

Contact interactions from high-mass Drell-Yan tails

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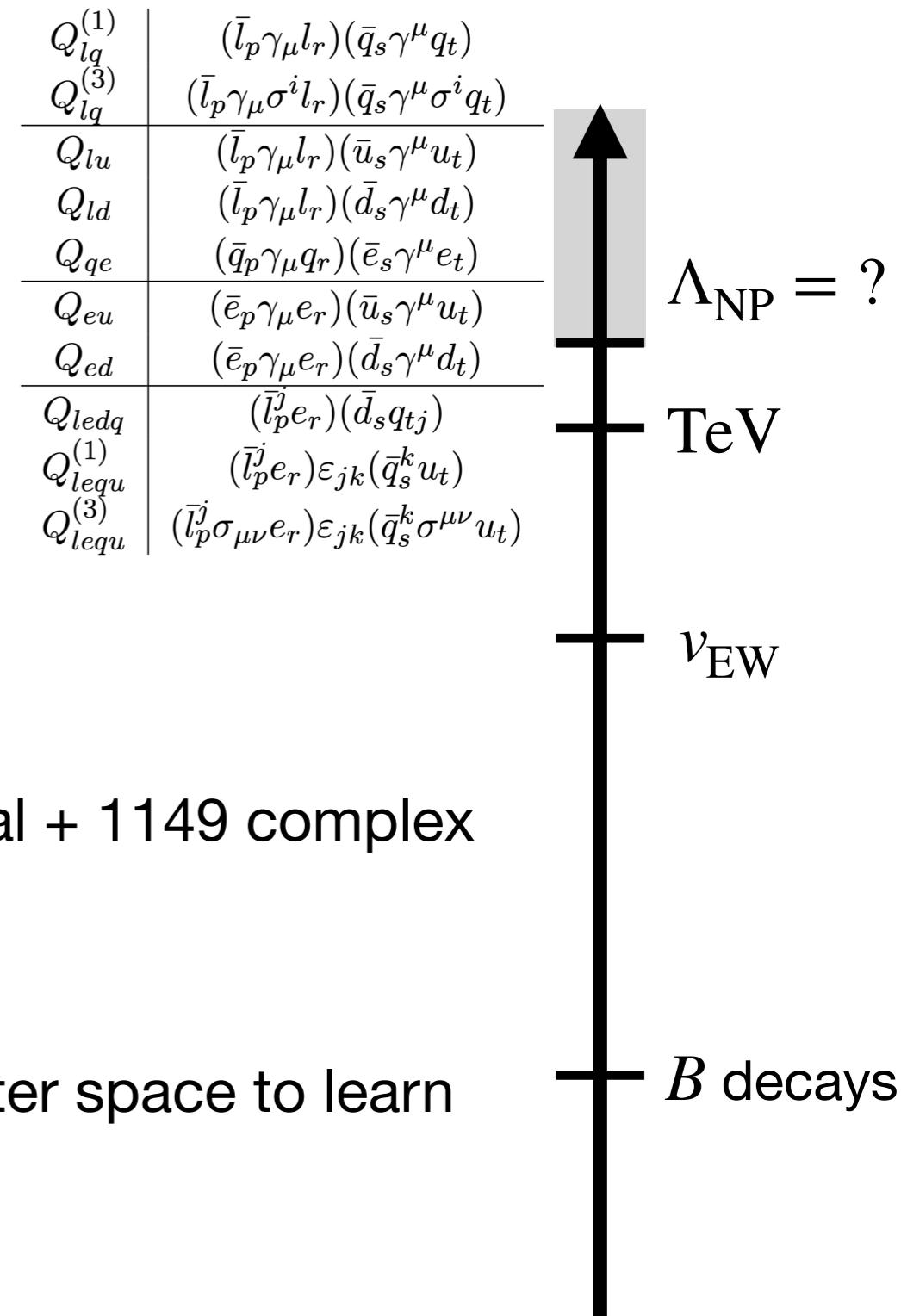
SMEFT

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \sum_i C_i Q_i$$

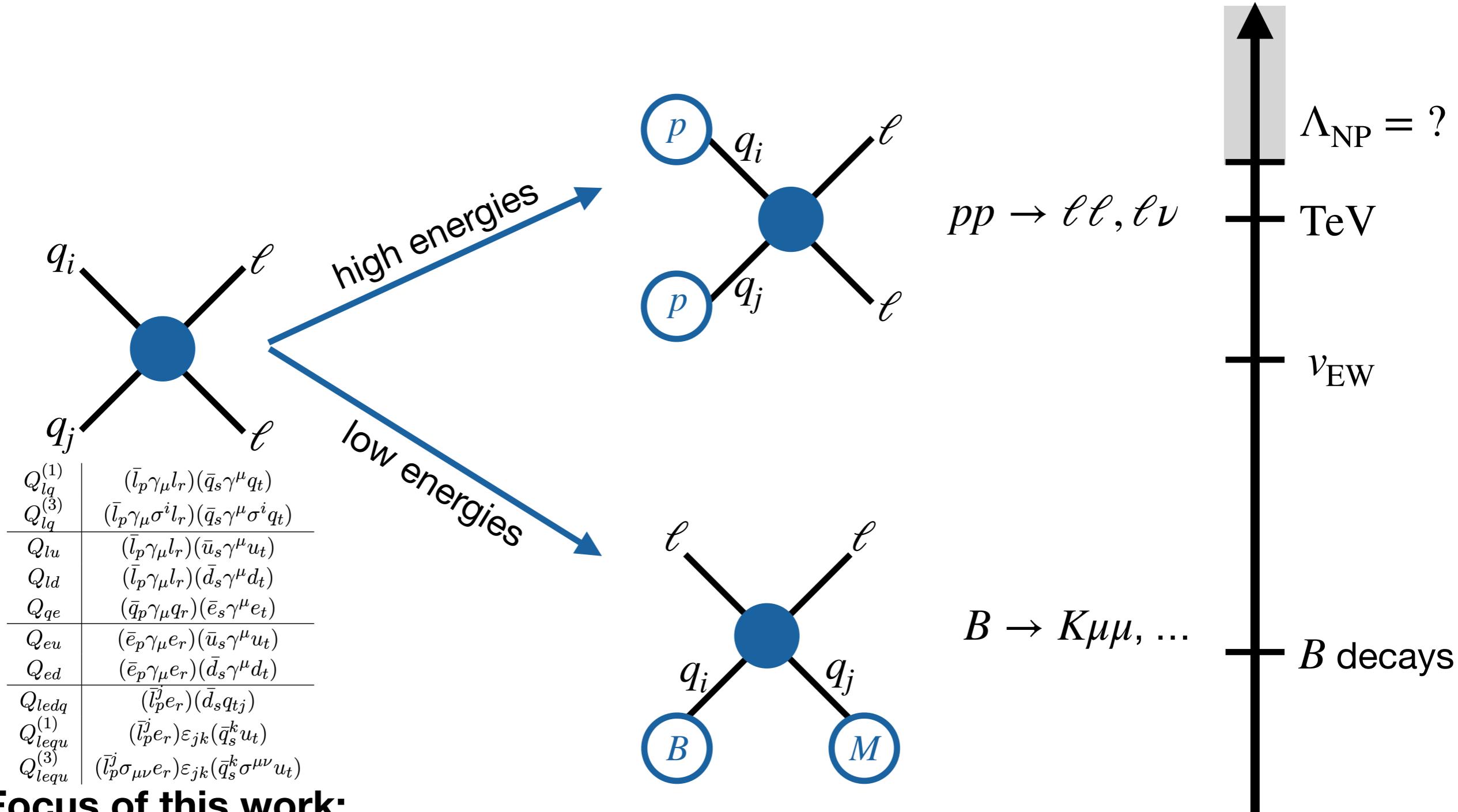
- Absence of NP in direct searches suggests a mass gap between v_{EW} and Λ_{NP}
- SMEFT: Low energy limit of generic heavy NP
- 59 operators for a single generation at dim. 6
- Flavour opens a huge parameter space: 1350 real + 1149 complex parameters
- Overarching goal:
efficient scanning, over-constraining the parameter space to learn about NP

Focus of this work:

Dim. 6 contact interactions entering high mass Drell-Yan tails
and interplay with low-energy B decays to light leptons



Drell-Yan vs B-decays



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A word about flavio

- Open-source python package flav-io.github.io
- Calculator of hundreds of observables in the SM and in dim. 6 EFTs
- Contains a large database of experimental measurements, allows for easy construction of likelihoods
- Interface with wilson for running and matching of WCs above and below the EW scale
- Basis of smelli - package providing a global SMEFT likelihood

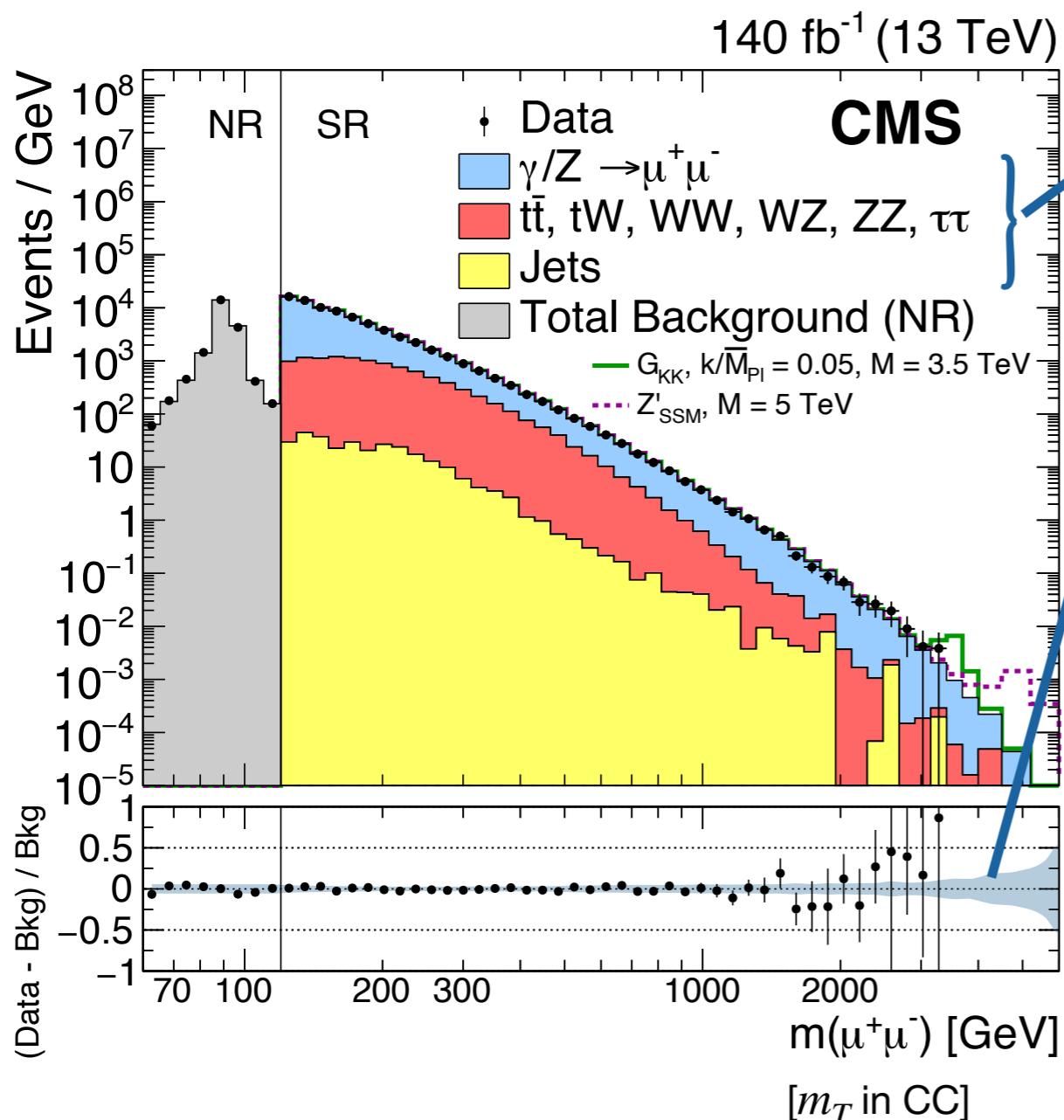
Our contributions:

- **Implementation of** theoretical predictions (dim. 6 in SMEFT) and latest experimental measurements of **neutral and charged current high-mass Drell-Yan** with light leptons $\ell = e, \mu$
- **Implementation and update of various $b \rightarrow d$ predictions and measurements**

Implementation of Drell-Yan

Experimental measurements:

We implement data ($\sim 140\text{fb}^{-1}$) from latest CMS and ATLAS searches:



	$pp \rightarrow \ell\ell$	$pp \rightarrow \ell\nu$
CMS	2103.02708	2202.06075
ATLAS	2006.12946	1906.05609

Expected # of events @ (N)NLO in QED(QCD)
including N_{DY}^{SM}

Systematic errors provided

Likelihood:

Poisson convolved with Normal (systematics)

For predictions:

We reweigh the reported expected number
of DY events with ratio

$$N_{DY}^{\text{NP+SM}} = \frac{\sigma^{\text{SM+NP}}}{\sigma^{\text{SM}}} N_{DY}^{\text{SM}}$$

⇒ implement analytical predictions
of $d\sigma(pp \rightarrow \ell\ell, \ell\nu)/dm$

A. Greljo and D. Marzocca: 1704.09015

Implementation of Drell-Yan

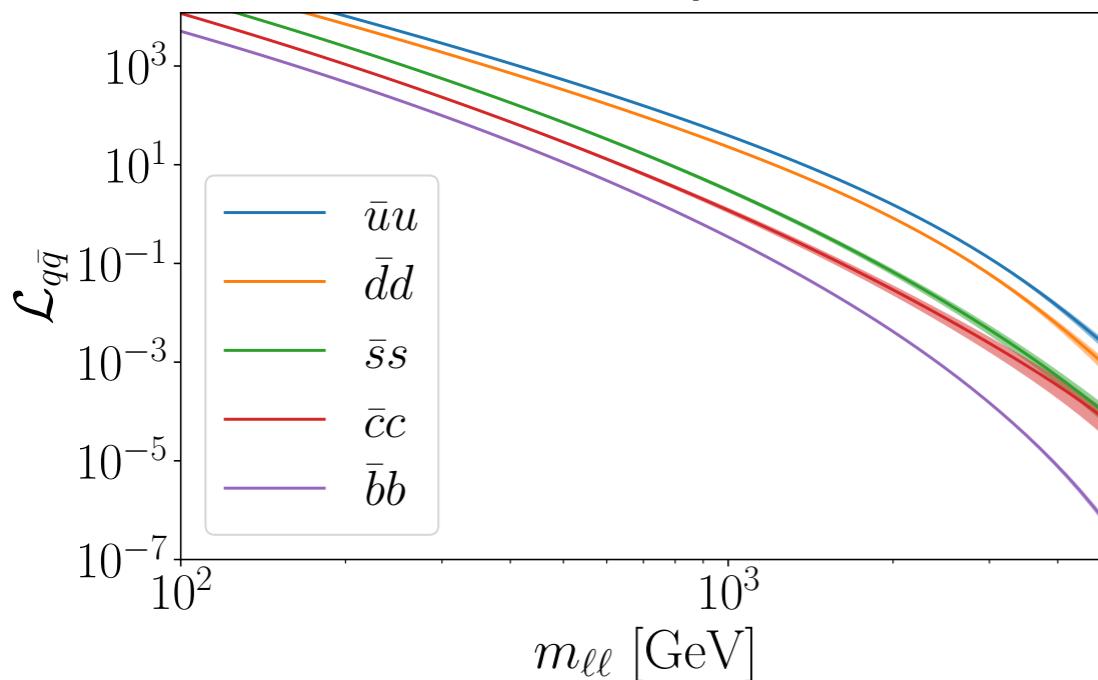
- Partonic cross-sections including all Lorentz structures, chiralities
- Implementation of PDFs in `flavio` (NNPDF 4.0), convolution of σ_{part} with luminosity functions

$$\sigma_{\text{part}}^{q\bar{q}} \sim \sum \left[\begin{array}{c} q_i \\ q_j \end{array} \right] \begin{array}{c} \nearrow \gamma/Z/W \\ \swarrow \ell \end{array} + \left[\begin{array}{c} q_i \\ q_j \end{array} \right] \begin{array}{c} \nearrow \ell \\ \text{blue circle} \end{array} \right] ^2$$

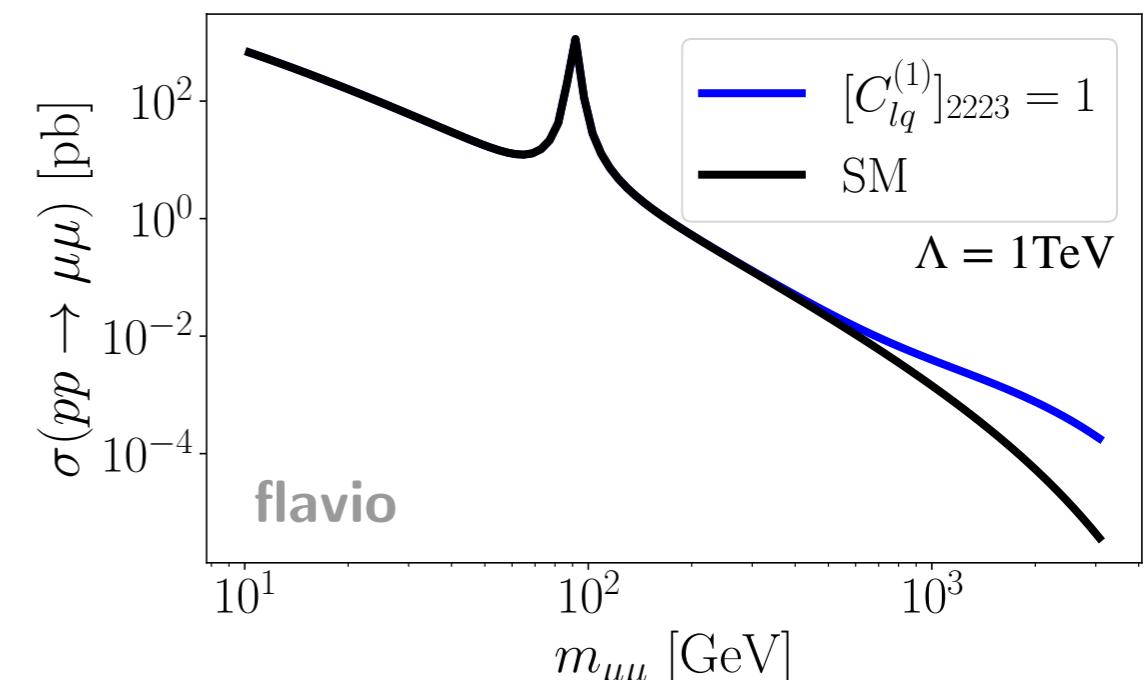
$$\sigma_{\text{had}} \sim \sum_{q\bar{q}} \mathcal{L}_{q\bar{q}} * \sigma_{\text{part}}^{q\bar{q}}$$

$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \sigma^i l_r)(\bar{q}_s \gamma^\mu \sigma^i q_t)$
Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$
Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s q_{tj})$
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$

sensitive to five quark flavours



EFT enhanced at high energies



Thorough validations against Monte Carlo simulations using Madgraph

Status of $b \rightarrow q\ell\ell$ in flavio

Current and **newly implemented** processes

$b \rightarrow s\mu\mu$

Various branching ratios, angular observables, and other observables, measured in many modes:
 $B \rightarrow K^{(*)}\mu\mu, B_s \rightarrow \phi\mu\mu, B_s \rightarrow \mu\mu, \Lambda_b \rightarrow \Lambda\mu\mu$

+ LFUV ratios $R_{K^{(*)}}$

$b \rightarrow see$

Upper limit on $B_s \rightarrow ee$ [LHCb 2020]
Inclusive $B \rightarrow X_s ee$ [BaBar 2013]
($B \rightarrow K^* ee$ at very low q^2 - LHCb)

$b \rightarrow d\mu\mu$

Upper limit on $B^0 \rightarrow \mu\mu$ [LHCb, ATLAS, CMS]

Branching ratio of $B \rightarrow \pi\mu\mu$ *[LHCb 2015]

Branching ratio of $B_s \rightarrow K^*\mu\mu$ **[LHCb 2018]

$b \rightarrow dee$

Upper limit on $B^0 \rightarrow ee$ [LHCb 2020]

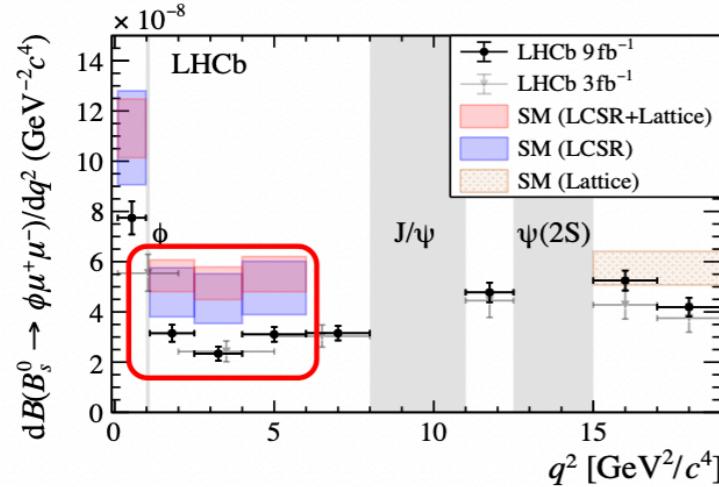
Upper limit on $B \rightarrow \pi ee$ *[Belle 2008]

*we update $B \rightarrow \pi$ **form factors** from *D. Leljak et al* (2102.07233) (lattice+LCSR)

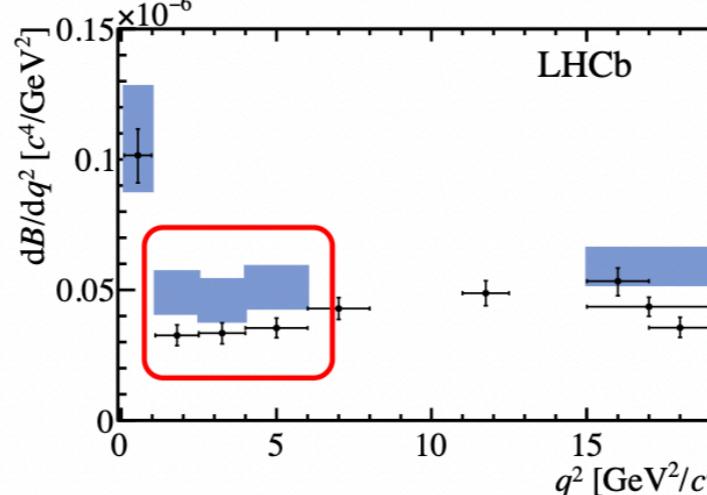
**we follow closely *R. Bause et al* (2209.04457) for treatment of resonant regions

$b \rightarrow s\mu\mu$ anomalies...

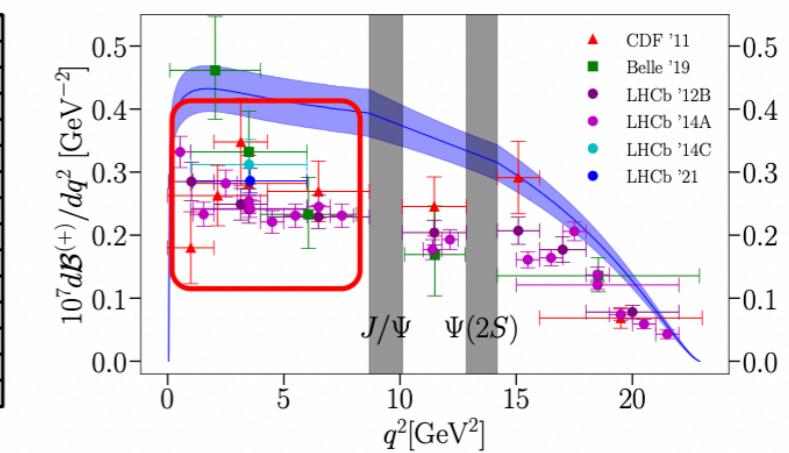
LHCb $B_s^0 \rightarrow \phi\mu^+\mu^-$ [PRL 127 (2021) 151801]



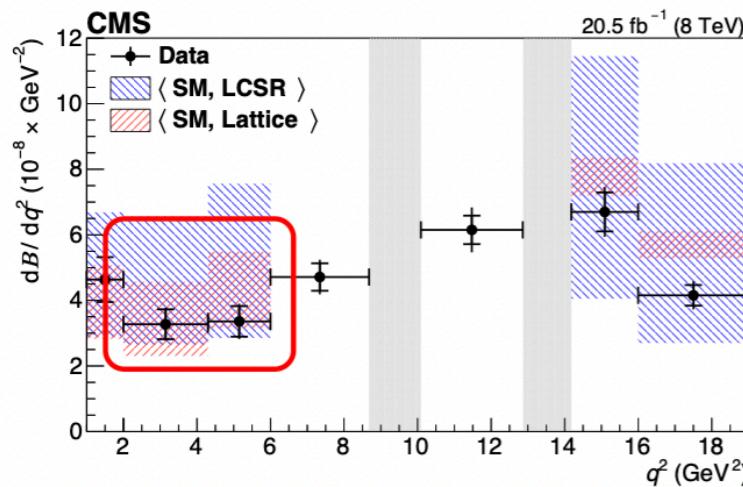
LHCb $B^0 \rightarrow K^{*0}\mu^+\mu^-$ [JHEP 11 (2016) 047]



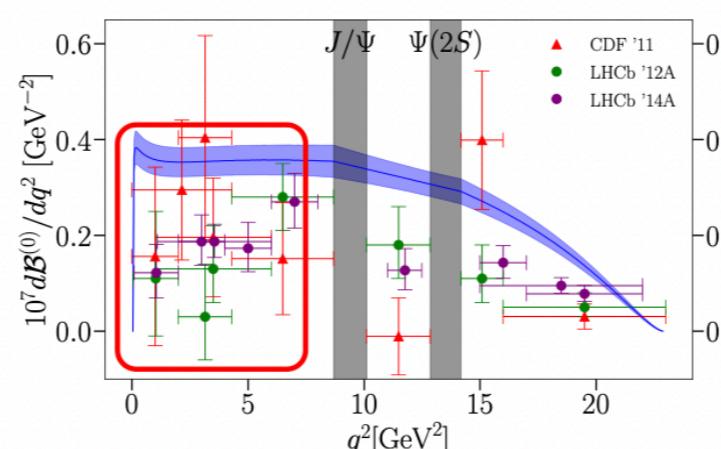
Lattice $B^+ \rightarrow K^+\mu^+\mu^-$ [arXiv:2207.13371]



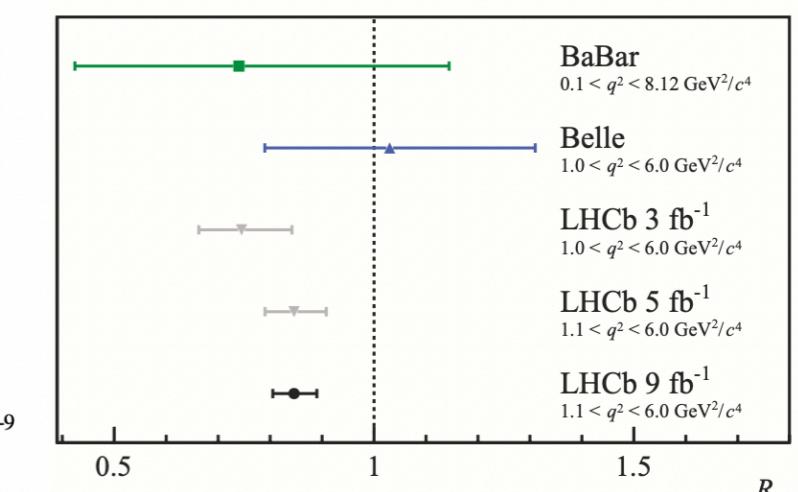
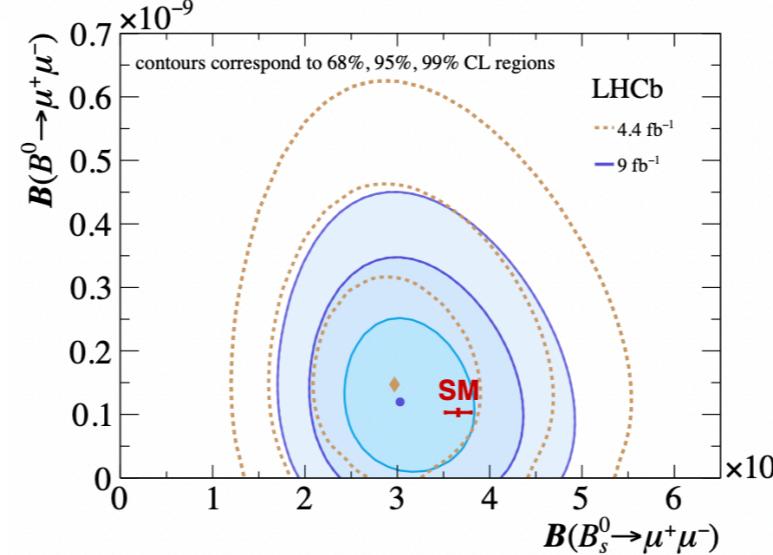
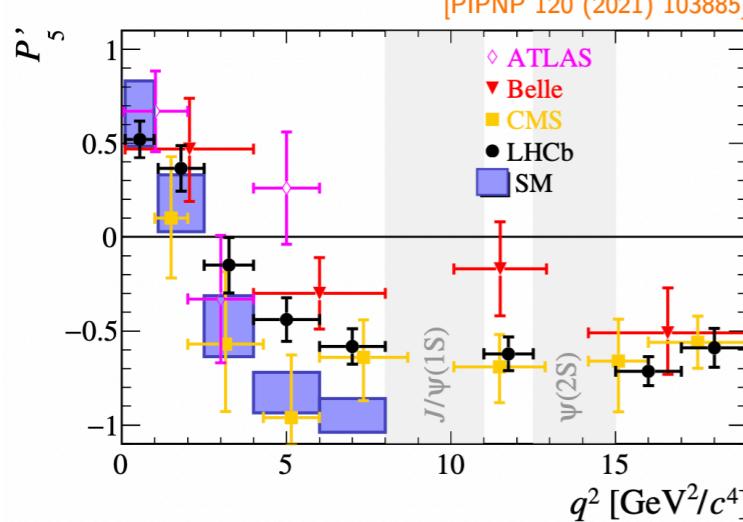
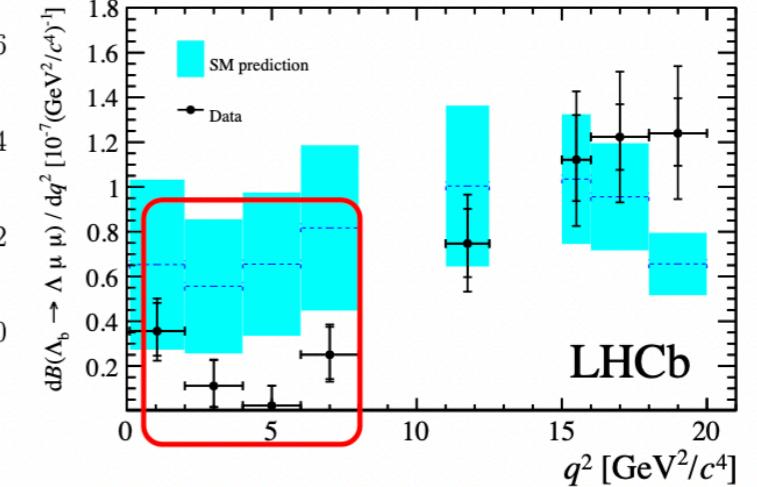
CMS $B^0 \rightarrow K^{*0}\mu^+\mu^-$ [PLB 753 (2016) 424]



Lattice $B^0 \rightarrow K^0\mu^+\mu^-$ [arXiv:2207.13371]



LHCb $\Lambda_b^0 \rightarrow \Lambda\mu^+\mu^-$ [JHEP 06 (2015) 115]

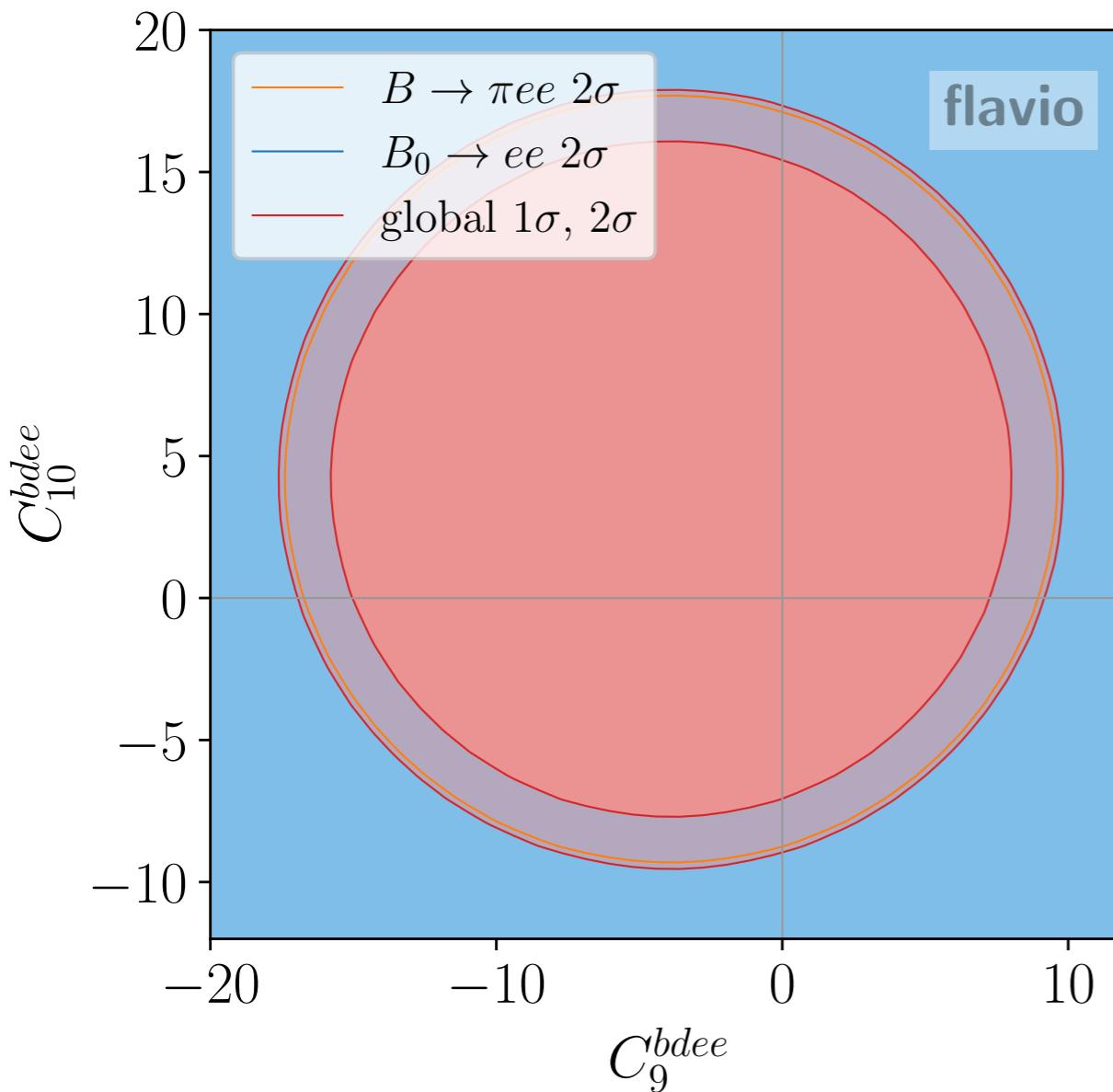


... but let us focus on $b \rightarrow dee$ for now

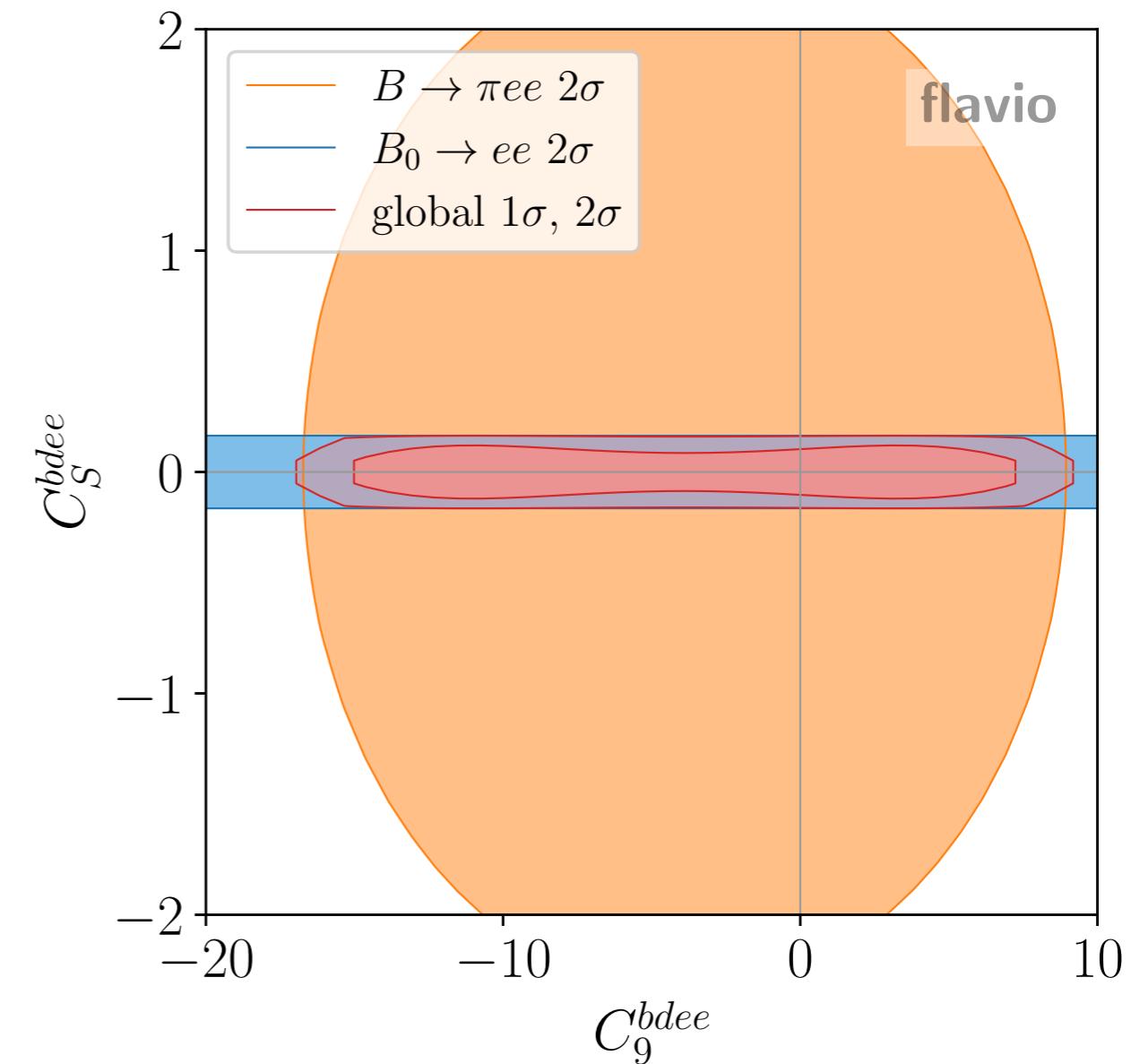
Minimalistic flavour scenario

consider only coefficients relevant for a certain decay

In WET:

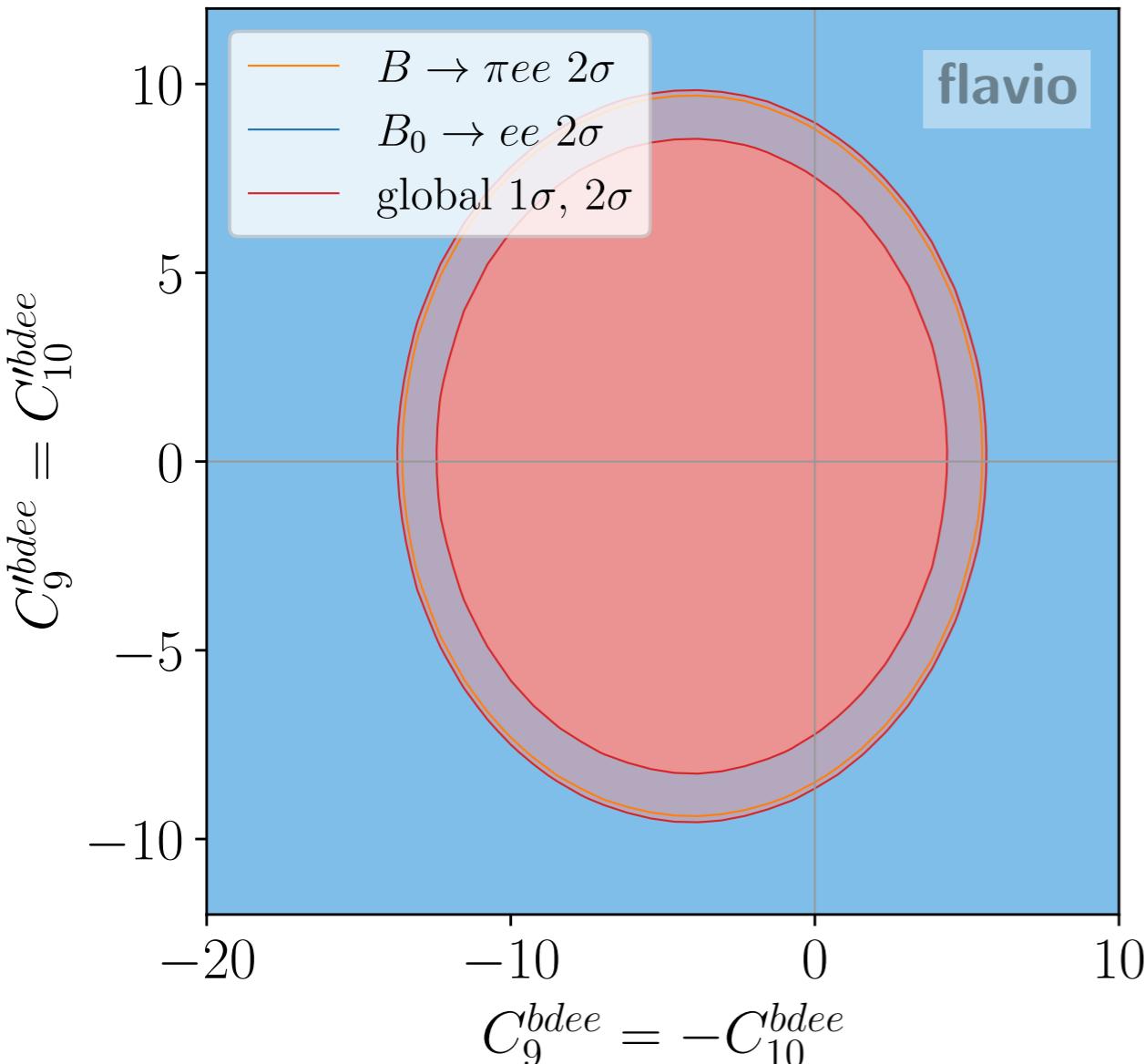


$$O_{9(10)}^{[']ijkk} \sim (\bar{q}_{L[R]}^j \gamma^\mu q_{L[R]}^i)(\bar{\ell}^k \gamma_\mu (\gamma_5) \ell^k)$$

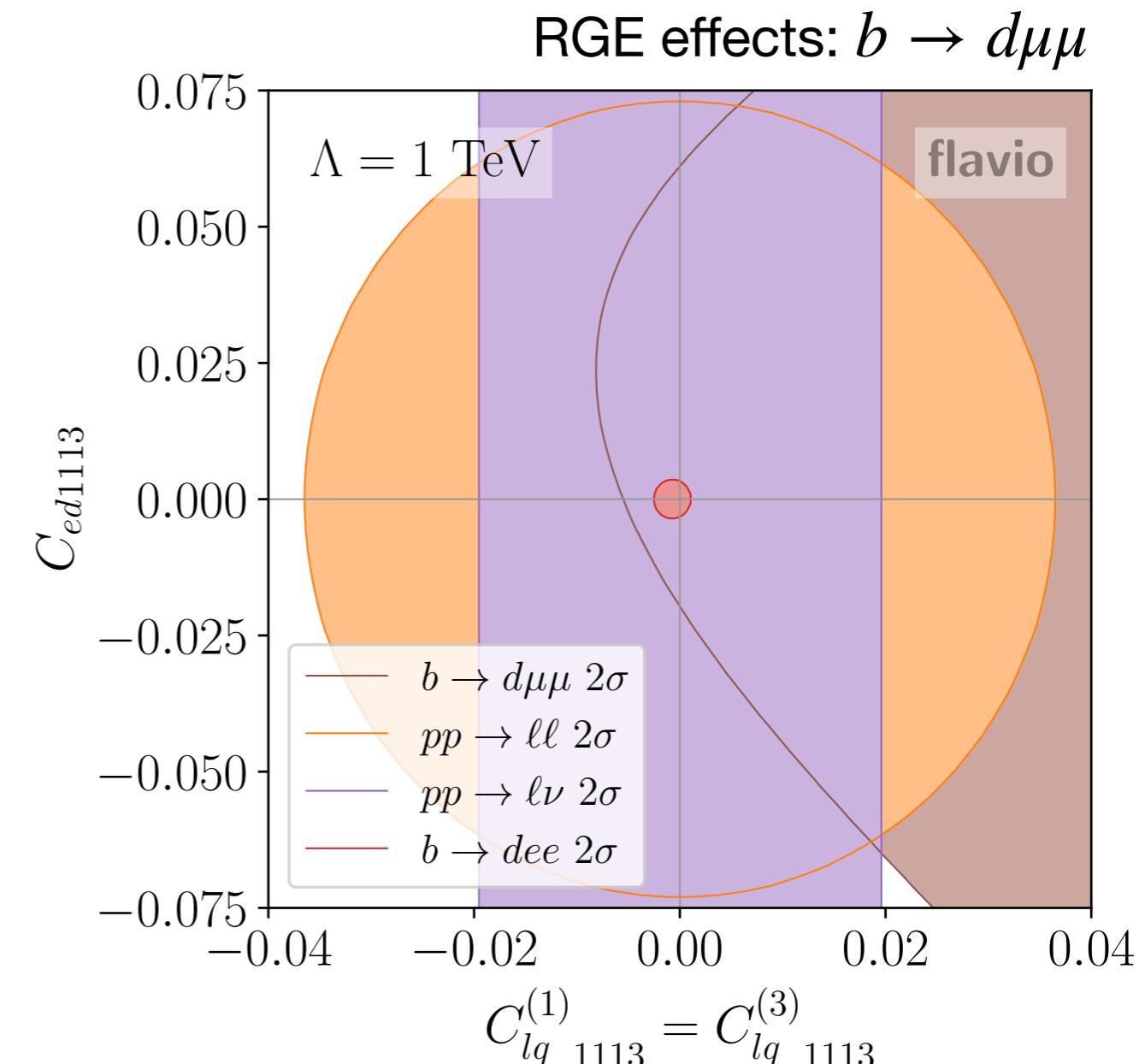


$$O_S^{ijkk} \sim (\bar{q}_L^j q_R^i)(\bar{\ell}^k \ell^k)$$

In SMEFT, confronting with high-mass Drell-Yan and $b \rightarrow q\nu\nu$ ($B \rightarrow K, \pi, \rho$)

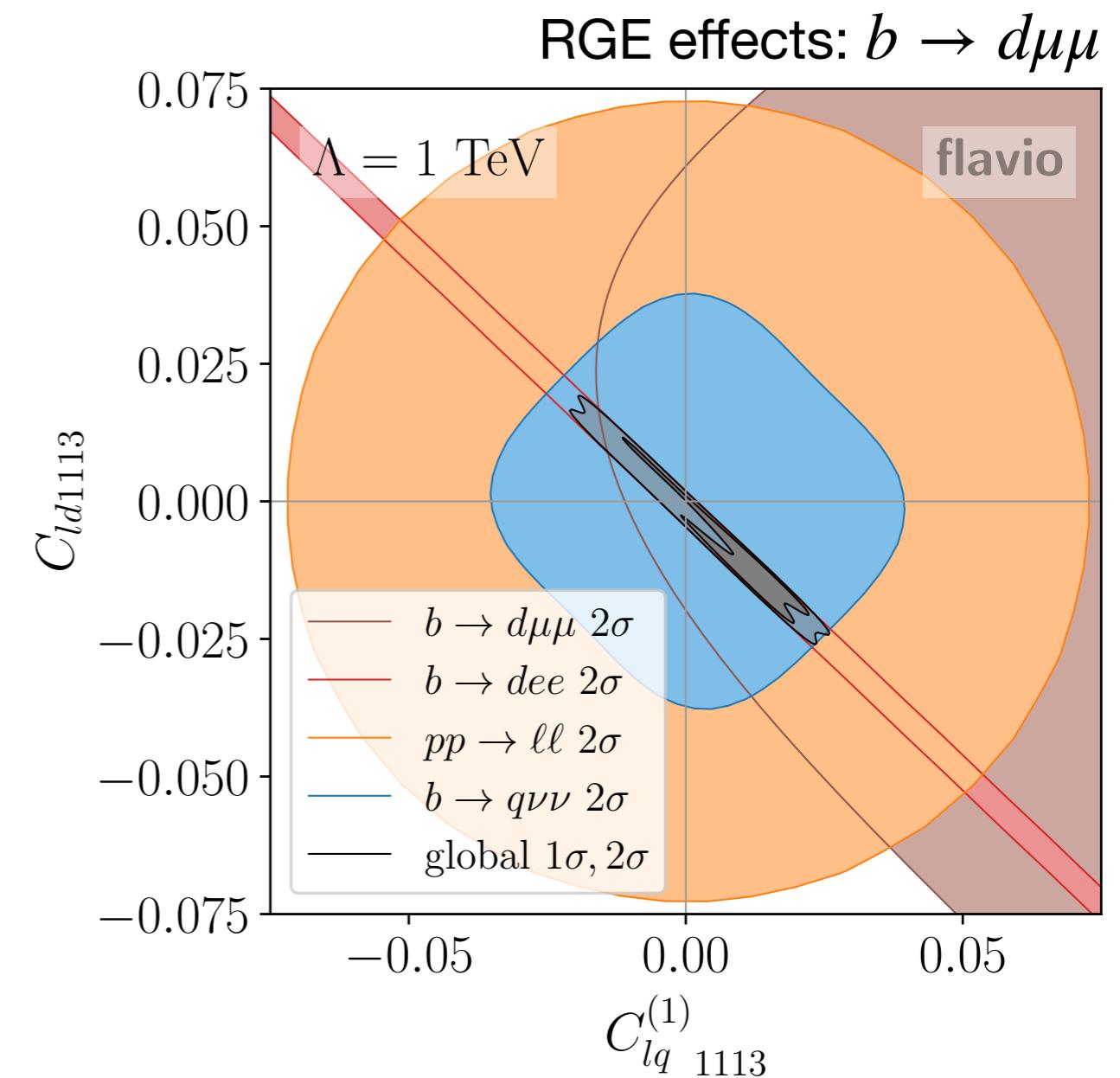
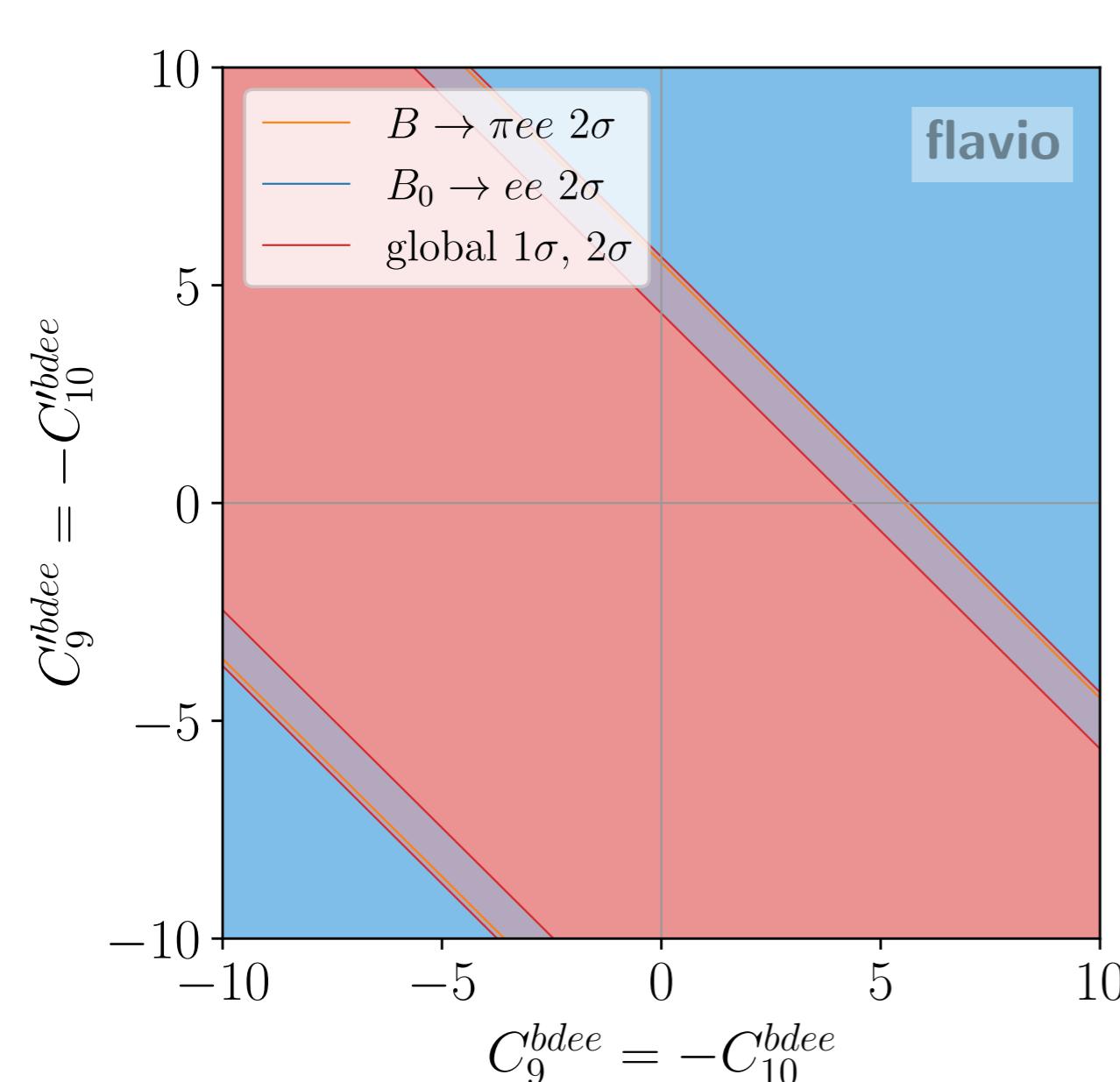


$$O_{9(10)}^{[']ijkk} \sim (\bar{q}_{L[R]}^j \gamma^\mu q_{L[R]}^i)(\bar{\ell}^k \gamma_\mu (\gamma_5) \ell^k)$$



Dominated by B decays

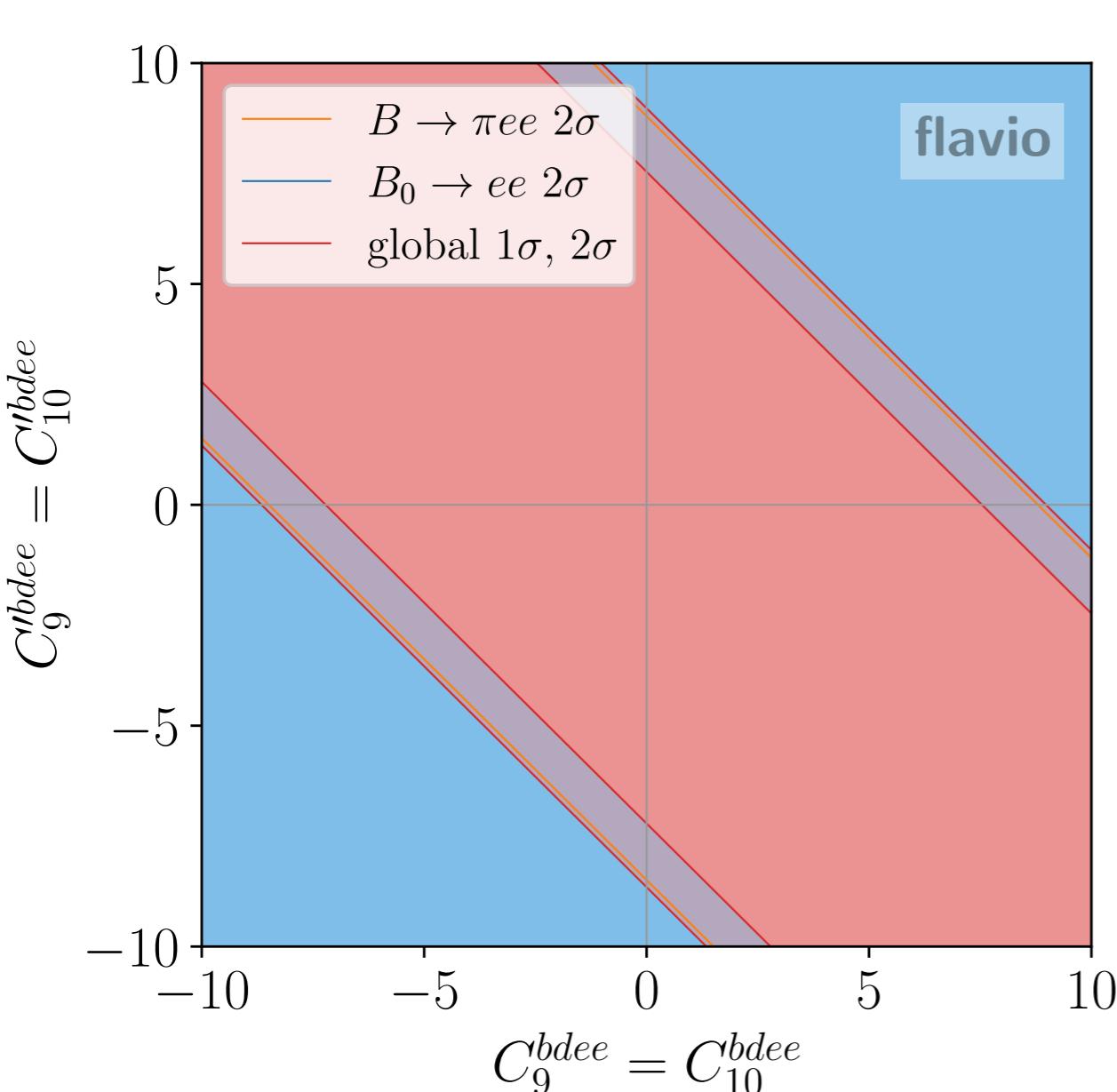
In SMEFT, confronting with high-mass Drell-Yan and $b \rightarrow q\nu\nu$ ($B \rightarrow K, \pi, \rho$)



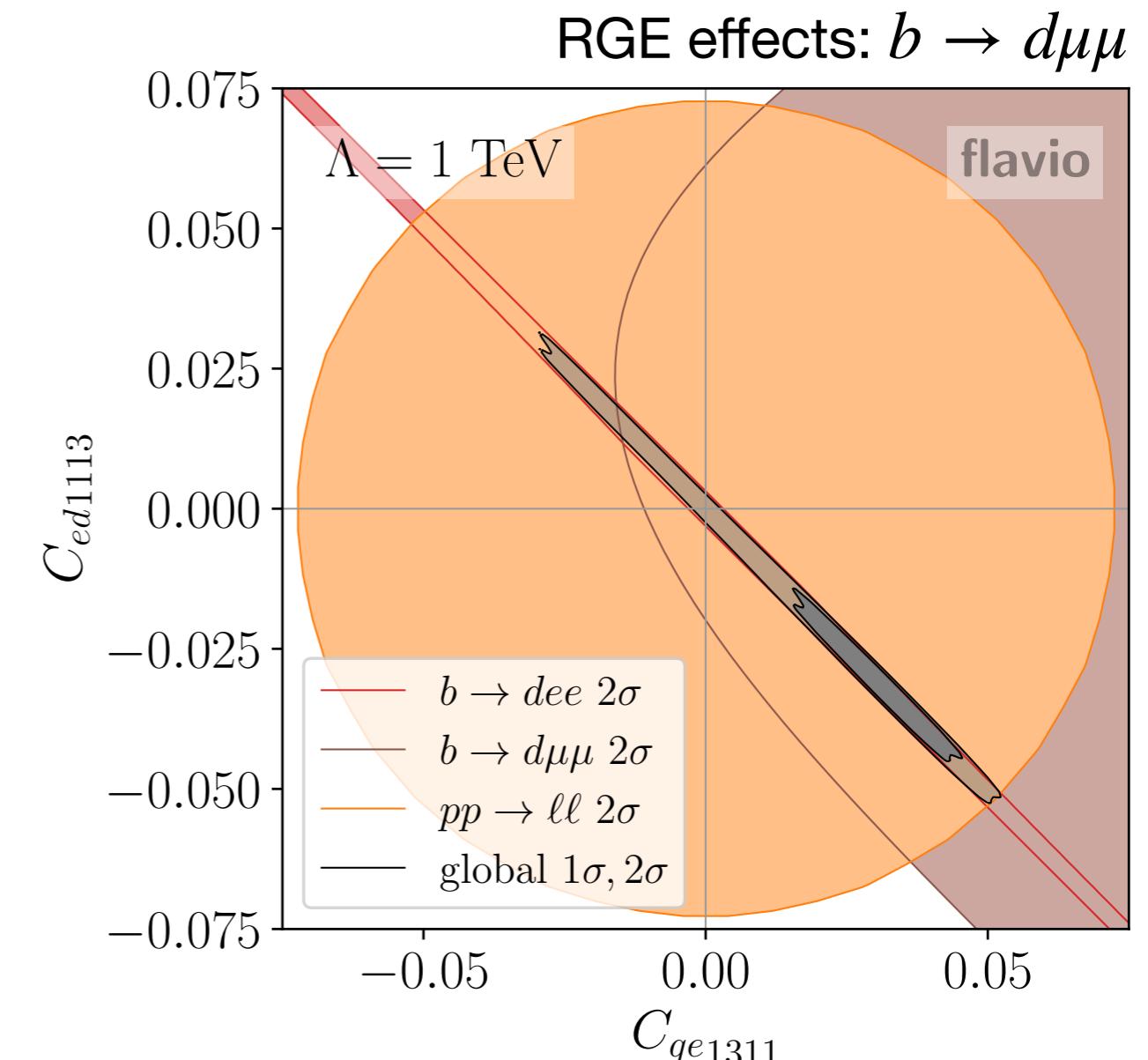
$$O_{9(10)}^{[']ijkk} \sim (\bar{q}_{L[R]}^j \gamma^\mu q_{L[R]}^i)(\bar{\ell}^k \gamma_\mu (\gamma_5) \ell^k)$$

Closed by $b \rightarrow q\nu\nu$

In SMEFT, confronting with high-mass Drell-Yan and $b \rightarrow q\nu\nu$ ($B \rightarrow K, \pi, \rho$)



$$O_{9(10)}^{[']ijkk} \sim (\bar{q}_{L[R]}^j \gamma^\mu q_{L[R]}^i)(\bar{\ell}^k \gamma_\mu (\gamma_5) \ell^k)$$



Closed by $b \rightarrow d\mu\mu$ and tails

Minimalistic flavour: B decays mostly more constraining, but now always

Towards realistic flavour

Similar to SM, NP most likely has a flavour structure

With a **flavour assumption** we correlate various SMEFT Wilson coefficients and decrease the number of free parameters

Minimal flavour violation:

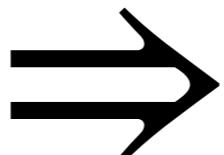
assume SM Yukawa couplings are the only source of flavour breaking also in NP

$$G_F = U(3)_Q \times U(3)_u \times U(3)_d$$

$$Q \sim (3,1,1)$$

$$u \sim (1,3,1)$$

$$d \sim (1,1,3)$$



$$Y_u \sim (3, \bar{3}, 1)$$

$$Y_d \sim (3, 1, \bar{3})$$

Then we can decompose coefficients with spurion insertions:

$$[C_{lq}^{(1)}]_{iist} \bar{L}_i \gamma_\mu L_i \bar{Q}_s \gamma^\mu Q_t \rightarrow [C_{lq}^{(1)}]_{iist} = \delta_{st} [C_{lq}^{(1)}]_\delta + (Y_u Y_u^\dagger)_{st} [C_{lq}^{(1)}]_{Y_u Y_u^\dagger}$$

(and similar for other operators involving $\bar{Q}Q$)

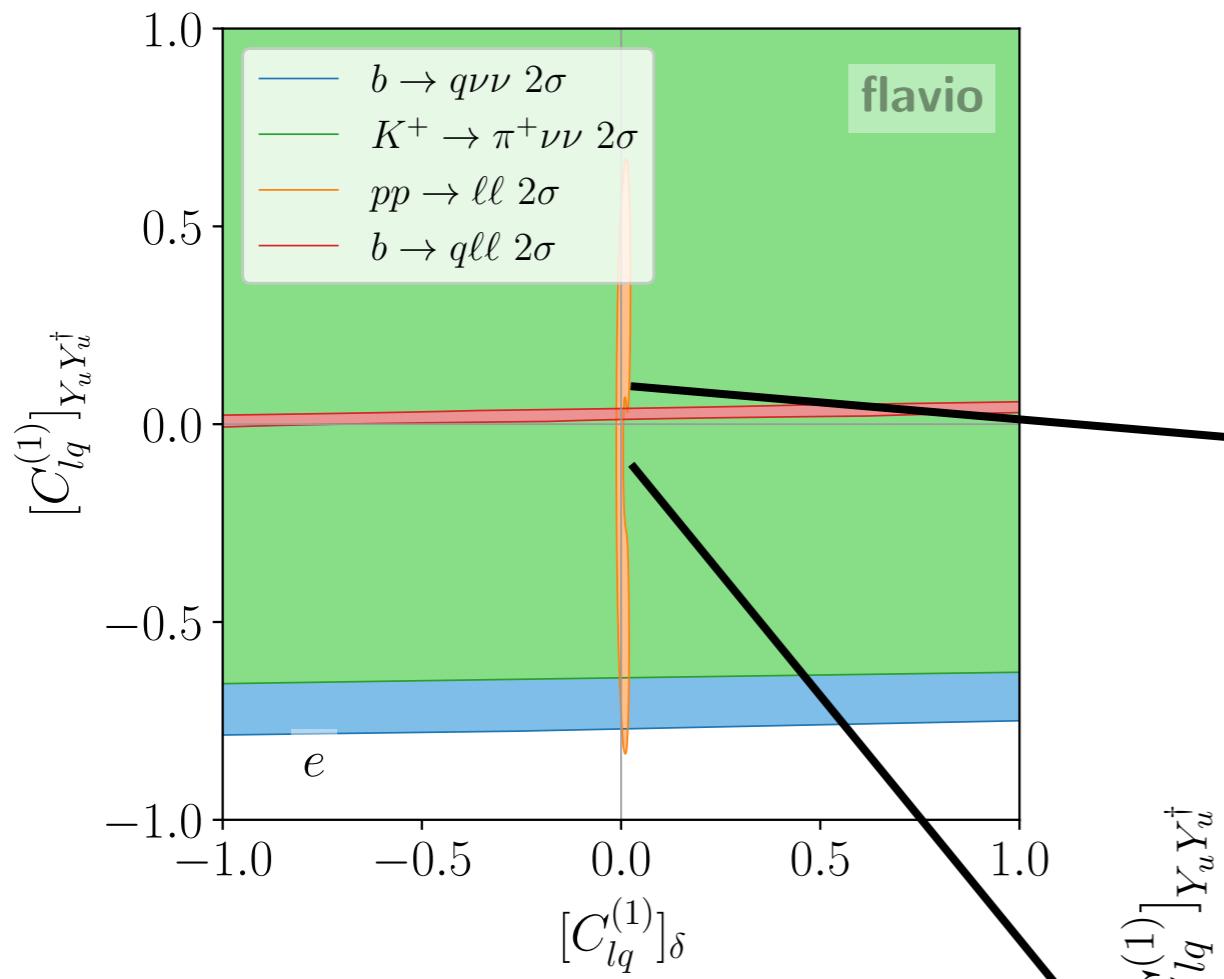
$$[C_{ledq}]_{iist} (\bar{L}_i e_i) (\bar{d}_s Q_t) \rightarrow [C_{ledq}]_{iist} = (Y_d^\dagger)_{st} [C_{ledq}]_{Y_d} + (Y_d^\dagger Y_u Y_u^\dagger)_{st} [C_{ledq}]_{Y_d^\dagger Y_u Y_u^\dagger}$$

$$\sim y_t^2 \begin{pmatrix} V_{td} V_{td}^* & V_{ts} V_{td}^* & V_{tb} V_{td}^* \\ V_{td} V_{ts}^* & V_{ts} V_{ts}^* & V_{tb} V_{ts}^* \\ V_{td} V_{tb}^* & V_{ts} V_{tb}^* & V_{tb} V_{tb}^* \end{pmatrix}$$

$$\sim y_b y_t^2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ V_{td} V_{tb}^* & V_{ts} V_{tb}^* & V_{tb} V_{tb}^* \end{pmatrix}$$

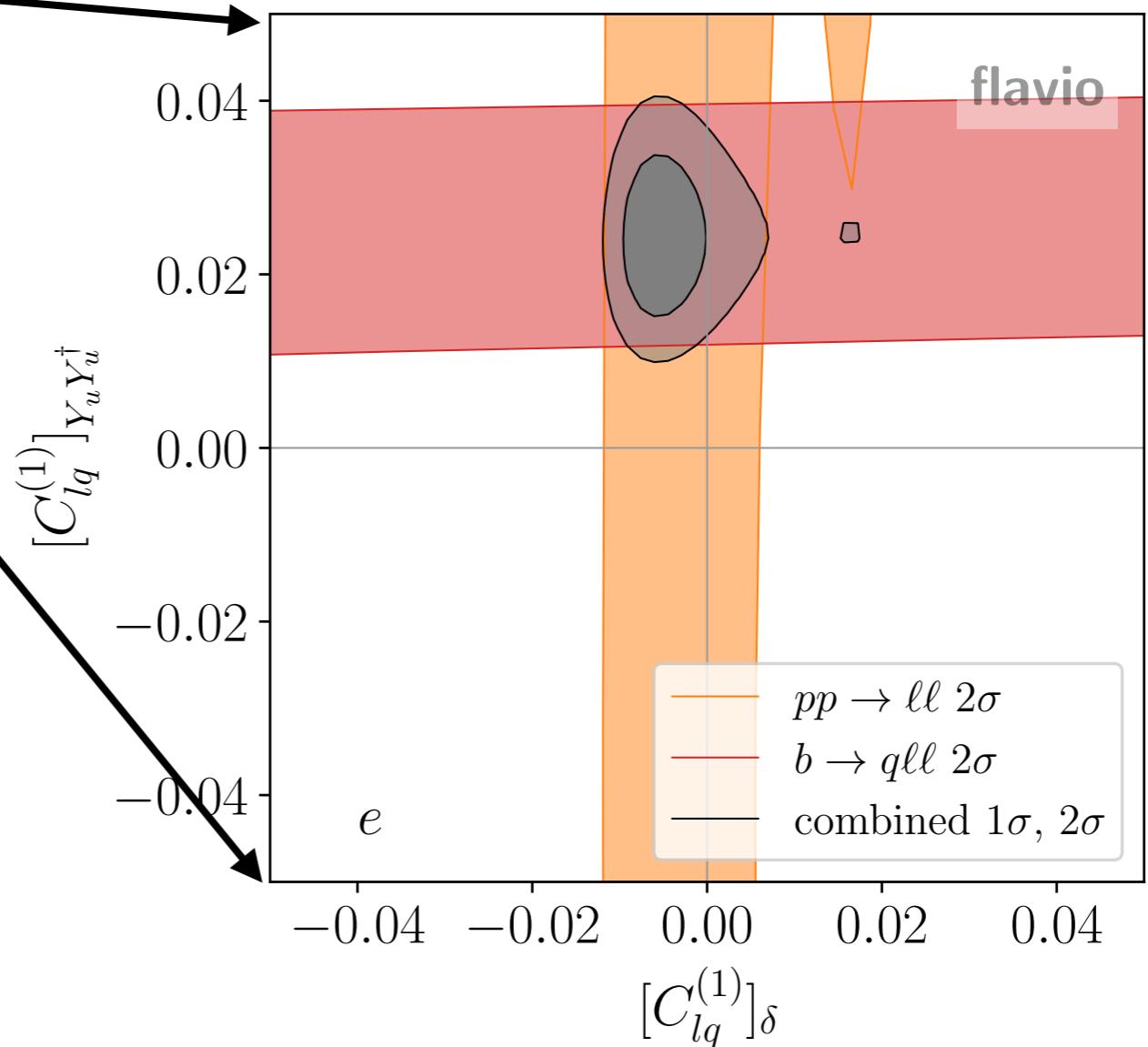
*all results in down-diagonal Warsaw basis

$$[C_{lq}^{(1)}]_{iist} \bar{L}_i \gamma_\mu L_i \bar{Q}_s \gamma^\mu Q_t \rightarrow [C_{lq}^{(1)}]_{iist} = \delta_{st} [C_{lq}^{(1)}]_\delta + (Y_u Y_u^\dagger)_{st} [C_{lq}^{(1)}]_{Y_u Y_u^\dagger}$$



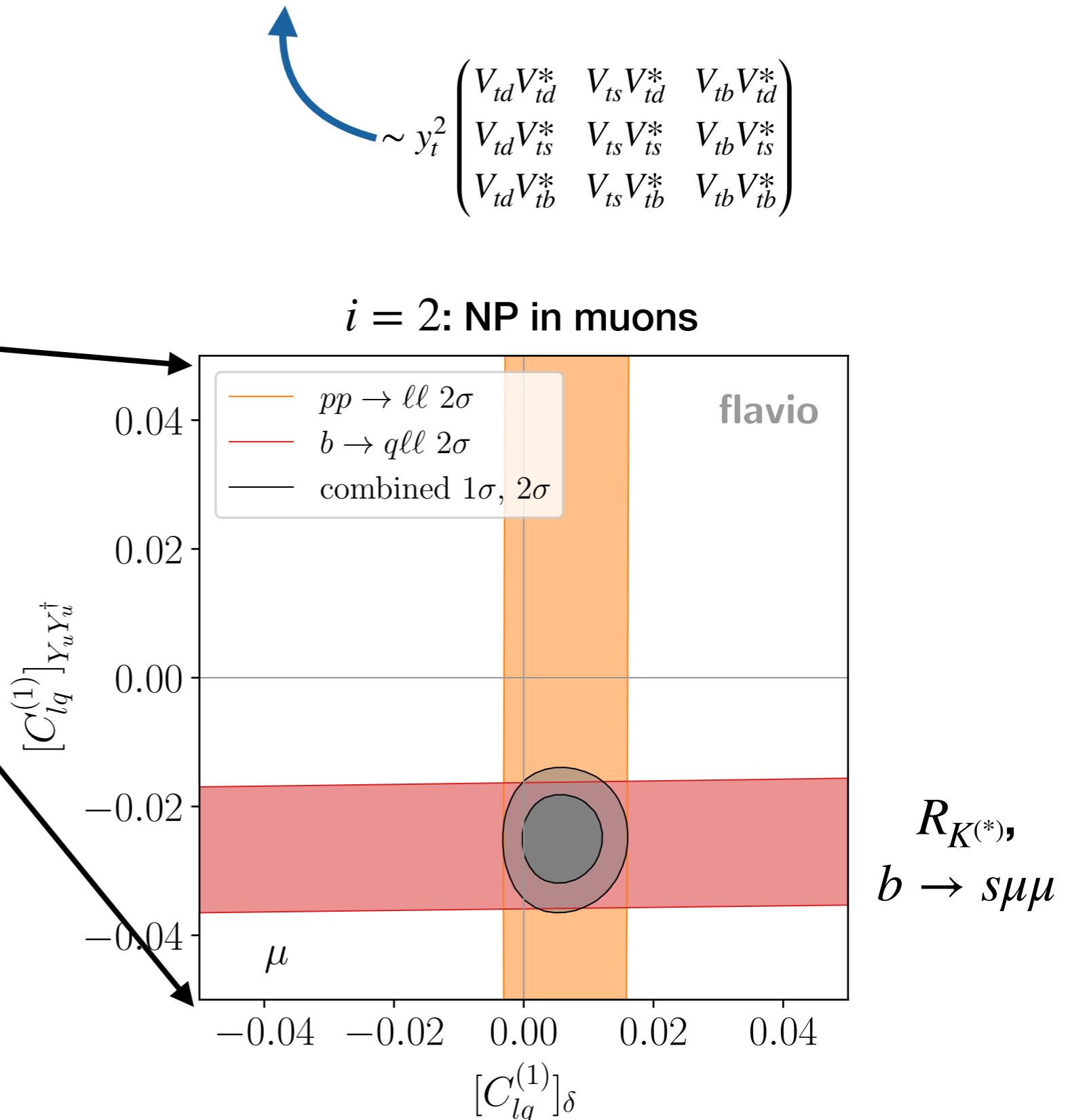
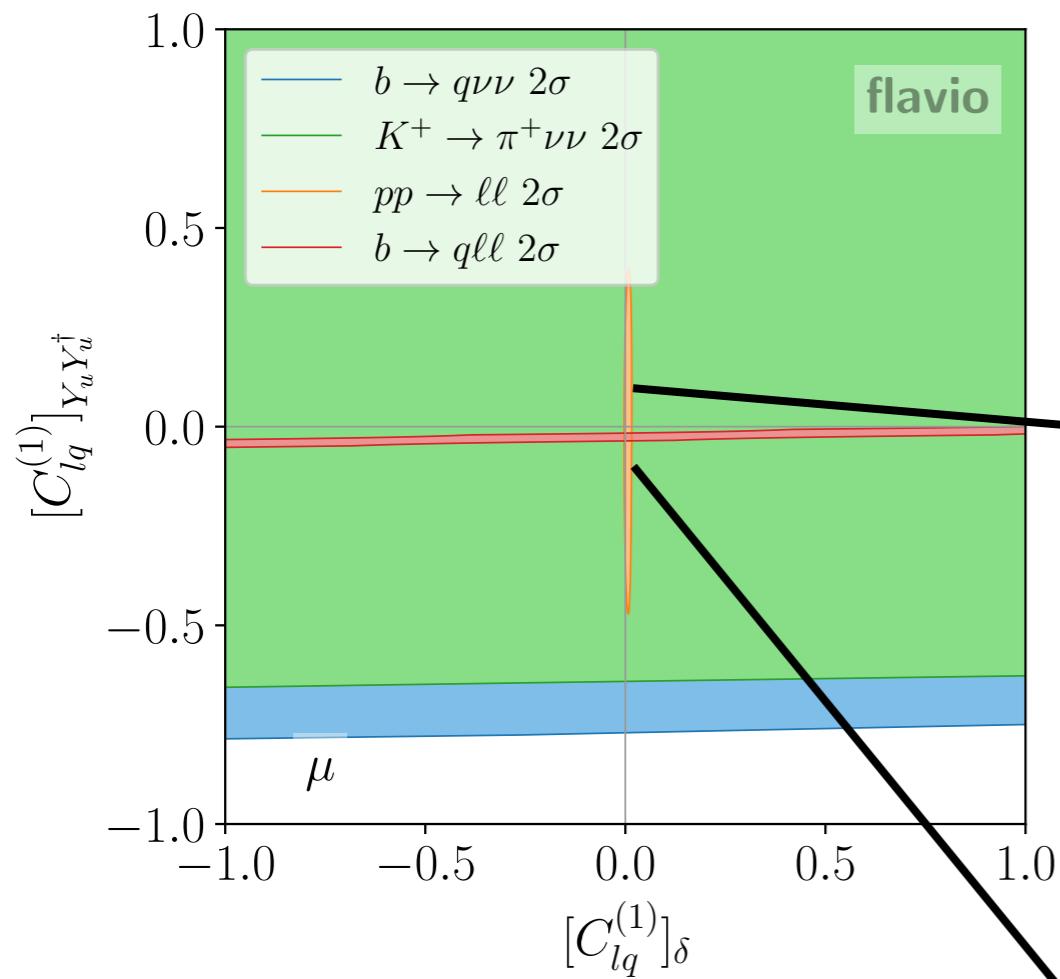
$$\sim y_t^2 \begin{pmatrix} V_{td}V_{td}^* & V_{ts}V_{td}^* & V_{tb}V_{td}^* \\ V_{td}V_{ts}^* & V_{ts}V_{ts}^* & V_{tb}V_{ts}^* \\ V_{td}V_{tb}^* & V_{ts}V_{tb}^* & V_{tb}V_{tb}^* \end{pmatrix}$$

$i = 1$: NP in electrons



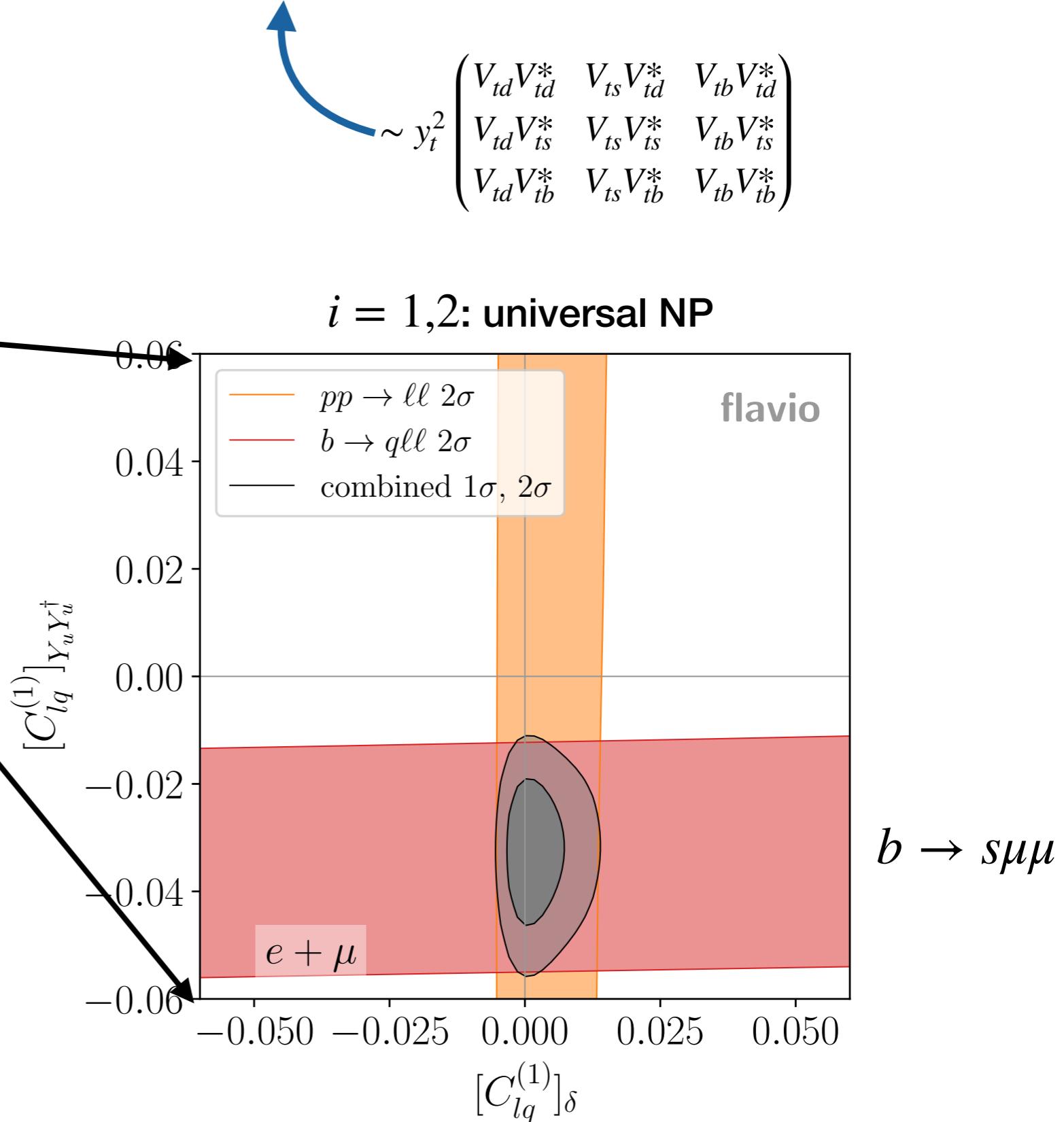
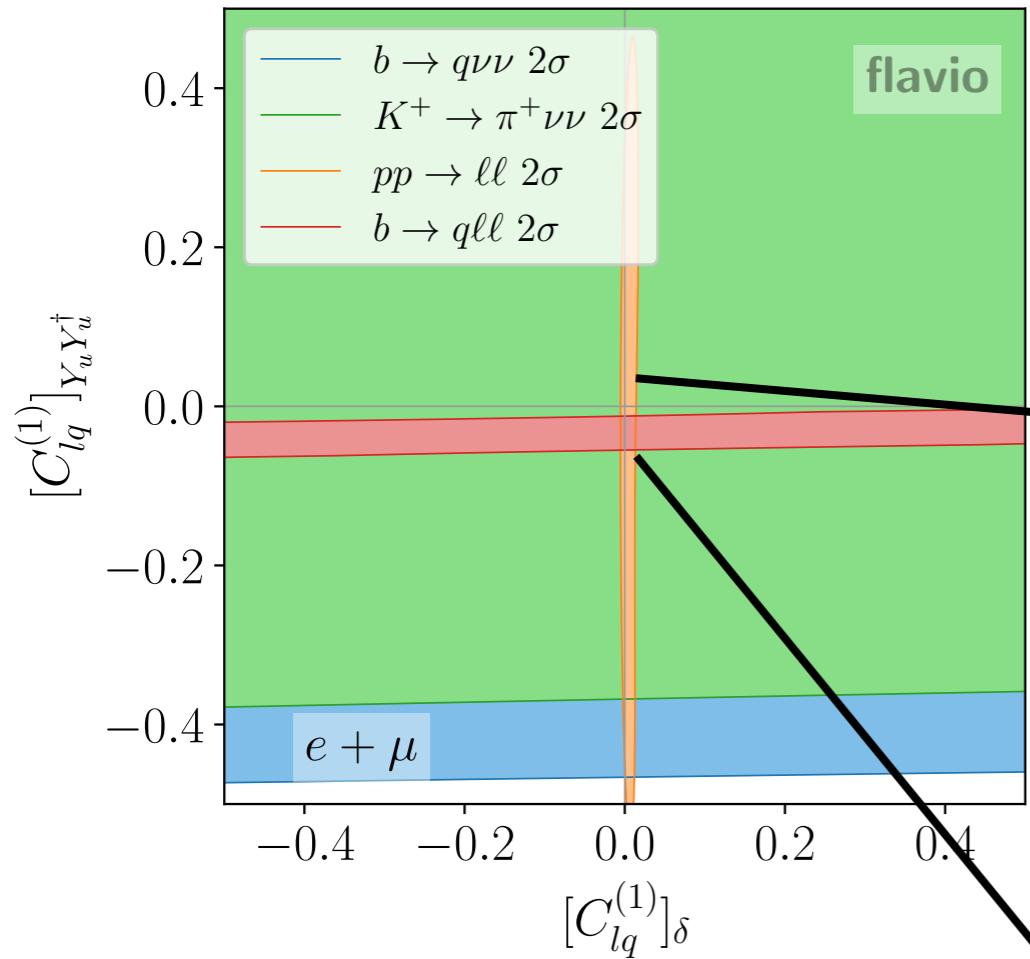
Realistic flavour: B decays and tails complementary

$$[C_{lq}^{(1)}]_{iist} \bar{L}_i \gamma_\mu L_i \bar{Q}_s \gamma^\mu Q_t \rightarrow [C_{lq}^{(1)}]_{iist} = \delta_{st} [C_{lq}^{(1)}]_\delta + (Y_u Y_u^\dagger)_{st} [C_{lq}^{(1)}]_{Y_u Y_u^\dagger}$$



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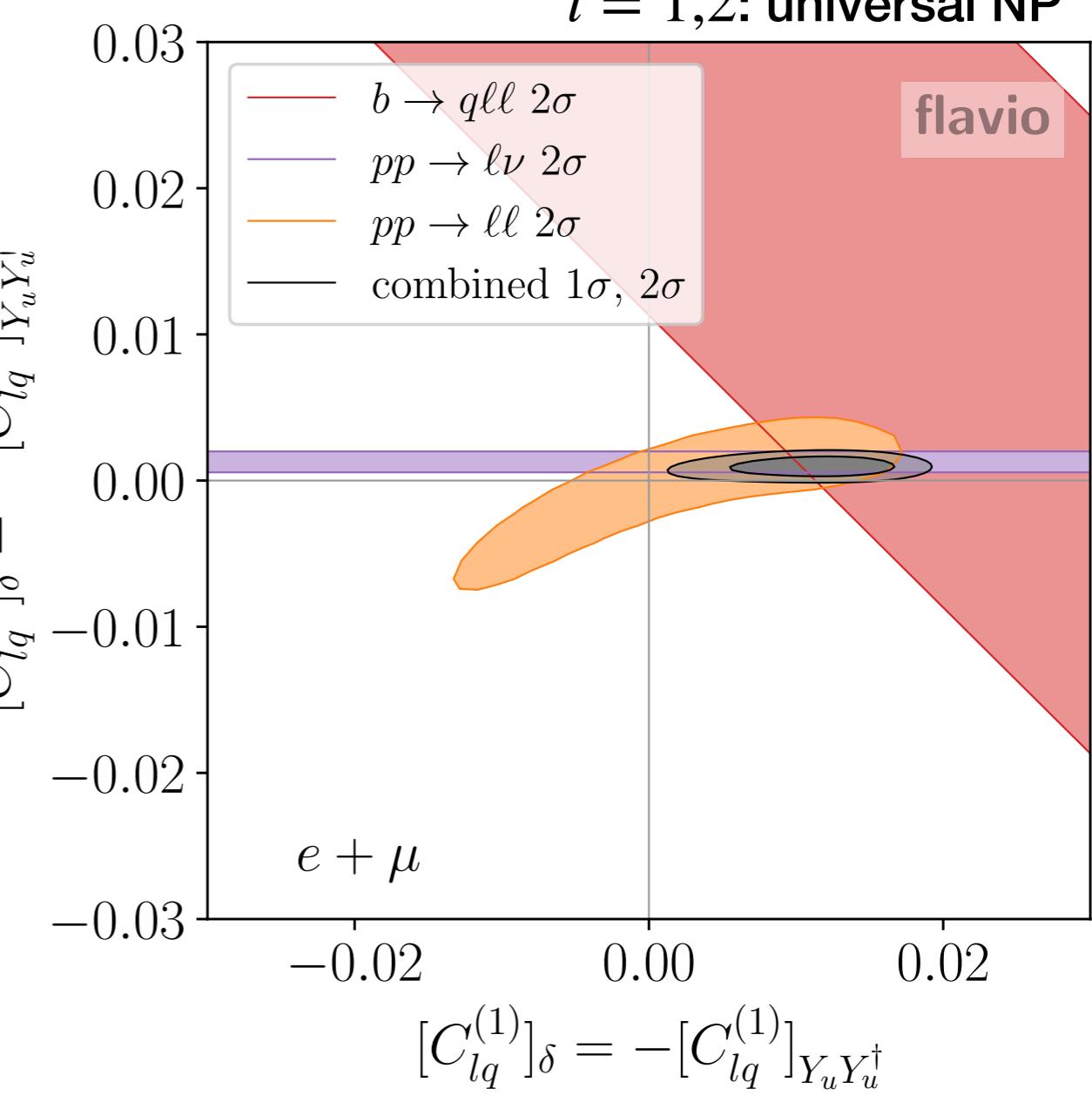
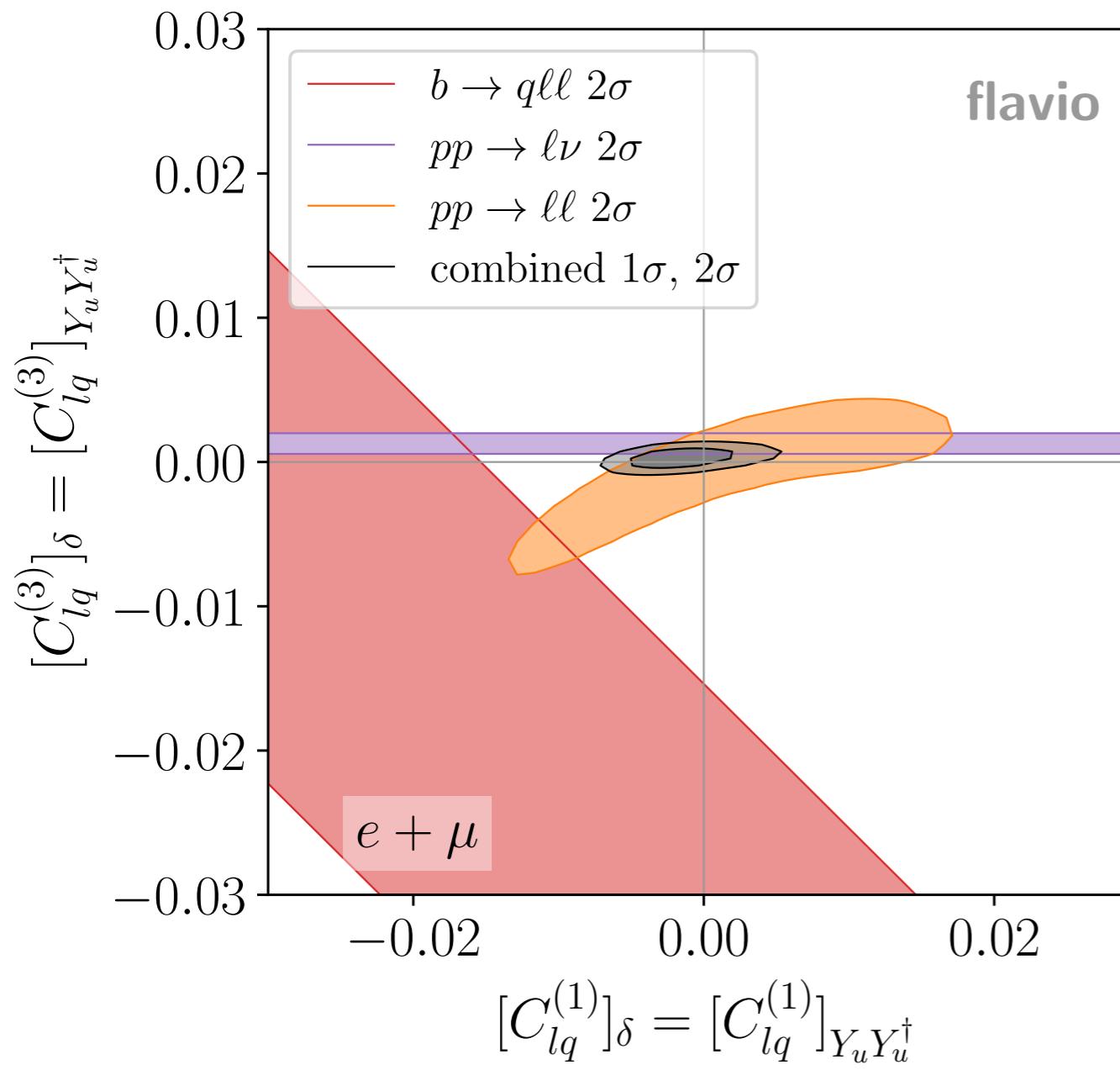
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Equating the coefficients with different signs

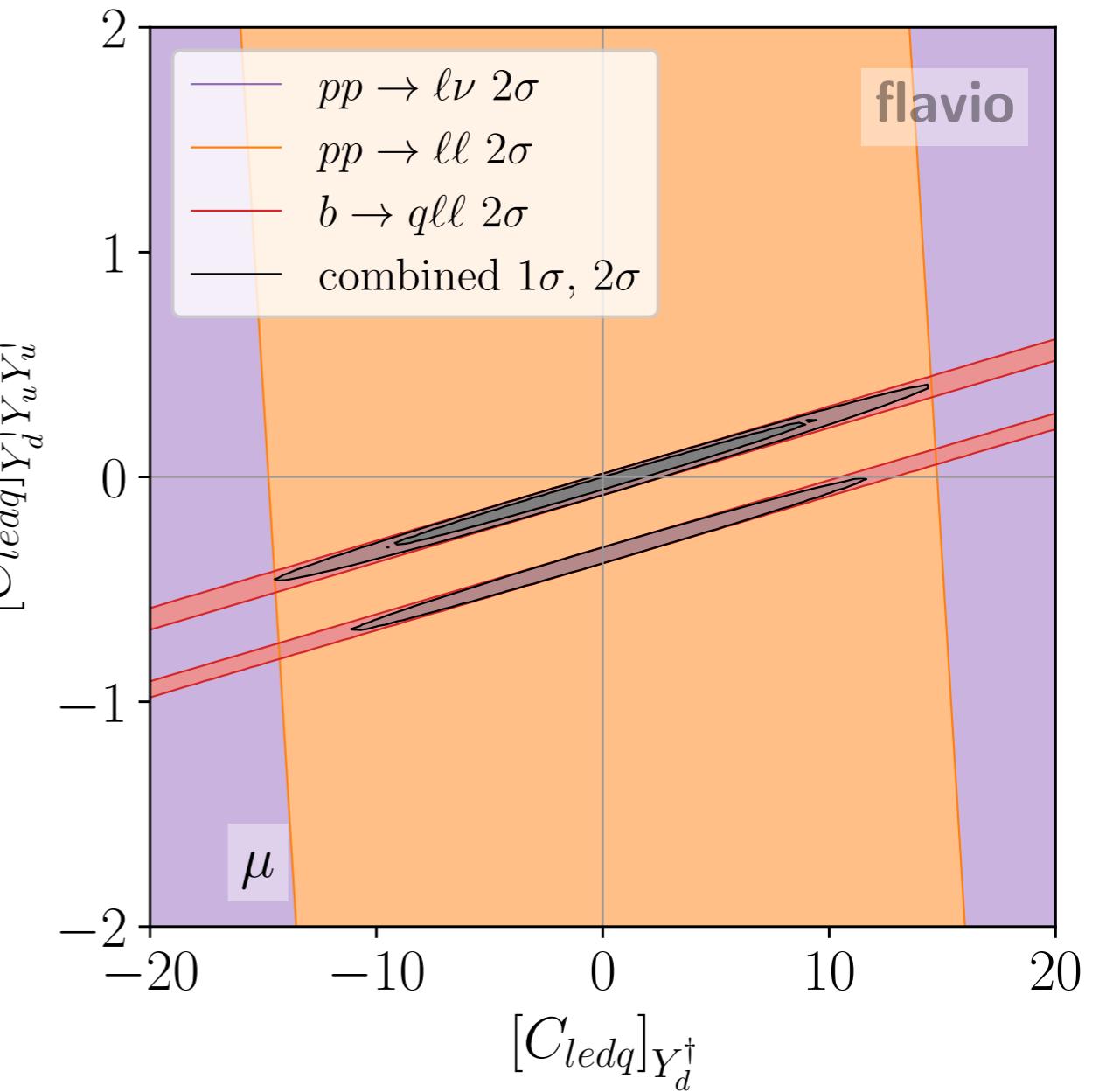
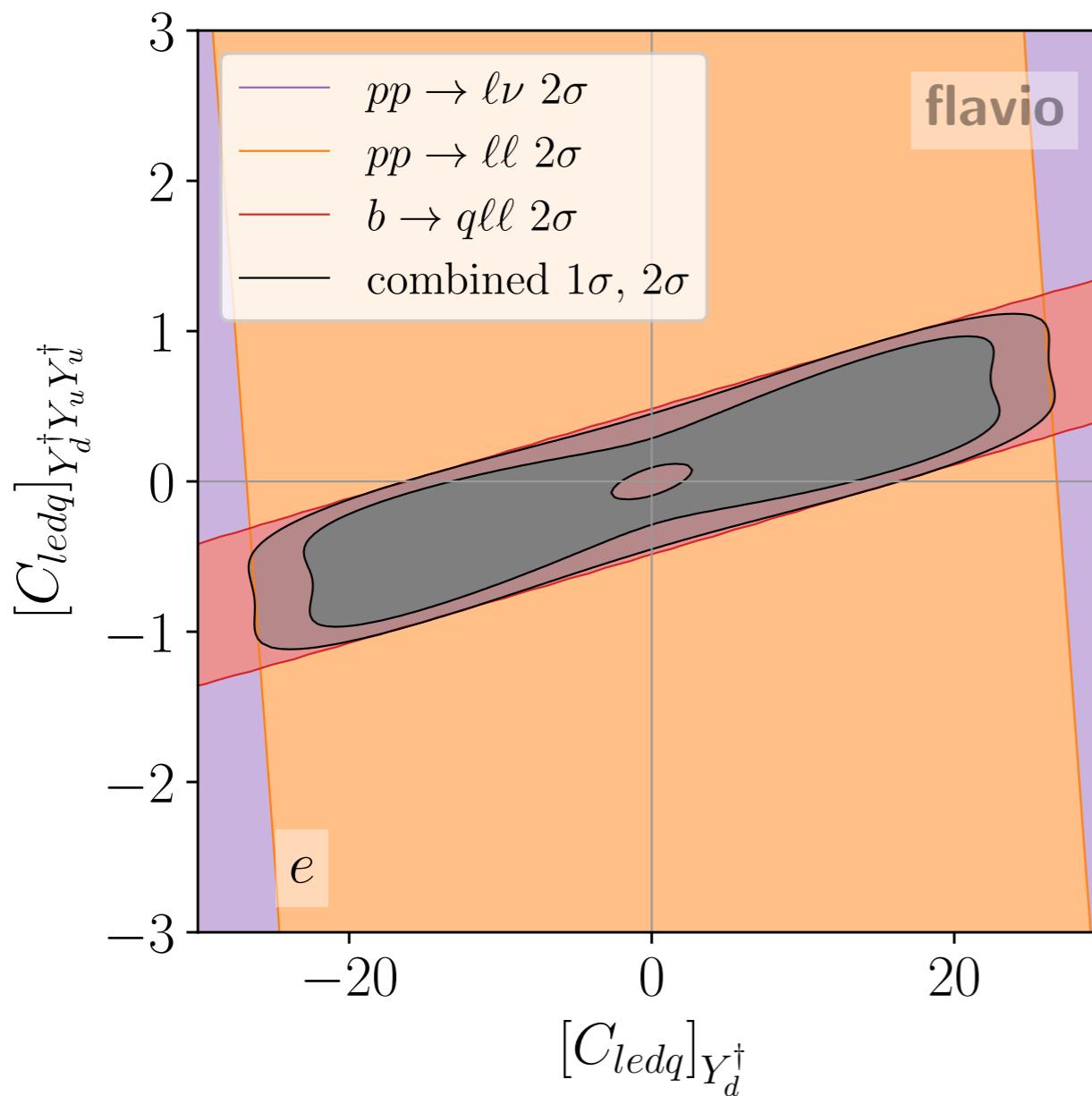
$b \rightarrow s\mu\mu$ anomalies
 $i = 1, 2$: universal NP



$$[C_{ledq}]_{iist}(\bar{L}_i e_i)(\bar{d}_s Q_t) \rightarrow [C_{ledq}]_{iist} = (Y_d^\dagger)_{st} [C_{ledq}]_{Y_d} + (Y_d^\dagger Y_u Y_u^\dagger)_{st} [C_{ledq}]_{Y_d^\dagger Y_u Y_u^\dagger}$$



$$\sim y_b y_t^2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ V_{td} V_{tb}^* & V_{ts} V_{tb}^* & V_{tb} V_{tb}^* \end{pmatrix}$$



At low energies dominated by $B_s \rightarrow ee, \mu\mu$

Summary and outlook

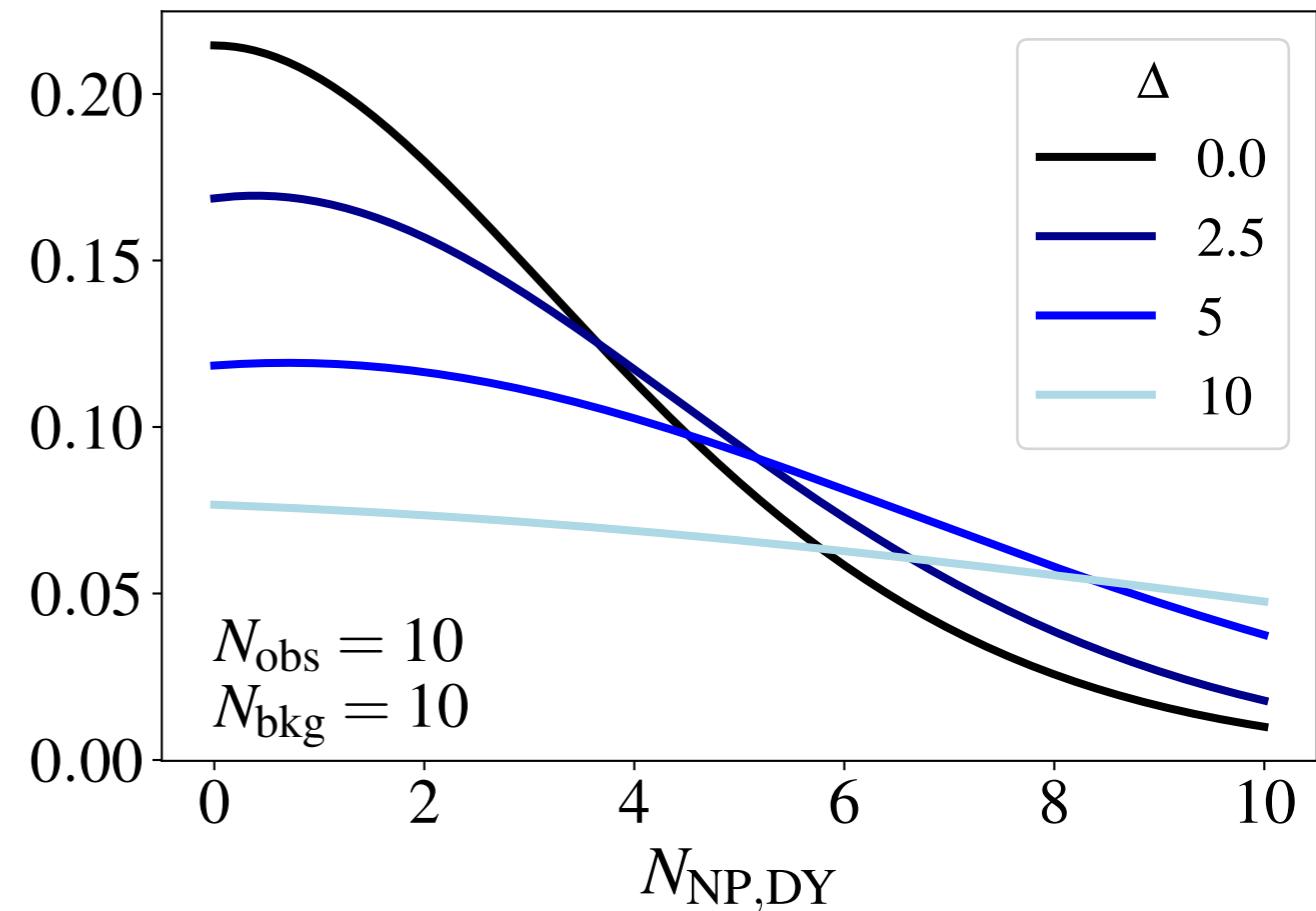
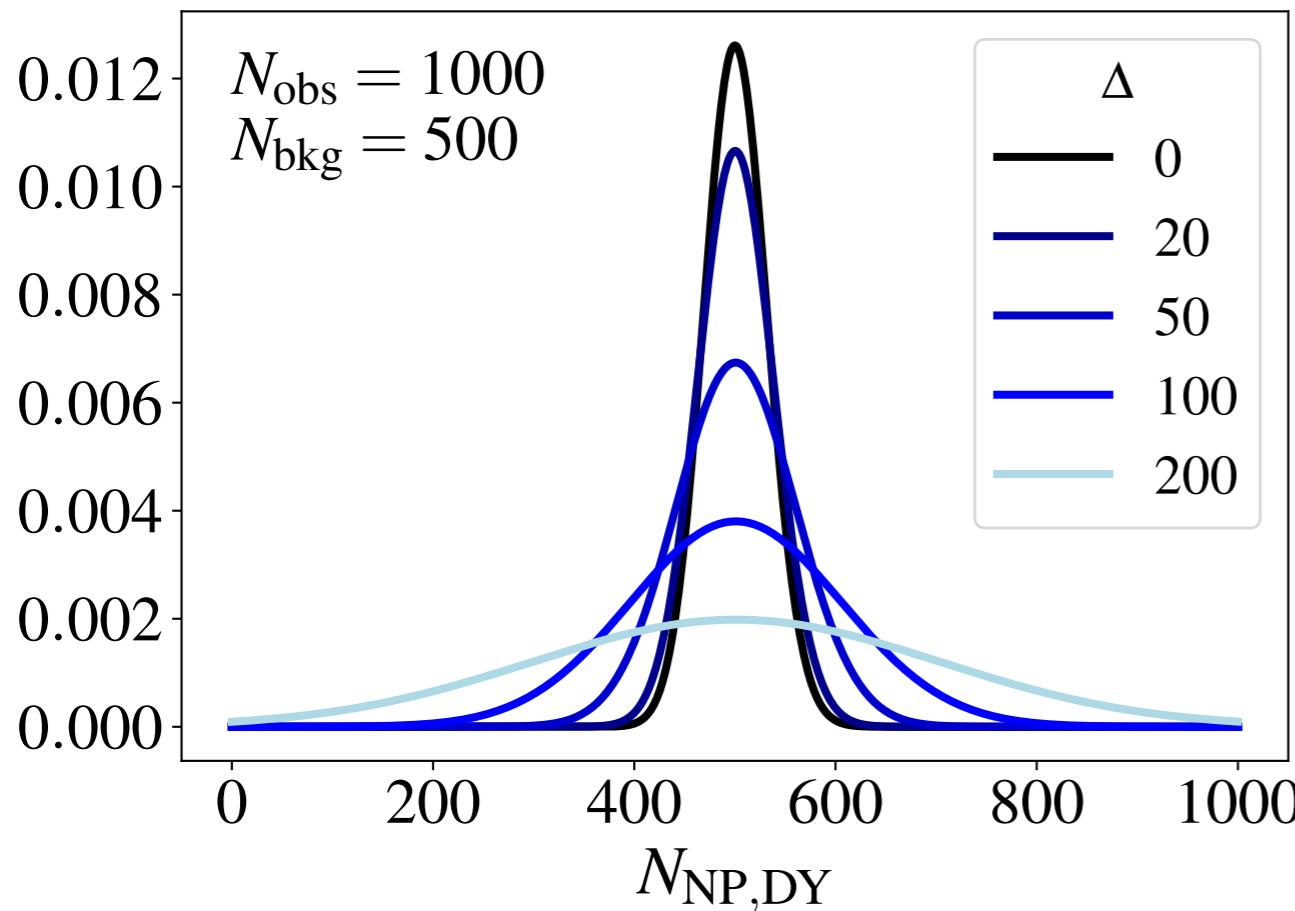
- We implement neutral and charged current high-mass Drell-Yan tails into the well established flavio framework
- We analyse the interplay between low-energy B -meson decays into light leptons and high-mass Drell-Yan constraints in the SMEFT
- B decays offer excellent constraining power, especially dominating in minimalistic flavour scenarios, more data in $b \rightarrow d\ell\ell$, $b \rightarrow see$ welcome
- High complementarity between Drell-Yan and meson decays in realistic flavour scenarios
- Stay tuned for more interesting examples and scenarios
Updated flavio coming soon!

Implementation of Drell-Yan

Experimental likelihood:

- Assume events in each bin independent, following the Poisson distribution
- Systematic uncertainties included by convolving with a Gaussian
- Number of expected events: $N_{\text{tot}} = N_{\text{NP,DY}} + N_{\text{bkg}}$
- Number of observed events: N_{obs}

$$f(N_{\text{NP,DY}} | N_{\text{obs}}) = \int d\tau \frac{(N_{\text{tot}} + \tau)^{N_{\text{obs}}} e^{-(N_{\text{tot}} + \tau)}}{N_{\text{obs}}!} \mathcal{N}(0, \Delta)(\tau)$$



Sensitivity of Drell-Yan

