

SUSY Conference 2023

Precise Estimate of Chargino Decay

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Based on

JHEP 01 (2023) 017 (arXiv: 2210.16035)

and work in progress

In Collaboration with

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July 19th, 2023, University of Southampton

Plan of Talk

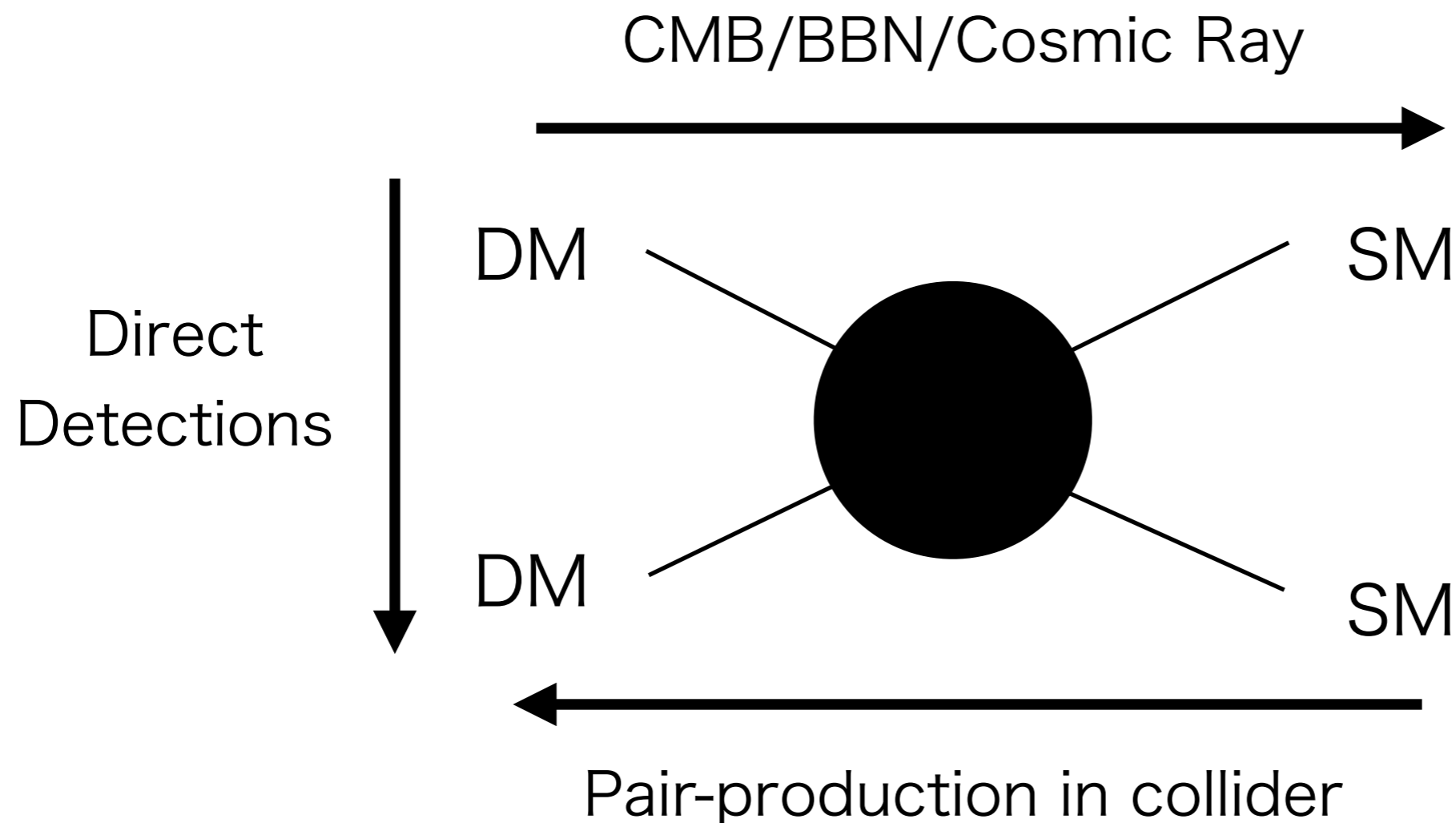
1. Introduction
2. Wino Decay and Mass Difference
3. EW Corrections in Single Pion Mode
4. Results and Summary

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Dark Matter in SUSY

- ▶ The MSSM contains an attractive candidate of freeze-out DM
- ▶ The supersymmetric DM can be tested with experiments in various way



Wino and Higgsino

Wino

- ▶ $SU(2)_L$ triplet
- ▶ Zero hypercharge
- ▶ One charged Dirac
+ One neutral Majorana

Higgsino

- ▶ $SU(2)_L$ doublet
- ▶ $Y = \pm 1/2$
- ▶ One charged Dirac
+ Two neutral Majorana

Today's discussion

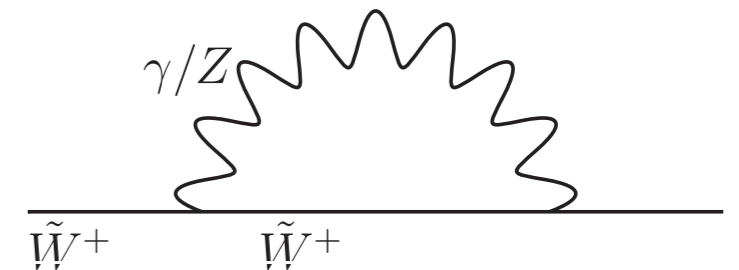
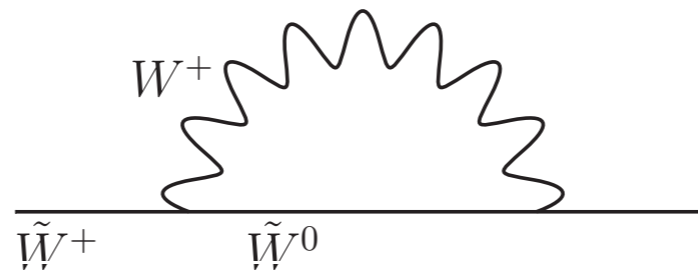
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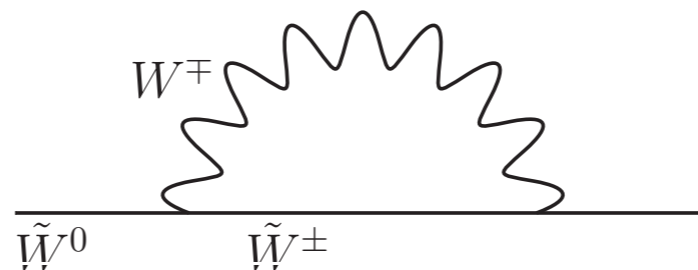
Mass Splitting and Decay

- ▶ The charged Wino becomes heavier than the neutral one because of EW radiative corrections
- ▶ At one-loop level,

Charged Wino

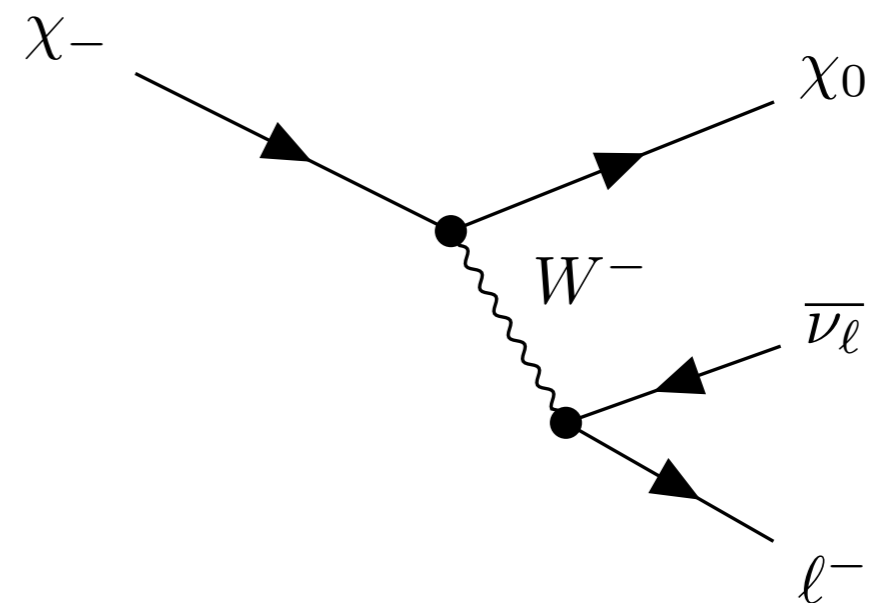
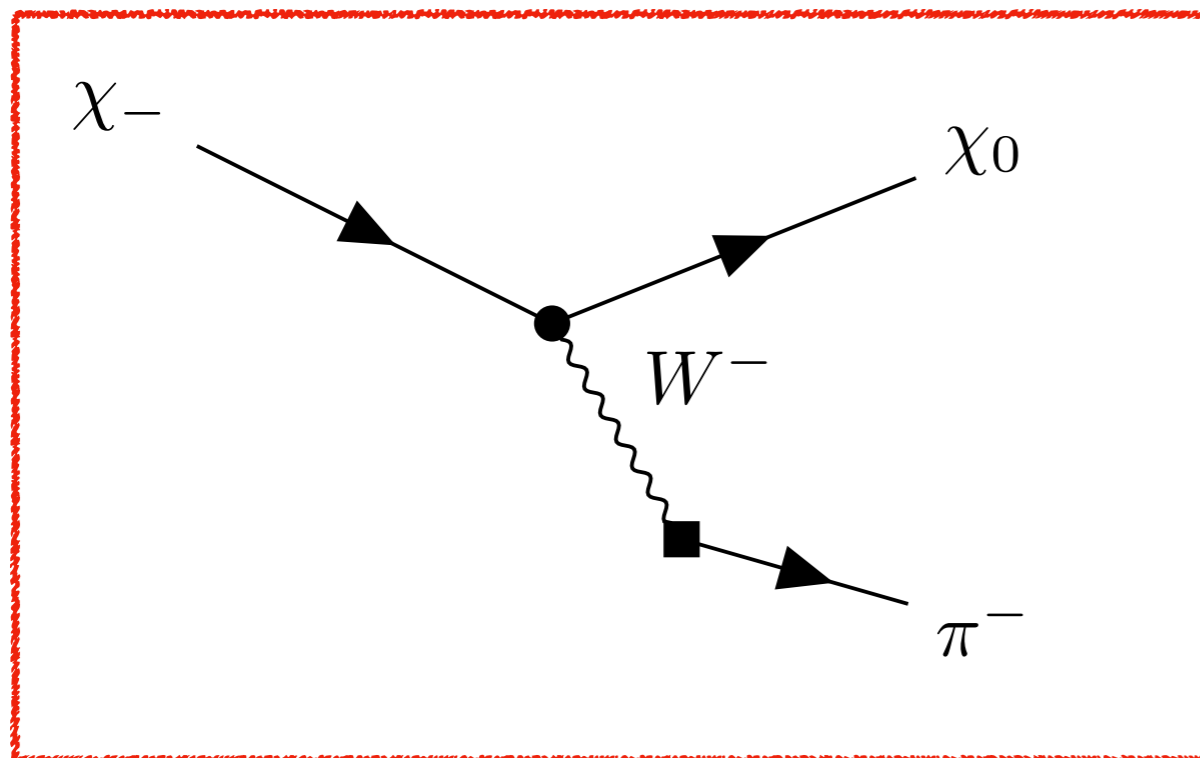


Neutral Wino



Mass Splitting and Decay

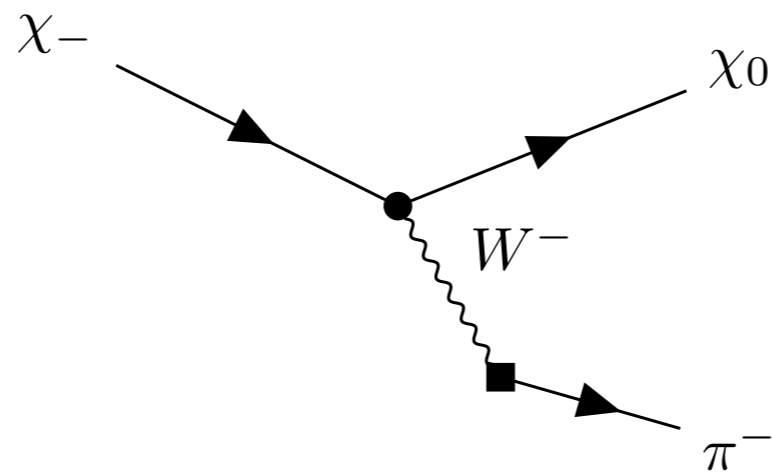
- The charged Wino can decay into the lighter neutral Wino through weak interaction



Dominant for $\Delta m \sim 160 \text{ MeV}$

Mass Splitting and Decay

- The theoretical prediction of the charged Wino decay rate is very sensitive to the mass difference

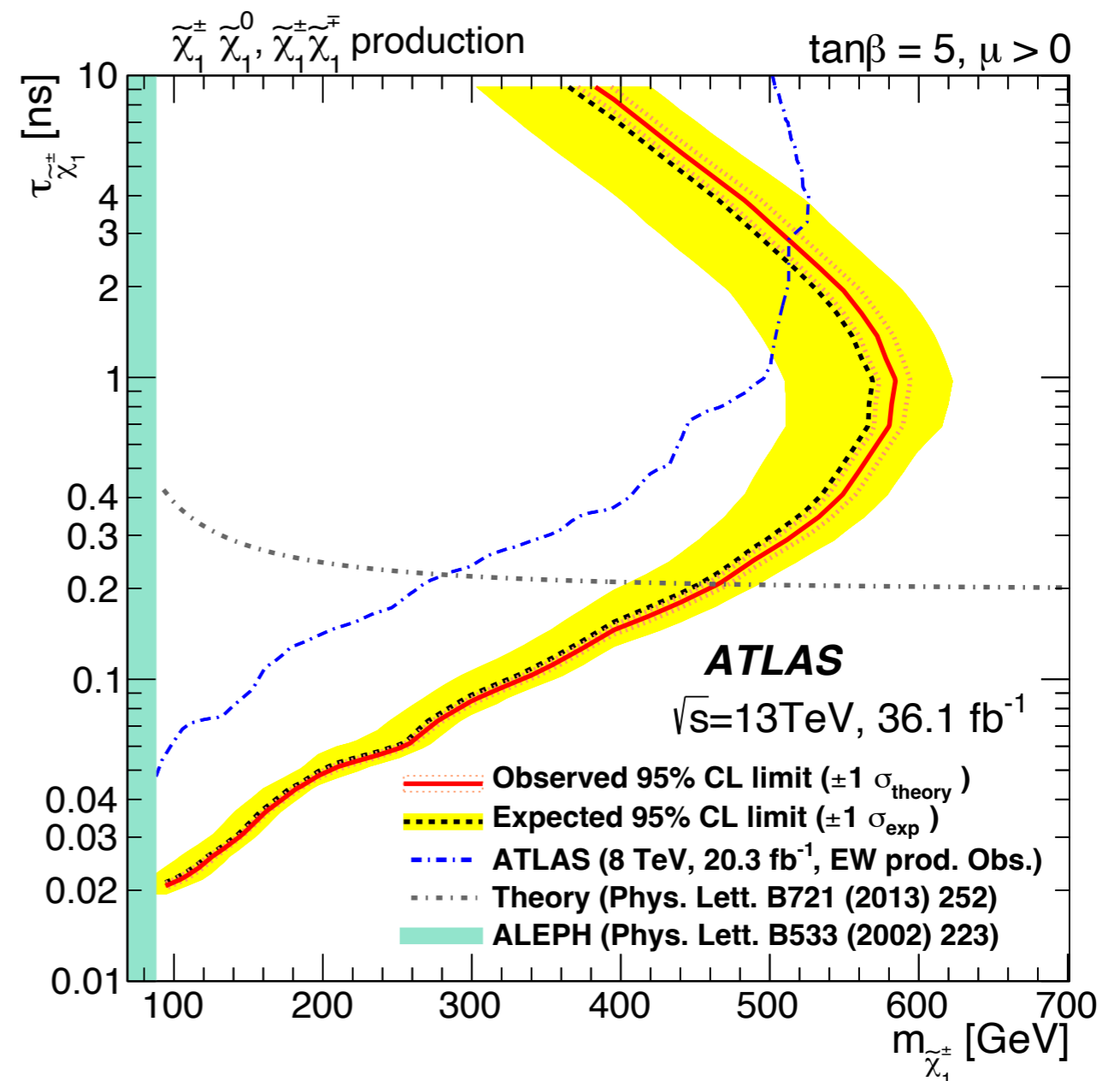
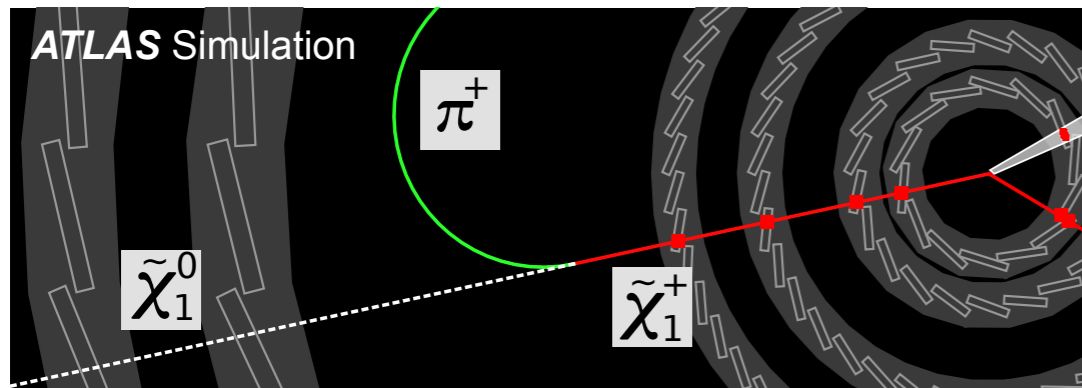


$$\Gamma_{\text{tree}}(\chi^- \rightarrow \pi^- + \chi^0) \simeq \frac{4}{\pi} F_{\pi}^2 (G_{\pi}^0)^2 \Delta m^3 \left(1 - \frac{m_{\pi}^2}{\Delta m^2} \right)^{1/2}$$

2% error in $\Delta m \rightarrow$ about $3 \times 2\% \simeq 6\%$ in the decay rate

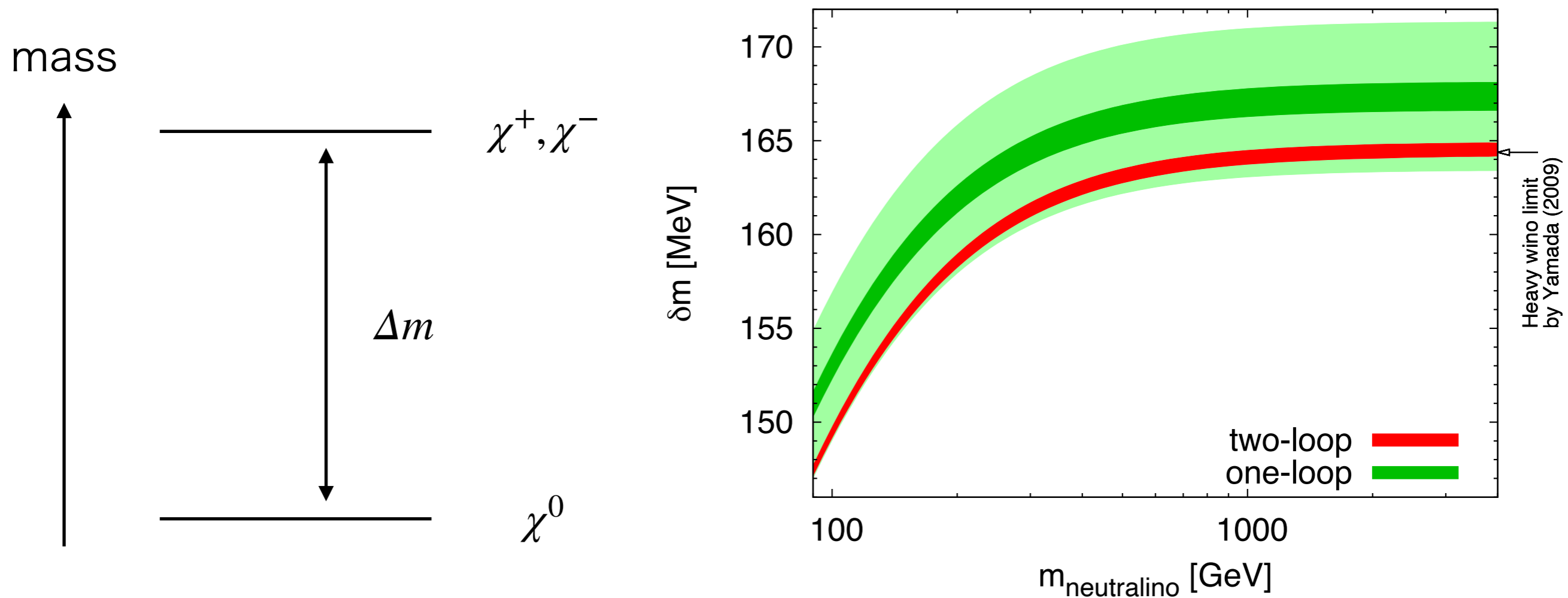
Experimental Constraint

- Collider constraint on the Wino mass strongly depends on the rate of the single pion mode



NNLO Wino Mass Difference

- ▶ Actually in literature the mass difference of SU(2) triplet has already been computed in two-loop level!



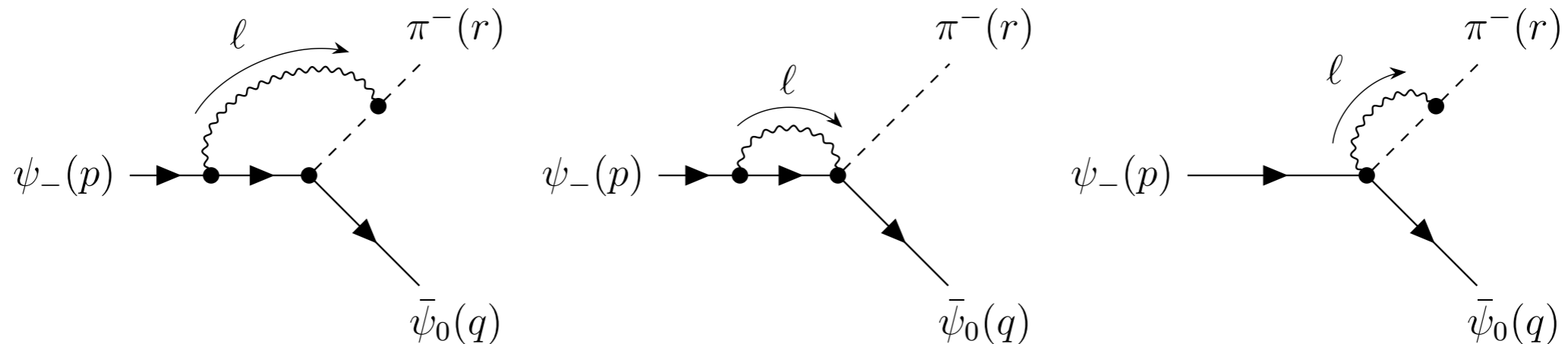
[M. Ibe, R. Sato, S. Matsumoto, Phys. Lett. B 721 (2013) 252-260]

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Remaining Corrections

How about EW corrections to the decay process itself?



Scales of the problem

$$m_\chi \gg m_W \gg \Delta m \gtrsim m_\pi$$

Question

Is there large contribution such as $\log(m_\chi/m_\pi)$?

How are experimental constraint changed?

Violation of Shift Symmetry

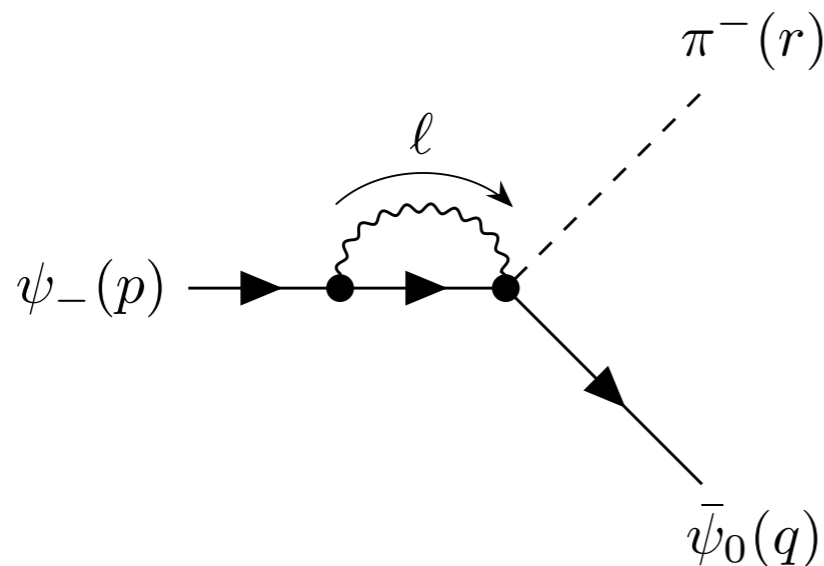
Pion's shift symmetry ensures that the decay amplitude is suppressed by Δm

$$\mathcal{L}_{\text{Wino-Pion}} = -2\sqrt{2}f_{\pi}G_{\pi}^0(\partial_{\mu}\pi^{-}) \times (\bar{\psi}_{-}\gamma^{\mu}\psi_0) + \text{h.c.}$$

$$p \sim \Delta m$$

But QED corrections break the symmetry:

$$\partial_{\mu} \rightarrow D_{\mu} = \partial_{\mu} - ieA_{\mu}$$



mass correction to the Wino

$m_{\chi}/\Delta m$ enhancement???

→ Cured by appropriate determination of CTs

Computational Scheme

Matching procedure à la Descotes-Genon and Moussallam

[S. Descotes-Genon and B. Moussallam, Eur. Phys. J. C 42, 403 (2005)]

EW theory w/ Wino

$$\mathcal{L} = \mathcal{L}_{\text{EW}} + \mathcal{L}_{\text{Wino}}$$

Match free-quark decay rate

$$\Gamma(\chi^- \rightarrow \chi^0 + \bar{u} + d) \text{ @ 1-loop}$$

Four-Fermi theory w/ Wino

$$\mathcal{L} = \mathcal{L}_{\text{Four-Fermi}} + \mathcal{L}_{\text{Wino}} + \mathcal{L}_{\text{CT}}$$

ChPT w/ Wino

$$\mathcal{L} = \mathcal{L}_{\text{ChPT}} + \mathcal{L}_{\text{Wino}} + \mathcal{L}_{\text{CT}}$$

Match the “current correlator”
@ 1-loop

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Single Pion Mode (pure-Wino)

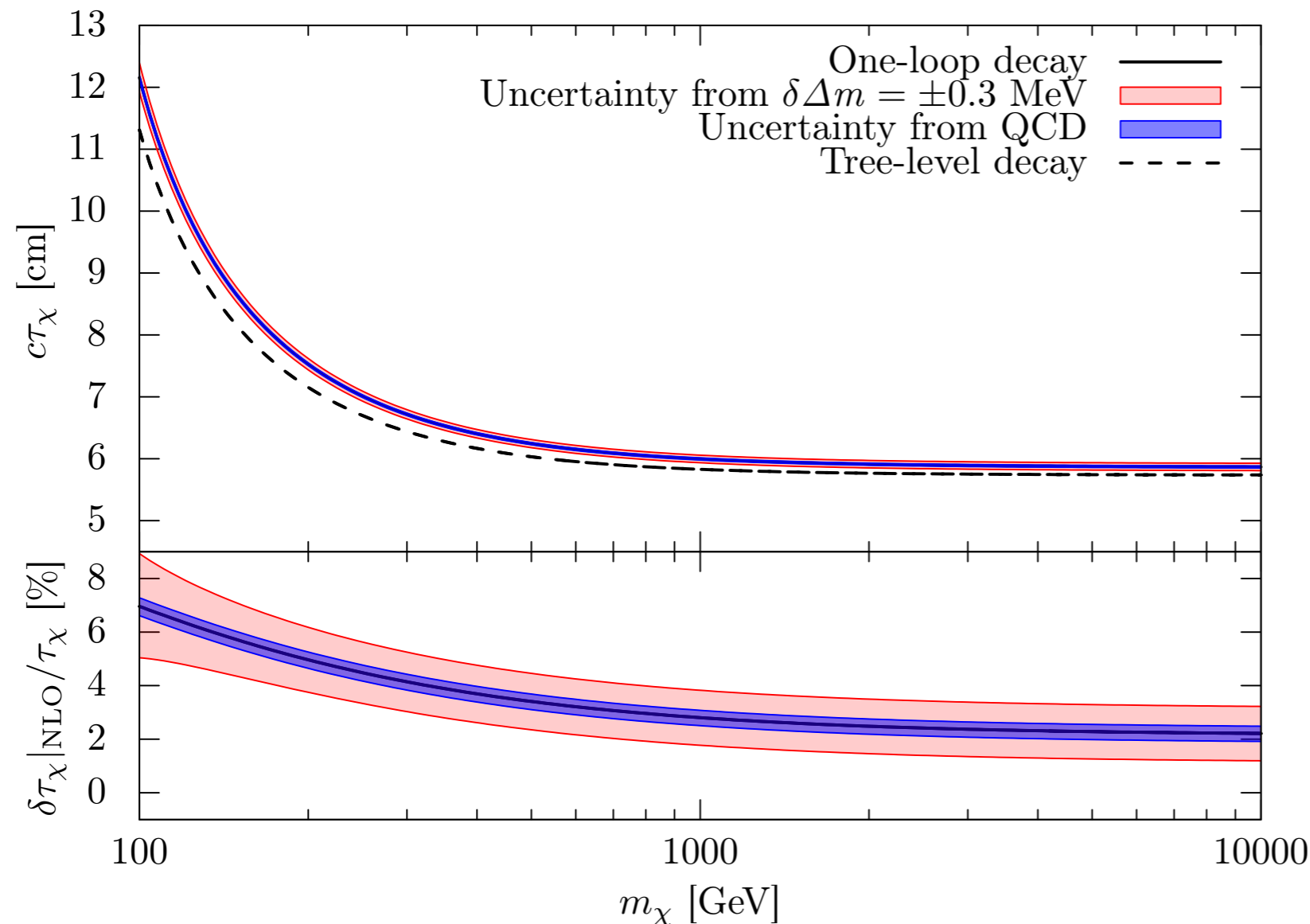
In total, ~ 0.03

$$\Gamma = \Gamma_{\text{tree}} \left\{ 1 + \frac{\alpha}{4\pi} \left[\underbrace{\sum_{n<0} c^{(n)} \left(\frac{\Delta m}{m_\chi} \right)^n}_{\text{Exactly canceled by CT}} + \underbrace{c^{(\log)} \log \left(\frac{m_\chi}{\Delta m} \right)}_{\text{Canceled between various contributions}} + \sum_{n=0}^{\infty} c^{(n)} \left(\frac{\Delta m}{m_\chi} \right)^n \right] \right\}$$

Exactly canceled by CT

Canceled between
various contributions

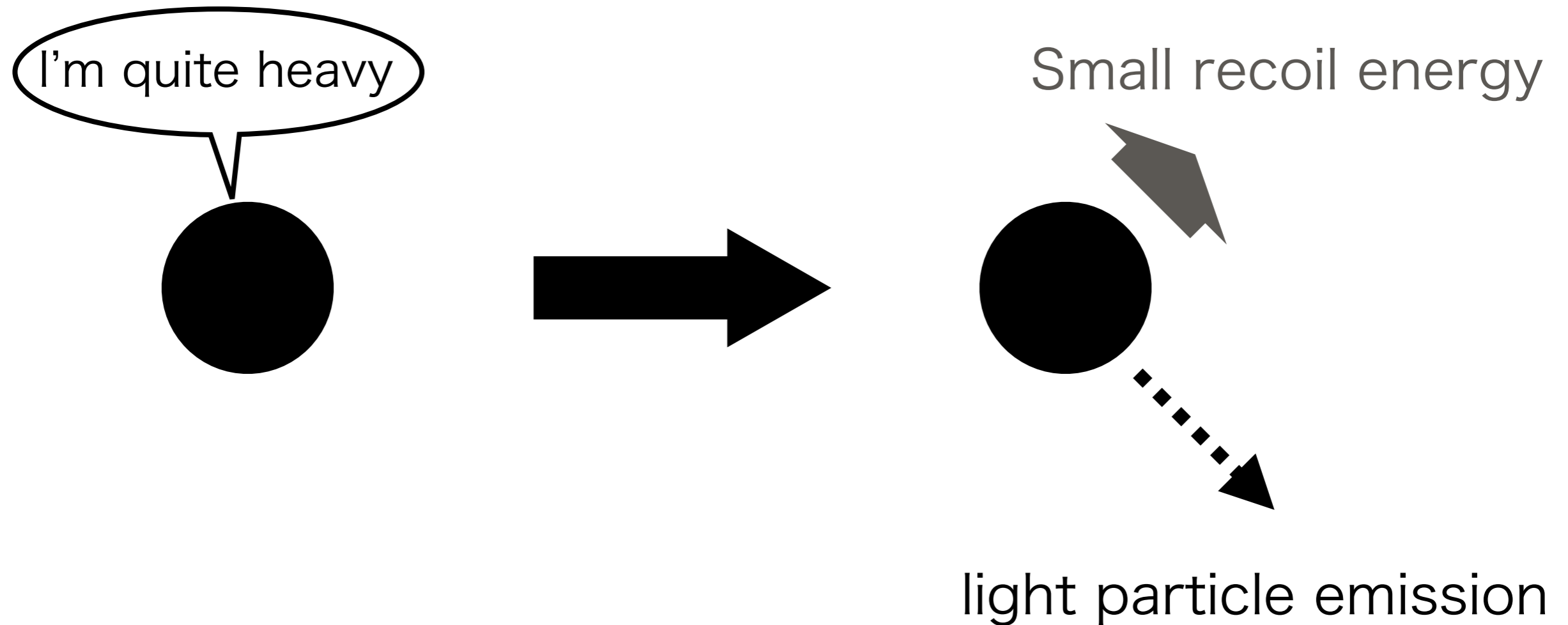
Single Pion Mode (pure-Wino)



$$\delta\tau_\chi|_{\text{NLO}} \equiv \tau_\chi|_{\text{NLO}} - \tau_\chi|_{\text{LO}}$$

- ▶ Apart from uncertainties, Wino decay length become around 2-7% longer, depending on m_χ
- ▶ Radiative correction tends to be a constant as $m_\chi \rightarrow \infty$
- ▶ The 3-loop effect on Δm dominates the uncertainty

A New Theorem?

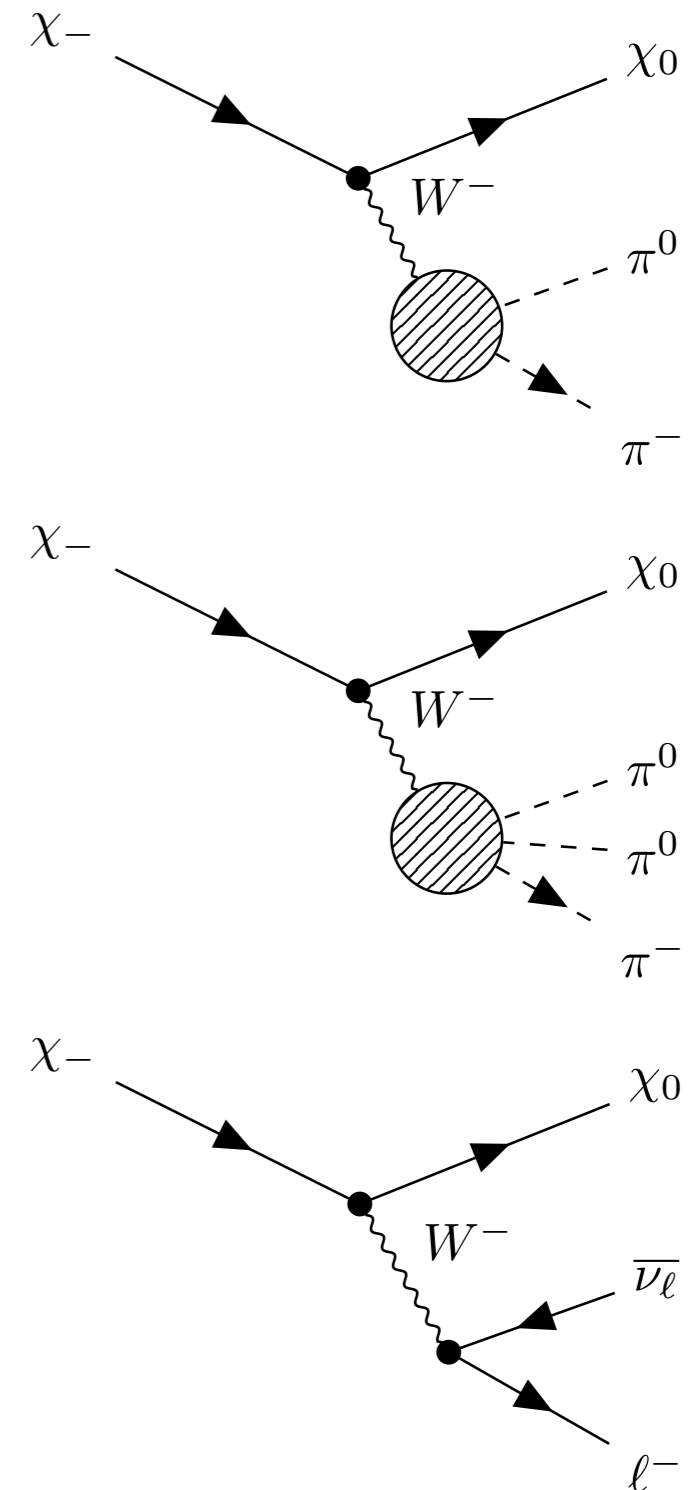
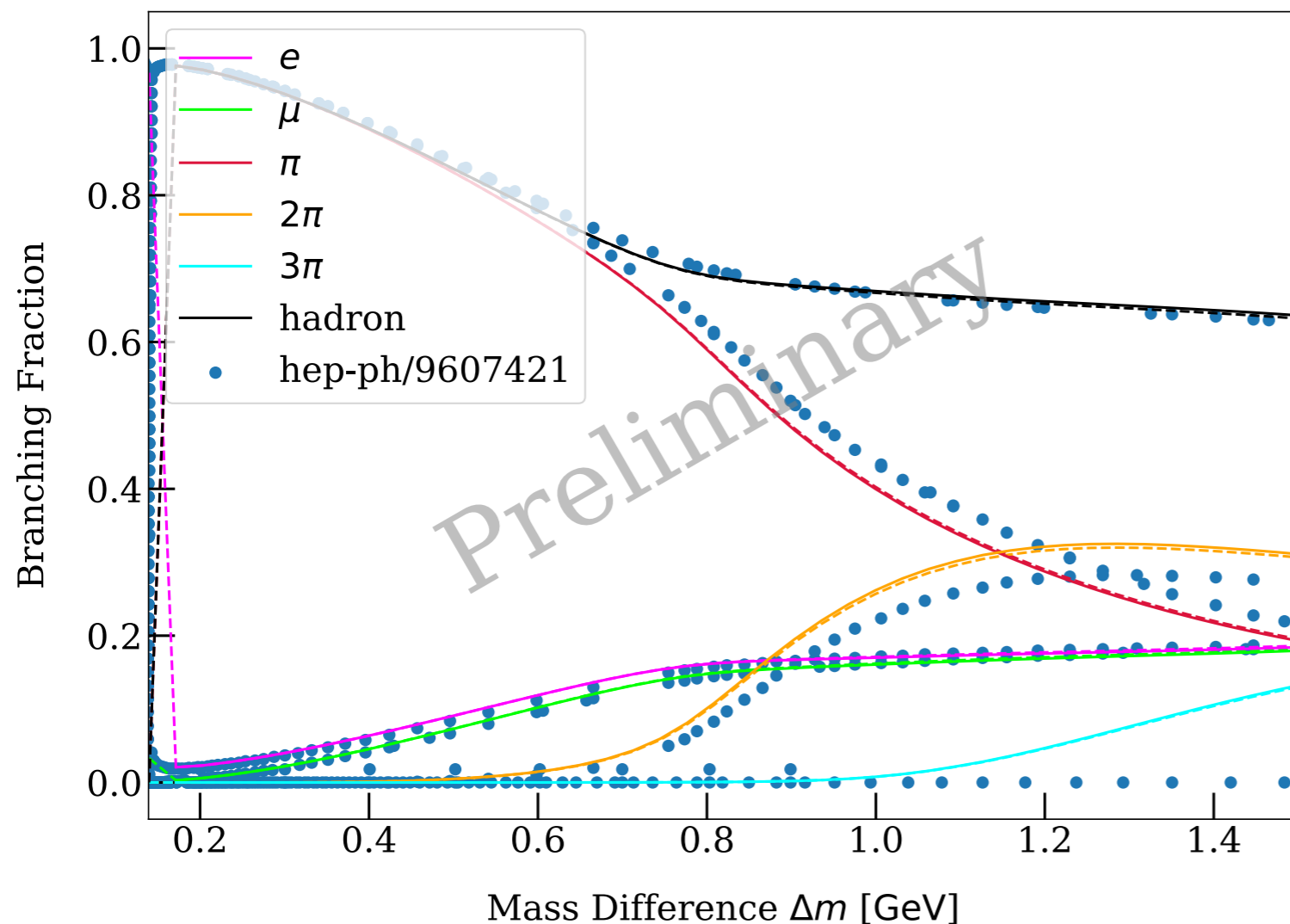


$$\Gamma = \sum_{n < 0} \left(\frac{\Delta m}{m_\chi} \right)^n \Gamma^{(n)} + \log \left(\frac{m_\chi}{\Delta m} \right) \Gamma^{(\log)} + \sum_{n=0}^{\infty} \left(\frac{\Delta m}{m_\chi} \right)^n \Gamma^{(n)}$$

should be zero at all order of α

Comment on Higgsino Case

Higgsino has larger mass difference,
so various modes should be considered
Updating the previous work in progress



Summary

- ▶ We computed the single pion mode at EW one-loop; and obtained $O(1)\%$ correction with 0.5% th. error
- ▶ No $m_\chi/\Delta m$ or $\log m_\chi$ enhancement at one-loop level;
We conjecture that heavy **external** particles decouple from physical quantities at any order, although it has not been proofed rigorously
- ▶ Higgsino can decay into heavier state due to the larger mass difference.
Stay tuned for our numerical estimate of Higgsino decay!

Back up

Mass Splitting

Chargino-neutralino mass matrices

$$M_{\tilde{N}} = \begin{pmatrix} M_1 & 0 & g_1 \langle H_d^0 \rangle / \sqrt{2} & -g_1 \langle H_u^0 \rangle / \sqrt{2} \\ 0 & M_2 & -g_2 \langle H_d^0 \rangle / \sqrt{2} & g_2 \langle H_u^0 \rangle / \sqrt{2} \\ g_2 \langle H_d^0 \rangle / \sqrt{2} & -g_2 \langle H_d^0 \rangle / \sqrt{2} & 0 & \mu \\ -g_2 \langle H_u^0 \rangle / \sqrt{2} & g_2 \langle H_u^0 \rangle / \sqrt{2} & \mu & 0 \end{pmatrix}$$

$$X = \begin{pmatrix} M_2 & -g_2 \langle H_u^0 \rangle \\ -g_2 \langle H_d^0 \rangle & -\mu \end{pmatrix}$$

induce tree-level mixing and mass splittings:

$$\mathcal{L}_{\text{MSSM}} \supset g_2 W_\mu^- (\bar{\Psi}_\chi^-)_i \gamma^\mu (O_{ij}^L P_L + O_{ij}^R P_R) (\Psi_\chi^0)_j$$

$$m_{\tilde{N}_1} < m_{\tilde{N}_2} < m_{\tilde{N}_3} < m_{\tilde{N}_4}, \quad m_{\tilde{C}_1} < m_{\tilde{C}_2}$$

Computational Scheme

The counterterm contribution to the Wino decay rate

Cancels the Δm enhancement

$$\frac{\delta\Gamma_\chi}{\Gamma_\chi}\bigg|_{K,Y} = e^2 \left[\frac{8}{3}(K_1 + K_2) + \frac{20}{9}(K_5 + K_6) + 4K_{12} \right. \\ \left. - \hat{Y}_6 - \frac{4}{3}(Y_1 + \hat{Y}_1) - 4 \left(Y_2 + \hat{Y}_2 - \frac{m_\chi}{\Delta m} Y_3 \right) \right]$$

Matching with the FF theory

(Pole of ChPT) + (Finite part of the FF theory's CT)

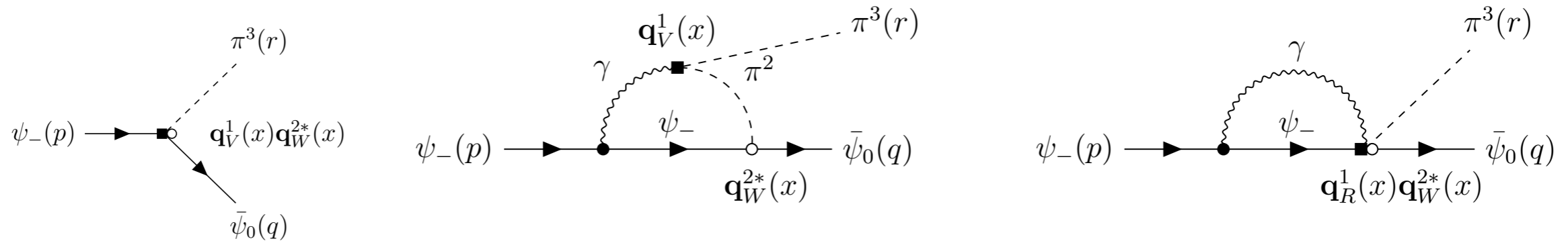
Matching with the EW theory

(Pole of ChPT) + (The EW theory's input parameters)

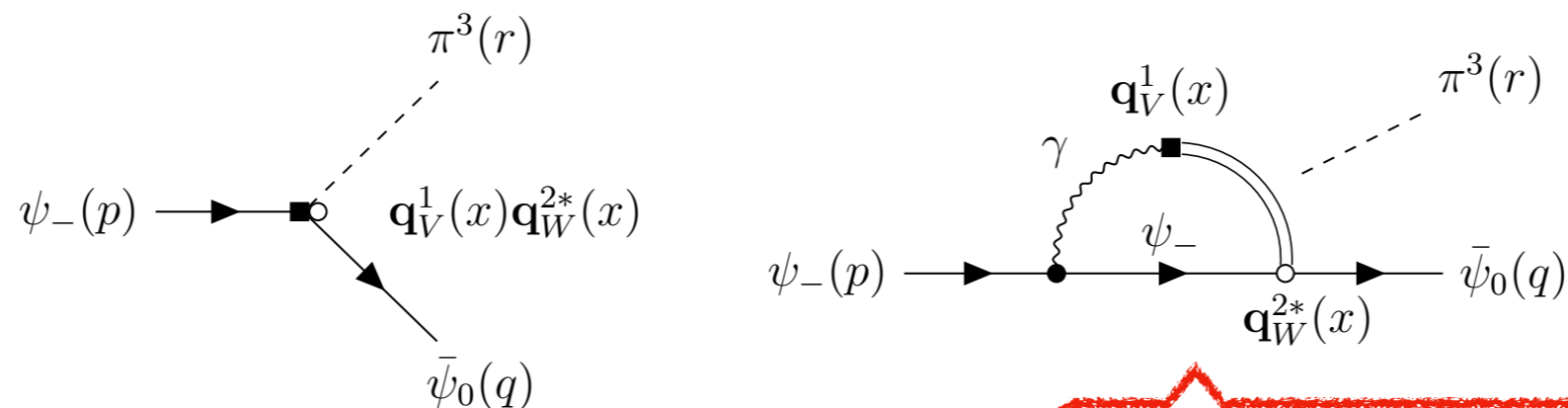
FF-ChPT Matching

Matching example $\langle 0|T\psi_-(x)\bar{\psi}_0(y)\pi^3(x)|0\rangle$

► In the ChPT w/ Wino:

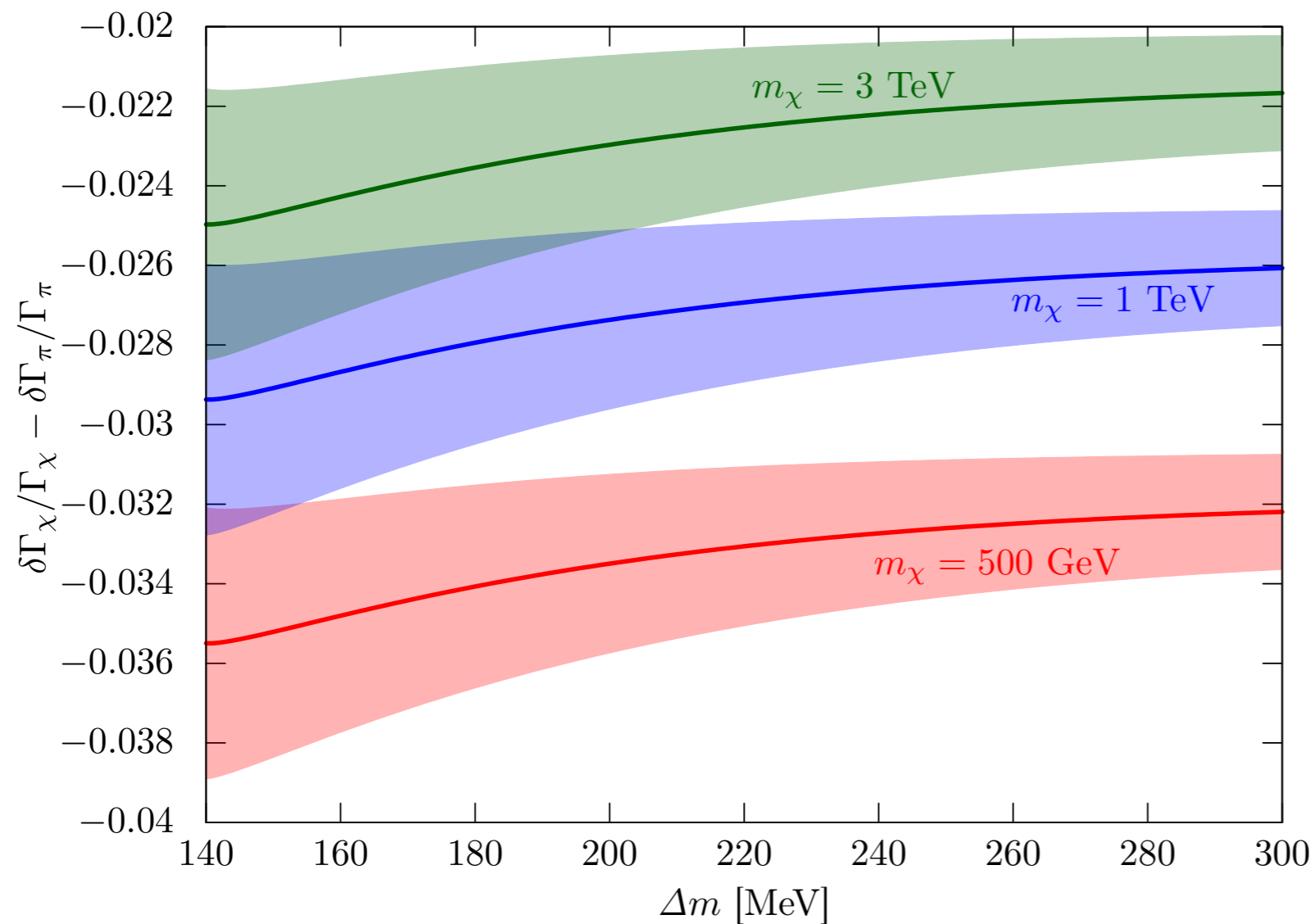


► In the FF theory w/ Wino:



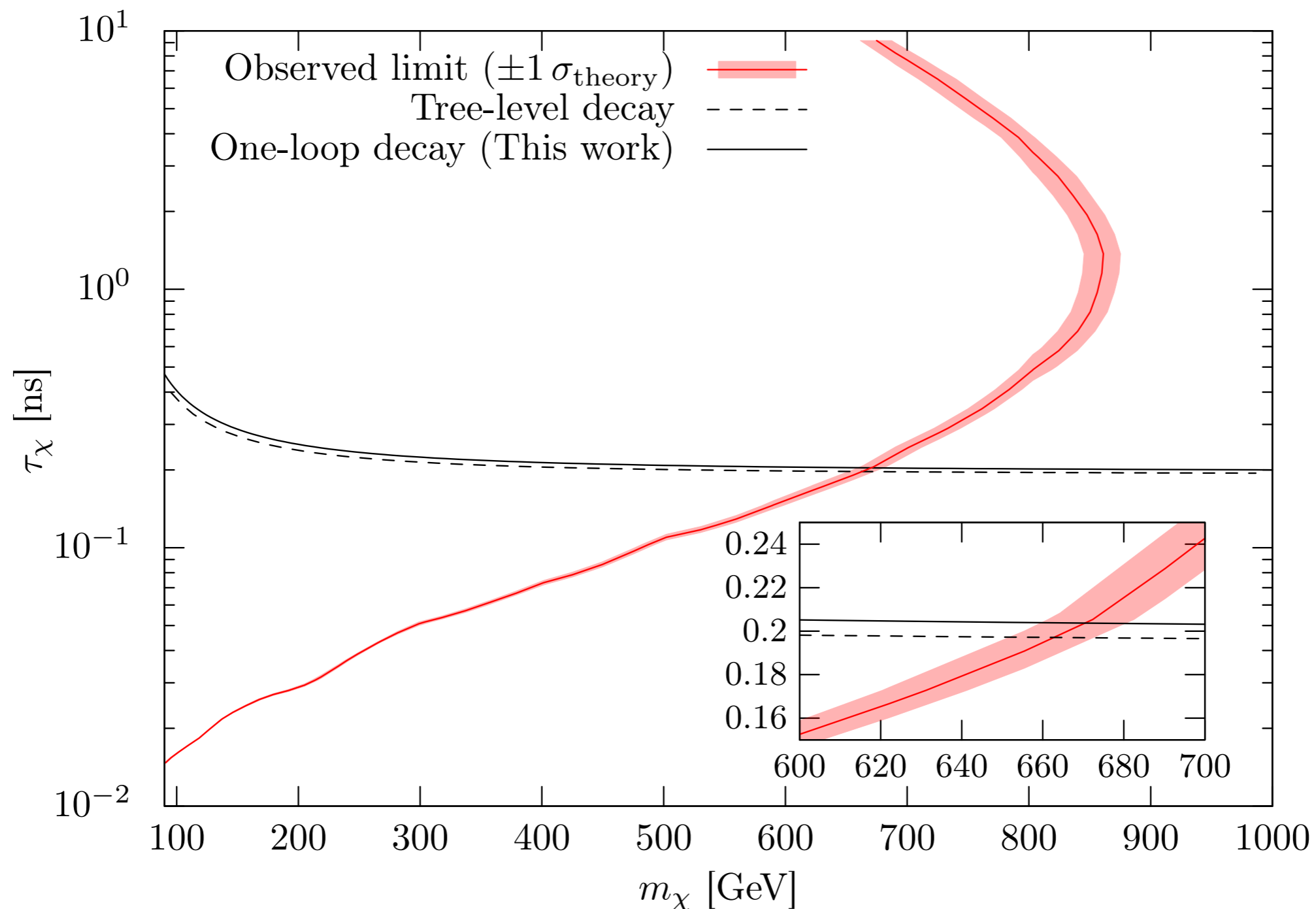
requires $\langle \pi^a(r)|T J_V^{b\mu}(x) J_A^{c\nu}(0)|0\rangle$

Theory Error



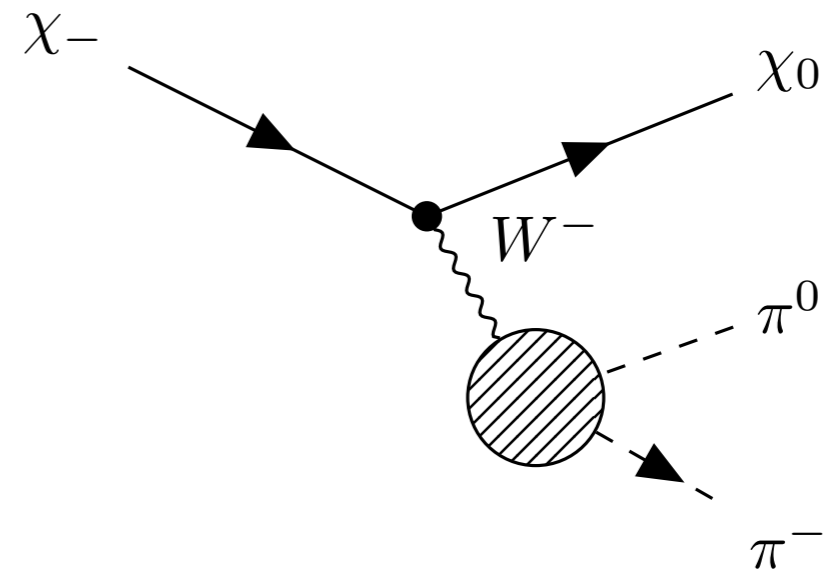
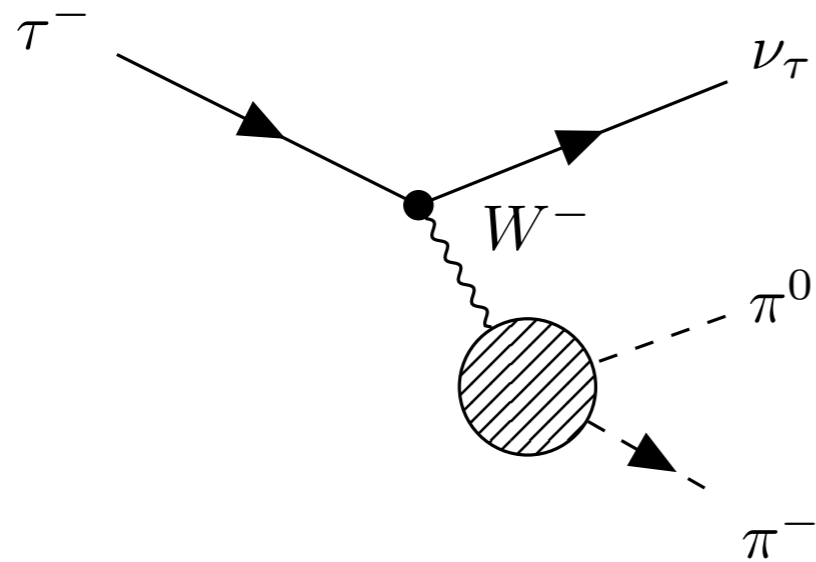
- Each band represents uncertainty from the minimal resonance model
- In the pure Wino theory, $\Delta m \sim 163 - 165 \text{ MeV}$ for $m_\chi > 600 \text{ MeV}$
- For this mass difference, one-loop contribution reduces Γ_χ/Γ_π about **2-4%**

Updated ATLAS Constraint



Treatment of Hadronization

We have very similar decay process; tau decay



Extract
mass distributions

Interpolate
or
parametrize
data

$$\langle \pi^- \pi^0 | \bar{d} \gamma^\mu u | 0 \rangle$$

Nonperturbative QCD encoded

Data of Mass Distributions

Experimental data of **non-strange** spectral functions is available **up to tau lepton mass squared**

[ALEPH Collaboration, Phys. Rept. 421 (2005) 191]

