

PART 4



WHAT WILL THIS LECTURE BE ABOUT?

INTRODUCTION

- Definitions and basic concepts

INPUT TO THE PHYSICS

- The data: trigger, data preparation
- The theory: Monte carlo simulations
- Reconstruction, or how to translate detector signals to particles

PHYSICS ANALYSES

- Through example, step-by-step
- Discussion of analysis methods



*Is there a topic you would like to add to this material?
If so: please let me know at the end of this lecture and I will see if I can add it!*

LHC PHYSICS
AN ANALYSIS
STEP-BY-STEP

PHYSICS ANALYSES

Measurements

- ⊙ Allow important tests of the consistency of the theory.
- ⊙ Typically limited by systematic uncertainties.

“Systematic” uncertainties are introduced by inaccuracies in the methods used to perform the measurement.

Searches

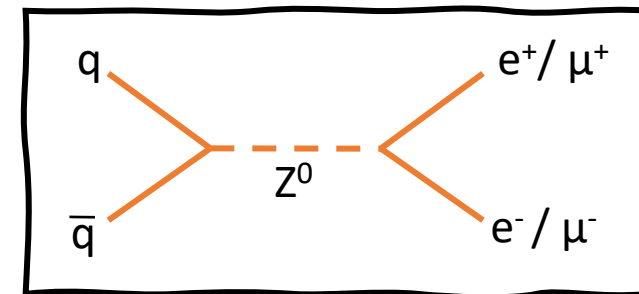
- ⊙ ... For new particles.
- ⊙ If no signal, set limits on some model.
- ⊙ If signal, a potential discovery!
- ⊙ More data typically improve a search.

SIMPLE EXAMPLE:

MEASURING THE Z^0 CROSS-SECTION AT LHC

⊙ Z^0 boson decays to lepton or quark pairs

⊙ We can reconstruct it in the e^+e^- or $\mu^+\mu^-$ decay modes



⊙ Discovery and study of the Z^0 boson was a critical part of understanding the electroweak force.



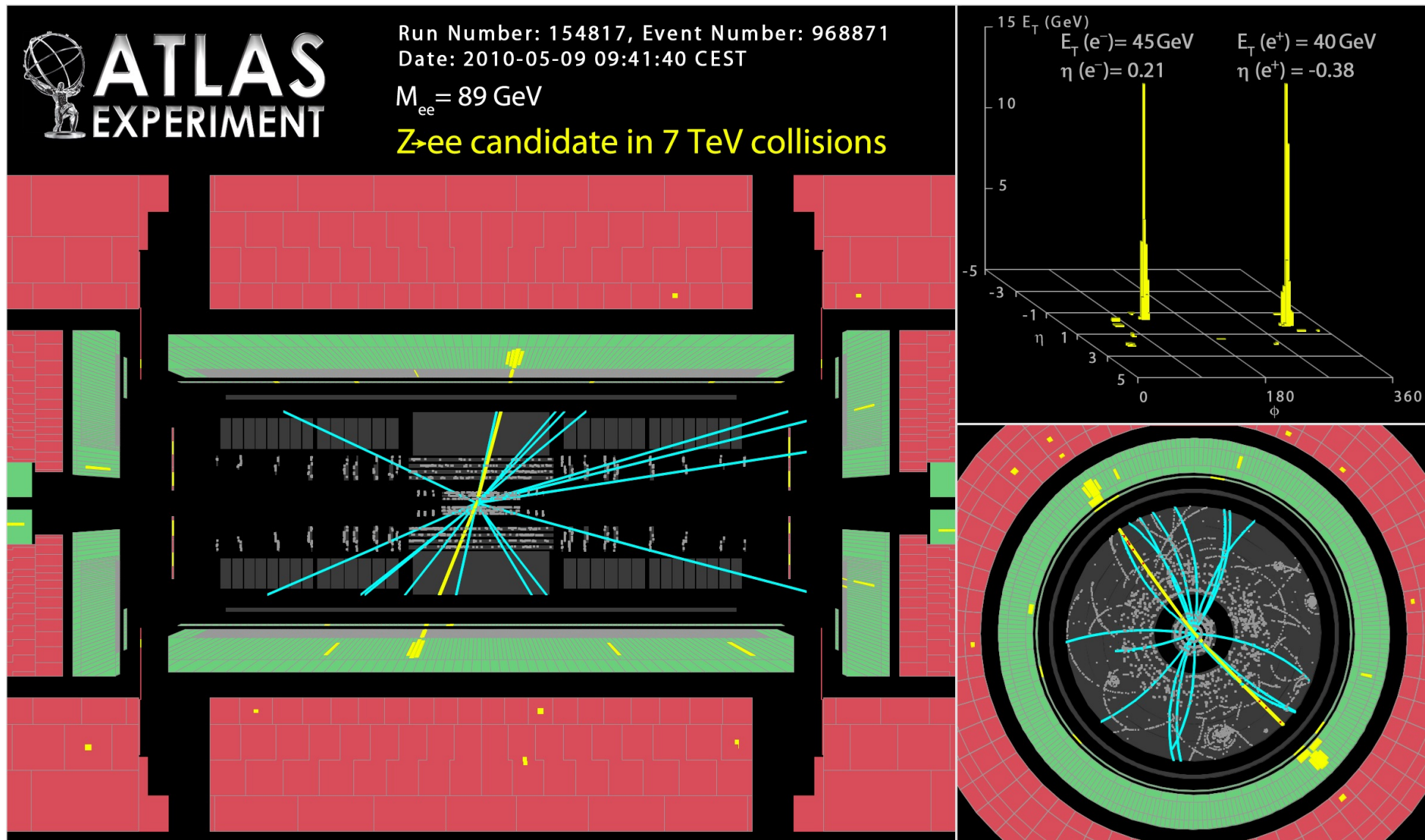
⊙ And now, at the LHC?

⊙ **Important test of theory:** does the measurement agree with the theoretical prediction at LHC collision energy?

⊙ **A standard candle** for studying reconstruction and deriving calibrations.

⊙ Can be used for luminosity determination!

MEASURING THE Z^0 CROSS-SECTION AT LHC

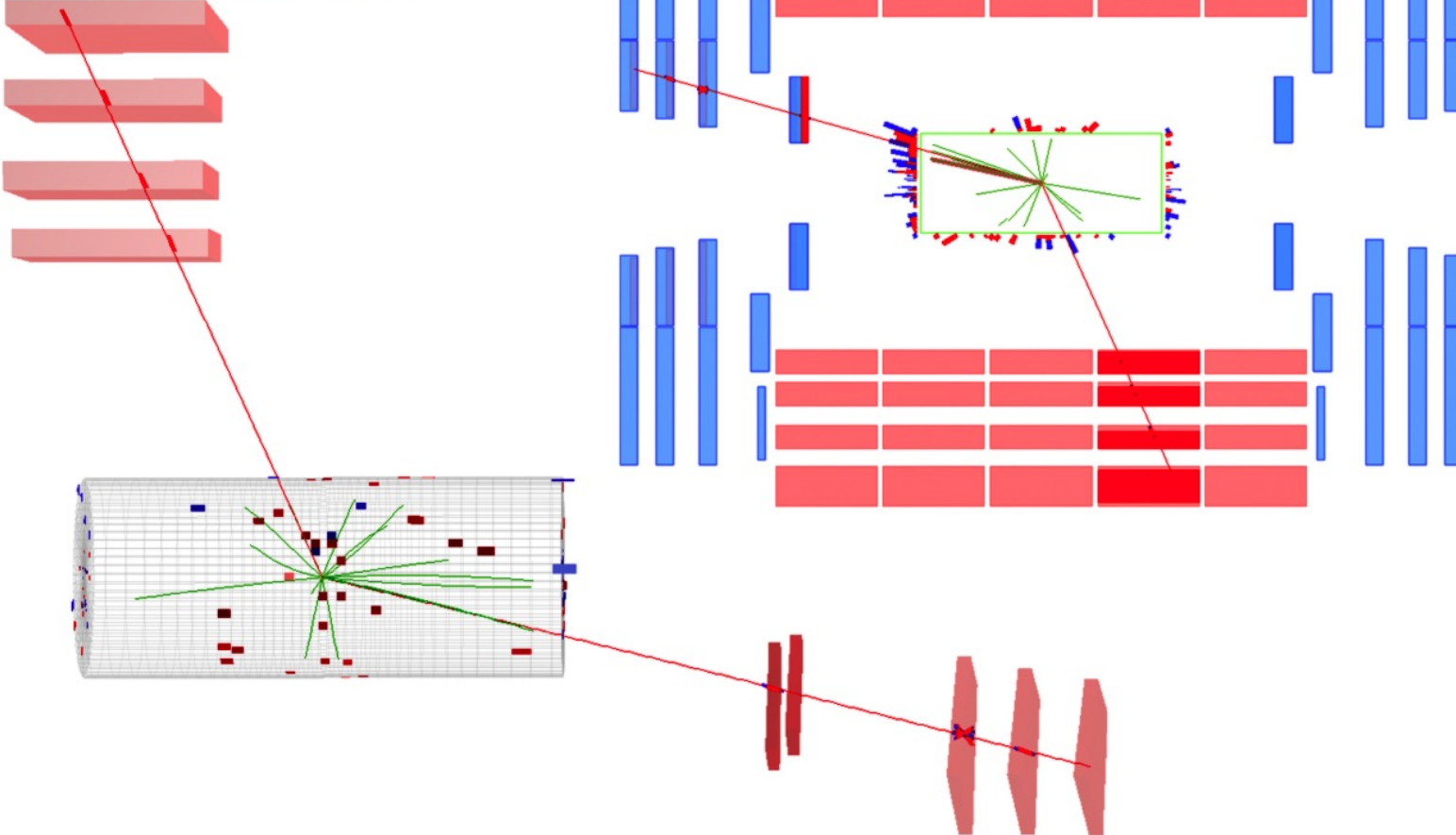




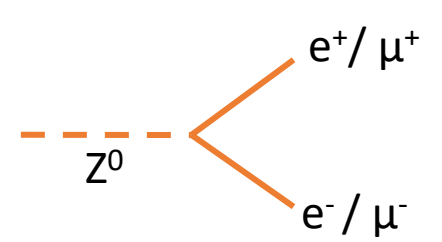
CMS Experiment at LHC, CERN
Run 136087 Event 39967482
Lumi section: 314
Mon May 24 2010, 15:31:58 CEST

MEASURING THE Z^0 CROSS-SECTION AT LHC

Muon $p_T = 27.3, 20.5 \text{ GeV}/c$
Inv. mass = $85.5 \text{ GeV}/c^2$



RECONSTRUCTING Z^0 'S



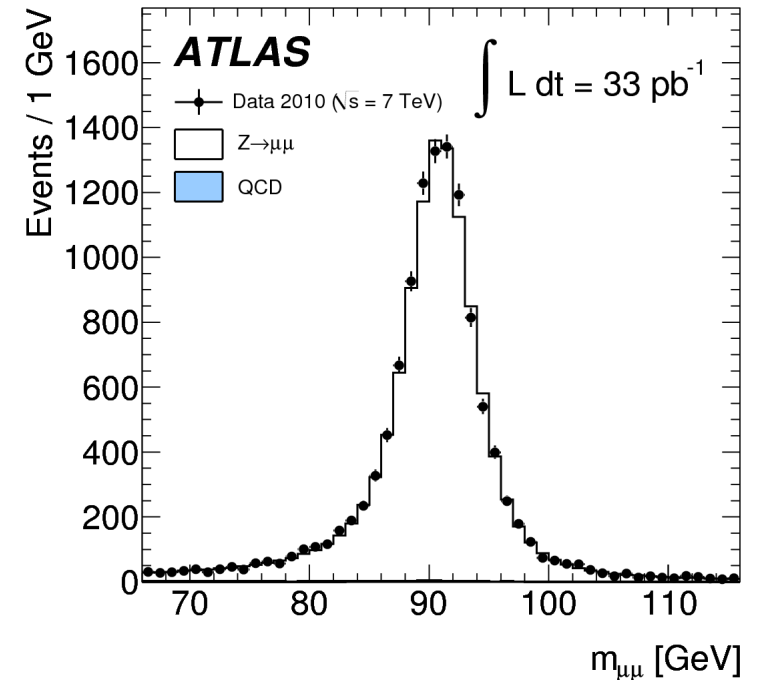
STEP-1: IDENTIFY THE OBSERVABLE OF INTEREST

- Identify Z decays using the invariant mass of the 2 leptons
 $M^2 = (L_1 + L_2)^2$ where $L_i = (E_i, \mathbf{p}_i)$ = 4-vector for lepton i
- Under assumption that lepton is massless compared to mass of Z^0
 $\Rightarrow M^2 = 2 E_1 E_2 (1 - \cos\theta_{12})$ where θ_{12} = angle between the leptons

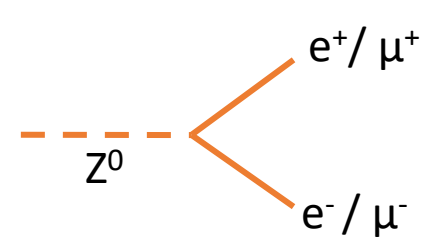
STEP-2: SELECT Z^0 EVENTS WITH 'ANALYSIS CUTS':

- Events with 2 high momentum electrons or muons
- Require the electrons or muons are of opposite charge
- With di-lepton mass close to the Z^0 mass
(e.g. $70 < m_{l+l^-} < 110$ GeV)

Very little background in Z^0 mass region!



RECONSTRUCTING Z^0 'S



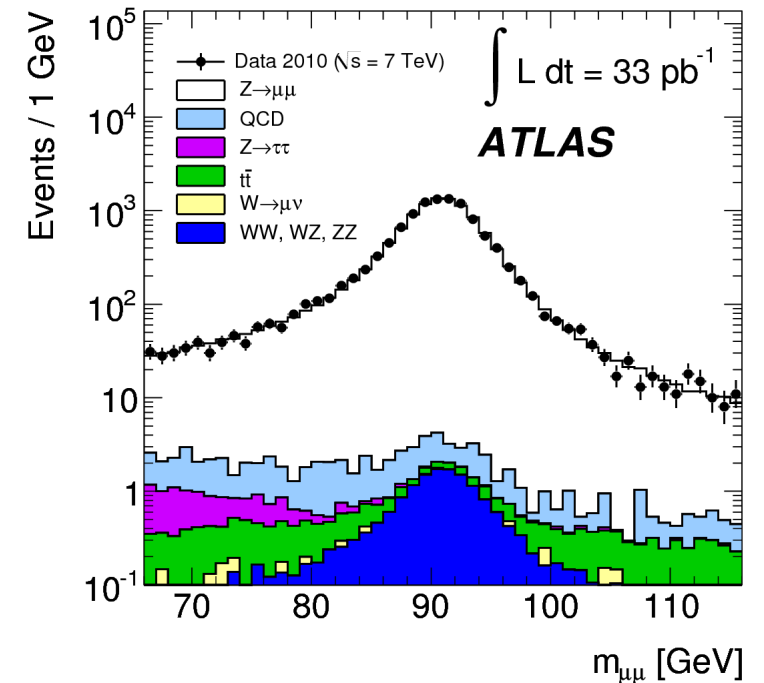
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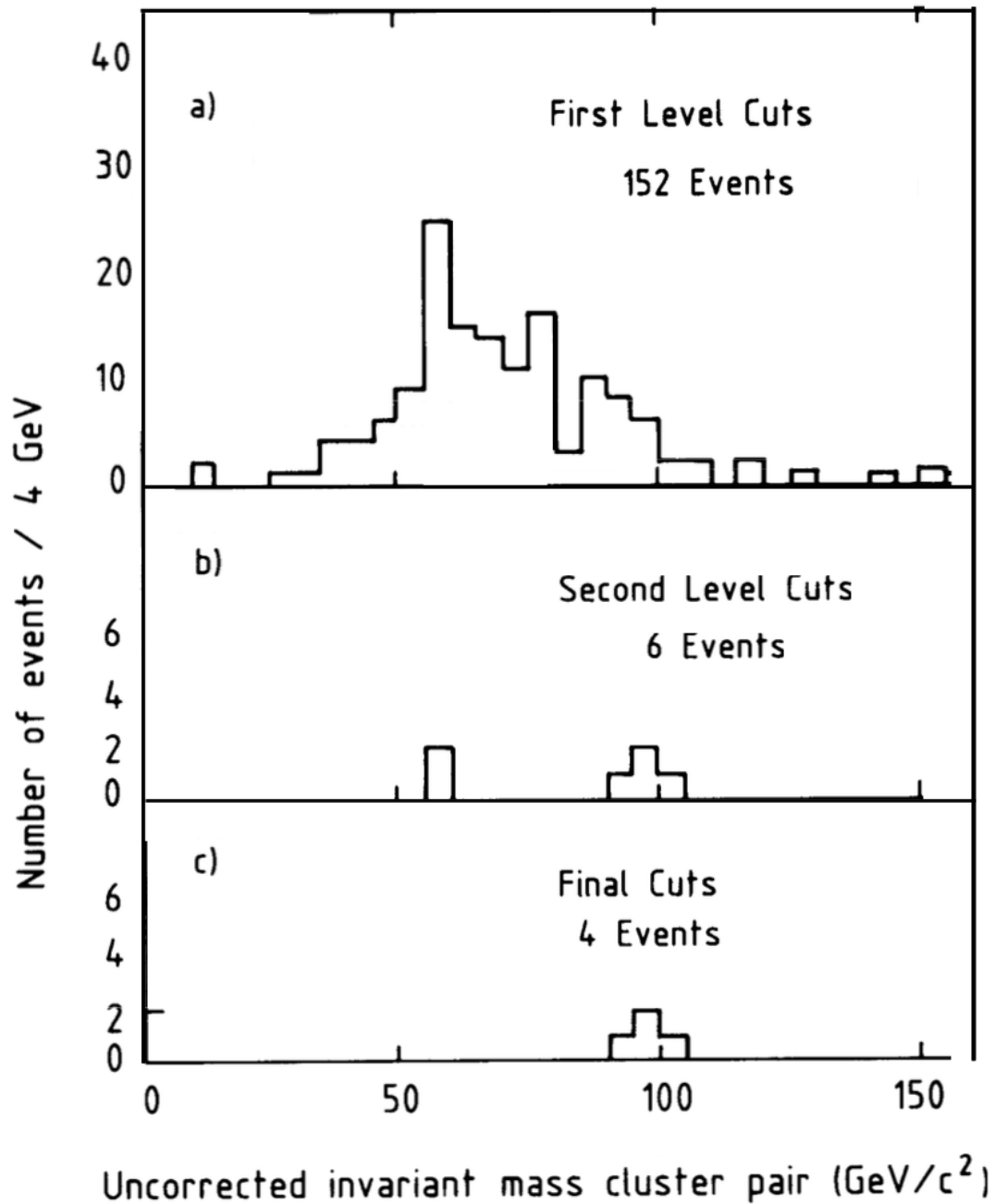
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A step back in time



Z->ee in UA1

Two EM clusters with E_T>25GeV.

As above plus a track with p_T>7GeV pointing to the cluster.
Hadronic and track isolation requirements applied.

A second cluster has also an isolated track.

MEASURING Z^0 CROSS-SECTION

THEORETICALLY

Cross-section calculated for:

- © Specific production mechanism (pp, pp, e^+e^-)
- © Centre-of-Mass of the collisions (7, 8, 13 TeV at LHC)

EXPERIMENTALLY

$$\sigma \cdot \text{BR} = \frac{\text{Number of events}}{\alpha \cdot \epsilon \cdot L}$$

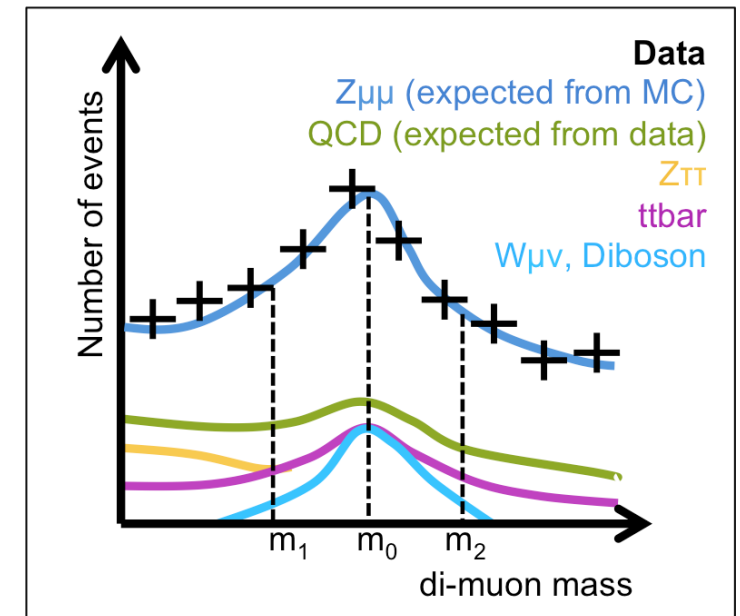
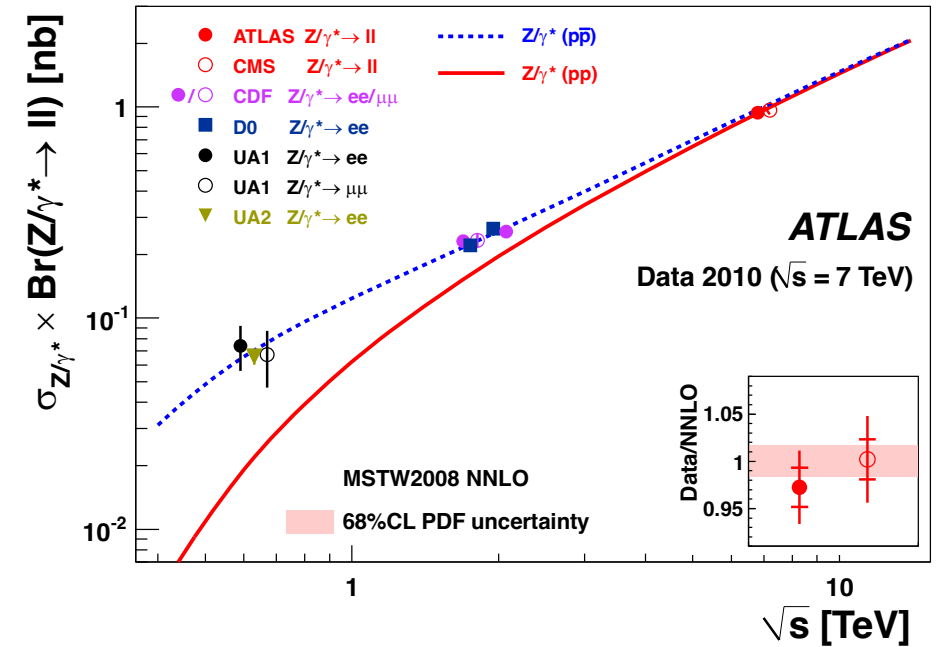
N of events: N of events on data – N of expected background events

α – acceptance: fraction of events passing selection requirements

ϵ – efficiency: reconstruction efficiency of relevant objects

L – luminosity

All numbers carry **uncertainties** –
both “**statistical**” and “**systematic**”!



MEASURING Z⁰ CROSS-SECTION

THEORETICALLY

Cross-section calculated for:

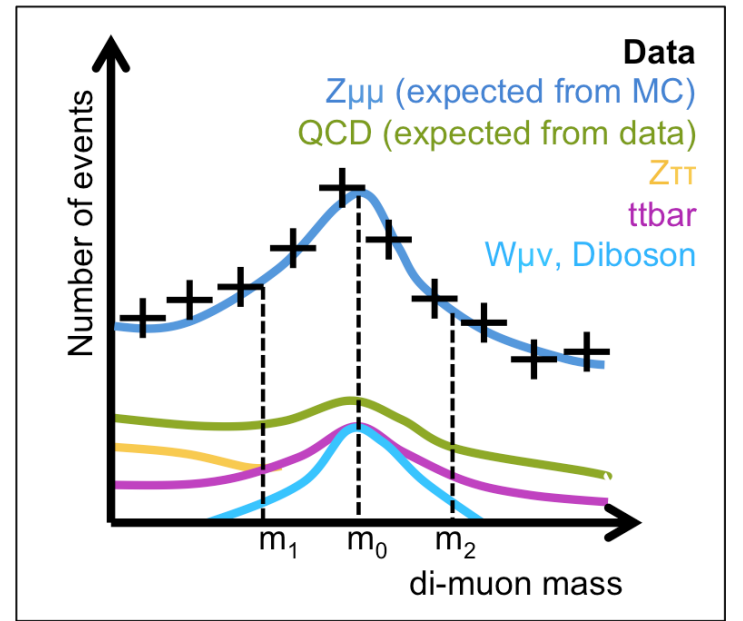
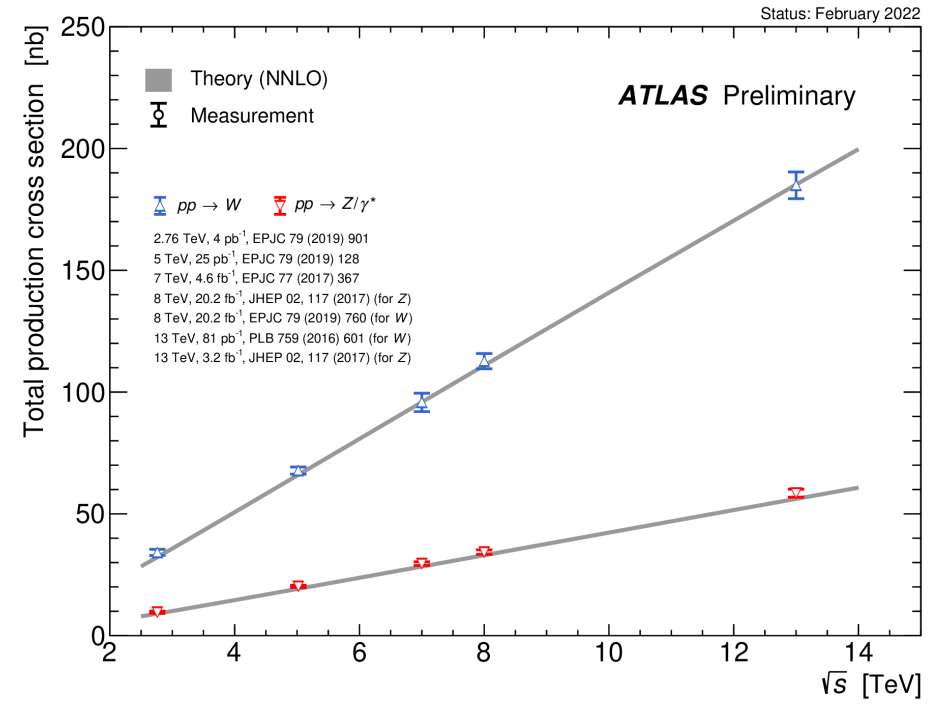
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EXPERIMENTALLY

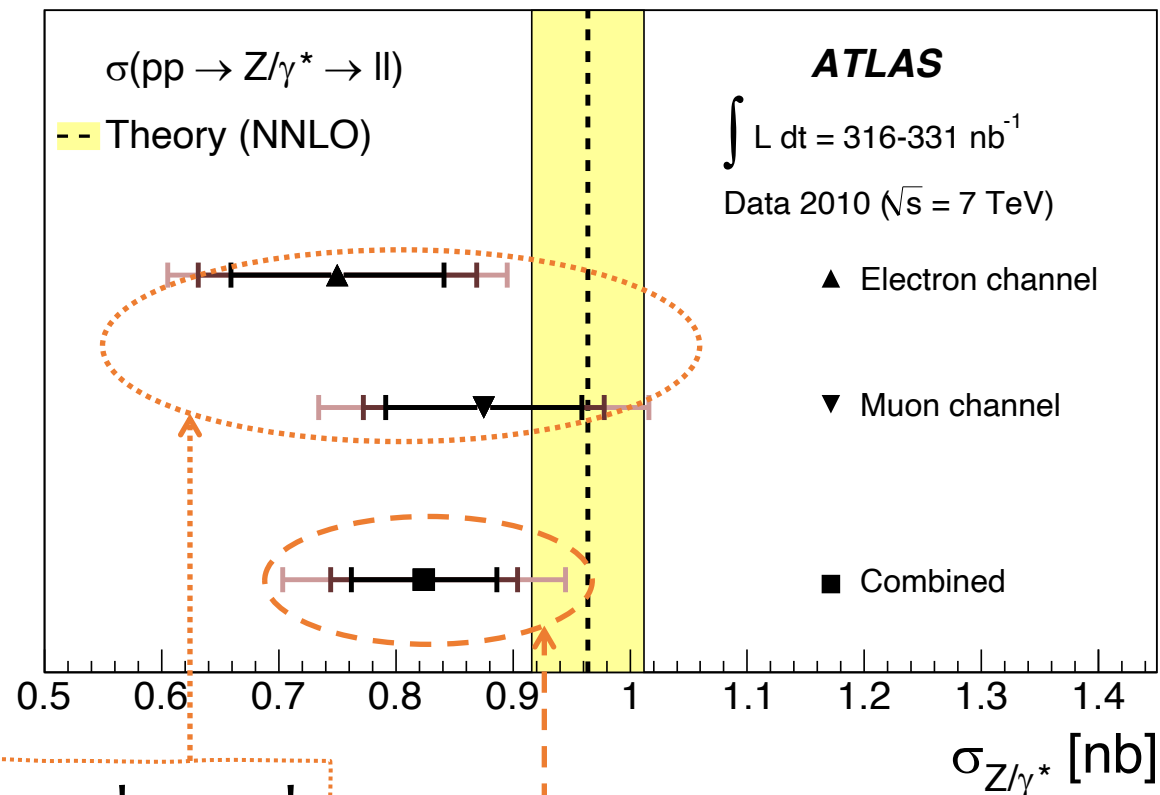
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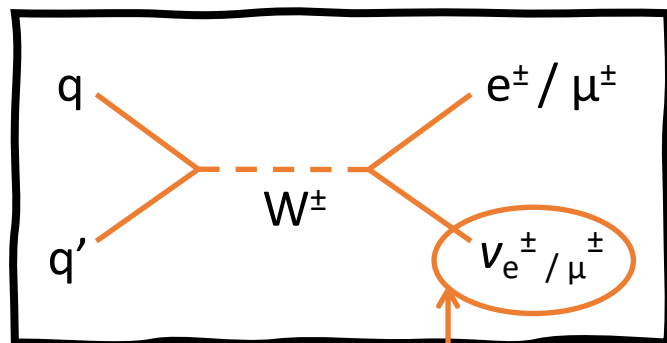
MEASURING Z^0 CROSS-SECTION



Electron and Muon channel agree within uncertainties

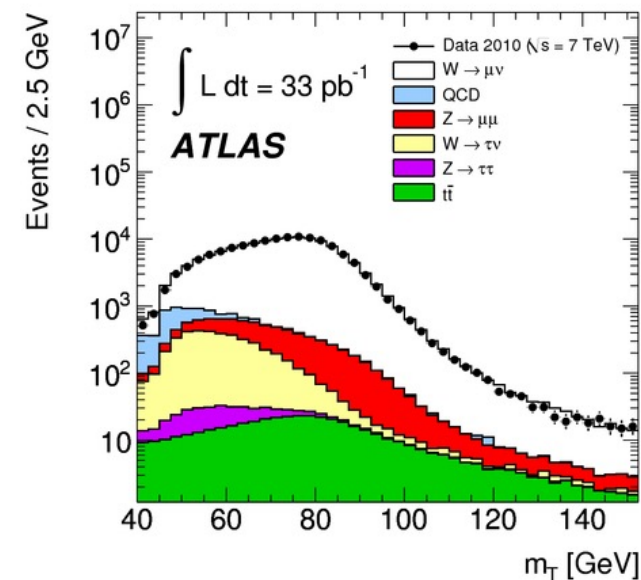
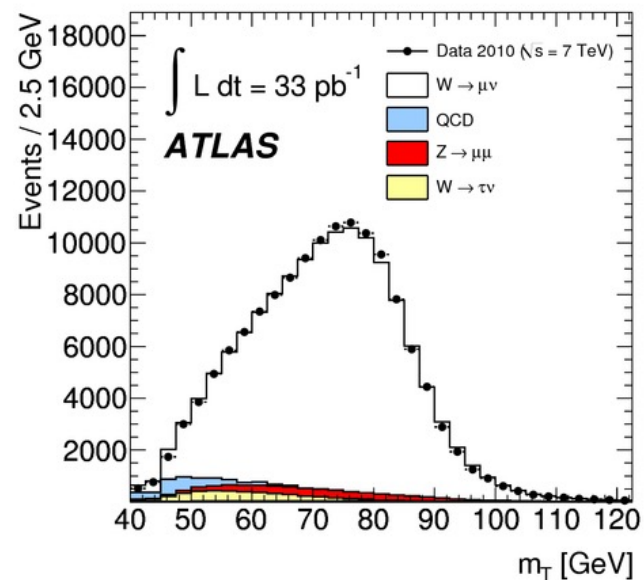
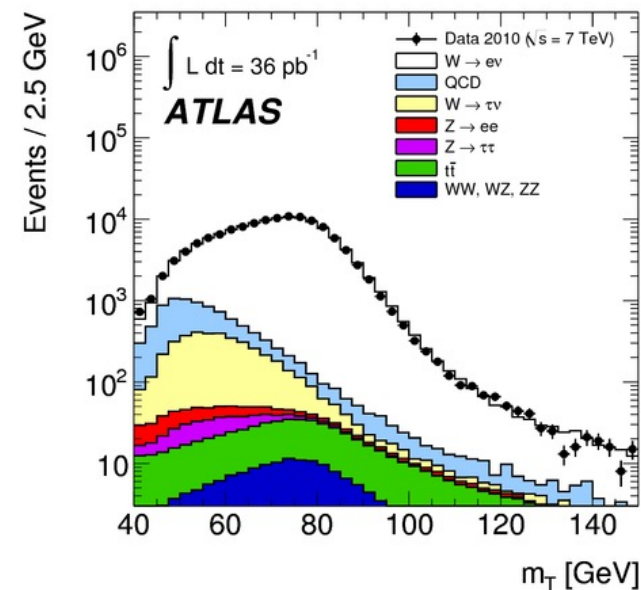
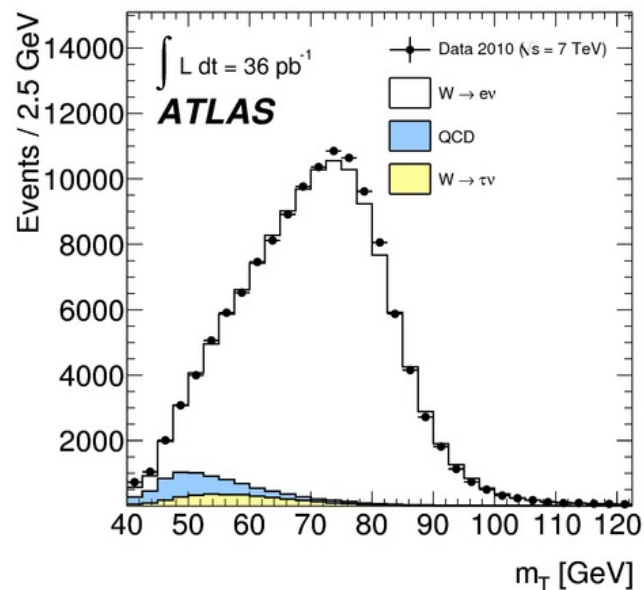
Measurement consistent with prediction within uncertainties

MEASURING W CROSS-SECTION

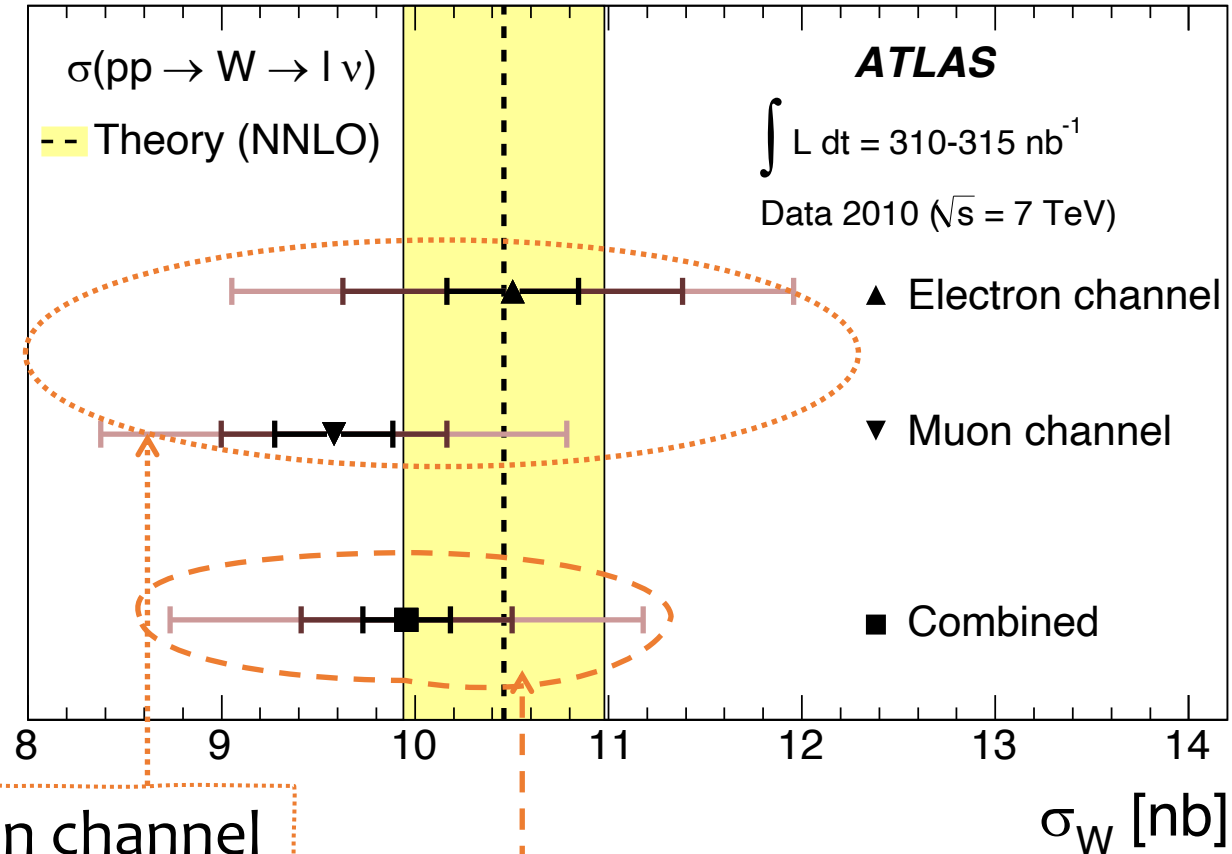


Available in the
transverse plane only!

$$M_T^2 = 2 E_{T1} E_{T2} (1 - \cos\theta_{12})$$



MEASURING W CROSS-SECTION

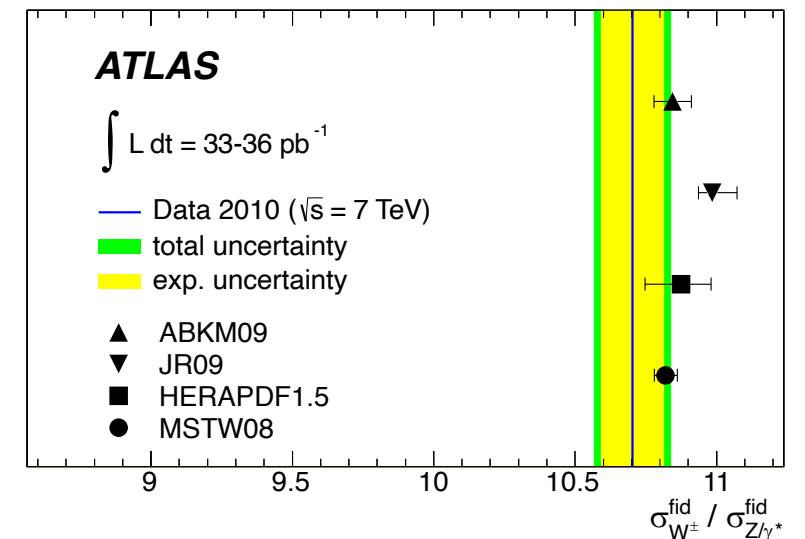
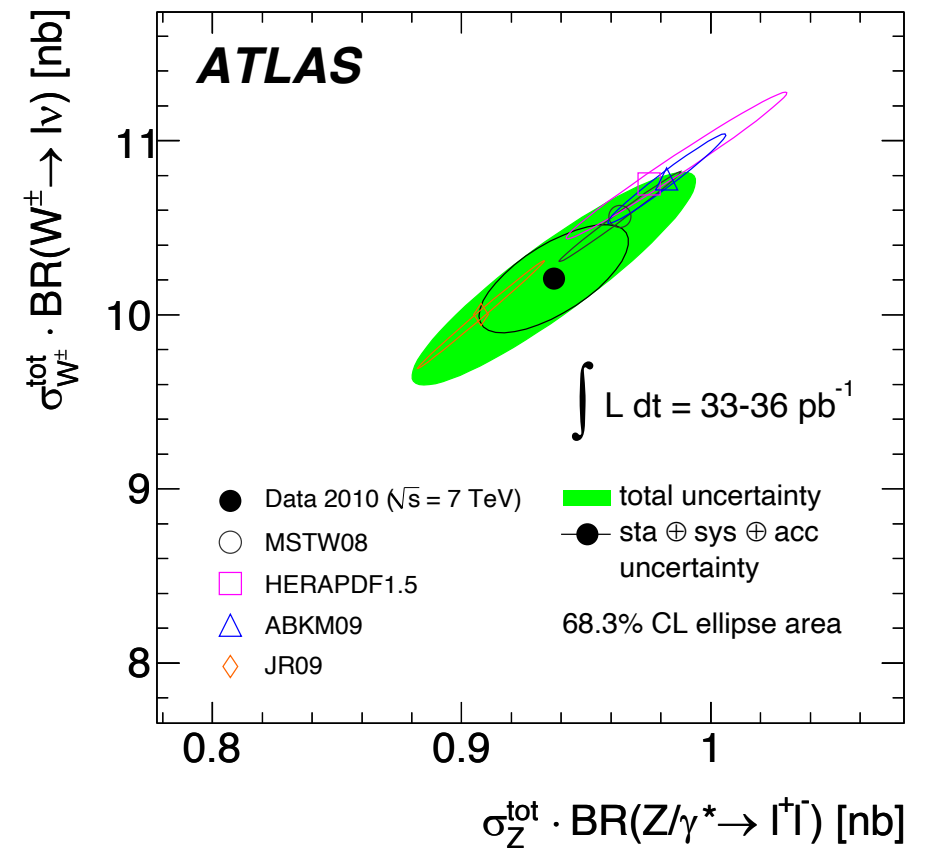
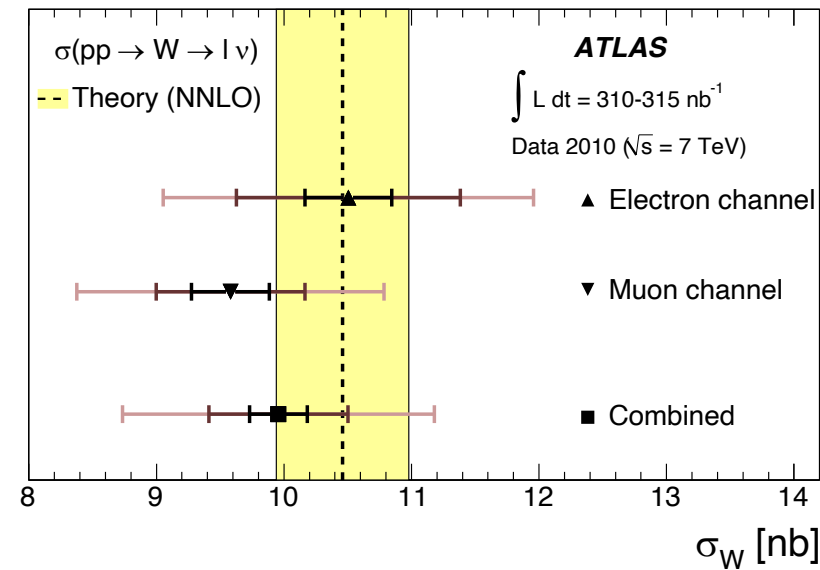
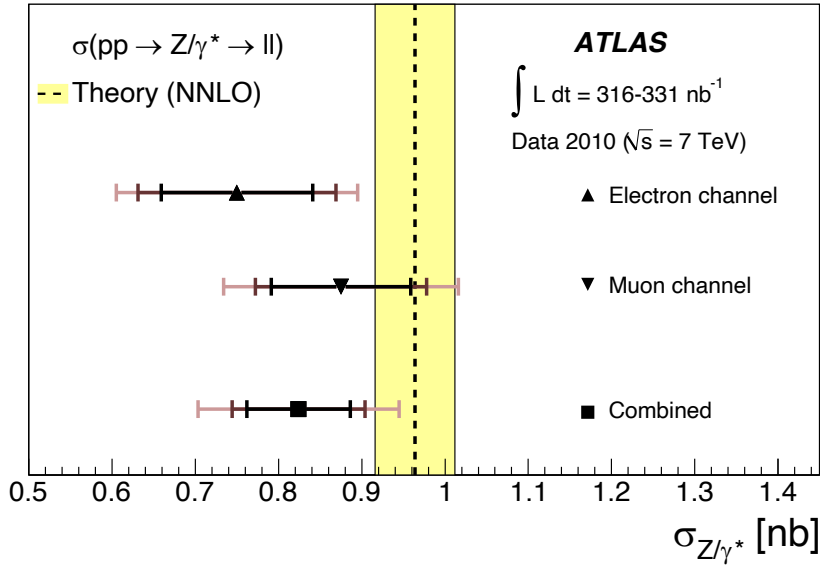


Electron and Muon channel agree within uncertainties

Measurement consistent with prediction within uncertainties

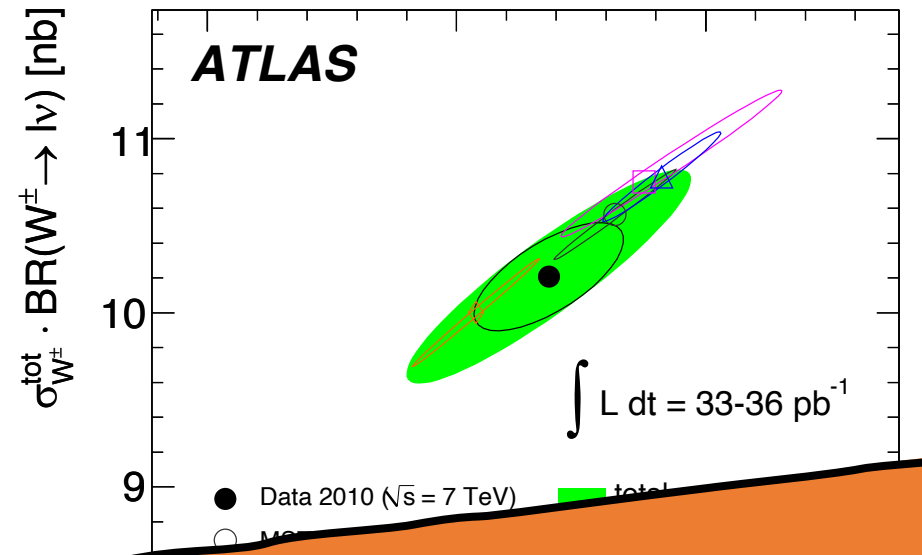
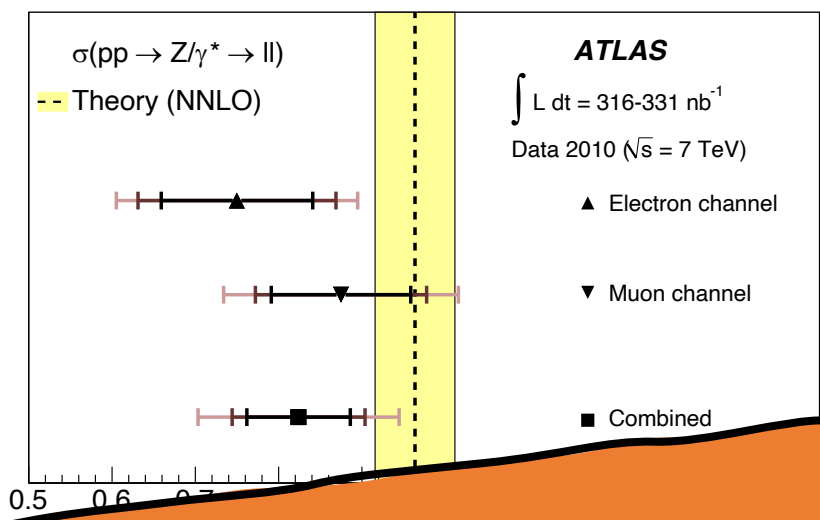
MEASURING CROSS-SECTIONS

RATIOS

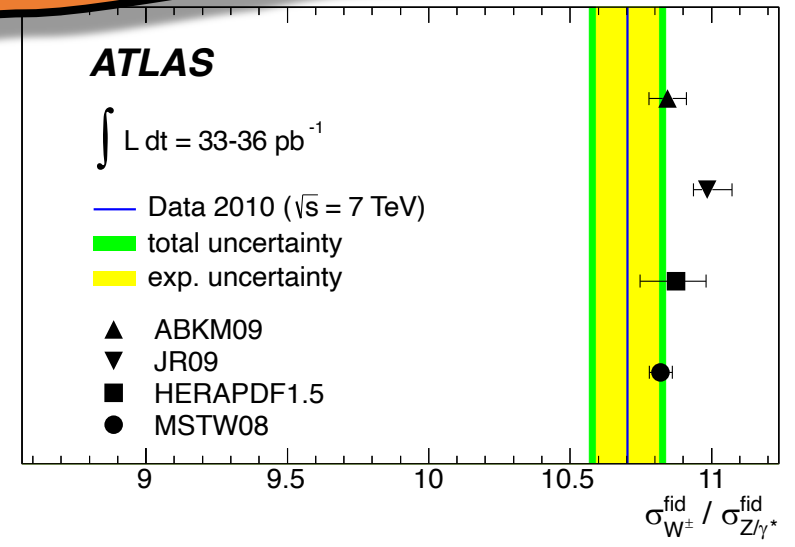
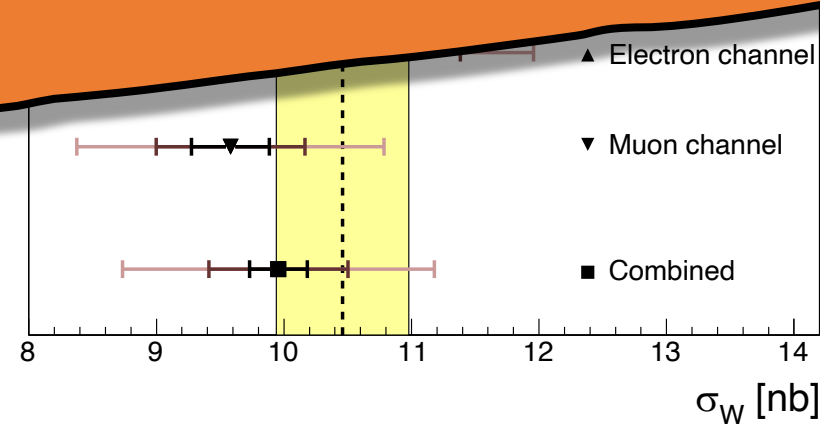


MEASURING CROSS-SECTIONS

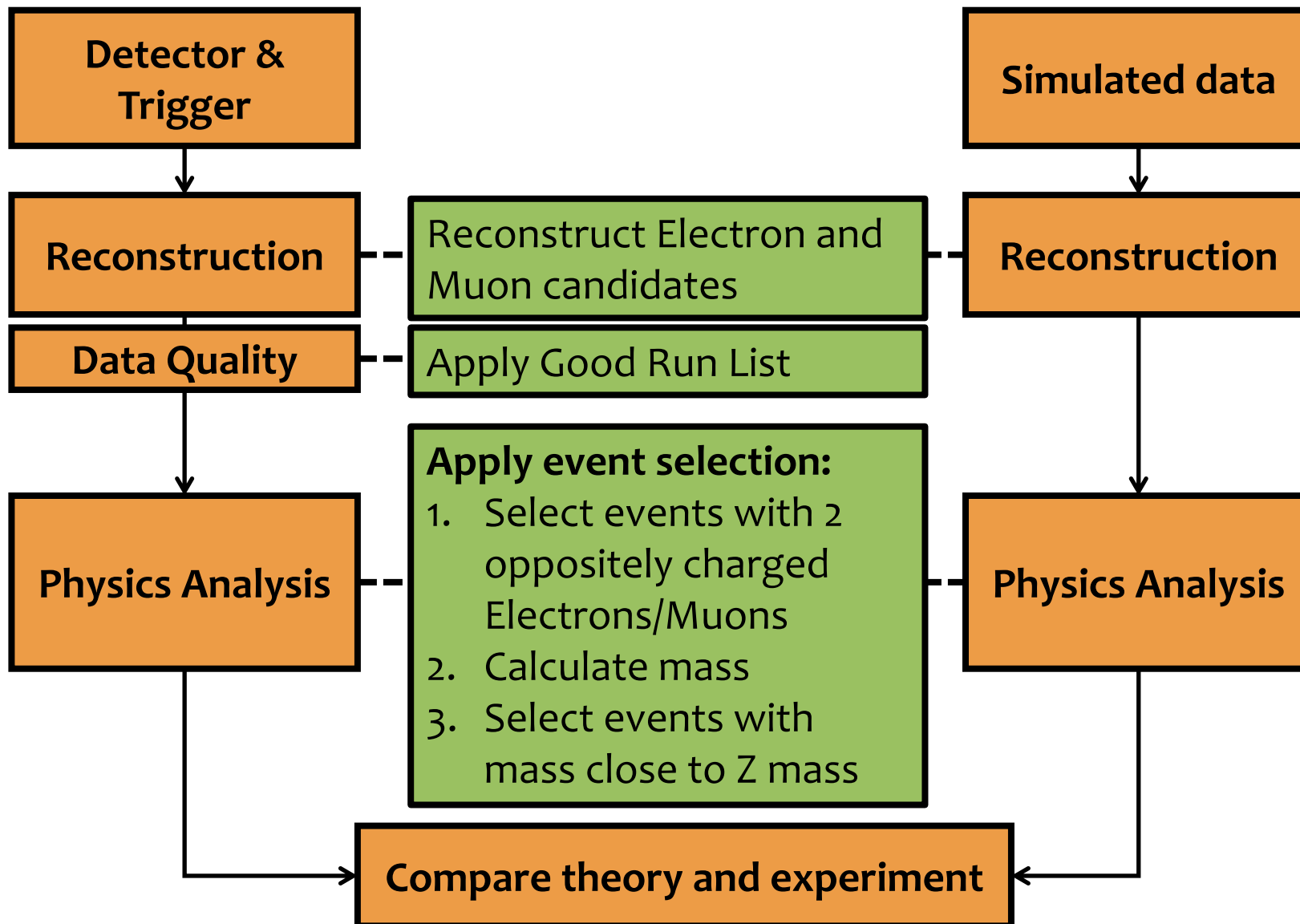
RATIOS



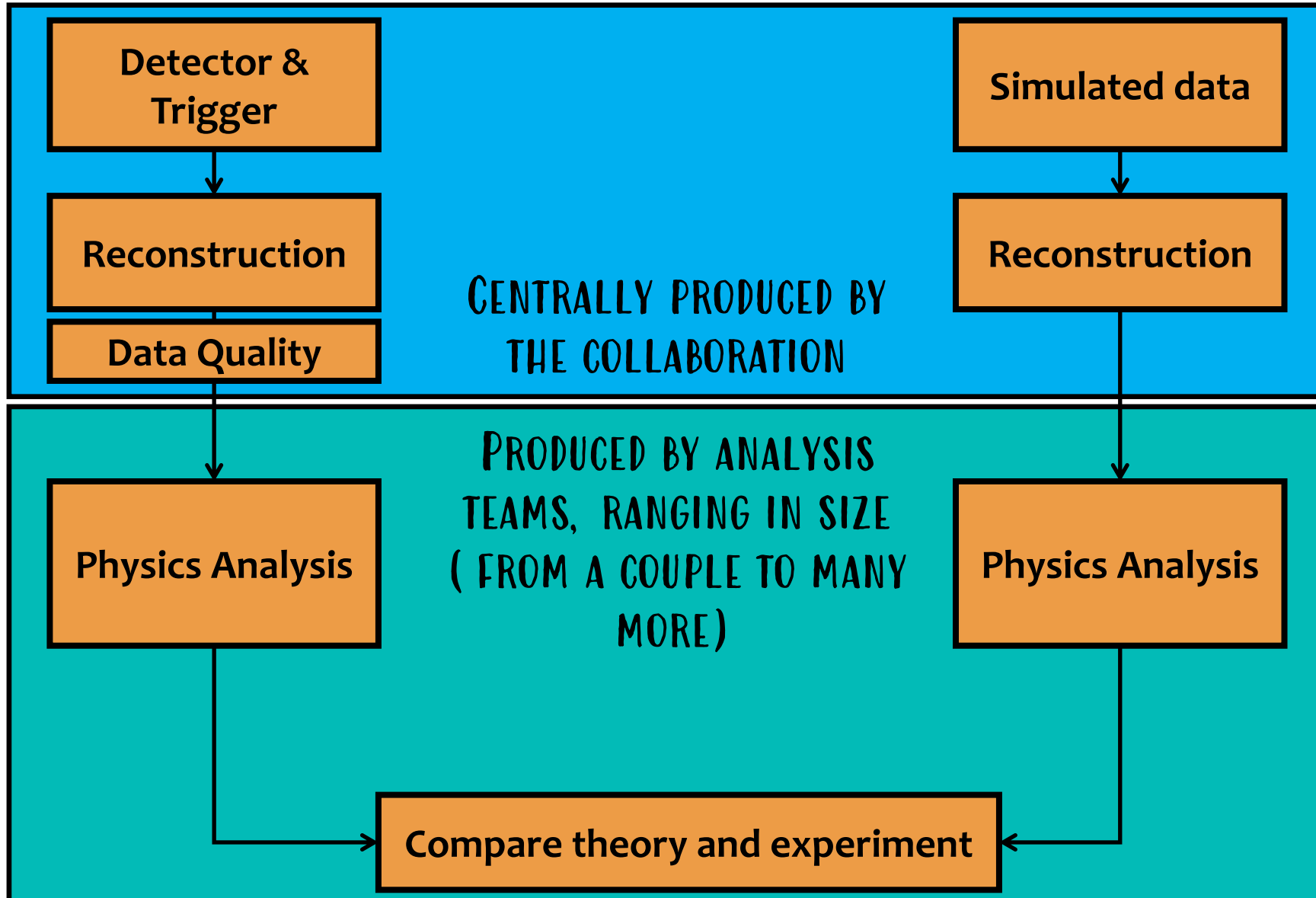
General aim:
 find ways to reduce the uncertainties!



ANALYSIS FLOW – E.G. CROSS-SECTION MEASUREMENT



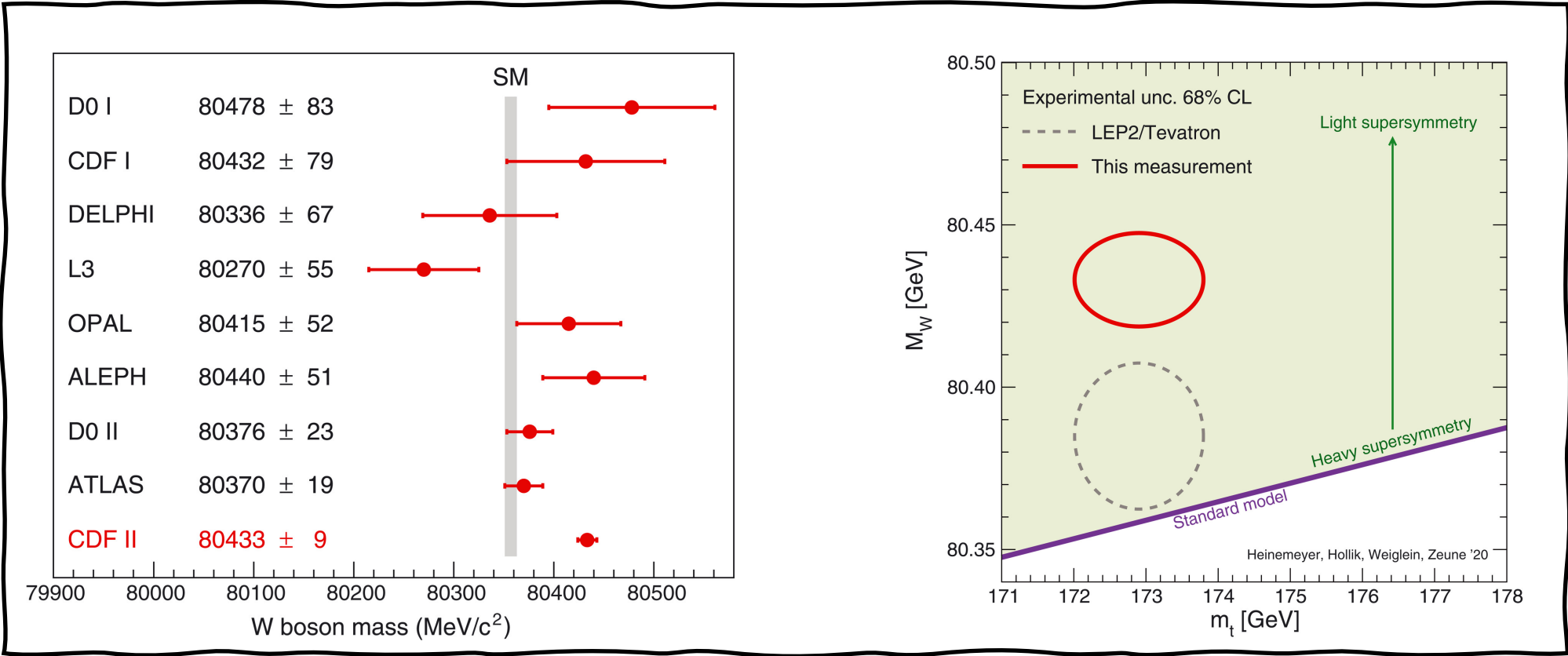
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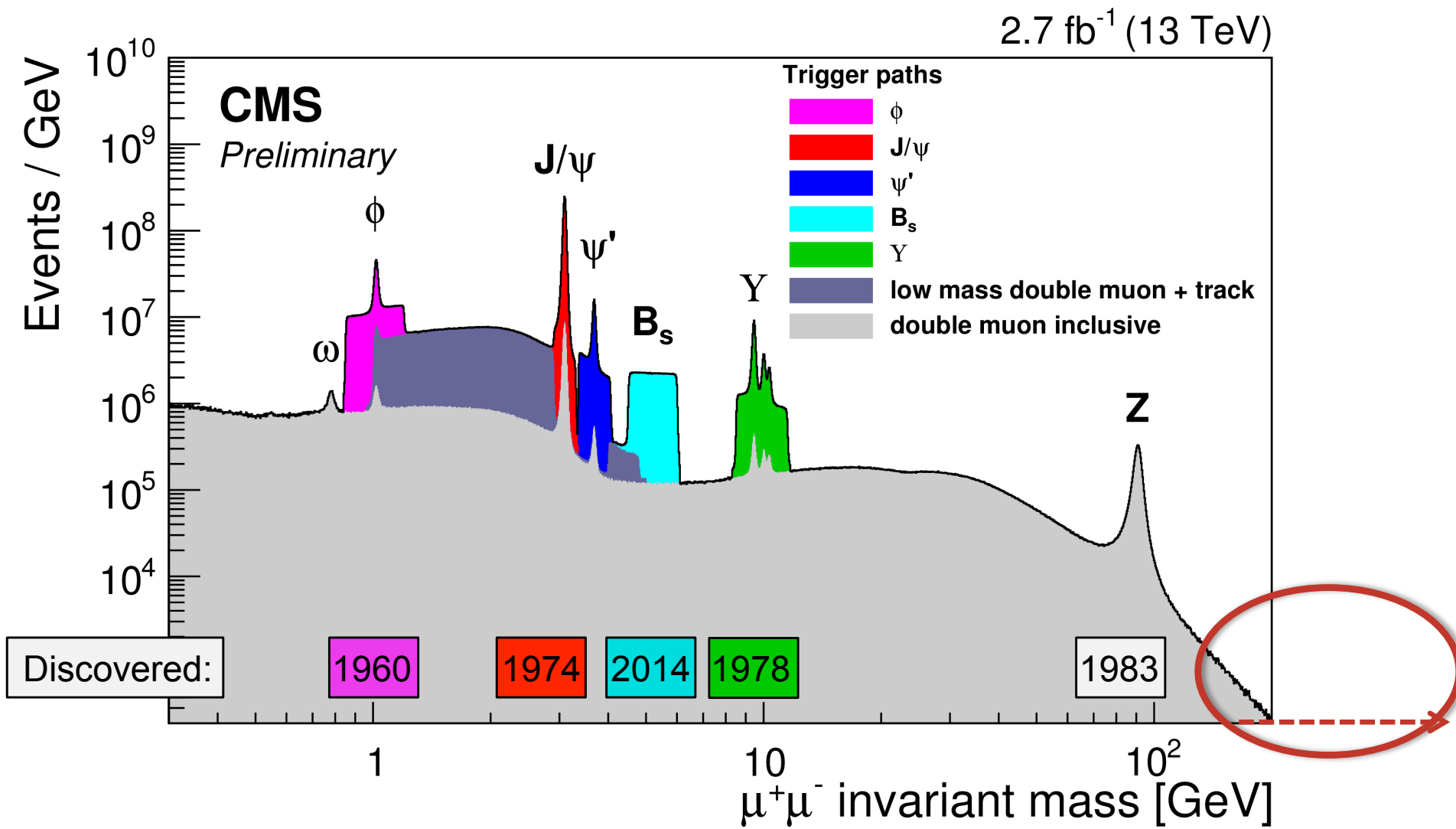


MASSES – RECENT NEWS! (NOT FROM THE LHC)

Following up from discussion in class...

“ This measurement, $M_W=80,433.5\pm 9.4$ MeV, is more precise than all previous measurements of M_W combined. A comparison with the SM expectation of $M_W=80,357\pm 6$ MeV [...] yields a difference with a significance of 7.0σ and suggests the possibility of improvements to the SM calculation or of extensions to the SM. ”



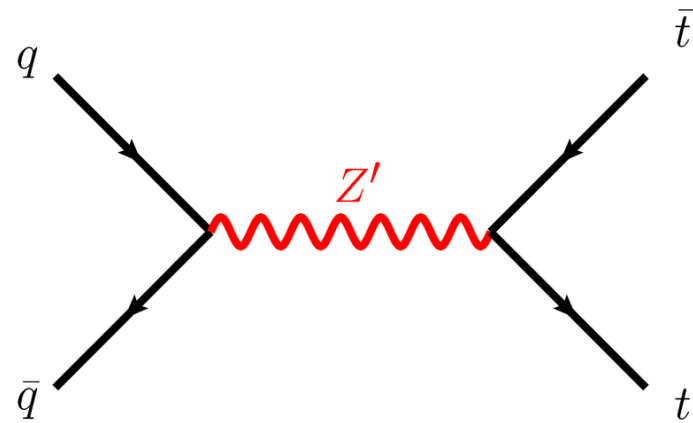


What's out there?

WHAT “NEW PHYSICS” CAN CREATE BUMPS?

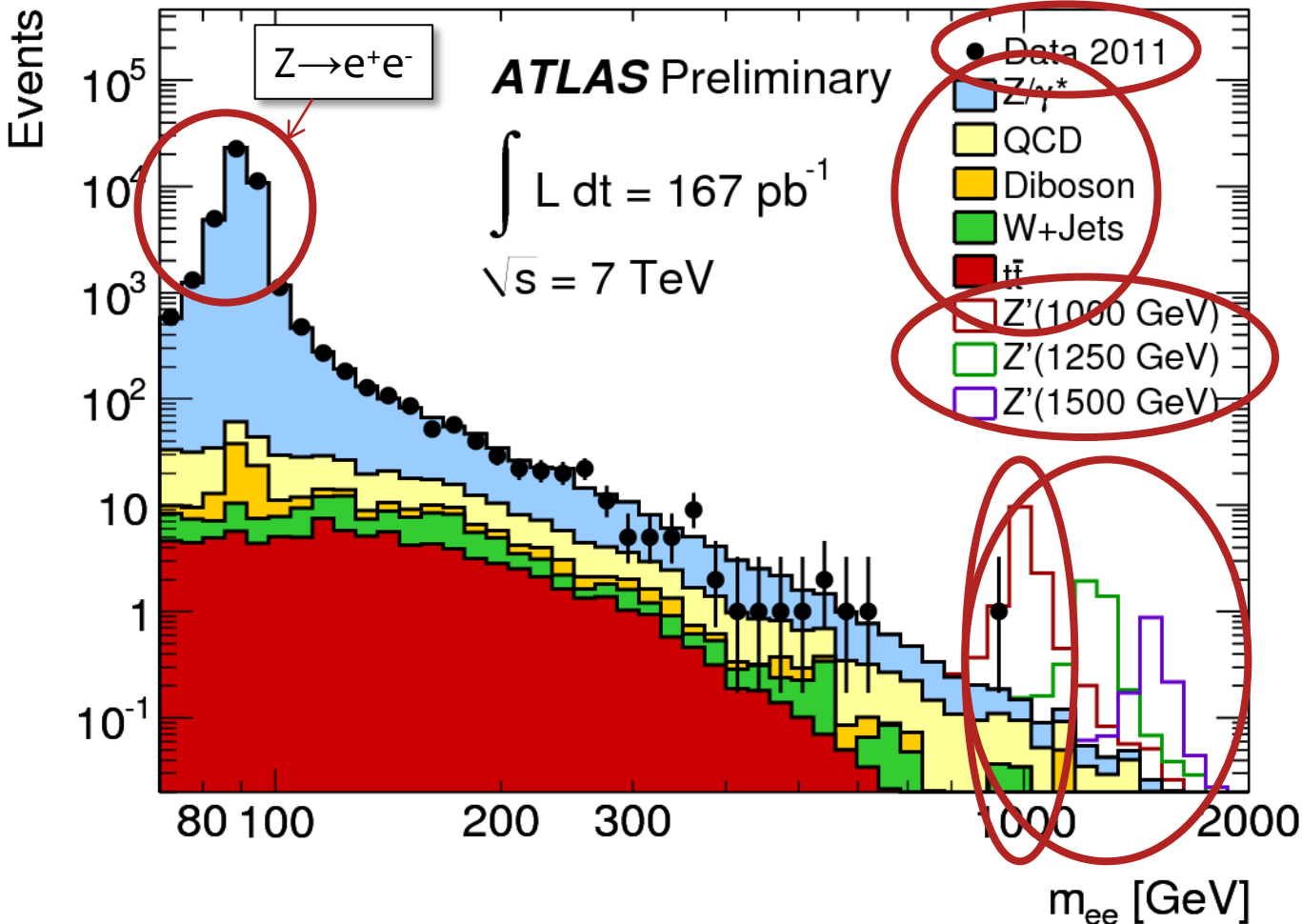
Many models giving answers to the SM problems include new heavy particles with short lifetime that appear as resonances

- Heavy bosons e.g. in grand unified and additional gauge symmetries
- High mass states, gravitons, e.g. in ‘randall-sundrum’ models
- Heavy quark partners, excited leptons, leptoquarks
- Composite Higgs
- Dark matter candidates



SIMPLE SEARCH EXAMPLE: SEARCH FOR A HEAVY Z'

© Like $Z \rightarrow ee$ but at higher mass



Select 2 electron candidates and plot their invariant mass for:

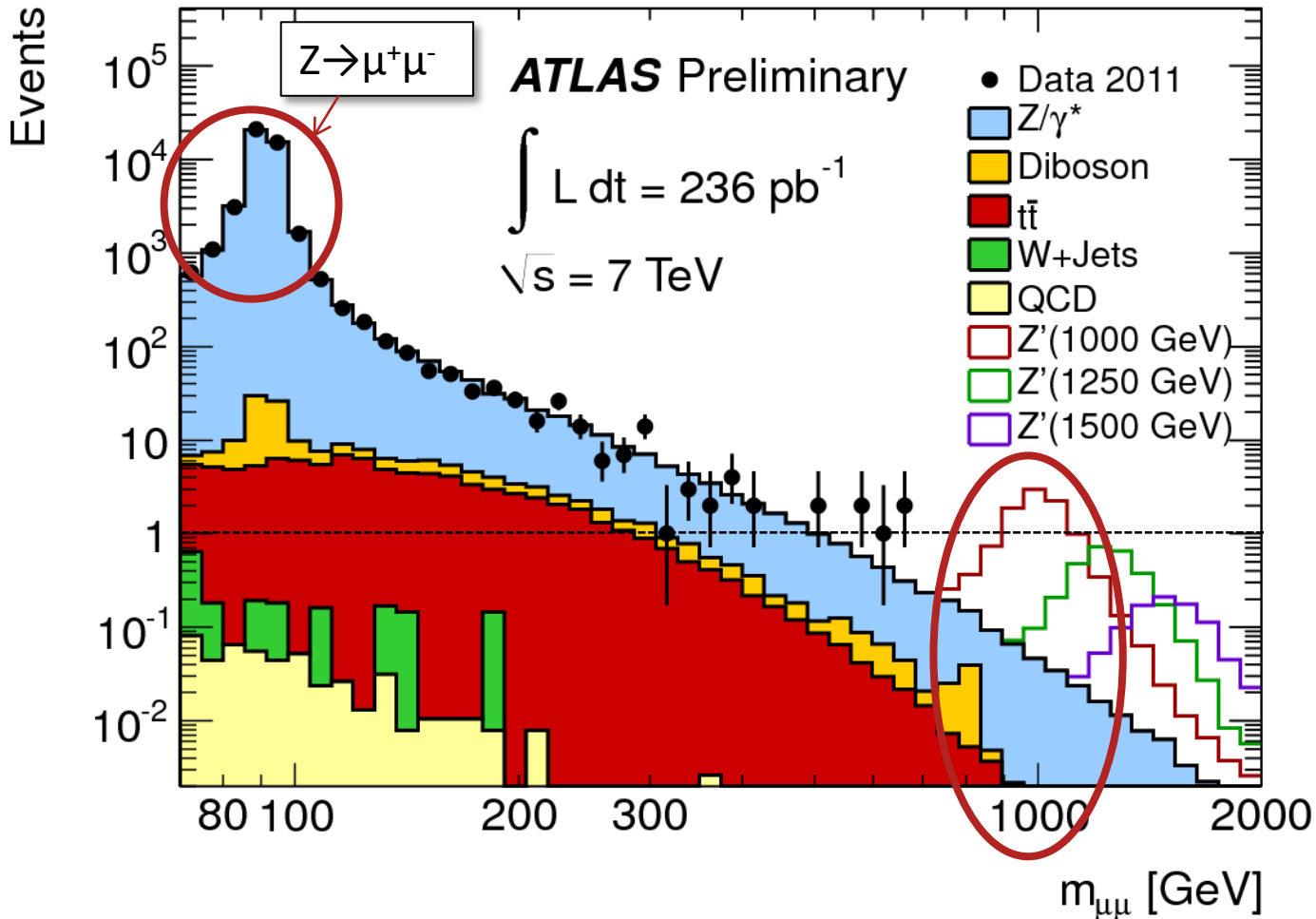
1. **Data**
2. **Simulated background events**
3. **Simulated signal with different masses**

Data inconsistent with a 1TeV Z'

Cross-section decreases with mass (higher the mass of the Z' , the more data needed to discover it)

SIMPLE SEARCH EXAMPLE: SEARCH FOR A HEAVY Z'

© And similar for muons



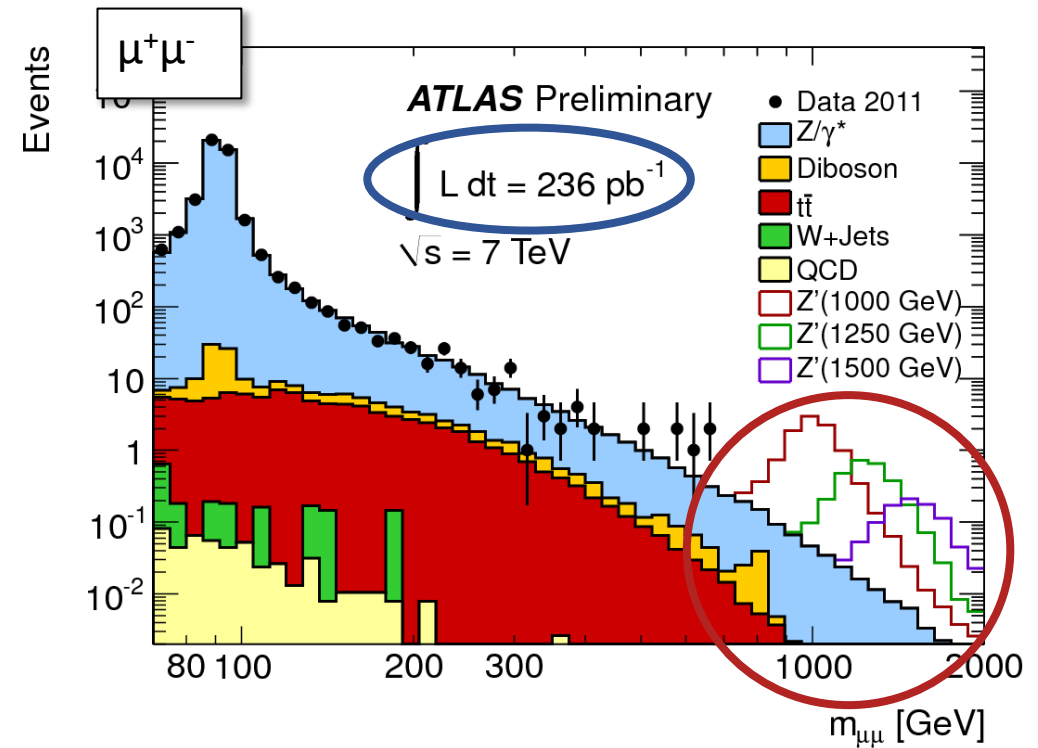
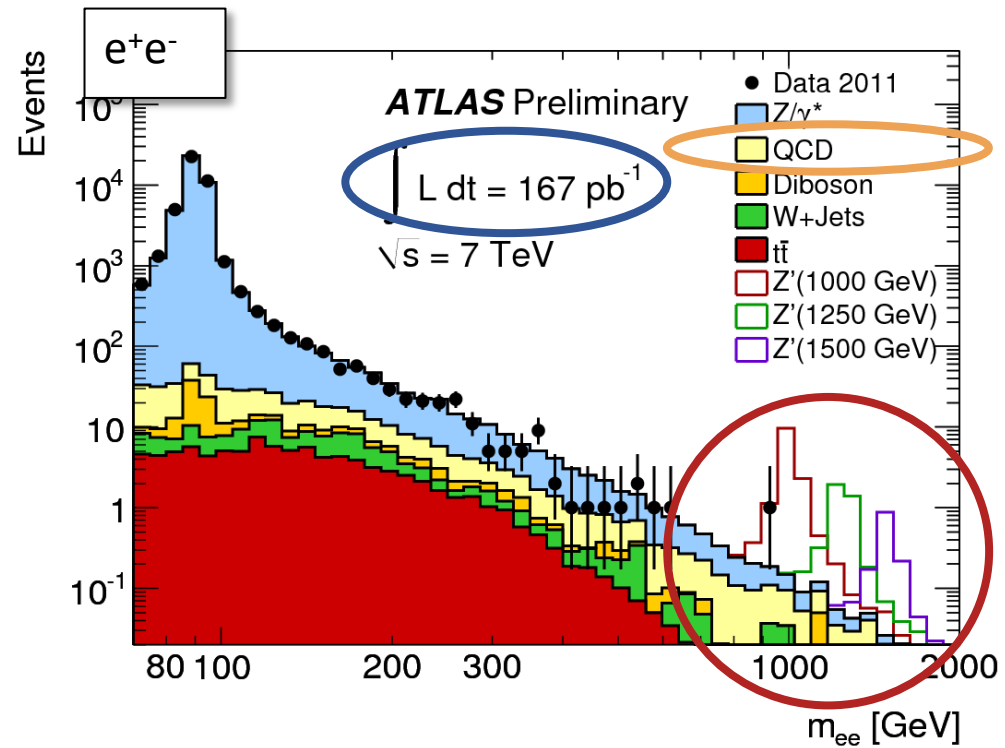
Select 2 electron candidates and plot their invariant mass for:

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3. **Simulated signal with different masses**

Data inconsistent with a 1TeV Z'

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A SMALL COMPARISON



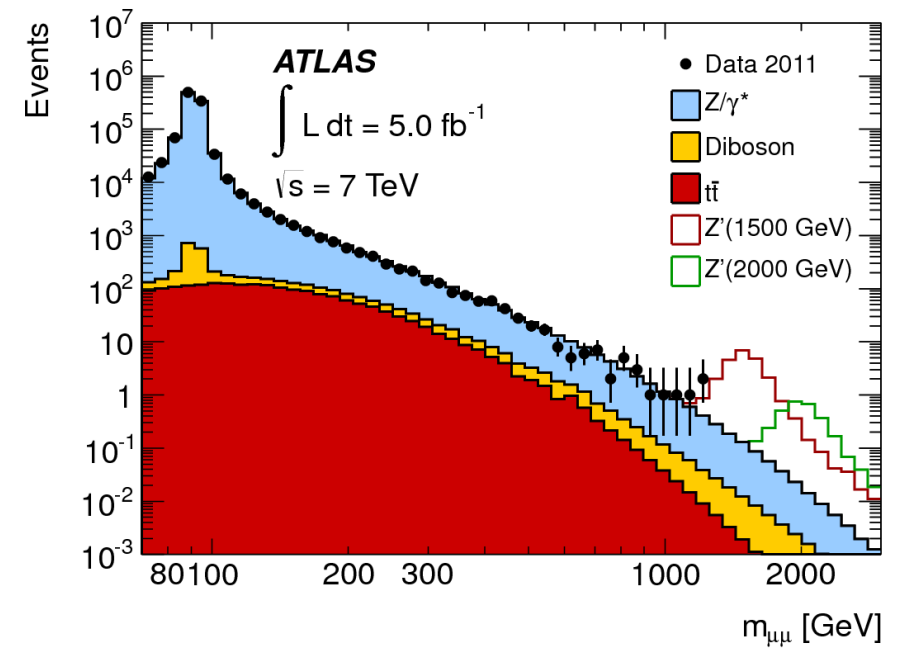
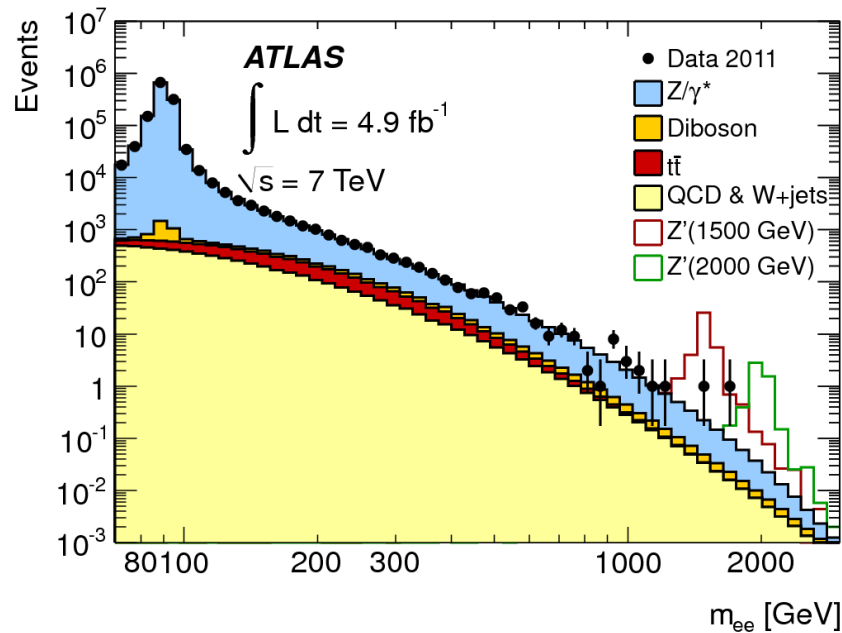
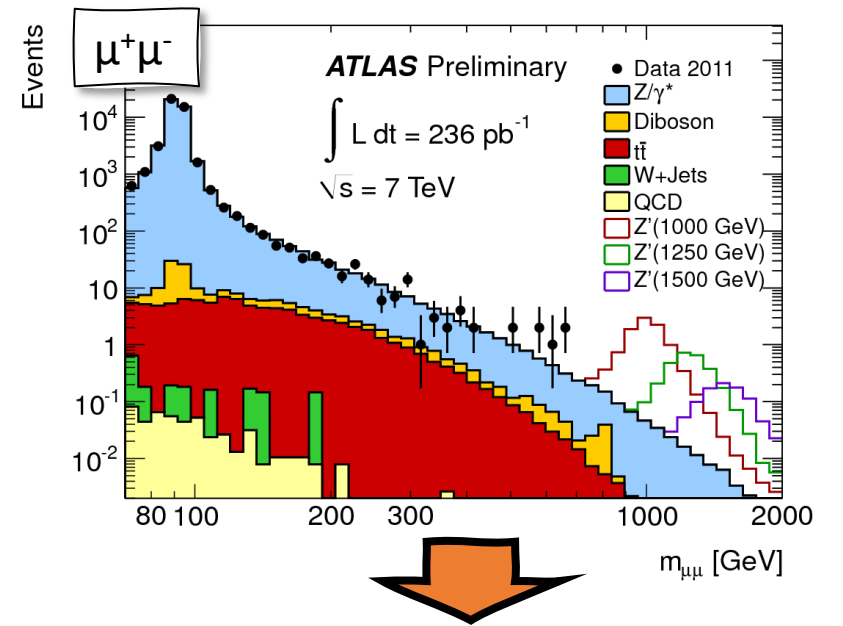
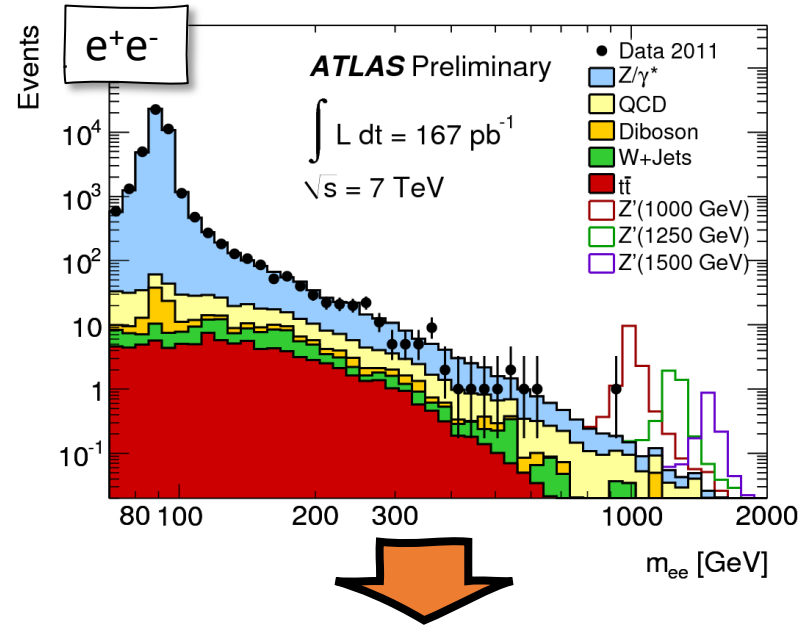
Differences in:

⊙ Resolution

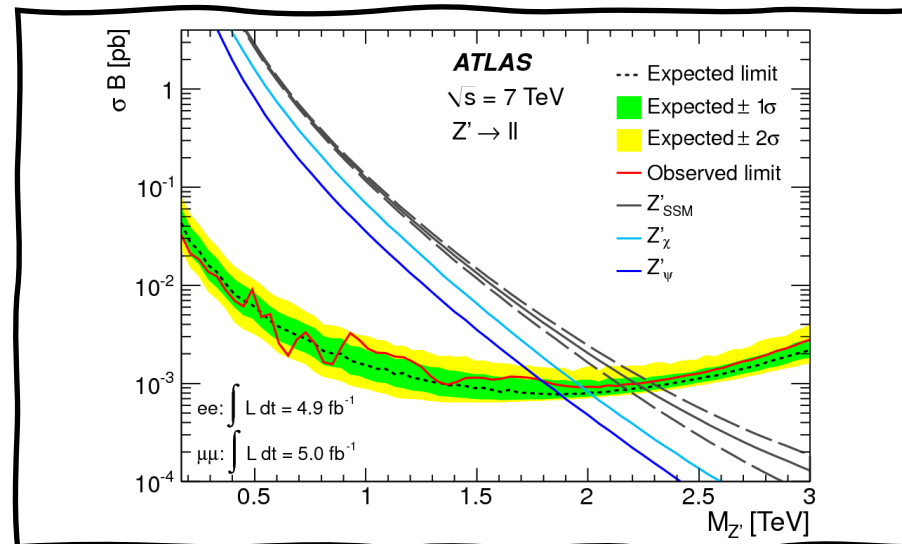
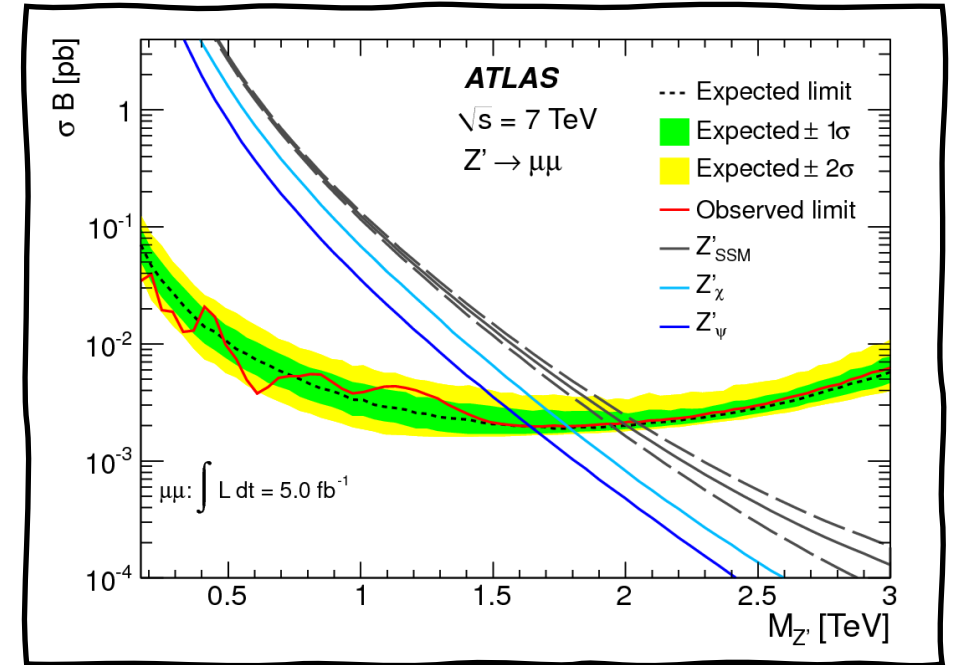
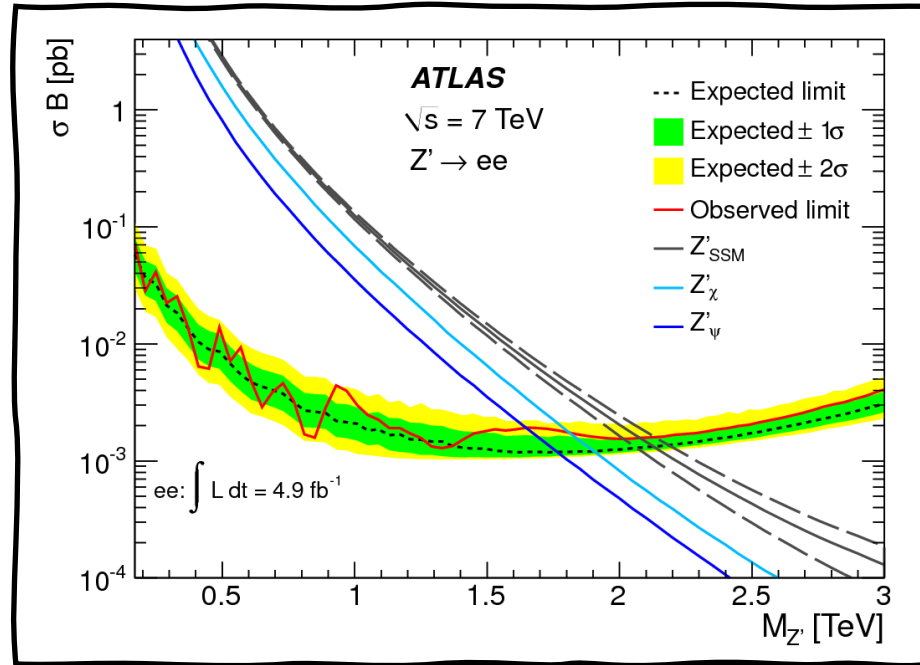
⊙ Background composition

⊙ Dataset

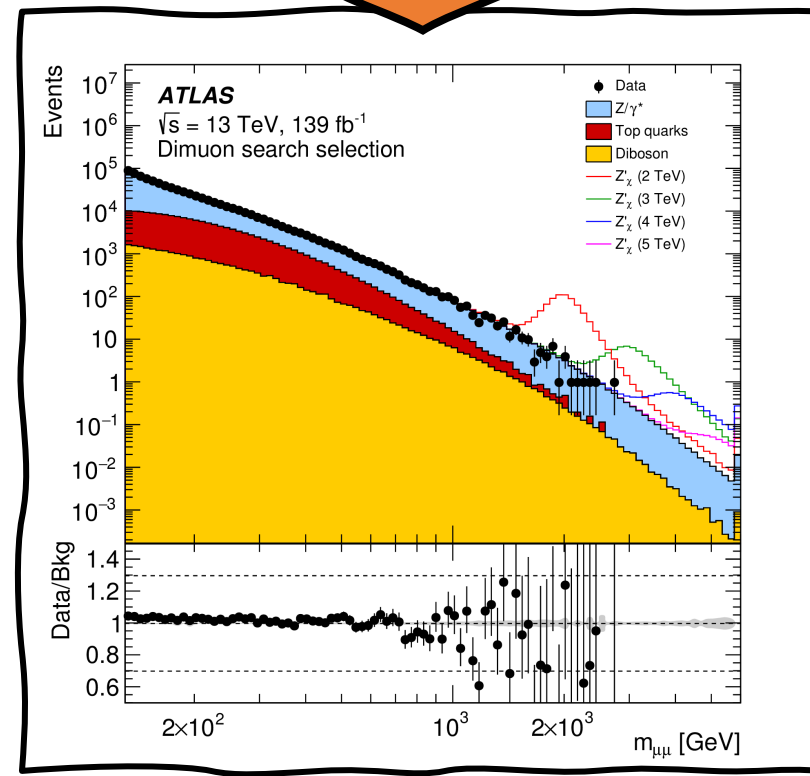
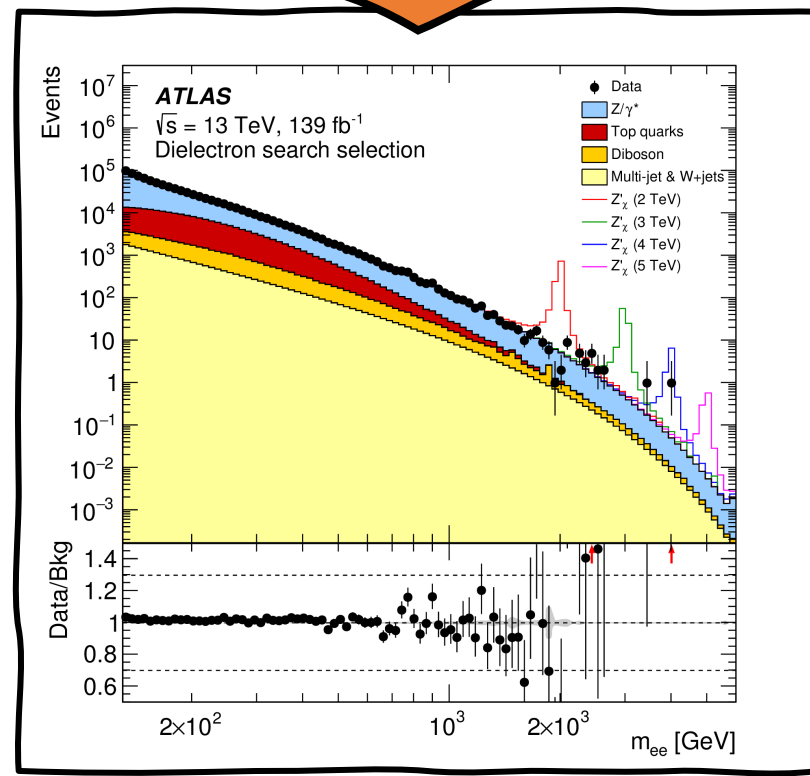
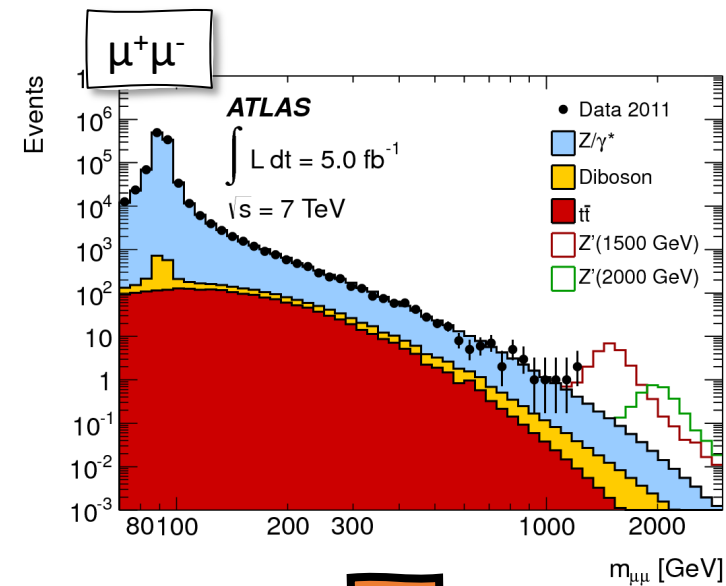
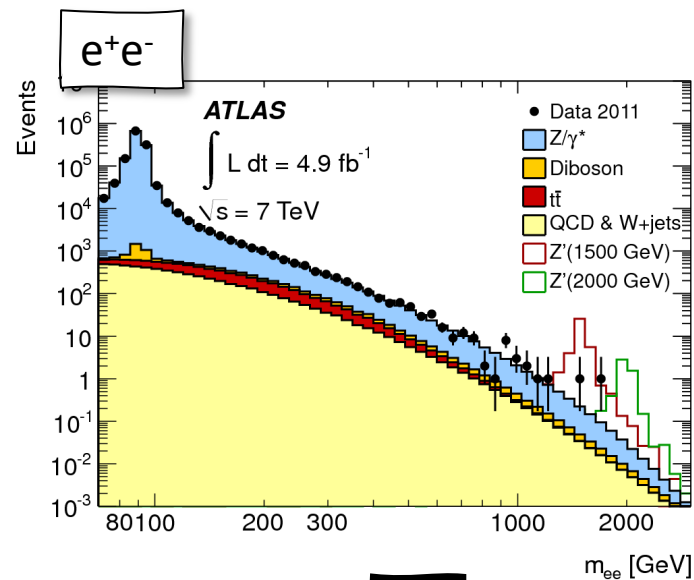
EVOLUTION...



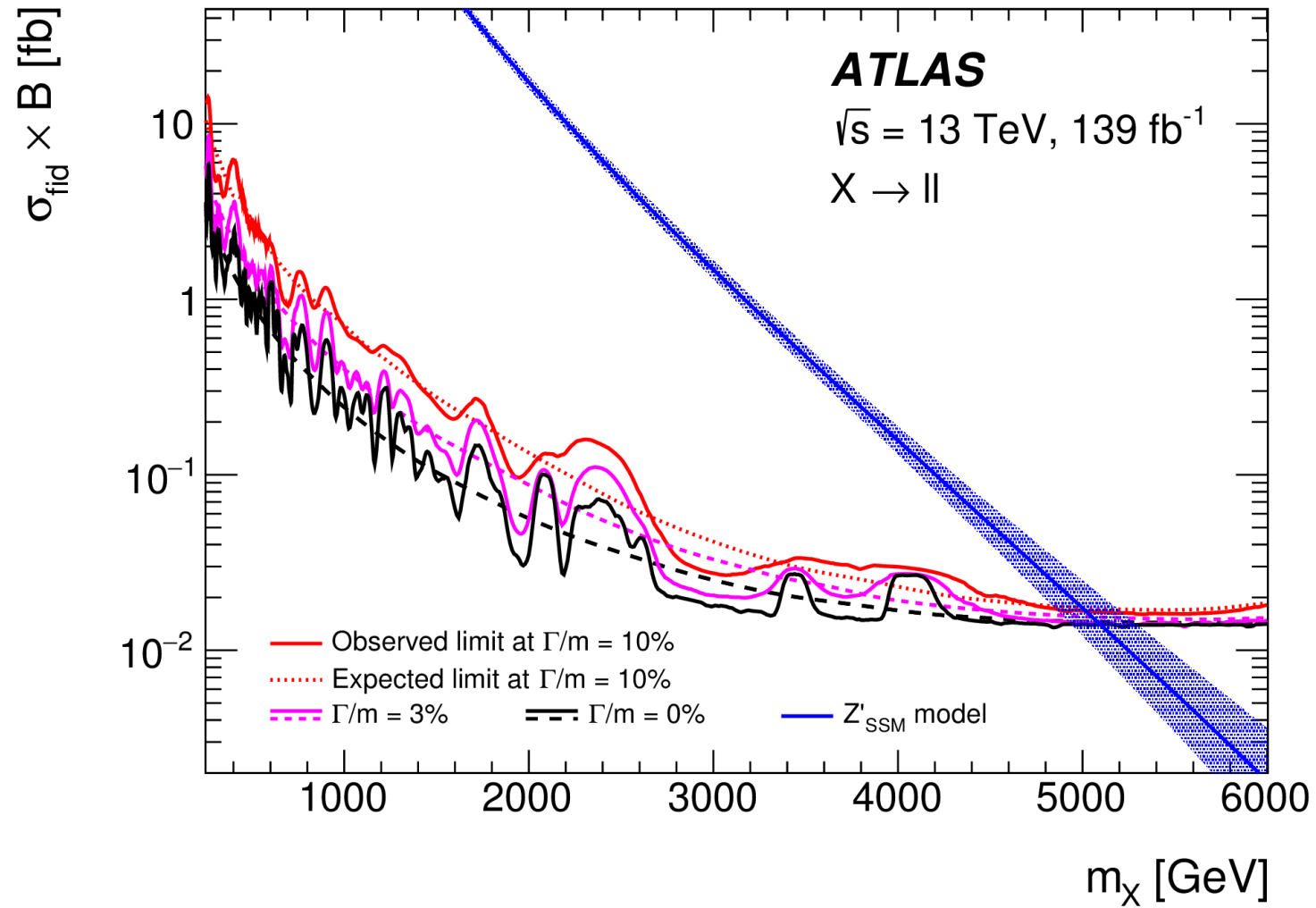
EXCLUSION LIMITS



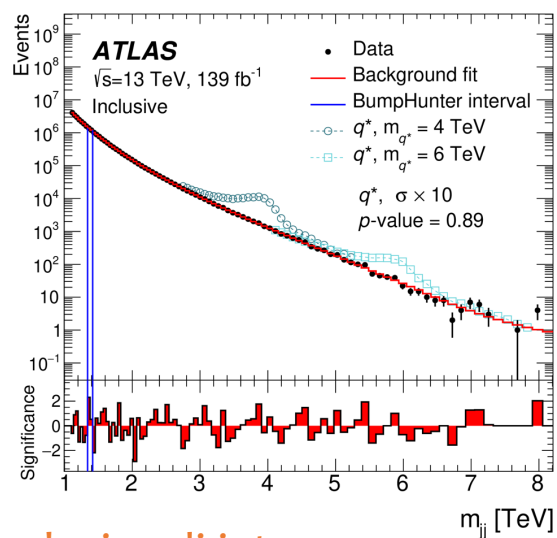
EVOLUTION...



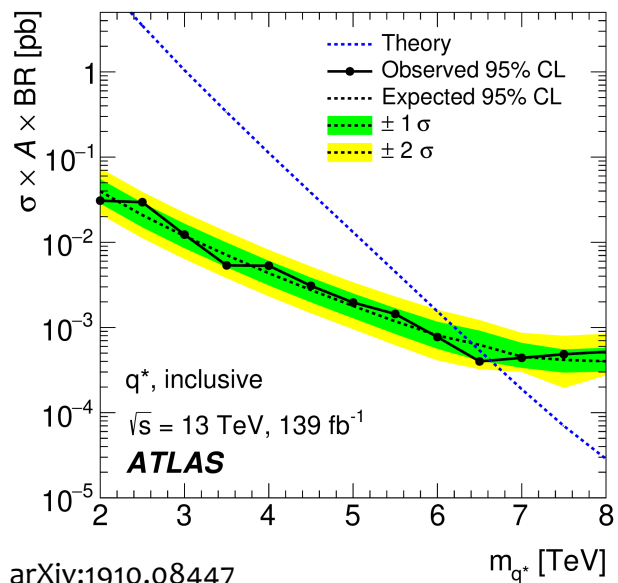
EXCLUSION LIMITS



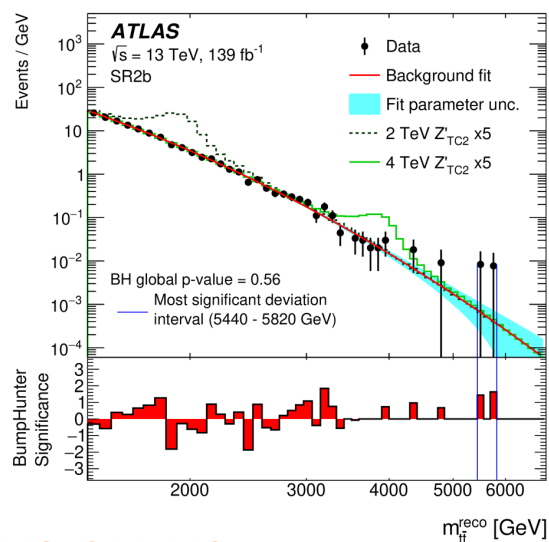
HIGH MASS RESONANCES



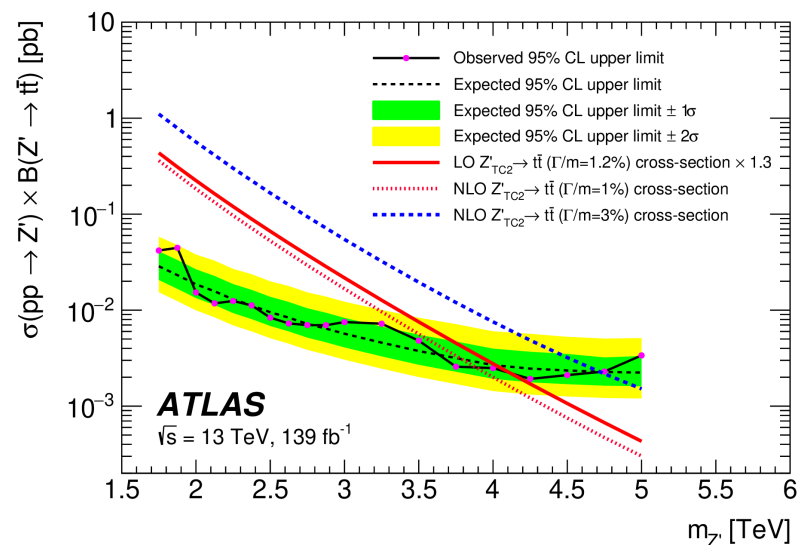
inclusive di-jet



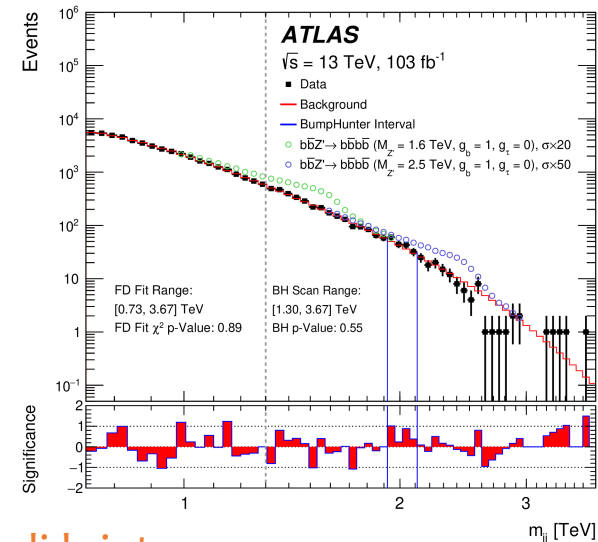
arXiv:1910.08447



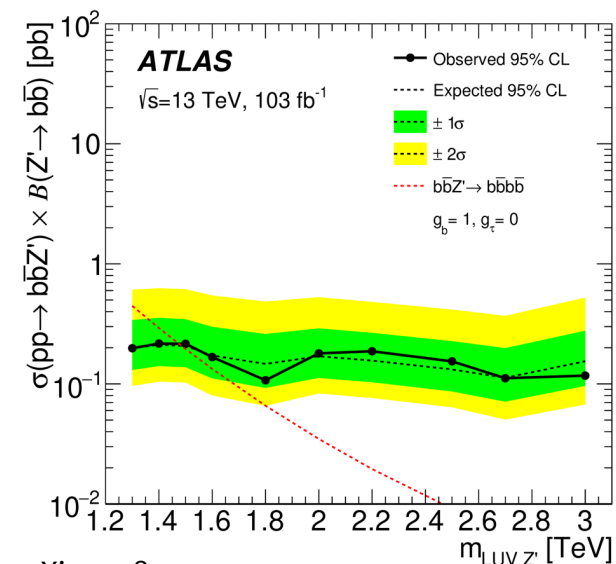
ttbar resonances



arXiv:2005.05138

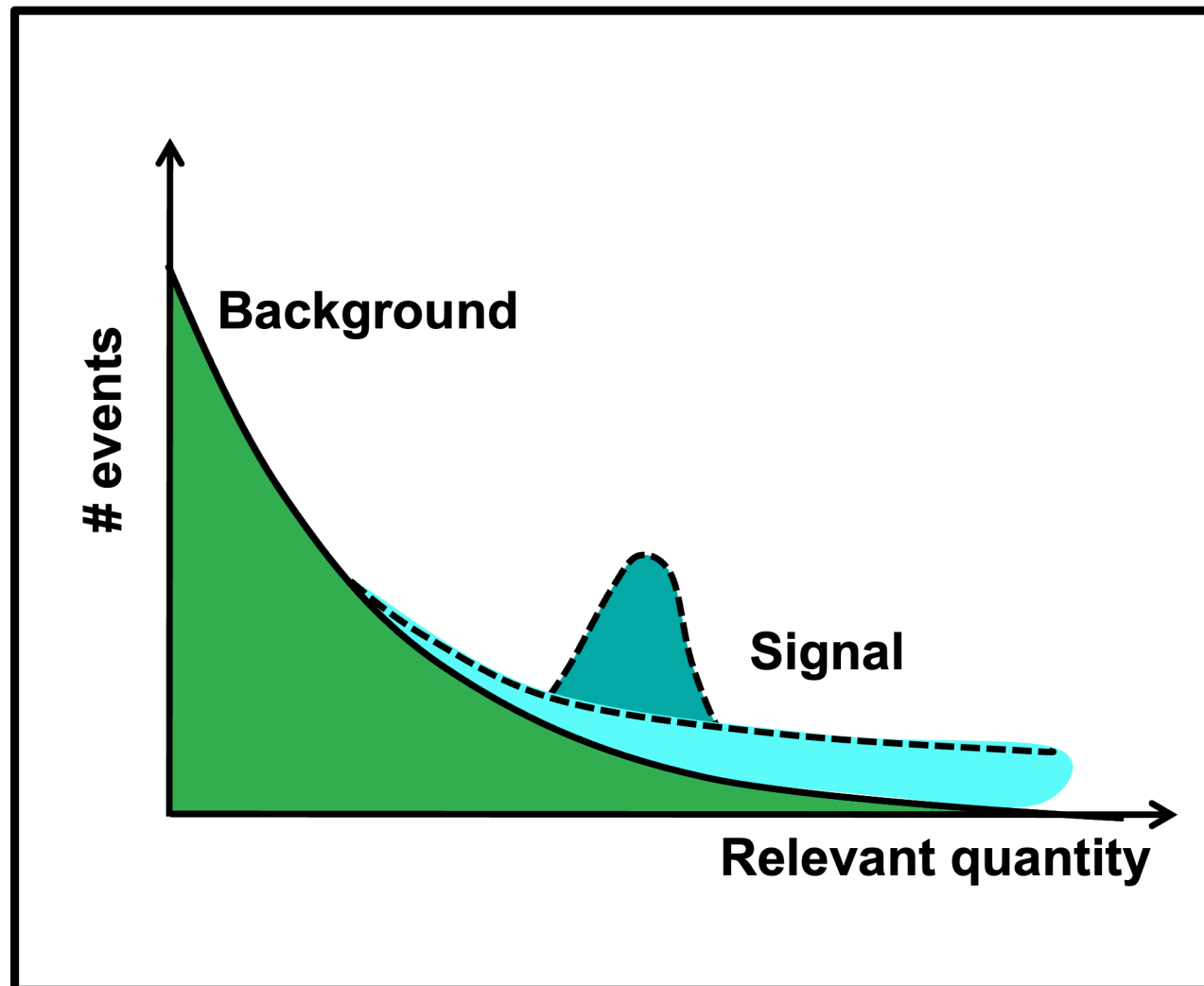


di-b-jet

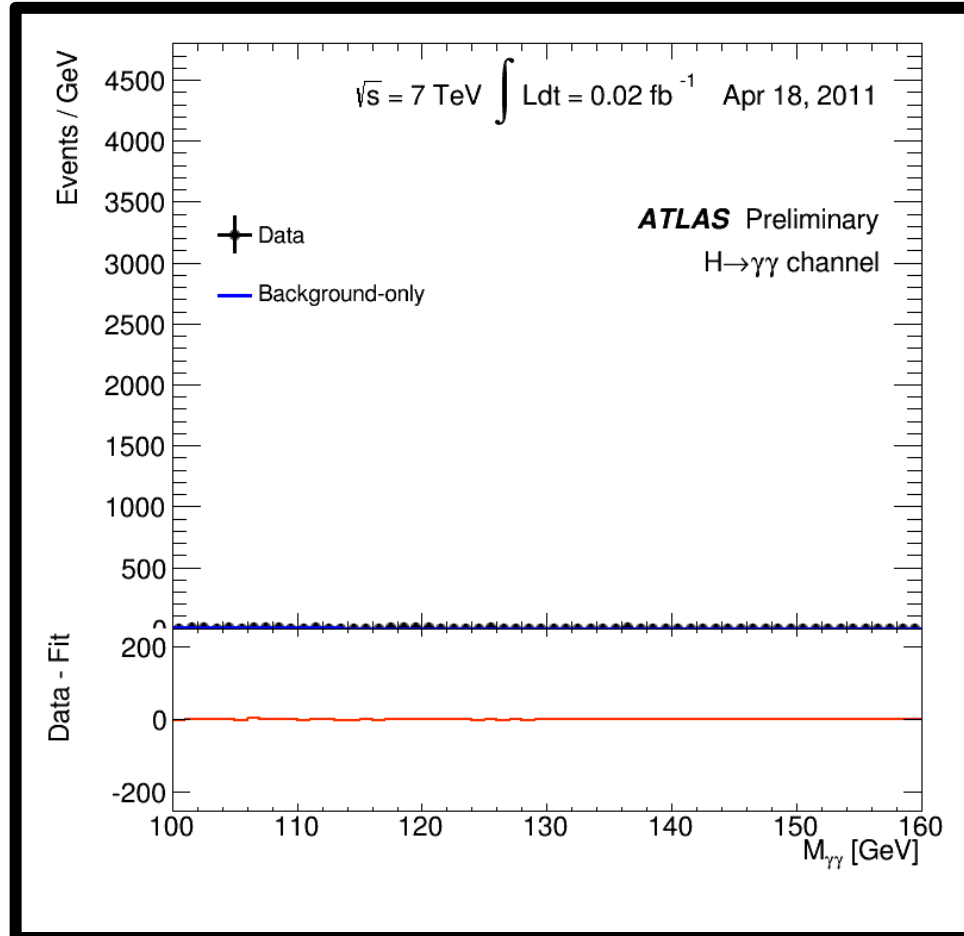


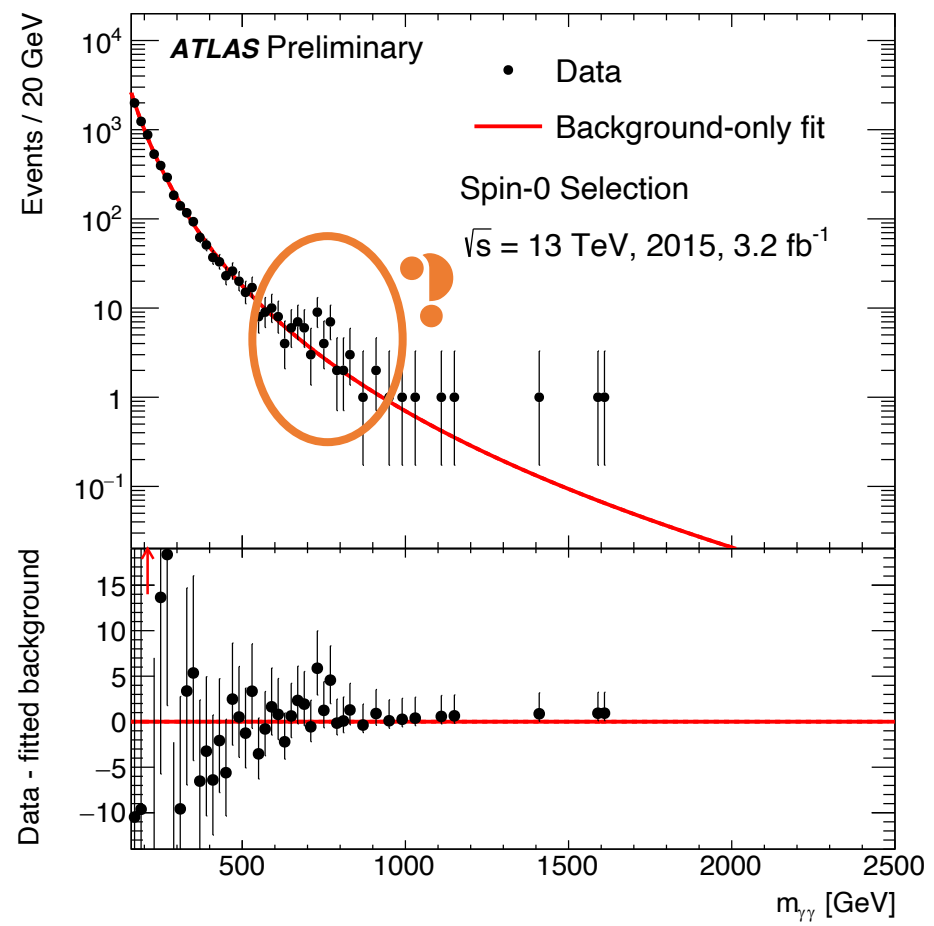
arXiv:2108.09059

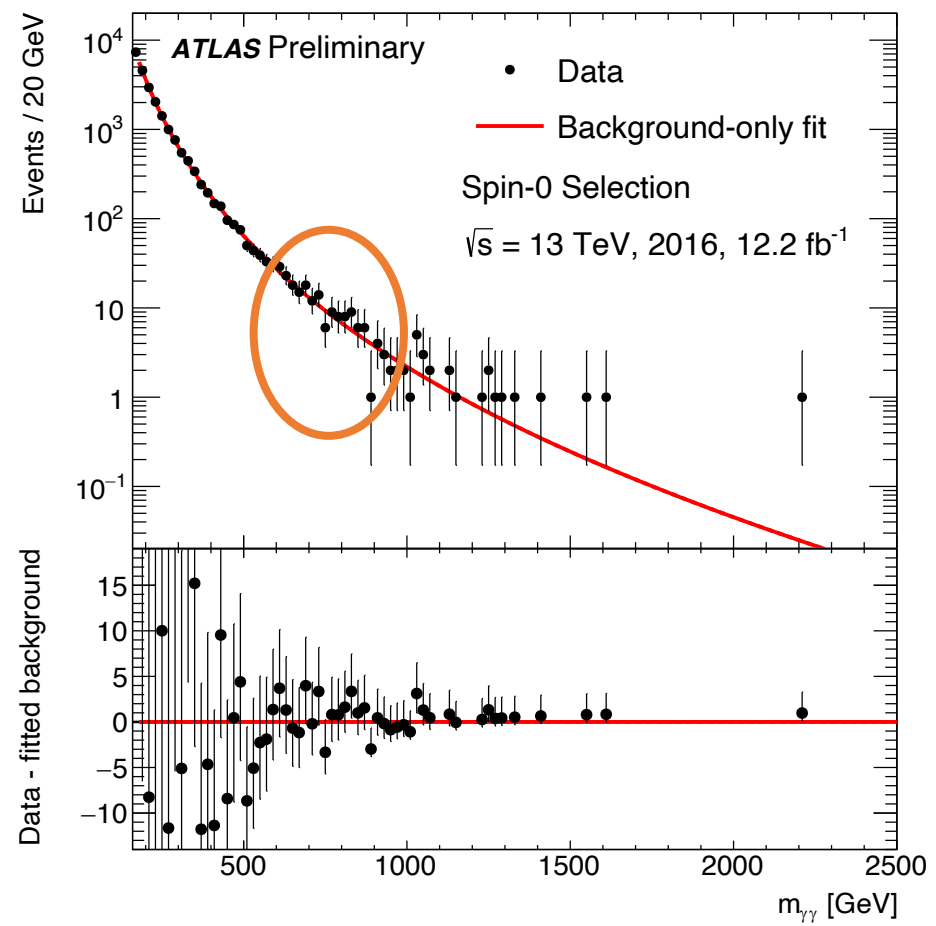
SEARCHES



A WELL-KNOWN BUMP SEARCH







MARIO
096950

● x23

WORLD
1-4

TIME
593

THANK YOU MARIO!

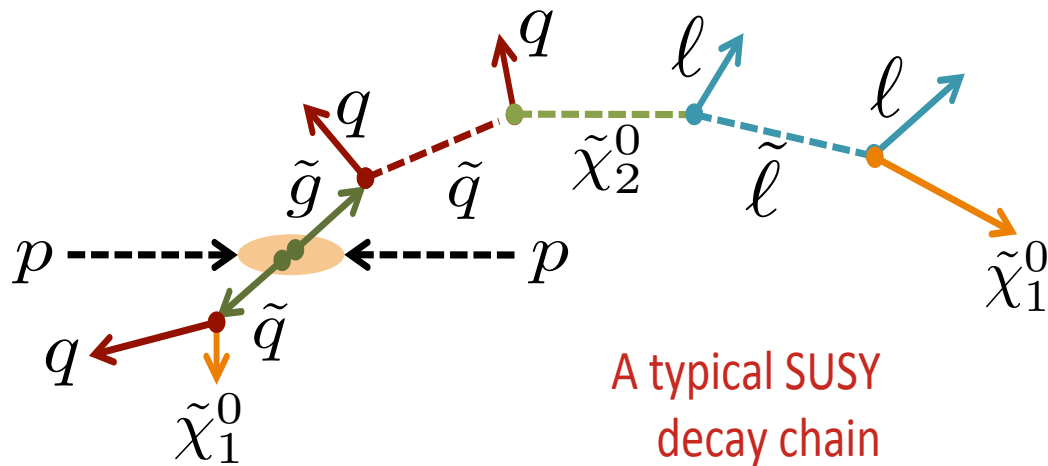
BUT OUR PRINCESS IS IN
ANOTHER CASTLE!



LHC PHYSICS

OTHER SEARCHES

ANOTHER SEARCH EXAMPLE: SEARCH FOR SUSY IN EVENTS WITH LARGE JET MULTIPLICITIES



Disclaimer:

This is only an example!

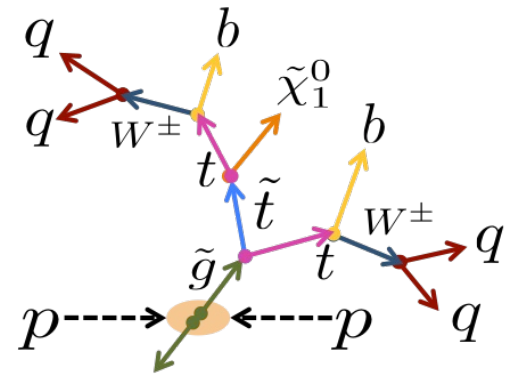
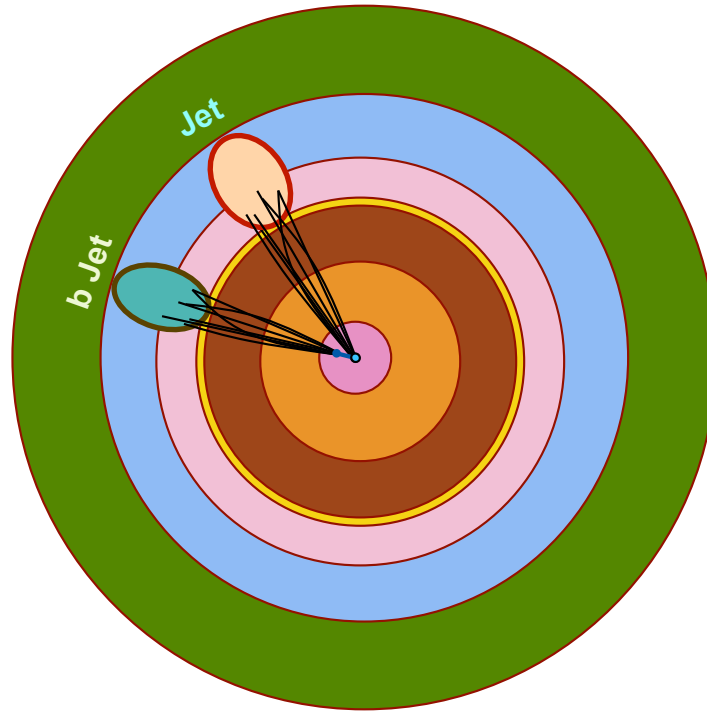
There are numerous such searches!
Each of them differs in

- event selections,
- background determinations,
- methodology

SEARCHING FOR NEW PHYSICS IS FUN!

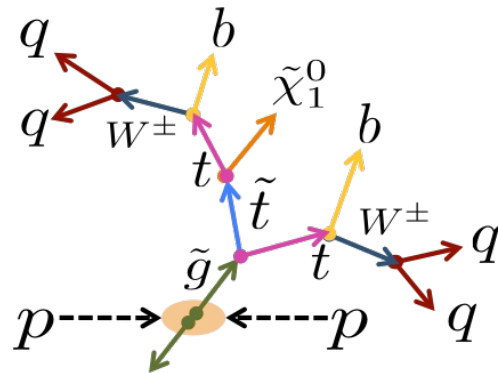
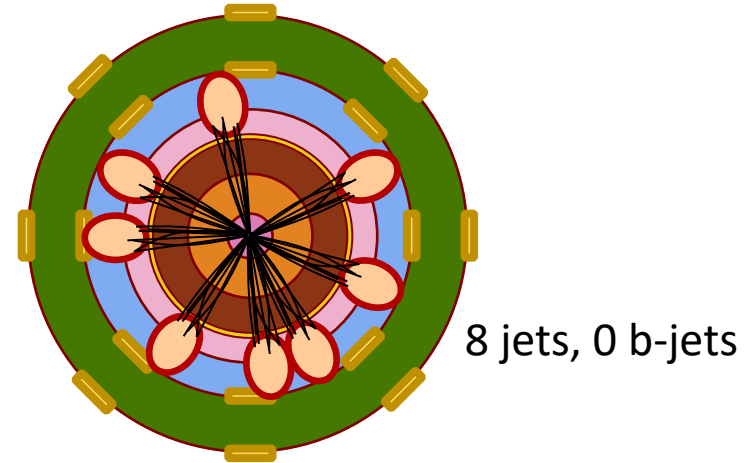
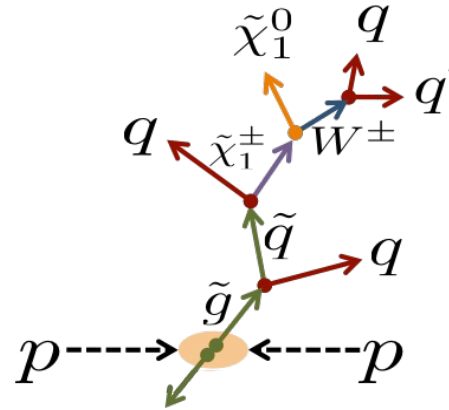
EVENT SELECTIONS

b-jets



EVENT SELECTIONS

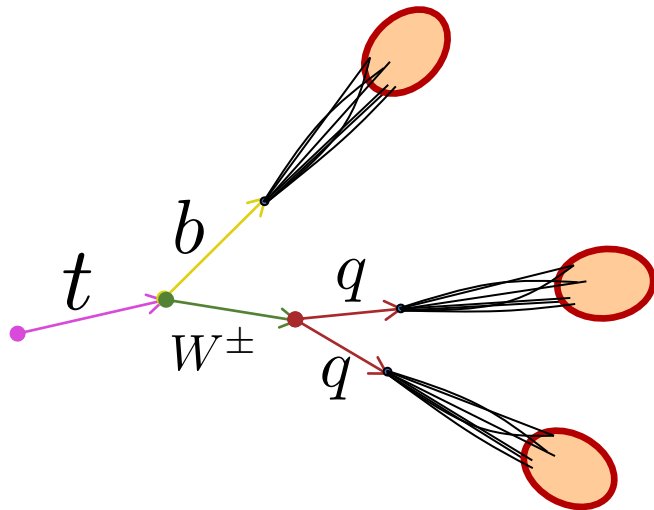
b-jets



Signal regions can range in jet p_T and jet & b-jet multiplicity.

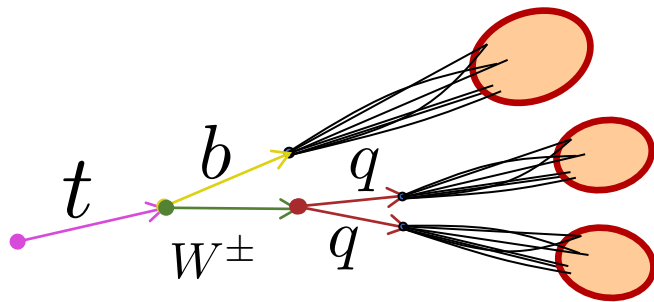
EVENT SELECTIONS

“fat-jets”



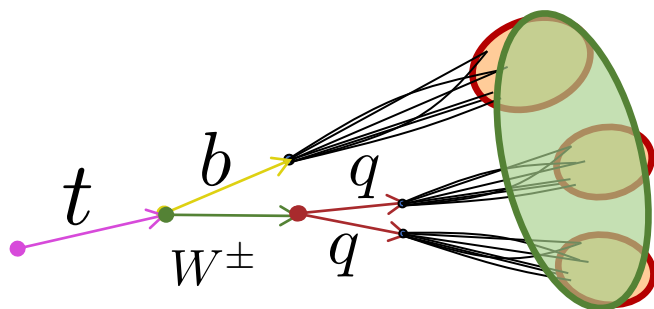
EVENT SELECTIONS

“fat-jets”



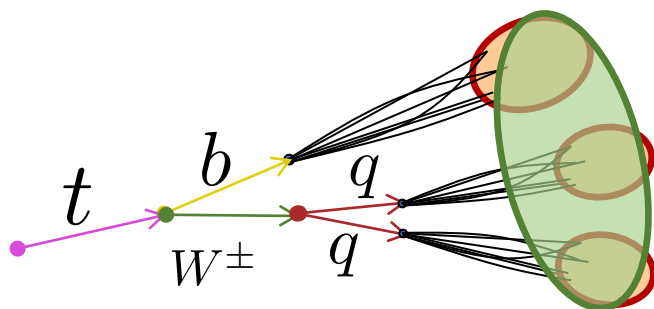
EVENT SELECTIONS

“fat-jets”



EVENT SELECTIONS

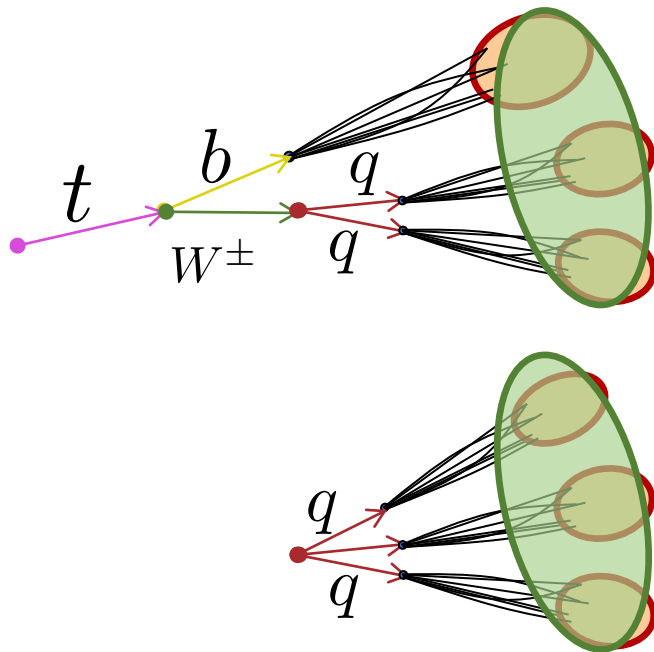
“fat-jets”



Fat-jets are a key signature in searches for boosted objects, e.g. boosted tops.

EVENT SELECTIONS

“fat-jets”

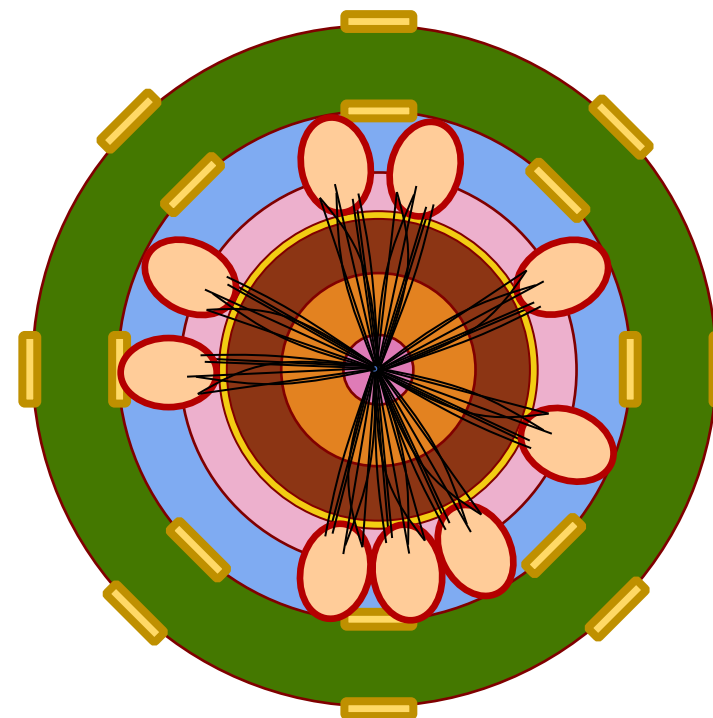
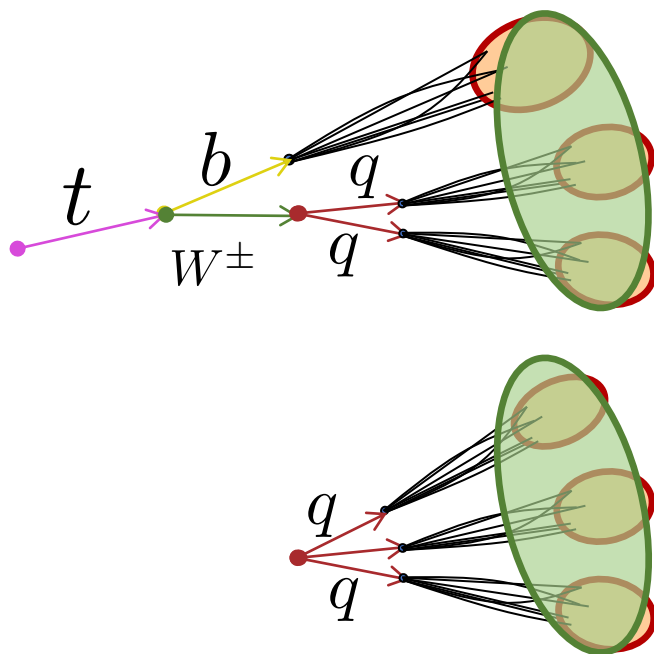


$$m_j \text{ (QCD)} < m_j \text{ (SUSY)}$$

Proposed in arXiv:1202.0558

EVENT SELECTIONS

“fat-jets”

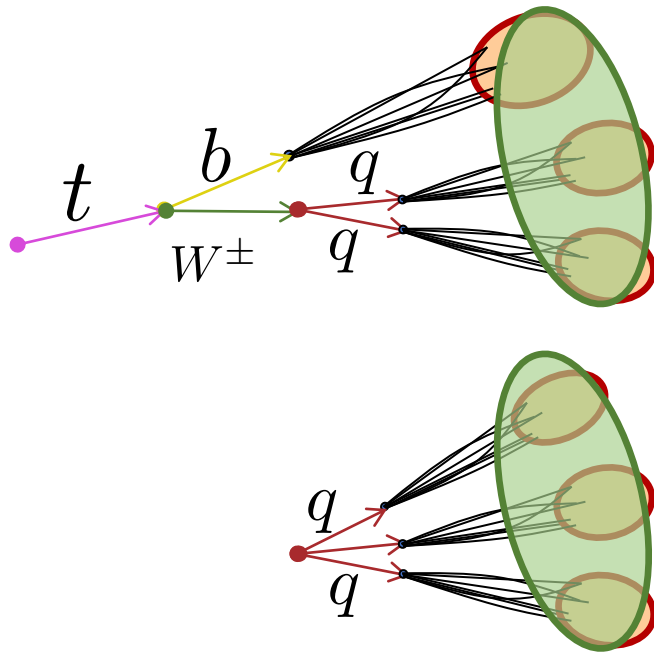


$$m_j \text{ (QCD)} < m_j \text{ (SUSY)}$$

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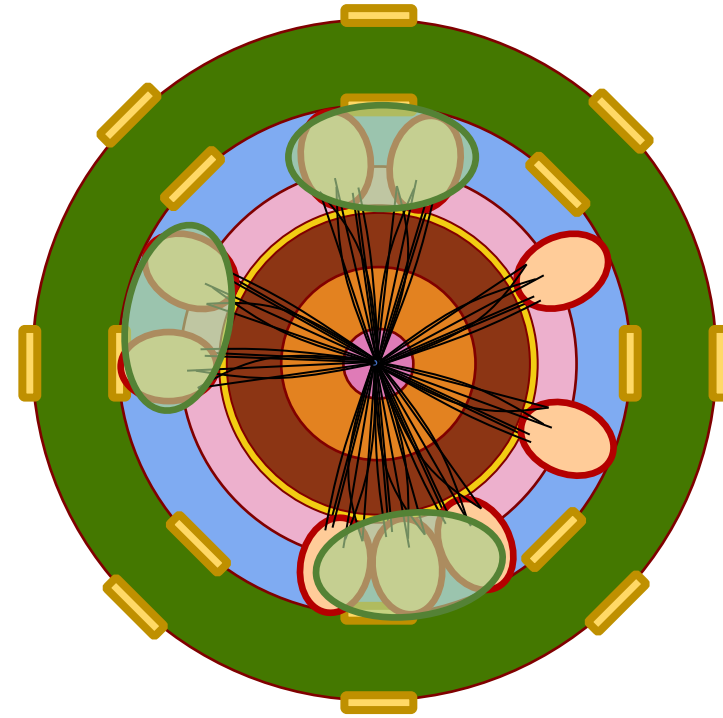
EVENT SELECTIONS

“fat-jets”



$$m_j \text{ (QCD)} < m_j \text{ (SUSY)}$$

Proposed in arXiv:1202.0558



$$M_J^\Sigma = \sum_{i=1}^{n_J} m_{j_i}$$

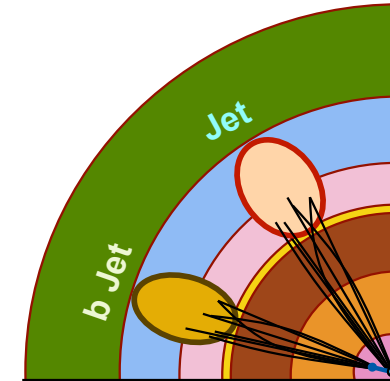
Signal regions can range in jet multiplicity and M_J^Σ cuts.

EVENT SELECTIONS

An example of a search

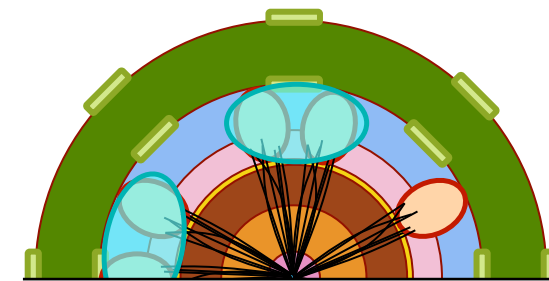
“b-jet stream”

ID	8j50			9j50			≥10j50	7j80			≥8j80		
Jet $ \eta $	< 2.0												
Jet p_T	50 GeV							80 GeV					
Jet count	=8			=9			≥10	=7			≥8		
b-jets	0	1	≥2	0	1	≥2	-	0	1	≥2	0	1	≥2
ME_T/VH_T	> 4 GeV ^½												



“fat-jet stream”

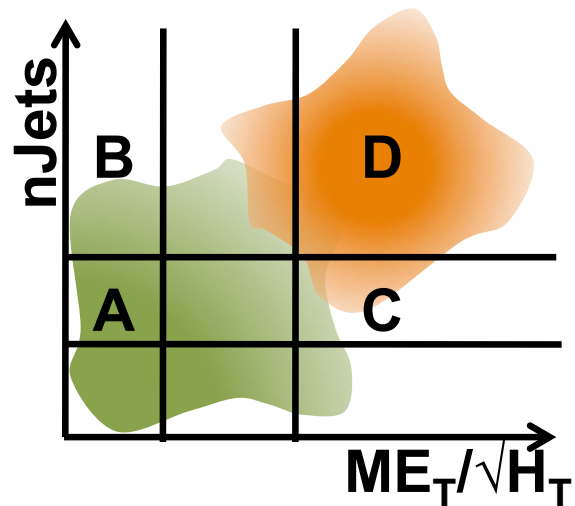
ID	≥8j50		≥9j50		≥10j50	
Jet $ \eta $	< 2.8					
Jet p_T	50 GeV					
Jet count	≥8		≥9		≥10	
M_J^Σ (GeV)	>340	>420	>340	>420	>340	>420
ME_T/VH_T	> 4 GeV ^½					



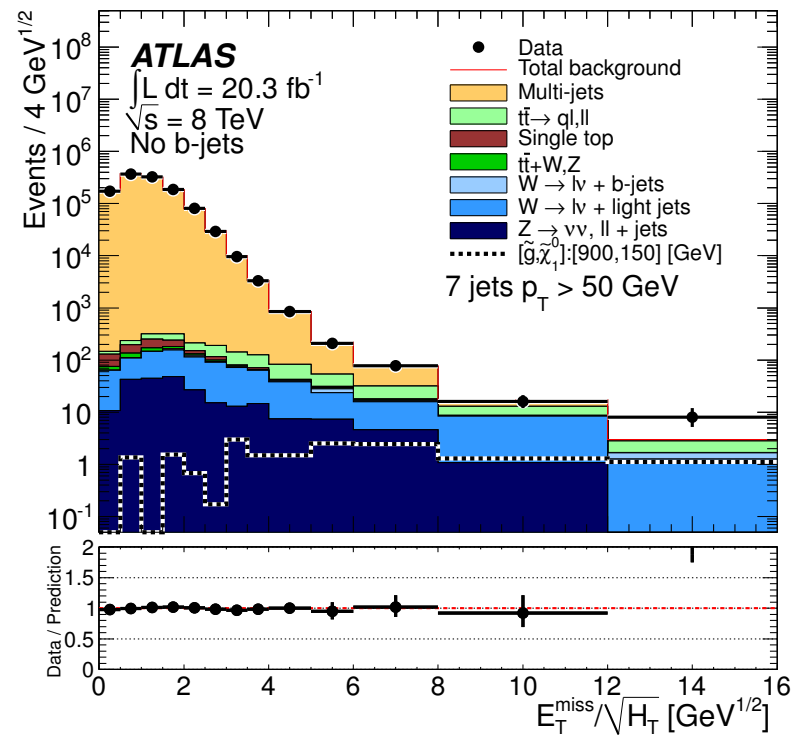
Proposed in arXiv:1202.0558

$$M_J^\Sigma = \sum_{i=1}^{n_J} m_{j_i}$$

BACKGROUND DETERMINATION



where $H_T = \sum_j p_T^j$

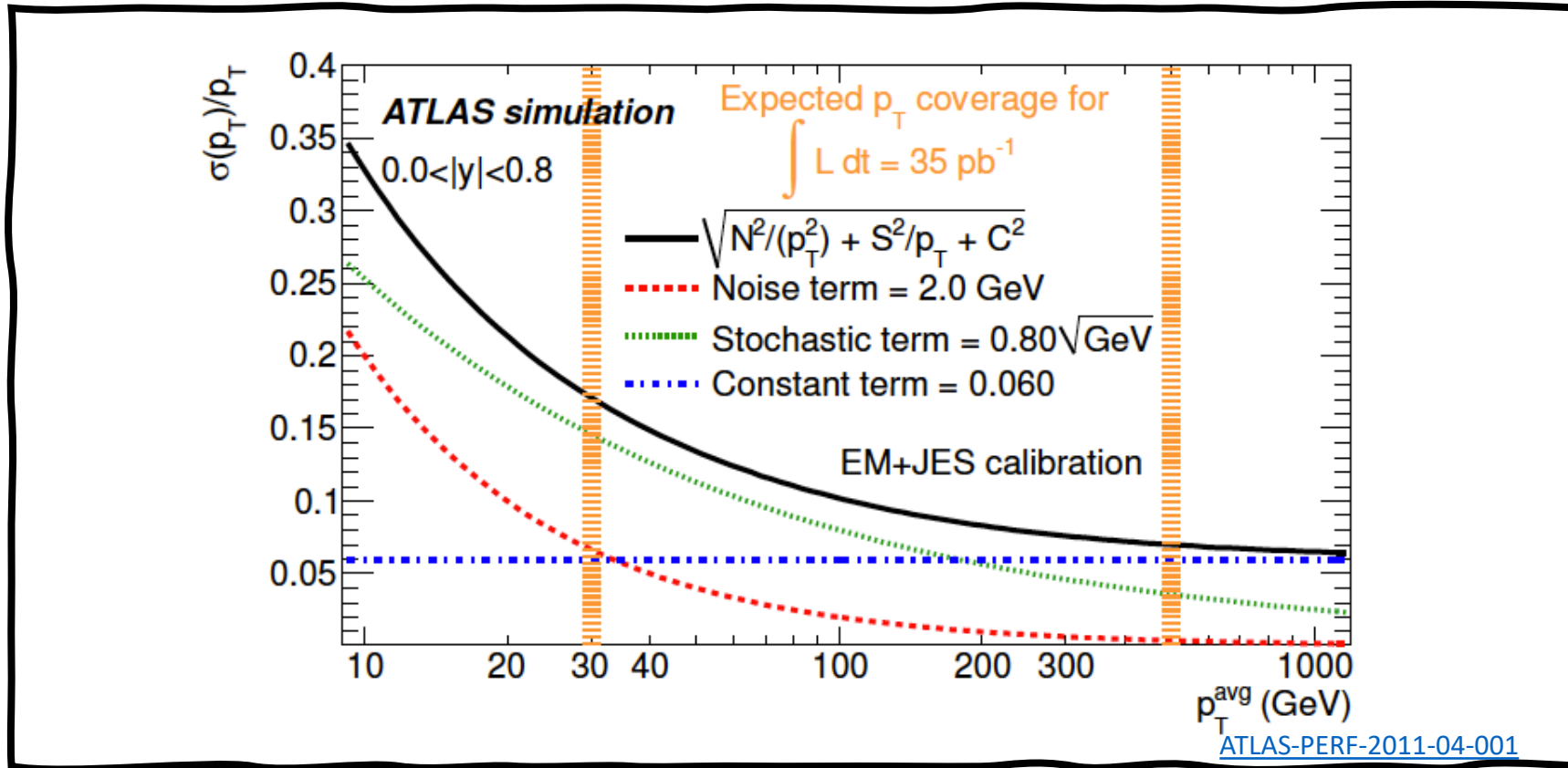


Why $ME_T / \sqrt{H_T}$?

\Rightarrow a measure of ME_T in units of standard deviations of the fake ME_T

$$\frac{\sigma_{p_T}}{p_T} = \frac{N}{p_T} \oplus \frac{S}{\sqrt{p_T}} \oplus C$$

BACKGROUND DETERMINATION



Why ME_T/VH_T ?

\Rightarrow a measure of ME_T in units of standard deviations of the fake ME_T

$$\frac{\sigma_{p_T}}{p_T} = \frac{N}{p_T} \oplus \frac{S}{\sqrt{p_T}} \oplus C$$

RESULTS

b-jet stream

ID	8j50			9j50			≥10j50
b-jets	0	1	≥2	0	1	≥2	0
Expected evts	35±4	40±10	50±10	3.3±0.7	6.1±1.7	8.0±2.7	1.37±0.35
Observed evts	40	44	44	5	8	7	3
Significance (σ)	0.7	-0.02	-0.6	0.8	0.6	-0.28	1.11

ID	7j80			≥8j80		
b-jets	0	1	≥2	0	1	≥2
Expected evts	11.0±2.2	17±6	25±10	0.9±0.6	1.5±0.9	3.3±2.2
Observed evts	12	17	13	2	1	3
Significance (σ)	0.05	-0.14	-1.0	0.9	-0.28	-0.06

fat-jet stream

ID	≥8j50		≥9j50		≥10j50	
M_T^Σ (GeV)	340	420	340	420	340	420
Expected evts	75±19	45±14	17±7	11±5	3.2±3.7	2.2±2.0
Observed evts	69	37	13	9	1	1
Significance (σ)	-0.27	-0.6	-0.6	-0.34	-0.8	-0.6

RESULTS

b-jet stream

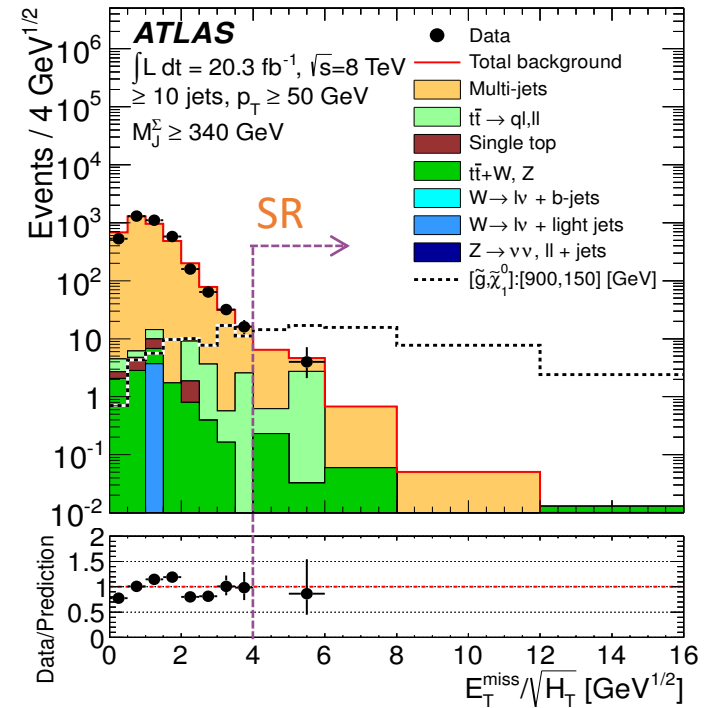
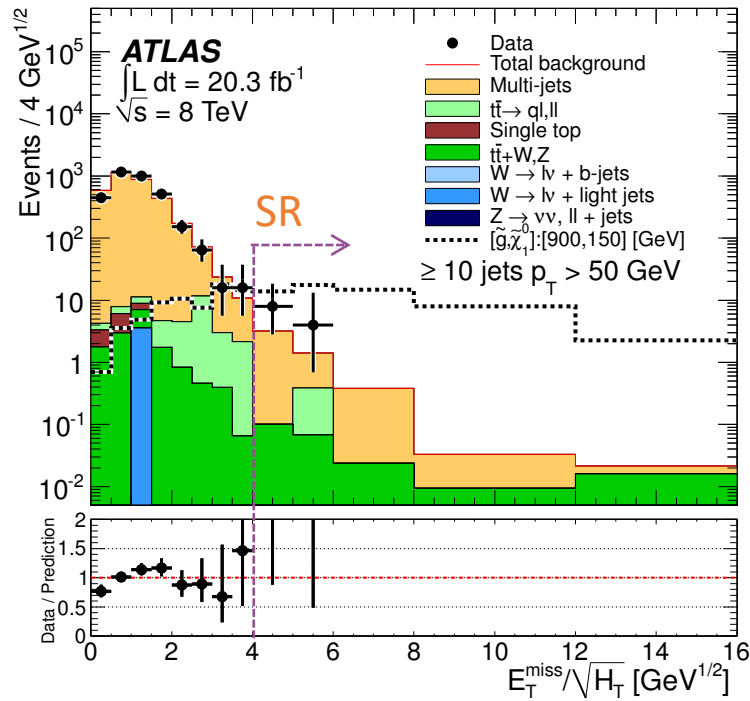
ID	8j50			9j50			≥10j50
b-jets	0	1	≥2	0	1	≥2	0
Expected evts	35±4	40±10	50±10	3.3±0.7	6.1±1.7	8.0±2.7	1.37±0.35
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fat-jet stream

ID	≥8j50		≥9j50		≥10j50	
M_T^Σ (GeV)	340	420	340	420	340	420
Expected evts	75±19	45±14	17±7	11±5	3.2±3.7	2.2±2.0
Observed evts	69	37	13	9	1	1
Significance (σ)	-0.27	-0.6	-0.6	-0.34	-0.8	-0.6

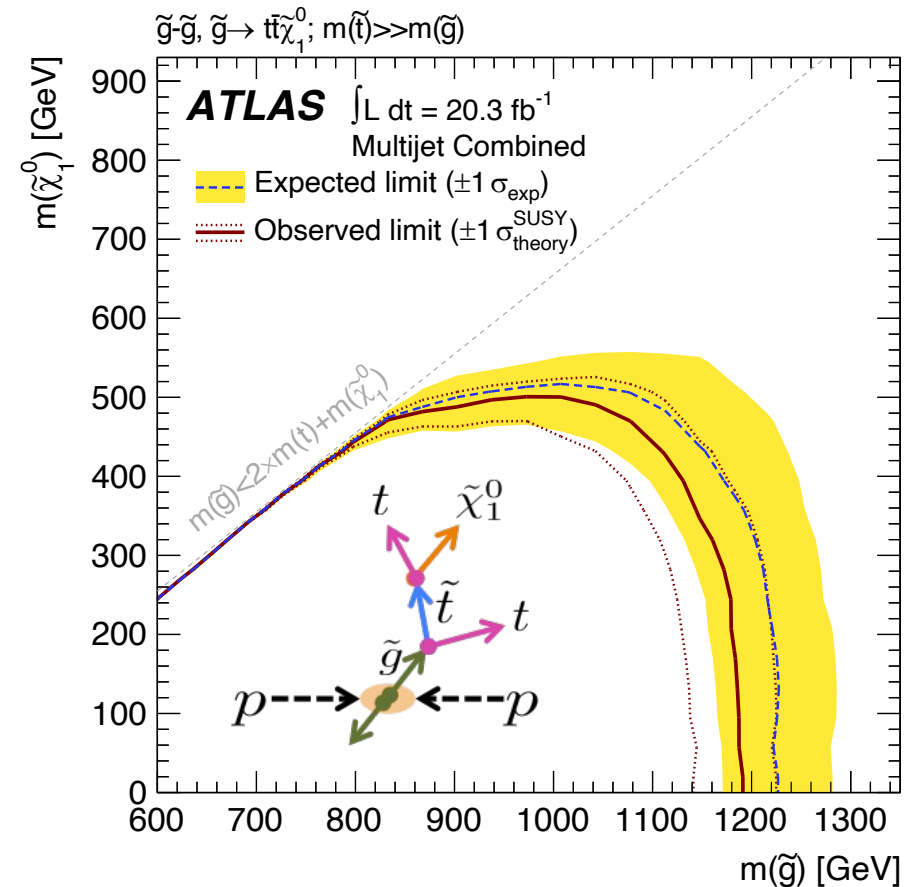
RESULTS



INTERPRETATIONS

Real or Simplified models

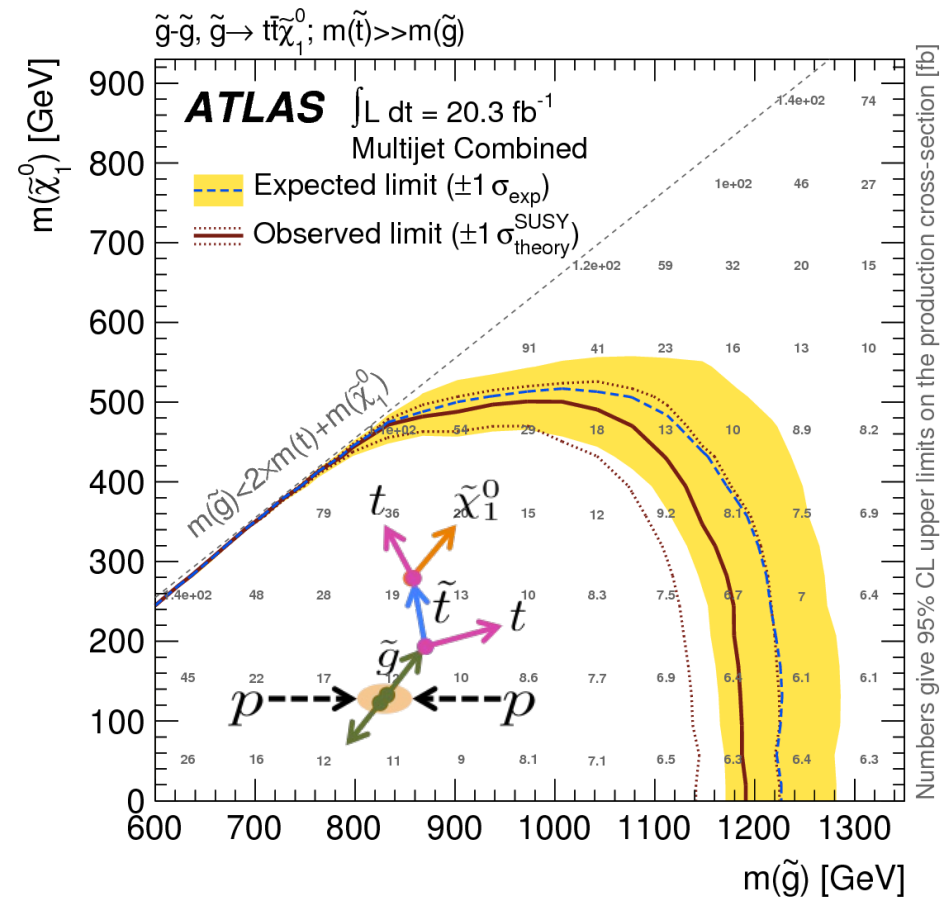
© Simplified topologies include typically one production and one decay process. Provide useful information for theorists.



INTERPRETATIONS

Real or Simplified models

- © Simplified topologies include typically one production and one decay process. Provide useful information for theorists.

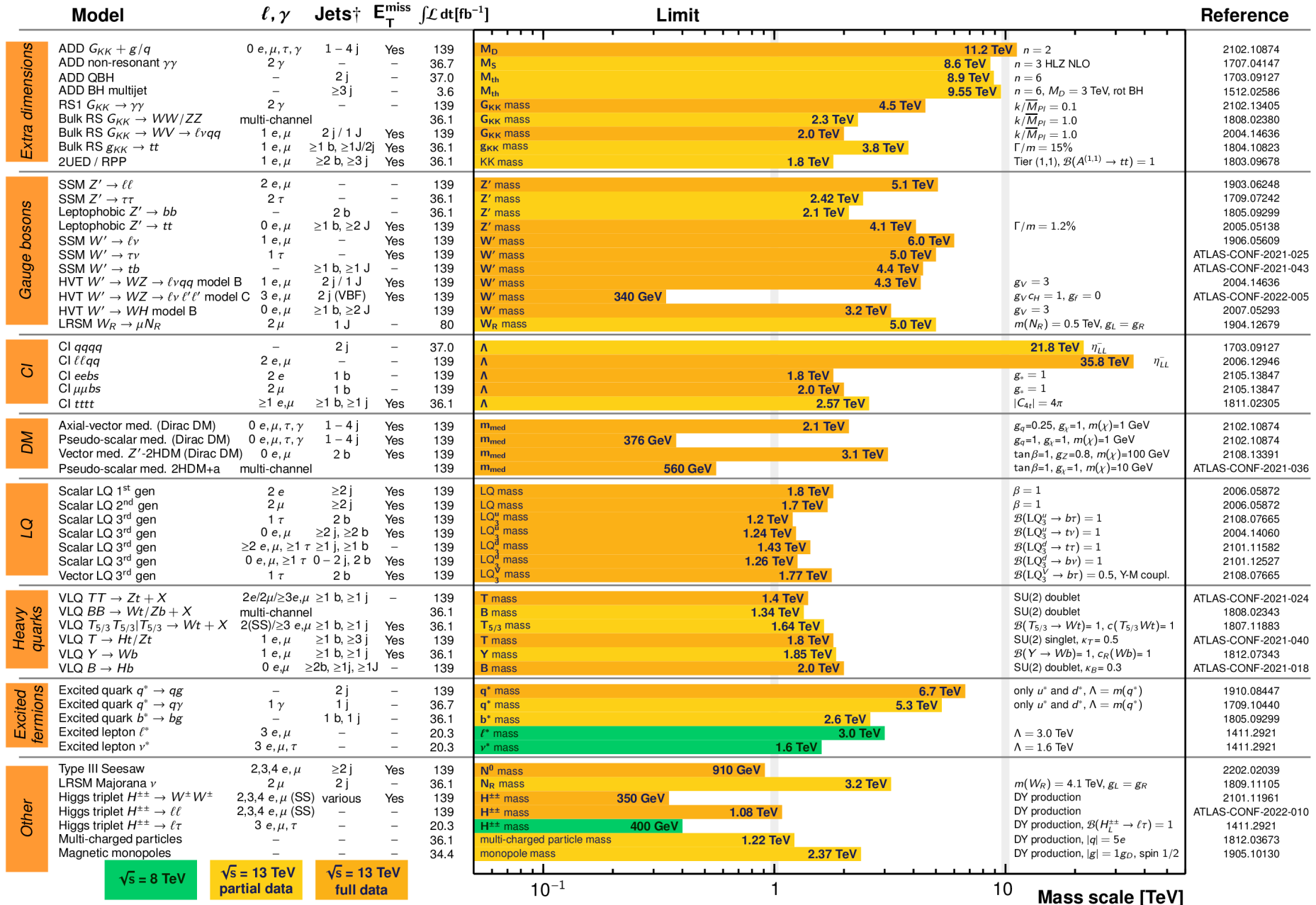


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$\sqrt{s} = 8 \text{ TeV}$ $\sqrt{s} = 13 \text{ TeV}$ partial data $\sqrt{s} = 13 \text{ TeV}$ full data

*Only a selection of the available mass limits on new states or phenomena is shown.

[†]Small-radius (large-radius) jets are denoted by the letter j (J).

Model	Signature	$\int \mathcal{L} dt$ [fb ⁻¹]	Mass limit	Reference			
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	E_T^{miss} 139 E_T^{miss} 139	\tilde{q} [1x, 8x Degen.] 1.0 \tilde{q} [8x Degen.] 0.9	$m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 5$ GeV	2010.14293 2102.10874	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ 2-6 jets	E_T^{miss} 139	\tilde{g} 2.3 Forbidden 1.15-1.95	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV	2010.14293 2010.14293	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, μ 2-6 jets	E_T^{miss} 139	\tilde{g} 2.2	$m(\tilde{\chi}_1^0) < 600$ GeV	2101.01629	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$ 2 jets	E_T^{miss} 139	\tilde{g} 2.2	$m(\tilde{\chi}_1^0) < 700$ GeV	CERN-EP-2022-014	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	0 e, μ SS e, μ	7-11 jets 6 jets	E_T^{miss} 139 E_T^{miss} 139	\tilde{g} 1.97 \tilde{g} 1.15	$m(\tilde{\chi}_1^0) < 600$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	2008.06032 1909.08457
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	3 b 6 jets	E_T^{miss} 79.8 E_T^{miss} 139	\tilde{g} 2.25 \tilde{g} 1.25	$m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	ATLAS-CONF-2018-041 1909.08457
	3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1$	0 e, μ 2 b	E_T^{miss} 139	\tilde{b}_1 1.255 \tilde{b}_1 0.68	$m(\tilde{\chi}_1^0) < 400$ GeV 10 GeV $< \Delta m(\tilde{b}_1, \tilde{\chi}_1^0) < 20$ GeV	2101.12527 2101.12527
$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$		0 e, μ 2 τ	6 b 2 b	E_T^{miss} 139 E_T^{miss} 139	Forbidden 0.23-1.35 \tilde{b}_1 0.13-0.85	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV	1908.03122 2103.08189
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$		0-1 e, μ ≥ 1 jet	E_T^{miss} 139	\tilde{t}_1 1.25	$m(\tilde{\chi}_1^0) = 1$ GeV	2004.14060, 2012.03799	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$		1 e, μ 3 jets/1 b	E_T^{miss} 139	Forbidden 0.65	$m(\tilde{\chi}_1^0) = 500$ GeV	2012.03799	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$		1-2 τ 2 jets/1 b	E_T^{miss} 139	Forbidden 1.4	$m(\tilde{\tau}_1) = 800$ GeV	2108.07665	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$		0 e, μ 2 c mono-jet	E_T^{miss} 36.1 E_T^{miss} 139	\tilde{c} 0.85 \tilde{t}_1 0.55	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 2102.10874	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$ $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$		1-2 e, μ 1-4 b 3 e, μ	1 b 1 b	E_T^{miss} 139 E_T^{miss} 139	\tilde{t}_1 0.067-1.18 \tilde{t}_2 Forbidden 0.86	$m(\tilde{\chi}_2^0) = 500$ GeV $m(\tilde{\chi}_1^0) = 360$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 40$ GeV	2006.05880 2006.05880
EW direct	$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WZ	Multiple ℓ /jets $ee, \mu\mu$	E_T^{miss} 139 E_T^{miss} 139	$\tilde{\chi}_1^\pm / \tilde{\chi}_2^0$ 0.96 $\tilde{\chi}_1^\pm / \tilde{\chi}_2^0$ 0.205	$m(\tilde{\chi}_1^0) = 0$, wino-bino $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV, wino-bino	2106.01676, 2108.07586 1911.12606	
	$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$ via WW	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 0.42	$m(\tilde{\chi}_1^0) = 0$, wino-bino	1908.08215	
	$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via Wh	Multiple ℓ /jets	E_T^{miss} 139	$\tilde{\chi}_1^\pm / \tilde{\chi}_2^0$ Forbidden 1.06	$m(\tilde{\chi}_1^0) = 70$ GeV, wino-bino	2004.10894, 2108.07586	
	$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, μ	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 1.0	$m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$	1908.08215	
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$	2 τ	E_T^{miss} 139	$\tilde{\tau}$ [$\tilde{\tau}_L, \tilde{\tau}_{R,L}$] 0.16-0.3 0.12-0.39	$m(\tilde{\chi}_1^0) = 0$	1911.06660	
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ ≥ 1 jet	E_T^{miss} 139 E_T^{miss} 139	$\tilde{\ell}$ 0.7 $\tilde{\ell}$ 0.256	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 10$ GeV	1908.08215 1911.12606	
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ 0 e, μ	≥ 3 b 0 jets ≥ 2 large jets	E_T^{miss} 36.1 E_T^{miss} 139 E_T^{miss} 139	\tilde{H} 0.13-0.23 \tilde{H} 0.55 \tilde{H} 0.29-0.88 \tilde{H} 0.45-0.93	$BR(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $BR(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$ $BR(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$	1806.04030 2103.11684 2108.07586
Long-lived particles	Direct $\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk 1 jet	E_T^{miss} 139	$\tilde{\chi}_1^\pm$ 0.66 $\tilde{\chi}_1^\pm$ 0.21	Pure Wino Pure higgsino	2201.02472 2201.02472	
	Stable \tilde{g} R-hadron	pixel dE/dx	E_T^{miss} 139	\tilde{g} 2.05	$m(\tilde{\chi}_1^0) = 100$ GeV	CERN-EP-2022-029	
	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	pixel dE/dx	E_T^{miss} 139	\tilde{g} [$\tau(\tilde{g}) = 10$ ns] 2.2	$m(\tilde{\chi}_1^0) = 100$ GeV	CERN-EP-2022-029	
	$\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell\tilde{G}$	Displ. lep	E_T^{miss} 139	$\tilde{\ell}, \tilde{\mu}$ 0.7 $\tilde{\tau}$ 0.34 $\tilde{\tau}$ 0.36	$\tau(\tilde{\ell}) = 0.1$ ns $\tau(\tilde{\ell}) = 0.1$ ns $\tau(\tilde{\ell}) = 10$ ns	2011.07812 2011.07812 CERN-EP-2022-029	
RPV	$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp / \tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow Z\ell \rightarrow \ell\ell\ell$	3 e, μ	139	$\tilde{\chi}_1^\pm / \tilde{\chi}_1^0$ [BR(Z τ)=1, BR(Z e)=1] 0.625 1.05	Pure Wino	2011.10543	
	$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp / \tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\nu\nu$	4 e, μ 0 jets	E_T^{miss} 139	$\tilde{\chi}_1^\pm / \tilde{\chi}_2^0$ [$\lambda_{133} \neq 0, \lambda_{12k} \neq 0$] 0.95 1.55	$m(\tilde{\chi}_1^0) = 200$ GeV	2103.11684	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow q\tilde{q}\tilde{\chi}_1^0$	4-5 large jets	36.1	\tilde{g} [$m(\tilde{\chi}_1^0) = 200$ GeV, 1100 GeV] 1.3 1.9	Large $\lambda'_{1,2}$	1804.03568	
	$\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow t\tilde{b}s$	Multiple	36.1	\tilde{t}_1 [$\lambda'_{323} = 2e-4, 1e-2$] 0.55 1.05	$m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	ATLAS-CONF-2018-003	
	$\tilde{u}, \tilde{t} \rightarrow b\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow b\tilde{b}s$	$\geq 4b$	139	\tilde{t}_1 Forbidden 0.95	$m(\tilde{\chi}_1^0) = 500$ GeV	2010.01015	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{s}$	2 jets + 2 b	36.7	\tilde{t}_1 [qq, bs] 0.42 0.61	$BR(\tilde{t}_1 \rightarrow b\ell/b\mu) > 20\%$	1710.07171	
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\tilde{\ell}$	2 e, μ 1 μ DV	2 b DV	36.1 136	\tilde{t}_1 0.4-1.45 \tilde{t}_1 [1e-10 < $\lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$] 1.0 1.6	$BR(\tilde{t}_1 \rightarrow q\mu) = 100\%, \cos\theta = 1$	1710.05544 2003.11956
$\tilde{\chi}_1^\pm / \tilde{\chi}_2^0 / \tilde{\chi}_1^0, \tilde{\chi}_{1,2}^0 \rightarrow t\tilde{b}s, \tilde{\chi}_1^\pm \rightarrow b\tilde{b}s$	1-2 e, μ ≥ 6 jets	139	$\tilde{\chi}_1^0$ 0.2-0.32	Pure higgsino	2106.09609		

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

FROM RAW DATA TO PHYSICS
INSTEAD OF SUMMARY:

COMPONENTS OF AN ANALYSIS

COMPONENTS OF A PHYSICS ANALYSIS

- ◎ Data-set and Monte Carlo samples
- ◎ Trigger
- ◎ Object definitions and event selections
- ◎ Background determination
- ◎ Systematic uncertainties
- ◎ Statistical methods
- ◎ Results
- ◎ [Interpretations]

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- ◎ Statistical methods
- ◎ Results
- ◎ [Interpretations]

The data and simulation samples used in the analysis. Data for the measurement / search, simulation to compare data to predictions.

Data-set specifics:

- ◎ Data quality \Rightarrow Good run list.
- ◎ Luminosity.

Monte Carlo sample specifics:

- ◎ Generator, tunes.
- ◎ Statistics.

COMPONENTS OF A PHYSICS ANALYSIS

© Data-set and Monte Carlo samples

© Trigger

© Object definitions

© Background definitions

© Systematic uncertainties

© Statistical methods

© Results

© [Interpretations]

The trigger used to collect the data with.

Trigger specifics:

© Prescales; typically unprescaled triggers are used, prescaled triggers for QCD / high stat measurements.

© Trigger (in)efficiencies.

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- ◎ [Interpretations]

The exact definition of objects (electrons, muon, jets, ...) and how these are combined in selecting events to be analyzed.

Object definition specifics:

- ◎ “Flavor” of the identification (loose, medium, tight).
- ◎ Calibrations.

Event selection specifics:

- ◎ Event cleaning (e.g. from noise and cosmics).
- ◎ Momentum, geom. acceptance and multiplicity of objects.
- ◎ Higher level cuts, such as invariant mass.
- ◎ “**Signal regions**”.

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Events that are imitating the signal we are searching for or measuring.

Background determination specifics:

- ◎ Can/must be **data-driven** or **simulation-based**.
- ◎ “**Validation regions**” and “**control regions**” required. These can use different triggers wrt signal regions.

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- ◎ Any ‘intermediate’ measurement we have performed carries uncertainties (statistical and systematic).
- ◎ **“Systematic” uncertainties are introduced by inaccuracies in the methods used to perform the measurement.**
- ◎ Efficiencies, acceptance, number of events, luminosity, cross sections used in Monte Carlo scaling...
- ◎ Some of them are “centrally” assessed by the performance groups of an experiment. Some of them are analysis-specific.

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Dealing with large data-sets, we use statistical methods to make sense of the numbers we measure.

Typical method:

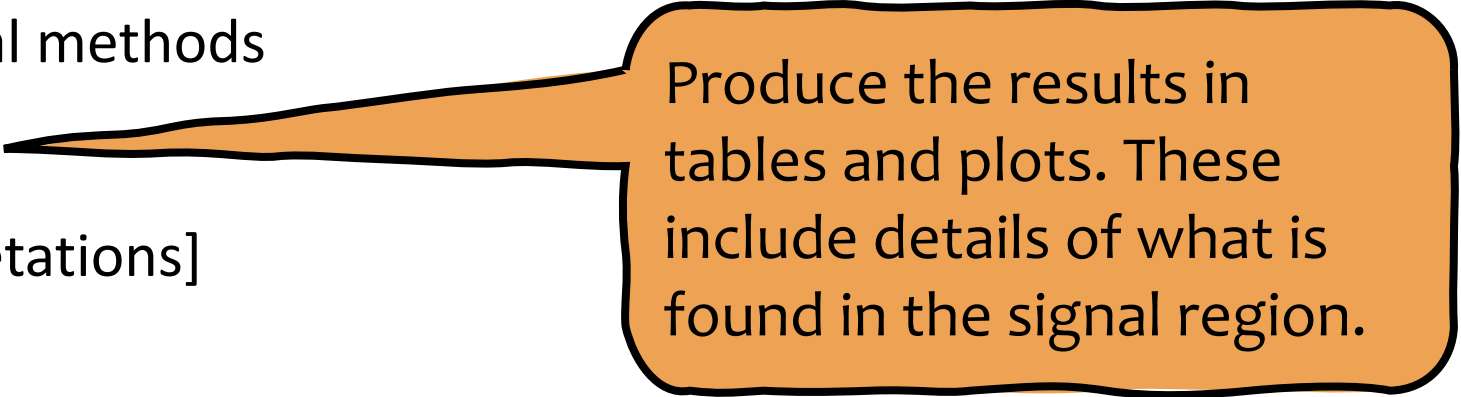
- ◎ Do a fit to extract signal from background.

Methodologies can vary a lot, but nowadays they are pretty unified within and across experiments.

Neural nets and other machine learning methods are broadly used, primarily to improve signal over background discrimination!

COMPONENTS OF A PHYSICS ANALYSIS

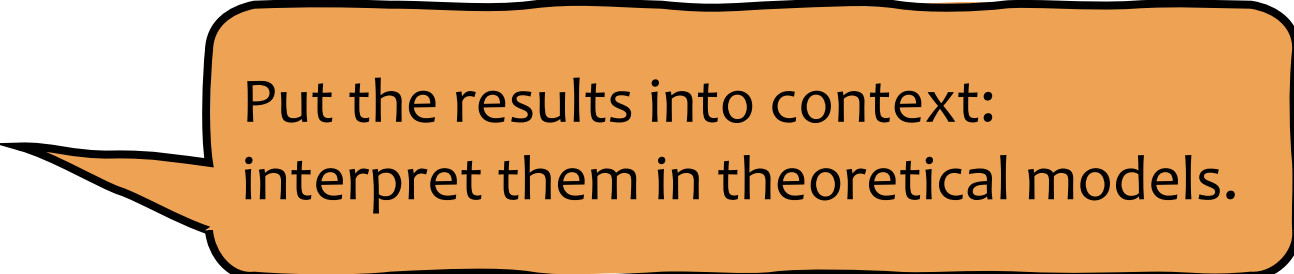
- ◎ Data-set and Monte Carlo samples
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Produce the results in tables and plots. These include details of what is found in the signal region.

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Put the results into context:
interpret them in theoretical models.



<https://xkcd.com/1838/>

EXCITING TIMES COMING UP IN HEP
GOOD LUCK IN YOUR RESEARCH!

Please get in touch for question, comments,
or simply feedback on this lecture

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