

# “My” Recent Progress on CP Violation

**Qin Qin**

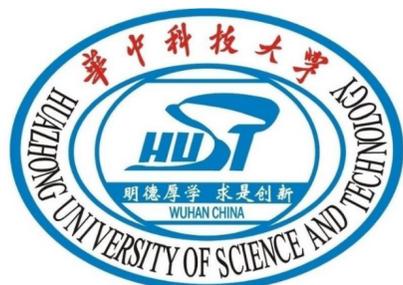
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arXiv: 2301.05848, 2211.07332

Collaborated with Y.F.Shen, W.J.Song, J.P.Wang, F.S.Yu

**IAS Program on High Energy Physics (HEP2023)**

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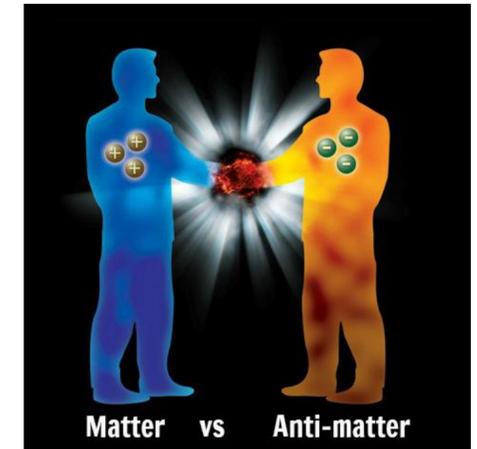


# Why CP violation?

- One of the Sakharov's criteria of matter-anti-matter asymmetry

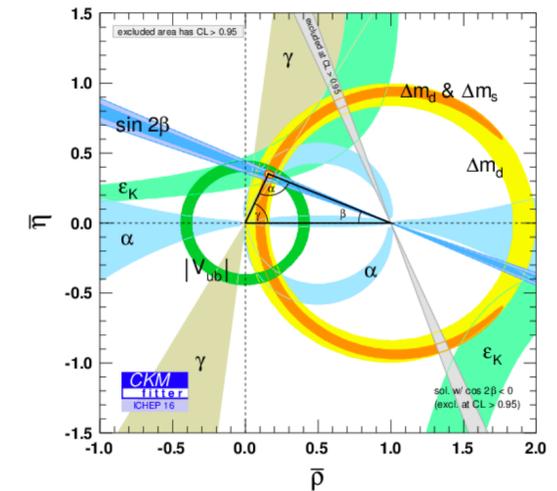
➔ Requires new source of CP violation

[Sakharov, '67]



- Determination of CKM matrix phase angles

➔ To test the unitarity of the CKM matrix



- Open windows to new dynamics beyond the SM

# CP violation in Flavor Physics

- **Discovery:** Mixing-induced CPV observed in Kaon decays

[Christenson, Cronin, Fitch, Turlay, '64]



Cronin Fitch

[Nobel Prize for Physics in 1980]

- The Kobayashi-Maskawa mechanism

[Kobayashi, Maskawa, '73]



Kobayashi Maskawa

[Nobel Prize for Physics in 2008]

- Direct CPV discovered in B meson decays

[BaBar & Belle, '01]

- **What's next?**

➡ More precise! **More observables!**

**More is better!**

# CP violation observables

- Common CPV observables

✓ CPV in decay (direct CPV)  $(M \rightarrow f) \neq (\bar{M} \rightarrow \bar{f})$

✓ CPV in mixing (indirect CPV)  $(M^0 \rightarrow \bar{M}^0) \neq (\bar{M}^0 \rightarrow M^0)$

✓ CPV in interference between a decay without and with **initial** mixing  $(M^0 \rightarrow f) + (M^0 \rightarrow \bar{M}^0 \rightarrow f)$

- CPV in interference between a decay without and with **final** mixing

$$(P \rightarrow M^0) + (P \rightarrow \bar{M}^0 \rightarrow M^0) \quad [\text{Yu, Wang, Li, PRL 119 (2017) 181802}]$$

- We propose/investigate two new types of CPV observables

➡ CPV induced by interference between two meson mixing [Shen, Song, QQ, 2301.05848]

$$(M_1^0 \rightarrow \bar{M}_2^0 \rightarrow M_2^0) + (M_1^0 \rightarrow \bar{M}_1^0 \rightarrow M_2^0)$$

➡ A special kind of time-reversal-odd (T-odd) observables  $TQ_- = -Q_-T$  [Wang, QQ, Yu, 2211.07332]

## Part 1. Double mixing CP violation

[Shen, Song, **QQ**, 2301.05848]

# Double mixing CP violation

- Visualization of CPV requires interference between different amplitudes

✓ CPV in decay (direct CPV): tree and penguin amplitudes

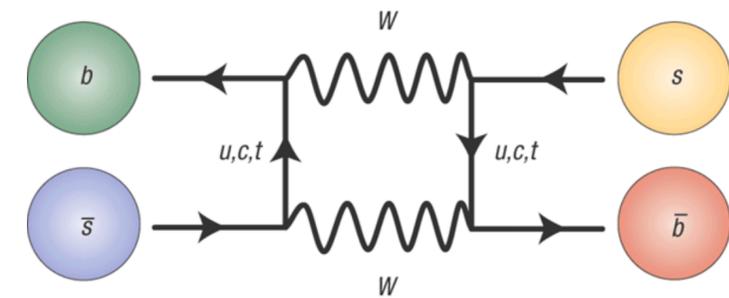
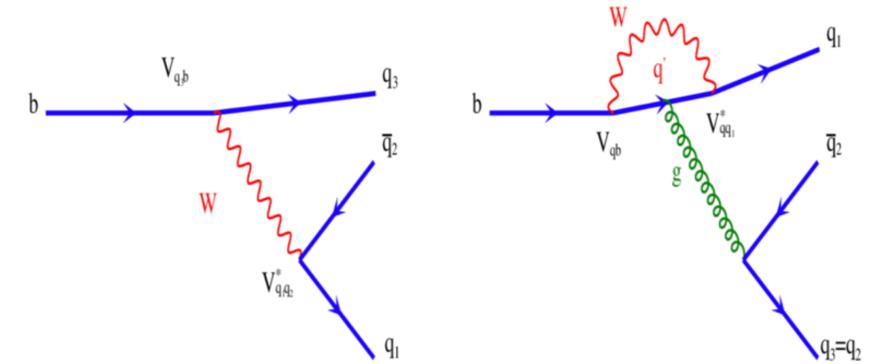
✓ CPV in mixing: different quark mediating box amplitudes

✓ CPV in interference between a decay without and with mixing

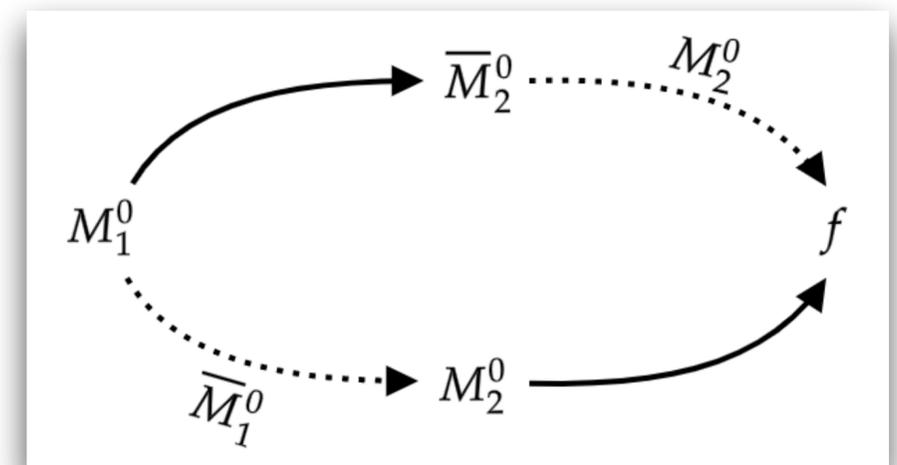
- Double mixing CP violation: induced by interference of different mixing paths of neutral mesons

$$B_s^0 \rightarrow \bar{B}_s^0 \rightarrow \rho^0 K^0 \rightarrow \rho^0 \pi^- e^+ \nu$$

$$B_s^0 \rightarrow \rho^0 \bar{K}^0 \rightarrow \rho^0 K^0 \rightarrow \rho^0 \pi^- e^+ \nu$$



$$M^0 \rightarrow f \quad M^0 \rightarrow \bar{M}^0 \rightarrow f$$



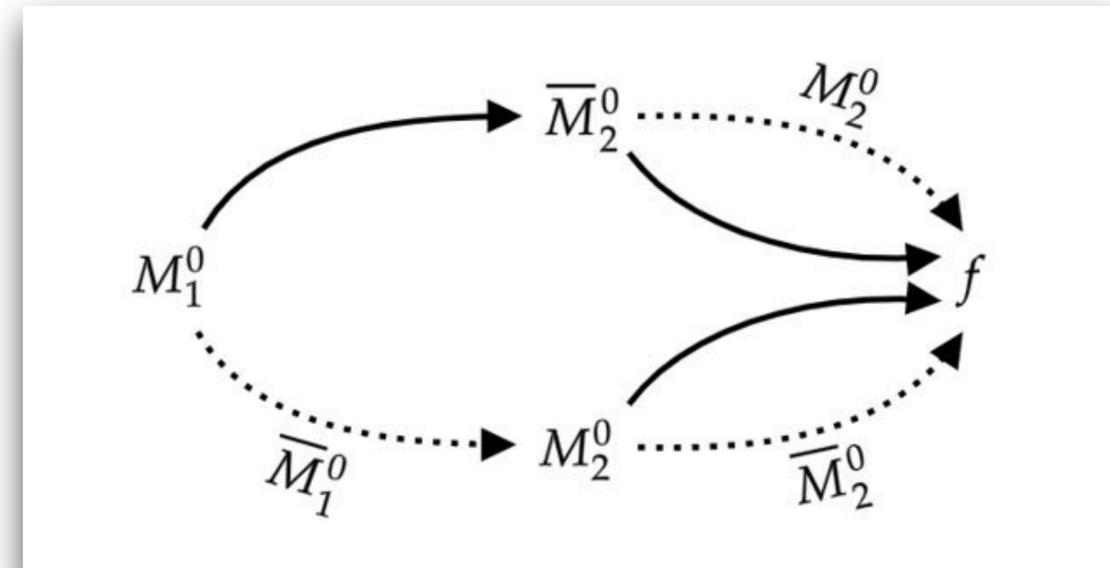
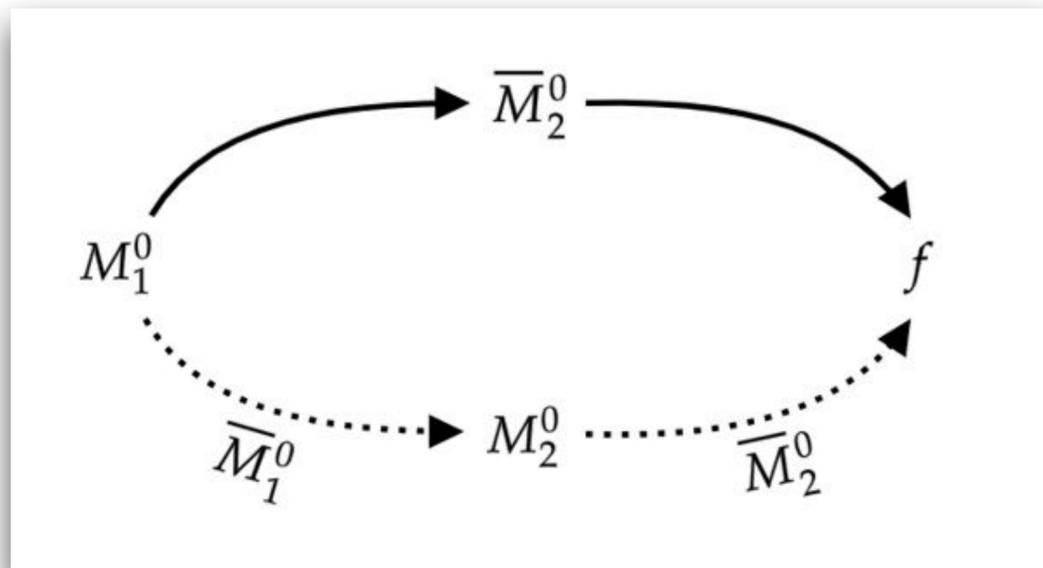
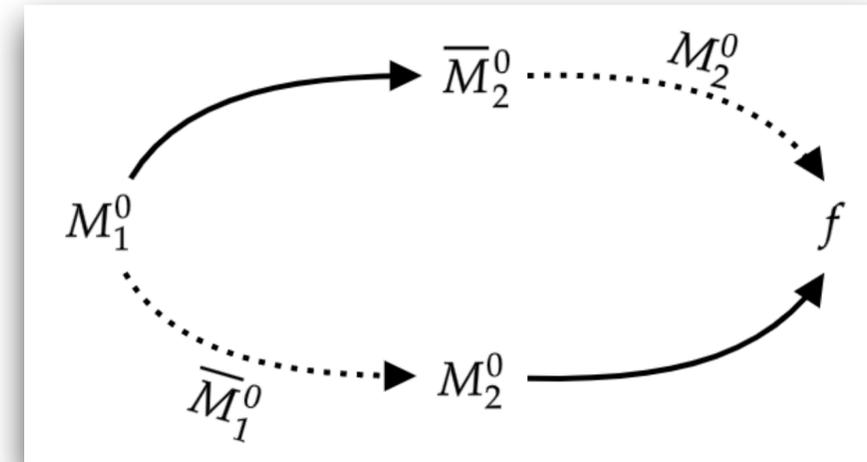
# Double mixing CP violation

- **Double mixing CP violation:** induced by interference of different mixing paths of neutral mesons

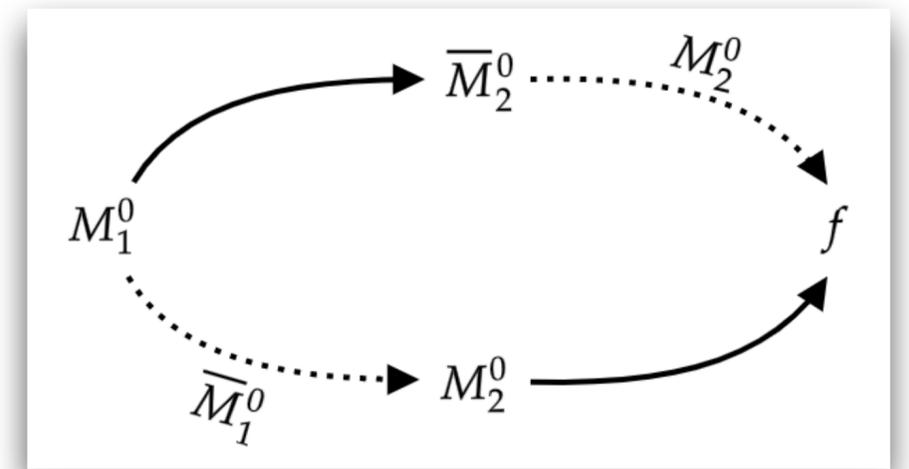
- ✓ At least two mixing mesons are involved

- ✓ At least two decays in the chain — — cascade decay

- More complicated cases



# Double mixing CP violation



- Take  $B_s^0(t_1) \rightarrow \rho^0 \bar{K}^0(t_2) \rightarrow \rho^0 \pi^- e^+ \nu$  as an example (penguin  $\approx 0$ )

$$M_1(t_1, t_2) \propto g_{1,+}(t_1) \left[ -\frac{p}{q} g_{2,-}(t_2) \right]$$

$\bar{K}^0 \rightarrow K^0$

$$M_2(t_1, t_2) \propto \left[ -\frac{q}{p} g_{1,-}(t_1) \right] g_{2,-}(t_2)$$

$B_s^0 \rightarrow \bar{B}_s^0$

$$|M^0(t)\rangle = g_+(t) |M^0\rangle - \frac{q}{p} g_-(t) |\bar{M}^0\rangle$$

$$|\bar{M}^0(t)\rangle = g_+(t) |\bar{M}^0\rangle - \frac{p}{q} g_-(t) |M^0\rangle$$

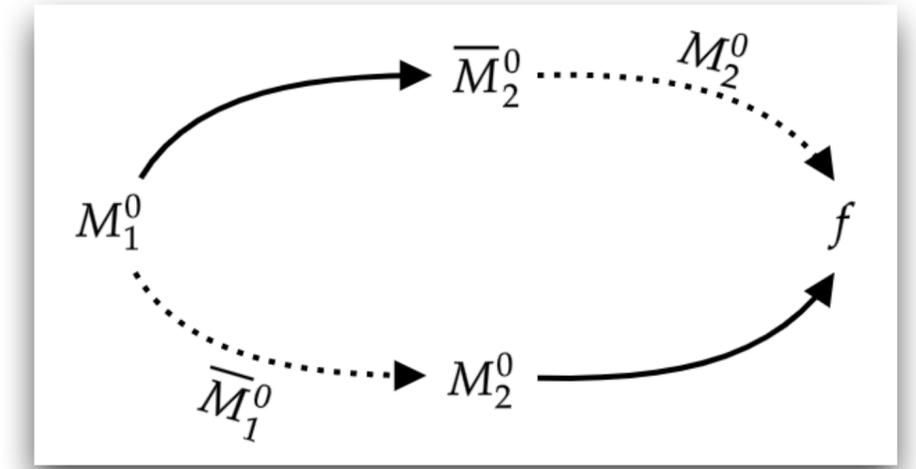
- Two-time dependent CPV:  $A_{CP}(t_1, t_2) = (|M|^2 - |\bar{M}|^2) / (|M|^2 + |\bar{M}|^2)$   $M \equiv M_1 + M_2$

Sources: 1.  $|M_1|^2$ ;  $\longrightarrow$  CP violation in  $K^0$  mixing

2.  $|M_2|^2$ ;  $\longrightarrow$  CP violation in  $B_s^0$  mixing

3.  $M_1^* M_2 + M_1 M_2^*$   $\longrightarrow$  CP violation in interference between  $K^0, B_s^0$  mixing

# Double mixing CP violation



- Take  $B_s^0(t_1) \rightarrow \rho^0 \bar{K}^0(t_2) \rightarrow \rho^0 \pi^- e^+ \nu$  as an example (penguin  $\approx 0$ )

$$A_{CP}(t_1, t_2) \propto \underbrace{|g_{1,+}(t_1)|^2}_{|M_1|^2} C_+(t_2) + \underbrace{|g_{1,-}(t_1)|^2}_{|M_2|^2} C_-(t_2) + \underbrace{e^{-\Gamma_1 t_1} \sinh \frac{\Delta\Gamma_1 t_1}{2} S_h(t_2) + e^{-\Gamma_1 t_1} \sin(\Delta m_1 t_1) S_n(t_2)}_{M_1^* M_2 + M_1 M_2^*}$$

$$C_+(t_2) = |g_{2,-}(t_2)|^2 \left( \left| \frac{p}{q} \right|_2^2 - \left| \frac{q}{p} \right|_2^2 \right),$$

$$S_h(t_2) = \frac{e^{-\Gamma_2 t_2}}{2} [-2 \sin \Delta m_2 t_2 \sin(\phi_1 + \phi_2 + 2\delta)]$$

$$C_-(t_2) = |g_{2,+}(t_2)|^2 \left( \left| \frac{q}{p} \right|_1^2 - \left| \frac{p}{q} \right|_1^2 \right),$$

$$S_n(t_2) = \frac{e^{-\Gamma_2 t_2}}{2} 2 \sinh \frac{\Delta\Gamma_2 t_2}{2} \sin(\phi_1 + \phi_2 + 2\delta)$$

Weak Phase

$$\begin{aligned} q_1/p_1 &= |q_1/p_1| e^{-i\phi_1} \\ q_2/p_2 &= |q_2/p_2| e^{-i\phi_2} \\ \langle \rho \bar{K} | B_s \rangle &= \langle \rho K | \bar{B}_s \rangle e^{2i\delta} \end{aligned}$$

$B_s$  &  $K$  mixing interference

# Double mixing CP violation

- Take  $B_s^0(t_1) \rightarrow \rho^0 \bar{K}^0(t_2) \rightarrow \rho^0 \pi^- e^+ \nu$  as an example (penguin  $\approx 0$ )

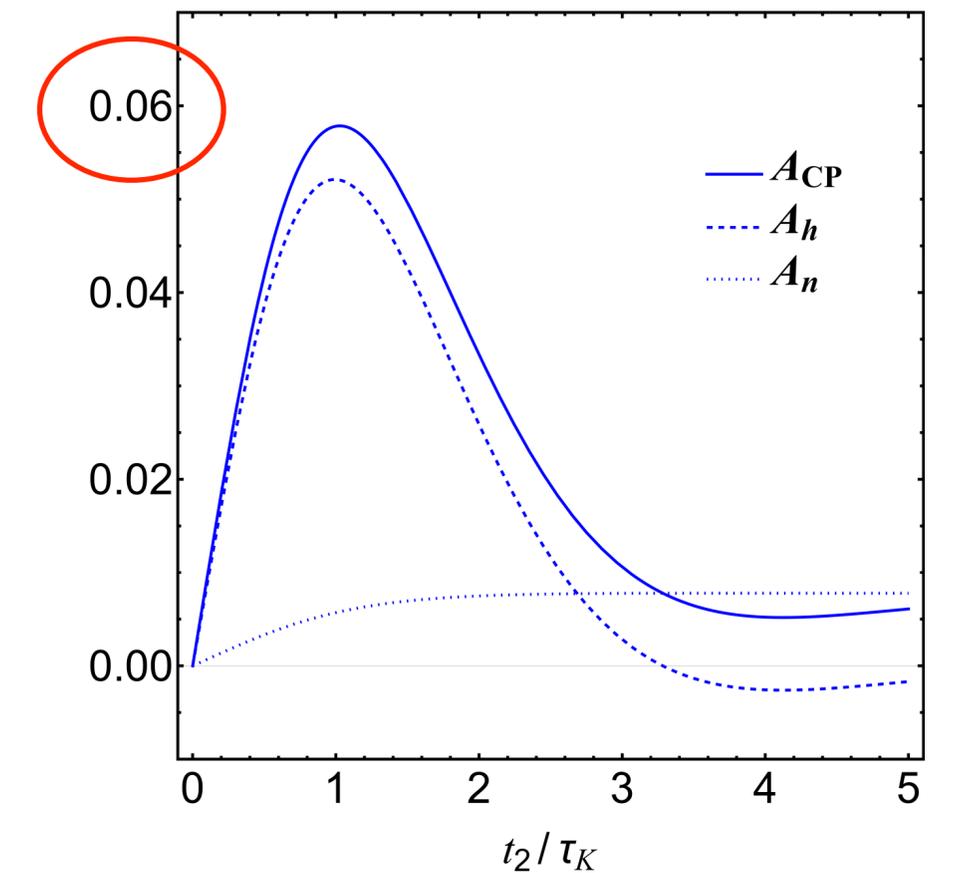
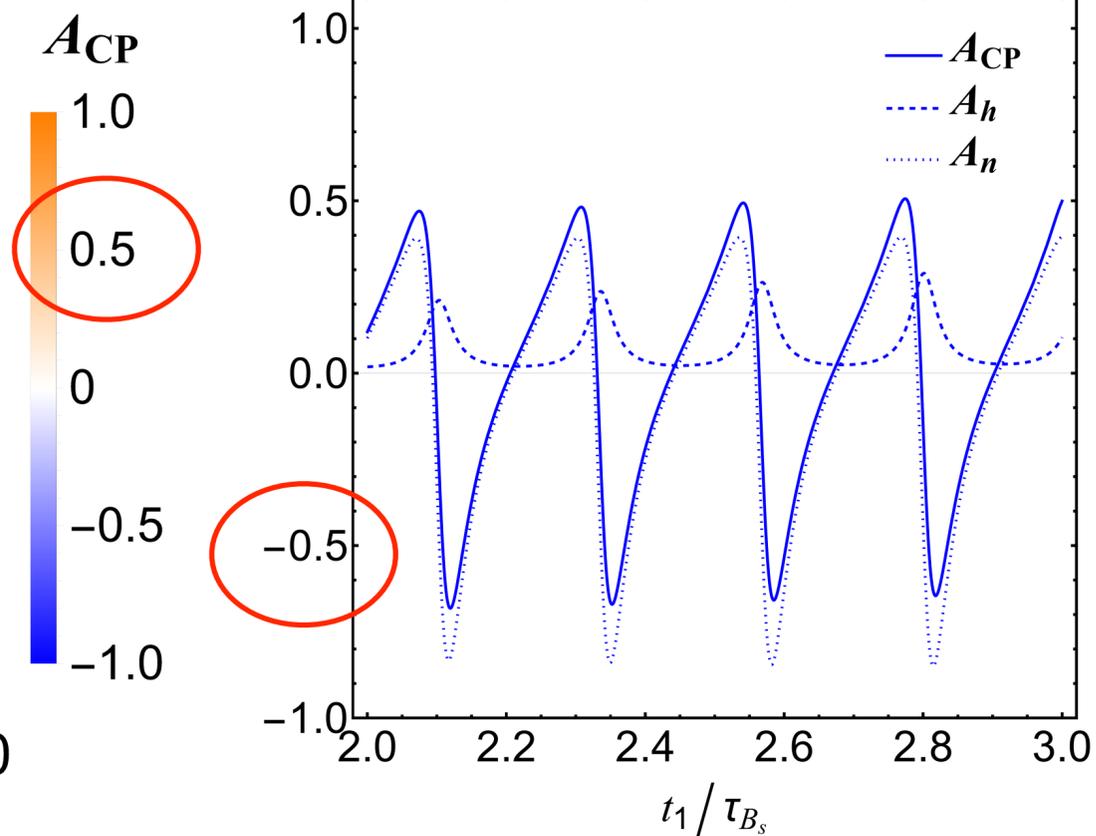
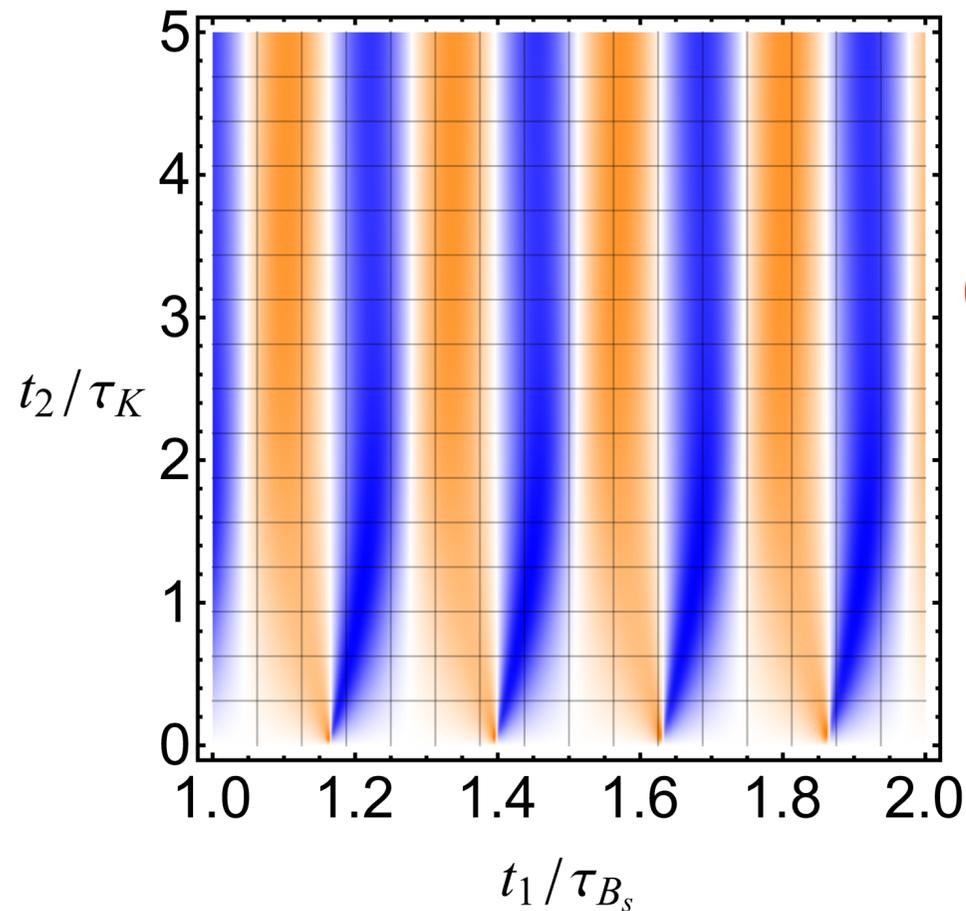
Time dependence:

$$S_n \propto \sin(\Delta m_1 t_1) \sinh \frac{\Delta \Gamma_2 t_2}{2} \sin(\phi_1 + \phi_2 + 2\delta)$$

$$S_h \propto \sinh \frac{\Delta \Gamma_1 t_1}{2} \sin(\Delta m_2 t_2) \sin(\phi_1 + \phi_2 + 2\delta)$$

$K_S + K_L$

$K_S, K_L$  interference



**Measurable at Z-factories like CEPC!**

## Double mixing CP violation

- The double-mixing CP asymmetry depends on two time variables and thus a **two-dimensional time dependence analysis** can be performed.
- It does not require nonzero strong phases, providing opportunities to directly **extract weak phases without strong pollution**.
- **Significant** in some channels, to be measured in experiments.
- To be investigated in **more channels** and to include **more accurate calculation**.

## Part 2. T-odd CP violation

[Wang, **QQ**, Yu, 2211.07332]

## Baryonic CPV

- T-odd CP violation is motivated in investigation of baryonic CPV.
- Is direct CP asymmetry the correct observable for baryonic CPV?

$$A_{CP} = \frac{|A|^2 - |\bar{A}|^2}{|A|^2 + |\bar{A}|^2}, \quad \begin{aligned} A &= A_1 e^{i\phi_1} e^{i\delta_1} + A_2 e^{i\phi_2} e^{i\delta_2} = A_1 e^{i\phi_1} e^{i\delta_1} (1 + r e^{i\phi} e^{i\delta}) \\ \bar{A} &= A_1 e^{-i\phi_1} e^{i\delta_1} + A_2 e^{-i\phi_2} e^{i\delta_2} = A_1 e^{-i\phi_1} e^{i\delta_1} (1 + r e^{-i\phi} e^{i\delta}) \end{aligned}$$



$$A_{CP} \propto 2r \sin \phi \sin \delta$$

Requirements:

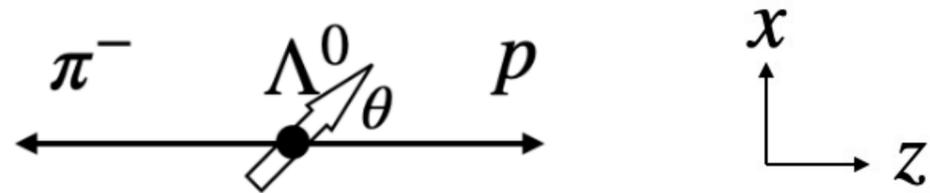
1. r is large
2. weak phase  $\phi$  is large
3. strong phase  $\delta$  is large

Far beyond control!

- Alternative observables to **satisfy** or **relax** the requirements?

## Polarization induced observables

- **Polarizations/helicities** of baryons provide fruitful observables.
- Lee-Yang parameters:  $\alpha, \beta, \gamma$



$$A(\Lambda^0 \rightarrow p\pi) = \bar{u}_p(S + P\gamma_5)u_\Lambda$$

**General Partial Wave Analysis of the  
 Decay of a Hyperon of Spin  $\frac{1}{2}$**   
 T. D. LEE\* AND C. N. YANG  
*Institute for Advanced Study, Princeton, New Jersey*  
 (Received October 22, 1957)

Theoretically, they are expressed by **partial wave amplitudes** (helicity amplitudes  $h_\pm = S \pm P$ ) as:

$$\alpha = \frac{2\text{Re}(S^*P)}{|S|^2 + |P|^2}, \quad \beta = \frac{2\text{Im}(S^*P)}{|S|^2 + |P|^2}, \quad \gamma = \frac{|S|^2 - |P|^2}{|S|^2 + |P|^2}$$

$$\frac{d\Gamma}{d\cos\theta} \propto 1 + \alpha \cos\theta$$

Experimentally, they are measured by **proton polarizations**:

$$P_p = \frac{(\alpha + \cos\theta)\hat{p} + \beta\hat{p} \times \hat{s} + \gamma(\hat{p} \times \hat{s}) \times \hat{p}}{1 + \alpha \cos\theta}$$

Spin measurements are difficult!

## Polarization induced observables

$$\frac{d\Gamma}{d\cos\theta} \propto 1 + \alpha \cos\theta$$

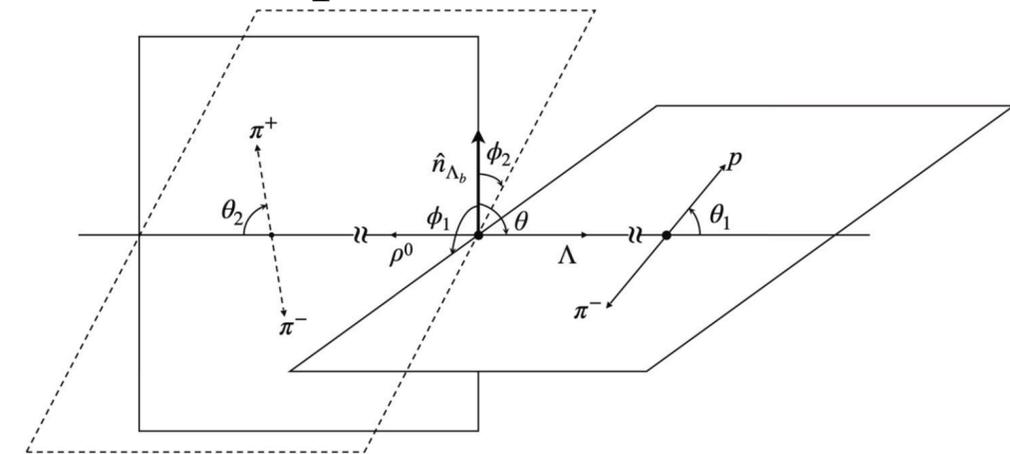
- **Key point:** particle spins are encoded in their decay products.
- With entangled  $\Xi^- \bar{\Xi}^+$  and  $\Xi^- \rightarrow \Lambda \pi^- \rightarrow p 2\pi^-$ , BESIII measure the Lee-Yang parameters and their induced CPV [BESIII, Nature 2022]

**Strong phase independent!**  $\leftarrow \Delta\phi_{\text{CP}} \approx \frac{\langle\alpha\rangle}{\sqrt{1-\langle\alpha\rangle^2}} \left( \frac{\beta + \bar{\beta}}{\alpha - \bar{\alpha}} \right)_{\Xi} = (-5 \pm 15) \times 10^{-3}$

- Application to more channels with Cascade decays (e.g.  $\Lambda_b \rightarrow \Lambda V \rightarrow p 3\pi$ )

1. Angular distribution encodes the **helicity amplitudes**
2. They induce CPVs with **different strong phase dependences**

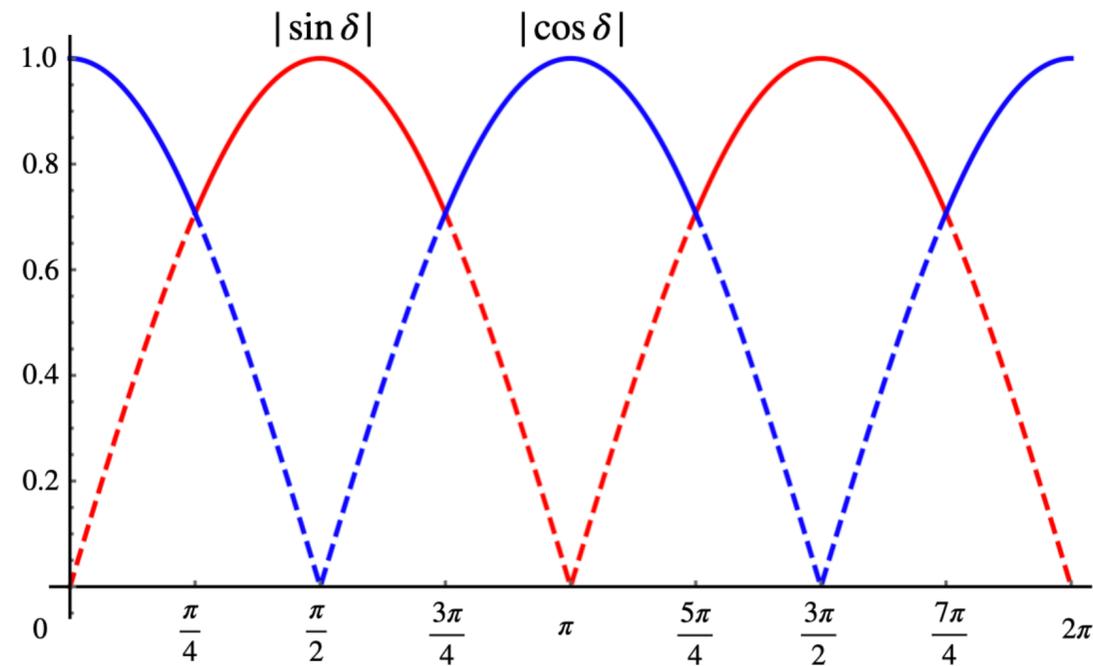
$$\sin\delta_s \text{ vs } \cos\delta_s$$



[Geng, Liu, Wei, et al, 2106.10628,2109.09524,2206.00348;Zhou, et al, 2210.15357]

## Polarization induced observables

- **Strong phase dependence:  $\sin \delta_s$  vs  $\cos \delta_s$**



Whatever the strong phase is, either  $|\sin \delta|$  or  $|\cos \delta|$  would be larger than 0.7.

- **Question:** does this complementarity generally exist?
- **Question:** if yes, how to find them systematically?

## T-odd correlation induced CP asymmetry

- T-odd correlation  $Q_-$  induced CPV have cosine dependence on strong phases

$$TQ_- = -Q_-T, \quad A_{CP}^{Q_-} \equiv \frac{\langle Q_- \rangle - \langle \bar{Q}_- \rangle}{\langle Q_- \rangle + \langle \bar{Q}_- \rangle} \propto \cos \delta_s$$

if it satisfies two conditions: (1) for the final-state basis  $\{|\psi_n\rangle, n=1,2,\dots\}$ , there is a unitary transformation  $U$ , s.t.  $UT|\psi_n\rangle = e^{-i\alpha}|\psi_n\rangle$ ; (2)  $UQ_-U^\dagger = Q_-$ .

### Proof:

$$\begin{aligned} \langle f|Q_-|f\rangle &= \langle i|S^\dagger Q_- S|i\rangle \\ &= \sum_{m,n} \langle \psi_i|S^\dagger|\psi_m\rangle \langle \psi_m|Q_-|\psi_n\rangle \langle \psi_n|S|\psi_i\rangle \\ &= \sum_{m,n} A_m^* A_n \langle \psi_m|Q_-|\psi_n\rangle. \end{aligned}$$

$$\begin{aligned} \langle \psi_m|Q_-|\psi_n\rangle &= \langle \psi_m|\mathcal{T}^\dagger \mathcal{T} Q_-|\psi_n\rangle^* \\ &= -\langle \psi_m|\mathcal{T}^\dagger Q_- \mathcal{T}|\psi_n\rangle^* \\ &= -\langle \psi_m|\mathcal{T}^\dagger U^\dagger U Q_- U^\dagger U \mathcal{T}|\psi_n\rangle^* \\ &= -\langle \psi_m|\mathcal{T}^\dagger U^\dagger Q_- U \mathcal{T}|\psi_n\rangle^* \\ &= -\langle \psi_m|Q_-|\psi_n\rangle^*, \end{aligned}$$



$$\langle f|Q_-|f\rangle \ni \text{Im}(A_m^* A_n)$$



$$A_{CP}^{Q_-} \propto \sin \delta_w \cos \delta_s$$

## T-odd correlation induced CP asymmetry

- Example 1. Triple product  $Q_1 \equiv (\vec{s}_1 \times \vec{s}_2) \cdot \hat{p}$  in  $P \rightarrow P_1 P_2$

$$T : \vec{p} \rightarrow -\vec{p}, h \rightarrow h; \quad U = R(\pi) : -\vec{p} \rightarrow \vec{p}, h \rightarrow h \quad \longrightarrow \quad \text{condition (i)}$$

$$T : Q_1 \rightarrow -Q_1; \quad U = R(\pi) : Q_1 \rightarrow Q_1 \quad \longrightarrow \quad \text{condition (ii)}$$



- Example 2. Triple product  $Q_p \equiv (\hat{p}_1 \times \hat{p}_2) \cdot \hat{p}_3$  in  $P \rightarrow P_1 P_2 P_3 P_4$

$$T : \vec{p} \rightarrow -\vec{p}; \quad U = P : -\vec{p} \rightarrow \vec{p} \quad \longrightarrow \quad \text{condition (i)}$$

$$T : Q_p \rightarrow -Q_p; \quad U = P : Q_p \rightarrow -Q_p \quad \not\longrightarrow \quad \text{condition (ii)}$$



[Wang, QQ, Yu, 2211.07332]

## T-odd correlation induced CP asymmetry

- For the decay  $\Lambda_b \rightarrow N^*(1520)K^*$ , three such T-odd correlations

**Triple product**

$$Q_1 \equiv (\vec{s}_1 \times \vec{s}_2) \cdot \hat{p} = \frac{i}{2}(s_1^+ s_2^- - s_1^- s_2^+)$$

**Hepta product**

$$Q_2 \equiv (\vec{s}_1 \cdot \hat{p})(\vec{s}_2 \cdot \hat{p})Q_1 + Q_1(\vec{s}_1 \cdot \hat{p})(\vec{s}_2 \cdot \hat{p}) = \frac{i}{2}s_1^z s_2^z (s_1^+ s_2^- - s_1^- s_2^+) + \frac{i}{2}(s_1^+ s_2^- - s_1^- s_2^+) s_1^z s_2^z$$

**Penta product**

$$Q_3 \equiv (\vec{s}_1 \cdot \vec{s}_2)Q_1 + Q_1(\vec{s}_1 \cdot \vec{s}_2) - Q_2 = \frac{i}{2}(s_1^+ s_1^+ s_2^- s_2^- - s_1^- s_1^- s_2^+ s_2^+)$$

- Their expectations are imaginary helicity amplitude interferences

$$\langle Q_3 \rangle = 2\sqrt{3} \operatorname{Im} (H_{+1,+\frac{3}{2}} H_{-1,-\frac{1}{2}}^* + H_{-1,-\frac{3}{2}}^* H_{+1,+\frac{1}{2}})$$

cos  $\delta_s$  vs sin  $\delta_s$

Exactly Complementary!

- Moreover, complementary T-even correlations are found

$$P_1 \equiv \vec{s}_1 \cdot \vec{s}_2 - (\vec{s}_1 \cdot \hat{p})(\vec{s}_2 \cdot \hat{p}), \quad P_2 \equiv (\vec{s}_1 \cdot \hat{p})(\vec{s}_2 \cdot \hat{p})P_1 + P_1(\vec{s}_1 \cdot \hat{p})(\vec{s}_2 \cdot \hat{p}),$$

$$P_3 \equiv P_1^2 - [\vec{s}_1^2 - (\vec{s}_1 \cdot \hat{p})^2][\vec{s}_2^2 - (\vec{s}_2 \cdot \hat{p})^2] - [(\vec{s}_1 \times \vec{s}_1) \cdot \hat{p}][(\vec{s}_2 \times \vec{s}_2) \cdot \hat{p}]$$

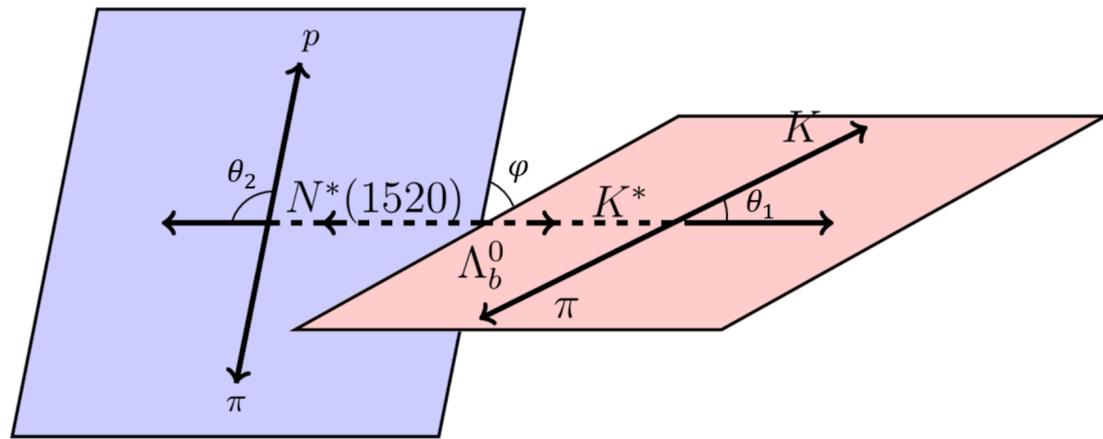
Real part

$$\langle P_3 \rangle \propto \operatorname{Re} (H_{+1,+\frac{3}{2}} H_{-1,-\frac{1}{2}}^* + H_{-1,-\frac{3}{2}}^* H_{+1,+\frac{1}{2}})$$

[Wang, QQ, Yu, 2211.07332]

## T-odd correlation induced CP asymmetry

- The expectations of the complementary T-odd and T-even correlations are both encoded in **angular distribution** of secondary decays of  $N^*(1520)K^*$



- Complementary CP asymmetries can thereby be measured, which depend on  $\cos \delta_s$  &  $\sin \delta_s$ .

$$\begin{aligned}
 \frac{d\Gamma}{dc_1 dc_2 d\varphi} &\propto s_1^2 s_2^2 \left( \left| \mathcal{H}_{+1,+\frac{3}{2}} \right|^2 + \left| \mathcal{H}_{-1,-\frac{3}{2}} \right|^2 \right) \\
 &+ s_1^2 \left( \frac{1}{3} + c_2^2 \right) \left( \left| \mathcal{H}_{+1,+\frac{1}{2}} \right|^2 + \left| \mathcal{H}_{-1,-\frac{1}{2}} \right|^2 \right) \\
 &+ 2c_1^2 \left( \frac{1}{3} + c_2^2 \right) \left( \left| \mathcal{H}_{0,-\frac{1}{2}} \right|^2 + \left| \mathcal{H}_{0,+\frac{1}{2}} \right|^2 \right) \\
 &- \frac{s_1^2 s_2^2}{\sqrt{3}} \text{Im} \left( \mathcal{H}_{+1,+\frac{3}{2}} \mathcal{H}_{-1,-\frac{1}{2}}^* + \mathcal{H}_{+1,+\frac{1}{2}} \mathcal{H}_{-1,-\frac{3}{2}}^* \right) \sin 2\varphi \quad \langle Q_3 \rangle \\
 &+ \frac{s_1^2 s_2^2}{\sqrt{3}} \text{Re} \left( \mathcal{H}_{+1,+\frac{3}{2}} \mathcal{H}_{-1,-\frac{1}{2}}^* + \mathcal{H}_{+1,+\frac{1}{2}} \mathcal{H}_{-1,-\frac{3}{2}}^* \right) \cos 2\varphi \quad \langle P_3 \rangle \\
 &- \frac{4s_1 c_1 s_2 c_2}{\sqrt{6}} \text{Im} \left( \mathcal{H}_{+1,+\frac{3}{2}} \mathcal{H}_{0,+\frac{1}{2}}^* + \mathcal{H}_{0,-\frac{1}{2}} \mathcal{H}_{-1,-\frac{3}{2}}^* \right) \sin \varphi \quad \langle Q_1 + 2Q_2 \rangle \\
 &+ \frac{4s_1 c_1 s_2 c_2}{\sqrt{6}} \text{Re} \left( \mathcal{H}_{+1,+\frac{3}{2}} \mathcal{H}_{0,+\frac{1}{2}}^* + \mathcal{H}_{0,-\frac{1}{2}} \mathcal{H}_{-1,-\frac{3}{2}}^* \right) \cos \varphi \quad \langle P_1 + 2P_2 \rangle
 \end{aligned}$$

[Wang, QQ, Yu, 2211.07332]

# Summary

# Summary

- **New CPV observables** are proposed
- including **double mixing CPV**, which could avoid strong pollution,
- and a type of **complementary T-odd CPV**, which would help discover **baryonic CPV**.
- Both of them awaits more investigation. **Collaboration is always welcome!**

**Thank you!**