# PART 3



### WHAT WILL THIS LECTURE BE ABOUT?

#### INTRODUCTION

• Definitions and basic concepts

#### **INPUT TO THE PHYSICS**

- The data: trigger, data preparation
- The theory: Monte carlo simulations
- Reconstruction, or how to translate detector signals to particles

#### **PHYSICS ANALYSES**

- Through example, step-by-step
- Discussion of analysis methods



Is there a topic you would like to add to this material? If so: please let me know at the end of this lecture and I will see if I can add it!

# RECONSTRUCTION



### WHAT DO WE RECONSTRUCT?

• Tracks and clusters

Combining those:
"objects", i.e. "particles"



#### **Simplified Detector Transverse View**



#### **RECONSTRUCTION - FIGURES OF MERIT**



#### **RECONSTRUCTION - FIGURES OF MERIT**

	DEFINITION	EXAMPLE		NEEDS BE:
EFFICIENCY	how often do we reconstruct the object we are interested in	electron identification efficiency = (number of reconstructed electrons) / (number of true electrons) in bins of transverse momentum	$\begin{array}{c} 0.95 \\ 0.95 \\ 0.85 \\ 0.8 \\ 0.75 \\ 0.7 \\ $	High
RESOLUTION	how accurately do we reconstruct the quantity	energy resolution = (measured energy – true energy)/(true energy)	$\sigma = (1.12 \pm 0.03)\%$	<b>Good</b> (a small number)
FAKE RATE	how often we reconstruct a different object as the object we are interested in	a jet faking an electron, fake rate = (Number of jets reconstructed as an electron) / (Number of jets) in bins of pseudorapidity	$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	Low

#### **RECONSTRUCTION - GOALS**

- High efficiency
- Good resolution
- Low fake rate

Noise

Dead regions of the detector

Increased pile-up

#### **©** Computing-friendly-

CPU time per event
Memory use

![](_page_6_Figure_11.jpeg)

![](_page_6_Figure_12.jpeg)

### WHAT DO WE RECONSTRUCT?

![](_page_7_Picture_1.jpeg)

- Combining those:
  - "objects", i.e. "particles"

![](_page_7_Figure_4.jpeg)

#### **©** For a track we measure:

- Its momentum;
- Its direction;
- Its charge;
- Its "perigee": the closest point to a reference line, transverse ( $d_0$ ) or longitudinal ( $z_0$ ).

Tracks are key ingredients of most of particle reconstruction.

### TRACKING IN A NUTSHELL

![](_page_8_Figure_8.jpeg)

![](_page_8_Picture_9.jpeg)

### TRACKING IN A NUTSHELL - TRACK FITTING

![](_page_9_Picture_1.jpeg)

Perfect measurement – ideal

![](_page_9_Picture_3.jpeg)

Imperfect measurement – reality

![](_page_9_Picture_5.jpeg)

**©** Small errors and more points help to constrain the possibilities

![](_page_9_Figure_7.jpeg)

- **©** Quantitatively:
  - Parameterize the track;
  - Find parameters by Least-Squares-Minimization;
  - Obtain also uncertainties on the track parameters.

### TRACKING IN A NUTSHELL - TRACK FITTING

#### **©** For a track we measure:

- Its momentum;
- Its direction;
- Its charge;
- Its "perigee": the closest point to a reference line, transverse ( $d_o$ ) or longitudinal ( $z_o$ ).

![](_page_10_Figure_6.jpeg)

### TRACKING IN A NUTSHELL - TRACK FITTING

#### **©** For a track we measure:

- ◎ Its momentum;
- Its direction;
- Its charge;
- Its "perigee": the closest point to a reference line, transverse ( $d_0$ ) or longitudinal ( $z_0$ ).
- And their uncertainty

#### **Small uncertainties are required.**

- <sup>(©)</sup> δdo is < O(10 $\mu$ m) and δθ < O(0.1mrad).
- Allows separation of tracks that come from different particle decays (which can be separated at the order of mm).

![](_page_11_Figure_10.jpeg)

#### Presence of Material

- Coulomb scattering off the core of atoms
- Energy loss due to ionization
- Bremsstrahlung
- Hadronic interaction

#### Misalignment

- Detector elements not positioned in space with perfect accuracy.
- Alignment corrections derived from data and applied in track reconstruction.

![](_page_12_Figure_9.jpeg)

#### **IMPACT OF GOOD ALIGNMENT**

Improving the tracker alignment description in the reconstruction gives better track momentum resolution which leads to better mass resolution.

![](_page_13_Figure_2.jpeg)

- Can see the reconstructed Z width gets narrower if we use better alignment constants. Very important for physics analysis to have good alignment.
- Alignment of detector elements can change with time, for example when the detector is opened for repair, or when the magnetic field is turned on and off.

### WHAT DO WE RECONSTRUCT?

![](_page_14_Picture_1.jpeg)

- Combining those:
  - "objects", i.e. "particles"

![](_page_14_Figure_4.jpeg)

#### A CALORIMETER VIEW

![](_page_15_Figure_1.jpeg)

### **CLUSTERING IN A NUTSHELL**

![](_page_16_Picture_1.jpeg)

Reconstruct energy deposited in the calorimeter by charged or neutral particles;
 electrons, photons and jets.

#### For a cluster we measure:

- ◎ The energy;
- ◎ The position of the deposit;
- The direction of the incident particles;

#### Calorimeters are segmented in cells.

Typically, a shower created by a particle interacting with the matter extends over several cells.

#### Various clustering algorithms, e.g.:

- Sliding window. Sum cells within a fixed-size rectangular window.
- Topo-clustering. Start with a seed cell and iteratively add to the cluster the neighbor of a cell already in the cluster.

#### CLUSTER FINDING - AN EXAMPLE

CMS crystal calorimeter – ECAL clusters
 electron energy in central crystal ~80%,

in 5x5 matrix around it ~96%.

![](_page_17_Figure_3.jpeg)

![](_page_17_Picture_4.jpeg)

#### CLUSTER FINDING - AN EXAMPLE

![](_page_18_Figure_1.jpeg)

#### Simple example of an algorithm

- Scan for seed crystals = local energy maximum above a defined seed threshold
- @ Starting from the seed position, adjacent crystals are examined, scanning first in  $\phi$  and then in  $\eta$
- - <sup>©</sup> The crystal's energy is above the noise level (lower threshold)
  - <sup>©</sup> The crystal has not been assigned to another cluster already

#### CLUSTER FINDING - AN EXAMPLE: DIFFICULTIES

Oreful tuning of thresholds needed.

- needs usually learning phase;
- adapt to noise conditions;
- ◎ too low : pick up too much unwanted energy;
- ◎ too high : loose too much of "real" energy. Corrections/Calibrations will be larger.

![](_page_19_Figure_6.jpeg)

### WHAT DO WE RECONSTRUCT?

![](_page_20_Figure_1.jpeg)

![](_page_21_Figure_0.jpeg)

### **ELECTRONS / PHOTONS**

Final Electron momentum measurement can come from tracking or calorimeter information (or a combination of both)

- Often have a final calibration to give the best electron energy
- Working points define categories
  - © E.g. loose, medium, tight
  - Trade-off: Efficiency vs Fakes
- Often want "isolated electrons"
  - Require little calorimeter energy or tracks in the region around the electron

![](_page_22_Figure_8.jpeg)

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![](_page_23_Picture_8.jpeg)

### ELECTRONS / PHOTONS - BACKGROUNDS

#### **Sources of backgrounds:**

Hadronic jets leaving energy in calorimeter

- While calorimeter clusters are much wider for jets than for electrons/photons there are many thousands more jets than electrons
   rate of jets faking an electron needs to be very small (~10<sup>-4</sup>)
- Complex identification algorithms are required to give the rejection whilst keeping a high efficiency

### **ELECTRONS – IDENTIFICATION ALGOS**

Signal

Signal

Signal

Background

Background

0 900 10 Δ E. (MeV)

Background

![](_page_25_Figure_1.jpeg)

Information can be exploited using multi-variate techniques such as **likelihood discriminants** or **boosted decision trees** or **other machine learning methods**.

Example of different calorimeter shower shape variables used to distinguish electron showers from jets in ATLAS

![](_page_26_Picture_0.jpeg)

 Combine the muon segments found in the muon detector with tracks from the tracking detector

- Momentum of muon determined from bending due to magnetic field in tracker and in muon system
  - Combine measurements to get best resolution
  - Need an accurate map of magnetic field in the reconstruction software
  - Alignment of the muon detectors also very important to get best momentum resolution

![](_page_26_Figure_6.jpeg)

### MUONS ON ATLAS

**Simplified Detector Transverse View** 

"MS" - Muon Spectrometer

![](_page_27_Figure_2.jpeg)

JETS

![](_page_28_Figure_1.jpeg)

#### **JET PRODUCTION PROCESSES**

![](_page_29_Figure_1.jpeg)

#### Jets are produced:

- by fragmentation of gluons and (light) quarks in QCD scattering
- by decays of heavy Standard Model particles, e.g. W & Z
- in association with particle
   production in Vector Boson Fusion,
   e.g. Higgs
- In decays of beyond the Standard Model particles, e.g. in SUSY

JETS

![](_page_30_Figure_1.jpeg)

![](_page_30_Figure_2.jpeg)

At low energy, jets are more likely produced by gluon fusion.

![](_page_30_Figure_4.jpeg)

### JET ALGORITHMS

• Theory requirements: infrared and collinear safe

![](_page_31_Figure_2.jpeg)

Soft gluon radiation should not merge jets

![](_page_31_Figure_4.jpeg)

![](_page_31_Picture_5.jpeg)

...and on signal split in two possibly below threshold

- Experimental requirements: Independent to detector technology and data taking conditions, easily implementable
- Jet algorithm commonly used at the LHC: 'anti-k<sub>t</sub>'. A 'recursive recombination' algorithm. Starts from (topo-)clusters. Hard stuff clusters with nearest neighbor. Various cone sizes (standard R=0.4/0.5, "fat" R=1.0).

![](_page_31_Figure_9.jpeg)

### JET CALIBRATION

- Correct the energy and position measurement and the resolution.
- Account for:

Instrumental effects Detector inefficiencies 'Pile-up' Electronic noise Clustering, noise suppression Dead material losses Detector response Algorithm efficiency

#### Physics effects

Algorithm efficiency 'Pile-up' 'Underlying event'

![](_page_32_Figure_6.jpeg)

#### JETS AND PILE-UP

![](_page_33_Picture_1.jpeg)

Multiple interactions from pile-up

![](_page_33_Figure_3.jpeg)

#### **B-JETS**

b-hadrons have a lifetime of ~ 10<sup>-12</sup> s.
They travel a small distance (fraction of mm) before decaying.
A "displaced vertex" creates a distinct jet, so b-jets can be tagged (b-tagged).
b-tagging uses sophisticated algorithms, mostly multi-variate (machine learning).

![](_page_34_Picture_3.jpeg)

![](_page_35_Picture_0.jpeg)

### MISSING TRANSVERSE MOMENTUM – ME<sub>T</sub>

![](_page_36_Figure_1.jpeg)

In the transverse plane:

$$\Sigma_i \vec{p}_{T,i} = 0$$

So for what we can't directly measure (e.g. neutrinos)

$$E_{\rm T}^{\rm miss} = -\Sigma_i \vec{p}_{T,i}$$

![](_page_36_Figure_6.jpeg)

### MISSING TRANSVERSE MOMENTUM – ME<sub>T</sub>

![](_page_37_Picture_1.jpeg)

In the transverse plane:

$$\Sigma_i \vec{p}_{T,i} = 0$$

OR DARK MATTER CANDIDATES!

So for what we can't directly measure (e.g. neutrinos)

$$E_{\rm T}^{\rm miss} = -\Sigma_i \vec{p}_{T,i}$$

#### **Simplified Detector Transverse View Muon Spectrometer Toroids** HadCAL **EMCAL** photon Solenoid electrol TRT SCT **Pixels** muon κv

## PARTICLE FLOW FOR HADRONIC RECONSTRUCTION

#### PARTICLE FLOW

![](_page_39_Picture_1.jpeg)

#### PARTICLE FLOW

![](_page_40_Picture_1.jpeg)

### PARTICLE FLOW

- Reconstruct and identify all particles, photons, electrons, pions, …
- Use best combination of all subdetectors for measuring the properties of the particles.
- First used at LEP (ALEPH) and then at the LHC (CMS).

![](_page_41_Figure_5.jpeg)

#### JETS IN PILE-UP

![](_page_42_Picture_1.jpeg)

Multiple interactions from pile-up

### JETS IN PILE-UP

![](_page_43_Figure_1.jpeg)

175

Multiple interactions from pile-up

![](_page_44_Figure_1.jpeg)

**Resolution**: the quality with which we measure the jet momentum.

![](_page_45_Figure_1.jpeg)

**Resolution**: the quality with which we measure the jet momentum.

![](_page_46_Figure_1.jpeg)

**Resolution**: the quality with which we measure the jet momentum.

![](_page_47_Figure_1.jpeg)

Significant improvement for low-pT jets. Similar for MET.

![](_page_48_Figure_1.jpeg)

In Jet Energy resolution and uncertainty, large improvements with respect to calo jets!

#### A COMPARISON

![](_page_49_Figure_1.jpeg)

PF jets (CMS) and calo jets (ATLAS) have similar performance.
 Particle reconstruction always needs to be optimized depending on the detector technologies and experimental requirements.

#### A COMPARISON

![](_page_50_Figure_1.jpeg)

 PF jets (CMS) and calo jets (ATLAS) have similar performance.
 Particle reconstruction always needs to be optimized depending on the detector technologies and experimental requirements.

![](_page_51_Picture_0.jpeg)

![](_page_51_Picture_1.jpeg)

Objective:Trigger ("online") reconstruction same as "offline".Problem:Time. Trigger decision needs to be taken fast.Solution:Simplification.Challenge:Clever simplification = good performance.

![](_page_52_Figure_2.jpeg)

E.g. track reconstruction in regions of interest and simplified MET calculation.

![](_page_53_Picture_0.jpeg)

• To profit fully from an improvement in reconstruction, the relevant algorithm has to be used at the relevant trigger selections to provide **optimal online-to-offline correlation**.

![](_page_53_Figure_2.jpeg)

Variable A: e.g. leading jet pT

#### **ONLINE RECONSTRUCTION**

trigger efficiency =  $\frac{\# \text{ events passing offline selection } \& \text{ trigger}}{\# \text{ events passing offline selection}}$ 

![](_page_54_Figure_2.jpeg)

Clever ideas need to be deployed to bring online closer to offline, making efficiency curves **sharper** and **plateau closer to 1**.

### **EFFICIENCY MEASUREMENTS**

Relevant beyond the trigger...

#### TAG AND PROBE

- Select events based on requirements on one object (tag) and study the response of the second object (probe), not used in the event selection, using some constraint such as the Z mass.
  - e.g.  $Z \rightarrow \tau \tau$  events.
  - Typically used for measurement of the identification efficiency

#### ORTHOGONAL SAMPL

- Measure directly the efficiency on an independent, orthogonal sample.
  - e.g. jet trigger efficiency on a sample triggered by muons,

#### BOOTSTRAP METHOD

• The efficiency,  $\varepsilon_B$ , of a selection B, inclusive compared to a selection A, can be determined in a sample of events passing selection A (provided that  $\varepsilon_{A}$  is measurable):  $\varepsilon_{B} = \varepsilon_{B|A} \times \varepsilon_{A}$ . ε ε<sub>ΒΙΑ</sub>

• e.g. trigger efficiencies, say B: tau50 loose & A: tau16 loose

## PHYSICS MENUS

Trigger selection	2015 offline threshold (GeV)	2016 offline threshold (GeV)	2017 offline threshold (GeV)	2022 offline threshold (GeV)	Representative physics case	
Peak Luminosity	5x10 <sup>33</sup> cm <sup>-2</sup> s <sup>-1</sup>	1.2x10 <sup>34</sup> cm <sup>-2</sup> s <sup>-1</sup>	1.7x10 <sup>34</sup> cm <sup>-2</sup> s <sup>-1</sup>	2.0x10 <sup>34</sup> cm <sup>-2</sup> s <sup>-1</sup>		
isolated single e	25	27	27	27	"Main" triggers. Thrs driven by Higgs (ZH, WH), Top, SUSY.	
isolated single μ	21	27	27	25		
di-γ	40, 30	40, 30	40, 30	40, 30	Higgs (H→γγ, HH→bbγγ).	
di-τ (+ jet)	40, 30	40, 30	40, 30	40, 30	Higgs (H→ττ, HH→bbττ), SUSY.	
four-jet (incl. HF)	45	45	45	45	SUSY, Higgs, exotics	
MET	180	200	200	200		

Offline selections from which the triggers are "usable", i.e. at efficiency plateau or highly efficient otherwise

#### **RECONSTRUCTING PARTICLES**

![](_page_57_Figure_1.jpeg)

![](_page_57_Figure_2.jpeg)

#### TAUS

Tau Decay Mode				
Leptonic		$\tau^{\pm} \rightarrow e^{\pm} + \nu + \nu$	17.8%	
		$\tau^{\pm} \rightarrow \mu^{\pm} + \nu + \nu$	17.4%	
Hadronic	1-prong	$\tau^{\pm} \rightarrow \pi^{\pm} + \nu$	11%	
		$\tau^{\pm} \rightarrow \pi^{\pm} + \nu + n\pi^{\circ}$	35%	
	3-prong	$\tau^{\pm} \rightarrow 3\pi^{\pm} + \nu$	9%	
		$\tau^{\pm} \rightarrow 3\pi^{\pm} + \nu + n\pi^{\circ}$	5%	
Other			~5%	

Hadronic tau reconstruction extremely challenging
 Using multi-variate (machine learning) techniques
 based on track multiplicity and shower shapes

![](_page_58_Figure_3.jpeg)

![](_page_59_Picture_0.jpeg)

![](_page_59_Figure_1.jpeg)

### AND THE HIGGS!

![](_page_60_Figure_1.jpeg)

### HOW ABOUT NEW PARTICLES?

• These decay to Standard Model particles or create  $\ensuremath{\mathsf{ME}_{\mathsf{T}}}$ 

![](_page_61_Figure_2.jpeg)

#### **PHYSICS ANALYSES**

![](_page_62_Figure_1.jpeg)