# HIGGS BALLS: NOVEL NON-TOPOLOGICAL SOLITONS VIA THERMAL CORRECTIONS

#### Lauren Pearce

Pennsylvania State University-New Kensington

Based on:

L. Pearce, G. White, and A. Kusenko

JHEP 08 (2022) 033 (arXiv:2205.13557)

Thank you to the organizers for allowing a remote talk!



• Extended (scalar) field configuration

- Extended (scalar) field configuration
- Carries conserved charge

- Extended (scalar) field configuration
- Carries conserved charge
- (Quasi-)stable if total energy is less than Qm, mass energy of Q free scalar quanta

- Extended (scalar) field configuration
- Carries conserved charge
- ullet (Quasi-)stable if total energy is less than Qm, mass energy of Q free scalar quanta
- Condition for global Coleman (thin wall) Q-balls to exist:

$$\sqrt{V(\phi)/\phi^2}$$
 is minimized at finite nonzero  $\phi_0$ 

- Extended (scalar) field configuration
- Carries conserved charge
- ullet (Quasi-)stable if total energy is less than Qm, mass energy of Q free scalar quanta
- Condition for global Coleman (thin wall) Q-balls to exist:

$$\sqrt{V(\phi)/\phi^2}$$
 is minimized at finite nonzero  $\phi_0$ 

Energy per unit charge  $\omega=\sqrt{2V(\phi_0)/\phi_0^2}$  is less than  $m=\sqrt{2V''}$ 

- Extended (scalar) field configuration
- Carries conserved charge
- ullet (Quasi-)stable if total energy is less than Qm, mass energy of Q free scalar quanta
- Condition for global Coleman (thin wall) Q-balls to exist:

$$\sqrt{V(\phi)/\phi^2}$$
 is minimized at finite nonzero  $\phi_0$ 

Energy per unit charge 
$$\omega=\sqrt{2V(\phi_0)/\phi_0^2}$$
 is less than  $m=\sqrt{2V''}$ 

Typically, to minimize  $\sqrt{V(\phi)/\phi^2}$  at nonzero  $\phi_0$  requires an attractive interaction (e.g.,  $\phi^3$ )...but is this necessary?

• Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature

- $\bullet$  Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature
- ullet Theory has bosons whose mass is proportional to the VEV,  $\emph{m} \sim \emph{g} \phi$

- Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature
- ullet Theory has bosons whose mass is proportional to the VEV,  $m\sim g\phi$
- At finite temperature, one-loop corrections to scalar potential:

$$V_{1-{
m loop}} \supset \sum_{{
m bosons}} rac{n_i T^4}{2\pi^2} J_B \left(rac{m_i^2}{T^2}
ight)$$

- Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature
- ullet Theory has bosons whose mass is proportional to the VEV,  $m\sim g\phi$
- At finite temperature, one-loop corrections to scalar potential:

$$V_{1-\mathrm{loop}} \supset \sum_{\mathrm{bosons}} \frac{n_i T^4}{2\pi^2} J_B\left(\frac{m_i^2}{T^2}\right)$$

At high temperatures:

$$J_B(x) \approx -\frac{\pi^4}{45} + \frac{\pi^2}{12}x - \frac{\pi}{6}x^{3/2}$$

- Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature
- ullet Theory has bosons whose mass is proportional to the VEV,  $m\sim g\phi$
- At finite temperature, one-loop corrections to scalar potential:

$$V_{1-\mathrm{loop}} \supset \sum_{\mathrm{bosons}} \frac{n_i T^4}{2\pi^2} J_B\left(\frac{m_i^2}{T^2}\right)$$

At high temperatures:

$$J_B(x) \approx -\frac{\pi^4}{45} + \frac{\pi^2}{12}x - \frac{\pi}{6}x^{3/2}$$

induces a term  $\sim -AT|\phi|^3$ .

◆□▶◆□▶◆□▶◆□▶ □□ ♥Q♥

- Renormalizable theory with a  $\phi \to -\phi$  symmetry: No Q-balls at zero temperature
- ullet Theory has bosons whose mass is proportional to the VEV,  $m\sim g\phi$
- At finite temperature, one-loop corrections to scalar potential:

$$V_{1-\mathrm{loop}} \supset \sum_{\mathrm{bosons}} \frac{n_i T^4}{2\pi^2} J_B\left(\frac{m_i^2}{T^2}\right)$$

At high temperatures:

$$J_B(x) \approx -\frac{\pi^4}{45} + \frac{\pi^2}{12}x - \frac{\pi}{6}x^{3/2}$$

induces a term  $\sim -AT|\phi|^3$ .

• Potential issues: Finite temperature corrections also affect the mass; High T not valid unless  $T\gg\phi_0$ 

◆ロト ◆部 ▶ ◆ 恵 ▶ ◆ 恵 ▶ ・ 恵 | 重 | 1 回 り へ ○

#### Complications with the SM Higgs:

• Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)

- Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)
- Solution: Adjust top pole mass to ensure V(h)>0 to  $h\gtrsim 10^{18}\,{
  m GeV}$

- Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)
- ullet Solution: Adjust top pole mass to ensure V(h)>0 to  $h\gtrsim 10^{18}\,{
  m GeV}$
- Q-balls made of Higgs quanta carry gauge charge- repulsive interactions mediated by gauge bosons increases energy

- Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)
- Solution: Adjust top pole mass to ensure V(h)>0 to  $h\gtrsim 10^{18}\,{
  m GeV}$
- Q-balls made of Higgs quanta carry gauge charge- repulsive interactions mediated by gauge bosons increases energy
- First consider "global" Higgs model: gauge bosons have masses  $\sim g \ \langle h \rangle$ , but ignore repulsive gauge boson interactions inside Q-ball

- Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)
- ullet Solution: Adjust top pole mass to ensure V(h)>0 to  $h\gtrsim 10^{18}\,{
  m GeV}$
- Q-balls made of Higgs quanta carry gauge charge- repulsive interactions mediated by gauge bosons increases energy
- First consider "global" Higgs model: gauge bosons have masses  $\sim g \, \langle h \rangle$ , but ignore repulsive gauge boson interactions inside Q-ball
- Then consider gauged SM Higgs model; extra energy from the gauge interactions prevents thermal Higgs balls from existing

#### Complications with the SM Higgs:

- Due to running quartic coupling, V(h) < 0 at large scales, which leads to solitosynthesis (phase transition to true vacuum)
- ullet Solution: Adjust top pole mass to ensure V(h)>0 to  $h\gtrsim 10^{18}\,{
  m GeV}$
- Q-balls made of Higgs quanta carry gauge charge- repulsive interactions mediated by gauge bosons increases energy
- First consider "global" Higgs model: gauge bosons have masses  $\sim g \ \langle h \rangle$ , but ignore repulsive gauge boson interactions inside Q-ball
- Then consider gauged SM Higgs model; extra energy from the gauge interactions prevents thermal Higgs balls from existing
- What about extensions of SM?

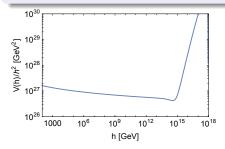
4 D > 4 P > 4 E > 4 E > 5 E = 4 9 Q P

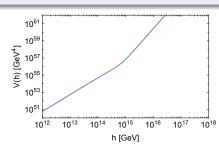
• Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball

- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects

- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects
- Do not make the high temperature expansion

- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects
- Do not make the high temperature expansion

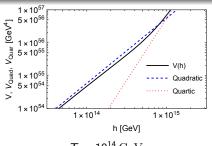




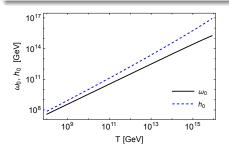
$$T=10^{14}\,\mathrm{GeV}$$

- 4 ロ ト 4 週 ト 4 夏 ト 4 夏 ト 夏 1 五 り 9 ( )

- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects
- Do not make the high temperature expansion

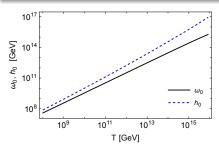


- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects
- Do not make the high temperature expansion

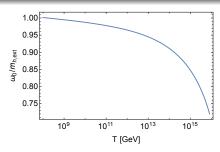


 $h_0, \omega_0 \propto T$ , high T never valid

- Treat the Higgs ball as if it carries global, not gauge charge: ignore gauge-boson-mediated interactions between charge inside the Q-ball
- Include in the Higgs potential: Zero-temperature one-loop effects, finite-temperature one-loop effects, and ring (daisy) finite temperature effects
- Do not make the high temperature expansion



 $h_0, \omega_0 \propto T$ , high T never valid



Q-balls for  $T\gtrsim 10^9\,{
m GeV}$ 

 $\bullet$  Gauged  $\mathrm{U}(1)$  Q-balls: Heeck et. al., Phys. Rev. D 103 (2021)

- ullet Gauged U(1) Q-balls: Heeck et. al., Phys. Rev. D 103 (2021)
- Generalized their approach to generic SU(N) and then  $SU(2) \times U(1)$ , but still making the static charge approximation:

$$\phi(x,t) = \frac{\phi_0}{\sqrt{2}}F(r)e^{i\omega t},$$

$$A_0^a(x,t) = A_0^r(t), \qquad A_i^a(x,t) = 0$$

- ullet Gauged U(1) Q-balls: Heeck et. al., Phys. Rev. D 103 (2021)
- Generalized their approach to generic SU(N) and then  $SU(2) \times U(1)$ , but still making the static charge approximation:

$$\phi(x,t) = \frac{\phi_0}{\sqrt{2}}F(r)e^{i\omega t},$$

$$A_0^a(x,t) = A_0^r(t), \qquad A_i^a(x,t) = 0$$

• Removes  $\mathrm{SU}(2)$  self-interactions between the gauge fields in the Q-ball (Higgs quanta still interact via  $\mathrm{SU}(2)$  interactions)

- ullet Gauged  $\mathrm{U}(1)$  Q-balls: Heeck et. al., Phys. Rev. D 103 (2021)
- Generalized their approach to generic SU(N) and then  $SU(2) \times U(1)$ , but still making the static charge approximation:

$$\phi(x,t) = \frac{\phi_0}{\sqrt{2}}F(r)e^{i\omega t},$$

$$A_0^a(x,t) = A_0^r(t), \qquad A_i^a(x,t) = 0$$

- Removes  $\mathrm{SU}(2)$  self-interactions between the gauge fields in the Q-ball (Higgs quanta still interact via  $\mathrm{SU}(2)$  interactions)
- Breakdown of static charge approximation  $\leftrightarrow$  confining nature of SU(N) gauge interactions

- ullet Gauged U(1) Q-balls: Heeck et. al., Phys. Rev. D 103 (2021)
- Generalized their approach to generic SU(N) and then  $SU(2) \times U(1)$ , but still making the static charge approximation:

$$\phi(x,t) = \frac{\phi_0}{\sqrt{2}}F(r)e^{i\omega t},$$

$$A_0^a(x,t) = A_0^r(t), \qquad A_i^a(x,t) = 0$$

- Removes  $\mathrm{SU}(2)$  self-interactions between the gauge fields in the Q-ball (Higgs quanta still interact via  $\mathrm{SU}(2)$  interactions)
- Breakdown of static charge approximation  $\leftrightarrow$  confining nature of  $\mathrm{SU}(N)$  gauge interactions
- ullet Use static charge approximation only if Q-ball radius is much less than  $\mathrm{SU}(2)$  confinement scale

### Properties of Gauged Higgs Balls

#### Results:

• Energy per unit charge  $\omega$ :

$$\omega = rac{1}{2} R extit{h}_0 \omega_0 \sqrt{g_W^2 + g_Y^2} \coth\left(rac{1}{2} R extit{h}_0 \sqrt{g_W^2 + g_Y^2}
ight)$$

where  $\omega_0$  is the "global" Q-ball energy per unit charge,  $h_0$  is the global Q-ball VEV, and R is the radius.

### Properties of Gauged Higgs Balls

#### Results:

• Energy per unit charge  $\omega$ :

$$\omega = \frac{1}{2} \textit{Rh}_0 \omega_0 \sqrt{g_W^2 + g_Y^2} \coth \left( \frac{1}{2} \textit{Rh}_0 \sqrt{g_W^2 + g_Y^2} \right)$$

where  $\omega_0$  is the "global" Q-ball energy per unit charge,  $h_0$  is the global Q-ball VEV, and R is the radius.

Charge:

$$Q_{Y} = -Q_{W} = \frac{8\pi R\omega_{0} \left(Rh_{0}\sqrt{g_{W}^{2} + g_{Y}^{2}} \coth\left(\frac{1}{2}Rh_{0}\sqrt{g_{W}^{2} + g_{Y}^{2}}\right) - 2\right)}{g_{W}^{2} + g_{Y}^{2}}$$

(Numerically invert to find radius for a given charge.)

### Properties of Gauged Higgs Balls

#### Results:

• Energy per unit charge  $\omega$ :

$$\omega = \frac{1}{2} R \textit{h}_0 \omega_0 \sqrt{g_W^2 + g_Y^2} \coth \left( \frac{1}{2} R \textit{h}_0 \sqrt{g_W^2 + g_Y^2} \right)$$

where  $\omega_0$  is the "global" Q-ball energy per unit charge,  $h_0$  is the global Q-ball VEV, and R is the radius.

Charge:

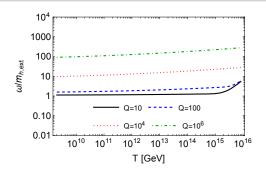
$$Q_{Y} = -Q_{W} = \frac{8\pi R\omega_{0} \left(Rh_{0}\sqrt{g_{W}^{2} + g_{Y}^{2}} \coth\left(\frac{1}{2}Rh_{0}\sqrt{g_{W}^{2} + g_{Y}^{2}}\right) - 2\right)}{g_{W}^{2} + g_{Y}^{2}}$$

(Numerically invert to find radius for a given charge.)

• VEV: Step function with  $h = h_0$  (global value) inside.

### Thermal Higgs Balls in the SM

No thermal Higgs balls in the Standard Model: Extra energy from gauge boson repulsion between charges makes the energy per unit charge  $\omega$  greater than the mass of a free Higgs quanta:



Does this hold for all extensions of the SM? (Or in non-SM sectors?)

• Idea: Modify the running of the gauge couplings  $g_W$  and  $g_Y$  (e.g., extra fermions)

- Idea: Modify the running of the gauge couplings  $g_W$  and  $g_Y$  (e.g., extra fermions)
- Why this might work: Decreases the energy contribution from repulsive gauge interactions

- Idea: Modify the running of the gauge couplings  $g_W$  and  $g_Y$  (e.g., extra fermions)
- Why this might work: Decreases the energy contribution from repulsive gauge interactions
- Why this doesn't work: The gauge boson masses  $\sim gh_0$ , so it also decreases the cubic term that makes the thermal balls

- Idea: Modify the running of the gauge couplings  $g_W$  and  $g_Y$  (e.g., extra fermions)
- Why this might work: Decreases the energy contribution from repulsive gauge interactions
- Why this doesn't work: The gauge boson masses  $\sim gh_0$ , so it also decreases the cubic term that makes the thermal balls
- Need to decouple these two things

Does this hold for all extensions of the SM? (Or in non-SM sectors?)

• Also introduce a scalar field:

$$V(H,S) = -\mu^2 H^\dagger H + \lambda_H (H^\dagger H)^2 + \frac{m_S^2}{2} S^2 + \lambda_S S^4 + \lambda_{HS} H^\dagger H S^2$$

which has mass:

$$m_{S,\mathrm{eff}} = \sqrt{m_S^2 + \lambda_{HS} h^2} \approx \sqrt{\lambda_{HS}} h$$



Does this hold for all extensions of the SM? (Or in non-SM sectors?)

Also introduce a scalar field:

$$V(H,S) = -\mu^2 H^{\dagger} H + \lambda_H (H^{\dagger} H)^2 + \frac{m_S^2}{2} S^2 + \lambda_S S^4 + \lambda_{HS} H^{\dagger} H S^2$$
 which has mass:

$$m_{S, ext{eff}} = \sqrt{m_S^2 + \lambda_{HS} h^2} pprox \sqrt{\lambda_{HS}} h$$

• Coefficient of  $-AT|h|^3$  term in potential controlled by  $\lambda_{HS}$ , not gauge couplings

Lauren Pearce (PSU-NK)

Does this hold for all extensions of the SM? (Or in non-SM sectors?)

Also introduce a scalar field:

$$V(H,S) = -\mu^2 H^{\dagger} H + \lambda_H (H^{\dagger} H)^2 + \frac{m_S^2}{2} S^2 + \lambda_S S^4 + \lambda_{HS} H^{\dagger} H S^2$$
 which has mass:

$$m_{S,\mathrm{eff}} = \sqrt{m_S^2 + \lambda_{HS}h^2} \approx \sqrt{\lambda_{HS}}h$$

- Coefficient of  $-AT|h|^3$  term in potential controlled by  $\lambda_{HS}$ , not gauge couplings
- This method can be used with global charge to make thermal Q-balls

Lauren Pearce (PSU-NK)

Does this hold for all extensions of the SM? (Or in non-SM sectors?)

Also introduce a scalar field:

$$V(H,S) = -\mu^2 H^{\dagger} H + \lambda_H (H^{\dagger} H)^2 + \frac{m_S^2}{2} S^2 + \lambda_S S^4 + \lambda_{HS} H^{\dagger} H S^2$$
 which has mass:

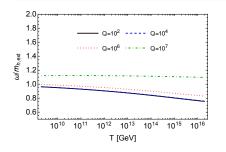
$$m_{S,\mathrm{eff}} = \sqrt{m_S^2 + \lambda_{HS} h^2} \approx \sqrt{\lambda_{HS}} h$$

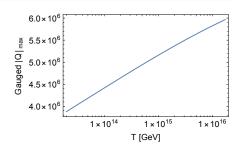
- Coefficient of  $-AT|h|^3$  term in potential controlled by  $\lambda_{HS}$ , not gauge couplings
- This method can be used with global charge to make thermal Q-balls
- SUSY has both: extra scalars & modified gauge coupling runnings

4 □ ▶ 4 □ ▶ 4 □ ▶ 4 □ ▶ 4 □ ▶

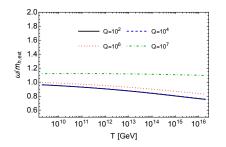
 $\lambda_{HS}=0.9$ , and  $g_Y,g_W$  running to one-tenth their SM values: (not fine-tuned since at one loop level, contributions to the  $\beta$  function  $\propto g^3$ )

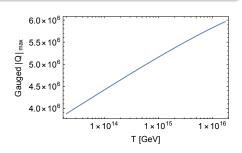
 $\lambda_{HS}=0.9$ , and  $g_Y,g_W$  running to one-tenth their SM values: (not fine-tuned since at one loop level, contributions to the  $\beta$  function  $\propto g^3$ )





 $\lambda_{HS}=0.9$ , and  $g_Y,g_W$  running to one-tenth their SM values: (not fine-tuned since at one loop level, contributions to the  $\beta$  function  $\propto g^3$ )

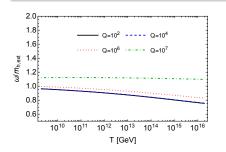


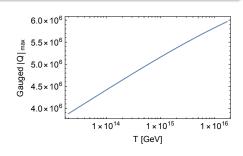


ullet Across a range of temperatures, Higgs balls exist up to charges  $\sim 10^6$ 

Lauren Pearce (PSU-NK)

 $\lambda_{HS}=0.9$ , and  $g_Y,g_W$  running to one-tenth their SM values: (not fine-tuned since at one loop level, contributions to the  $\beta$  function  $\propto g^3$ )





- ullet Across a range of temperatures, Higgs balls exist up to charges  $\sim 10^6$
- Quasi-stable: Higgs quanta can decay

Thermal Q-balls

#### Thermal Q-balls

 Non-topological solitons that exist at finite temperature due to induced attractive interactions

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

### Example: Higgs Balls

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

### Example: Higgs Balls

No Higgs balls in SM, but may exist in BSM models (SUSY?)

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

### Example: Higgs Balls

- No Higgs balls in SM, but may exist in BSM models (SUSY?)
- Stable at high temperatures

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

### Example: Higgs Balls

- No Higgs balls in SM, but may exist in BSM models (SUSY?)
- Stable at high temperatures
- Cosmological implications?

#### Thermal Q-balls

- Non-topological solitons that exist at finite temperature due to induced attractive interactions
- Requirement: Charged scalar field whose VEV determines bosonic masses
- Can have gauge charge (e.g., Higgs mechanism) or global charge (e.g., scalar interactions)

### Example: Higgs Balls

- No Higgs balls in SM, but may exist in BSM models (SUSY?)
- Stable at high temperatures
- Cosmological implications?

Thank you! Any questions?

• Higgs Balls are only *quasi-stable*: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.

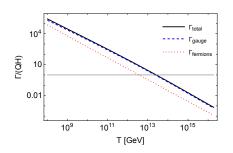
- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$

- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$
- In a cosmological context:

- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$
- In a cosmological context:
  - Higgs balls decay efficiently when  $\Gamma_{\rm Q-ball} \sim QH$  (where H is Hubble coefficient)

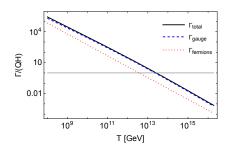
- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$
- In a cosmological context:
  - Higgs balls decay efficiently when  $\Gamma_{\mathrm{Q-ball}} \sim QH$  (where H is Hubble coefficient)
  - ▶ Equivalently, when  $\Gamma_h \sim H$

- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$
- In a cosmological context:
  - Higgs balls decay efficiently when  $\Gamma_{\mathrm{Q-ball}} \sim QH$  (where H is Hubble coefficient)
  - ightharpoonup Equivalently, when  $\Gamma_h \sim H$



 $Q=10^6$ , assuming radiation domination

- Higgs Balls are only quasi-stable: can't fall apart into individual Higgs quanta, but the Higgs quanta themselves can decay.
- ullet Decays occur throughout the volume of the Q-ball, so  $\Gamma_{\mathrm{Q-ball}} = Q \Gamma_h$
- In a cosmological context:
  - Higgs balls decay efficiently when  $\Gamma_{
    m Q-ball} \sim QH$  (where H is Hubble coefficient)
  - ightharpoonup Equivalently, when  $\Gamma_h \sim H$



 $Q=10^6$ , assuming radiation domination

Exist for wide range of temperatures, but stable against decay for  $T \gtrsim 10^{13}~{
m GeV}$ 

4 D > 4 A > 4 B > 4 B > 4 B = 900

Just because these states exists doesn't mean they were produced

Just because these states exists doesn't mean they were produced

No known production mechanism for gauged Q-balls!

Just because these states exists doesn't mean they were produced

- No known production mechanism for gauged Q-balls!
- Because of repulsive interaction, difficult to get large scale condensate of non-zero charge

Just because these states exists doesn't mean they were produced

- No known production mechanism for gauged Q-balls!
- Because of repulsive interaction, difficult to get large scale condensate of non-zero charge
- Therefore the usual production via fragmentation technique isn't (trivally) applicable

Just because these states exists doesn't mean they were produced

- No known production mechanism for gauged Q-balls!
- Because of repulsive interaction, difficult to get large scale condensate of non-zero charge
- Therefore the usual production via fragmentation technique isn't (trivally) applicable

Area for future work...