

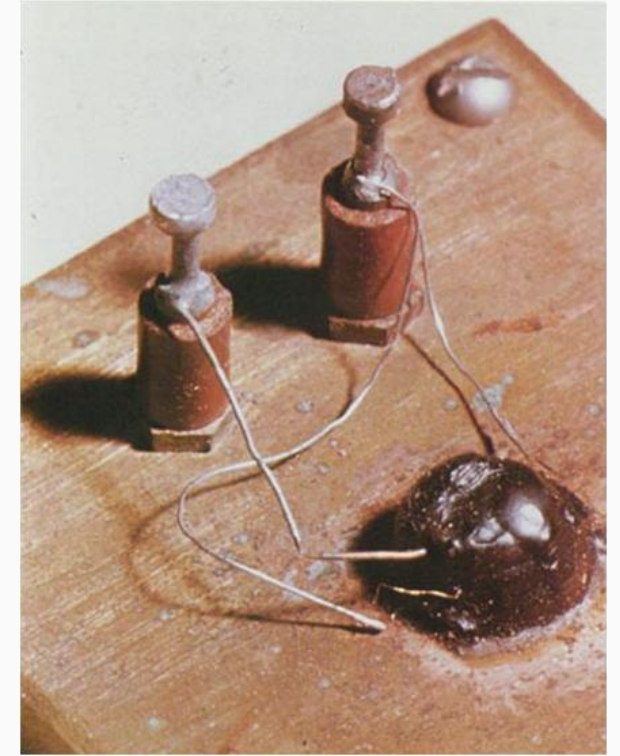
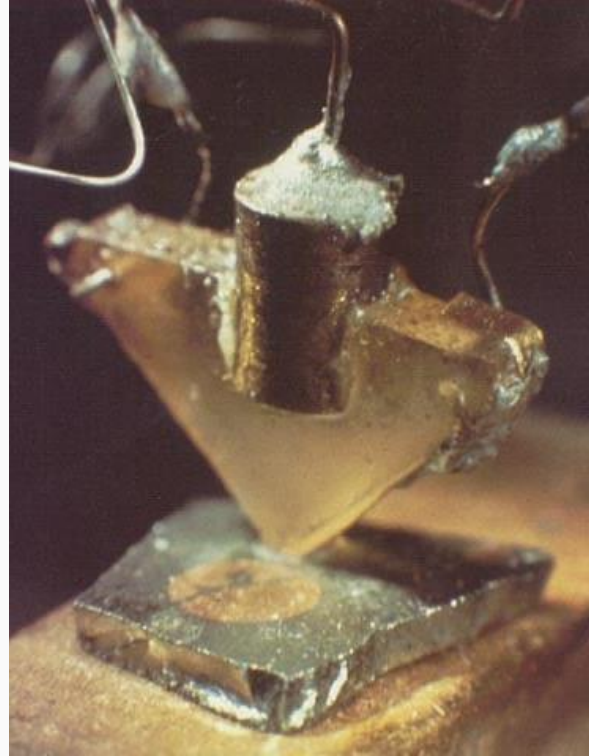
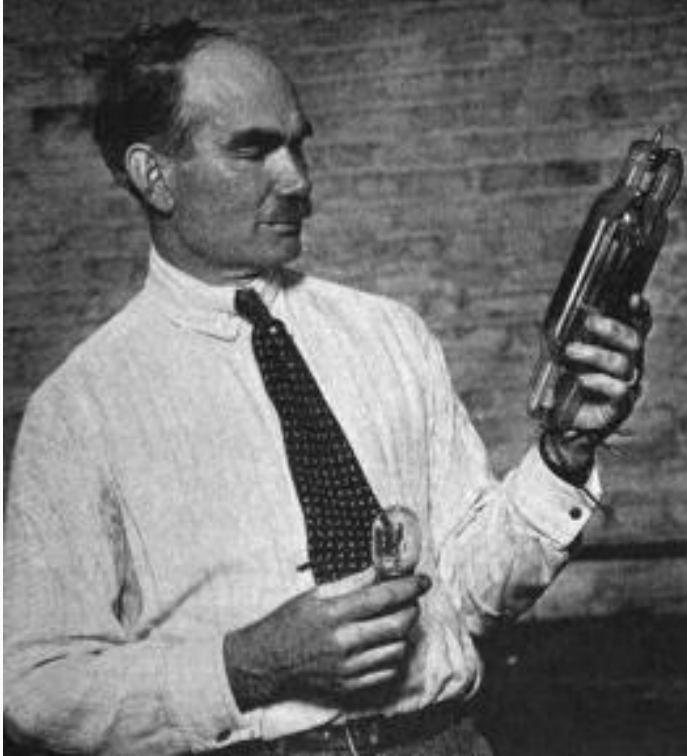
Quantum Technologies for Neutrino Mass



Superconducting Parametric Amplifiers

Songyuan Zhao

History of amplifiers

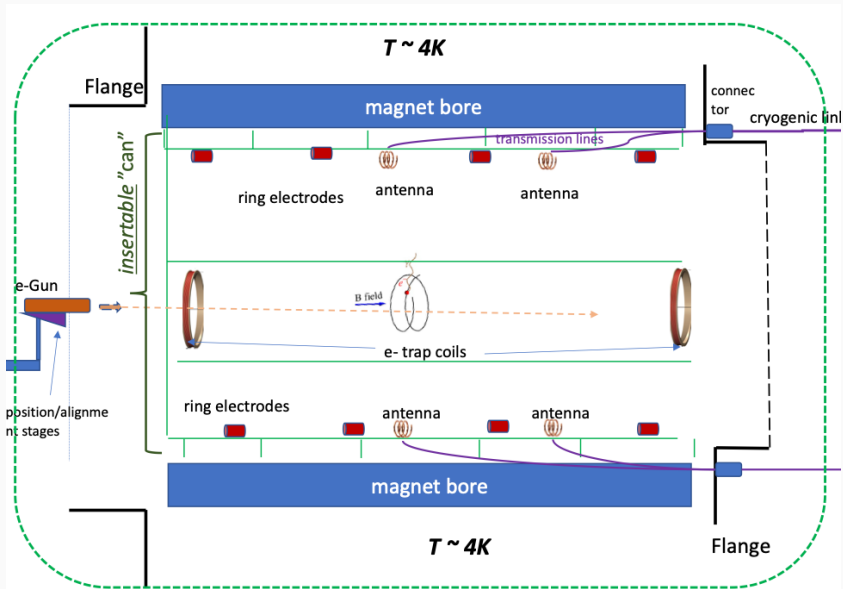


Vacuum tube / valve amplifier
(since 1906)

Transistor / solid state amplifier
(since 1950s)

Amplifiers in science

Experiment



cold environment
small noise
small signal

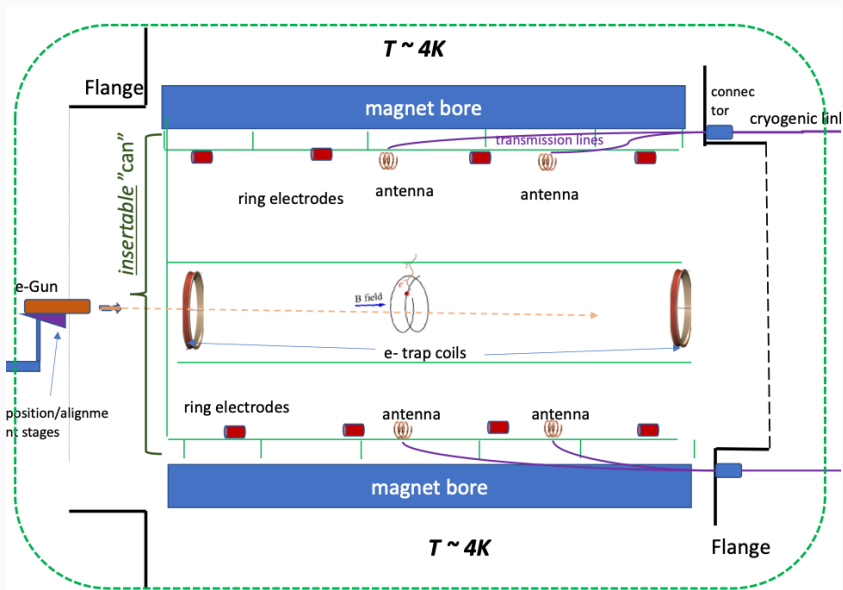
Data



warm environment
big noise
big signal

Amplifiers in science

Experiment



cold environment
small noise
small signal

Amplifier

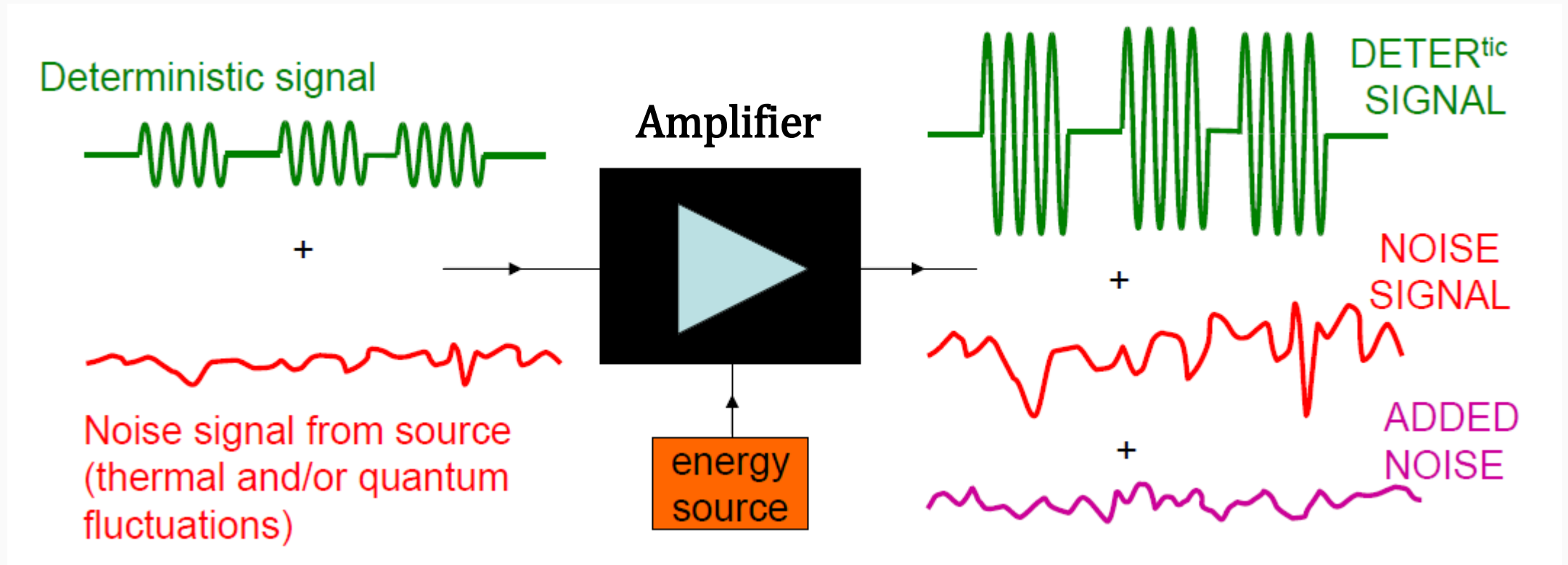


Data



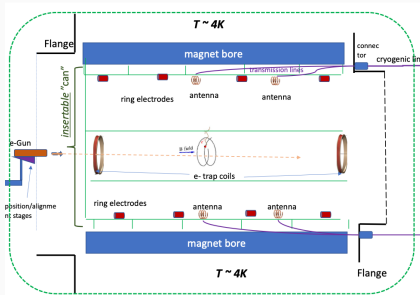
warm environment
big noise
big signal

Amplifiers in science



Example

Experiment



Noise: 3 K

Data

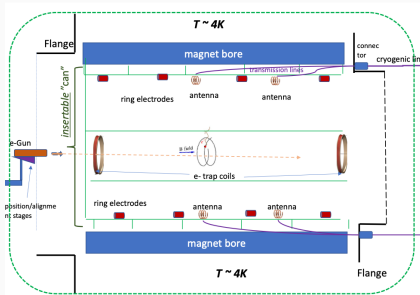


Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3

Example

Experiment



Noise: 3 K

Data



Noise: 300 K

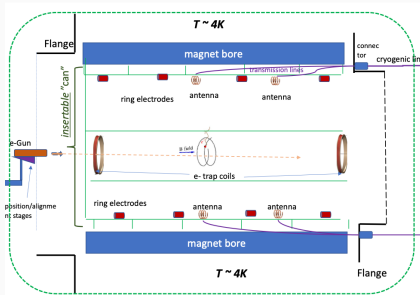
Signal: 10 K
Noise: 3 K
SNR: 3.3



Signal: 10 K
Noise: 303 K
SNR: 0.03

Example

Experiment



Noise: 3 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data

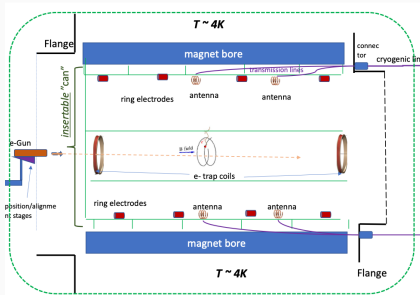


Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3

Example

Experiment



Noise: 3 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data



Noise: 300 K

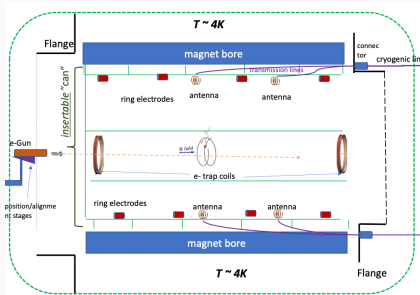
Signal: 10 K
Noise: 3 K
SNR: 3.3



Signal: 1000 K
Noise: 600 K
SNR: 1.7

Example

Experiment



Noise: 3 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data



Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3



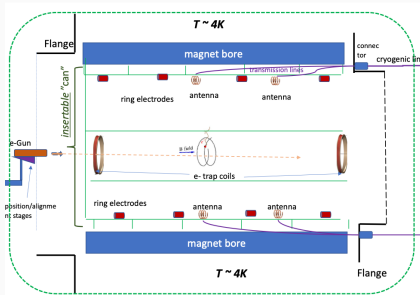
Signal: 1000 K
Noise: 600 K
SNR: 1.7



Signal: 1000 K
Noise: 900 K
SNR: 1.1

Example

Experiment



Noise: 3 K

Parametric Amplifier



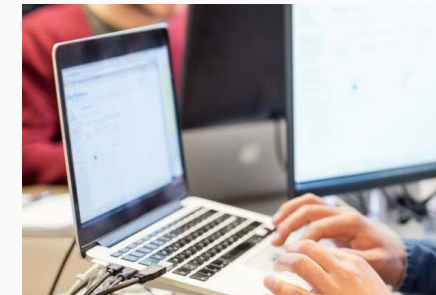
Gain: 10 (or 10 dB)
Noise: 1 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data

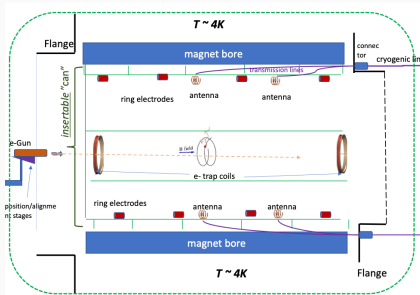


Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3

Example

Experiment



Noise: 3 K

Parametric Amplifier



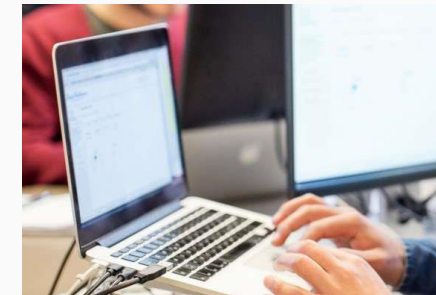
Gain: 10 (or 10 dB)
Noise: 1 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data



Noise: 300 K

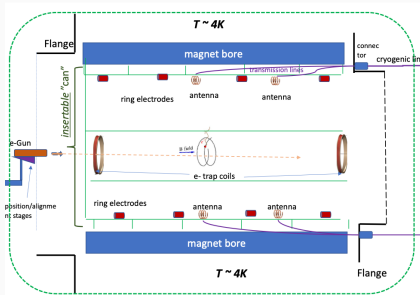
Signal: 10 K
Noise: 3 K
SNR: 3.3



Signal: 100 K
Noise: 40 K
SNR: 2.5

Example

Experiment



Noise: 3 K

Parametric Amplifier



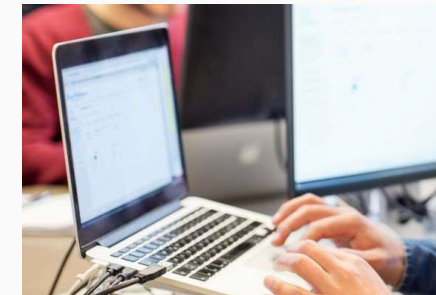
Gain: 10 (or 10 dB)
Noise: 1 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data



Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3



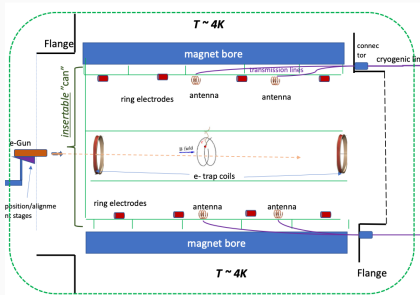
Signal: 100 K
Noise: 40 K
SNR: 2.5



Signal: 10000 K
Noise: 4300 K
SNR: 2.3

Example

Experiment



Noise: 3 K

Parametric Amplifier



Gain: 10 (or 10 dB)
Noise: 1 K

HEMT Amplifier



Gain: 100 (or 20 dB)
Noise: 3 K

Data



Noise: 300 K

Signal: 10 K
Noise: 3 K
SNR: 3.3



Signal: 100 K
Noise: 40 K
SNR: 2.5



Signal: 10000 K
Noise: 4300 K
SNR: 2.3



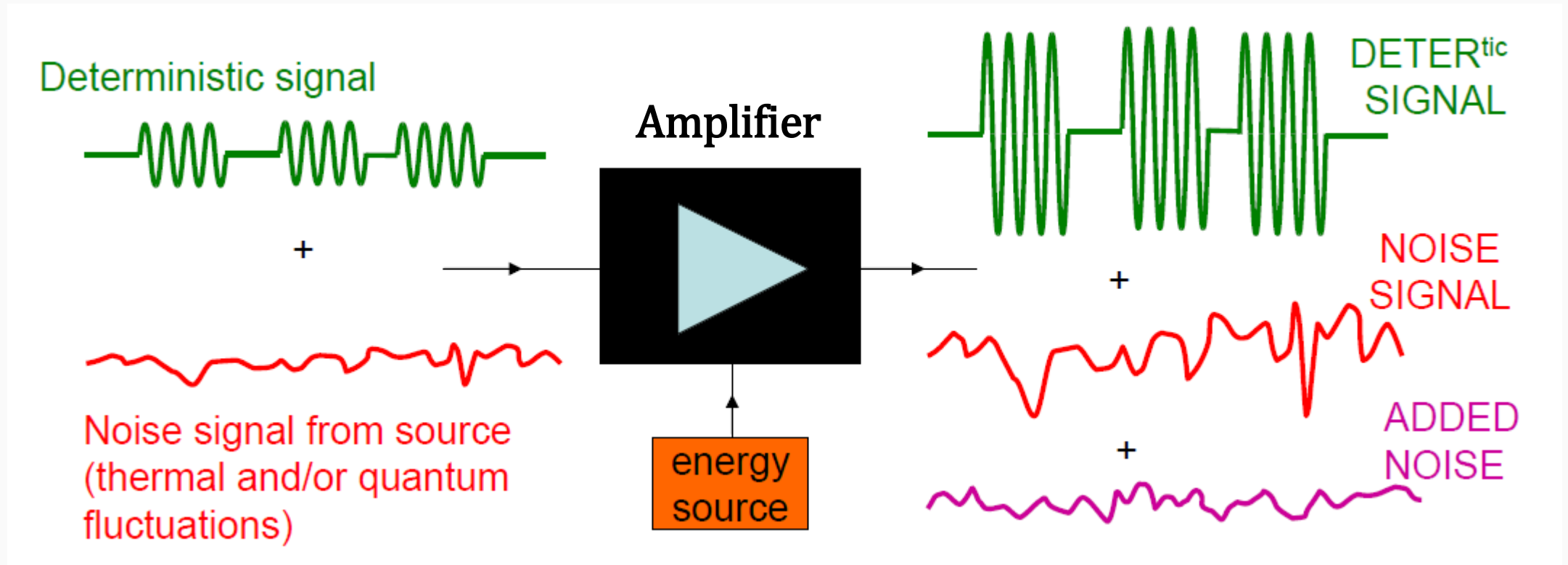
Signal: 10000 K
Noise: 4600 K
SNR: 2.2

Standard Quantum Limit

$$\hat{a}_{\text{out}} = g \hat{a}_{\text{in}} + \hat{N}$$

$$\frac{\langle |\Delta \hat{a}_{\text{out}}|^2 \rangle}{g^2} \geq \langle |\Delta \hat{a}_{\text{in}}|^2 \rangle + \frac{1}{2}$$

Standard Quantum Limit



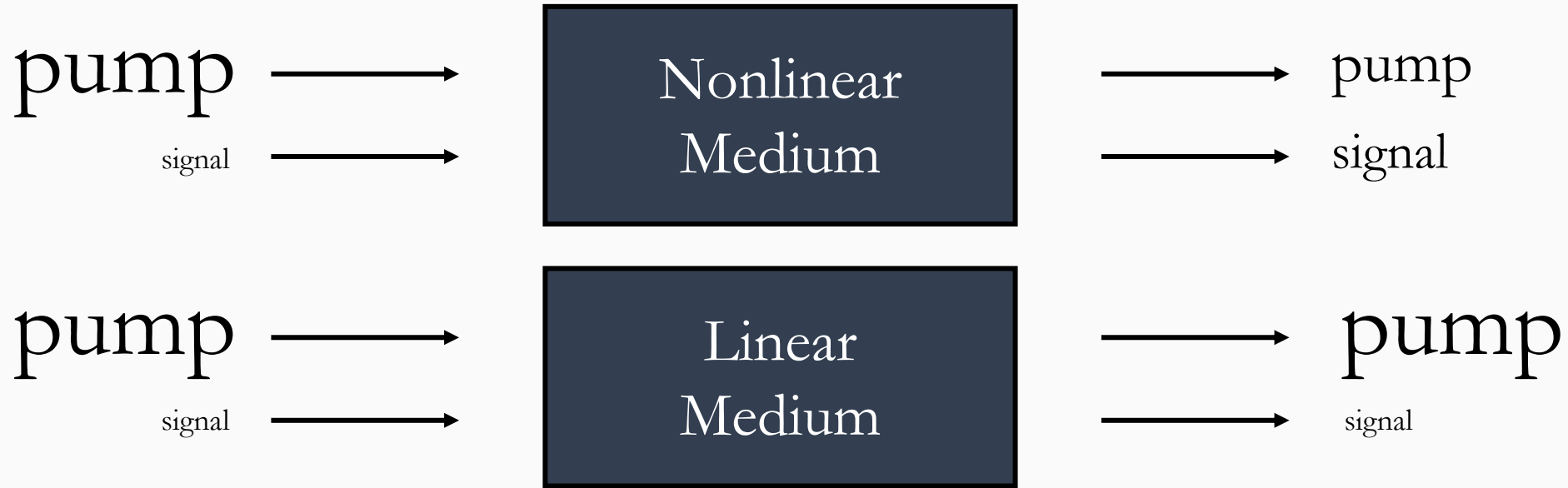
Nonlinearity



The Superposition Principle:

*For all **linear** systems, the net response caused by two or more stimuli is the sum of the responses that would have been caused by each stimulus **individually**.*

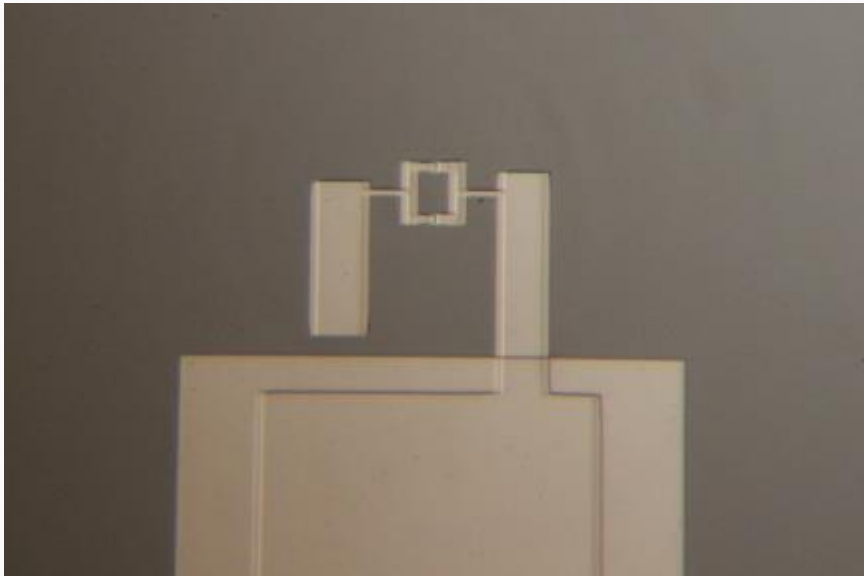
Nonlinearity



The Superposition Principle:

*For all **linear** systems, the net response caused by two or more stimuli is the sum of the responses that would have been caused by each stimulus **individually**.*

Types of nonlinearities



Josephson Junction Inductance

$$L = L_0 (\cos \phi)^{-1}$$
$$I = I_c \sin \phi$$



Kinetic Inductance of Superconductors

$$L \sim L_0 \left(1 + \frac{I^2}{I_*^2} \right)$$

Nonlinear kinetic inductance

$$L \sim L_0 \left(1 + \frac{I^2}{I_*^2} \right)$$

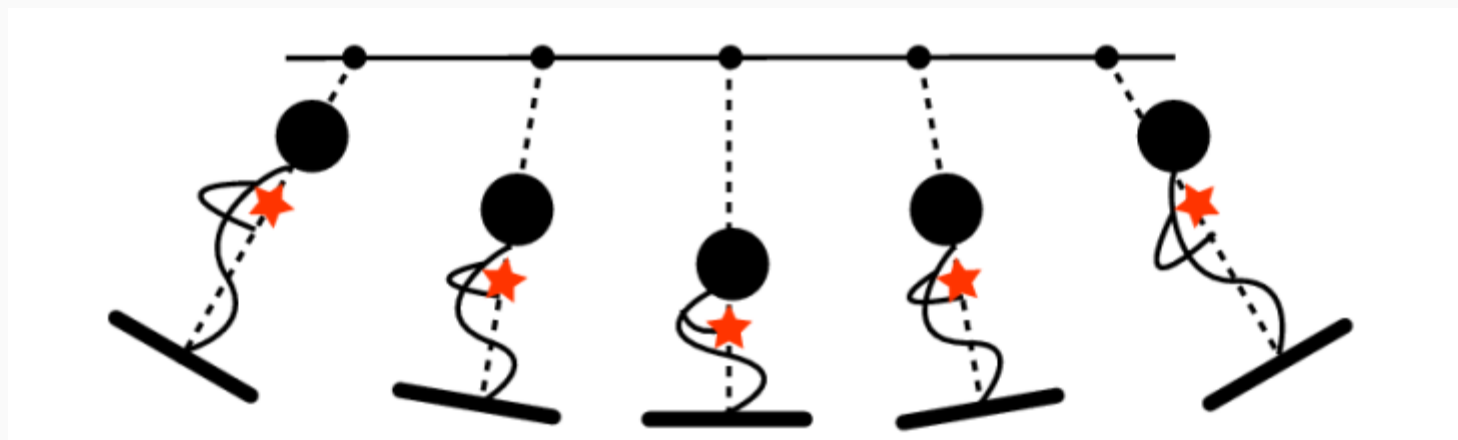
- Current results in a change in Cooper pair density
 - Phonon-mediated pair breaking
 - Superfluid velocity suppresses equilibrium pair density

$$F = F_0 + \alpha |\psi|^2 + \frac{\beta}{2} |\psi|^4 + \frac{1}{2m} |(-i\hbar\nabla - q\mathbf{A})\psi|^2 + \frac{|B|^2}{2\mu_0}$$

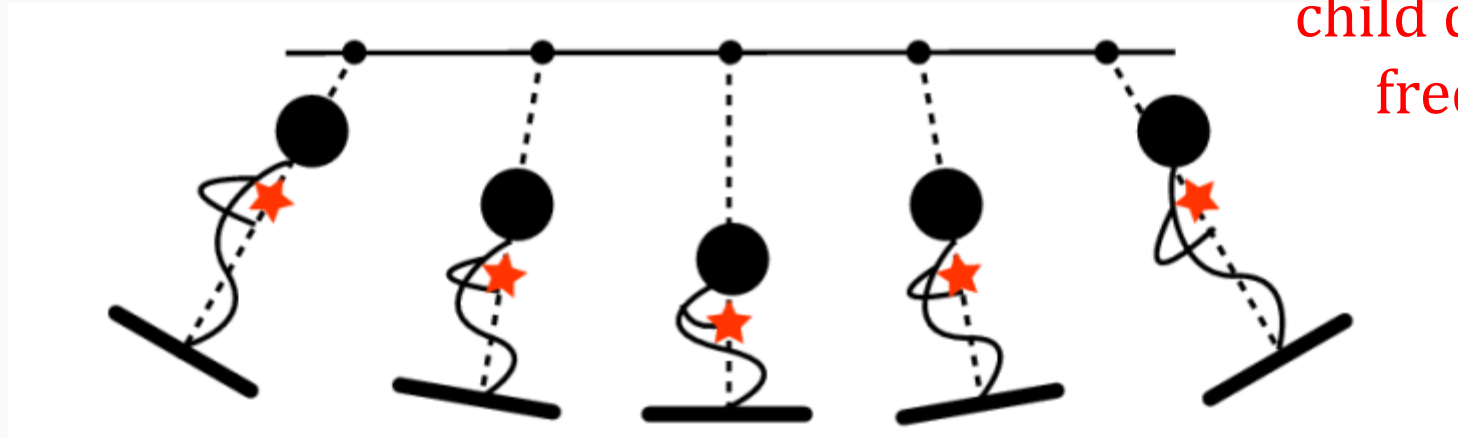
$$|\psi|^2 = \frac{1}{\beta} \left[\alpha_0 (T_c - T) - \frac{1}{2} m \mathbf{v}_s^2 \right]$$

- Other unknown mechanisms...

Amplification

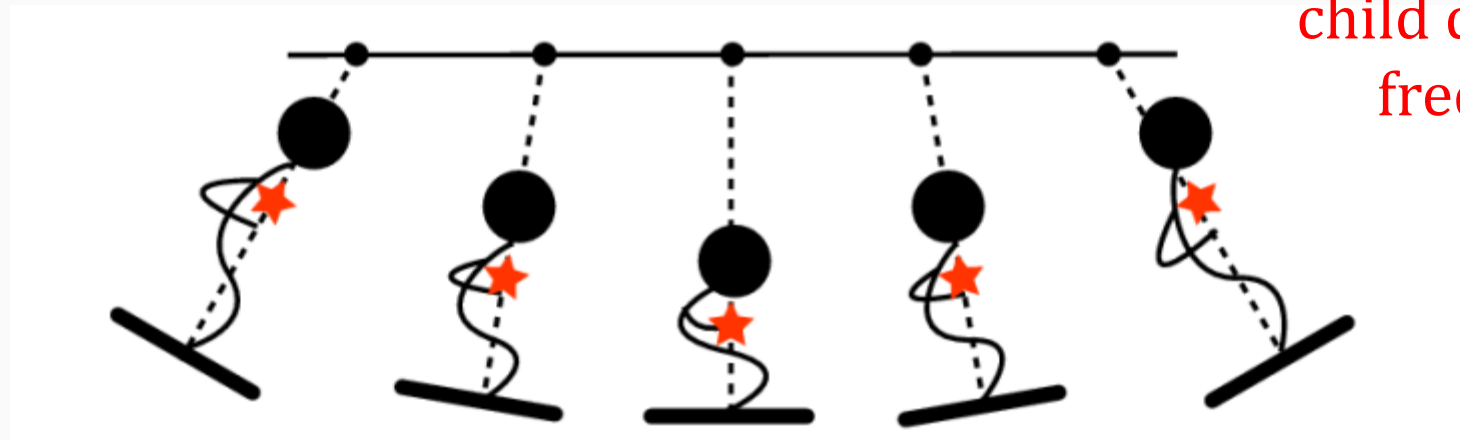


Amplification



child drives swing at
frequency $2\omega_0$

Amplification



child drives swing at
frequency $2\omega_0$

swing amplified at frequency ω_0

Amplification

- Four wave mixing

$$L \sim L_0 \left(1 + \frac{I^2}{I_*^2} \right)$$

- Important frequencies:

- Strong pump at ω_p
- Signal of interest at ω_s
- By-product idler at $\omega_i = 2\omega_p - \omega_s$

Amplification

$$L \sim L_0 \left(1 + \frac{I^2}{I_*^2} \right)$$

- Heisenberg picture

$$i\hbar \frac{d\hat{A}}{dt} = [\hat{A}, \hat{H}]$$

- Transmission line Hamiltonian

$$\hat{H} \propto \frac{1}{2} C \hat{V}^2 + \frac{1}{2} L \hat{I}^2$$

- Expansion into Fourier modes

$$\hat{I} \propto \sum_m (\hat{a}_m + \hat{a}_m^\dagger)$$

$$[\hat{a}_m, \hat{a}_n^\dagger] = \delta_{m,n}$$

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and Hermitian conjugate

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$$\frac{d\hat{a}_i}{dt} \propto \hat{a}_p \hat{a}_p \hat{a}_s^\dagger$$

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$$i\hbar \frac{d\hat{A}}{dt} = [\hat{A}, \hat{H}]$$

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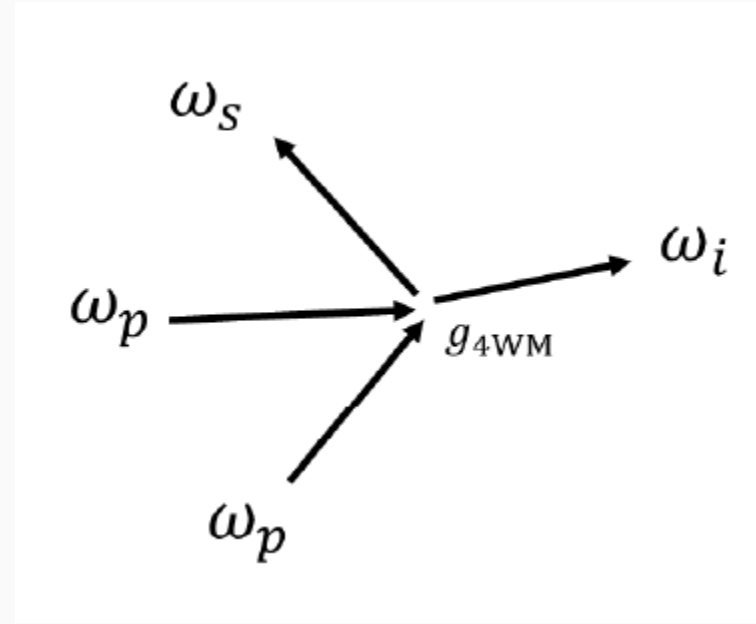
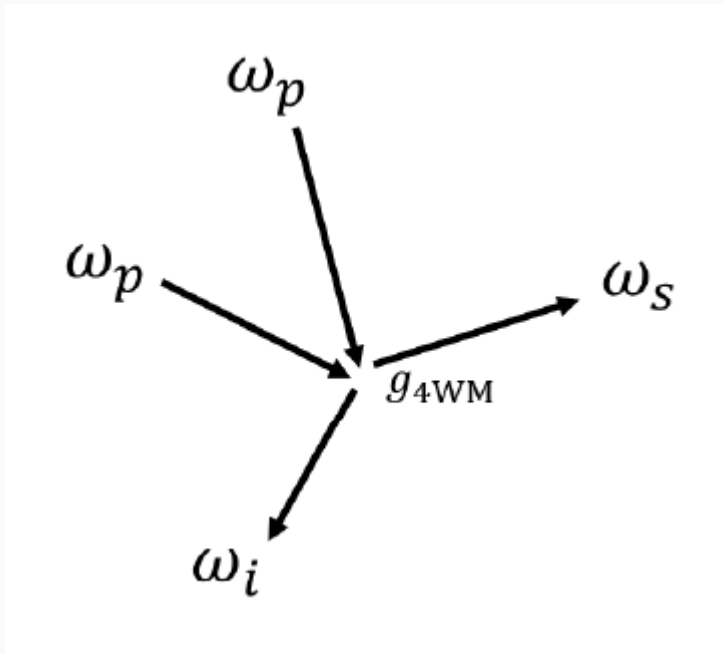
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- Strong pump \hat{a}_p can be approximated by $a_{p,0}$

$$\hat{a}_s = \hat{a}_{s,0} \cosh(kt) - i \hat{a}_{i,0}^\dagger \sinh(kt)$$

Amplification

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- Strong pump \hat{a}_p can be approximated by $a_{p,0}$

$$\hat{a}_s = \hat{a}_{s,0} \cosh(kt) - i \hat{a}_{i,0}^\dagger \sinh(kt)$$

- Set $\langle \hat{a}_i \rangle = 0$ as the starting condition

$$\langle \hat{a}_s \rangle = \langle \hat{a}_{s,0} \rangle \cosh(kt)$$

Squeezing amplification

- Introduce DC to current $I \rightarrow I_{\text{DC}} + I_{\text{RF}}$

$$L_0 \left(1 + \frac{I^2}{I_*^2} \right) \rightarrow L_0 \left(1 + \frac{I_{\text{DC}}^2}{I_*^2} + 2 \frac{I_{\text{DC}} I_{\text{RF}}}{I_*^2} + \frac{I_{\text{RF}}^2}{I_*^2} \right)$$

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- If $I_{\text{DC}} \gg I_{\text{RF}}$, the dominant nonlinear term is

$$L_0 \left(2 \frac{I_{\text{DC}} I_{\text{RF}}}{I_*^2} \right)$$

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Three wave mixing 

Squeezing amplification

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- If $I_{\text{DC}} \gg I_{\text{RF}}$, the dominant nonlinear term is

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Three wave mixing

- For squeezing
 $\omega_s = \omega_i = \omega_p/2$

Squeezing amplification

- Hamiltonian contains $\hat{a}_p \hat{a}_s^\dagger \hat{a}_s^\dagger$ and $\hat{a}_p^\dagger \hat{a}_s \hat{a}_s$

Squeezing amplification

▪ Hamiltonian contains $\hat{a}_p \hat{a}_s^\dagger \hat{a}_s^\dagger$ and $\hat{a}_p^\dagger \hat{a}_s \hat{a}_s$

▪ Strong pump \hat{a}_p can be approximated by $a_{p,0}$

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Squeezing amplification

- Hamiltonian contains $\hat{a}_p \hat{a}_s^\dagger \hat{a}_s^\dagger$ and $\hat{a}_p^\dagger \hat{a}_s \hat{a}_s$

- Strong pump \hat{a}_p can be approximated by $a_{p,0}$

$$\hat{a}_s = \hat{a}_{s,0} \cosh(kt) - i \hat{a}_{s,0}^\dagger \sinh(kt)$$

- Two quadrature of signal

$$\hat{X}_s = \hat{a}_s + \hat{a}_s^\dagger = \exp(kt) \hat{X}_{s,0}$$

$$\hat{Y}_s = \hat{a}_s - \hat{a}_s^\dagger = \exp(-kt) \hat{Y}_{s,0}$$

Squeezing amplification

- Hamiltonian contains $\hat{a}_p \hat{a}_s^\dagger \hat{a}_s^\dagger$ and $\hat{a}_p^\dagger \hat{a}_s \hat{a}_s$

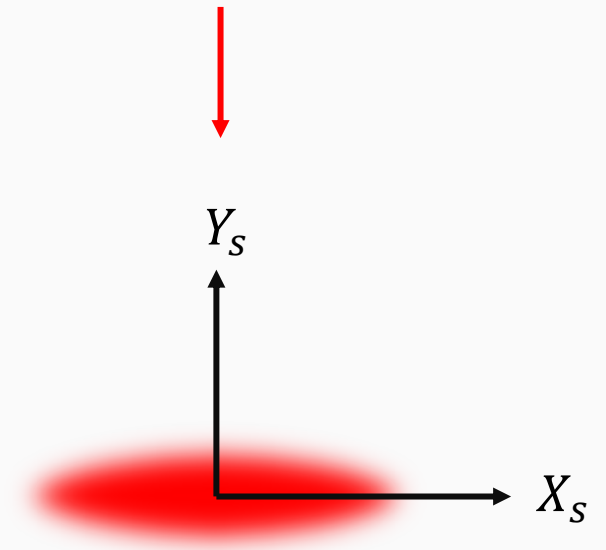
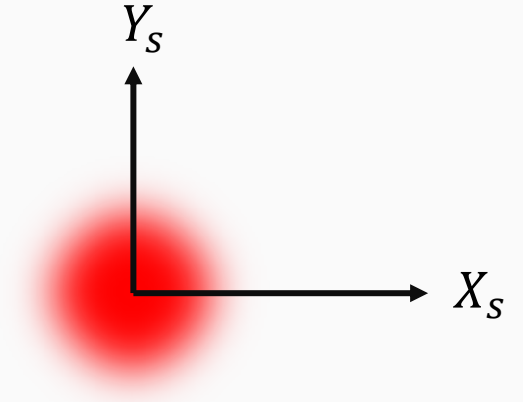
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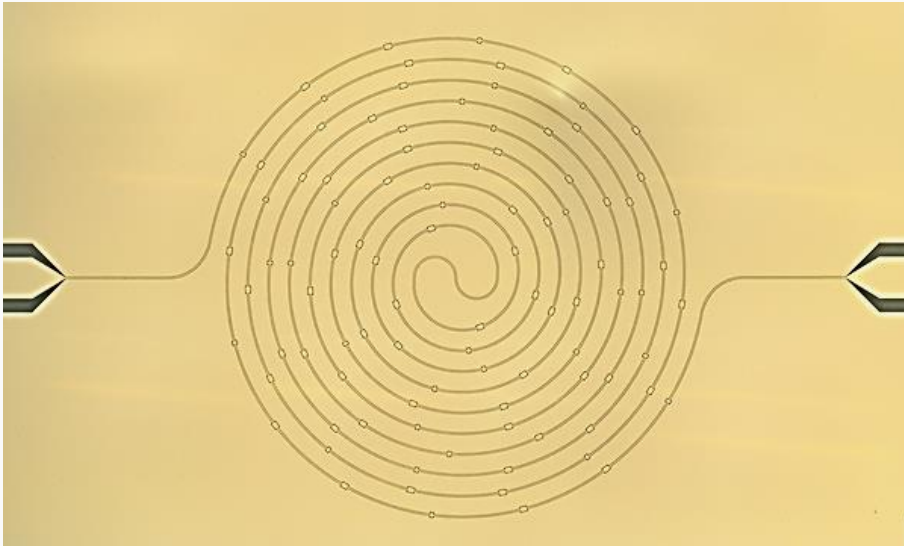
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Types of SPAs geometries



Travelling-Wave Parametric Amplifier
(TWPA)



Resonator Parametric Amplifier

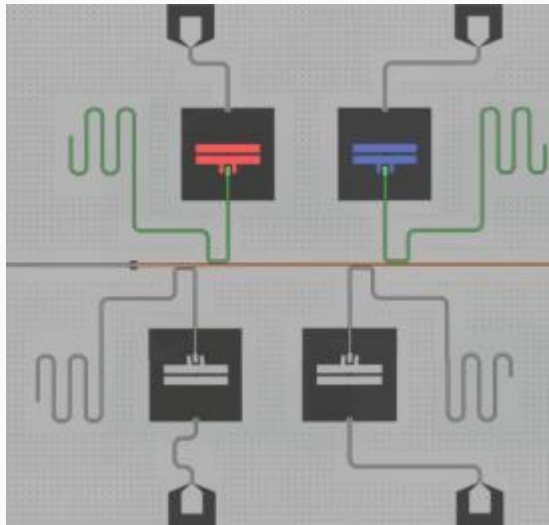
Research

- **Material** investigation – minimising unwanted nonlinearities, lower pump power requirements
- **Performance** optimisation – noise, operation frequency, bandwidths, operation temperature
- **Squeezing** amplification – generate squeezed states, amplification without adding noise

Research

Quantum information systems:

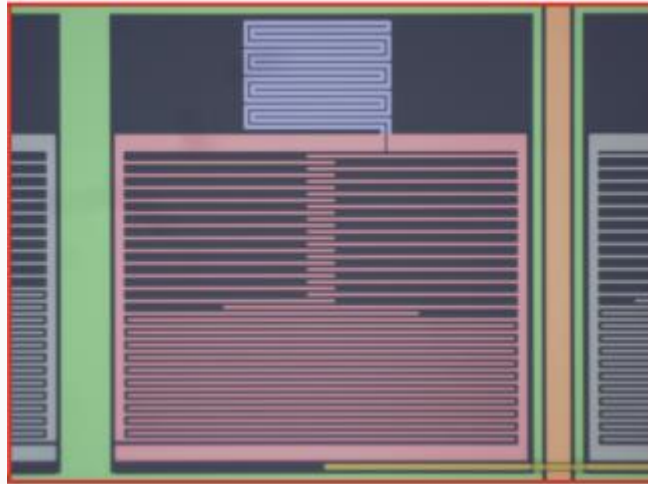
Qubit readout



(Ranzani, 2018)

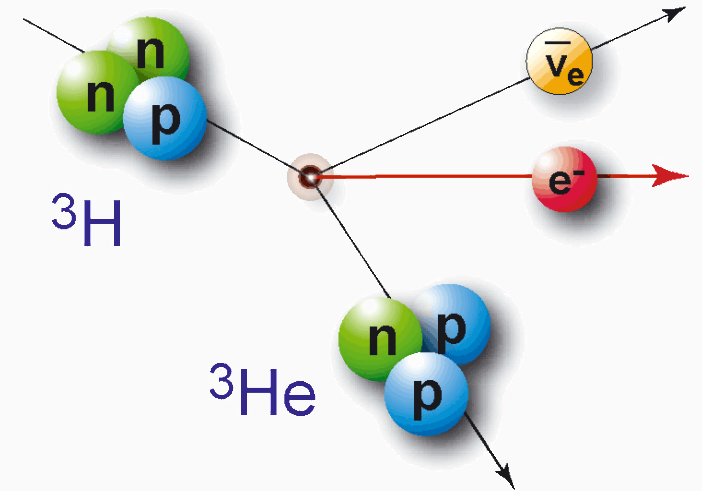
Observational astronomy:

Detector readout



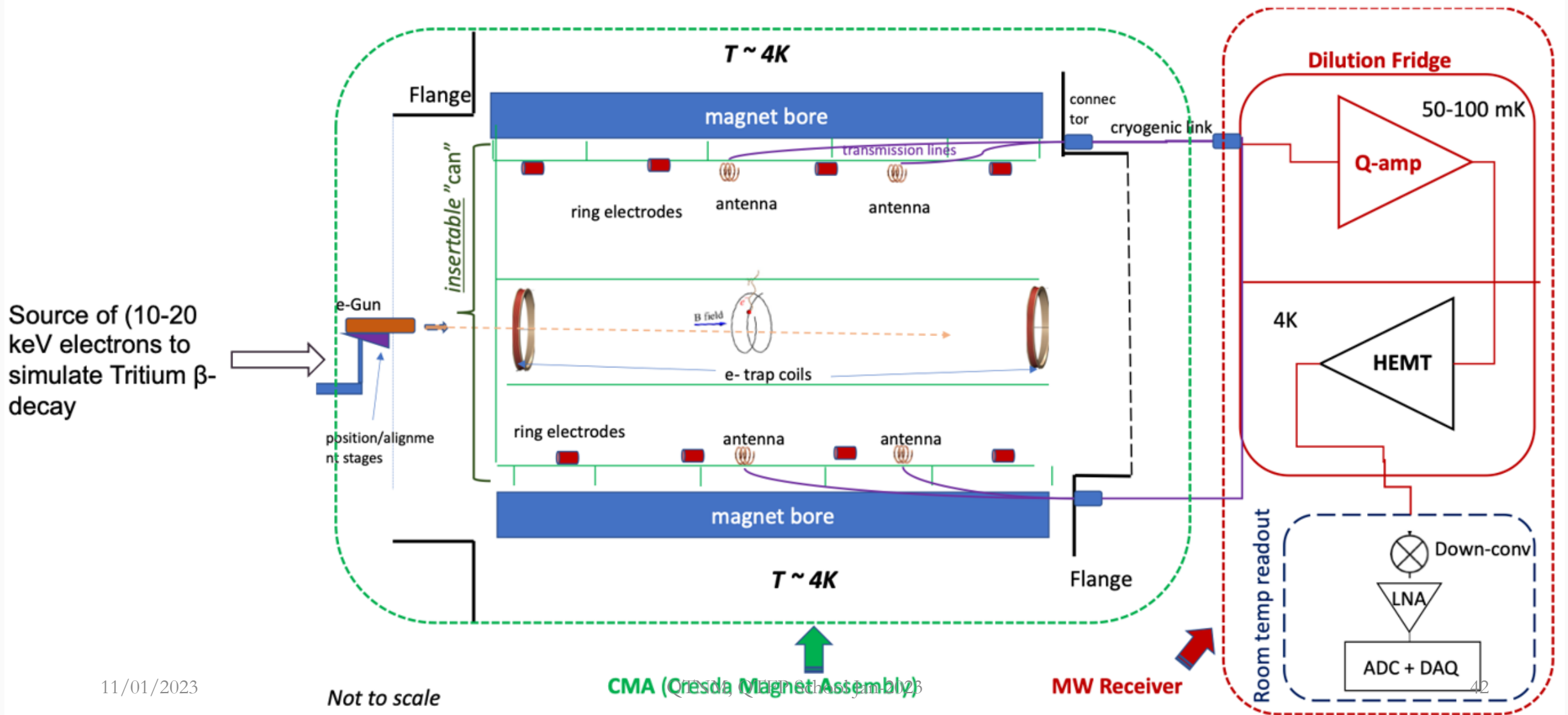
(Zobrist, 2022)

Fundamental physics:
Neutrino mass measurement



(QTNM, 2022)

QTNM: Quantum Technologies for Neutrino Mass



QTNM: Quantum Technologies for Neutrino Mass

- Quantum Technologies for Neutrino Mass

