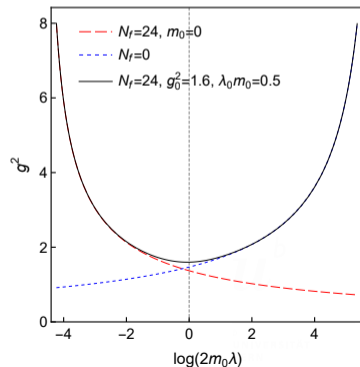
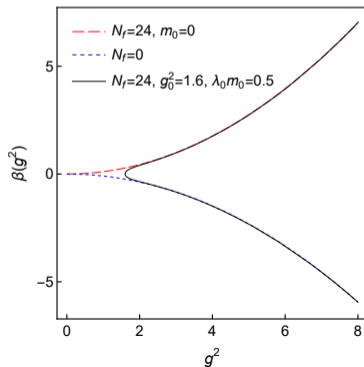
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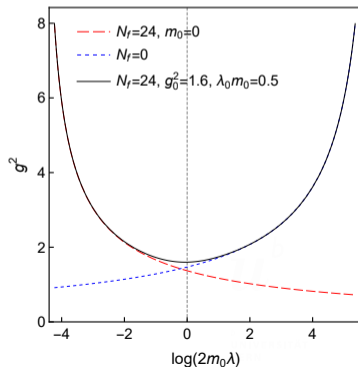
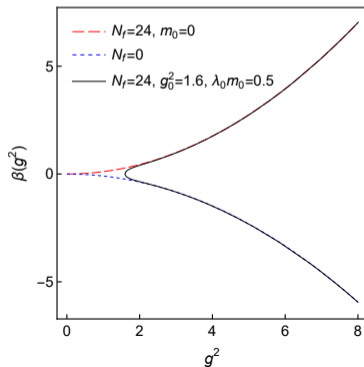
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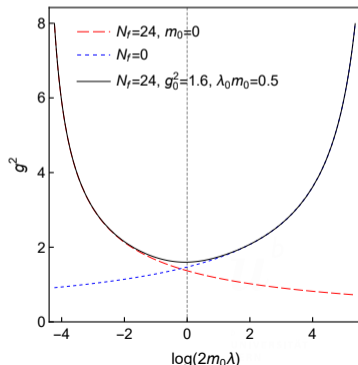
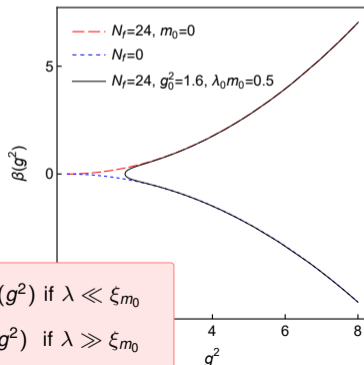
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- 2-loop BF-MOM scheme:

[Jegerlehner & Tarasov NPB549(1999)]

- > pole mass  $m = m_0$  (not running)

$$> \beta_{\text{BFM}, N_f=24}^{(2)}(g^2, \lambda m_0) = \begin{cases} \beta_{\overline{\text{MS}}, N_f=24}^{(2)}(g^2) & \text{if } \lambda \ll \xi_{m_0} \\ \beta_{\overline{\text{MS}}, N_f=0}^{(2)}(g^2) & \text{if } \lambda \gg \xi_{m_0} \end{cases}$$



# 1. Motivation

## Effect of non-zero fermion mass

- $SU(N)$  with  $N_f$  degenerate, fundamental fermions of mass  $m \Rightarrow$  additional scale  $m$ , resp.  $\xi_m \sim 1/m$  in system.

Expected asymptotic behavior:

- > for  $\lambda \ll \xi_m$  fermion mass irrelevant  $\Rightarrow \beta_{N_f, m>0}(g^2) \sim \beta_{N_f, m=0}(g^2)$
  - > for  $\lambda \gg \xi_m$  fermions decouple  $\Rightarrow \beta_{N_f, m>0}(g^2) \sim \beta_{N_f=0, m=0}(g^2)$
- } massive scheme interpolates

- General mass-dependent scheme:

$$\frac{dg^2}{d \log(\lambda)} = -\beta(g^2, \lambda m)$$

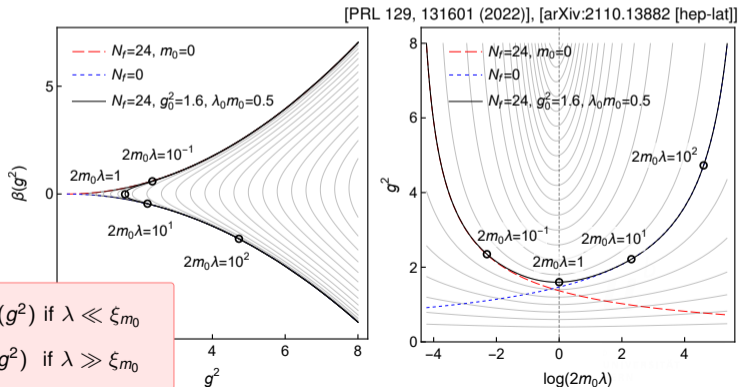
$$\frac{d \log(m)}{d \log(\lambda)} = \gamma(g^2, \lambda m)$$

- 2-loop BF-MOM scheme:

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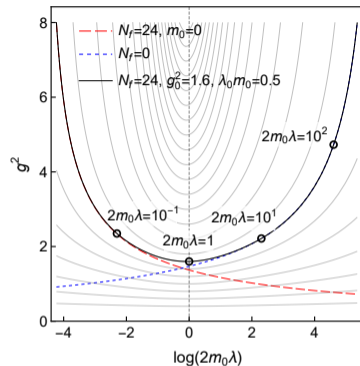
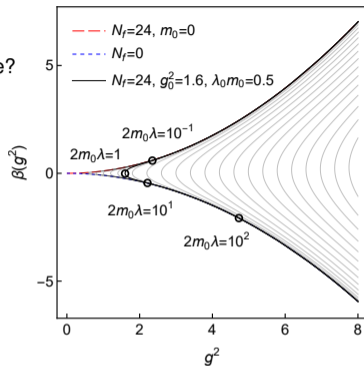
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## 2. Lattice methods

### Running coupling on the lattice

- Can this behavior be observed on the lattice?



[PRL 129, 131601 (2022)], [arXiv:2110.13882 [hep-lat]]

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### Running coupling on the lattice

#### ■ Can this behavior be observed on the lattice?

→ can use HMC data from previous study:

[Rantaharju et al. PRD 104, 114504 (2021)]

> SU(2) lattice gauge theory with  $N_f = 24$  massive Wilson fermions

> lattices of size  $V = L^4$  ( $L = 32a, 40a, 48a$ ) with periodic (gauge), resp. time-antiperiodic (fermions) BC

> lattice action:  $S = S_G(U) + S_F(V) + c_{\text{SW}} S_{\text{SW}}(V)$

→  $U$ : fundamental, unsmeared gauge link matrices

→  $V$ : HEX-smearred gauge link matrices (corresponding to  $U$ )

[Capitani et al. JHEP11(2006)028]

→  $S_G, S_F$ : Wilson gauge and Wilson fermion actions

→  $c_{\text{SW}} S_{\text{SW}}$ : clover term with Sheikholeslami-Wohlert coefficient  $c_{\text{SW}} = 1$

> 3 values of inverse gauge coupling,  $\beta \in \{ -0.25, 0.001, 0.25 \}$

> for each value of  $\beta$  4-6 different values of fermion hopping parameter  $\kappa$

> PCAC quark mass:  $a m_q = \frac{(\partial_4^* + \partial_4) f_A(x_4)}{4 f_P(x_4)} \Big|_{x_4=L/2}$  (axial and pseudo-scalar current correlators  $f_A, f_P$ )

> gradient flow running coupling (clover-energy, Lüscher-Weisz action)



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> GF scale  $\lambda = \sqrt{8t}$ , after flow time  $t$

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→ slightly modified: summing over lattice instead of continuum momenta

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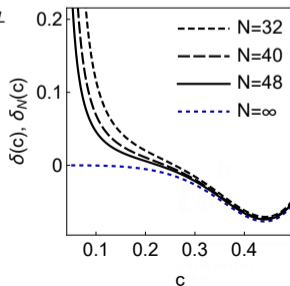
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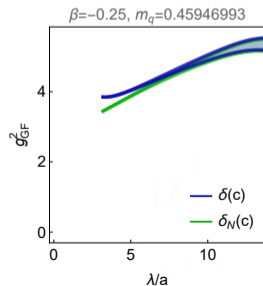
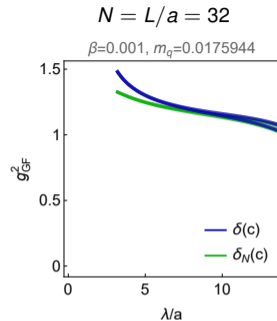
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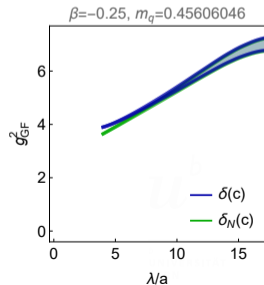
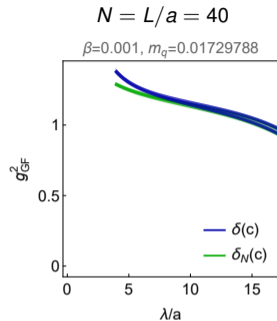
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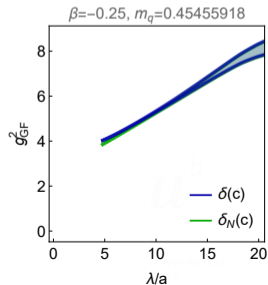
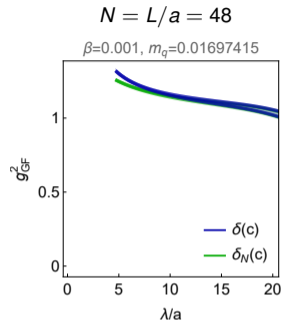
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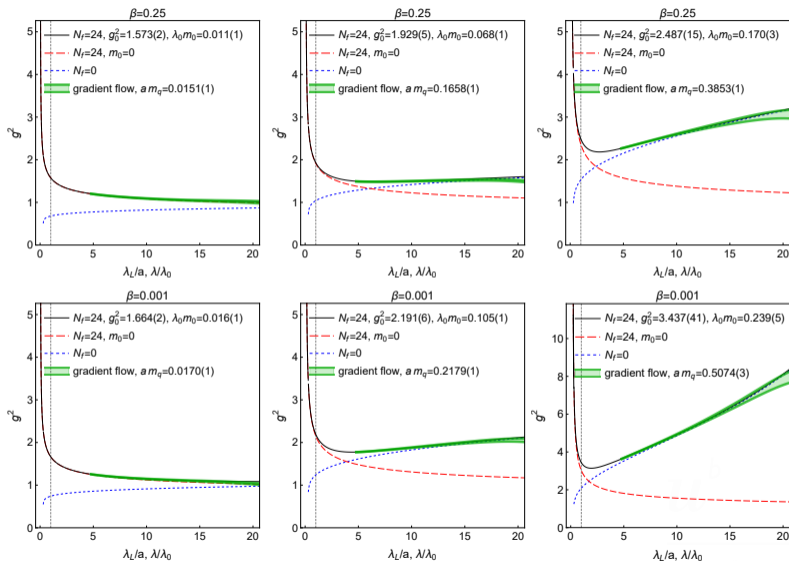
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- Will consider  $g_{\text{GF}}^2(\lambda_L, L)$  as function of  $\lambda_L$  at fixed  $L$  (no step scaling).

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#### Lattice GF-data

■  $V = (48a)^4$  lattice:





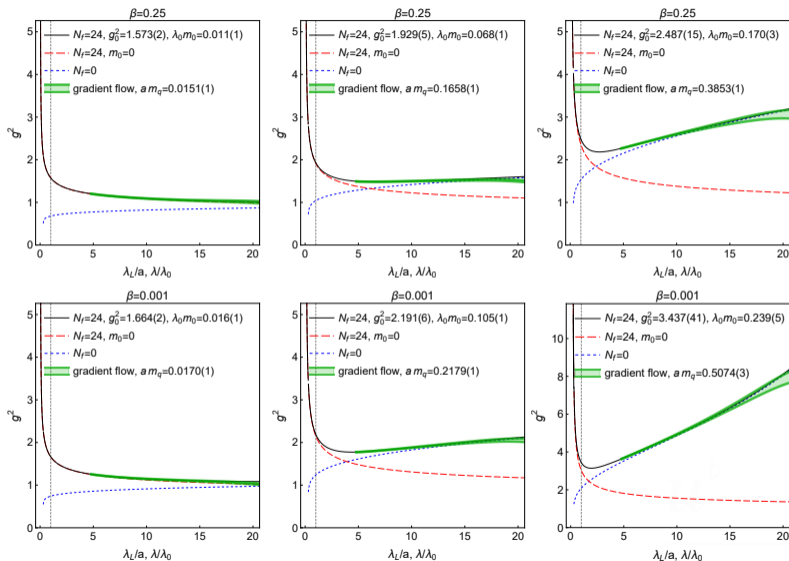
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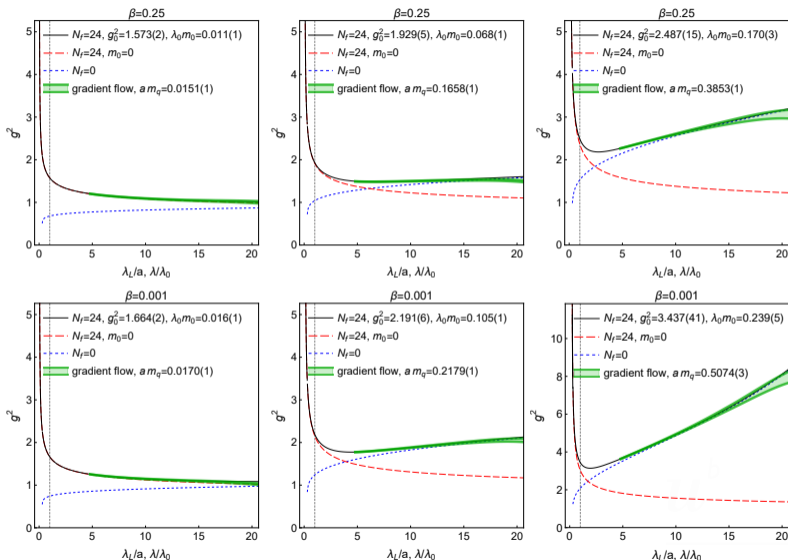


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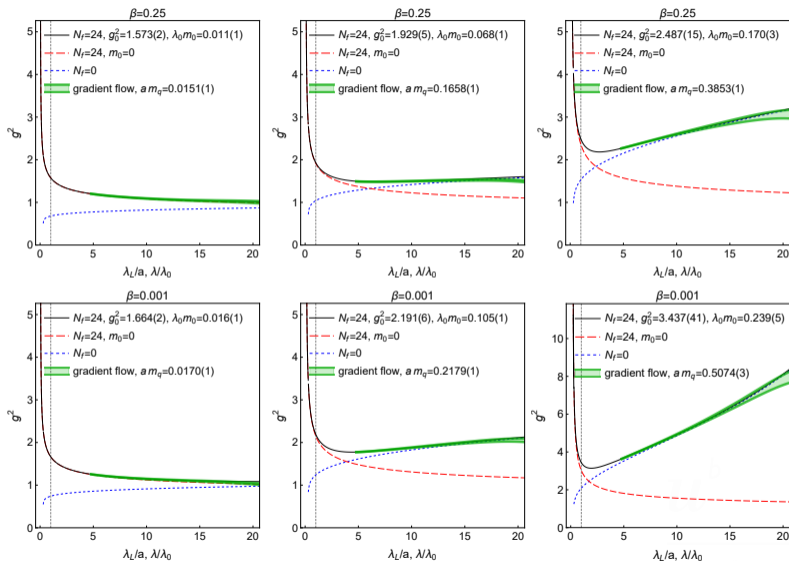
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( $a m_q$  unren. (no pole mass))



## 4. Using different renormalization schemes

### Perturbative massive GF scheme beta function

- Continuum GF running coupling:  $u_{\text{GF}}(\lambda) = g_{\text{GF}}^2(\lambda) = \frac{2\pi^2\lambda^4 \langle E(\lambda) \rangle}{3(N^2 - 1)}$  with  $\lambda = \sqrt{8t}$

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■ use 1-loop  $\overline{\text{MS}}$  result for leading  $m_q$ -contribution to  $\langle E(t) \rangle$  [Harlander & Neumann JHEP06(2016)161]

$$\rightarrow u_{\text{GF}}(\lambda) = u_{\overline{\text{MS}}}(\lambda_0) \left( 1 + k_1(\lambda/\lambda_0) \frac{u_{\overline{\text{MS}}}(\lambda_0)}{(4\pi)^2} + l_1(\lambda m(\lambda_0)) \frac{u_{\overline{\text{MS}}}(\lambda_0)}{(4\pi)^2} + \mathcal{O}(u_{\overline{\text{MS}}}^2(\lambda_0)) \right) (*)$$

> old:  $k_1(\lambda/\lambda_0) = (2 \log(\lambda/\lambda_0) + \gamma_E) \beta_0 + N \left( \frac{52}{9} - 3 \log(3) \right) - N_f \left( \frac{4}{9} - \frac{4}{3} \log(2) \right)$  [Lüscher JHEP08(2010)071]

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■ define GF beta function in terms of  $\lambda$ -derivative of (\*): [Harlander PoS(LATTICE2021)489]

$$\rightarrow \beta_{\text{GF}} = -\lambda \frac{du_{\text{GF}}}{d\lambda} = -\left( \lambda \frac{dk_1(\lambda/\lambda_0)}{d\lambda} + \lambda \frac{dl_1(\lambda m(\lambda_0))}{d\lambda} \right) \frac{u_{\overline{\text{MS}}}^2(\lambda_0)}{(4\pi)^2} + \mathcal{O}(u_{\overline{\text{MS}}}^3) (**)$$

> invert (\*):  $u_{\overline{\text{MS}}}(\lambda_0) = u_{\text{GF}}(\lambda) \left( 1 - k_1(\lambda/\lambda_0) \frac{u_{\text{GF}}(\lambda)}{(4\pi)^2} - l_1(\lambda m(\lambda_0)) \frac{u_{\text{GF}}(\lambda)}{(4\pi)^2} + \mathcal{O}(u_{\text{GF}}^2(\lambda)) \right)$

> and use it in (\*\*)  $\Rightarrow \beta_{\text{GF}}(u_{\text{GF}}, \lambda m) = -\underbrace{\left( \lambda \frac{dk_1(\lambda/\lambda_0)}{d\lambda} + \lambda \frac{dl_1(\lambda m(\lambda_0))}{d\lambda} \right)}_{2\beta_{0,\text{GF}}} \frac{u_{\text{GF}}^2}{(4\pi)^2} + \mathcal{O}(u_{\text{GF}}^3)$

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$$\blacksquare \beta_{\text{GF}}(u_{\text{GF}}, \lambda m) = -2 \beta_{0,\text{GF}}(x) \frac{u_{\text{GF}}^2}{(4\pi)^2} + \mathcal{O}(u_{\text{GF}}^3)$$

$$\text{with } \beta_{0,\text{GF}}(x) = \beta_0 + \frac{4}{3} T_R N_f x \frac{d\Omega_{1q}(x)}{dx} \quad \text{and} \quad x = -1/(2\lambda m)^2$$

> at 1-loop:  $m = m_0$  (const.)

$$> \beta_{0,\text{GF},N_f}(x) = \begin{cases} \beta_{0,\overline{\text{MS}},N_f=24} & \text{if } \lambda \ll 1/m_0 \\ \beta_{0,\overline{\text{MS}},N_f=0} & \text{if } \lambda \gg 1/m_0 \end{cases}$$

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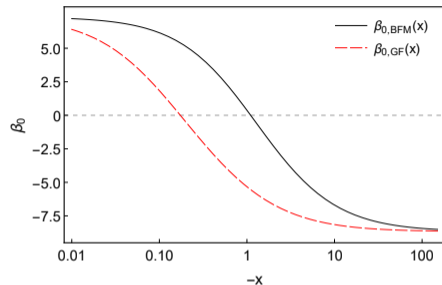
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> decoupling not near  $2\lambda m_0 \approx 1 \Rightarrow$  mismatch between  $\lambda$  and  $m_0$



[PRL 129, 131601 (2022)], [arXiv:2110.13882 [hep-lat]]



## 4. Using different renormalization schemes

### Relating GF and BF-MOM (BFM) scheme

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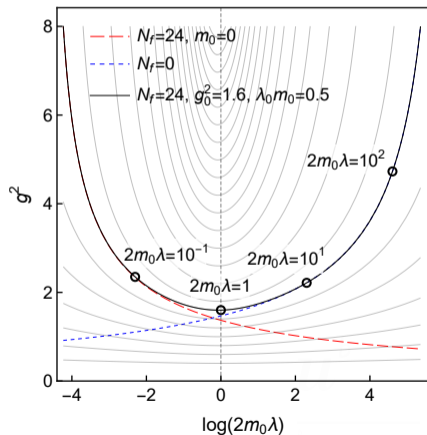
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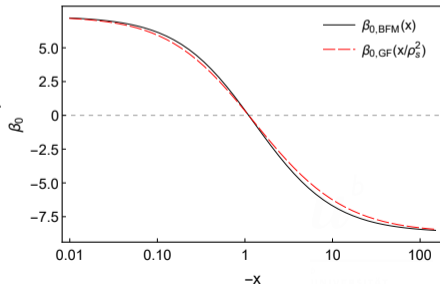
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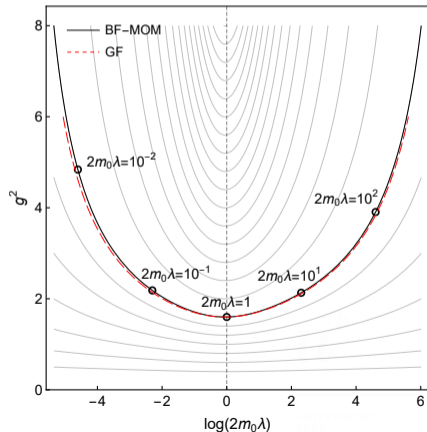
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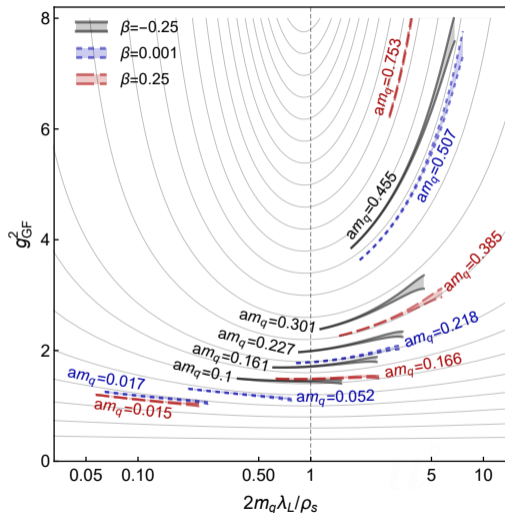
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### Lattice GF-data compared to 2-loop BF-MOM (no fit)

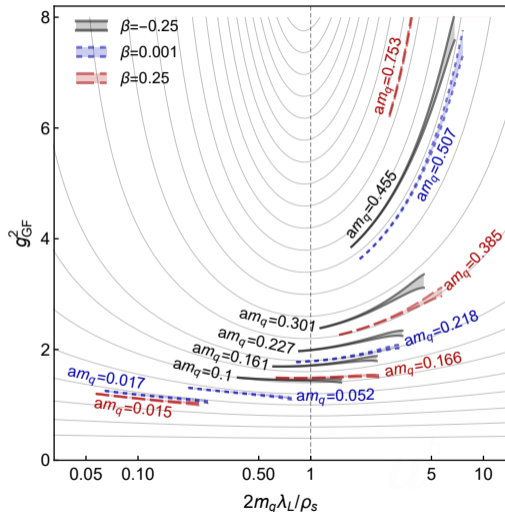
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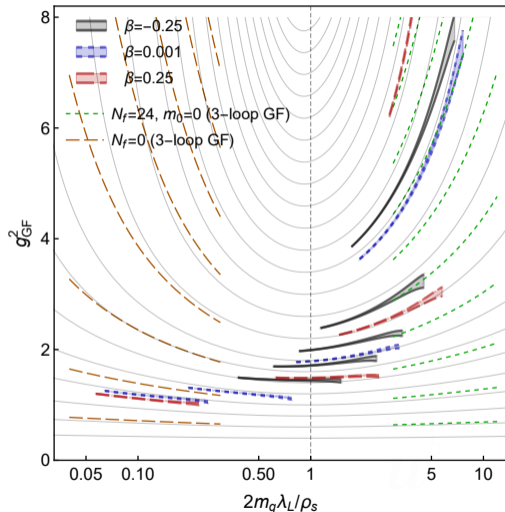


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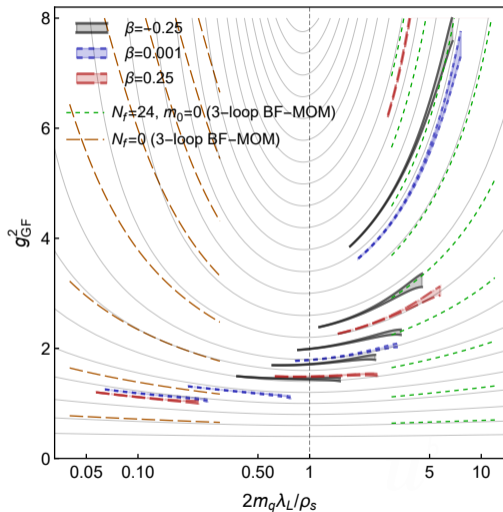


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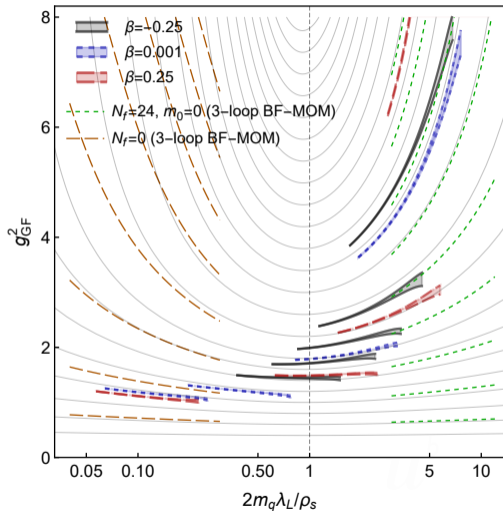


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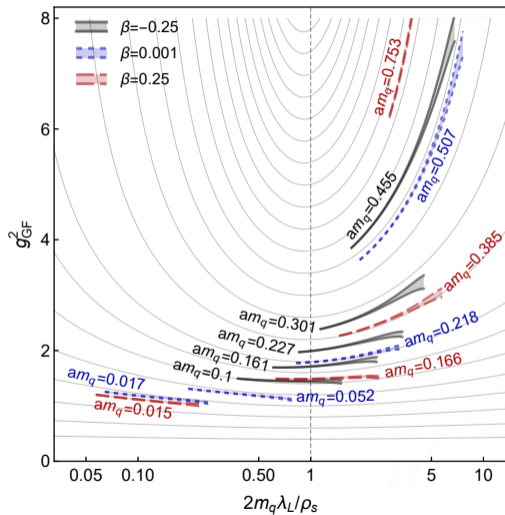


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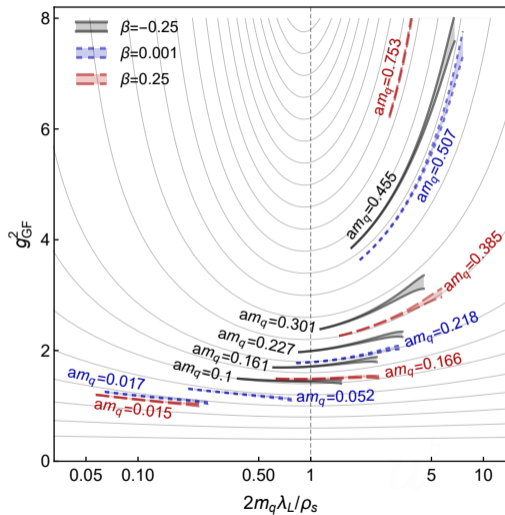




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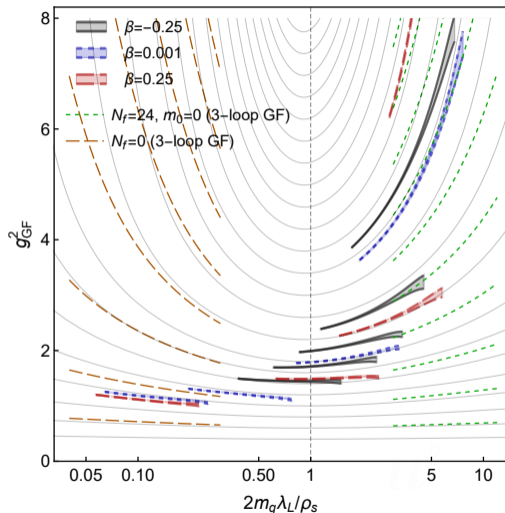
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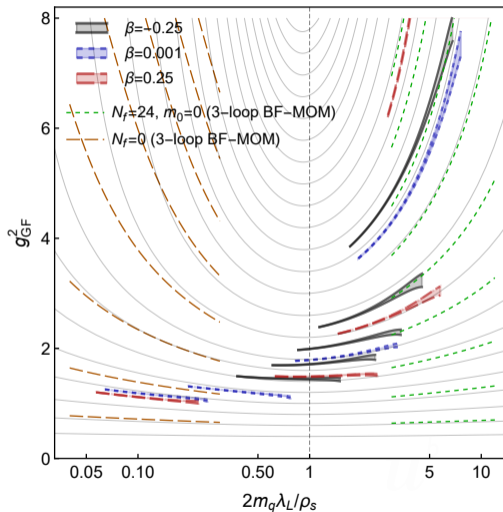
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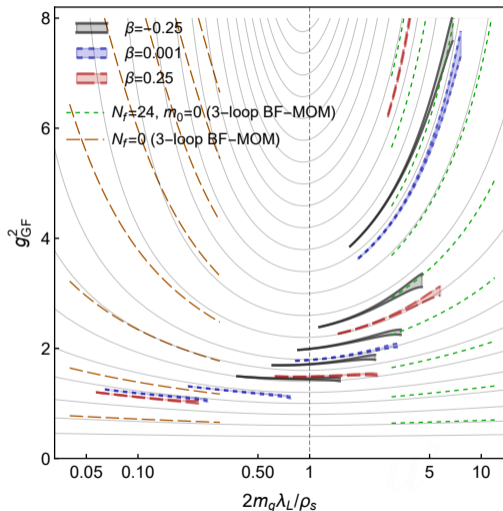
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## 7. Backup

### 3-loop non-universal effects

- Comparisons of the 3-loop beta function for SU(2) gauge theory with  $N_f = 24$  massless flavors (left) and  $N_f = 0$  flavors (right) in the  $\overline{\text{MS}}$  (solid, black), BF-MOM (long dashes, red) and GF scheme (short dashes, blue). For comparison also the corresponding 2-loop beta function is shown (dot-dashed, black).

