# Exploring mixed lepton-quark interactions in non-resonant leptoquark production at the LHC

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Work done in collaboration with: Antonio P. Morais, Roman Pasechnik, António Onofre and Werner Porod

Based on arXiv:2206.01674 and arXiv:2306.15460;

https://github.com/Mrazi09/LQ\_collider\_project

https://github.com/Mrazi09/Leptoquark-project

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Despite its successes, the SM either **fails** or is **insufficient** in tackling various observations, with **anomalous results** popping up in recent years . . .

• R<sub>D/D\*</sub>

$$R_{D^{(*)}} \equiv \frac{\operatorname{Br}\left(\bar{B} \to D^{(*)} \tau \bar{\nu}_{\tau}\right)}{\operatorname{Br}\left(\bar{B} \to D^{(*)} l \bar{\nu}_{l}\right)} \quad \text{with} \quad l = \mu, e,$$

 $2.3\sigma$  tension with the SM [Eur. Phys. J. C 81, 226 (2021)];

- $a_{\mu}$  4.2 $\sigma$  tension with the SM [Phys.Rev.Lett. 126, 141801 (2021)]. Although lattice results indicate consistency [Nature 593 (2021) 7857, 51-55].
- $M_W$  3.7 $\sigma$  tension with the SM [Phys.Rev.Lett. 129 (2022) 27, 271801]. Tension primarily driven by the new CDF result [Science 376 (2022) 170–176].
- U<sub>PMNS</sub> + m<sub>ν</sub> SM can not account for neutrino physics. Requires non-renormalisable 5D Weinberg operator [Phys. Lett. B 91 (1980) 51];

As shown in our previous work [2206.01674 [hep-ph]], a minimal LQ model is sufficient to address all of these open questions.

#### The model:

- One SU(2) singlet:  $S \sim (\mathbf{3}, \mathbf{1}, 1/3)^-$ ;
- One SU(2) doublet:  $R \sim (\bar{\bf 3}, {\bf 2}, 1/6)^-$ ;
- Yukawas that generate proton decay  $\to y_1 \bar Q^c Q S^\dagger + y_2 \bar d u S + \mathrm{h.c.}$

$$\mathbb{Z}_2 \text{ symmetry: } \mathbb{P}_B = (-1)^{3B+2S} \quad \Longrightarrow \quad S^-, \, R^-, \, q_{L,R}^+, \, \ell_{L,R}^-, \, H^-$$

Group charges and particles are emergent in a gauge-flavour unification model based on  $E_6\times SU(3)_F$  [Eur. Phys. J. C 80, 1162 (2020)].

- Yukawa Lagrangian:  $\mathcal{L}_Y = \Theta_{ij} \bar{Q}^c_j L_i S + \Omega_{ij} \bar{L}_i d_i R^\dagger + \Upsilon_{ij} \bar{u}^c_j e_i S + \mathrm{h.c.}$
- Relevant scalar potential:

$$V \supset -\mu^{2}|H|^{2} + \mu_{S}^{2}|S|^{2} + \mu_{R}^{2}|R|^{2} + \lambda(H^{\dagger}H)^{2} + g_{HR}(H^{\dagger}H)(R^{\dagger}R) + g'_{HR}(H^{\dagger}R)(R^{\dagger}H) + g_{HS}(H^{\dagger}H)(S^{\dagger}S) + (a_{1}RSH^{\dagger} + \text{h.c.}) .$$

ullet After EWSB, two physical (1/3)e appear in the spectrum

$$M_{LQ^{1/3}}^2 = \begin{bmatrix} \mu_S^2 + \frac{g_{HS}v^2}{2} & \frac{va_1}{\sqrt{2}} \\ \frac{va_1}{\sqrt{2}} & \mu_R^2 + \frac{(g_{HR} + g'_{HR})v^2}{2} \end{bmatrix}$$

Trade physical masses,  $m_{S_1^{1/3}}$  and  $m_{S_2^{1/3}}$ , for the  $\mu$  terms

$$\mu_{S}^{2} = \frac{1}{2} \left( m_{S_{1}^{1/3}}^{2} + m_{S_{2}^{1/3}}^{2} - g_{HS}v^{2} + \sqrt{(m_{S_{1}^{1/3}}^{2} - m_{S_{2}^{1/3}}^{2})^{2} - 2a_{1}^{2}v^{2}} \right),$$

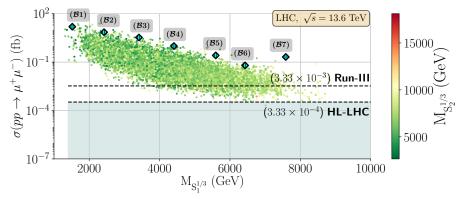
$$\mu_{R}^{2} = \frac{1}{2} \left( m_{S_{1}^{1/3}}^{2} + m_{S_{2}^{1/3}}^{2} - (g_{HR} + g'_{HR})v^{2} - \sqrt{(m_{S_{1}^{1/3}}^{2} - m_{S_{2}^{1/3}}^{2})^{2} - 2a_{1}^{2}v^{2}} \right).$$

Mixing controlled by the relation  $\sin(2\theta) = (\sqrt{2}va_1)(m_{S_1^{1/3}}^2 - m_{S_2^{1/3}}^2)^{-1}$ 

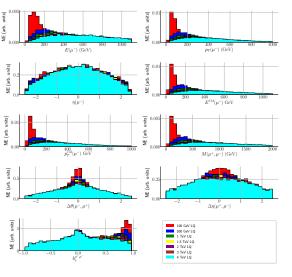
 $\bullet$  A single (2/3)e state appears and does not mix with  $m_{S^{2/3}}^2 = \mu_R^2 + \frac{g_{HR}v^2}{2}$  .

	u	c	t		d	s	b			,,-
e	$ \lambda_{eu}  \lesssim 0.22$ @ 1.5 TeV	$ \lambda_{ec}  \lesssim 2.0$ @ 1.5 TeV	$ \lambda_{et}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	e	$ \lambda_{ed}  \lesssim 0.22$ @ 1.5 TeV	$ \lambda_{es}  \lesssim 1.75$ @ 1.5 TeV	$ \lambda_{eb}  \lesssim 2.65$ @ 1.5 TeV		Υ	μ
$\mu$	$ \lambda_{\mu u}  \lesssim 0.8$ @ 1.5 TeV (This work)	$ A\mu c  \gtrsim 1.1$	$ \lambda_{\mu t}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	$\mu$	$ \lambda_{\mu d}  \lesssim 0.5$ @ 1.5 TeV	$ \lambda_{\mu s}  \lesssim 1.5$ @ 1.5 TeV	$ \lambda_{\mu b}  \lesssim 2.1$ @ 1.5 TeV		LQ	
au	$ \lambda_{\tau u}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	$ \lambda_{\tau c}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	$ \lambda_{\tau t}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	au	$ \lambda_{\tau d}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	$ \lambda_{\tau s}  \lesssim \sqrt{4\pi}$ @ 1.5 TeV	$ \lambda_{\tau b}  \lesssim 3.0$ @ 1.5 TeV	$\overline{u}$ $\overline{u}$	Υ	$\mu^+$

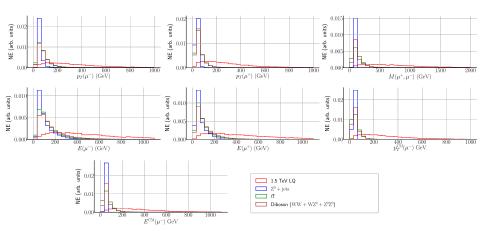
- Limits based on the combination of various channels (pair prod., single prod., recast) [JHEP 01 (2019) 132];
- Most searches tend towards diagonal channels. Current constraints are not strong for off-diagonal contributions. Particularly for 1st and 3rd generation couplings; Substantial phase-space left unprobed.
- $\bullet$  We focus our analysis for the  $\mu c$  coupling and study its implications at future runs of the LHC.



- Vast region of the parameter space in reach of both run-III and HL-LHC;
- Cross-section driven primarily by  $S_1^{1/3}$  due to small mixing  $\mathbf{a_1} \sim \mathcal{O}(1~{\rm GeV}) \Rightarrow \sin(2\theta) \ll 1;$
- Various benchmarks are defined for different mass scales of the lightest  $(1/3e)\ \text{LQ}.$

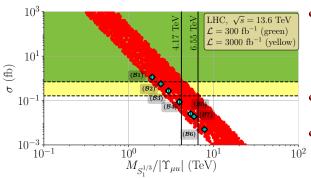


- Drell-Yan t-channel type process with two muons in the final state;
- LQ mass cannot be reconstructed. Although, kinematic distributions affected by it.
- Events generated at LO in MadGraph, with Delphes detector simulation.
   Backgrounds reweighted based on higher-order corrections;
- Selection criteria
  - $p_T(\mu^{\pm}) > 25 \text{ GeV}$
  - $|\eta(\mu^{\pm})| < 2.5$



- Dimension-full distributions offer the greater discriminating power. Angular distributions closely follow SM backgrounds;
- The signal tends to populate the high-energy regions, as opposed to the SM expectation;

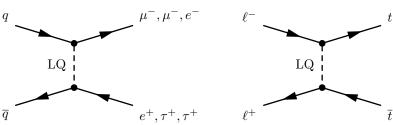
$M_{LQ}$ (GeV)	$M(\mu^+, \mu^-)$ $\mathcal{L} = 300 \text{ fb}^{-1}$	$M(\mu^+, \mu^-)$ $\mathcal{L} = 3000 \text{ fb}^{-1}$	$\mathcal{L} = 300 \text{ fb}^{-1}$	$\begin{array}{c c} E(\mu^+) \\ \mathcal{L} = 3000 \text{ fb}^{-1} \end{array}$	Combined $(300, 3000) \text{ fb}^{-1}$
1.5 TeV (B1)	$1.75\sigma$	$5.20\sigma$	$0.891\sigma$	$2.72\sigma$	$(3.06\sigma, 9.72\sigma)$
2.5 TeV (B2)	$0.573\sigma$	$2.28\sigma$	$0.744\sigma$	$2.33\sigma$	$(1.15\sigma, 4.97\sigma)$
3.5 TeV (B3)	$0.128\sigma$	$0.912\sigma$	$0.225\sigma$	$1.04\sigma$	$(0.288\sigma, 1.97\sigma)$



- The combination of the various distributions can extend the discovery threshold of our model's LQs;
- $M/\Upsilon < 4.17 \text{ TeV @ 95\%}$ CL for run-III;
- $M/\Upsilon < 6.55~{
  m TeV}$  @ 95% 2 CL for HL-LHC.

#### To conclude . . .

- Discussed a simple economical model and how it solves current open questions;
- Showed that there are regions of parameter space that can be probed at LHC (run-III and/or HL-LHC)



• Flavour off-diagonal channels also possible. LQ Yukawa couplings with top more easily measured in lepton colliders.

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## Thank you for your attention

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