EFT analysis of New Physics at COHERENT

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In collaboration with

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Two big questions

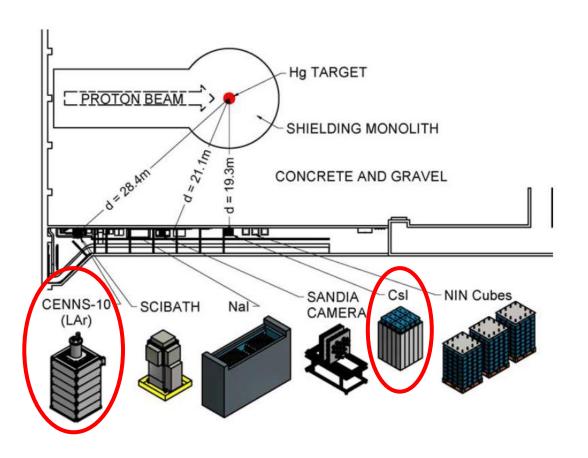
-Can we use the effective field theory approach to study nonstandard effects both in neutrino production and detection at COHERENT?

-Does COHERENT have a place in the electroweak EFT precision global fits?

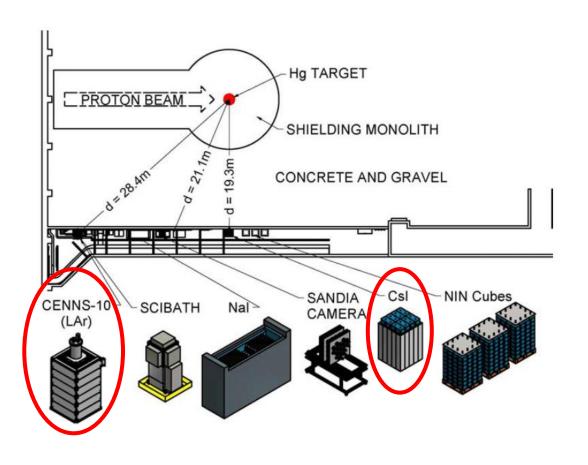
Two big questions

-Can we use the effective field theory approach to study non-standard effects both in neutrino production and detection at COHERENT? YES

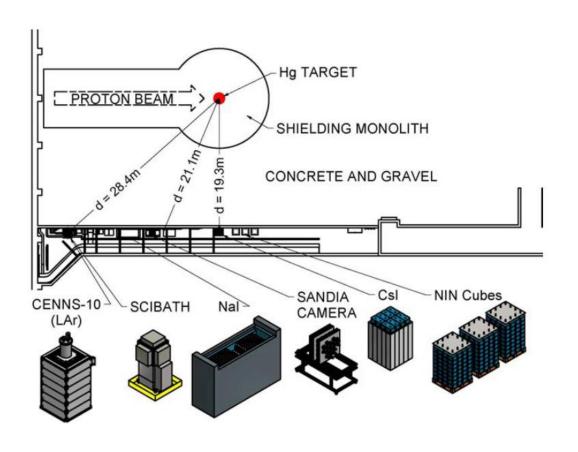
-Does COHERENT have a place in the electroweak EFT precision global fits? YES



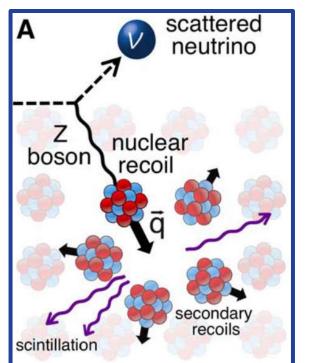
-The experiment consists of a set of detectors built around nuclear targets (CsI, Ar) exposed to neutrinos generated by the Spallation Neutron Source (SNS)



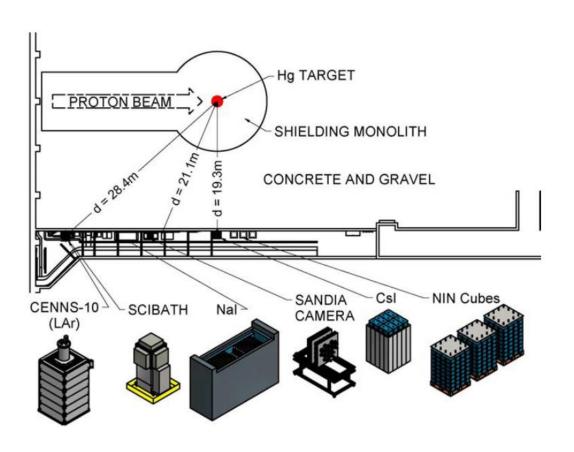
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- -Built to observe coherent elastic neutrino scattering off nuclei (CEvNS)



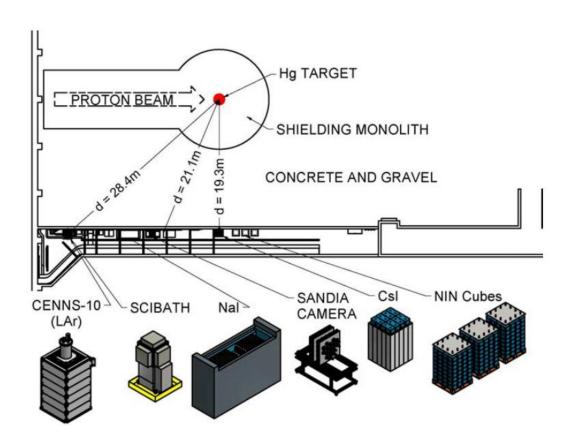
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- -Interaction enhanced for heavy nuclei
- -Excellent probe for NP at the neutrino sector

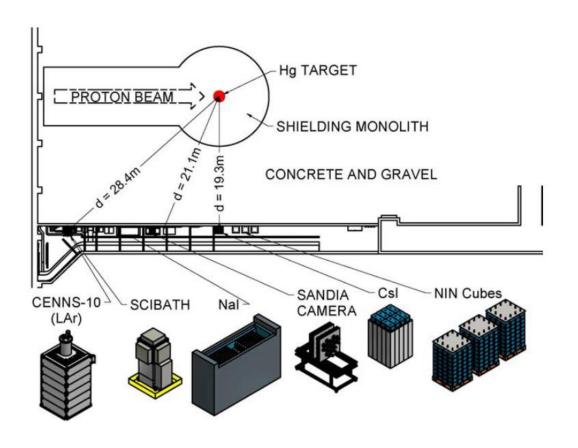


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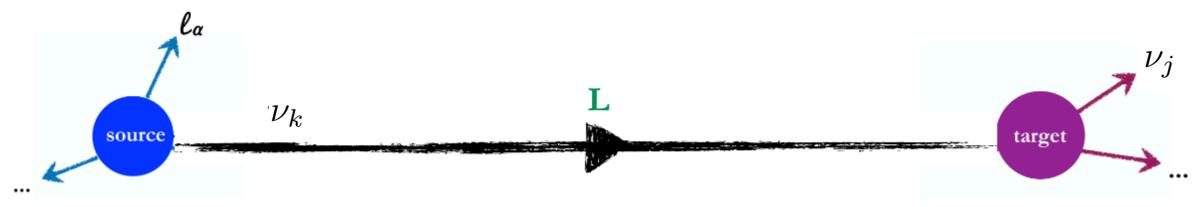


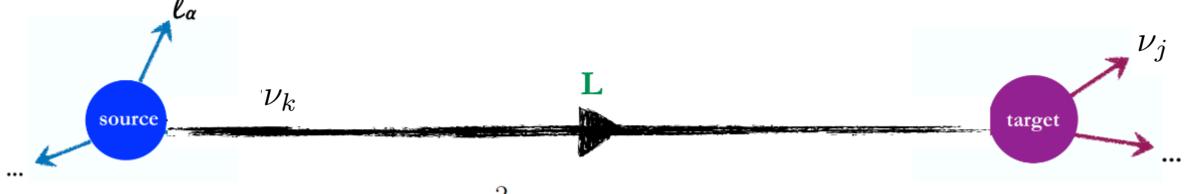
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- -First experiment to measure CEvNS and to describe its energy and time distributions [1708.01294,

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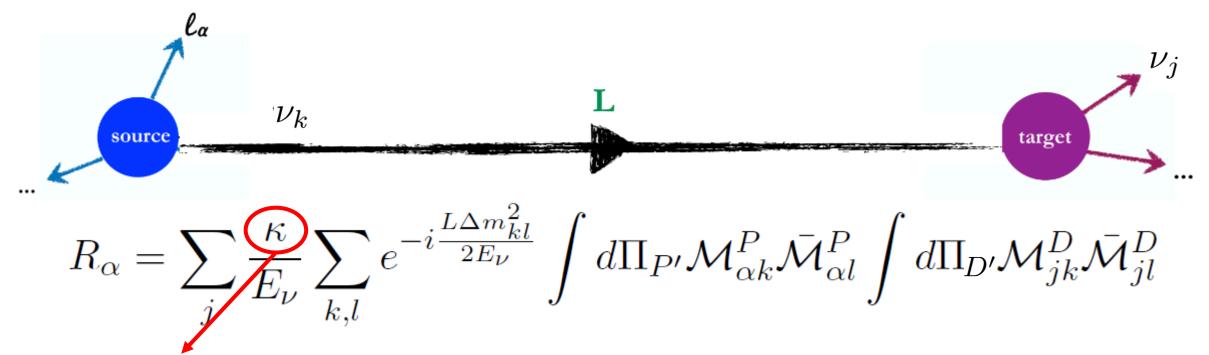


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- -First experiment to measure CEvNS and to describe its energy and time distributions [1708.01294, 2003.10630, 2110.07730]
- -How to introduce NP into this setup? We borrow from a recent theoretical description of neutrino oscillation observables in a QFT framework [Falkowski, González-Alonso, & Tabrizi, '20].



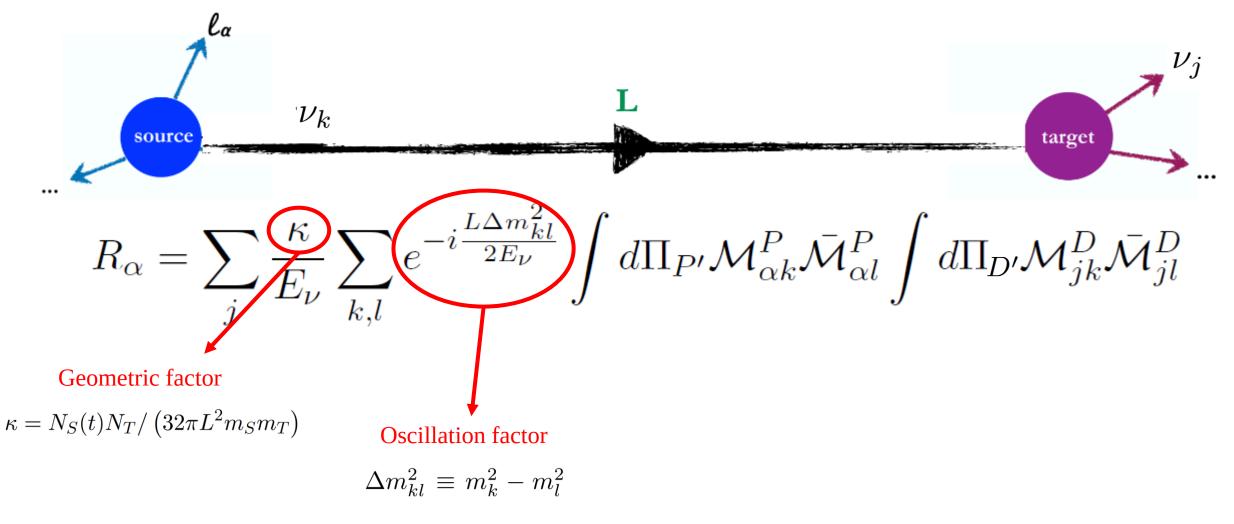


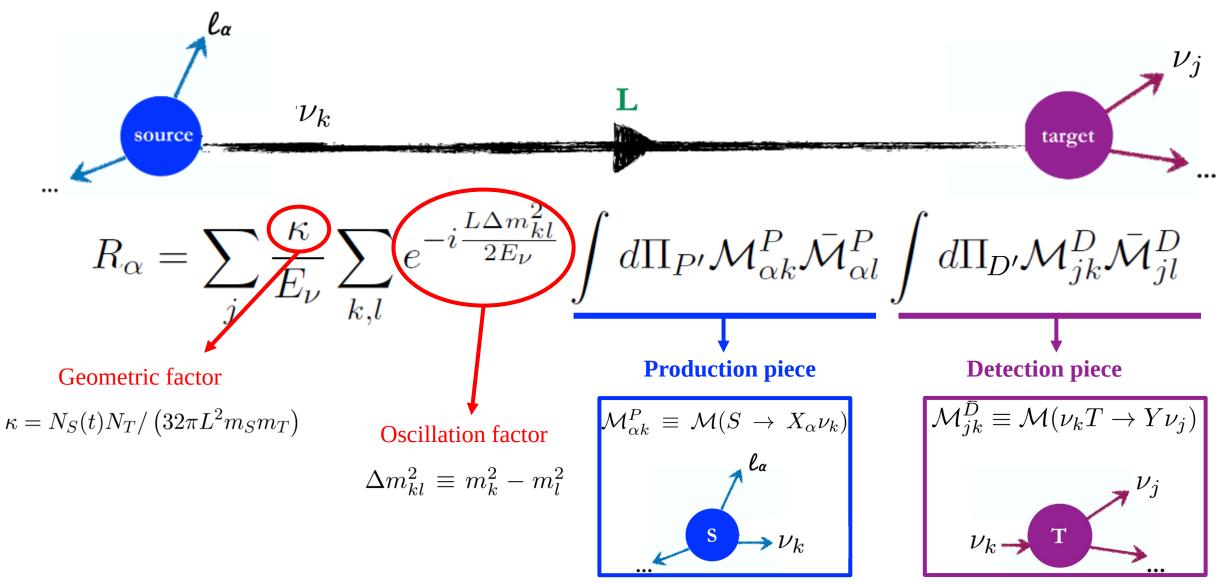
$$R_{\alpha} = \sum_{j} \frac{\kappa}{E_{\nu}} \sum_{k,l} e^{-i\frac{L\Delta m_{kl}^{2}}{2E_{\nu}}} \int d\Pi_{P'} \mathcal{M}_{\alpha k}^{P} \bar{\mathcal{M}}_{\alpha l}^{P} \int d\Pi_{D'} \mathcal{M}_{jk}^{D} \bar{\mathcal{M}}_{jl}^{D}$$

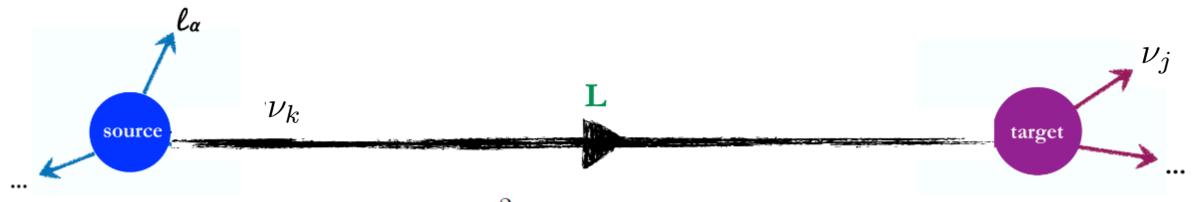


Geometric factor

$$\kappa = N_S(t)N_T / \left(32\pi L^2 m_S m_T\right)$$



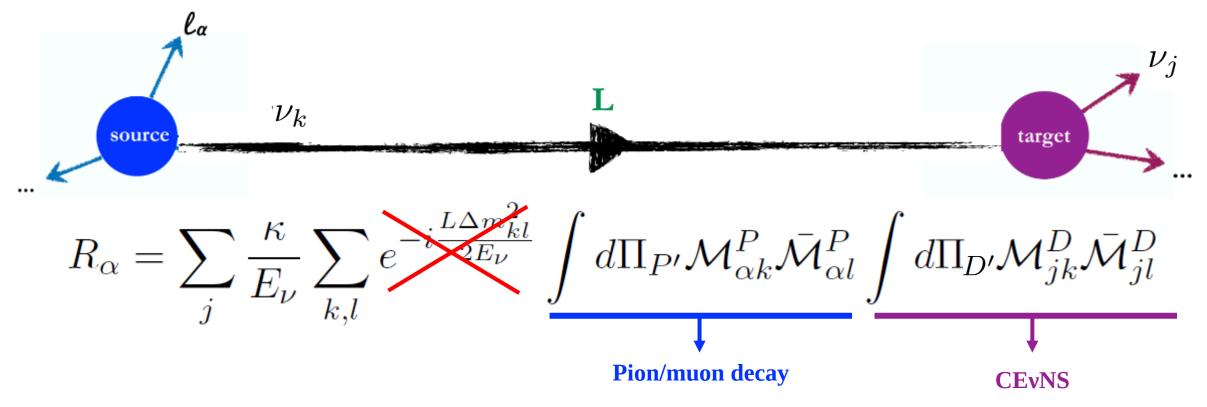


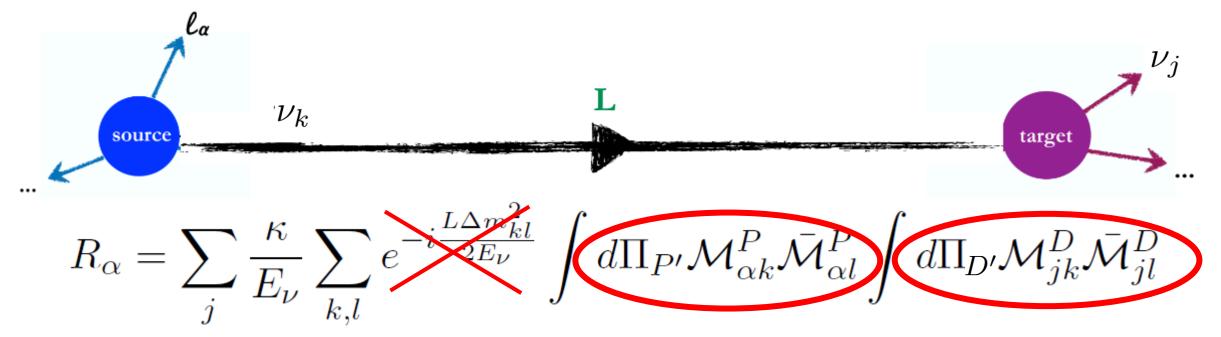


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Unique features when:

- -NP effects present both on production and detection
- -Flavor-violating couplings





NP effects easy to implement in the production/detection amplitudes!



What NP contributions enter the COHERENT observables?

-Pion decay production piece:

$$\mathcal{L}_{\text{WEFT}} \subset -\frac{2V_{ud}}{v^2} \Big\{ [1 + \epsilon_L^{ud}]_{\alpha\beta} \left(\bar{u}\gamma^{\mu} P_L d \right) \left(\bar{\ell}_{\alpha} \gamma_{\mu} P_L \nu_{\beta} \right) + [\epsilon_R^{ud}]_{\alpha\beta} \left(\bar{u}\gamma^{\mu} P_R d \right) \left(\bar{\ell}_{\alpha} \gamma_{\mu} P_L \nu_{\beta} \right) \\ -\frac{1}{2} \left[\epsilon_P^{ud} \right]_{\alpha\beta} \left(\bar{u}\gamma^5 d \right) \left(\bar{\ell}_{\alpha} P_L \nu_{\beta} \right) \Big\}$$

-Muon decay production piece:

$$\mathcal{L}_{\text{WEFT}} \subset -\frac{2}{v_0^2} \left\{ \left(\delta_{\alpha a} \delta_{\beta b} + [\rho_L]_{a\alpha\beta b} \right) \left(\bar{\ell}_a \gamma^{\mu} P_L \nu_{\alpha} \right) \left(\bar{\nu}_{\beta} \gamma_{\mu} P_L \ell_b \right) - 2 \left[\rho_R \right]_{a\alpha\beta b} \left(\bar{\ell}_a P_L \nu_{\alpha} \right) \left(\bar{\nu}_{\beta} P_R \ell_b \right) \right\}$$

-Detection piece:

$$\mathcal{L}_{\text{WEFT}} \subset -\frac{1}{v^2} \sum_{q=u,d} \left\{ \left[g_V^{qq} \, \mathbb{1} + \epsilon_V^{qq} \right]_{\alpha\beta} \left(\bar{q} \gamma^{\mu} q \right) \left(\bar{\nu}_{\alpha} \gamma_{\mu} P_L \nu_{\beta} \right) + \left[g_A^{qq} \, \mathbb{1} + \epsilon_A^{qq} \right]_{\alpha\beta} \left(\bar{q} \gamma^{\mu} \gamma^5 q \right) \left(\bar{\nu}_{\alpha} \gamma_{\mu} P_L \nu_{\beta} \right) \right\}$$

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NEW!
Indirect NP effects
from SM inputs

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-Main observable at COHERENT:

$$\frac{dN}{dt \, dT} = g_{\pi}(t) \, \frac{dN^{\text{prompt}}}{dT} + g_{\mu}(t) \, \frac{dN^{\text{delayed}}}{dT}$$

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Extracted from our knowledge about the pion and muon decay laws

$$N_S \rightarrow N_S(t)$$

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Calculated to include NP contributions, nuclear effects and experimental corrections

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-Linear marginalized limits → only (flavor diagonal) detection WCs remain:

$$\begin{pmatrix} 0.63 & -0.70 & -0.22 & 0.24 \\ 0.21 & -0.24 & 0.63 & -0.70 \\ -0.68 & -0.61 & 0.30 & 0.27 \\ 0.30 & 0.27 & 0.68 & 0.61 \end{pmatrix} \begin{pmatrix} \epsilon_{ee}^{dd} \\ \epsilon_{ee}^{uu} \\ \epsilon_{\mu\mu}^{dd} \\ \epsilon_{\mu\mu}^{uu} \end{pmatrix} = \begin{pmatrix} 2.0 \pm 5.7 \\ -0.2 \pm 1.7 \\ -0.037 \pm 0.042 \\ -0.004 \pm 0.013 \end{pmatrix}$$

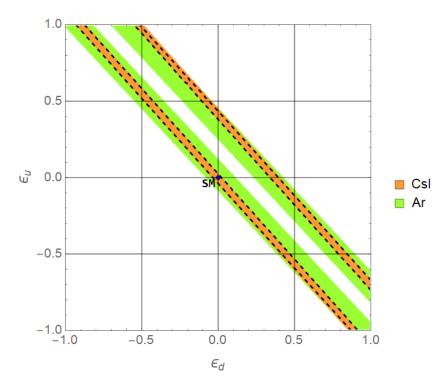
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$$\epsilon_{ee}^{uu}=\epsilon_{\mu\mu}^{uu}\equiv\epsilon_{u}$$
 (Lepton-flavor $\epsilon_{ee}^{dd}=\epsilon_{\mu\mu}^{dd}\equiv\epsilon_{d}$ universal limit)

$$0.67\epsilon_u + 0.74\epsilon_d = -0.002 \pm 0.010$$



-Flavor blind SMEFT matching $(U(3)^5)$ scheme):

 $0.67\epsilon_u + 0.74\epsilon_d = -0.002 \pm 0.010$

$$\epsilon_{u} = \delta g_{L}^{Zu} + \delta g_{R}^{Zu} + \left(1 - \frac{8s_{\theta}^{2}}{3}\right) \delta g_{L}^{Z\nu} - \frac{1}{2} \left(c_{lq}^{(1)} + c_{lq}^{(3)} + c_{lu}\right)$$

$$\epsilon_{d} = \delta g_{L}^{Zd} + \delta g_{R}^{Zd} - \left(1 - \frac{4s_{\theta}^{2}}{3}\right) \delta g_{L}^{Z\nu} - \frac{1}{2} \left(c_{lq}^{(1)} - c_{lq}^{(3)} + c_{ld}\right)$$

$$0.71c_{lq}^{(1)} - 0.04c_{lq}^{(3)} + 0.34c_{lu} + 0.37c_{ld} - 0.67(\delta g_L^{Zu} + \delta g_R^{Zu}) - 0.74(\delta g_L^{Zd} + \delta g_R^{Zd}) + 0.26\delta g_L^{Z\nu} = -0.003 \pm 0.010$$

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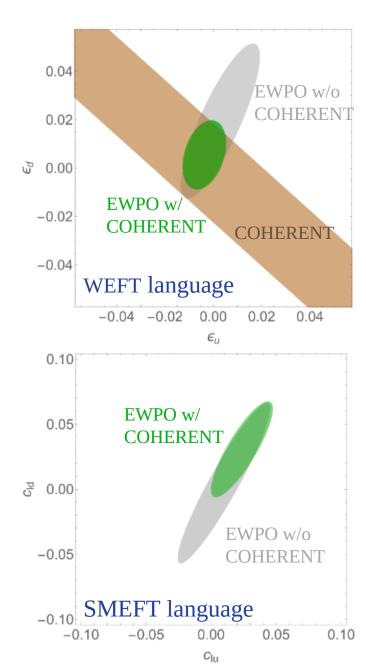
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How do these results fit into the EWPO global fit? (Z & W-pole data, atomic PV, tau decays...)

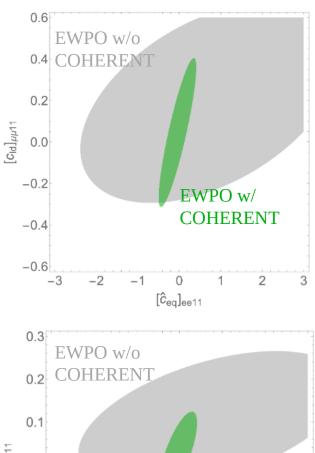
-Impact on the SMEFT EW global fit:

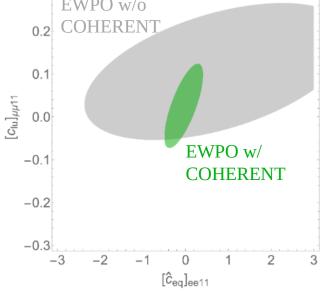
$$\begin{pmatrix} \delta g_L^{W\ell} \\ \delta g_L^{Ze} \\ \delta g_R^{Ze} \\ \delta g_R^{Zu} \\ \delta g_R^{Zu} \\ \delta g_L^{Zu} \\ \delta g_R^{Zd} \\ \delta g_L^{Zd} \\ \delta g_R^{Zd} \\ \delta g_L^{Zd} \\ \delta g_R^{Zd} \\ \delta g_R^{Zd} \\ \delta g_L^{Zd} \\ \delta g_R^{Zd} \\ c_{10} \\ c_{10} \\ c_{10} \\ c_{2d} \\ c_{2d} \\ c_{eq} \\ c_{eu} \\ c_{ed} \\ c_{11} \\ c_{eq} \\ c_{eu} \\ c_{ed} \\ c_{11} \\ c_{2d} \\$$



-Impact on the SMEFT EW global fit (general scheme):

$$\begin{pmatrix} [c_{lq}^{(3)}]_{ce11} \\ [\hat{c}_{eq}]_{ce11} \\ [\hat{c}_{lu}]_{ce11} \\ [\hat{c}_{lu}]_{ce11} \\ [\hat{c}_{eu}]_{ee11} \\ [\hat{c}_{eu}]_{ee11} \\ [\hat{c}_{eu}]_{ee11} \\ [\hat{c}_{eu}]_{ee11} \\ [\hat{c}_{eu}]_{ee11} \\ [\hat{c}_{leq}]_{ee11} \\ [c_{leq}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee11} \\ [c_{lequ}]_{ee22} \\ [c_{lu}]_{ee22} \\ [c_{lu}]_{ee23} \\ [c_{lu}]_{ee33} \\ [c_{lu}]_{ee$$



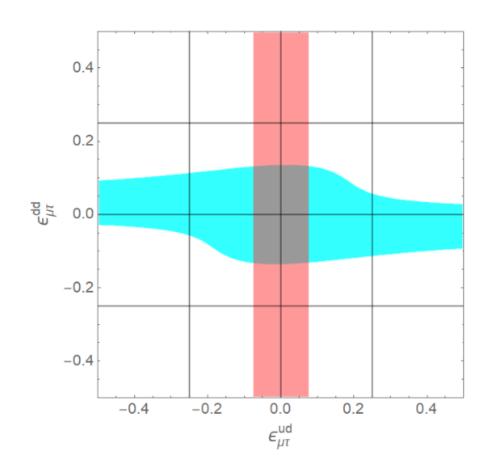


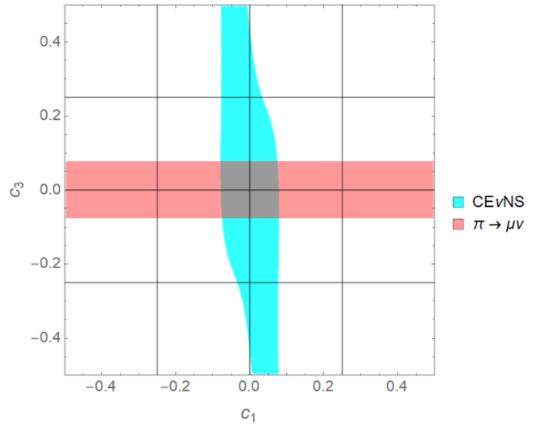
Relevant setup for models with explicit EW symmetry:

$$[\epsilon_L^{ud}]_{\mu\tau} = c_3 , \qquad \epsilon_{\mu\tau}^{uu} = c_1 - c_3 , \qquad \epsilon_{\mu\tau}^{dd} = c_1 + c_3$$

Results

-Production and detection WCs together:

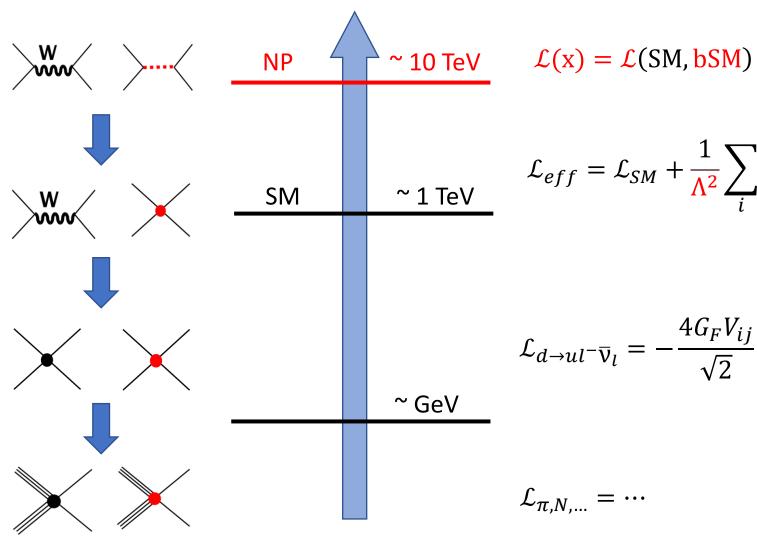




Conclusions

- We have successfully applied an EFT-based formalism for the description of **BSM physics at the COHERENT experiment**:
 - understand the UV meaning and limitations of the production/detection NSIs
 - take into account NP in production & detection
 - take into account NP affecting SM input
 - connect with specific NP models or interactions (e.g. leptoquarks)"
- We have quantitatively determined the impact of NP coming from production
- We have incorporated the COHERENT information into the SMEFT electroweak global fit

Extra: EFTs for low energy observables



$$\mathcal{L}(x) = \mathcal{L}(SM, bSM)$$

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{1}{\Lambda^2} \sum_i c_i O_i^{d=6}$$
 Standard Model EFT

[Buchmuller & Wyler'86, Leung et al.'86, Grzadkowksi et al., 10, Jenkins et al'13, ...]

$$\mathcal{L}_{d\to ul^{-}\overline{\nu}_{l}} = -\frac{4G_{F}V_{ij}}{\sqrt{2}} \left[\overline{l}_{L}\gamma_{\mu}\nu \cdot \overline{u}\gamma^{\mu}d_{l} + \sum_{\rho\delta\Gamma} \epsilon_{\rho\delta}^{\Gamma} \overline{l}_{\rho}\Gamma\nu \cdot \overline{u}\Gamma d_{\delta} \right]$$

Low energy EFT

[Cirigliano et al'09, Aebischer al.'15, Jenkins et al'18, ...]

$$C_{\pi,N,...} = \cdots$$

$$\frac{dN^{\text{prompt}}}{dT} = N_T \int dE_{\nu} \, \frac{d\Phi_{\nu_{\mu}}}{dE_{\nu}} \, \frac{d\tilde{\sigma}_{\nu_{\mu}}}{dT} ,$$

$$\frac{dN^{\text{delayed}}}{dT} = N_T \int dE_{\nu} \left(\frac{d\Phi_{\nu_{e}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\nu_{e}}}{dT} + \frac{d\Phi_{\bar{\nu}_{\mu}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\bar{\nu}_{\mu}}}{dT} \right)$$

$$\frac{d\tilde{\sigma}_f}{dT} = (m_{\mathcal{N}} + T) \frac{(\mathcal{F}(T))^2}{2v^4 \pi} \left(1 - \frac{(m_{\mathcal{N}} + 2E_{\nu}) T}{2E_{\nu}^2}\right) \tilde{Q}_f^2$$

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 $SM \rightarrow 1$ weak charge (per target nucleus) EFT $\rightarrow 3$ weak charges (per target nucleus)

$$\frac{d\tilde{\sigma}_f}{dT} = (m_{\mathcal{N}} + T) \frac{(\mathcal{F}(T))^2}{2v^4 \pi} \left(1 - \frac{(m_{\mathcal{N}} + 2E_{\nu}) T}{2E_{\nu}^2} \right) \tilde{Q}_f^2 \longrightarrow \tilde{\mathcal{Q}}_f^2 \equiv \mathcal{Q}_{\text{SM}}^2 + g_f \left(\epsilon_{CC}, \epsilon_{NC} \right)$$

-Generalized weak charges:

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\left[\mathcal{P}\mathcal{Q}^{2}\mathcal{P}^{\dagger}\right]_{\mu\mu}}{\left(\mathcal{P}\mathcal{P}^{\dagger}\right)_{\mu\mu}},$$

$$\tilde{Q}_{e}^{2} = \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{*}\mathcal{Q}^{2}P_{L}^{T} + \mathcal{P}_{R}^{T}\mathcal{Q}^{2}P_{R}^{*}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)},$$

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{T}\mathcal{Q}^{2}P_{L}^{*} + \mathcal{P}_{R}^{*}\mathcal{Q}^{2}P_{R}^{T}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)}$$

$$\left[\mathcal{P}_{\alpha\beta} \equiv \delta_{\alpha\beta} + [\epsilon_{L}]_{\alpha\beta} - [\epsilon_{R}]_{\alpha\beta} - [\epsilon_{R}$$

$$\frac{dN^{\text{prompt}}}{dT} = N_T \int dE_{\nu} \frac{d\Phi_{\nu_{\mu}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\nu_{\mu}}}{dT} ,$$

$$\frac{dN^{\text{delayed}}}{dT} = N_T \int dE_{\nu} \left(\frac{d\Phi_{\nu_e}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\nu_e}}{dT} + \frac{d\Phi_{\bar{\nu}_{\mu}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\bar{\nu}_{\mu}}}{dT} \right)$$

 $SM \rightarrow 1$ weak charge (per target nucleus) EFT $\rightarrow 3$ weak charges (per target nucleus)

$$\frac{d\tilde{\sigma}_f}{dT} = (m_{\mathcal{N}} + T) \frac{(\mathcal{F}(T))^2}{2v^4 \pi} \left(1 - \frac{(m_{\mathcal{N}} + 2E_{\nu}) T}{2E_{\nu}^2} \right) \tilde{Q}_f^2 \longrightarrow \tilde{\mathcal{Q}}_f^2 \equiv \mathcal{Q}_{\text{SM}}^2 + g_f \left(\epsilon_{CC}, \epsilon_{NC} \right)$$

-Generalized weak charges:

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\left[\mathcal{P}\mathcal{Q}^{2}\mathcal{P}^{\dagger}\right]_{\mu\mu}}{\left(\mathcal{P}\mathcal{P}^{\dagger}\right)_{\mu\mu}},$$

$$\tilde{Q}_{e}^{2} = \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{*}\mathcal{Q}^{2}P_{L}^{T} + \mathcal{P}_{R}^{T}\mathcal{Q}^{2}P_{R}^{*}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)},$$

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{T}\mathcal{Q}^{2}P_{L}^{*} + \mathcal{P}_{R}^{*}\mathcal{Q}^{2}P_{R}^{T}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)}$$

$$\left[\mathcal{P}_{\alpha\beta} \equiv \delta_{\alpha\beta} + [\epsilon_{L}]_{\alpha\beta} - [\epsilon_{R}]_{\alpha\beta} - [\epsilon_{R}$$

$$\frac{dN^{\text{prompt}}}{dT} = N_T \int dE_{\nu} \frac{d\Phi_{\nu_{\mu}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\nu_{\mu}}}{dT},$$

$$\frac{dN^{\text{delayed}}}{dT} = N_T \int dE_{\nu} \left(\frac{d\Phi_{\nu_e}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\nu_e}}{dT} + \frac{d\Phi_{\bar{\nu}_{\mu}}}{dE_{\nu}} \frac{d\tilde{\sigma}_{\bar{\nu}_{\mu}}}{dT} \right)$$

$$\begin{split} g_{\alpha\beta}^{\nu p} &= 2 \, \left[\left(2 \, g_V^{uu} + g_V^{dd} \right) \mathbb{1} + \left(2 \, \epsilon^{uu} + \epsilon^{dd} \right) \right]_{\alpha\beta} \\ g_{\alpha\beta}^{\nu n} &= 2 \, \left[\left(g_V^{uu} + 2 g_V^{dd} \right) \mathbb{1} + \left(\epsilon^{uu} + 2 \epsilon^{dd} \right) \right]_{\alpha\beta} \end{split}$$

 $SM \rightarrow 1$ weak charge (per target nucleus) EFT \rightarrow 3 weak charges (per target nucleus)

$$\frac{d\tilde{\sigma}_f}{dT} = (m_{\mathcal{N}} + T) \frac{(\mathcal{F}(T))^2}{2v^4 \pi} \left(1 - \frac{(m_{\mathcal{N}} + 2E_{\nu})T}{2E_{\nu}^2}\right) \tilde{Q}_f^2 -$$

$$\longrightarrow \tilde{\mathcal{Q}}_f^2 \equiv \mathcal{Q}_{\mathrm{SM}}^2 + g_f \left(\epsilon_{CC}, \epsilon_{NC} \right)$$

-Generalized weak charges:

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\left[\mathcal{P}\mathcal{Q}^{2}\mathcal{P}^{\dagger}\right]_{\mu\mu}}{\left(\mathcal{P}\mathcal{P}^{\dagger}\right)_{\mu\mu}},$$

$$\tilde{Q}_{e}^{2} = \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{*}\mathcal{Q}^{2}P_{L}^{T} + \mathcal{P}_{R}^{T}\mathcal{Q}^{2}P_{R}^{*}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)},$$

$$\tilde{Q}_{\mu}^{2} \equiv \frac{\operatorname{Tr}\left(\mathcal{P}_{L}^{T}\mathcal{Q}^{2}P_{L}^{*} + \mathcal{P}_{R}^{*}\mathcal{Q}^{2}P_{R}^{T}\right)}{\operatorname{Tr}\left(\mathcal{P}_{L}\mathcal{P}_{L}^{\dagger} + \mathcal{P}_{R}\mathcal{P}_{R}^{\dagger}\right)}$$

$$\left[\mathcal{P}_{\alpha\beta} \equiv \delta_{\alpha\beta} + [\epsilon_{L}]_{\alpha\beta} - [\epsilon_{R}]_{\alpha\beta} - [\epsilon_{R}$$

$$[\mathcal{P}]_{lphaeta} \equiv \delta_{lphaeta} + [\epsilon_L]_{lphaeta} - [\epsilon_R]_{lphaeta} - [\epsilon_P]_{lphaeta} rac{m_{\pi^{\pm}}^2}{m_{\ell_{lpha}}(m_u + m_d)} \; , \ [\mathcal{P}_L]_{lphaeta} \equiv \delta_{lpha\mu}\delta_{eta e} + [
ho_L]_{\mulphaeta e} \; , \ [\mathcal{P}_R]_{lphaeta} \equiv [
ho_R]_{\mulphaeta e} \; . \quad [\mathcal{Q}]_{lphaeta} = Zq_{lphaeta}^{
u p} + (A - Z)q_{lphaeta}^{
u n} \; .$$