

Probing Heavy Neutral Leptons at Future Experiments

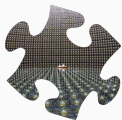
Juraj Klarić

February 13th, 2023



Some puzzles for physics beyond the Standard Model

Neutrino masses



The Baryon Asymmetry of the Universe

$$n_B/n_\gamma = 6.05(7) \times 10^{-10}$$

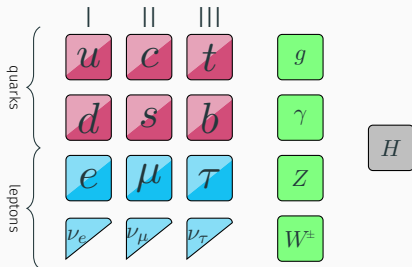
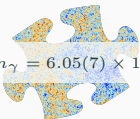


Image credits: Kamioka Observatory, ICRR, U. Tokyo; ESA and the Planck Collaboration

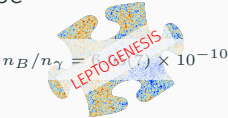
Some puzzles for physics beyond the Standard Model

Neutrino masses



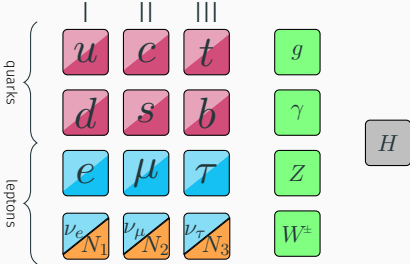
[Minkowski 1977...]

The Baryon Asymmetry of the Universe



[Fukugita/Yanagida '86...]

Image credits: Kamioka Observatory, ICRR, U. Tokyo; ESA and the Planck Collaboration

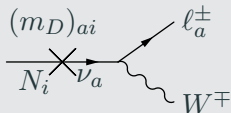


Where to look for HNLs?

Active neutrino masses

$$m_\nu = -m_D M_M^{-1} m_D^T$$

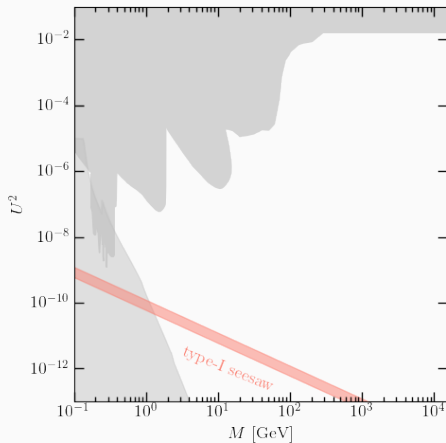
RHN mixing



$$U_{ai}^2 \equiv |(m_D M_M^{-1})_{ai}|^2$$

$$U^2 = \sum_{a,i} U_{ai}^2$$

$$U^2 \gtrsim m_\nu / M$$



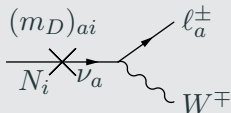
[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

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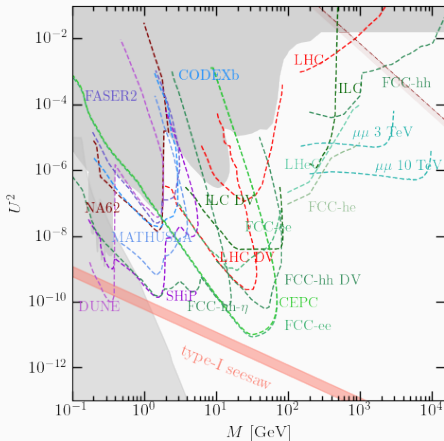
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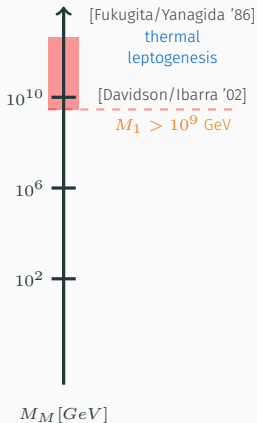
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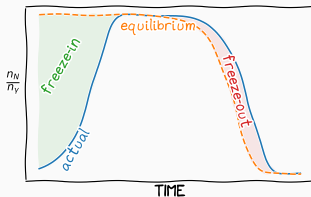
Low-scale leptogenesis

Leptogenesis mechanisms



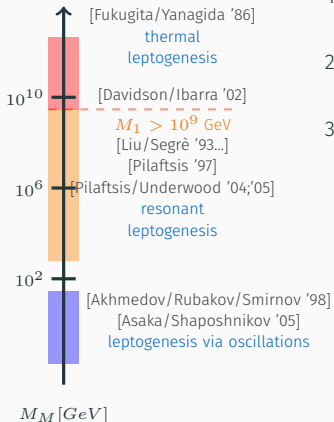
Sakharov conditions

1. Baryon number violation
sphaleron processes
2. C and CP violation
RHN decays and oscillations
3. Deviation from thermal equilibrium
freeze-in and freeze-out of RHN



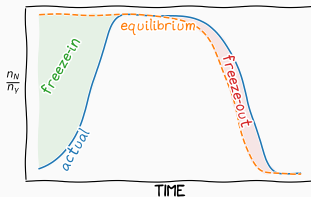
- for hierarchical RHN $M_1 \gtrsim 10^9 \text{ GeV}$

Leptogenesis mechanisms



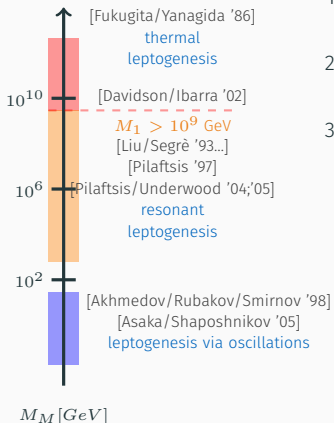
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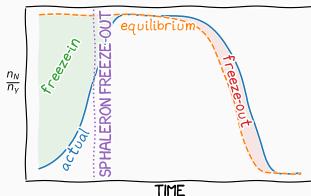
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Leptogenesis mechanisms



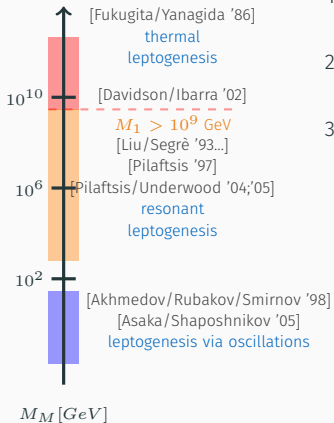
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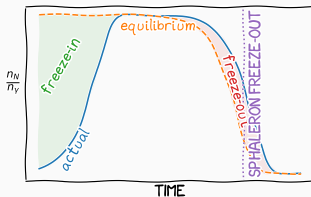
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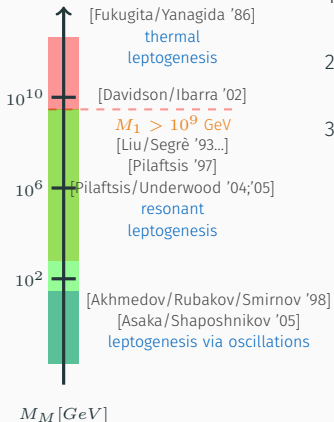
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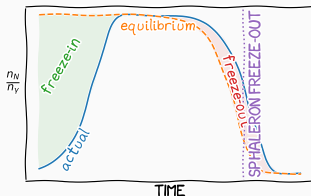
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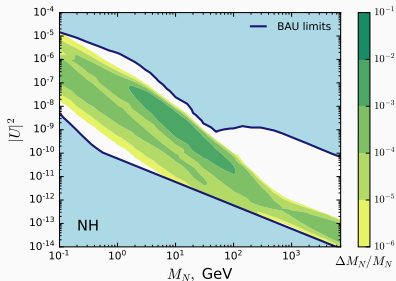
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- leptogenesis works in a wide range of RHN masses
- how are the low-scale mechanisms connected?

Results: The minimal model with 2 RHNs

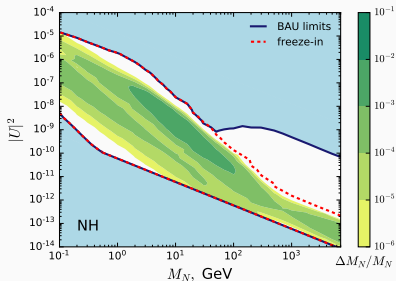


[JK/Timiryasov/Shaposhnikov 2103.16545]

- baryogenesis possible for all masses above 100 MeV!
- two main contributions to the BAU, from freeze-in and freeze-out
- there is significant overlap of the two regimes

- in resonant leptogenesis freeze-out (HNL decays) dominates, we can start with thermal initial conditions
- leptogenesis via oscillations is freeze-in dominated, we neglect HNLs falling out of equilibrium
- results depend on low-energy CP phases:
 - optimal phases $\delta = 0$ and $\eta = \pi/2$
 - less overlap for e.g. $\delta = \pi$ and $\eta = 0$
 - maximal $\Delta M/M \lesssim 10^{-1} \rightarrow 10^{-3}$

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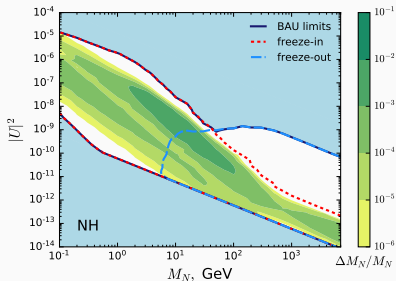


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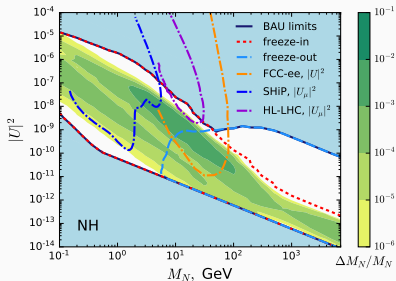


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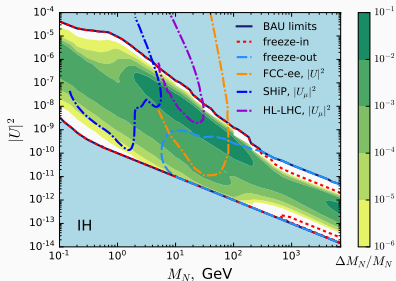


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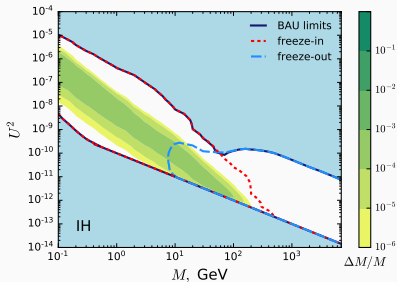


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How is $3 \neq 2$?

- asymmetry can be generated even without washout

[Akhmedov/Rubakov/Smirnov hep-ph/9803255]

- Sakharov II: CP
 - more CP phases than in the case with two RHN

- large hierarchy in the washout is possible

[Canetti/Drewes/Garbrecht 1404.7144]

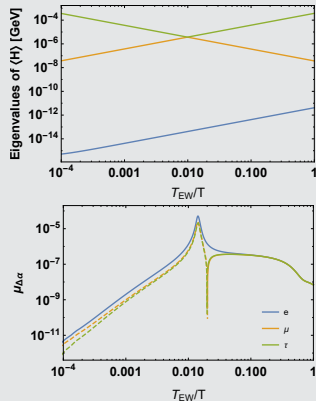
- Sakharov III: non-equilibrium

- level crossing between the heavy neutrinos

[Abada/Arcadi/Domcke/Drewes/JK/Lucente 1810.12463]

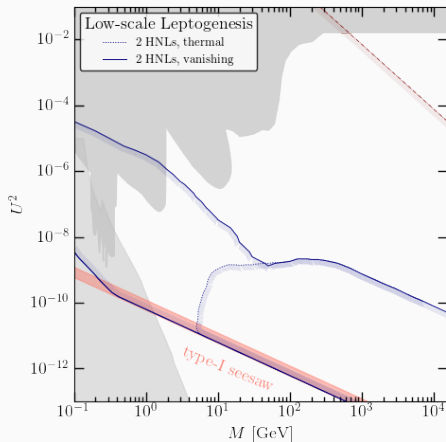
- Sakharov II: CP

Enhancement by level crossing



Results: Leptogenesis with 3 RHNs

- both **freeze-in** and **freeze-out** leptogeneses within reach of existing experiments
- all U^2 are allowed for experimentally accessible masses
- the maximal value of U^2 depends on m_1



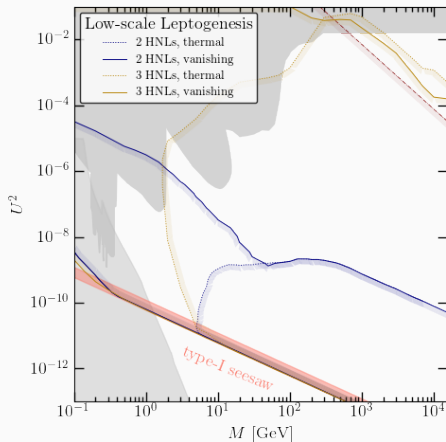
[figure adapted from Snowmass WPs 2203.08039 and 2203.05502]

[leptogenesis bounds from JK/Timiryasov/Shaposhnikov 2103.16545

and Drewes/Georis/JK 2106.16226]

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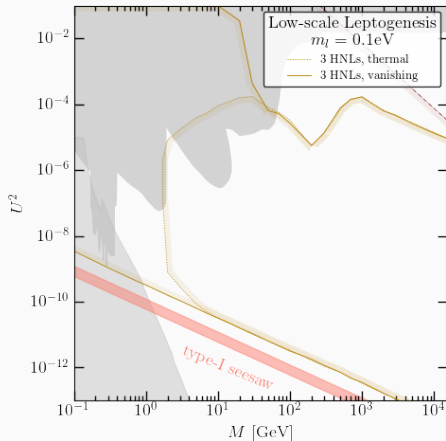
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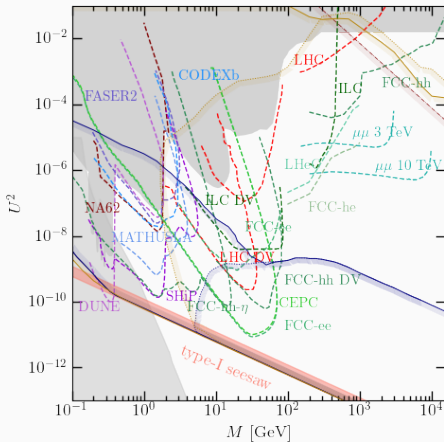
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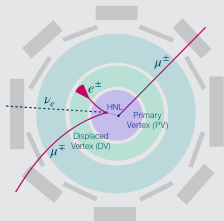
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Discovering Heavy Neutral Leptons

How to find HNLs?

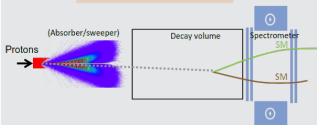
Displaced Vertices



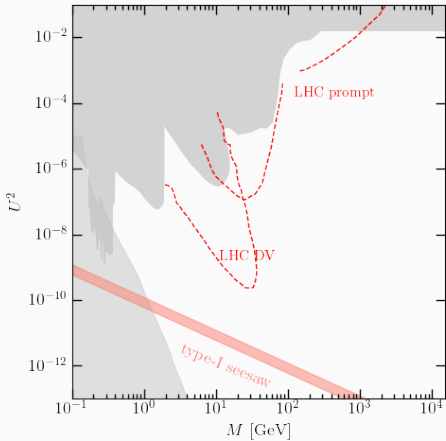
[graphic by D. Trischuk]

LLP searches

HS decay to SM particles



[graphic by A. Golutvin]

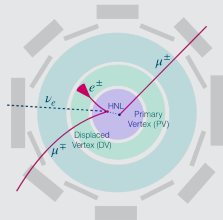


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[also see the the talk by S. Kulkarni]

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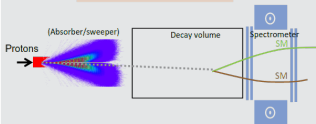
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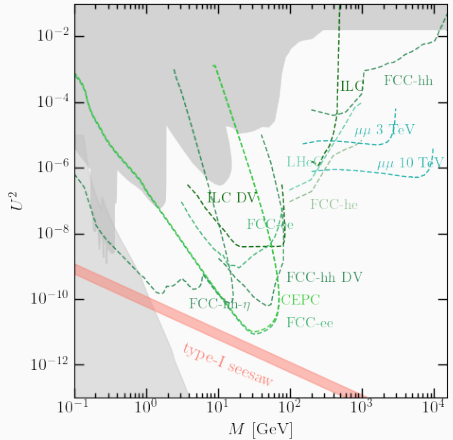
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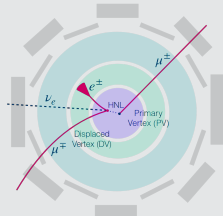


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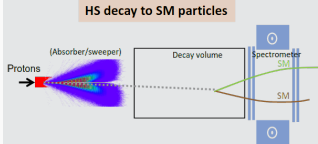
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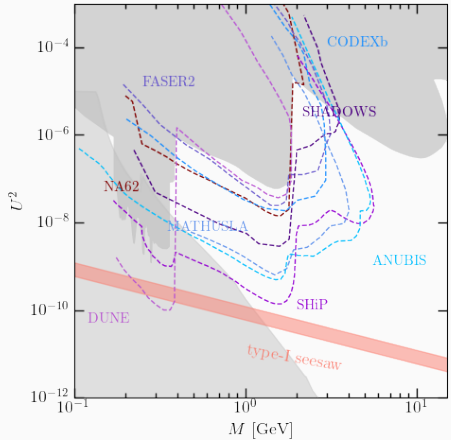


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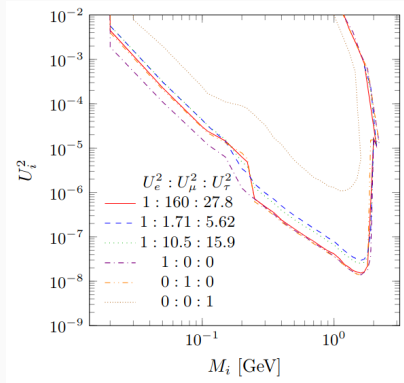
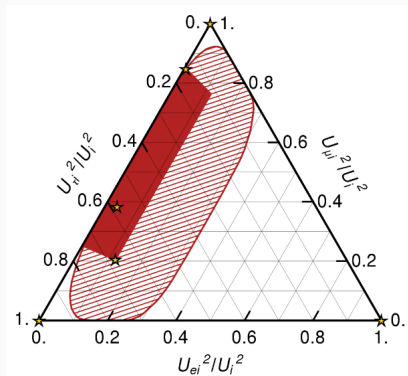
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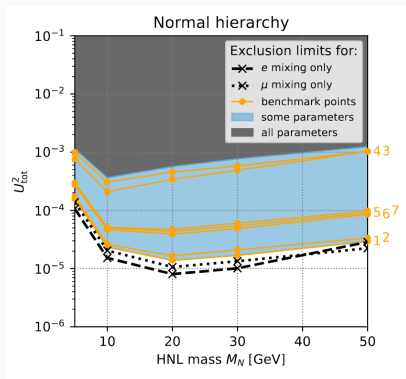
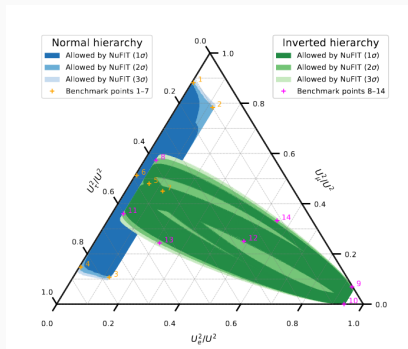
Sensitivity of experiments highly depends on mixing ratios: NA62 in beam dump mode



[Drewes/Hajer/JK/Lanfranchi 1801.04207]

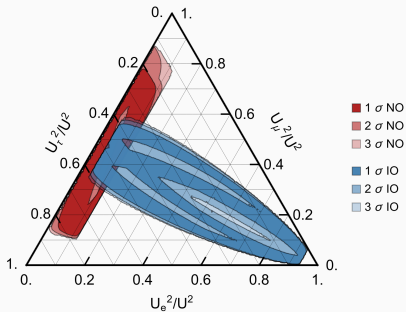
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ATLAS



[Tastet/Ruchayskiy/Timiryasov 2107.12980]

Constraints from the seesaw mechanism

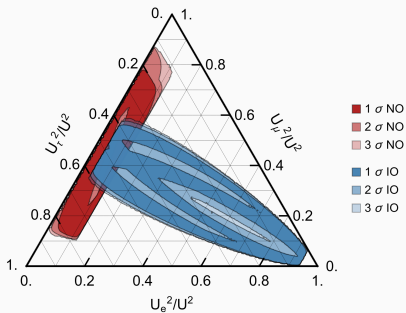


[Drewes/JK/Lopez-Pavon 2207.02742]

[using nuFIT 5.1 2007.14792]

- in the minimal seesaw model the flavour ratios are completely determined by U_{PMNS}
- uncertainty dominated by Majorana phase η , Dirac phase δ and θ_{23}
- allowed ratios become smaller as we pin down the PMNS parameters
- How to choose future-proof benchmarks?

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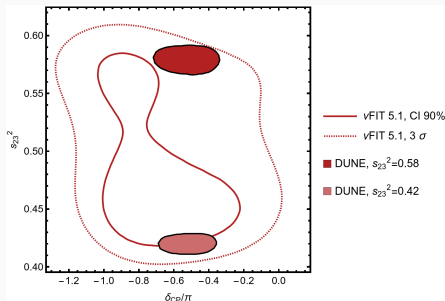
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Future sensitivity?

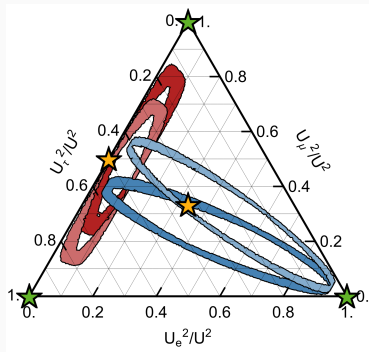
- significant improvement expected with DUNE and HyperK
- we can use the sensitivity estimates to estimate how the allowed flavor ratios change



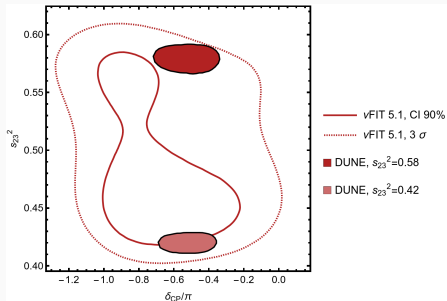
[nuFIT 5.1 2007:14792]

[DUNE TDR 2002.03005]

Future sensitivity?



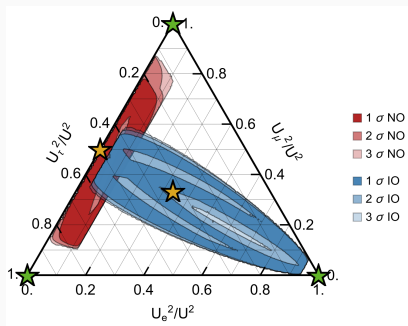
[Drewes/JK/Lopez-Pavon 2207.02742]



[nuFIT 5.1 2007.14792]

[DUNE TDR 2002.03005]

New Benchmark Points

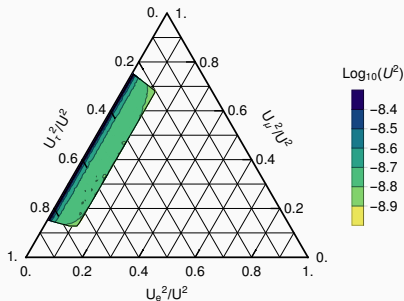


[Figure from 2207.02742]

- new benchmarks prepared for the HNL WG of the FIPs physics centre
- selection criteria:
 1. consistency with ν -osc. data
 2. added value
 3. symmetry considerations
 4. simplicity
 5. leptogenesis
- in addition to the single flavor benchmarks, we propose the new points:
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 0 : 1 : 1$
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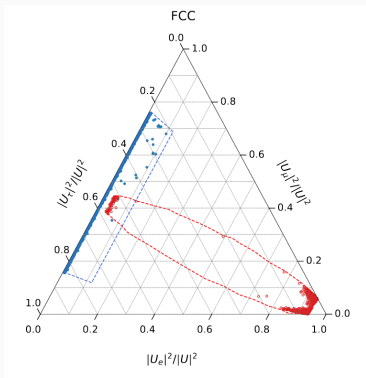


[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

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$$\Delta M/M = 10^{-2}$$

[Hernandez/Lopez-Pavon/Rius/Sandner 2207.01651]

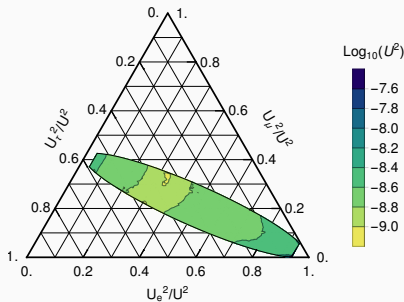
- new benchmarks prepared for the HNL WG of the FIPs physics centre
- selection criteria:
 1. consistency with ν -osc. data
 2. added value
 3. symmetry considerations
 4. simplicity
 5. leptogenesis
- in addition to the single flavor benchmarks, we propose the new points:
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 0 : 1 : 1$
 - $U_e^2 : U_\mu^2 : U_\tau^2 = 1 : 1 : 1$
- Common benchmarks can be used to compare the reach of different searches

From discovery to tests

Measuring flavor ratios at experiments

- the HNL branching ratios are constrained for a fixed U^2
- large number of HNLs possible at FCC-ee allow for measurement of U_e^2/U^2
- similar sensitivity @ SHiP

IO, $M = 30$ GeV

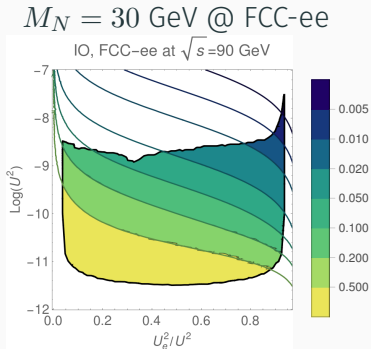


[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter]/JK

1710.03744]

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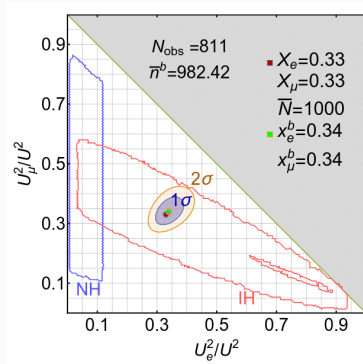
[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

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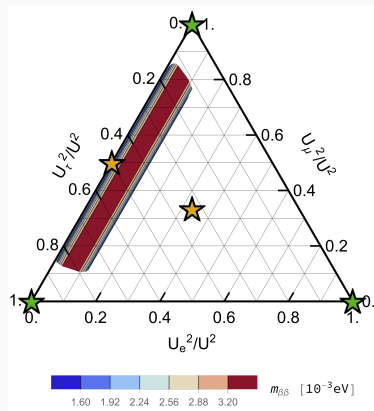
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$M_N = 1 \text{ GeV @ SHiP}$



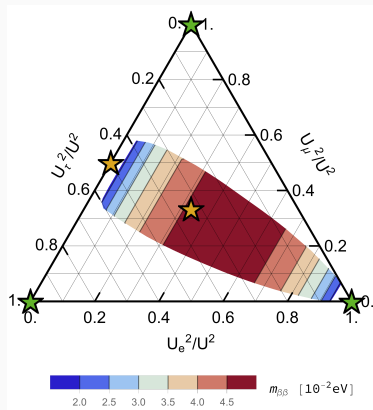
[Snowmass HNL WP 2203.08039]

Complementarity with neutrinoless double beta decay



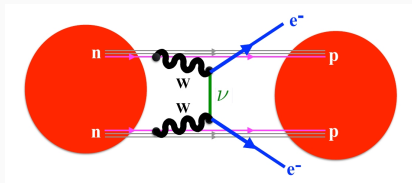
- $m_{\beta\beta}$ is a complementary probe of the flavor mixing ratios for $M_N \gg 100 \text{MeV}$
- excluding $m_{\beta\beta}$ limits allowed flavour ratios

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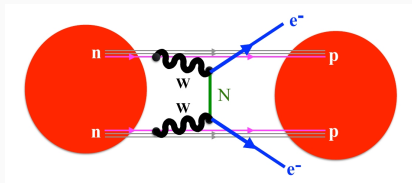
HNL contribution to neutrinoless double β decay



[figure from 1910.04688]

- RHN can contribute to $m_{\beta\beta}$
- large mass splitting is required to have an observable effect (not always compatible with leptogenesis)
- some leptogenesis scenarios can already be excluded by current results

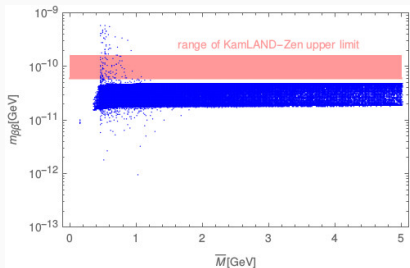
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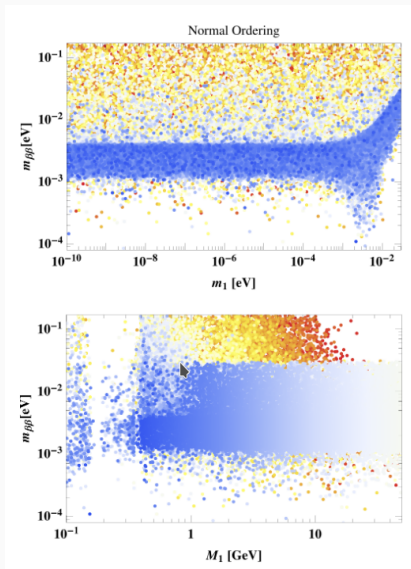


[Eijima/Drewes 1606.06221,

Hernández/Kekic/López-Pavón/Salvado 1606.06719]

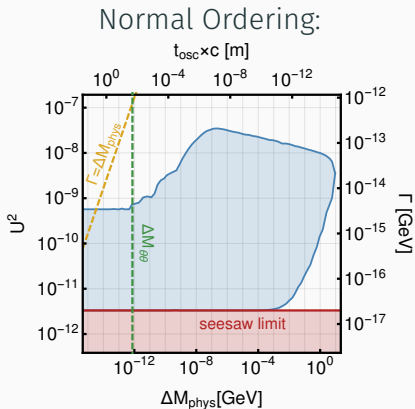
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Measuring the mass splitting in model with 2 HNLs



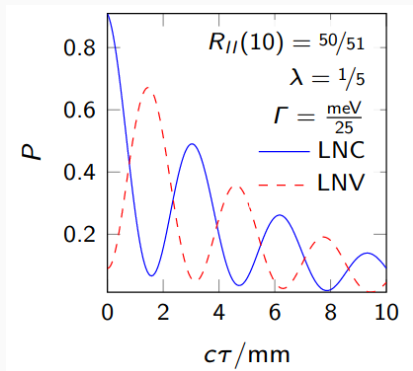
$$M = 30 \text{ GeV}$$

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

- large range of ΔM consistent with leptogenesis
- energy resolution of planned experiments - $\Delta M/M \sim \mathcal{O}(\text{few}\%)$
- Higgs vev contribution to RHN mass difference $\Delta M_{\theta\theta}$ practically implies lower limit on the mass splitting

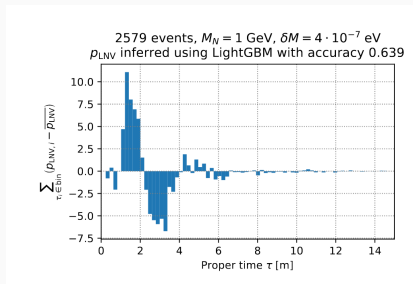
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[Antusch/Hajer/Roskopp 2210.10738]

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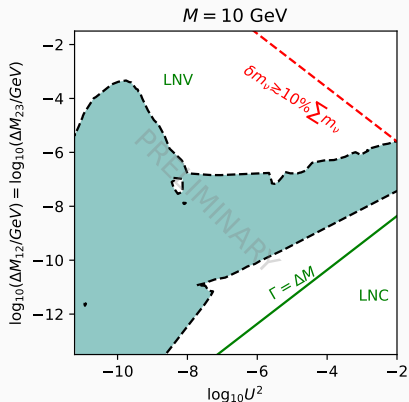
Measuring the mass splitting in model with 2 HNLs



[Tastet/Timiryasov 1912.05520]

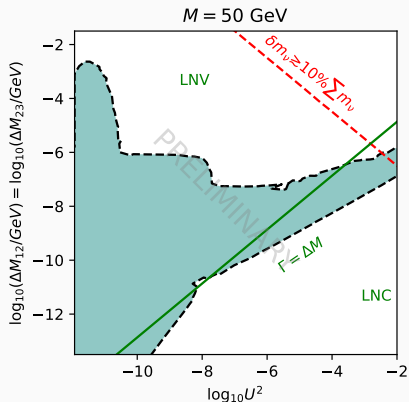
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Mass splittings with 3 HNLs



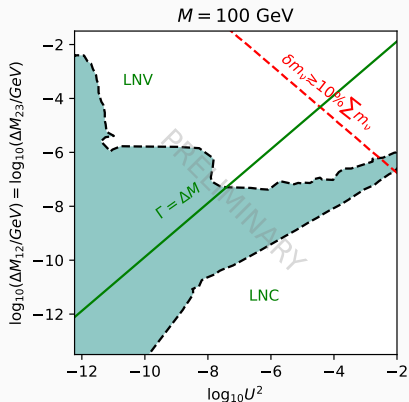
- benchmark with fixed $U_{\alpha I}^2/U^2$
- upper bound on U^2 arises through a combination of baryogenesis + fine tuning constraints
- leptogenesis consistent with both LNV and LNC RHN decays
- nontrivial LNV/LNC ratios can further constrain the RHN parameters

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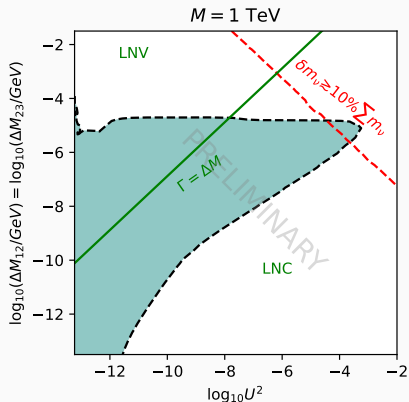
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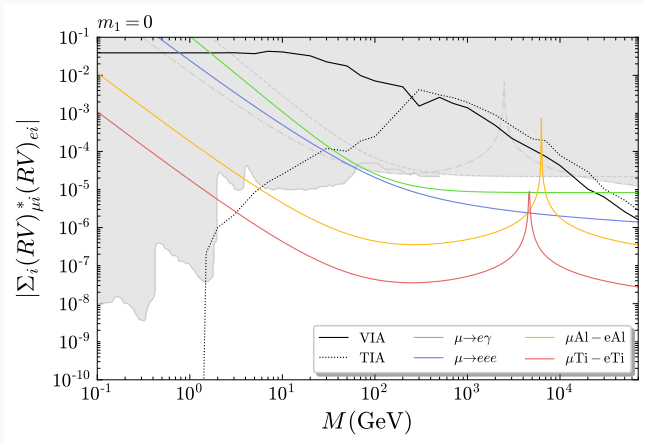
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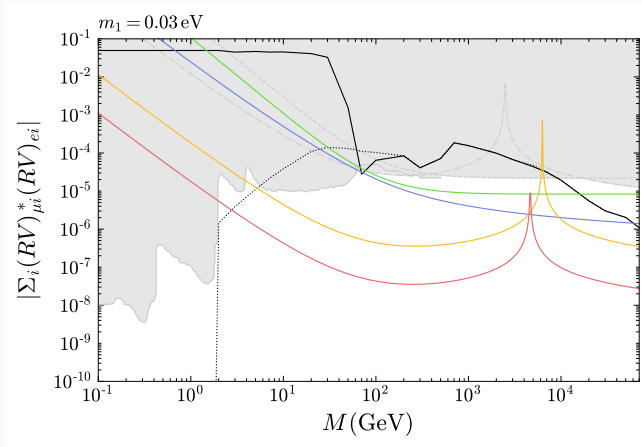
Indirect probes: Charged LFV



[Granelli/JK/Petcov 2206.04342]

- parameters space in the TeV region already severely constrained by cLFV observables
- future $\mu \rightarrow e$ conversion experiments can probe a large part of the $N = 3$ parameter space

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Conclusions

- **right-handed neutrinos** can offer a minimal solution to the origins of **neutrino masses** and the **baryon asymmetry of the Universe**
- the existence right-handed neutrinos can be **tested** at existing and near-future experiments
 - there is synergy between **high-energy** and **high-intensity** experiments!
 - together they will cover a large portion of the low-scale leptogenesis parameter space
- leptogenesis is a viable baryogenesis mechanism for **all heavy neutrino masses** above the $\mathcal{O}(100)$ MeV scale
- indirect searches can offer further insight

Thank you!

Large mixing angles and approximate B-L symmetry

- large U^2 require cancellations between different entries of the Yukawa matrices F
- this cancellation can be associated with an approximate lepton number symmetry

[Shaposhnikov hep-ph/0605047, Kersten Smirnov

0705.3221, Moffat Pascoli Weiland 1712.07611]

- symmetry broken by small parameters $\epsilon, \epsilon', \mu, \mu'$

Pseudo-Dirac pairs

$$N_s = \frac{N_1 + iN_2}{\sqrt{2}}, N_w = \frac{N_1 - iN_2}{\sqrt{2}}$$

B-L parametrisation

$$M_M = \bar{M} \begin{pmatrix} 1 - \mu & 0 & 0 \\ 0 & 1 + \mu & 0 \\ 0 & 0 & \mu' \end{pmatrix}$$

$$F = \frac{1}{\sqrt{2}} \begin{pmatrix} F_e(1 + \epsilon_e) & iF_e(1 - \epsilon_e) & F_e\epsilon'_e \\ F_\mu(1 + \epsilon_\mu) & iF_\mu(1 - \epsilon_\mu) & F_\mu\epsilon'_\mu \\ F_\tau(1 + \epsilon_\tau) & iF_\tau(1 - \epsilon_\tau) & F_\tau\epsilon'_\tau \end{pmatrix}$$

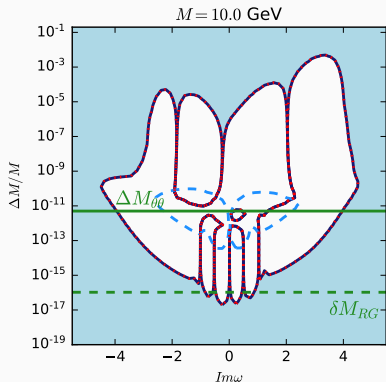
Fine tuning

- if present, symmetries are manifest to all orders in p.t.
- in the case of a large B-L breaking, radiative corrections can cause large neutrino masses
- we can use the size of radiative corrections to the light neutrino masses to quantify tuning

Fine Tuning

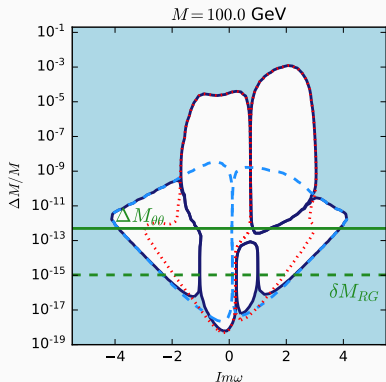
$$f.t.(m_\nu) = \sqrt{\sum_{i=1}^3 \left(\frac{m_i^{\text{loop}} - m_i^{\text{tree}}}{m_i^{\text{loop}}} \right)^2}$$

Slices of the parameter space



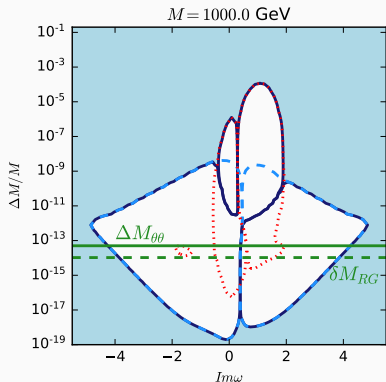
- two characteristic mass splittings
- mass splitting induced by the Higgs $\Delta M_{\theta\theta}$
- mass splitting induced by RG running δM_{RG}

Slices of the parameter space



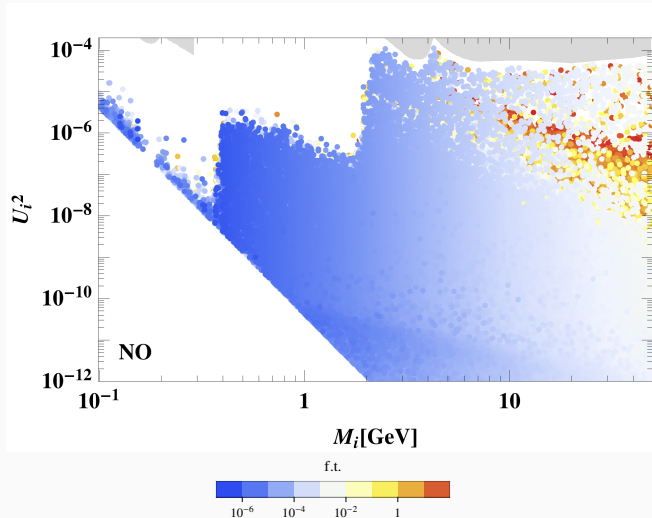
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Results: Leptogenesis with 3 RHN (Normal Ordering)



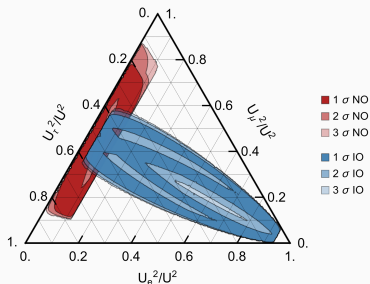
Hierarchy in the washout

- lepton asymmetry can survive washout if hidden in a particular flavor
- washout suppression

$$f \equiv \frac{\Gamma_a}{\Gamma} \sim \frac{U_a^2}{U^2}$$

- for 2 RHN $f > 5 \times 10^{-3}$
- for 3 RHN $f \ll 1$ possible

2 RHNs:



[Snowmass White Paper 2203.08039]

[Drewes/Garbrecht/Gueter]/JK 1609.09069]

[Caputo/Hernandez/Lopez-Pavon/Salvado 1704.08721]

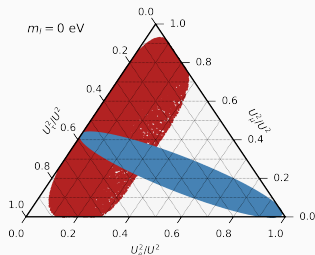
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[Drewes/Georis/JK 220x.xxxx]

[Chrzaszcz/Drewes/Gonzalo/Harz/Krishnamurthy/Weniger 1908.02302]

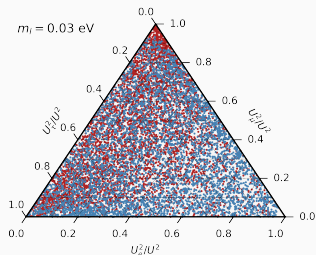
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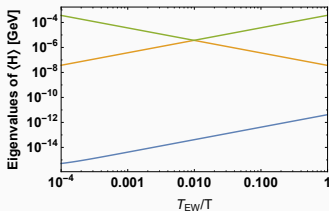
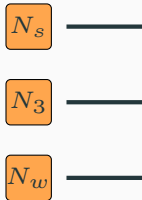
[Drewes/Georis/JK 220x.xxxx]

[Chrzaszcz/Drewes/Gonzalo/Harz/Krishnamurthy/Weniger 1908.02302]

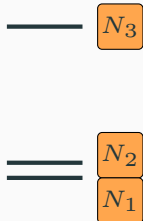
Enhancement due to level crossing

- in the $B - L$ symmetric limit two heavy neutrinos form a pseudo-Dirac pair
- the “3rd” heavy neutrino can be heavier than the pseudo-Dirac pair
- for $T \gg T_{EW}$, the pseudo-Dirac pair also has a thermal mass

$T \gg T_{EW}$

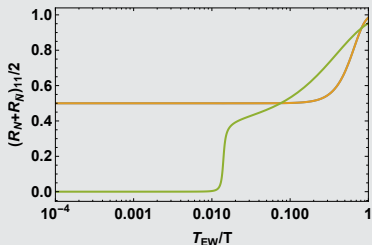


$T \ll T_{EW}$

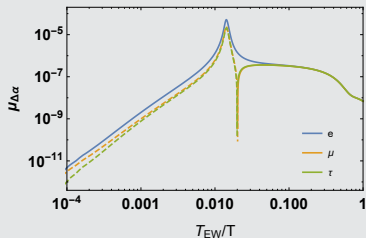


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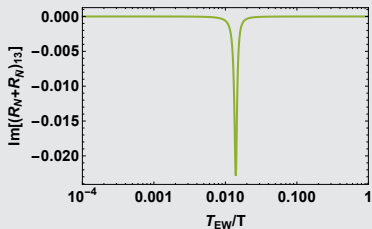
Heavy Neutrino Densities



Lepton flavour asymmetries



Heavy Neutrino correlations



Lepton number asymmetry

